

Cite this: *Mater. Adv.*, 2024,  
5, 3824

## *In situ* atomic-resolution study of transformations in double polymorph $\gamma/\beta$ -Ga<sub>2</sub>O<sub>3</sub> structures†

J. García-Fernández, \*<sup>a</sup> S. B. Kjeldby, <sup>a</sup> L. J. Zeng, <sup>b</sup> A. Azarov, <sup>a</sup> A. Pokle, <sup>a</sup>  
P. D. Nguyen, <sup>a</sup> E. Olsson, <sup>b</sup> L. Vines, <sup>a</sup> A. Kuznetsov \*<sup>a</sup> and Ø. Prytz \*<sup>a</sup>

Disorder-induced formation of metastable Ga<sub>2</sub>O<sub>3</sub> polymorphs as well as the recovery of the stable state upon annealing attract attention because of the fundamental novelty and rapidly increasing interest in the use of Ga<sub>2</sub>O<sub>3</sub> in practical applications. In this study, double polymorph  $\gamma/\beta$ -Ga<sub>2</sub>O<sub>3</sub> structures fabricated by the radiation-induced disorder approach were used as a starting point for systematic *in situ* annealing electron microscopy experiments. We show that, under the conditions of the TEM *in situ* annealing, double  $\gamma/\beta$ -Ga<sub>2</sub>O<sub>3</sub> polymorph structures remained stable up to 300 °C, when onsets of the  $\gamma$ -to- $\beta$  transformation become traceable, leading to a prominent  $\gamma$ - and  $\beta$ -mixture already at 500 °C. Interestingly, the recrystallization of the  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> occurs throughout the whole  $\gamma$ -film and the preferential alignments at the newly emerging  $\gamma/\beta$ -interfaces are different from that of the initial  $\gamma/\beta$ -interface formed as a result of the disorder-induced ordering. The alignments of the two polymorphs are maintained as a function of temperature – with a reduction in the volumetric ratio of  $\gamma$ -domains for increasing annealing temperature. Finally, at 1100 °C,  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> fully transforms into  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>, without dominating crystallographic relationships or preferred orientations, indicating that energy barriers are not any longer implied limiting factors, because of a sufficiently high thermal energy supply. Thus, these TEM *in situ* measurements enable a new level of accuracy for assessing polymorphic transformations in Ga<sub>2</sub>O<sub>3</sub>.

Received 15th November 2023,  
Accepted 8th March 2024

DOI: 10.1039/d3ma01011b

rsc.li/materials-advances

## 1. Introduction

Polymorphism contributes to the diversity of materials, so that even materials with the same chemical compositions exhibit variations in properties because of different crystal structures. Mechanistically, switching from one crystal structure to another depends on the ability of atoms to move out of their sites, which can be achieved by changing temperature and/or pressure.<sup>1</sup> Moving atoms from their equilibrium positions by accelerated particle beams is another option; even though such polymorph transitions are rare due to a competing amorphization process. Nevertheless, disorder-induced formation of metastable polymorphs by irradiation has been observed in several materials,<sup>2–5</sup> including gallium oxide (Ga<sub>2</sub>O<sub>3</sub>).<sup>6–13</sup> At present, the mechanisms of polymorph transitions in Ga<sub>2</sub>O<sub>3</sub> are under intensive investigation, because of the rapidly increasing interest in Ga<sub>2</sub>O<sub>3</sub> for its use in power electronics and UV photonics.<sup>14–19</sup> Indeed, since ion beam processing is often an inevitable part of

the device fabrication, it is essential to establish full control over the polymorphism of Ga<sub>2</sub>O<sub>3</sub>, in particular upon ion implantation and annealing. Concurrently, there is a fundamental interest in studying reconstruction of the stable polymorph, since the rich polymorphism of Ga<sub>2</sub>O<sub>3</sub> provides several potential reconstruction routes.

The present understanding is that upon the introduction of the lattice disorder by ion implantation, the thermodynamically stable  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> (monoclinic,  $C2/m^{20,21}$ ) transforms into a metastable spinel  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> (cubic,  $Fd\bar{3}m^{20,21}$ ), forming a continuous and homogeneous  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> layer on top of the  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> wafer; *i.e.* in the region containing the highest disorder induced by implanted ions and recoils.<sup>6–12</sup> In the rest of the paper, we refer to this structure as a “double polymorph  $\gamma/\beta$ -Ga<sub>2</sub>O<sub>3</sub> structure” and use it as a starting point for systematic *in situ* annealing experiments with an electron microscope, which is currently missing in the literature. Notably, until the disorder-induced  $\beta$ -to- $\gamma$  transformation in Ga<sub>2</sub>O<sub>3</sub> was explained,<sup>6–12</sup> the interpretations of the annealing experiments of ion implanted  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> samples were carried out in terms of the reduction of the “radiation disorder” in an unchanged lattice – as is conventionally done for other materials. Thus, on top of the missing *in situ* studies, even *ex situ* annealing experiments of the ion implanted  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> – until very recently – neglected the aspect of the polymorph reconstruction. Moreover, even though

<sup>a</sup> Department of Physics and Centre for Materials Science and Nanotechnology, University of Oslo, Oslo 0315, Norway. E-mail: j.g.fernandez@smn.uio.no, andrej.kuznetsov@fys.uio.no, oystein.prytz@fys.uio.no

<sup>b</sup> Department of Physics, Chalmers University of Technology, Gothenburg 41296, Sweden

† Electronic supplementary information (ESI) available. See DOI: <https://doi.org/10.1039/d3ma01011b>



the most recent *ex situ* studies revealed that  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> eventually transforms back into thermodynamically stable  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> upon annealing,<sup>6,7,22–24</sup> the kinetics of the process have not been clearly articulated. Indeed, apart from the transition temperature uncertainties, there are unresolved – and even previously unasked in literature – questions related *e.g.* to the nucleation sites of the “recrystallized”  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> as well as to its crystallographic alignment with the “initial”  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> matrix and eventually remaining  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> inclusions. Another important question is whether the chemistry of the introduced ions affects the transformation upon annealing. The question is not trivial, since it has been shown that the  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> formation as a result of the  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> disordering occurs independently of the chemical nature of the implanted ions,<sup>9,10</sup> because of the prime role of the radiation-induced atomic displacements instead of the chemical effects, at least up to a certain impurity concentration.<sup>7,11,12</sup> However, this may not necessarily hold for the transformations in double polymorph  $\gamma/\beta$ -Ga<sub>2</sub>O<sub>3</sub> structures induced by temperature.

Thus, in the present work, we monitored structural transformations happening in double  $\gamma/\beta$ -Ga<sub>2</sub>O<sub>3</sub> polymorph structures during annealing performed *in situ* with an electron microscope using samples fabricated with either Si or Au ion implants. Si implants were selected because of their technological relevance for n-type doping in Ga<sub>2</sub>O<sub>3</sub>,<sup>25</sup> while Au implants were used for comparison, investigating possible chemical effects at the background of similar levels of the initial lattice disorder in both samples. We show that, under the conditions of the *in situ* electron microscope annealing, double  $\gamma/\beta$ -Ga<sub>2</sub>O<sub>3</sub> polymorph structures remained stable up to 300 °C, when onsets of the  $\gamma$ -to- $\beta$  transformation become traceable, leading to a prominent  $\gamma$ - and  $\beta$ -mixture already at 500 °C. Interestingly, the recrystallization of the  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> occurs throughout the whole  $\gamma$ -film and the preferential alignments of the newly emerging  $\gamma/\beta$  interfaces is different from that at the initial  $\gamma/\beta$  interface formed as a result of the disorder-induced ordering. This transformation pattern is maintained as a function of temperature – with a gradually reduced portion of the  $\gamma$ -domains. Finally, at 1100 °C,  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> fully transforms into  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>, without dominating crystallographic relationships or preferred orientations, indicating that energy barriers are not any longer implied as limiting factors because of the sufficiently high thermal energy supply. No apparent trends in the transformations were attributed to the chemical effects of the implanted species. The *in situ* TEM methodology enabled a new level of accuracy in the assessments of the transformations occurring in double polymorph  $\gamma/\beta$ -Ga<sub>2</sub>O<sub>3</sub> structures.

## 2. Experimental section

Double  $\gamma/\beta$ -Ga<sub>2</sub>O<sub>3</sub> polymorph structures used in this work were fabricated by implantation of the (010)-oriented  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> single crystal wafers with either <sup>28</sup>Si<sup>+</sup> or <sup>197</sup>Au<sup>+</sup>. Specifically, 300 keV <sup>28</sup>Si<sup>+</sup> and 1.2 MeV <sup>197</sup>Au<sup>+</sup> ions were implanted with fluences of  $2 \times 10^{16}$  Si/cm<sup>2</sup> and  $3 \times 10^{15}$  Au/cm<sup>2</sup>, respectively. The implants were performed at room temperature and at 7° from the surface

normal to minimize possible channeling effects. The Si and Au implants result in 31 and 40 displacements per atom (dpa) in the maximum of the disorder peaks, respectively (see Fig. S1 and S2 in the ESI†). Similar implantation conditions were used previously in the literature for fabrication of double  $\gamma/\beta$ -Ga<sub>2</sub>O<sub>3</sub> polymorph structures.<sup>9,10,22</sup>

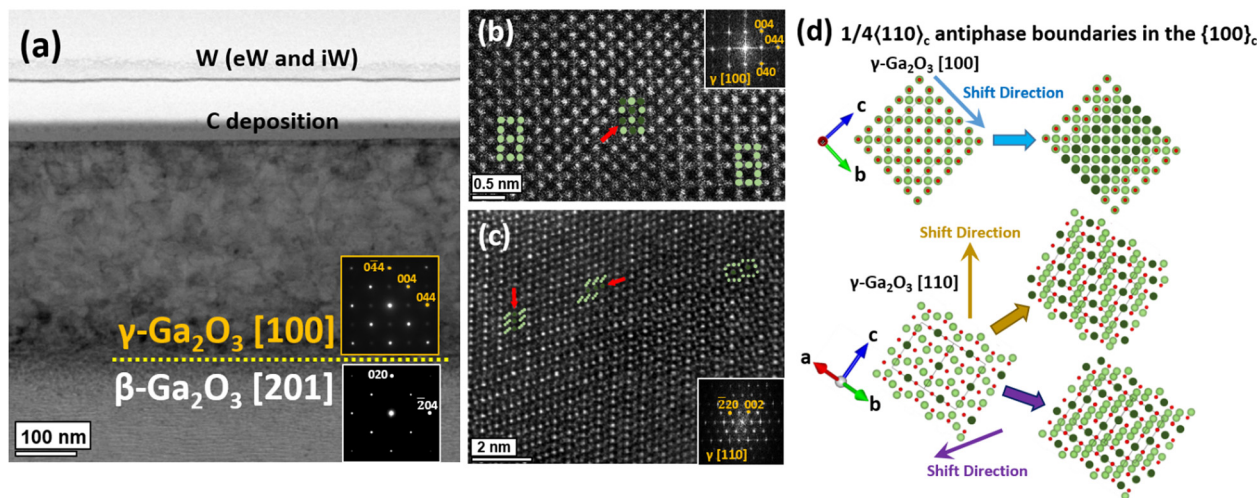
Subsequently, TEM lamellas were prepared by focused ion beam (FIB) using a JEOL JIB 4500 and a FEI Versa 3D DualBeam systems operating at 1–30 kV for the electron and Ga-ion beam. Final polishing of the TEM lamellas was done with a low ion beam energy (2 kV) and current (~20 pA) to minimize the effect of the Ga-ion beam. The crystal wafers were first sputter-coated by carbon (~20 nm) using a Leica EM ACE600 sputter system. Further ~2  $\mu$ m protective layers of C/W were deposited using an e-beam (eW), followed by a Ga-ion beam (iW). High-resolution scanning transmission electron microscopy (HRTEM/STEM) and selected area electron diffraction (SAED) were conducted using a Thermo Fisher Scientific Cs-corrected Titan G2 60–300 kV microscope operated at 300 kV and equipped with a Gatan GIF Quantum 965 spectrometer. High angle annular dark field (HAADF) and bright field (BF) STEM images were recorded using a probe convergence semi-angle of 23 mrad, a nominal camera length of 60 mm, corresponding to an inner and outer collection semi-angle of 100 and 200 mrad for HAADF and 0–22 mrad for BF, respectively. TEM *in situ* heating experiments (under vacuum) for the Si-implanted sample from RT to 850 °C in steps of 50 °C were done using a Gatan double tilt heating holder (model 652), keeping the sample at the target temperature for 30 minutes in each step. Additionally, we used a Protochips Fusion holder to repeat the same experiments for the Au implanted sample but with improved temperature control, stability, and a higher temperature range from RT to 1100 °C. As an example, Fig. S3 in the ESI† shows a low magnification BF STEM image of the prepared lamella. SAED patterns were acquired at each annealing temperature step during heating. Additional SAED patterns were acquired when cooling to RT to ensure that no crystal structure changes occurred during such cooling steps. All HAADF STEM images were recorded at RT. All the structural models for Ga<sub>2</sub>O<sub>3</sub> polymorphs in this manuscript were displayed using VESTA software.<sup>26</sup> In all lattice schematics used in the manuscript, gallium and oxygen atoms are represented in green and red, respectively. Higher color densities, darker green for gallium, are used for indicating higher atomic column density.

## 3. Results and discussion

### 3.1. As-fabricated double $\gamma/\beta$ -Ga<sub>2</sub>O<sub>3</sub> polymorph sample

Fig. 1 summarizes the TEM data collected from the as-fabricated samples. In particular Fig. 1(a) shows the low magnification BF STEM image and the corresponding SAED pattern of the Si-implanted sample, revealing a ~380 nm thick  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> layer formed as a consequence of the radiation-induced disorder. The  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> layer formed as a result of Au-implants looks similar, see Fig. S3 in the ESI† as expected because of the similar radiation disorder produced by both ions (see the





**Fig. 1** Microscopy analysis of as-implanted samples. (a) Low magnification BF STEM image and SAED patterns of the Si-implanted sample showing the formation of the double  $\gamma/\beta$ -Ga<sub>2</sub>O<sub>3</sub> polymorph structure. (b) HAADF STEM image and corresponding FFT of  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> along the [100] zone axis for the Si-implanted sample. (c) HAADF STEM image and corresponding FFT of  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> along the [110] zone axis for the Au-implanted sample. (d) Schematic representation of the lattices of  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> along [100] and [110] zone axes, together with the projected model structures and the lattice shifts as a consequence of the APB. Color code: Ga green (dark green represents high density occupancy sites), O red.

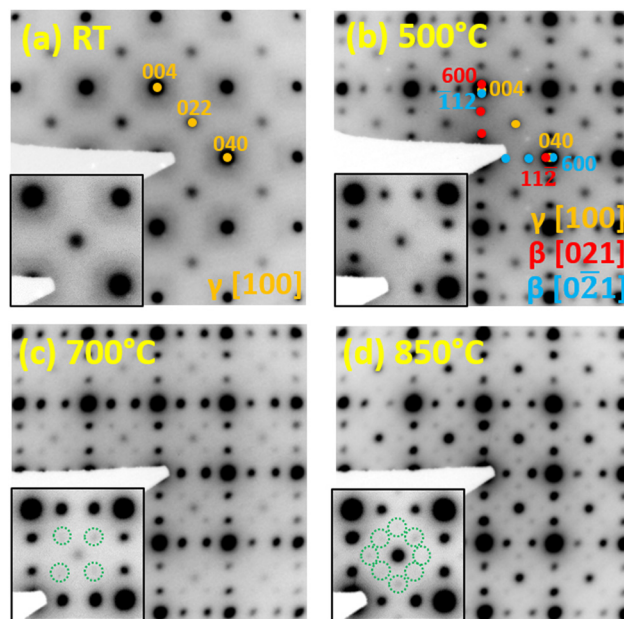
explanation for the disorder calculations in the ESI†). This is in accordance with previous literature reports concerning the disorder-induced  $\beta$ - to  $\gamma$ -polymorph transition.<sup>6–12</sup> Notably, the newly formed  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> is aligned with respect to the original  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> substrate following the crystallographic relationship: [100]  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub>//[201]  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> and [110]  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub>//[132]  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>. This crystallographic orientation is probably due to the in-plane matching (for planes vertical to the interface) between the newly formed  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> and original  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>.

There is a stacking similarity between the interplanar distances  $d_{(044)} = 1.48$  Å and  $d_{(-111)} = 4.83$  Å of  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> with  $d_{(-204)} = 1.47$  Å and  $d_{(-201)} = 4.82$  Å of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> phases, respectively. Furthermore, Fig. 1(b) and (c) show HAADF STEM images and corresponding fast Fourier transforms (FFT) of  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> along the [100] zone axis for the Si-implanted sample, and along the [110] zone axis for the Au-implanted sample. Both images show the presence of antiphase boundaries (APB), indicated by red arrows, that are characterized by the displacement vector  $1/4\langle 110 \rangle c$  in the  $\{100\}c$  planes of the cubic spinel structure. Consequently, Fig. 1(d) shows two schematic representations of the  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> lattice (viewed along the [100] and [110] directions) including the formation of APBs as a result of the atomic shifts. These schematics are in good agreement with the contrast observed in the HAADF STEM images in Fig. 1(b) and (c). Similar APBs in the  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> spinel structure have been reported in the literature.<sup>8,27–29</sup> Thus, in the initial stage, the samples used in this study are representative examples of the double  $\gamma/\beta$ -Ga<sub>2</sub>O<sub>3</sub> polymorph structures. No apparent differences were detected in the quality of the  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> produced either by Si or Au implants. We used both samples in the *in situ* study to check for possible differences at the annealing stage.

### 3.2. *In situ* monitoring of the $\gamma$ - to $\beta$ -polymorph transition

Fig. 2(a)–(d) display the evolution in the SAED patterns for the Si-implanted sample, at RT and after *in situ* annealing at 500 °C,

700 °C, and 850 °C, respectively. The  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> layer was initially tilted to the [100] zone axis at RT, as shown in the indexed SAED pattern in Fig. 2(a). Notably, each SAED pattern is accompanied by an inset showing the scaled features around the (022) spot of the  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> to highlight the structural changes. We observed an onset  $\gamma$ - to  $\beta$ -polymorph transition already at  $\sim 300$  °C (not observed at 250 °C, see Fig. S4 in the ESI†) and prominently occurring at 500 °C. The formation of the  $\beta$ -phase in the



**Fig. 2** Comparison of the SAED patterns from the Si-implanted sample in (a) the as-fabricated state, with that upon *in situ* annealing at (b) 500 °C, (c) 700 °C, and (d) 850 °C. The insets show the scaled-up features around the (022) spot. Color code: orange:  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> along the [100] zone axis, red and blue:  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> along the [021] and  $[0\bar{2}1]$  zone axes, respectively.



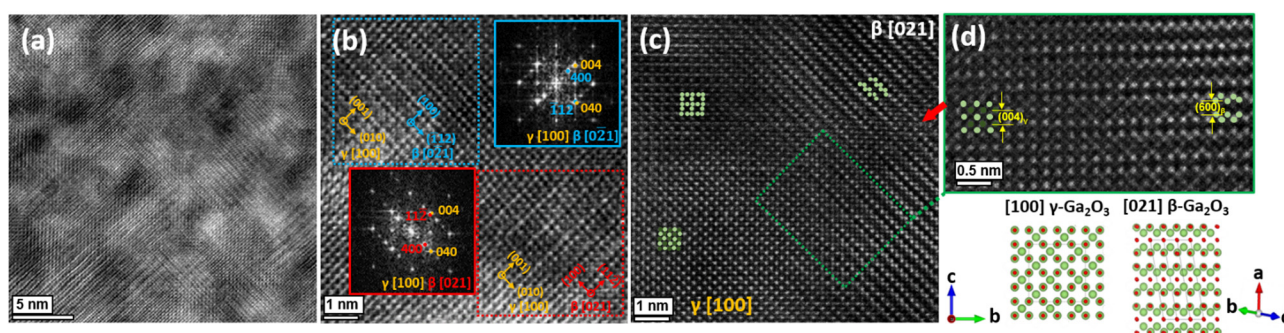


$\gamma$ -Ga<sub>2</sub>O<sub>3</sub> layer is confirmed by the emergence of new diffraction spots in the SAED pattern in Fig. 2(b). These new spots can be indexed as  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> phase domains (represented by red and blue colors) along the [021] and [0 $\bar{2}$ 1] zone axes, which are related by a 90° rotation. This could be induced by the four-fold symmetry of  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> along the [100] direction, which may direct the formation of  $\beta$ -domains with a 90° rotational relationship. Subsequently, upon reaching 700 °C, weaker spots surrounding the (022) spot become visible, as indicated by green dashed circles in the inset in Fig. 2(c). The number of these spots increases upon reaching 850 °C, see the inset in Fig. 2(d). Additionally, the transformation from  $\gamma$ - to  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> is confirmed by EELS spectra in Fig. S5 in the ESI† Fig. S5 (ESI†) shows the evolution of the O-K edge for the as-implanted  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> and after annealing at 500 °C and 850 °C, together with bulk  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>, as a reference. The O K-edge is characterized by two peaks at 537 eV and 543 eV,<sup>10,24</sup> and the relative intensity between these peaks can be used as a fingerprint of the corresponding polymorph. Importantly, we did not observe any change in the sample while cooling the samples to room temperature, confirming the stability of the transformation, and enabling accurate high-resolution measurements.

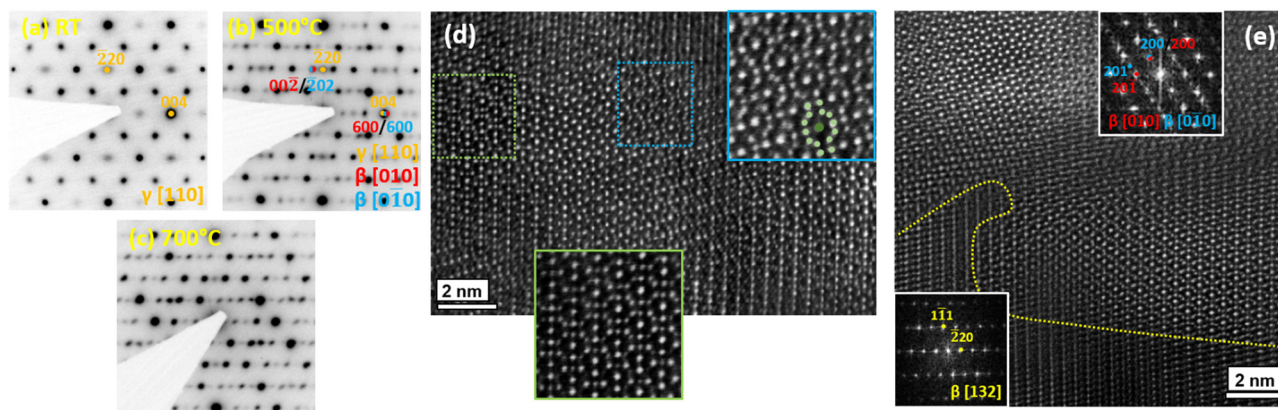
Indeed, Fig. 3 provides an example of such atomic-resolution measurements for samples annealed at 500 °C and 850 °C. Specifically, Fig. 3(a) and (b) depict low and high magnification HAADF STEM images of the Si-implanted sample after heating at 500 °C. In both images, two types of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> domains can be observed, oriented along the [021] and [0 $\bar{2}$ 1] directions and related to each other by a 90° rotation, marked in blue and red. Importantly, Fig. 3(b) is accompanied by two extra panels, showing the FFTs obtained from the areas marked by red and blue dashed rectangles. These FFTs confirm the presence of remaining  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> in the [100] orientation and reveal the presence of individual  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> domains exhibiting the [021]/[0 $\bar{2}$ 1] orientations with a 90° rotational relationship as described above. Fig. 3(c) displays a HAADF STEM image of the sample after heating at 850 °C. On the left-hand side in Fig. 3(c), the remaining  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> is still visible, together with the presence of APBs, identified already in the as-fabricated samples. On the

right-hand side of Fig. 3(c), the contrast associated with  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> along the [021] zone axis is clearly observed. Additionally, the presence of APBs between two  $\beta$ -phase regions is visible in Fig. 3(c), as indicated by the red arrow. Thus, Fig. 3(c), in combination with the SAED pattern (Fig. 2(d)), confirms the presence of remaining  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> upon 850 °C annealing. For clarity, Fig. 3(d) provides a detailed view of the interface between the [100]  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> (left) and [021]  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> (right) phases, taken from the region marked with a green rectangle in Fig. 3(c), accompanied by schematic representations of both phases. As can be deduced from Fig. 3(d), the  $\gamma$ - to  $\beta$ -polymorph transition proceeds *via* the stacking of {400}  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> to {600}  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> planes. We suggest that the newly formed strong spots observed in Fig. 2(b)–(d) correspond to the formation of new  $\beta$ -phase along the [021] direction, while the weaker spots marked with dashed green in Fig. 2(c) and (d) could be a consequence of the presence of domains and extended defects resulting from the recrystallization of the  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> phase, consistently with the literature.<sup>30–32</sup> Recently, the presence of four rotational domains of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> related by a 90° rotation has been reported for  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> grown on top of  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub>.<sup>33</sup> Thus, in our samples, cubic  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> may play a similar role, promoting the domains formation.

To control the generic character of the  $\gamma$ - to  $\beta$ -polymorph transition observed in Fig. 2 and 3, we repeated the experiment with the Au-implanted sample. This also allows us to explore the different crystallographic orientation between  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> and the newly formed  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>. Firstly, Fig. S4 in the ESI† shows an HRTEM image and corresponding FFT from different areas after annealing at 300 °C. Upon careful inspection in some of the FFTs, we observe the formation of additional weak spots assigned to  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> inclusions, confirming the  $\gamma$ - to  $\beta$ -polymorph transition starts already ~300 °C (these inclusions were not observed at 250 °C). Furthermore, Fig. 4(a)–(c) display the SAED patterns of the Au-implanted sample acquired at RT as well as after heating at 500 °C and 700 °C. The SAED pattern acquired at RT can be indexed to the  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> phase along the [110] zone axis (see the orange indexing in Fig. 4(a)). After heating, additional spots emerge in the SAED pattern marked by red and blue spots in Fig. 4(b). Similar to the Si-implanted



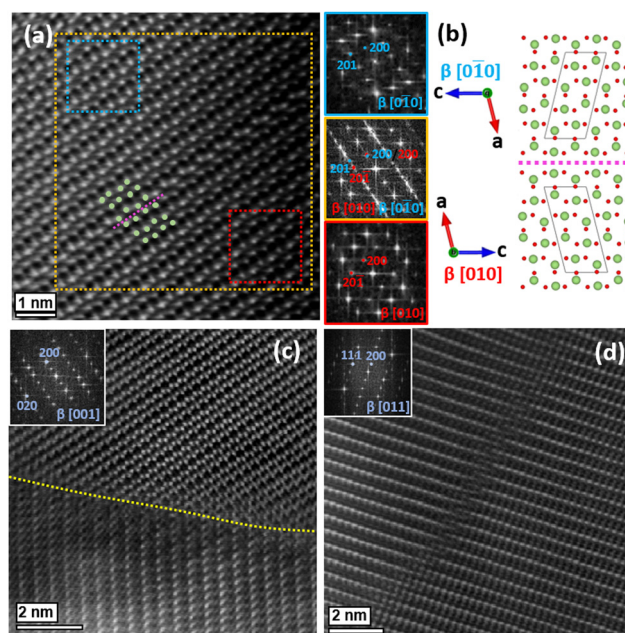
**Fig. 3** Atomic-resolution measurements of the Si-implanted sample upon annealing at 500 °C and 850 °C. (a) Low and (b) high magnification HAADF STEM images and corresponding FFTs of the Si-implanted sample after heating at 500 °C. (c) Atomic resolution HAADF STEM images of the same sample after heating at 850 °C showing the presence of [100]  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> (left) and [021]  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> (right) with the presence of extensive defects such as APBs (red arrow). (d) HAADF STEM image at the interface between the two polymorphs. Schematic representation of [100]  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> and [021]  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> polymorphs. Color code: Ga green (dark green represents high density occupancy sites), O red.



**Fig. 4** Complementary microscopy data measured on the Au-implanted sample. SAED patterns for (a) as-fabricated state, with that upon *in situ* annealing at (b) 500 °C and (c) 700 °C. Color code: orange:  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> along the [110] zone axis, red and blue:  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> along the [010] and  $[0\bar{1}0]$  zone axes. (d) and (e) HAADF STEM images and corresponding FFT after heating at 500 °C and 700 °C, respectively. The green and blue boxes in figure (d) correspond to two domains of the [010]  $\beta$ -phase. Color code: Ga green (dark green represents high density occupancy sites). The dashed yellow line in figure (e) indicates the interface between the original (untransformed)  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> and the newly formed  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>.

sample, these new spots can be indexed as two separate  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> domains, in this case, oriented along the [010] and  $[0\bar{1}0]$  zone axes. This observation is in full agreement and correlation with the data obtained for the Si-implanted sample acquired from a different zone axis. This is because the angle between [100] and [110] directions in the  $\gamma$ -phase (45°) is similar to that for [021] and [010] directions in the  $\beta$ -phase (43.68°). As shown in the SAED pattern, like in the previous case, there is perfect overlapping between the (004) plane of  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> with the (600) plane of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>. The HAADF STEM image recorded after heating at 500 °C (Fig. 4(d)) displays the remaining  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> phase, as well as the newly formed [010]  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> domains in the presence of the stacking faults which have been observed previously in the literature.<sup>34</sup> The green and blue squares in Fig. 4(d) show details of the defect associated with the newly formed  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> phase. Finally, Fig. 4(e) shows a representative HAADF STEM image and corresponding FFTs taken at the interface, indicated by a dashed yellow line, between the newly formed  $\beta$ -phase along the [010] zone axis and the original (untransformed)  $\beta$ -phase along the [132] zone axis after heating at 700 °C, confirming that the crystallographic relationship between  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> and the original  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> is different from the crystallographic relationship between  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> and the recrystallized  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>. Thus, taking into account data in Fig. 2–4, we conclude that the  $\gamma$ - to  $\beta$ -polymorph transition in the double  $\gamma/\beta$ -polymorph structure fabricated *via* Si and Au ion implantation occurs in a generic and analogous manner, at least for the dose range and *in situ* TEM annealing conditions (under vacuum) used in our experiment.

Furthermore, to investigate whether the  $\beta$ -matrix could be recovered with high temperature annealing, we performed *in situ* heating of the Au-implanted sample up to 1100 °C. Notably, after annealing at 1100 °C, the  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> has been completely transformed to  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> in other orientations. For example, Fig. 5(a) shows an area of the TEM lamella containing newly formed  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> viewed along the [010] direction. In addition, the corresponding FFT of two domains (blue and



**Fig. 5** HAADF STEM images of the Au-implanted sample after heating at 1100 °C. (a)  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> along the [010] direction with corresponding global FFT (orange) and local FFT taken from different areas (red and blue boxes) corresponding to two different domains. (b) Schematic representation of both domains and the twin boundary indicated by the pink dashed line. Ga atomic columns are represented in green and O in red. (c) Image recorded at the interface (indicated by a dashed yellow line) between newly formed [010]  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> and original (untransformed) [132]  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>. (d) Image recorded from another area of the lamella corresponding to newly formed [111]  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> containing APB.

red) and the global FFT from the region marked in orange are included as insets. The FFTs obtained from the blue and red areas confirm the presence of two domains of [010] and  $[0\bar{1}0]$   $\beta$ -Ga<sub>2</sub>O<sub>3</sub> separated by a twin boundary indicated by the pink dashed line. For clarity, Fig. 5(b) displays a schematic representation of [010]  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> resulting in the twin boundary observed





in Fig. 5(a). Similar results were previously observed in  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> domains related by a twin boundary in the literature.<sup>31,35,36</sup>

Nevertheless, we found that the crystallographic relationship deduced from the SAED patterns in Fig. 4(b) and (c) was no longer maintained. Fig. 5(c) shows a HAADF STEM image of the interface between the original and the newly formed  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>, with the interface marked by the dashed yellow line. Here, the original (untransformed)  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> is viewed along the [132] zone axis, while the newly formed  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> is now viewed along the [001] zone axis instead. Additionally, Fig. 5(d) displays a HAADF STEM image of the newly formed  $\beta$ -phase from another area of the lamella at the same tilting angle, oriented in the [011] direction, and revealing the presence of an APB. In both cases, this crystallographic relationship is different from that initially observed for lower annealing temperatures, as in Fig. 4(e) at 700 °C.

### 3.3. Discussion

The data in Fig. 1–5 provide new insights on how the metastable  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub>, formed by the disorder-induced transformation, converts back into  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> upon annealing. In our interpretation, under the annealing conditions used in this work,  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> forms domains exhibiting the crystallographic relationship of [021]  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>//[100]  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> and [010]  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>//[110]  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub>. Importantly, looking accurately into the SAED patterns, *e.g.* in Fig. 2(b) and 3(b), we observe that some of the interplanar distances in the  $\gamma$ - and  $\beta$ -polymorphs are rather similar. Specifically, we see that the lattice spacing of 2.06 Å for (400) planes in the  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> is close to that of 1.98 Å and 2.10 Å for (600) and (112) planes in the  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>, respectively. Such similarity in distances may contribute to the reduction of the strain barrier associated with the recrystallization of the  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> domains inside the  $\gamma$ -film. Our results also showed that the nucleation of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> domains takes place all over the  $\gamma$ -layer, and not in a preferential location, such as at the  $\gamma/\beta$ -interface. Spectacularly, these newly formed  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> domains exhibit different preferential crystallographic stacking with the  $\gamma$ -lattice as compared to that in the initial  $\gamma/\beta$ -interface. Recent reports<sup>37,38</sup> suggest the existence of similarities between the two polymorphs: (i) both phases exhibit a face-centered cubic oxygen lattice, and (ii) the similarity in the cations positions (similar Ga sites in both lattices and Ga sites of the  $\gamma$ -phase that correspond to Ga interstitial sites in the  $\beta$ -phase). Considering this, the  $\gamma$ -phase was described as a disordered version of the  $\beta$ -phase, where cations from regular sites on the  $\beta$ -lattice are displaced to interstitial sites, leaving a vacancy. Therefore, this phase transition can be considered as a disorder–order

transition induced by temperature. For clarity, Fig. 6 illustrates the dominating orientation relationships observed in our samples.

Notably, the formation of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> domains inside the  $\gamma$ -matrix upon *ex situ* annealing was proposed in the literature too, specifically suggesting that the defects available in the  $\gamma$ -matrix may induce the formation of such  $\beta$ -domains at temperatures below those require for a massive  $\gamma$ -to- $\beta$  transition.<sup>23</sup> This hypothesis is in agreement with our data, where high concentrations of APBs in the  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> triggers the formation of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> already at 300 °C. Further on, for the annealing in the range of 900–1100 °C, the remaining  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> back-transforms to  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> without preferential orientations, indicating that energy barriers are no longer implied as limiting factors because of a high enough thermal energy supplied.

Importantly, the parameters controlling the polymorphism in a macroscopic sample (typically used in the *ex situ* annealing experiments in the literature<sup>8,22</sup>) and in a TEM-lamella sized sample (as used in this work) – may be different. Consequently, there could be potential discrepancies between the *in situ* and *ex situ* observations. In this sense, J. Wang *et al.*<sup>39</sup> showed experimentally and by first-principles calculations that similar transitions occur in Ga<sub>2</sub>O<sub>3</sub> nanowires. From their results, the calculated surface energy of {600} $\beta$  is smaller than that of {400} $\gamma$ . Considering that a TEM-lamella has a higher surface area/volume ratio than a bulk sample, the observed transition in this work could be influenced by the reduction of the surface energy. Moreover, all *in situ* annealing in the electron microscope was performed in high vacuum, featuring quite different conditions as compared to that in air or gas flow frequently used in the literature for *ex situ* annealing. In its turn, it is well-known that the differentiation between the oxygen rich or lean conditions may radically change properties of oxides.<sup>40</sup> Thus, this factor might be considered while comparing *in situ* and *ex situ* observations too. For example, in our previous study of *ex situ* air annealing of the double polymorph  $\gamma/\beta$ -Ga<sub>2</sub>O<sub>3</sub> structure fabricated using similar Si ion implantation conditions,<sup>22</sup> the subsequent crystallographic relationship between the two phases was different. Moreover, in the present work, we have not detected precipitations correlated with the implants, *e.g.* in the form of SiO<sub>x</sub> nanoparticles nor voids, which were also observed during air annealing in our previous work.<sup>22</sup> All-in-all, the results in Fig. 1–5, taken in the context of the literature data, confirm the complexity of the processes occurring in the double polymorph  $\gamma/\beta$ -Ga<sub>2</sub>O<sub>3</sub> structures upon annealing.

## 4. Conclusions

In conclusion, we performed *in situ* TEM annealing studies of the double  $\gamma/\beta$ -Ga<sub>2</sub>O<sub>3</sub> polymorph structures fabricated by the radiation induced disorder approach and revealed several new interesting trends for the  $\gamma$ -to- $\beta$  polymorph transition. The *in situ* TEM study reveals the emergence of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> domains at temperatures around 300 °C. The nucleation of these newly formed  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> domains take place according to the crystallographic

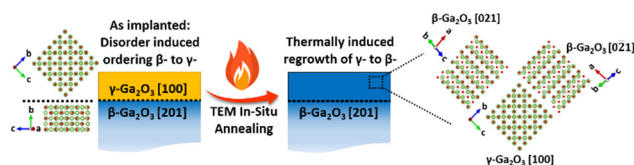


Fig. 6 Schematics of the dominating orientation relationships observed in our samples.



relationship: [100]  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub>/[021]  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> and [110]  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub>/[010]  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>, possibly explicable by the reduction of the surface energy and strain minimization due to very similar lattice plane spacings in the polymorphs. This was interpreted in terms of the reduction of the strain barrier associated with the formation of the  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> domains inside the  $\gamma$ -film. Importantly, these orientations are different from that established at the initial as-implanted  $\gamma/\beta$ -interface. However, at high temperatures (900–1100 °C),  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> fully recrystallizes into  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> in other orientations, indicating that energy barriers are not any longer implied as limiting factors because of a sufficiently high thermal energy supply. Thus, the *in situ* observation methodology enabled a novel unprecedented accuracy in the assessments of the  $\gamma$ -to- $\beta$  polymorph transition.

## Author contributions

J. G. F. and S. B. K.: conceptualization, data curation, formal analysis, investigation, methodology, visualization, writing original draft, writing – review and editing. L. J. Z., A. A., A. P., P. D. N.: formal analysis, investigation, methodology, visualization, writing original draft, writing – review and editing. E. O., L. V., A. K. and Ø. P.: formal analysis, funding acquisition, investigation, project administration, resources, supervision, validation, writing – review & editing.

## Conflicts of interest

There are no conflicts to declare.

## Acknowledgements

The Research Council of Norway is acknowledged for support through the projects NANO2021 researcher project Functionalization of conducting oxides by ion beam and defect engineering (FUNCTION, no. 287729) and the M-ERA.NET Program *via* GOFIB project (no. 337627). We also acknowledge the Norwegian Center for Transmission Electron Microscopy (NORTEM, no. 197405/F50), the Norwegian Micro- and Nanofabrication Facility (NorFab, no. 295864), and NordForsk is acknowledged for support to the NordTEMhub project no. 96971. Viktor Mihaly Bobal is acknowledged for performing ion implantation on the samples studied herein. This work was performed in part at the Chalmers Material Analysis Laboratory, CMAL. The authors acknowledge the financial support from Swedish Research Council (VR) and Swedish Foundation for Strategic Research (SSF) for the access to ARTEMI, the Swedish National Infrastructure in Advanced Electron Microscopy (2021-00171 and RIF21-0026).

## References

- W. Li, X. Qian and J. Li, *Nat. Rev. Mater.*, 2021, **6**, 829.
- M. Ishimaru, I. V. Afanasyev-Charkin and K. E. Sickafus, *Appl. Phys. Lett.*, 2000, **76**, 2556.
- S. Dhara, A. Datta, C. T. Wu, Z. H. Lan, K. H. Chen, Y. L. Wang, C. W. Hsu, C. H. Shen, L. C. Chen and C. C. Chen, *Appl. Phys. Lett.*, 2004, **84**, 5473.
- A. Benyagoub, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2005, **72**, 094114.
- G. Sattonnay and L. Thomé, *J. Nucl. Mater.*, 2006, **348**, 223.
- K. R. Gann, N. Pieczulewski, C. A. Gorsak, K. Heinselman, T. J. Asel, B. A. Noesges, K. T. Smith, D. M. Dryden, H. G. Xing, H. P. Nair, D. A. Muller and M. O. Thompson, *J. Appl. Phys.*, 2024, **135**, 015302.
- H.-L. Huang, C. Chae, J. M. Johnson, A. Senckowski, S. Sharma, U. Singiseti, M. H. Wong and J. Hwang, *APL Mater.*, 2023, **11**, 061113.
- T. Yoo, X. Xia, F. Ren, A. Jacobs, M. J. Tadjer, S. Pearton and H. Kim, *Appl. Phys. Lett.*, 2022, **121**, 072111.
- J. García-Fernández, S. B. Kjeldby, P. D. Nguyen, O. B. Karlsen, L. Vines and Ø. Prytz, *Appl. Phys. Lett.*, 2022, **121**, 191601.
- A. Azarov, J. G. Fernández, J. Zhao, F. Djurabekova, H. He, R. He, Ø. Prytz, L. Vines, U. Bektas, P. Chekhonin, N. Klingner, G. Hlawacek and A. Kuznetsov, *Nat. Commun.*, 2023, **14**, 4855.
- H.-L. Huang, J. M. Johnson, C. Chae, A. Senckowski, M. H. Wong and J. Hwang, *Appl. Phys. Lett.*, 2023, **122**, 251602.
- J. Zhao, J. G. Fernández, A. Azarov, R. He, Ø. Prytz, K. Nordlund, M. Hua, F. Djurabekova and A. Kuznetsov, *arXiv*, 2024, preprint, arXiv:2401.07675, DOI: [10.48550/arXiv.2401.07675](https://doi.org/10.48550/arXiv.2401.07675).
- A. Azarov, C. Bazioti, V. Venkatachalapathy, P. Vajeeston, E. Monakhov and A. Kuznetsov, *Phys. Rev. Lett.*, 2022, **128**, 015704.
- S. J. Pearton, F. Ren, M. Tadjer and J. Kim, *J. Appl. Phys.*, 2018, **124**, 220901.
- M. J. Tadjer, *Science*, 2022, **378**, 724.
- J. Xu, W. Zheng and F. Huang, *J. Mater. Chem. C*, 2019, **7**, 8753.
- W. Choi, D.-W. Jeon, J.-H. Park, D. Lee, S. Lee, K. H. Baik and J. Kim, *Mater. Adv.*, 2023, **4**, 4520.
- B. Zhao, F. Wang, H. Chen, L. Zheng, L. Su, D. Zhao and X. Fang, *Adv. Funct. Mater.*, 2017, **27**, 1700264.
- D. Zhang, H. Yu, G. You, G. Shao, Z. Fang, Z. Liang, T. Zhang, H. Hou, L. Wang, Q. Chen, J. Teng and W. Yang, *J. Mater. Sci. Technol.*, 2023, **163**, 150.
- M. Bosi, P. Mazzolini, L. Seravalli and R. Fornari, *J. Mater. Chem. C*, 2020, **8**, 10975.
- H. Y. Playford, A. C. Hannon, E. R. Barney and R. I. Walton, *Chem. – Eur. J.*, 2013, **19**, 2803.
- S. B. Kjeldby, A. Azarov, P. D. Nguyen, V. Venkatachalapathy, R. Mikšová, A. Macková, A. Kuznetsov, Ø. Prytz and L. Vines, *J. Appl. Phys.*, 2022, **131**, 125701.
- P. Castro-Fernández, M. V. Blanco, R. Verel, E. Willinger, A. Fedorov, P. M. Abdala and C. R. Müller, *J. Phys. Chem. C*, 2020, **124**, 20578.
- A. Sharma, M. Varshney, H.-J. Shin, K. H. Chae and S. O. Won, *RSC Adv.*, 2017, **7**, 52543.
- A. T. Neal, S. Mou, S. Rafique, H. Zhao, E. Ahmadi, J. S. Speck, K. T. Stevens, J. D. Blevins, D. B. Thomson, N. Moser, K. D. Chabak and G. H. Jessen, *Appl. Phys. Lett.*, 2018, **113**, 062101.



- 26 K. Momma and F. Izumi, *J. Appl. Crystallogr.*, 2011, **44**, 1272.
- 27 M. Mitome, S. Kohiki, T. Nagai, K. Kurashima, K. Kimoto and Y. Bando, *Cryst. Growth Des.*, 2013, **13**, 3577.
- 28 C. S. Chang, N. Tanen, V. Protasenko, T. J. Asel, S. Mou, H. G. Xing, D. Jena and D. A. Muller, *APL Mater.*, 2021, **9**, 051119.
- 29 J. Tang, K. Jiang, S. D. House, C. Xu, K. Xiao, L. M. Porter and R. F. Davis, *Appl. Phys. Lett.*, 2023, **123**, 012103.
- 30 P. Mazzolini, A. Falkenstein, C. Wouters, R. Schewski, T. Markurt, Z. Galazka, M. Martin, M. Albrecht and O. Bierwagen, *APL Mater.*, 2020, **8**, 011107.
- 31 G. Wagner, M. Baldini, D. Gogova, M. Schmidbauer, R. Schewski, M. Albrecht, Z. Galazka, D. Klimm and R. Fornari, *Phys. Status Solidi A*, 2014, **211**, 27.
- 32 S. Gao, Y. Wu, R. Kang and H. Huang, *Mater. Sci. Semicond. Process.*, 2018, **79**, 165.
- 33 K. Jiang, J. Tang, C. Xu, K. Xiao, R. F. Davis and L. M. Porter, *J. Vac. Sci. Technol., A*, 2023, **41**, 062702.
- 34 D. Gogova, G. Wagner, M. Baldini, M. Schmidbauer, K. Irmscher, R. Schewski, Z. Galazka, M. Albrecht and R. Fornari, *J. Cryst. Growth*, 2014, **401**, 665.
- 35 R. Schewski, K. Lion, A. Fiedler, C. Wouters, A. Popp, S. V. Levchenko, T. Schulz, M. Schmidbauer, S. Bin Anooz, R. Grüneberg, Z. Galazka, G. Wagner, K. Irmscher, M. Scheffler, C. Draxl and M. Albrecht, *APL Mater.*, 2019, **7**, 022515.
- 36 S. Vura, U. Ul Muazzam, V. Kumar, S. C. Vanjari, R. Muralidharan, N. Digbijoy, P. Nukala and S. Raghavan, *ACS Appl. Electron. Mater.*, 2022, **4**, 1619.
- 37 C. Wouters, M. Nofal, P. Mazzolini, J. Zhang, T. Remmele, A. Kwasniewski, O. Bierwagen and M. Albrecht, *APL Mater.*, 2024, **12**, 011110.
- 38 J. Tang, K. Jiang, C. Xu, M. J. Cabral, K. Xiao, L. M. Porter and R. F. Davis, *APL Mater.*, 2024, **12**, 011109.
- 39 J. Wang, X. Guan, H. Zheng, L. Zhao, R. Jiang, P. Zhao, Y. Zhang, J. Hu, P. Li, S. Jia and J. Wang, *Nano Lett.*, 2023, **23**, 7364.
- 40 T. M. Børseth, B. G. Svensson, A. Y. Kuznetsov, P. Klason, Q. X. Zhao and M. Willander, *Appl. Phys. Lett.*, 2006, **89**, 262112.

