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The effect of co-substitution of heterovalent ions Ga³⁺ and Sb⁵⁺ on nonlinear optical properties of phosphate crystals†

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The strong second harmonic response of nonlinear optical (NLO) materials is of great scientific significance for laser technology. The second harmonic generation (SHG) response in NLO crystals is usually enhanced by introducing non-centrosymmetric active structural units with permanent hyperpolarizability into the structure. We have designed and grown a series of Ga³⁺:Sb⁵⁺:KTIPOPO₄ crystals based on traditional KTIPOPO₄ by co-doping Ga and Sb ions. The chemical compositions of the crystals were determined through inductively coupled plasma atomic emission spectroscopy (ICP-AES). The electronic structure was analyzed by X-ray photoelectron spectroscopy (XPS), and it was determined that the chemical strengths of Ti–O K–O and P–O bonds are related to the binding energy difference. After the introduction of doping ions, the frequency doubling conversion efficiency increased from 1064 to 532 nm, up to approximately 22%. And the SHG responses of Ga³⁺:Sb⁵⁺:KTP crystals have significantly enhanced, with a NLO coefficient close to 1.25 times that of KTP crystals. First-principles calculations and structural analyses show that the improved SHG response of Ga³⁺:Sb⁵⁺:KTP crystals is mainly influenced by the [SbO₆] groups. This study provides a new illustration for improving the NLO performance of crystals.

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1. Introduction

In 1961, Franken *et al.*¹ first discovered the laser frequency doubling phenomenon in a quartz crystal, marking the beginning of NLO materials and promoting their rapid development. NLO crystals are one of the important photoelectric information functional materials, and their development degree is closely related to the development of laser technology. They can be used for converting laser frequency and expanding the wavelength of lasers; modulating the intensity and phase of lasers; realizing holographic storage of laser signals, and so on. They are indispensable materials in high-tech and modern military technology, and have important applications in the fields of photoelectric communication, laser processing, laser weapons and laser medicine. Metal phosphates have attracted widespread attention from the academic community due to their excellent physical and chemical properties and structural

advantages.^{2–5} The structure of a PO₄ tetrahedron is conducive to ensuring a shorter UV cutoff edge, making it a good candidate material for UV or deep UV NLO crystals.^{2,6–12} However, due to the highly symmetric tetrahedral geometries of [PO₄]^{3–} units, metal phosphates exhibit a weak SHG response.¹³ Up to now, the largest SHG response of an alkali- or alkaline-earth metal phosphate is 2.6 × KDP (LiCs₂PO₄).² As a critical optical performance of nonlinear optical crystals, obtaining a large NLO response to achieve high laser output efficiency remains a core issue.

In order to improve the SHG response, it is necessary to introduce additional structural NLO-active units, that is, introducing non-centrosymmetric active structural units with permanent hyperpolarizability into the structure, which are generally divided into the following three types: (1) octahedrally coordinated d⁰ transition metal cations (Zr⁴⁺, Ti⁴⁺, Nb⁵⁺, Mo⁶⁺, W⁶⁺, V⁵⁺, Ta⁵⁺, and Sb⁵⁺) with a second-order Jahn Teller (SOJT) effect,^{14–17} (2) stereoactive lone pair cations (Pb²⁺, Bi³⁺, Sn²⁺, Se⁴⁺, and Te⁴⁺),^{18–22} and (3) d¹⁰ cations (Zn²⁺, Cd²⁺, and Hg²⁺) with large polar displacement.^{18,23–25} The influence of different nonlinear active units on the NLO performance of crystals varies. Taking octahedrally coordinated d⁰ transition metal cations as an example, it was found that different metal polyhedra have different degrees of distortion, and the average

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magnitude of the intraoctahedral distortion scales are as follows: $\text{Mo}^{6+} > \text{V}^{5+} > \text{W}^{6+} > \text{Nb}^{5+} > \text{Ta}^{5+} > \text{Ti}^{4+}$.¹⁴ For example, two new phosphate crystals: KCsWP_2O_9 and $\text{KCsMoP}_2\text{O}_9$ crystals,^{15,16} both of which contain noncentrosymmetric active ions: W^{6+} and Mo^{6+} . Their NLO coefficients are 16.8 times and 27 times that of KDP crystals, respectively. The large NLO responses of both crystals originate from the octahedron formed by d^0 transition metal cations: WO_6 and MO_6 octahedrons, and it is confirmed from the literature that the distortion of Mo^{6+} ions is greater than that of W^{6+} ions,¹⁴ which leads to a stronger NLO response. Due to the small NLO coefficient of LiCs_2PO_4 ,² some researchers have modified it *via* ion substitution, for example, by introducing NLO-active ions such as Hg^{2+} and Pb^{2+} ions and obtaining two new crystals: LiHgPO_4 ²³ and LiPbPO_4 .²⁶ The PbO_8 polyhedra and the highly distorted HgO_6 octahedron are the important sources of the SHG effect. The NLO coefficients were successfully enhanced, which are 11 times and 3 times that of KDP crystals, respectively.

KTP crystals have excellent NLO properties and have been a research hotspot in recent years. In order to further improve the NLO response and frequency doubling conversion efficiency of KTP crystals, we consider doping them with NLO-active ions. From the structural aspect, it can be inferred that Sb has multiple coordination modes and is able to form various polyhedra with high distortion and large anisotropy.²⁷ However, single doping of NLO-active ions in KTP crystals can lead to difficulties in crystal growth,²⁸ and the NLO performance of the crystals will not be improved.²⁹ Therefore, we consider introducing the Ga element to balance charges, as Ga^{3+} , Sb^{5+} and Ti^{4+} ions have extremely similar ionic radius and are prone to substitution. In this study, we chose the Sb^{5+} NLO-active ion with a d^{10} electronic configuration and Ga^{3+} ions for co-doping in a KTP crystal. Based on the analysis above, we successfully synthesized a series of KTP-type crystals $\text{Ga}^{3+}:\text{Sb}^{5+}:\text{KTP}$, explored their optical and NLO properties and analyzed the correlation between their structure and performance.

2. Experimental

2.1 Crystal growth

Ga^{3+} and Sb^{5+} co-doped KTP crystals with the composition of $\text{KTi}_{1-2x}\text{Ga}_x\text{Sb}_x\text{OPO}_4$ ($x = 0.025$ and 0.05) were grown through the top-seeded solution growth method. K_2CO_3 , KH_2PO_4 , TiO_2 , Ga_2O_3 , and Sb_2O_5 with 99.9% purity were used as raw materials. The starting mixtures were placed in a platinum crucible with a diameter of 70 mm and a height of 90 mm. Then, the reagents were mixed evenly and heated to generate the flux ($\text{K}_6\text{P}_4\text{O}_{13}$) and the solute ($\text{Ga}^{3+}:\text{Sb}^{5+}:\text{KTP}$). After the saturation temperature of the solution was measured by the seed crystal method, a KTP seed was introduced into the solution at about 1 °C above the saturation temperature for a period of time to dissolve the impurities on the seed surface. After this, the temperature of the solution decreased slowly to grow the crystal. The rotation rate was set to 30 rpm and the crystal growth period was about 40 days. After the growth period, the crystal was taken out of the

solution and cooled to ambient temperature. The grown $\text{Ga}^{3+}:\text{Sb}^{5+}:\text{KTP}$ crystals were processed into dimensions of $6 \times 13 \times 10 \text{ mm}^3$ and $5 \times 10 \times 8 \text{ mm}^3$, respectively.

2.2 Powder X-ray diffraction (PXRD)

The powder X-ray diffraction (RXP) patterns were recorded using a Bruker-AXS D8 Advance X-ray diffractometer ($\text{Cu K}\alpha = 1.5406 \text{ \AA}$) produced by Bruker-AXS company in German. The measurement was accomplished at room temperature with an angle range of $2\theta = 10\text{--}70^\circ$ and a scan step of 0.02° .

2.3 Thermal stability

Thermogravimetric (TG) and differential thermal analysis (DTA) data were recorded using a NETZSCH STA8000 simultaneous analyzer in air. The heating rate was $10 \text{ }^\circ\text{C min}^{-1}$ in the test range of 30–1300 °C.

2.4 X-ray photoelectron spectroscopy

X-ray photoelectron spectroscopy (XPS) measurements of $\text{Ga}^{3+}:\text{Sb}^{5+}:\text{KTP}$ crystals were performed using an ESCALAB 250 with Al $\text{K}\alpha$ radiation and using C 1s (284.8 eV) as a reference to investigate the elemental composition and valence of samples.

2.5 Inductively coupled plasma atomic emission spectroscopy (ICP-AES)

The elemental contents of $\text{Ga}^{3+}:\text{Sb}^{5+}:\text{KTP}$ crystals were analyzed using a PerkinElmer NEXION350X inductively coupled plasma atomic emission spectrometer (ICP-AES). Five different concentrations of $5 \text{ }\mu\text{g L}^{-1}$, $10 \text{ }\mu\text{g L}^{-1}$, $20 \text{ }\mu\text{g L}^{-1}$, $50 \text{ }\mu\text{g L}^{-1}$, and $100 \text{ }\mu\text{g L}^{-1}$ calibration solution samples of Ga^{3+} and Sb^{5+} ions were used. The crystal powder samples were digested in sulfuric acid before the analysis.

2.6 UV-vis-NIR transmission spectroscopy

The room temperature optical transmittance spectra of $\text{Ga}^{3+}:\text{Sb}^{5+}:\text{KTP}$ crystals were measured using an Agilent Cary 5000 UV-vis-NIR spectrophotometer in the wavelength range from 200 to 700 nm and a Thermo Nicolet NEXUS 670 FTIR spectrometer from 2500 to 6000 nm.

2.7 Second harmonic generation measurements

The power SHG responses of $\text{Ga}^{3+}:\text{Sb}^{5+}:\text{KTP}$ were measured by means of the Kurtz–Perry method.³⁰ The fundamental wavelength of 1064 nm was generated using a Q-switched Nd:YAG laser. The crystals were ground and sieved into distinct particle size ranges: 38–48, 48–75, 75–96, 96–150, 150–250 and 250–400 μm . To judge the strength of SHG, pure KTP crystals were also ground and sieved into the same particle size range and used as a reference. The SHG intensity of the $\text{Ga}^{3+}:\text{Sb}^{5+}:\text{KTP}/\text{KTP}$ ratio was evaluated based on the same size (200–300 μm).

2.8 First-principles calculations

The electronic structure and optical properties of $\text{Ga}^{3+}:\text{Sb}^{5+}:\text{KTP}$ and KTP crystals were calculated using CASTEP software through density functional theory (DFT).³¹ The generalized



gradient approximation (GGA) PBE exchange-correction functional and norm-conserving pseudopotential are adopted. In order to simulate the doping system, we replaced a Ti atom with a doped Ga atom and Sb atom, with a doping concentration of 2%. The model was fully optimized using the BFGS method. The convergence thresholds for energy change and force were set to 10^{-5} eV and 0.01 eV \AA^{-1} , respectively. The energy cutoff values were set to 800 eV and 750 eV, respectively. A convergence threshold of 1×10^{-6} eV per atom was set for the self-consistent-field of the total electronic energy. The numerical integration of the Brillouin zones was performed using $1 \times 1 \times 1$ and $1 \times 2 \times 1$ Monkhorst–Pack k -point meshes, respectively. Norm-conserving pseudopotentials were used with the following valence electron configurations: K, $3s^2 3p^6$; Ti, $3s^2 3p^6 3d^2 4s^2$; O, $2s^2 2p^4$; P, $3s^2 3p^3$; Ga, $3s^2 3p^6 3d^{10}$; Sb, and $5s^2 5p^3$.

3. Results and discussion

3.1 Crystal growth

The crystal phase for $\text{Ga}^{3+}:\text{Sb}^{5+}:\text{KTP}$ crystals was detected using powder X-ray diffraction (PXRD). The results and crystal photographs are shown in Fig. 1. It is established that the PXRD patterns of the grown crystals match with those of pure KTP crystals very well, indicating that the grown crystals were pure-phase and the incorporation of Ga and Sb hardly changes the crystal phase of the original compound. However, by partially amplifying the PXRD spectrum, it can be observed that some peak positions show a slight shift toward a lower diffraction

angle when Ga and Sb entering the crystal structure, indicating that the lattice parameters of these doped crystals have changed a bit.

3.2 Thermal stability

The thermal stabilities of pure KTP and $\text{Ga}^{3+}:\text{Sb}^{5+}:\text{KTP}$ crystals were investigated through TGA and DTA. As shown in Fig. 2, it can be observed that all three crystals have only one endothermic peak, which represents the melting point of the crystals. With the increase of the doping concentration, the melting point of the crystal gradually increased from 1158°C to 1162°C , and there are also mass loss on the TG curves. The results show that the thermal stability of KTP crystals improved after doping with Ga^{3+} and Sb^{5+} , and this phenomenon also occurs in the RTiOPO_4 crystal.³²

3.3 X-ray photoelectron spectroscopy

The electronic structures of $\text{Ga}^{3+}:\text{Sb}^{5+}:\text{KTP}$ crystals were studied using X-ray photoelectron spectroscopy (XPS) measurements at 300 K. The XPS spectra of Ti 2p, K 2p P 2p and Sb 3d core levels are displayed in Fig. 3. The characteristic peaks of Ga elements cannot be accurately analyzed possibly due to excessive noise. The characteristic peaks of Sb $3d_{5/2}$ and Sb $3d_{3/2}$ were observed to be located at 532.4 and 539.8 eV, respectively, indicating that the valence of the doped Sb element is SbV .³³ The binding energy of Ti 2p doublet peaks shows a shift to higher binding energy with the increase of the doping concentration in $\text{Ga}^{3+}:\text{Sb}^{5+}:\text{KTP}$ crystals. To study the influence of Ga and Sb co-doping on the chemical bonding properties of KTP crystals, after obtaining the binding energy (BE) value of the representative levels, we use the parameter of binding energy difference: $\Delta\text{BE}(\text{O}-\text{M}) = \text{BE}(\text{O } 1s) - \text{BE}(\text{M})$ to characterize the chemical bond properties of metal and oxygen in the compound, and M is a representative metal core-level. For $\text{Ga}^{3+}:\text{Sb}^{5+}:\text{KTP}$ crystals, the calculated binding energy differences are shown in Table 1. The values of $\Delta\text{BE}(\text{O}-\text{Ti})$ increase as the doping concentration increases, and the values of both $\Delta\text{BE}(\text{O}-\text{P})$ and $\Delta\text{BE}(\text{O}-\text{K})$ increase with the increase of Sb content. From the view of the strength of the bonds,^{34,35} we consider that in the case of a small amount of doping, the ionic strength of O–Ti bonds increased, and the covalent strength of O–P and O–K bonds also increased in $\text{Ga}^{3+}:\text{Sb}^{5+}:\text{KTP}$. It is worth noting that according to Fig. 3d and Fig. S1 (ESI[†]), the ratio of the peak area of oxygen vacancies to the peak area of O 1s before and after doping decreased from 8% to 6%, which proved that the oxygen vacancies in the $\text{Ga}^{3+}:\text{Sb}^{5+}:\text{KTP}$ crystal decreased. It is

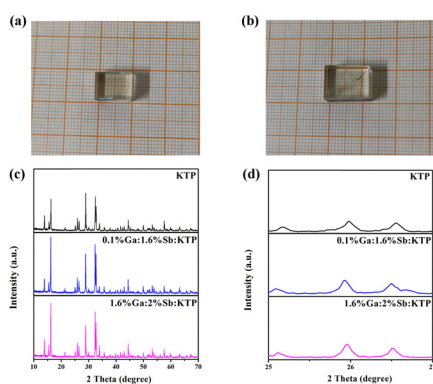


Fig. 1 Single crystal photographs of (a) 0.1% Ga:1.6% Sb:KTP and (b) 1.6% Ga:2% Sb:KTP. (c) Powder X-ray diffraction patterns of KTP and $\text{Ga}^{3+}:\text{Sb}^{5+}:\text{KTP}$ crystals. (d) Enlarged section of the 2θ region from 25 to 27° .

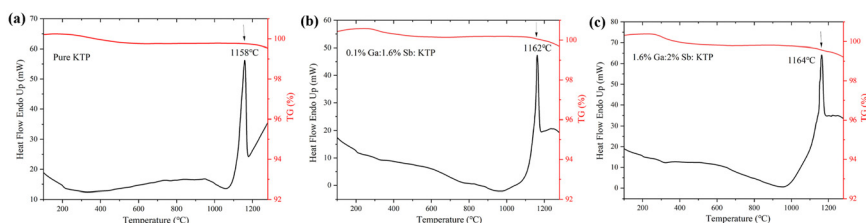


Fig. 2 TG and DTA curves of the (a) pure KTP, (b) 0.1% Ga:1.6% Sb:KTP, and (c) 1.6% Ga:2% Sb:KTP crystals.



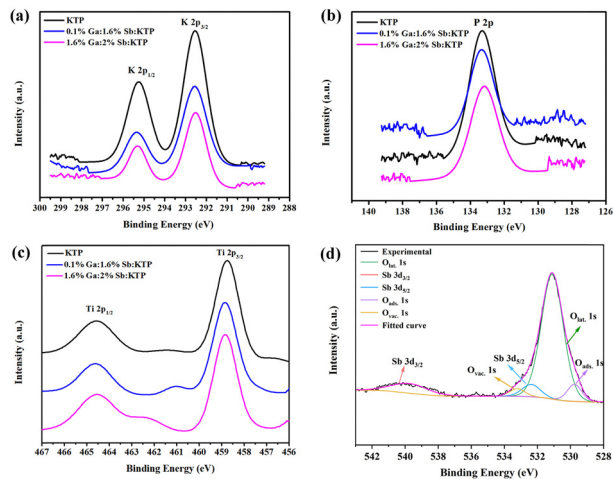


Fig. 3 High-resolution X-ray photoelectron spectra of $\text{Ga}^{3+}:\text{Sb}^{5+}:\text{KTP}$ crystals: (a) K 2p core level, (b) P 2p core level, (c) Ti 2p core level, and (d) Sb 3d and O 1s core levels.

Table 1 Binding energy values (eV) of pure and $\text{Ga}^{3+}:\text{Sb}^{5+}:\text{KTP}$ crystals

	KTP	0.1% Ga:1.6% Sb:KTP	1.6% Ga:2% Sb:KTP
K 2p _{3/2}	292.51	292.56	292.75
Ti 2p _{3/2}	458.74	458.85	458.84
Ti 2p _{1/2}	464.55	464.59	464.56
P 2p	133.27	133.3	133.1
O 1s	530.78	530.88	531.08
ΔBE (O 1s–Ti 2p _{3/2})	72.04	72.03	72.24
ΔBE (O 1s–Ti 2p _{1/2})	66.23	66.29	66.52
ΔBE (O 1s–P 2p)	397.51	397.58	397.98
ΔBE (O 1s–K 2p _{3/2})	238.27	238.32	238.33

attributable to the Sb^{5+} ions, in which extra electrons will neutralize the positive charges of the oxygen vacancy, resulting in the decrease of oxygen vacancies.

3.4 Inductively coupled plasma atomic emission spectroscopy (ICP-AES)

To obtain the chemical composition of the $\text{Ga}^{3+}:\text{Sb}^{5+}:\text{KTP}$ crystal, the concentrations of elements were measured through ICP-AES. Through calculations, the molecular compositions of the two crystals were found to be 0.1% Ga:1.6% Sb:KTP and 1.6% Ga:2% Sb:KTP, respectively.

3.5 UV-vis-IR transmission spectroscopy

The UV-vis-IR transmission spectra of KTP and $\text{Ga}^{3+}:\text{Sb}^{5+}:\text{KTP}$ compounds were collected and are shown in Fig. 4. As can be seen, the ultraviolet absorption edges of KTP, 0.1% Ga:1.6% Sb:KTP and 1.6% Ga:2% Sb:KTP were at about 353 nm, 357 nm and 360 nm, respectively, which increase with the doping concentration. The infrared cutoff edge of crystals was all located near 4.4 μm with almost no change. Based on the transmission results, the absorption coefficient can be estimated according to the following equation:

$$\alpha = \ln(100/T)/d \quad (1)$$

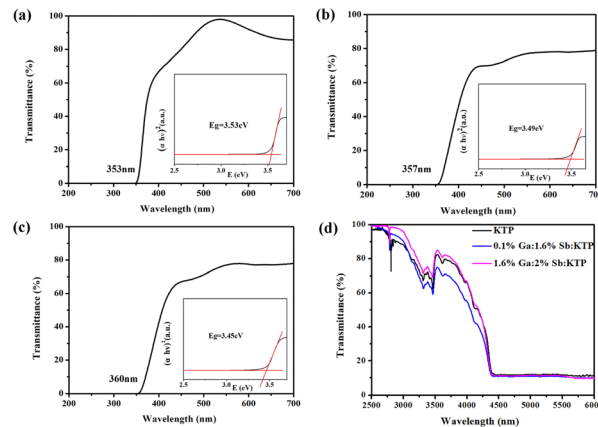


Fig. 4 UV-vis transmission spectra of crystals: (a) KTP crystal, (b) 0.1% Ga:1.6% Sb:KTP, and (c) 1.6% Ga:2% Sb:KTP. (d) Mid-IR transmission spectrum.

$$(\alpha h\nu)^{1/n} = A(h\nu - E_g) \quad (2)$$

where α , T , d , h , ν and E_g represent the absorption coefficient, transmittance (%), crystal thickness, Planck constant, frequency and bandgap energy, respectively; A is a constant, and n is a constant that has a value of 1/2 for direct bandgaps and 2 for indirect bandgaps. Thus, we can plot a figure according to eqn (2). By drawing $(\alpha h\nu)^{1/n}$ as the vertical axis and $h\nu$ as the horizontal axis, we can obtain the bandgap values. The curves show that the bandgap values of KTP, 0.1% Ga:1.6% Sb:KTP and 1.6% Ga:2% Sb:KTP are about 3.53 eV, 3.49 eV and 3.45 eV, respectively. It is worth noting that the bandgap values are gradually reduced with the increase of Sb^{5+} and Ga^{3+} . This phenomenon indicates that the doped Sb^{5+} and Ga^{3+} have changed the electronic structure of crystals to varying degrees. Nevertheless, the bandgap change caused by co-doping is relatively small, only reducing by 0.08 eV, which has significant advantages compared to single doped crystals.³⁶

3.6 Second harmonic generation measurements

The optical SHG responses of $\text{Ga}^{3+}:\text{Sb}^{5+}:\text{KTP}$ were measured by means of the Kurtz–Perry method.³⁰ Pure KTP crystals were used as the reference. The measurement results indicated that $\text{Ga}^{3+}:\text{Sb}^{5+}:\text{KTP}$ crystals exhibited a strong NLO response. It can be seen from Fig. 5a and b that when the doping concentrations were 0.1% Ga and 1.6% Sb, the SHG response is the highest, reaching 1.25 times that of KTP crystals. When the doping concentrations were 1.6% Ga and 2% Sb, the SHG response actually slightly decreased. In addition, it can be seen from Fig. 5a that the SHG intensity increases with the increase of the particle size, indicating that all the $\text{Ga}^{3+}:\text{Sb}^{5+}:\text{KTP}$ crystals exhibit type I phase-matching behavior under 1064 nm laser irradiation. Furthermore, the SHG performance of bulk single crystals was also characterized. The frequency doubling conversion efficiencies from 1064 nm to 532 nm of bulk single crystals with different doping concentrations were obtained. The experimental configuration is shown in Fig. 5g. The base



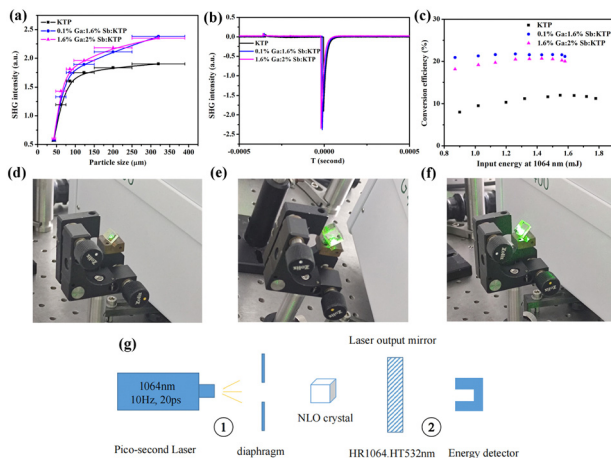


Fig. 5 NLO properties of $\text{Ga}^{3+}:\text{Sb}^{5+}:\text{KTP}$. (a) Phase-matching curves of KTP, 0.1% Ga:1.6% Sb:KTP and 1.6% Ga:2% Sb:KTP by powder SHG measurements at 1064 nm. (b) The oscilloscope traces of the SHG signals for the KTP, 0.1% Ga:1.6% Sb:KTP and 1.6% Ga:2% Sb:KTP at 1064 nm. (c) Frequency doubling conversion efficiencies from 1064 nm to 532 nm of KTP, 0.1% Ga:1.6% Sb:KTP and 1.6% Ga:2% Sb:KTP crystals. The green laser output of the bulk (d) KTP, (e) 0.1% Ga:1.6% Sb:KTP, and (f) 1.6% Ga:2% Sb:KTP crystals using a 1064 nm laser. (g) Schematic diagram of the frequency doubling experimental configuration.

frequency source is a pulsed picosecond 1064 nm laser, the repetition frequency is 10 Hz, and the pulse width is 20 ps. The fundamental frequency laser spot is reduced by the diaphragm and incident perpendicular to the crystal, with a crystal thickness of 5 mm. The 1064 nm fundamental frequency laser passing through the frequency doubling crystal generated a 532 nm laser, leaving a portion of the 1064 nm laser that has not been doubled. This part of the 1064 nm laser was filtered out through the 1064 nm high reflection 532 nm high transmittance laser output mirror, leaving only 532 nm laser in the end. The output energy was measured using an energy detector. When the detector was located in position 1, the energy of the initial 1064 nm laser was measured; when the detector was located in position 2, the output energy of the 532 nm laser was measured. As shown in Fig. 5c, it can be seen that as the input energy increases, the conversion efficiency increases and eventually tends to be saturated. The maximum conversion efficiencies of KTP, 0.1% Ga:1.6% Sb:KTP and 1.6% Ga:2% Sb:KTP are 12%, 21.74% and 20.67%, respectively, which correspond to the results of powder frequency doubling. This improvement is greatly significant. The intensity of the output green laser can be seen from the crystal photographs; pure KTP has the weakest green laser and $\text{Ga}^{3+}:\text{Sb}^{5+}:\text{KTP}$ has a strong green laser, which are consistent with the experimental data. The above results indicate that the co-doping of Ga and Sb ions does contribute to improving the NLO response of the crystal. However, while the concentration of the Sb^{5+} ion remains almost invariable, as the concentration of the Ga ion increased from 0.1% to 1.6%, the SHG response decreased a little. This phenomenon further indicates that only NLO-active ions can improve the NLO performance of crystals, and the introduction of the Ga ion will only play a role in balancing charges and does not improve

the NLO performance of the crystal. In brief, the co-substitution of Ga^{3+} and Sb^{5+} successfully improved the NLO properties of traditional KTP crystals.

3.7 Electronic structure analysis

In order to explore the correlation between the crystal structure and the performance, we calculated the electronic structures of pure KTP and $\text{Ga}^{3+}:\text{Sb}^{5+}:\text{KTP}$ crystals based on first principles. Based on previous research on the heterovalent co-doping of RTP, which is isomorphic to KTP, we assume that doped ions occupy the Ti site for theoretical calculations.^{34,35} The band structure indicates that the pure KTP crystal has an indirect band gap of 3.335 eV, which is similar to the experimental result. The $\text{Ga}^{3+}:\text{Sb}^{5+}:\text{KTP}$ crystal has an indirect band gap of 2.878 eV, which is slightly smaller than the experimental result. The GGA usually underestimates the band gap, which has been often observed in previous studies.^{6,37} The calculated total densities of states and partial densities of states (TDOS and PDOS) are plotted in Fig. 6b and d. The density of states (DOS) shows that, for pure KTP crystals, the bottom region of the valence band (VB) is composed of K 3p states, which is located at about -11 eV. The region of the VB from -8 eV to the Fermi level is mainly contributed by O 2p, with minor contribution from P 3p and Ti 3d orbitals. The bottom region of the conductive bands (CB) is mainly composed of Ti 3d and O 2p orbitals. The region of the CB from 8 eV to 13 eV is mainly contributed by K, P 3p, O 2p orbitals and minor contribution from Ti 3p orbitals. For $\text{Ga}^{3+}:\text{Sb}^{5+}:\text{KTP}$, the top of the VB was dominated by Ti 3d, Sb 5p, P 2p, and O 2p orbitals, and the bottom of the CB was dominated by Ti 3d and Sb 5s orbitals. The introduction of Sb ion orbitals near the Fermi energy level leads to the increase of electron activity at the bottom of the conduction band, which makes the movement of the conduction band larger than that of the valence band, resulting in the down-shift of the conduction band.

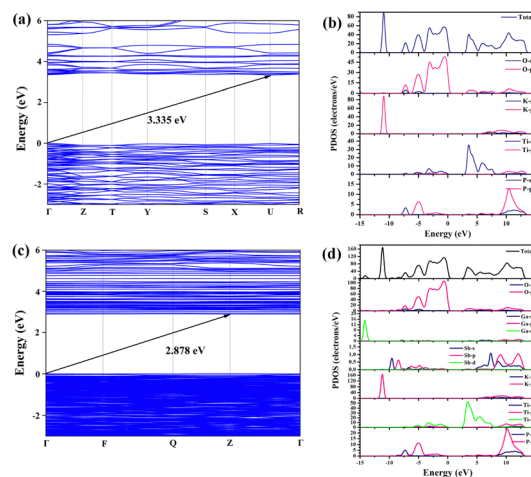


Fig. 6 Band structure and DOS of pure KTP and $\text{Ga}^{3+}:\text{Sb}^{5+}:\text{KTP}$. (a) Band structure of KTP. (b) DOS of KTP. (c) Band structure of $\text{Ga}^{3+}:\text{Sb}^{5+}:\text{KTP}$. (d) DOS of $\text{Ga}^{3+}:\text{Sb}^{5+}:\text{KTP}$.



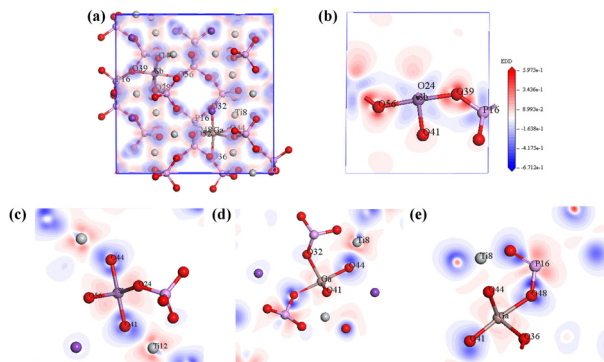


Fig. 7 Calculated electron-density difference (EDD) plots of the local structure: (a) total EDD of $\text{Ga}^{3+}:\text{Sb}^{5+}:\text{KTP}$, (b) $\text{Sb}-\text{O}-\text{P}$ EDD, (c) $\text{Sb}-\text{O}-\text{Ti}$ EDD, (d) $\text{Ga}-\text{O}-\text{Ti}$ EDD, and (e) $\text{Ga}-\text{O}-\text{P}$ EDD.

Consequently, it can be understood that the introduction of Sb ions changes the densities of states near the bottom of the CB and the top of the VB in KTP, causing the bottom of the CB to move downwards and reduce the band gap. Since the optical properties of crystals mainly rely on the optical transitions between the electronic states near the bandgap, while the 3d orbitals introduced by Ga ions are located around -14 eV. Therefore, the Ga ions have a minimal effect on the densities of states and only play a role in balancing charges within the structure. The improved NLO properties of $\text{Ga}^{3+}:\text{Sb}^{5+}:\text{KTP}$ mainly originate from the contribution of TiO_6 , PO_4 and SbO_6 groups. These results correspond to the calculated dipole moment (Table S1, ESI[†]), and the MO_6 octahedron with the second-order Jahn Teller (SOJT) effect dominates the NLO response.

To observe the polarization degree of the bridged oxygen electronic structure of NLO-active units (SbO_6 , TiO_6 , and PO_4 groups) in $\text{Ga}^{3+}:\text{Sb}^{5+}:\text{KTP}$ crystals, the calculated electron-density difference (EDD) plots of the local structure are shown in Fig. 7. From the schematics of the local structure, it can be seen that the electron cloud distribution around Ga atoms is relatively uniform, while the electron cloud around Sb and Ti atoms tends to be far away from the O side, which has obvious polarity distribution characteristics. This polarity distribution is conducive to improve the microscopic second-order susceptibility of NLO-active units. According to the anionic group theory, the overall SHG coefficient of the crystal is the geometrical superposition of the microscopic second-order susceptibility tensors of the relevant ionic groups³⁸ (SbO_6 , TiO_6 and PO_4 groups), and ultimately reflected in the significant increase of the SHG coefficient of the $\text{Ga}^{3+}:\text{Sb}^{5+}:\text{KTP}$ crystals. And, according to the trend of theoretical calculation results, it can be predicted that if the concentration further increases, the NLO performance of the crystal will continue to improve.

4. Conclusions

In summary, a series of co-doped $\text{Ga}^{3+}:\text{Sb}^{5+}:\text{KTP}$ crystals with KTP structures were designed and prepared in this study. The transparent, high quality $\text{Ga}^{3+}:\text{Sb}^{5+}:\text{KTP}$ single crystals were

successfully grown using the top seed solution growth method. With an increase in the doping concentration, the UV cutoff edge shifted from 353 nm to 360 nm, and the optical bandgap decreased from 3.53 eV to 3.45 eV. The change in the bandgap is significantly reduced compared to that of single doped crystals, which is beneficial for maintaining their positive optical performance. The chemical bond properties of $\text{Ga}^{3+}:\text{Sb}^{5+}:\text{KTP}$ crystals were studied by using binding energy difference between the metals and oxygen. The ionic strength of the O-Ti bonds increased, and the covalent strength of the O-P and O-K bonds also increased in the case of a small amount of doping. By comparing the proportion of peak areas of oxygen vacancies, it has been proved that the number of oxygen vacancies has decreased after co-doping. The frequency doubling conversion efficiency and NLO coefficient of $\text{Ga}^{3+}:\text{Sb}^{5+}:\text{KTP}$ crystals have significantly improved compared to those of pure KTP crystals (1.25 times that of KTP crystals) at 1064 nm. Through DFT calculations, it has been confirmed that the improvement of NLO performance was due to the contribution of TiO_6 , SbO_6 and PO_4 groups. The contribution of Ga ions is almost negligible, only serving to balance charges. The NLO performance of KTP type crystals was significantly improved through co-doping without significantly changing the crystal bandgap. This provides a good strategy for the introduction of non-centrosymmetric active structural units to improve crystal performance and explore new NLO crystals.

Author contributions

The article was written through contributions from all authors. All authors have given approval to the final version of the article.

Conflicts of interest

There are no conflicts to declare.

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Notes and references

- 1 A. Franken, A. E. Hill, C. W. Peters and G. Weinreich, *Phys. Rev. Lett.*, 1961, 7, 118.
- 2 L. Li, Y. Wang, B. H. Lei, S. J. Han, Z. H. Yang, K. R. Poeppelmeier and S. L. Pan, *J. Am. Chem. Soc.*, 2016, 138, 9101–9104.
- 3 L. Li, Y. Wang, B. H. Lei, S. J. Han, Z. H. Yang, H. Y. Lia and S. L. Pan, *J. Mater. Chem. C*, 2017, 5, 269–274.
- 4 S. W. Liu, B. B. Zhang, H. P. Wu, H. W. Yu, Z. G. Hu, J. Y. Wang and Y. C. Wu, *Inorg. Chem. Front.*, 2021, 8, 2061–2067.
- 5 W. F. Zhou and S. P. Guo, *Acc. Chem. Res.*, 2024, 57, 648–660.



- 6 H. W. Yu, J. S. Young, H. P. Wu, W. G. Zhang, J. M. Rondinelli and P. S. Halasyamani, *Chem. Mater.*, 2017, **29**, 1845–1855.
- 7 Z. Qian, H. P. Wu, H. W. Yu, Z. G. Hu, J. Y. Wang and Y. C. Wu, *New J. Chem.*, 2020, **44**, 6771–6777.
- 8 L. Li, S. J. Han, B. H. Lei, Y. Wang, H. Y. Li, Z. H. Yanga and S. L. Pan, *Dalton Trans.*, 2016, **45**, 3936–3942.
- 9 Y. F. Li, F. Liang, H. M. Song, W. Liu, Z. S. Lin, G. C. Zhang and Y. C. Wu, *Solid State Sci.*, 2019, **91**, 23–27.
- 10 T. Baiheti, H. Li, Z. H. Yang, S. J. Han, Y. Wang and S. L. Pan, *J. Solid State Chem.*, 2018, **266**, 150–154.
- 11 W. H. Liu, M. H. Lee, R. X. Guo and J. Y. Yao, *Inorg. Chem.*, 2023, **62**, 2480–2488.
- 12 F. F. Yuan, L. Hu, Z. Y. Bai, L. H. Liu, Y. S. Huang, L. Z. Zhang, D. S. Wei and Z. B. Lin, *Inorg. Chem.*, 2021, **60**, 8103–8110.
- 13 Y. Yang, Y. Qiu, P. F. Gong, L. Kang, G. M. Song, X. M. Liu, J. L. Sun and Z. S. Lin, *Chem. – Eur. J.*, 2019, **25**, 5648–5651.
- 14 P. S. Halasyamani, *Chem. Mater.*, 2004, **16**, 3586–3592.
- 15 T. Baiheti, S. Han, A. Tudi, Z. H. Yang and S. L. Pan, *Chem. Mater.*, 2020, **32**, 3297–3303.
- 16 T. Baiheti, S. J. Han, A. Tudi, Z. H. Yang and S. L. Pan, *J. Mater. Chem. C*, 2020, **8**, 11441–11448.
- 17 C. F. Sun, C. L. Hu, X. Xu, B. P. Yang and J. G. Mao, *J. Am. Chem. Soc.*, 2011, **133**, 5561–5572.
- 18 Z. Q. Li, Y. Chen, P. F. Zhu, N. J. Ji, X. L. Duan and H. D. Jiang, *RSC Adv.*, 2017, **7**, 53111–53116.
- 19 J. K. Wang, B. K. Xiong, H. P. Wu, H. W. Yu, Z. G. Hu, J. Y. Wang and Y. C. Wu, *Inorg. Chem. Front.*, 2021, **8**, 344–351.
- 20 K. C. Chen, C. S. Lin, G. Peng, Y. Chen, H. Z. Huang, E. Z. Chen, Y. X. Min, T. Yan, M. Luo and N. Ye, *Chem. Mater.*, 2022, **34**, 399–404.
- 21 H. W. Yu, W. G. Zhang, J. S. Young, J. M. Rondinelli and P. S. Halasyamani, *J. Am. Chem. Soc.*, 2016, **138**, 88–91.
- 22 S. H. Kim, J. Yeon and P. S. Halasyamani, *Chem. Mater.*, 2009, **21**, 5335–5342.
- 23 B. L. Wu, C. L. Hu, F. F. Mao, R. L. Tang and J. G. Mao, *J. Am. Chem. Soc.*, 2019, **141**, 10188–10192.
- 24 S. G. Zhao, J. H. Luo, P. Zhou, S. Q. Zhang, Z. H. Sun and M. C. Hong, *RSC Adv.*, 2013, **3**, 14000–14006.
- 25 Y. L. Sun, G. X. Liu, Y. L. Lv, L. Ma, Y. L. Yao and R. L. Tang, *J. Solid State Chem.*, 2023, **325**, 124171.
- 26 G. P. Han, Q. Liu, Y. Wang, X. Su, Z. H. Yang and S. L. Pan, *Phys. Chem. Chem. Phys.*, 2016, **18**, 19123–19129.
- 27 C. Hu, B. B. Zhang, B. H. Lei, S. L. Pan and Z. H. Yang, *ACS Appl. Mater. Interfaces*, 2018, **10**, 26413–26421.
- 28 J. Y. Wang, Y. Y. Liu, J. Q. Wei, M. H. Jiang, Z. S. Shao, W. J. Liu and S. S. Jiang, *Cryst. Res. Technol.*, 1997, **32**, 319–327.
- 29 S. Sadhasivam, R. N. Perumal and P. Ramasamy, *Spectrochim. Acta, Part A*, 2015, **149**, 183–189.
- 30 S. K. Kurtz and T. T. Perry, A Powder Technique for the Evaluation of Nonlinear Optical Materials, *J. Appl. Phys.*, 1968, **39**, 3798–3813.
- 31 S. J. Clark, M. D. Segall, C. J. Pickard, P. J. Hasnip, M. I. J. Probert, K. Refson and M. C. Payne, *Z. Kristallogr.-Cryst. Mater.*, 2005, **220**, 567–570.
- 32 J. J. Carvajal, J. Gavalda, J. Massons, F. Diaz and M. Aguilo, *Chem. Mater.*, 2003, **15**, 2730–2736.
- 33 F. Garbassi, *Surf. Interface Anal.*, 1980, **2**, 165–169.
- 34 J. Liu, Z. Q. Li, P. F. Zhu and X. L. Duan, *Opt. Mater.*, 2018, **75**, 38–43.
- 35 Z. Q. Li, P. F. Zhu, J. M. Ding, N. J. Ji, Y. Chen, Z. Y. Wang, X. L. Duan and H. D. Jiang, *Cryst. Growth Des.*, 2019, **19**, 7143–7152.
- 36 P. F. Zhu, Z. Q. Li, Z. Y. Wang, L. Ai, J. F. Zhao, C. Li, X. L. Duan, X. Zhao and F. P. Yu, *Adv. Opt. Mater.*, 2022, **10**, 2201643.
- 37 H. N. Liu, H. P. Wu, H. W. Yu, Z. G. Hu, J. Y. Wang and Y. C. Wu, *J. Mater. Chem. C*, 2021, **9**, 15321–15328.
- 38 C. T. Chen, Y. C. Wu and R. K. Li, *Int. Rev. Phys. Chem.*, 1989, **8**, 65–91.

