

Cite this: *Mater. Adv.*, 2024,  
5, 5096

# Room temperature chemical transformation of SnSe to Ag<sub>2</sub>Se nanocrystals via cation exchange†‡

Yiqiao Huang and Pierre F. P. Poudeu  \*

Atomic-scale control of the chemical composition of semiconductor nanocrystals through a cation exchange reaction affords greater tunability in the design of multifunctional semiconductor composite nanocrystals. Here, we report a facile route to SnSe–Ag<sub>2</sub>Se composite nanocrystals using cation exchange at room temperature. Starting from freshly synthesized SnSe nanorods, we leverage the strong distortion of the Sn<sup>2+</sup> octahedral coordination in SnSe and the hard–soft acid–base (HSAB) principle, to promote the exchange of an Sn<sup>2+</sup> ion with two Ag<sup>+</sup> ions in methanol leading to Ag<sub>2</sub>Se/SnSe nanocomposites. The morphology and chemistry of the nanocrystals evolve from nanorods with SnSe@Ag<sub>2</sub>Se (core@shell) structures for SnSe-rich composites to nanorods with a random distribution of Sn<sup>2+</sup> and Ag<sup>+</sup> ions for nearly equimolar composites, and finally to irregular fragmented nanocrystals for Ag<sub>2</sub>Se-rich composites. A mechanistic understanding of the observed morphology evolution is discussed using the change in the cation coordination from octahedral (Sn<sup>2+</sup>) to tetrahedral (Ag<sup>+</sup>) geometry and the accompanying expansion of the hcp Se<sup>2–</sup> sublattice. Interestingly, the synthesized composite nanocrystals exhibit an optical band gap value tunable within a wide energy range by increasing the Ag<sub>2</sub>Se/SnSe ratio. This work provides a useful and facile strategy to modify the optical behavior of semiconductor nanomaterials, which can be leveraged for the design of better optical and/or photovoltaic devices.

Received 16th April 2024,  
Accepted 19th April 2024

DOI: 10.1039/d4ma00394b

rsc.li/materials-advances

## Introduction

Semiconductor nanocomposites have received considerable research attention in recent years due to their potential in multiple applications including thermoelectrics, photovoltaics, optoelectronics, photocatalysis, electrochemistry, and medical technology.<sup>1–11</sup> Achieving semiconductor nanocomposites with targeted properties necessitates the utilization of synthesis methods that enable control over the materials' chemistry and structure, which are critical parameters for performance optimization. Cation exchange has emerged as an elegant approach for the modification of the chemistry and microstructure of pre-synthesized semiconductor nanocrystals, which enables greater tunability of functional behavior and paves the way for the design of multifunctional semiconductor composite nanocrystals.<sup>1,12–14</sup> In a typical cation exchange reaction, the mobile cations in the parent nanocrystals tend to be replaced by the guest cations in solution, while the anion sublattice usually remains rigid.<sup>13,15</sup> The cation exchange process generally uses solvents and soluble salts to provide

guest cations. The cation exchange reaction that results in the retention of the anion sublattice requires some degrees of similarity in the atomic structure of the parent and the targeted phases. Typically, both the outgoing and the incoming cations adopt a similar coordination geometry within the anion sublattice and the exchange process results in a marginal change in the unit cell volume, which facilitates the conservation of the shape of the original nanocrystal.<sup>16,17</sup> Such a cation exchange process has been widely explored for the synthesis of metastable phases.<sup>16</sup> The cation exchange reaction that results in changes in the anion sublattice has also been reported.<sup>18–20</sup> This is typically observed when the ion exchange product either has a different crystal structure from the parent phase,<sup>19</sup> or has a similar structure with a significant change in the unit cell volume.<sup>18</sup> In such ion-exchange reactions, the shape of the original nanocrystal is not conserved due to the rearrangement of the anion sublattice.

One common structural feature of many cation exchange reactions reported to date, regardless of whether the cation exchange process changes or maintains the anion sublattice, is that the outgoing ion in the parent phase and the incoming cation in the targeted phase essentially adopt similar coordination geometry albeit marginal distortions. There are a few examples of cation exchange reactions that result in a change in the coordination of the exchanged cations. For instance, a reversible transformation between CdQ zinc blende nanoparticles and PbQ rock-salt

Laboratory for Emerging Energy and Electronic Materials, Department of Materials Science and Engineering, University of Michigan, Ann Arbor, Michigan 48109, USA.  
E-mail: ppoudeup@umich.edu

† Dedicated to Professor Michael Ruck on the occasion of his 60th birthday.

‡ Electronic supplementary information (ESI) available. See DOI: <https://doi.org/10.1039/d4ma00394b>



nanoparticles (Q = S, Se, Te) has been demonstrated.<sup>21–26</sup> However, these reactions require thermal activation as well as complex organic solvents/ligands such as octadecene (ODE), 1-dodecanethiol (DDT), trioctylphosphine oxide (TOPO), trioctylphosphine (TOP), and oleylamine (OLA)<sup>27–31</sup> to facilitate the change in the coordination of Cd<sup>2+</sup> and Pb<sup>2+</sup> ions between the tetrahedral and octahedral sites in the face-centered cubic Q<sup>2–</sup> sublattice. Moreover, these reactions are typically conducted under oxygen-free conditions, which poses a constraint on their scalability for industrial applications.

Here, we report a facile cation exchange between the octahedrally coordinated Sn<sup>2+</sup> ion in SnSe and the tetrahedrally coordinated Ag<sup>+</sup> ion in Ag<sub>2</sub>Se, which enables the fabrication of a range of Ag<sub>2</sub>Se/SnSe nanocomposites starting from the SnSe nanorod (NR) precursor. It is remarkable that the transformation of SnSe to Ag<sub>2</sub>Se proceeds at room temperature and does not use complex organic solvents, unlike other cation exchange reactions that result in a change in the coordination of the exchanged cations.<sup>27–31</sup> The Ag<sub>2</sub>Se/SnSe ratio in the final nanocomposites can be tuned by controlling the Ag/Sn molar ratio in the starting mixture. Interestingly, such control of chemical composition also enables the modification of the morphology as well as the elemental distribution within the final nanoparticles. We discuss the underlying atomic-scale mechanism by using X-ray powder diffraction (XRD), transmission electron microscopy (TEM), scanning transmission electron microscopy (STEM), and X-ray photoelectron spectroscopy (XPS) data on various SnSe/Ag<sub>2</sub>Se composites to systematically track the chemical transformation process. The optical band gap of various composites extracted from UV-vis-NIR diffuse reflectance spectroscopy data is also reported.

## Results and discussion

### Synthesis and structure

SnSe nanocrystals were synthesized following a modified surfactant-free method<sup>32</sup> by mixing freshly prepared solutions of NaHSe and Na<sub>2</sub>SnO<sub>2</sub> at room temperature followed by 2 h magnetic stirring. The phase purity and crystal structure of the synthesized SnSe precursor were confirmed by X-ray diffraction (XRD) and transmission electron microscopy. The experimental XRD pattern (Fig. 1(a)) of the synthesized SnSe nanoparticles can be indexed to the orthorhombic crystal structure (space group *Pnma* (#62), *a* = 11.502 Å, *b* = 4.153 Å, and *c* = 4.45 Å) of SnSe (ICSD# 50542). Transmission electron microscopy (TEM) revealed that the SnSe precursor adopts a nanorod-like morphology with dimensions ranging from 80 to 200 nm in length and 10 to 20 nm in diameter that are agglomerated into clusters of various sizes (Fig. 2(a)). The crystallinity of the SnSe nanorod was confirmed by high resolution TEM (HRTEM), which reveals the 0.36 nm interplanar distance of the (201) plane of the SnSe crystal (Fig. 2(f)).

Cation exchange reactions towards the synthesis of various Ag<sub>2</sub>Se/SnSe composite nanocrystals were performed by mixing the SnSe nanorods dispersed in methanol with varying

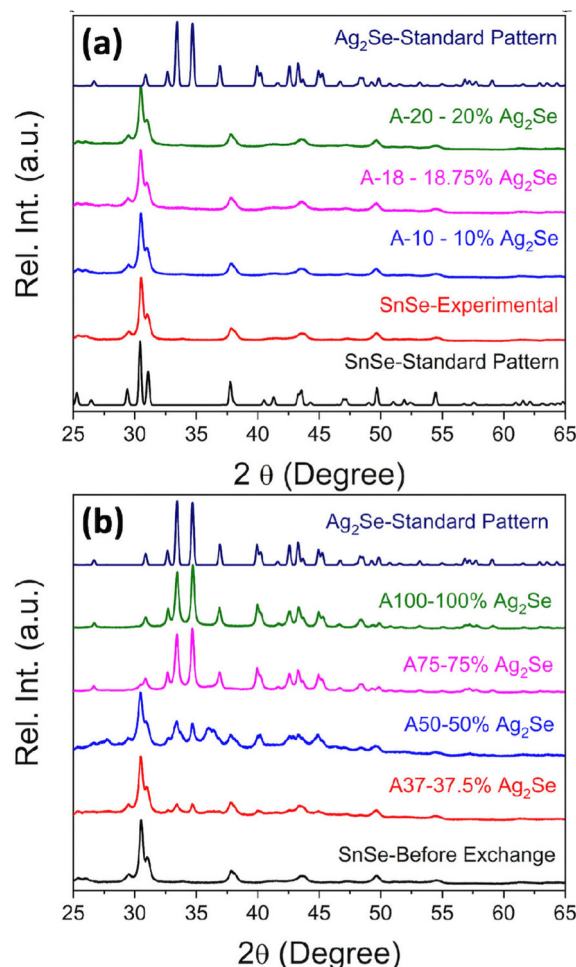
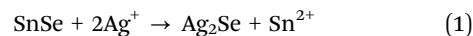


Fig. 1 (a) XRD pattern of the synthesized SnSe precursor compared with the standard XRD pattern of SnSe (ICSD#50542), and XRD patterns of Ag<sub>2</sub>Se/SnSe nanocomposites with low Ag/Sn ratio showing only peaks from the SnSe phase; (b) XRD patterns of Ag<sub>2</sub>Se/SnSe nanocomposites with high Ag/Sn ratio displaying diffraction peaks of both SnSe and Ag<sub>2</sub>Se phases.

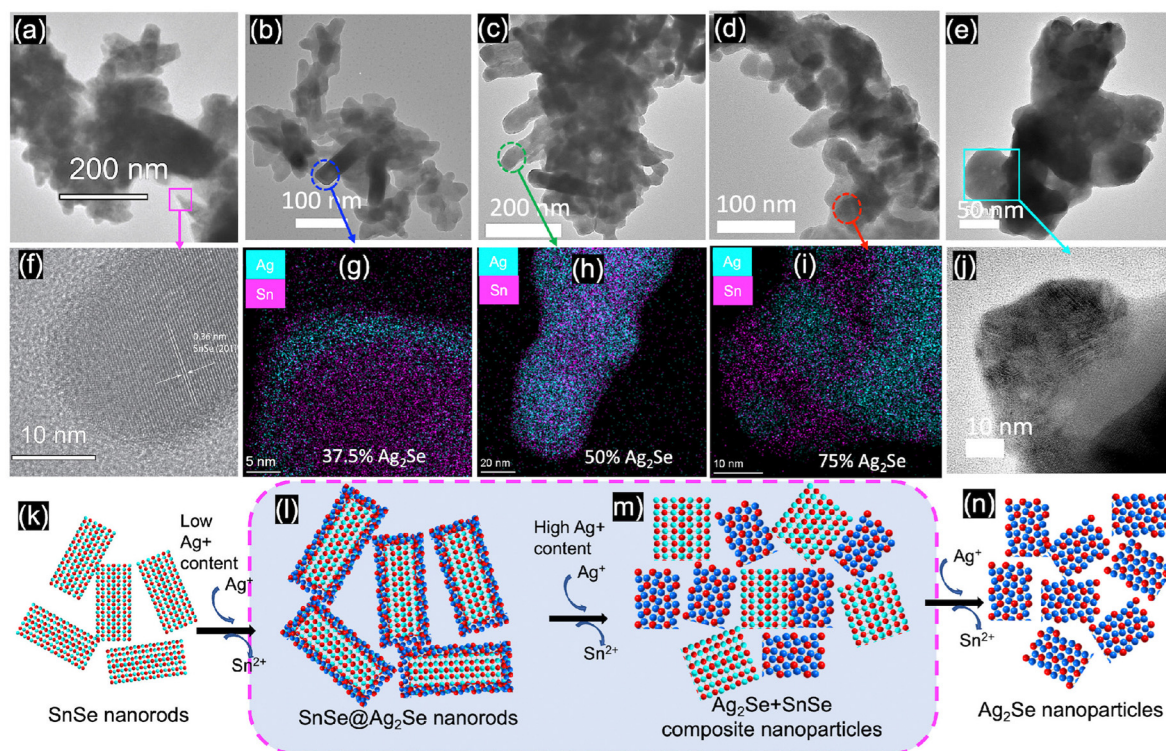
amounts of Ag<sup>+</sup> solutions prepared by dissolution of AgNO<sub>3</sub> crystals in methanol. The mixture underwent ultrasonic processing for 4 hours followed by 2 days of magnetic stirring at room temperature. The resulting cation exchange reaction can be described using the following chemical eqn (1).



The nominal composition of the synthesized Ag<sub>2</sub>Se/SnSe composites was controlled by adjusting the Ag/Sn molar ratio in the starting mixture (Table 1).

Fig. 1(a) and (b) show the XRD patterns of the Ag<sub>2</sub>Se/SnSe nanocomposites synthesized through cation exchange reactions. For nanocomposites with a low Ag/Sn ratio (Ag/Sn ≤ 0.4/1), only diffraction peaks from the SnSe phase are observed on the XRD patterns. This suggests either the formation of interstitial solid solution through the incorporation of Ag<sup>+</sup> ions within the crystal lattice of SnSe nanorods, or that the Ag<sub>2</sub>Se phase formed in these samples are likely amorphous with no





**Fig. 2** STEM images and the distribution of Ag and Sn in selected  $\text{Ag}_2\text{Se}/\text{SnSe}$  samples: (a) and (f) SnSe precursor; (b) and (g) 37.5%  $\text{Ag}_2\text{Se}$  (A-37); (c) and (h) 50%  $\text{Ag}_2\text{Se}$  (A-50); (d) and (i) 75%  $\text{Ag}_2\text{Se}$  (A-75); and (e) and (j)  $\text{Ag}_2\text{Se}$  nanocrystals. (k)–(n) Schematic illustration of the chemical transformation from the SnSe nanorod precursor to the  $\text{Ag}_2\text{Se}$  nanoparticles.

**Table 1** Nominal fraction of  $\text{Ag}_2\text{Se}$  and SnSe phases in SnSe/ $\text{Ag}_2\text{Se}$  nanocomposites

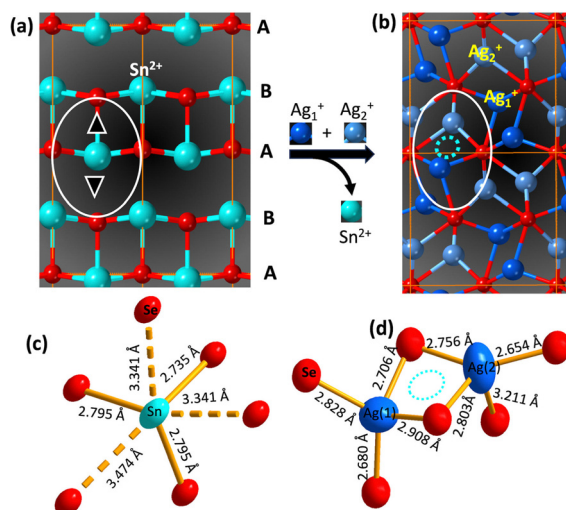
Sample code	Ag/Sn ratio	$\text{Ag}_2\text{Se}$ mole%	SnSe mole%
SnSe	0	0	100
A-10	0.2/1	10	90
A-18	0.375/1	18.75	81.25
A-20	0.4/1	20	80
A-37	0.75/1	37.5	62.5
A-50	1/1	50	50
A-75	1.5/1	75	25
A-100	2/1	100	0

long-range ordered structure. Upon increasing the Ag/Sn ratio to above 0.75/1 (62.5% SnSe/37.5%  $\text{Ag}_2\text{Se}$ ), diffraction peaks associated with the  $\text{Ag}_2\text{Se}$  phase can be observed on the XRD patterns along with peaks from the SnSe phase. Interestingly, the relative intensity of diffraction peaks from the SnSe phase decreases with the increasing Ag/Sn ratio whereas that of diffraction peaks from the  $\text{Ag}_2\text{Se}$  phase increases suggesting a gradual chemical transformation of the SnSe nanocrystals into  $\text{Ag}_2\text{Se}$ . Astonishingly, only diffraction peaks corresponding to the  $\text{Ag}_2\text{Se}$  phase are observed on the XRD patterns of nanocomposites with the Ag/Sn ratio exceeding 1.5/1 (25% SnSe/75%  $\text{Ag}_2\text{Se}$ ). This suggests that the exchange of a large fraction of  $\text{Sn}^{2+}$  ions within the SnSe nanocrystals by  $\text{Ag}^+$  ions resulted in the destruction of the original crystal into amorphous smaller particles.

The diffraction peaks from the experimental XRD patterns of samples with 75%  $\text{Ag}_2\text{Se}$  and 100%  $\text{Ag}_2\text{Se}$  can be indexed to the orthorhombic crystal structure ( $P2_12_12_1$  (#19))  $a = 4.336 \text{ \AA}$ ,  $b = 7.0774 \text{ \AA}$ , and  $c = 7.774 \text{ \AA}$  of  $\text{Ag}_2\text{Se}$ . This indicates a successful chemical transformation from the orthorhombic SnSe ( $Pnma$  (#62)) to the lower symmetry orthorhombic structure of  $\text{Ag}_2\text{Se}$  ( $P2_12_12_1$  (#19)) (Fig. 3). A careful comparison of both crystal structures reveals that the cation exchange reaction results in a change in the coordination geometry of the exchanged ions. Indeed, the removal of one  $\text{Sn}^{2+}$  ion from the severely distorted  $[1 + 2 + 2 + 1]$  octahedral coordination (Fig. 3d) in the SnSe structure is electronically compensated by two  $\text{Ag}^+$  ions, which are accommodated within the two tetrahedral voids ( $\text{Ag}(1)$  and  $\text{Ag}(2)$ ), adjacent to the octahedral void, in the deformed hexagonal close packed (hcp)  $\text{Se}^{2-}$  sublattice of SnSe (Fig. 3a and b). This substitution reaction results in the formation, by the  $\text{Ag}^+$  ions, of edge-sharing tetrahedral clusters (Fig. 3d) in the orthorhombic structure of  $\text{Ag}_2\text{Se}$ . Interestingly, the chemical transformation from SnSe to  $\text{Ag}_2\text{Se}$  retains the hcp  $\text{Se}^{2-}$  sublattice but is accompanied by an  $\sim 12\%$  increase in the unit cell volume from  $212 \text{ \AA}^3$  (SnSe) to  $238 \text{ \AA}^3$  ( $\text{Ag}_2\text{Se}$ ). This implies that the cation exchange between  $\text{Sn}^{2+}$  and  $\text{Ag}^+$  requires a local expansion of the  $\text{Se}^{2-}$  framework in SnSe. Such expansion of the  $\text{Se}^{2-}$  hcp sublattice is likely necessary not only to facilitate the diffusion of  $\text{Ag}^+$  ions within the tetrahedral voids, but also to accommodate the repulsion between  $\text{Ag}^+(1)$  and  $\text{Ag}^+(2)$  ions in the edge-sharing tetrahedral cluster. Interestingly, the lattice







**Fig. 3** Schematic illustration of the atomic-scale cation exchange process leading to the transformation from SnSe to Ag<sub>2</sub>Se orthorhombic structure. (a) Crystal structure of SnSe highlighting the hcp Se<sup>2−</sup> sublattice with Sn<sup>2+</sup> in octahedral void; the two empty tetrahedral voids (triangles) adjacent to the octahedrally coordinated Sn<sup>2+</sup> are occupied by Ag<sup>+</sup> ions in the crystal structure of Ag<sub>2</sub>Se (b). (c) distorted octahedral coordination geometry of Sn<sup>2+</sup>; (d) edge-sharing tetrahedral coordination of Ag<sup>+</sup> ions in Ag<sub>2</sub>Se.

**Table 2** Lattice parameters of Ag<sub>2</sub>Se in selected SnSe/Ag<sub>2</sub>Se nanocomposites

Lattice parameters	Ag <sub>2</sub> Se (standard pattern)	A-75	A-100
<i>a</i> (Å)	4.336	4.338	4.331
<i>b</i> (Å)	7.070	7.074	7.064
<i>c</i> (Å)	7.774	7.789	7.777
<i>V</i> (Å <sup>3</sup> )	238.32	239.01	237.95

parameters (Table 2 and Table S2, ESI†) of the Ag<sub>2</sub>Se phase in the SnSe/Ag<sub>2</sub>Se composites with 75% Ag<sub>2</sub>Se (A-75) as well as in the fully transformed sample with 100% Ag<sub>2</sub>Se (A-100) are highly consistent with the reported values for pristine Ag<sub>2</sub>Se demonstrating the effectiveness of the cation exchange process for the synthesis of the highly crystalline Ag<sub>2</sub>Se phase at 300 K.

The chemical composition of the SnSe precursor and the synthesized Ag<sub>2</sub>Se/SnSe nanocomposites was determined using EDX analysis. Table S1 and Fig. S1(a)–(c) (ESI†) show the nominal atomic percentages (at%) of Se, Sn, and Ag as well as the experimentally determined values for all samples. The Ag content in all samples is very close to the nominal value. Interestingly, the experimental atomic percentages of Se are always higher than the theoretical value especially in the nanocomposites with high Ag/Sn molar ratios, while the experimental values for the Sn element are always lower than the theoretical value. This result suggests the possibility of partial hydrolysis and/or oxidation of Sn<sup>2+</sup> upon extraction from the SnSe crystal lattice.

### Driving force for cation exchange

Given the large volume change ( $\Delta V/V \sim 12\%$ ) of the Se<sup>2−</sup> hcp sublattice and the change in the cation coordination, it is

remarkable that the chemical transformation from SnSe to Ag<sub>2</sub>Se occurs at room temperature, which prompts the question, what is the driving force for the cation exchange process? Compared to the solid state synthesis of Ag<sub>2</sub>Se from the elements, which typically requires high temperatures and long annealing, the relatively fast cation exchange at room temperature between SnSe and Ag<sup>+</sup> ion in methanol suggests that the formation of Ag<sub>2</sub>Se is thermodynamically favorable under these conditions. This implies a relatively low activation energy for the removal of Sn<sup>2+</sup> from the distorted octahedral position within the hcp Se<sup>2−</sup> sublattice of SnSe and the incorporation of Ag<sup>+</sup> ions at the tetrahedral positions. Indeed, a close examination of the cation coordination geometry in SnSe and Ag<sub>2</sub>Se revealed a severe distortion of Sn–Se bonding (3 short and 3 long) within the octahedral coordination with bond distances ranging from 2.735 Å to 3.474 Å ( $\Delta x/x \sim 37\%$ ), whereas the tetrahedral coordination of Ag<sup>+</sup> ions in Ag<sub>2</sub>Se shows weaker distortions,  $\Delta x/x \sim 21\%$  for Ag(2) and  $\Delta x/x \sim 8.5\%$  for Ag(1). Such severe distortion of the Sn<sup>2+</sup> coordination in SnSe is a manifestation of the structural instability of SnSe, presumably due to the relatively small size of the Sn<sup>2+</sup> ion (ionic radius (IR): 81 pm<sup>33</sup> to 93 pm<sup>34</sup> for CN = 6) for the large octahedral void. Conversely, the weaker distortion of the Ag<sup>+</sup> (IR = 100 pm<sup>35,36</sup> for CN = 4) tetrahedral coordination in Ag<sub>2</sub>Se points to a relatively more stable structure compared to SnSe.

The relative stability of Ag<sub>2</sub>Se over SnSe can also be assessed by comparing their solubility product constant ( $K_{sp}$ ) in methanol/water. For instance, the  $K_{sp}$  of Ag<sub>2</sub>Se in water ranges from  $2 \times 10^{-64}$  to  $3 \times 10^{-54}$ , which is much lower than that of SnSe ( $5 \times 10^{-34}$ )<sup>13,37,38</sup> indicating that Ag<sub>2</sub>Se is more stable in water than SnSe. Following a similar trend, one can anticipate that Ag<sub>2</sub>Se would also possess a relatively lower  $K_{sp}$  than SnSe in polar solvents like methanol, indicating that Ag<sub>2</sub>Se is more thermodynamically stable than SnSe in the cation exchange environment, which is consistent with the observed relatively easy chemical transformation of SnSe to Ag<sub>2</sub>Se through exchange of Sn<sup>2+</sup> by Ag<sup>+</sup> at 300 K.

The relatively facile transformation of SnSe to Ag<sub>2</sub>Se can also be rationalized by considering the close spatial proximity of the octahedral (O) and tetrahedral (T<sup>+</sup>, T<sup>−</sup>) voids within the hcp Se<sup>2−</sup> sublattice in the SnSe structure as well as their 1:2 ratio that enables the accommodation of both Ag<sup>+</sup> ions used to electronically compensate for the outgoing Sn<sup>2+</sup> ions. Upon removal of Sn<sup>2+</sup> ions, the large octahedral voids left behind likely serve as the point of entry of Ag<sup>+</sup> (IR  $\sim 100$  pm<sup>35,36</sup>) ions, which are slightly larger ( $(\Delta(IR))/(IR) \sim 23\%$ ) than the outgoing Sn<sup>2+</sup> ion, but still quite small for the large octahedral coordination within the Se<sup>2−</sup> hcp sublattice in SnSe. This results in an unstable intermediate structure as the octahedral voids are too large for the Ag<sup>+</sup> ion. Such structural instability is likely resolved by the subsequent diffusion of Ag<sup>+</sup> ions into the much smaller tetrahedral voids that are located  $\sim 4$  Å to  $\sim 5$  Å away from the octahedral site (Fig. 3b).

The driving force for the cation exchange reaction was also likely improved by the choice of a solvent that can both dissolve AgNO<sub>3</sub> solids to provide free Ag<sup>+</sup> cations and promote the



cation exchange process by extracting  $\text{Sn}^{2+}$  from the original SnSe lattice. In this work, we leverage the Pearson's hard and soft acids and bases (HSAB) theory, which can be used to establish a qualitative affinity between metal ions and ligands/solvents based on the concept of absolute hardness ( $\eta$ ),<sup>13</sup> for the selection of methanol as the suitable solvent for the cation exchange process. According to the HSAB theory, a hard base tends to bind with a hard acid while a soft base prefers to bind with a soft acid.<sup>13,39,40</sup> Therefore, one can anticipate a successful cation exchange reaction between SnSe and  $\text{Ag}^+$  ions using a hard base such as methanol or water as the solvent, since the guest cation ( $\text{Ag}^+$ ) is softer ( $\eta = 6.96$ ) than the host cation  $\text{Sn}^{2+}$  ( $\eta = 7.94$ ). Because of the relatively easy hydrolysis of  $\text{Sn}^{2+}$  ions in water, methanol was selected as the most appropriate solvent for the ion exchange process. Under this condition, the hard base molecule (methanol) tends to bind with the hard acid cation ( $\text{Sn}^{2+}$ ), which should facilitate its removal from the SnSe nanorods, creating empty cation sites (octahedral voids) for  $\text{Ag}^+$  ions to hop into. This is presumably the mechanism crucial for the successful fabrication of SnSe/ $\text{Ag}_2\text{Se}$  composites through the cation exchange process at 300 K.

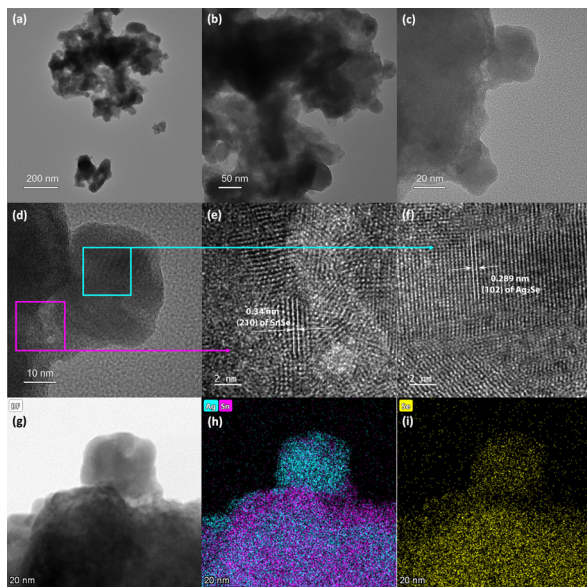
### Mechanism of cation exchange

To probe the mechanism of the chemical transformation of SnSe nanocrystals into  $\text{Ag}_2\text{Se}$ , the evolution of the composition and morphology of nanocrystals in selected  $\text{Ag}_2\text{Se}/\text{SnSe}$  nanocomposites was investigated by transmission electron microscopy (TEM) and scanning transmission electron microscopy (STEM). A careful examination of the dark field (DF) and bright field (BF) images of various samples and the associated EDX elemental mappings revealed that both the nanocrystal morphology and the distribution of Ag and Sn within the composites change with the increasing Ag/Sn ratio (Fig. 2(b) to (e)). For instance, the TEM BF image of sample A-37 (37.5%  $\text{Ag}_2\text{Se}$ ) showed the formation of a core-shell structure (Fig. 2(b)). The corresponding elemental mappings revealed the concentration of Ag atoms on the shell part of the nanorod while Sn atoms are mainly located within the core (Fig. 2(g) and (l)). This suggests that for samples with low Ag/Sn ratio, the cation exchange proceeds from the extraction of  $\text{Sn}^{2+}$  on the surface of the initial SnSe nanorod to form a solid solution or a thin amorphous  $\text{Ag}_2\text{Se}$  shell, which could explain the lack of diffraction peaks for the  $\text{Ag}_2\text{Se}$  phase on the XRD patterns (Fig. 1a). Upon increasing the Ag/Sn ratio to 1/1 (50%  $\text{Ag}_2\text{Se}$ ), the TEM-BF image (Fig. 2(c)) along with the corresponding elemental mapping (Fig. 2(h)) of the resulting sample (A-50) revealed a random distribution of nano-domains of the  $\text{Ag}_2\text{Se}$  and SnSe phases throughout the nanorod. Such a change in the distribution of Ag and Sn indicates that as the Ag/Sn molar ratio increases, the initial core-shell heterostructure becomes unstable, presumably due to the large lattice mismatch at the  $\text{Ag}_2\text{Se}/\text{SnSe}$  interface and the large change in the unit cell volume ( $\Delta V/V \sim 12\%$ ) arising from the expansion of the  $\text{Se}^{2-}$  hcp sublattice (Fig. 3). During the exchange process, one  $\text{Sn}^{2+}$  must be replaced by two  $\text{Ag}^+$  ions to maintain the charge balance. Considering that the ionic radius of  $\text{Sn}^{2+}$  ( $\sim 81 \text{ pm}^{33}$  to  $93 \text{ pm}^{34}$  for CN = 6) is slightly

smaller than that of  $\text{Ag}^+$  ( $100 \text{ pm}^{35,36}$  for CN = 4), the insertion of two  $\text{Ag}^+$  ions in replacement for one  $\text{Sn}^{2+}$  ion is anticipated to introduce intense internal stress to the original hcp  $\text{Se}^{2-}$  sublattice in the SnSe structure. The internal stress is alleviated in the  $\text{Ag}_2\text{Se}$  structure by the expansion of the  $\text{Se}^{2-}$  framework leading to the formation of incoherent SnSe/ $\text{Ag}_2\text{Se}$  interfaces. Such a high-energy interface likely facilitates the separation of the thin  $\text{Ag}_2\text{Se}$  shell from the SnSe nanorod enabling additional diffusion of  $\text{Ag}^+$  ions into the remaining SnSe crystal. This process presumably continues until depletion of  $\text{Ag}^+$  ions from the solution, which results in the fragmentation of the initial SnSe nanorod into small clusters of nanocrystals with SnSe and  $\text{Ag}_2\text{Se}$  compositions as observed in Fig. 2(c) and (h). As the  $\text{Ag}_2\text{Se}$  molar percentage further increases to 75% (A-75), the initial nanorods are broken into multiple smaller Ag-rich nanoparticles connected by Sn-rich bridges, as shown in Fig. 2(d) and (i), and illustrated in Fig. 2(m). This morphological evolution from a core-shell structure to random distribution of  $\text{Ag}_2\text{Se}$  and SnSe nanodomains and eventually to complete Ag-rich nanoparticles connected by Sn-rich bridges as schematically illustrated in Fig. 2(k) to (n) can be rationalized by considering that the incoherent SnSe/ $\text{Ag}_2\text{Se}$  high energy interfaces lead to the separation of  $\text{Ag}_2\text{Se}$  domains from the SnSe nanorod, which in turn facilitates the diffusion of  $\text{Ag}^+$  ions into the freshly exposed surface of the SnSe nanorod. The continuation of this cation-exchange process results in random distribution of  $\text{Ag}_2\text{Se}$  and SnSe nano domains in the remaining nanorods for equimolar ratio and finally to the formation of small nanocrystals with  $\text{Ag}_2\text{Se}$  composition and structure (Fig. 2(e), (j) and (n)).

To elucidate the crystallinity near the interface between the SnSe and  $\text{Ag}_2\text{Se}$  phases, additional high-resolution TEM (HRTEM) images were acquired for sample A-75 (Fig. 4). It can be seen from Fig. 4(a) and (b) that the sample A-75 features agglomerated nano-sized particles with irregular shapes. The higher magnification TEM image (Fig. 4(c)) shows multiple small nanoparticles on the surface of the larger piece. The corresponding STEM-BF image (Fig. 4(g)) and elemental mapping (Fig. 4(h) and (i)) reveal that the small nanoparticles are  $\text{Ag}_2\text{Se}$  phase, while the larger piece is mainly the original SnSe phase with some degree of intermixing with  $\text{Ag}_2\text{Se}$ . The high resolution TEM (HRTEM) image of the interface between the SnSe-rich particles and the  $\text{Ag}_2\text{Se}$  nanoparticles (Fig. 4(d)) depicts the crystallinity of the SnSe phase (large piece) as well as that of the  $\text{Ag}_2\text{Se}$  phase (small particle). As shown in Fig. 4(e), the local area at the interface between SnSe and  $\text{Ag}_2\text{Se}$  shows a clear lattice pattern, on the SnSe side, with an interplanar spacing of 0.34 nm corresponding to the (210) plane of the SnSe crystal. The interface area shows no long-range ordered structure, and it is difficult to acquire the lattice information. Such a local amorphous region may be caused by the process of  $\text{Sn}^{2+}$  extraction and  $\text{Ag}^+$  insertion at the interface. For each  $\text{Sn}^{2+}$  extracted, there are 2  $\text{Ag}^+$  ions that must enter the original  $\text{Se}^{2-}$  sublattice to maintain the charge balance. This non-isovalent exchange process likely causes internal stress to the original lattice leading to severe deformation. As a result, the original

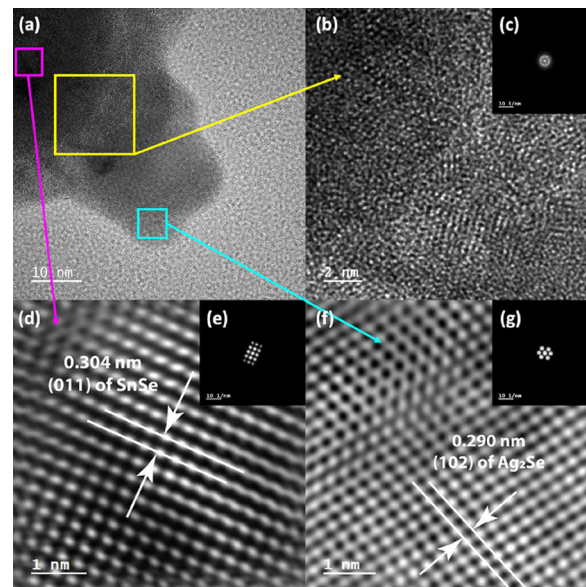




**Fig. 4** (a)–(c) TEM images of the nanocomposites A-75 under various magnifications, showing the agglomeration feature and interface between the small  $\text{Ag}_2\text{Se}$  nanoparticles and large piece SnSe matrix; (d) HRTEM image of the interface between SnSe precursor and  $\text{Ag}_2\text{Se}$  nanoparticles. The two highlighted areas show the lattice feature of the interface area (e) and the small  $\text{Ag}_2\text{Se}$  nanoparticle area (f); (e) zoomed-in HRTEM image of the interface area, showing the lattice feature of SnSe (210) plane at the left side of the interface; (f) zoomed-in HRTEM image of the small nanoparticles, showing the lattice feature of the  $\text{Ag}_2\text{Se}$  (102) plane; (g) STEM-BF image of the same SnSe/ $\text{Ag}_2\text{Se}$  interface shown in (d) with corresponding elemental mapping results ((h), and (i)), confirming that the large matrix is SnSe-rich while the small nanoparticle is  $\text{Ag}_2\text{Se}$ -rich.

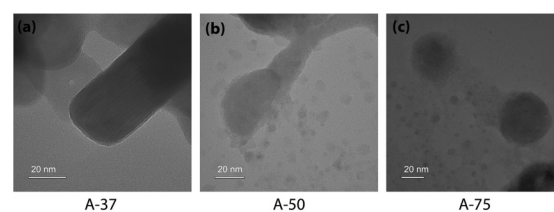
long-range ordered SnSe structure is broken to form the Sn–Ag intermixing interface with poor crystallinity. Interestingly, the HRTEM image of the small  $\text{Ag}_2\text{Se}$  nanoparticle region (Fig. 4(f)) revealed a clear lattice pattern with an interplanar spacing of 0.289 nm corresponding to the (102) plane of the  $\text{Ag}_2\text{Se}$  phase. This suggests that a long-range ordered structure is formed upon complete exchange of  $\text{Sn}^{2+}$  by  $\text{Ag}^+$  from the original SnSe crystal.

Further analysis of the SnSe/ $\text{Ag}_2\text{Se}$  interface at a different area of the sample A-75 (Fig. 5(a)) also revealed a well-ordered structure at both the SnSe-rich and  $\text{Ag}_2\text{Se}$ -rich ends whereas the phase boundary showed a disordered structure. As can be seen from Fig. 5(d) and (f), the HRTEM and the corresponding fast-Fourier transformed (FFT) pattern for the SnSe-rich (Fig. 5(e)) and  $\text{Ag}_2\text{Se}$ -rich (Fig. 5(g)) areas show a highly periodic lattice structure, whereas the phase boundary (interface) looks amorphous without any long-range ordered structure (Fig. 5(b)) and the FFT pattern only shows circular rings feature (Fig. 5(c)), instead of patterned points feature, which also indicates that the interface area contains no long-range ordered structure. Fig. 5(d) shows the crystal lattice feature on the SnSe-rich side with an interplanar distance of 0.304 nm, which corresponds to the (011) plane of SnSe. The lattice feature on the small  $\text{Ag}_2\text{Se}$  nanoparticle side (Fig. 5(f)) with an interplanar distance of 0.290 nm can be ascribed to the (102) plane of  $\text{Ag}_2\text{Se}$ .



**Fig. 5** (a) HRTEM image of the lower side of the small nanoparticle shown in Fig. 4(c). The 3 highlighted areas were selected to explore the lattice features at the interface (b), on the SnSe side of the interface (d), and on the  $\text{Ag}_2\text{Se}$  side of the interface (f); (b) and (c) zoomed-in HRTEM image of the two-phase interface with the corresponding FFT pattern, showing an amorphous (no long-range-ordered) structure; (d) and (e) zoomed-in HRTEM image of the SnSe-rich side of the interface with the corresponding FFT pattern, showing the (011) plane of SnSe; (f) and (g) zoomed-in HRTEM image of the  $\text{Ag}_2\text{Se}$ -rich side of the interface with the corresponding FFT pattern, showing the (102) plane of  $\text{Ag}_2\text{Se}$ .

Another noticeable feature of the synthesized nanocomposites is the morphology change of the SnSe nanorod along with various compositions. Fig. 6 displays the morphology of nanoparticles in A-37 (a), A-50 (b), and A-75 (c) SnSe/ $\text{Ag}_2\text{Se}$  nanocomposites. The original shape of the SnSe nanorod precursor is preserved in sample A-37 (Fig. 6(a)). However, when it comes to sample A-50 (Fig. 6(b)), the nanorod begins to break into smaller nanoparticles. Furthermore, the nanorod almost completely breaks into smaller nanoparticles in samples A-75 (Fig. 6(c)). There is a clear trend that with a higher  $\text{Ag}_2\text{Se}$  percentage in the nanocomposites (from 37.5% to 75%), the degree of cation exchange is higher, and as a result, the disintegration of the original SnSe nanorod into smaller particles becomes more prominent. This can be rationalized by the local expansion of the  $\text{Se}^{2-}$  sublattice that results from the replacement of a relatively smaller  $\text{Sn}^{2+}$  (81–93 pm) ion by two



**Fig. 6** HRTEM images of samples A-37 (a), A-50 (b), and A-75 (c), showing the morphology change from nanorods to smaller nanoparticles due to the cracking of the initial nanorod.





relatively larger  $\text{Ag}^+$  (100 pm) ions during the cation exchange process to maintain the charge neutrality. When the degree of cation exchange is low (for example, in A-37), this tension stress is not strong enough to undermine the integrity of the original nanorod. However, as the degree of cation exchange increases (for example, A-50 and A-75), the build-up tension stress can reach a considerable level, which can damage the original lattice and break the nanorod into smaller nanoparticles.

To further probe the mechanism of the ion exchange process, we have examined the oxidation state of Sn, Ag and Se atoms in various samples using X-ray photoelectron spectroscopy (XPS). Fig. 7(a) shows the Sn-3d XPS spectra of the SnSe precursor and various SnSe/ $\text{Ag}_2\text{Se}$  nanocomposites. For the SnSe precursor, the Sn-3d spectrum shows two peaks at 485.4 eV and 493.8 eV, corresponding to the binding energy of Sn-3d<sub>5/2</sub> and Sn-3d<sub>3/2</sub> states in divalent tin ( $\text{Sn}^{2+}$ ), respectively.<sup>41</sup> Upon partial exchange of  $\text{Sn}^{2+}$  by  $\text{Ag}^+$  in SnSe nanorods, broadening of both Sn-3d<sub>5/2</sub> and Sn-3d<sub>3/2</sub> peaks with a shift of the central position to higher binding energy is observed on the Sn-3d XPS spectra of various composites. A careful analysis of the peak shape indicates the appearance of a shoulder peak at 487.2 eV and 495.6 eV, which are respectively associated with the Sn-3d<sub>5/2</sub> and Sn-3d<sub>3/2</sub> states in tetravalent tin ( $\text{Sn}^{4+}$ ) in  $\text{SnO}_2$ .<sup>42,43</sup> This result suggests the oxidation of a small fraction of  $\text{Sn}^{2+}$  ions upon removal from the SnSe crystal lattice.<sup>44,45</sup>

Interestingly, the intensity of the peaks at 485.4 eV and 493.8 eV decreases with increasing  $\text{Ag}_2\text{Se}$  content whereas that of the shoulder peaks at 487.2 eV and 495.6 eV increases. For composition with 75%  $\text{Ag}_2\text{Se}$  (sample A-75), the peaks at 485.4 eV and 493.8 eV completely vanished with only the higher energy peaks remaining. A similar Sn-3d XPS spectrum is observed for the sample A-100, which nominally corresponds to the sample with complete cation exchange (*i.e.* 100%  $\text{Ag}_2\text{Se}$ ). The observed data further imply that during the cation exchange process, a fraction of  $\text{Sn}^{2+}$  is oxidized to  $\text{Sn}^{4+}$  presumably from surface oxidation of the SnSe nanorods, whereas the remaining  $\text{Sn}^{2+}$  ions are gradually replaced by  $\text{Ag}^+$  demonstrating the completeness of the cation exchange process. Indeed, the Ag-3d XPS spectra (Fig. 7(b)) of various composites (A-10 to A-100) revealed the presence of two isolated peaks at 373.6 eV and 367.6 eV,

which correspond, respectively, to the Ag-3d<sub>3/2</sub> and Ag-3d<sub>5/2</sub> with a spin-orbit splitting of 6.0 eV, suggesting that the valence of Ag was +1.<sup>46</sup> As the  $\text{Ag}_2\text{Se}$  concentration increases, the peak positions of Ag-3d<sub>3/2</sub> and Ag-3d<sub>5/2</sub> states remain relatively unchanged. For the sample A-100, both isolated peaks appear at 374.2 eV (Ag-3d<sub>3/2</sub>) and 368.2 eV (Ag-3d<sub>5/2</sub>), which is still within the range reported for monovalent (+1) Ag in silver chalcogenides.<sup>47</sup> Fig. 7(c) depicts the Se-3d XPS spectra of the SnSe precursor and all SnSe/ $\text{Ag}_2\text{Se}$  nanocomposites. One can notice the relative stability in both the SnSe precursors and the nanocomposites of the peak at 53.6 eV that is associated with the Se 3d<sub>5/2</sub> state implying the stability of the  $\text{Se}^{2-}$  anion during the cation exchange process. As the fraction of  $\text{Ag}_2\text{Se}$  increases to 100%, two strong peaks are observed at 54.4 eV and 53.6 eV, which can be attributed to Se-3d<sub>3/2</sub> and Se-3d<sub>5/2</sub> states, indicating the presence of  $\text{Se}^{2-}$  anions in the fully transformed  $\text{Ag}_2\text{Se}$  sample.<sup>48</sup>

### Optical properties

To assess the effect of  $\text{Ag}^+$  to  $\text{Sn}^{2+}$  cation exchange on the optical band gap, we have conducted diffuse reflectance measurements at 300 K (Fig. S2, ESI†) on the SnSe precursor as well as on various SnSe/ $\text{Ag}_2\text{Se}$  nanocomposites using a UV-vis-NIR spectrometer. Direct bandgap values extracted from the Tauc plots of the diffuse reflectance spectra (DRS) are displayed in Fig. 8 and the change in the band gap with increasing  $\text{Ag}_2\text{Se}$  content is shown in Fig. S3 (ESI†). The SnSe precursor showed a direct bandgap value of 1.02 eV, which is in agreement with previously reported values.<sup>49</sup> The partial exchange of  $\text{Sn}^{2+}$  by  $\text{Ag}^+$  ions results in an initial marginal increase of the optical band gap of the resulting nanocomposites. For instance, the band gap value for the composite with 10%  $\text{Ag}_2\text{Se}$  (sample A-10) is 1.09 eV, which corresponds to a 7% increase compared to that of the SnSe nanorod precursor. A further increase in the  $\text{Ag}_2\text{Se}$  content led to a marginal drop in the band gap to 1.07 eV for the sample A-20 and to 1.06 eV for sample A-37, which are still larger than the band gaps of the SnSe precursor (Fig. 8(a)). This trend is quite unexpected since the band gap of pristine

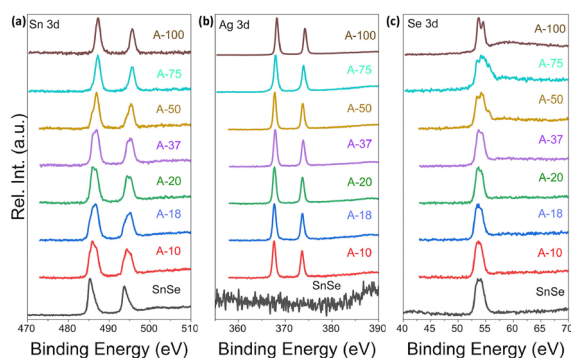


Fig. 7 XPS spectra of (a) Sn-3d, (b) Ag-3d, and (c) Se-3d shell electrons in SnSe and selected  $\text{Ag}_2\text{Se}/\text{SnSe}$  nanocomposites.

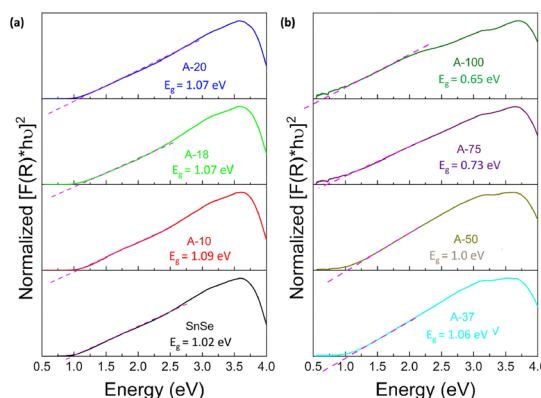


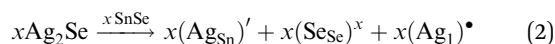
Fig. 8 Tauc plots of the UV-vis-NIR diffuse reflectance spectra showing the optical band gap of  $\text{Ag}_2\text{Se}/\text{SnSe}$  nanocomposites with (a) low  $\text{Ag}_2\text{Se}$  content; and (b) high  $\text{Ag}_2\text{Se}$  content.



**Table 3** Lattice parameters of SnSe in selected SnSe/Ag<sub>2</sub>Se nano-composites

Lattice parameters	SnSe	A-10	A-18	A-20
<i>a</i> (Å)	11.518	11.522	11.522	11.524
<i>b</i> (Å)	4.163	4.148	4.146	4.142
<i>c</i> (Å)	4.409	4.440	4.453	4.454
<i>V</i> (Å <sup>3</sup> )	211.415	212.216	212.751	212.632

Ag<sub>2</sub>Se (~0.1 eV) is considerably lower than that of SnSe, which is in the range 1.0–1.4 eV.<sup>50–55</sup> One plausible explanation for this trend is that the initial substitutions of Sn<sup>2+</sup> by two Ag<sup>+</sup> ions lead to the formation of Ag<sup>+</sup> interstitial solid-solution, Sn<sub>1-x</sub>Ag<sub>2x</sub>Se, with an SnSe-structure type according to the defect eqn (2).



This conjecture is supported by the presence of diffraction peaks from only the SnSe phase on the XRD patterns of composites with low Ag<sub>2</sub>Se content (A-10 to A-20). Furthermore, a careful analysis of the lattice parameters (Table 3 and Table S2, ESI†) of SnSe in selected samples with low Ag content shows an initial small expansion of the unit cell volume, which is consistent with the formation of interstitial Ag<sup>+</sup> ions in the SnSe crystal lattice. Within this defected SnSe structure picture, the incorporation of Ag<sup>+</sup> into the SnSe lattice could result in the formation of increasingly ionic bonding, M–Se (M = Sn/Ag) owing to the slightly higher electronegativity value of Ag (1.9) compared to Sn (1.8).<sup>56,57</sup> This could explain the observed increase in the energy band gap of composites with 10% and 20% Ag<sub>2</sub>Se. The Sn<sub>1-x</sub>Ag<sub>2x</sub>Se solid solution likely collapses with the substitution of a large fraction of Sn<sup>2+</sup> ions (A-37 to A-75) resulting in a mixture of Sn-rich (SnSe) and Ag-rich (Ag<sub>2</sub>Se) phases (Fig. 2(g), (h), and (i)). This analysis is consistent with the observed rapid drop in the band gap value to 1.00 eV for sample A-50 and to 0.73 eV for sample A-75 (Fig. 8(b) and Fig. S3, ESI†). Interestingly, the band gap value of 0.65 eV observed for sample A-100 (100% Ag<sub>2</sub>Se) is significantly larger than the reported value around 0.1 eV.<sup>50,58</sup> We attribute such a large increase in the band gap to the nanometer scale particle size of Ag<sub>2</sub>Se obtained through the cation exchange reaction (Fig. 2e).

## Conclusions

In summary, we have developed a novel cation exchange method for a facile synthesis of Ag<sub>2</sub>Se/SnSe nanocomposites at room temperature. Using methanol as a solvent, we demonstrate the formation at 300 K of a highly crystalline Ag<sub>2</sub>Se phase through chemical transformation of SnSe nanorods with cation exchange of one Sn<sup>2+</sup> ion by two Ag<sup>+</sup> ions. The fraction of Ag<sub>2</sub>Se in the composites can be tuned by adjusting the Ag/Sn ratio in the starting mixture. XRD patterns and STEM images of selected composites revealed four stages in the transformation of SnSe nanorods to Ag<sub>2</sub>Se nanoparticles. For samples with low

Ag content (A-10 to A-20), the ion exchange process results in the incorporation of Ag<sup>+</sup> ions into the SnSe structure leading to the formation of a partial solid-solution with SnSe-structure type. A further increase in the Ag content results in the formation of a SnSe/Ag<sub>2</sub>Se core/shell structure for the sample A-37 followed by a random distribution of Sn-rich and Ag-rich nanodomains in composite with 50% Ag<sub>2</sub>Se (A-50) and finally to the phase segregation between Ag<sub>2</sub>Se and SnSe domains (A-75). The ability to achieve complete transformation of SnSe nanorods into highly crystalline Ag<sub>2</sub>Se nanoparticles at 300 K was rationalized using a combination of factors such as the relatively superior structural stability in polar solvent of Ag<sub>2</sub>Se compared to SnSe, and the strong binding affinity of the methanol molecule to the Sn<sup>2+</sup> ion, which promotes its removal from octahedral voids in the SnSe crystal lattice paving the way for the incorporation of Ag<sup>+</sup> ions.

UV/vis diffuse reflectance measurement revealed an initial increase in the optical band gap for samples with low Ag content (A-10 to A-20), which is attributed to the formation of Ag<sup>+</sup> interstitial solid-solution, Sn<sub>1-x</sub>Ag<sub>2x</sub>Se, as well as the slightly enhanced (Sn/Ag)–Se ionic bonding. A further increase in the Ag<sub>2</sub>Se content results in a rapid drop in the band gap, which is consistent with the mixture between the wide band gap SnSe and the narrower band gap Ag<sub>2</sub>Se phase. This work demonstrates the potential of the gradual cation exchange process as a simple approach to engineer the optical properties of semiconductors, which can be leveraged for future material property engineering, especially in the fields of solar cells and optoelectronics.

## Experimental section

### Chemicals

Elemental selenium (Se) was purchased from Sigma-Aldrich. Sodium borohydride (NaBH<sub>4</sub>), tin(II) chloride dihydrate (SnCl<sub>2</sub>·2H<sub>2</sub>O), and silver nitrate (AgNO<sub>3</sub>) were purchased from Alfa Aesar. Sodium hydroxide (NaOH), methanol, and isopropyl alcohol (IPA) were purchased from Fisher Scientific. All reagents were used as received without additional purification.

### Synthesis of SnSe nanorods

The SnSe nanorods were synthesized using a modified surfactant-free route.<sup>32</sup> First, Se powder, NaBH<sub>4</sub>, and deionized (DI) water were added in sequence in an empty beaker. The mixture experienced a fierce reaction which is manifested by the release of bubbles. This reaction can be described using the following chemical eqn (3):

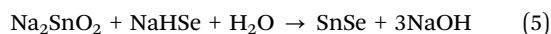


In another empty beaker, a certain amount of NaOH was dissolved in DI water followed by the addition of SnCl<sub>2</sub>·2H<sub>2</sub>O into the NaOH solution. This mixture initially yielded some white precipitates, which gradually redissolved. The reaction can be described using the following chemical eqn (4):





Next, the two solutions were mixed, which resulted in the precipitation of black SnSe nanoparticles. The reaction can be described using the following chemical eqn (5):



The as-synthesized product was washed in sequence and with centrifugation using diluted NaOH solution, DI water, and methanol. The purified SnSe product was dried in a vacuum at room temperature and stored for future use.

In a typical reaction towards the synthesis of 2 grams of SnSe precursor, 0.7989 g of Se was mixed with 0.7655 g NaBH<sub>4</sub> in a beaker. Then, 60 ml of DI water was added, and magnetic stirring was performed. The mixture immediately turned into black and finally became a transparent and colorless solution (NaHSe). In a separate beaker, 4.856 g of NaOH was dissolved in 50 ml DI water under magnetic stirring. Then 2.2831 g of SnCl<sub>2</sub>·2H<sub>2</sub>O powder was added into the NaOH solution to make Na<sub>2</sub>SnO<sub>2</sub>. When the solid fully dissolved and became transparent, the as-synthesized NaHSe solution was injected to the Na<sub>2</sub>SnO<sub>2</sub> solution. The mixture was kept under magnetic stirring for 2 hours at room temperature and then the synthesized SnSe precursor was collected by centrifugation. The product was washed several times using NaOH solution, DI water, and methanol and dried under vacuum.

### Synthesis of SnSe–Ag<sub>2</sub>Se composite nanocrystals

For cation exchange reactions, a certain amount of SnSe (~0.1 g) nanorods was added into a glass vial and dispersed in 4 ml of methanol. In a separate vial, the desired amount of AgNO<sub>3</sub> powder was dissolved in methanol. The AgNO<sub>3</sub> solution was subsequently added into the vial containing SnSe suspension and sealed. The mixture was placed under ultrasonic processing for 4 hours and underwent magnetic stirring for 2 days at room temperature. The final Ag<sub>2</sub>Se/SnSe nanocomposite was collected by centrifugation and washed with methanol several times. The final products were dried and kept under vacuum.

### Characterization

X-ray diffracted data were collected on a Rigaku Smart Lab, using monochromatic CuKα radiation (anode voltage 40 kV, emission current 44 mA) and a (D/teX 250) detector. XRD powder patterns were analyzed using PDXL software. Scanning electron microscopy (SEM) and energy dispersive X-Ray analyses (SEM-EDX) were carried out on a JEOL IT500 SEM. Further characterization on the nano-sized particles was performed using transmission electron microscopy (TEM), scanning transmission electron microscopy (STEM), STEM-EDX, and high-resolution TEM (HRTEM) with a ThermoFisher Scientific Talos F200X G2 S/TEM, operating with an accelerating voltage of 200 kV. For the TEM and HRTEM modes, the images were collected using a OneView camera, while for STEM and STEM-EDX, the images and elemental mappings were collected using 4 detectors: high-angle annular dark-field (HAADF), dark field-2 (DF2), dark field-4 (DF4), and bright field (BF). Samples for TEM/STEM experiments were prepared by dispersing powder

samples of the Ag<sub>2</sub>Se/SnSe nanocomposite in isopropyl alcohol (IPA) using ultrasonic processing followed by the deposition of droplets on copper grids with carbon films. The loaded grids were dried in a vacuum with mild heating to remove the remaining IPA. HRTEM images were analyzed using Digital Micrograph software and STEM-EDX results were processed using Velox software. X-ray photoelectron spectroscopy (XPS) data were collected on a Kratos Axis Ultra XPS using a mono Al source (Al Kα, 1486.6 eV). The collected data were analyzed using Casa XPS software. Finally, the optical data (reflection spectrum) were collected on a JASCO V-770 UV-vis-NIR spectrophotometer.

## Author contributions

The manuscript was written through contributions from all authors. All authors have given approval to the final version of the manuscript.

## Conflicts of interest

The authors declare no competing financial interest.

## Acknowledgements

We gratefully acknowledge the financial support from the Department of Energy, Office of Basic Energy Science under Award # DE-SC-0018941. This work made use of the TEM and XPS from the Michigan Materials Characterization Center, (MC)<sup>2</sup>, purchased with funds from the University of Michigan College of Engineering and the National Science Foundation Awards DMR-0420785.

## References

- 1 N. A. Moroz, A. Olvera, G. M. Willis and P. F. P. Poudeu, *Nanoscale*, 2015, **7**, 9452–9456.
- 2 A. A. Olvera, N. A. Moroz, P. Sahoo, P. Ren, T. P. Bailey, A. A. Page, C. Uher and P. F. P. Poudeu, *Energy Environ. Sci.*, 2017, **10**, 1668–1676.
- 3 L. Maddi, K. Vinukonda, T. R. Gurugubelli and R. Koutavarapu, *Electronics*, 2023, **12**, 1245.
- 4 S. K. Pandey, P. K. Mishra and D. Tiwary, *J. Environ. Chem. Eng.*, 2022, **10**, 107459.
- 5 M. Revathi, A. Pricilla Jeyakumari, R. Sridarane, M. Shkir, E. El Sayed Massoud and V. Sreedevi Gedi, *J. Cluster Sci.*, 2022, 1–12.
- 6 X. Song, X. Li, Z. Chen and Z. Wang, *Mater. Lett.*, 2020, **275**, 128109.
- 7 M. Wang, P. F. Fang, Y. Chen, X. Y. Leng, Y. Yan, S. B. Yang, P. Xu and C. Yan, *Adv. Funct. Mater.*, 2023, 2213902.
- 8 L. Zeng, H. Zhao, Y. Zhu, S. Chen, Y. Zhang, D. Wei, J. Sun and H. Fan, *J. Mater. Chem. B*, 2020, **8**, 4093–4105.
- 9 A. Olvera, T. P. Bailey, C. Uher and P. F. P. Poudeu, *J. Mater. Chem. A*, 2018, **6**, 6997–7004.



- 10 C. Zhou, Y. K. Lee, Y. Yu, S. Byun, Z.-Z. Luo, H. Lee, B. Ge, Y.-L. Lee, X. Chen and J. Y. Lee, *Nat. Mater.*, 2021, **20**, 1378–1384.
- 11 S. Chandra, U. Bhat, P. Dutta, A. Bhardwaj, R. Datta and K. Biswas, *Adv. Mater.*, 2022, **34**, 2203725.
- 12 S. Bano, S. I. Raj, A. Khalilullah, A. Jaiswal and I. Uddin, *J. Photochem. Photobiol., A*, 2021, **405**, 112925.
- 13 L. De Trizio and L. Manna, *Chem. Rev.*, 2016, **116**, 10852–10887.
- 14 T. Ling, M. Jaroniec and S. Z. Qiao, *Adv. Mater.*, 2020, **32**, 2001866.
- 15 B. J. Beberwyck, Y. Surendranath and A. P. Alivisatos, *J. Phys. Chem. C*, 2013, **117**, 19759–19770.
- 16 A. E. Powell, J. M. Hodges and R. E. Schaak, *J. Am. Chem. Soc.*, 2016, **138**, 471–474.
- 17 D. H. Ha, A. H. Caldwell, M. J. Ward, S. Honrao, K. Mathew, R. Hovden, M. K. A. Koker, D. A. Muller, R. G. Hennig and R. D. Robinson, *Nano Lett.*, 2014, **14**, 7090–7099.
- 18 E. A. Hernandez-Pagan, A. O'Hara, S. L. Arrowood, J. R. McBride, J. M. Rhodes, S. T. Pantelides and J. E. Macdonald, *Chem. Mater.*, 2018, **30**, 8843–8851.
- 19 D. H. Son, S. M. Hughes, Y. D. Yin and A. P. Alivisatos, *Science*, 2004, **306**, 1009–1012.
- 20 Z. Z. Li, M. Saruyama, T. Asaka, Y. Tatetsu and T. Teranishi, *Science*, 2021, **373**, 332.
- 21 J. M. Pietryga, D. J. Werder, D. J. Williams, J. L. Casson, R. D. Schaller, V. I. Klimov and J. A. Hollingsworth, *J. Am. Chem. Soc.*, 2008, **130**, 4879–4885.
- 22 M. Casavola, M. A. van Huis, S. Bals, K. Lambert, Z. Hens and D. Vanmaekelbergh, *Chem. Mater.*, 2012, **24**, 294–302.
- 23 K. Lambert, B. De Geyter, I. Moreels and Z. Hens, *Chem. Mater.*, 2009, **21**, 778.
- 24 Y. Justo, P. Geiregat, K. Van Hoecke, F. Vanhaecke, C. D. Donega and Z. Hens, *J. Phys. Chem. C*, 2013, **117**, 20171–20177.
- 25 J. B. Zhang, J. B. Gao, C. P. Church, E. M. Miller, J. M. Luther, V. I. Klimov and M. C. Beard, *Nano Lett.*, 2014, **14**, 6010–6015.
- 26 J. B. Zhang, B. D. Chernomordik, R. W. Crisp, D. M. Kroupa, J. M. Luther, E. M. Miller, J. B. Gao and M. C. Beard, *ACS Nano*, 2015, **9**, 7151–7163.
- 27 A. G. Butterfield, L. T. Alameda and R. E. Schaak, *J. Am. Chem. Soc.*, 2021, **143**, 1779–1783.
- 28 F. Huang, J. Ning, Z. Duan, A. A. Sergeev, A. Portniagin, S. V. Kershaw, J. Tian and A. L. Rogach, *Chem. Mater.*, 2021, **33**, 2398–2407.
- 29 G. Gariano, V. Lesnyak, R. Brescia, G. Bertoni, Z. Dang, R. Gaspari, L. De Trizio and L. Manna, *J. Am. Chem. Soc.*, 2017, **139**, 9583–9590.
- 30 L. De Trizio, H. Li, A. Casu, A. Genovese, A. Sathya, G. C. Messina and L. Manna, *J. Am. Chem. Soc.*, 2014, **136**, 16277–16284.
- 31 H. Li, R. Brescia, R. Krahne, G. Bertoni, M. J. Alcocer, C. D'Andrea, F. Scotognella, F. Tassone, M. Zanella and M. De Giorgi, *ACS Nano*, 2012, **6**, 1637–1647.
- 32 G. Han, S. R. Popuri, H. F. Greer, J. W. G. Bos, W. Zhou, A. R. Knox, A. Montecucco, J. Siviter, E. A. Man and M. Macauley, *Angew. Chem., Int. Ed.*, 2016, **55**, 6433–6437.
- 33 V. Sidey, *J. Phys. Chem. Solids*, 2022, **171**, 110992.
- 34 L. H. Ahrens, *Geochim. Cosmochim. Acta*, 1952, **2**, 155–169.
- 35 R. D. Shannon, *Acta Crystallogr., Sect. A: Cryst. Phys., Diffraction, Theor. Gen. Crystallogr.*, 1976, **32**, 751–767.
- 36 S. S. Ralsanov, *Acta Crystallogr., Sect. B: Struct. Sci., Cryst. Eng. Mater.*, 2020, **76**, 38–40.
- 37 M. S. Shirazi, A. Foroumadi, I. Saberikia and M. M. Farimani, *Ultrason. Sonochem.*, 2022, **87**, 106037.
- 38 B. Pejjaei, V. R. M. Reddy, K. Seku, M. R. Pallavolu and C. Park, *New J. Chem.*, 2018, **42**, 4843–4853.
- 39 R. G. Parr and R. G. Pearson, *J. Am. Chem. Soc.*, 1983, **105**, 7512–7516.
- 40 H. Lee, B. Yoo, D. Kim, J. Cha, Y. K. Kang, S.-P. Cho, T. Hyeon, M.-G. Kim, M. G. Kanatzidis and I. Chung, *J. Am. Chem. Soc.*, 2023, **145**, 15951–15962.
- 41 S. H. Heo, S. Jo, H. S. Kim, G. Choi, J. Y. Song, J.-Y. Kang, N.-J. Park, H. W. Ban, F. Kim and H. Jeong, *Nat. Commun.*, 2019, **10**, 864.
- 42 L. Huang, J. Lu, D. Ma, C. Ma, B. Zhang, H. Wang, G. Wang, D. H. Gregory, X. Zhou and G. Han, *J. Mater. Chem. A*, 2020, **8**, 1394–1402.
- 43 G. Zhang, G. Li, J. Wang, H. Tong, J. Wang, Y. Du, S. Sun and F. Dang, *Adv. Energy Mater.*, 2022, **12**, 2103910.
- 44 Y. K. Lee, Z. Luo, S. P. Cho, M. G. Kanatzidis and I. Chung, *Joule*, 2019, **3**, 719–731.
- 45 L.-D. Zhao, C. Chang, G. Tan and M. G. Kanatzidis, *Energy Environ. Sci.*, 2016, **9**, 3044–3060.
- 46 D. Lee, W. Park, Y. A. Kang, H. T. Lim, S. Park, Y. Mun, J. Kim and K.-S. Jang, *ACS Appl. Mater. Interfaces*, 2023, **15**, 3047–3053.
- 47 V. Krylova, S. Žalėnienė, N. Dukstienė and J. Baltrusaitis, *Appl. Surf. Sci.*, 2015, **351**, 203–208.
- 48 C. Wang, H. Yang, Q. Wang, L. Qiao, X. Peng, J. Li, J. Han, Q. Wang, X. Li and Z. Wang, *CrystEngComm*, 2020, **22**, 5296–5301.
- 49 L. Qiu, X. Lai and J. Jian, *Mater. Charact.*, 2021, **172**, 110864.
- 50 P. Naumov, O. Barkalov, H. Mirhosseini, C. Felser and S. Medvedev, *J. Phys.: Condens. Matter*, 2016, **28**, 385801.
- 51 C.-N. Zhu, P. Jiang, Z.-L. Zhang, D.-L. Zhu, Z.-Q. Tian and D.-W. Pang, *ACS Appl. Mater. Interfaces*, 2013, **5**, 1186–1189.
- 52 Z. Li, L. Peng, Y. Fang, Z. Chen, D. Pan and M. Wu, *Radiat. Phys. Chem.*, 2011, **80**, 1333–1336.
- 53 V. Drozd, I. Nikiforova, V. Bogevoilnov, A. Yafyasov, E. Filatova and D. Papazoglou, *J. Phys. D: Appl. Phys.*, 2009, **42**, 125306.
- 54 W. Shi, M. Gao, J. Wei, J. Gao, C. Fan, E. Ashalley, H. Li and Z. Wang, *Adv. Sci.*, 2018, **5**, 1700602.
- 55 R. Wang, M. Wei, G. Jiang, W. Liu and C. Zhu, *Chem. Lett.*, 2014, **43**, 693–695.
- 56 M. Pachecka, J. M. Sturm, R. W. E. van de Kruijs, C. J. Lee and F. Bijkerk, *AIP Adv.*, 2016, **6**, 075222.
- 57 S. S. Batsanov, *Russ. Chem. Rev.*, 1982, **51**, 684.
- 58 M.-A. Langevin, D. Lachance-Quirion, A. M. Ritcey and C. N. Allen, *J. Phys. Chem. C*, 2013, **117**, 5424–5428.

