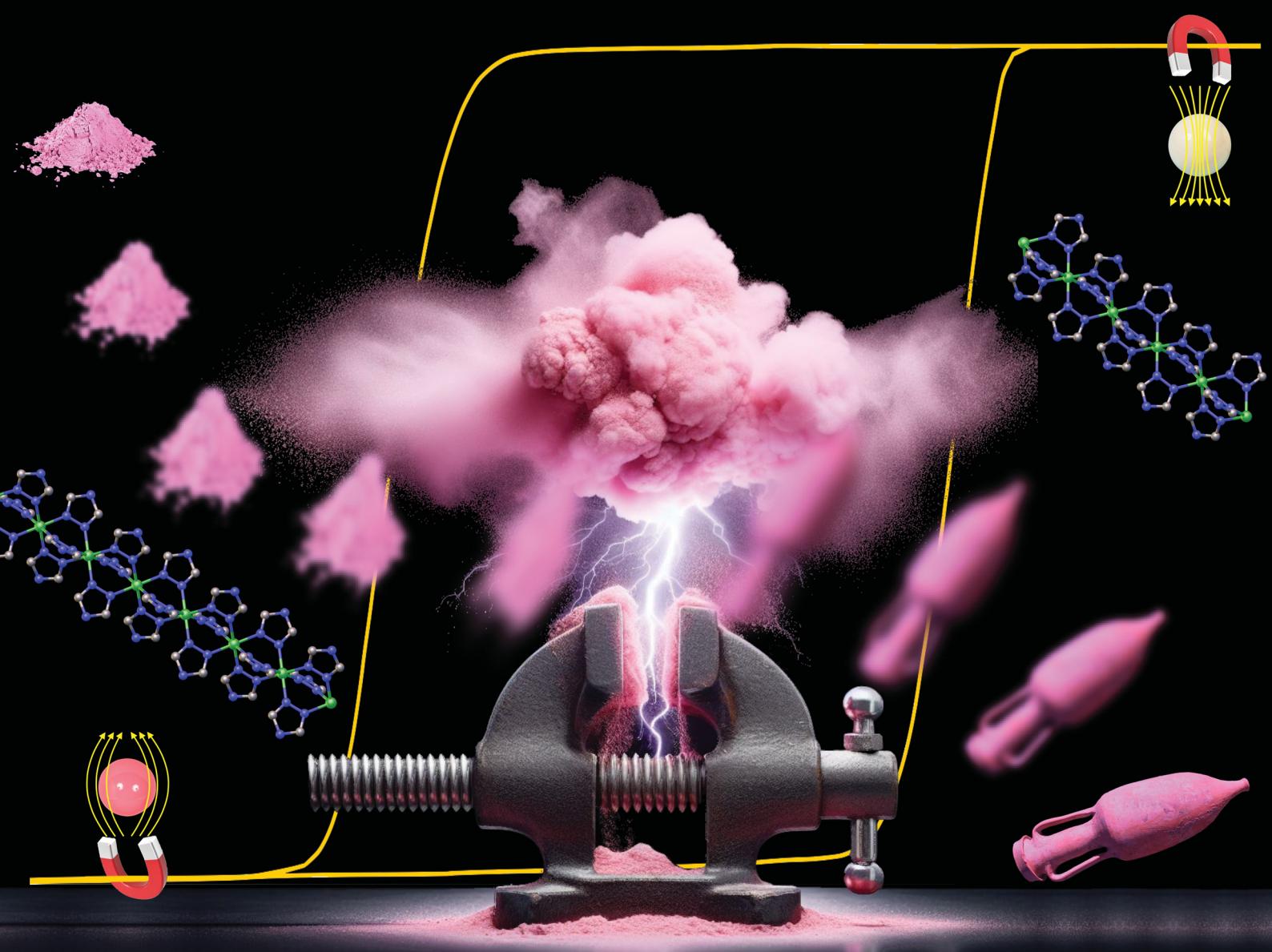


Materials Advances

rsc.li/materials-advances



ISSN 2633-5409

PAPER

Michael Josse, Mathieu Marchivie *et al.*
Spin crossover molecular ceramics by Cool-SPS:
consequences on switching features beyond
microstructural effects



Cite this: *Mater. Adv.*, 2024,
5, 1502

Spin crossover molecular ceramics by Cool-SPS: consequences on switching features beyond microstructural effects[†]

Liza El-Khoury,^a Nathalie Daro,^{ID a} Guillaume Chastanet,^{ID a} Patrick Rosa,^{ID a} Dominique Denux,^a Laetitia Etienne,^a Vincent Mazel,^b Michael Josse^{*a} and Mathieu Marchivie^{ID *a}

The sintering of spin crossover materials using spark plasma sintering at low temperature (Cool-SPS) provides a new way of shaping such compounds into functional molecular ceramics. These ceramics reach a high relative density of 95%, which may address several issues for using spin crossover materials in barocaloric devices. Starting from the reference complex $[\text{Fe}(\text{Htrz})_2(\text{trz})](\text{BF}_4)$, we first investigated the magnetic, structural, and microstructural properties as well as the fatigability behavior of the starting powder using multiple magnetic measurements, X-ray diffraction and calorimetry to compare them with the elaborated ceramics. The best conditions of pressure and temperature during the SPS process to obtain reproducible molecular ceramics with high relative density were found to be between 250 and 300 °C, and 300 and 400 MPa. The same complete set of characterization performed on a molecular ceramic of 95% of relative density reveals that the crystal structures, as well as the abrupt hysteretic SCO of $[\text{Fe}(\text{Htrz})_2(\text{trz})]\text{BF}_4$, are perfectly conserved after sintering. However, ceramics present faster stabilization of their microstructural and magnetic properties upon cycling and a higher cooperativity at the macroscopic level was observed compared to the starting powder.

Received 9th September 2023,
Accepted 18th October 2023

DOI: 10.1039/d3ma00688c

rsc.li/materials-advances

Introduction

Ceramics have been widely used since antiquity to transform ores and compounds into functional materials. Nowadays, ceramics are probably one of the most used form to shape and investigate solids into functional materials. They are present in numerous high technology applications such as microelectronics and are widely studied to design smart materials in the emergent green economy through energy storage applications, solar cells or caloric materials.^{1–4} Ceramics are mainly based on tough compounds such as inorganic metal oxides. Fragile and low density materials, on the other hand, are very scarcely elaborated as ceramics, despite great interest, mainly due to their limited thermal stability preventing the application of conventional sintering processes. During the last decade, however, low temperature sintering methods started to grow^{5–8} and

very recently ceramics of thermodynamically fragile materials were obtained using Cool-SPS^{9–11} (low temperature SPS, where SPS = Spark Plasma Sintering), leading to the conception of dense and cohesive polycrystalline molecular materials, *i.e.* molecular ceramics. While the sintering process is not fully understood,¹⁰ this novel approach opens a new way for shaping compounds belonging to the very wide and versatile world of functional molecular materials into ceramics.

Among them, molecular switches constitute a promising class of materials that can be switched between two or more different states upon an external stimulus. Due to their high versatility interest in them keeps growing in the scientific community owing to numerous potential applications like memory devices, sensors, molecular motors, actuators or quantum materials.^{12–16} Spin cross-over (SCO) complexes represent one of the very promising families of such molecular switches where the switching states, separated by a small energy gap, correspond to two different electronic configurations with distinct spin repartitions.^{17–22} Considering octahedral complexes of the Fe^{2+} cation, which represent the most reported cases, the two configurations correspond to a low spin state (LS) of electronic configuration $t_{2g}^6e_g^0$ ($S = 0$) and a high spin state (HS) of electronic configuration $t_{2g}^4e_g^2$ ($S = 2$). The crossover between these two states can be triggered by numerous stimuli

^a Univ. Bordeaux, CNRS, Bordeaux INP, ICMCB, UMR 5026, F-33600 Pessac, France. E-mail: mathieu.marchivie@icmcb.cnrs.fr, michael.josse@icmcb.cnrs.fr

^b Univ. Bordeaux, CNRS, I2M Bordeaux, 146 rue Léon Saignat, F-33000 Bordeaux, France

† Electronic supplementary information (ESI) available: Additional information such as sintering maps, diametral compression test of pellets, and complementary figures can be found in an additional pdf file. See DOI: <https://doi.org/10.1039/d3ma00688c>



such as temperature, pressure, magnetic field, light irradiation, electric field^{23–28} or chemical sorption.^{9–31} The SCO signature can be probed by several physical property changes such as magnetic susceptibility, optical measurements, vibrational spectroscopies, structural changes,³² thermodynamic properties³³ and more scarcely studied changes in conductivity or dielectric constant.^{34–36} These last effects may be very interesting concerning potential applications in electronic devices but require difficult measurements on single crystals, compressed powders or thin layers.

Due to the high versatility of their possible chemical composition and their wide structural diversity, SCO materials show a very wide panel of switching behaviours from gradual to very abrupt transitions with or without a hysteresis loop. Besides, the salient role of the crystal structure on the SCO feature is now well established.^{32,37,38} Notably, intermolecular interactions have been widely studied and the hydrogen bonding network, pi-pi stacking³⁹ and even H-H van der Waals contacts have been proven to play a crucial role in cooperativity and hysteresis loop width.^{40,41} Beyond the crystal structure, the shaping of SCO complexes was investigated for several years to form nanomaterials,^{19,42} thin or thick films,^{43,44} submonolayers,^{45,46} composites with polymers.⁴⁷ Most of these material forms lead to modification of the SCO features, which shed light on the multi-scale nature of the SCO properties and triggered new approaches to this phenomenon.⁴⁸ From atoms to bonds, to coherent domains, links were established that related macroscopic features of SCO to the contributions of the aforementioned microscopic phenomena. However, throughout the various characteristic scales of materials, the influence of the microstructure, with features such as grain boundaries, is less investigated which may be due to the lack of controlled ways to modify microstructural characteristics in molecular systems. Molecular ceramics could contribute to addressing this issue. Additionally, SCO materials have been recently identified as very promising systems for barocaloric applications provided their low density and low thermal conductivity problems are overcome.^{49,50} Shaping SCO compounds into ceramics could address these concerns. Nevertheless, although ceramics are ubiquitous in solid state chemistry, it is worth noting that ceramics were almost never explored concerning molecular compounds and SCO materials have never been elaborated as ceramics to date.

The aim of this communication concerns the preparation of a ceramic of a SCO coordination polymer using Cool-SPS and the investigation of its impact on the functional properties. We chose the well-known $[\text{Fe}(\text{Htrz})_2(\text{trz})](\text{BF}_4)$ complex (where Htrz = 1H-1,2,4-triazole) as the starting material because of its well-controlled synthesis route, its wide hysteresis loop above room temperature and its good stability up to *ca.* 300 °C. Additionally, microstructural effects have been already studied in this compound at the nanoscale,^{42,51–54} after integration in a polymer matrix,^{55,56} under hydrostatic pressure^{57,58} or during grinding experiments suggesting a possible dependence of the SCO feature upon shaping.^{59,60}

We propose here to explore the effect of sintering of $[\text{Fe}(\text{Htrz})_2(\text{trz})](\text{BF}_4)$ on the structural, microstructural and SCO features by comprehensive comparison between powder and molecular ceramics obtained by Cool-SPS.

Experimental section

Synthesis procedure of $[\text{Fe}(\text{Htrz})_2(\text{trz})](\text{BF}_4)$ (1)

All the chemicals and solvents were used as purchased. The iron salt $(\text{Fe}(\text{BF}_4)_2 \cdot 6\text{H}_2\text{O}$; 97%), triazole (1H-1,2,4-triazole, 98%) and ascorbic acid (99%) were used.

The synthesis consists of first dissolving the iron salt $(\text{Fe}(\text{BF}_4)_2 \cdot 6\text{H}_2\text{O}$) (3 mmol) with a few mg of ascorbic acid into 5 mL of demineralized water on one side and the triazole (1H-1,2,4-triazole, Htrz) (9 mmol) into 5 mL of ethanol on the other side. The two solutions were heated at 80 °C. Then, the solution of Htrz was added dropwise to the iron solution with magnetic stirring at 80 °C under a N_2 atmosphere. The two solutions were allowed to react with mixing for 15 min at 80 °C. The resulting purple suspension was then allowed to cool down at 5 °C for 1 hour to promote precipitation. After centrifugation (12 000 rcf, 5 minutes) of the suspension, the precipitate was washed three times with ethanol and once with diethyl ether (50 mL of the solvent was used each time). For these washing steps, the precipitate was separated from the solvent by the same procedure of centrifugation. The precipitate was finally allowed to dry at RT under an ambient atmosphere overnight and a purple powder was obtained.

Sintering using Cool-SPS

An SPS machine (Dr Sinter SPS-515S) operating under vacuum ($P \approx 10$ Pa) was used. To allow for the application of large pressures (up to 600 MPa), a tungsten carbide (WC) mould was used. To protect both the sample and the mould, a carbon foil (Papyex®) of 0.2 mm thickness separated them. Approximately 200 mg of starting powder were used to elaborate a ceramic. Once prepared, the die assembly containing the powder was placed in the SPS chamber, which was evacuated, and the selected sintering sequence was subsequently applied, using a 12:2 pulse sequence. Throughout sintering, several instrumental parameters were monitored including voltage, current, temperature, applied force, SPS chamber pressure and displacement. Cold-pressing experiments were run in the SPS environment without the use of any electrical power, thus using the same die and applying the pressure in an identical way. Samples were recovered by dismounting a three-part mould, to preserve the integrity of obtained ceramics.

X-ray diffraction

Powder X-ray diffraction (PXRD) patterns were collected on a PANalytical X'pert PRO MPD diffractometer in Bragg–Brentano θ -2 θ geometry equipped with a primary monochromator (oriented Ge(111) crystal) to filter the $\text{K}\alpha_2$ radiation and X'Celerator multi-strip detector. Each measurement was made within an angular range of $2\theta = 5\text{--}50^\circ$ and lasted overnight. The $\text{Cu}-\text{K}\alpha_1$ radiation was generated at 45 KV and 40 mA ($\lambda = 1.5059$ Å).

The samples were placed on sample holders made of Si wafer to reduce background and flattened with a razor blade.



DSC measurements

DSC experiments were performed on a PerkinElmer DSC8000 system coupled with CLN2 for thermalisation of the calorimetric bloc. The samples were introduced into a standard aluminium sample pan and put inside the calorimeter. We took particular care considering sample shaping and thermal transfer for both ceramic and powder by using similar mass (*ca.* 15–20 mg) and thickness of samples in order to ensure the best comparability of results.

Cycles were performed at different rates from 0.5 K min⁻¹ to 30 K min⁻¹ under an argon flux of 40 mL min⁻¹.

Magnetic measurements

Magnetic measurements were performed using a vibrating sample magnetometer (VSM MicroSense EZ7) operating at 15 kOe of the applied magnetic field. The weighted powder samples (\approx 5–10 mg) were sealed in a tin capsule then glued onto a quartz fibre and placed in a nitrogen flux furnace between the magnets. The diamagnetic contribution of the quartz and the tin were subtracted from the data using a blank experiment. The procedure consisted of cycling the samples between 300 K and 420 K twenty times for each of them. The thermal cycles were performed at 5 K min⁻¹ scan rate for all the measurements.

Diametral compression test

Diametral compression tests were performed using a TA.HD-plus texture analyser (Stable Microsystems, United Kingdom). Compacts were compressed between two flat surfaces at a constant speed of 0.10 mm s⁻¹ with an acquisition rate of 500 Hz. The breaking force Fr of each tablet was recorded. The maximum tensile stress inside the sample during test σ_T of each tablet was evaluated using the equation:

$$\sigma_T = \frac{2F}{\pi Dh}$$

where h and D are respectively the tablet height and diameter.

Result and discussion

Starting material characterization

The SCO complex $[\text{Fe}(\text{Htrz})_2(\text{trz})]\text{BF}_4$ (**1**) was synthesized following the previously described procedure by addition under continuous stirring at 80 °C of an ethanolic solution of Htrz into an aqueous solution of $\text{Fe}(\text{BF}_4)_2$.⁶¹ After drying, the resulting pink powder corresponds to the desired compound (Elem. anal. calcd for $\text{C}_6\text{H}_8\text{N}_9\text{FeBF}_4$ wt%: C: 20.6; H: 2.3; N: 36.1; B: 3.1; Fe: 16.0, Found: C: 20.1; H: 2.3; N: 34.5; B: 3.4, Fe: 14.6). X-ray powder diffraction pattern of **1** reveals a well crystallized compound corresponding to the previously described polymorph I made of interacting $[\text{Fe}(\text{Htrz})_2(\text{trz})]^+$ chains where Htrz and trz- act as bridging ligands.^{62,63} Rietveld refinement of **1** has been performed using JANA2006⁶⁴ and Fullprof softwares⁶⁵ using the previously described crystal structure as a starting point.⁶² The calculated profile fits very well the experimental

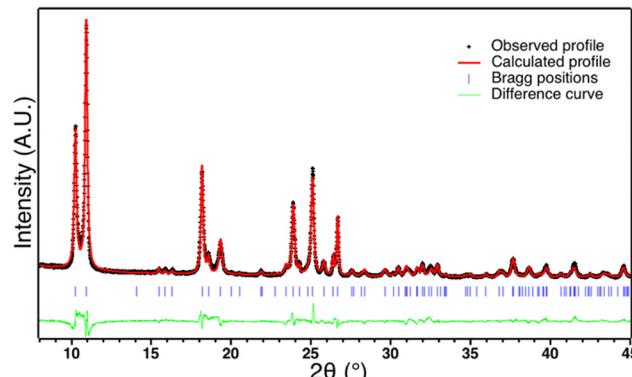


Fig. 1 Rietveld refinement results on powder X-ray diffraction pattern of **1** at 293 K. Experimental pattern (black cross), calculated pattern using polymorph I data (red line), Difference curve (green line), Bragg positions (vertical lines).

pattern ($R_{wp} = 0.10$, $R_p = 0.07$, $R_{\text{Bragg}} = 0.07$), which confirms the crystal structure of **1** (Fig. 1). After subtraction of instrumental broadening, sample contribution to peak broadening could be estimated to a mean coherent domain size of *ca.* 50 nm and micro-deformations of *ca.* 20.10^{-4} . Furthermore, the coherent domain sizes are anisotropic with a strong elongation along the *b* parameter corresponding to the chain direction (260 nm in the *b* direction and 34 nm in the *ac* direction). These observations are in line with the particle size observed by TEM (Fig. S1, ESI[†]) and with the previously reported results on the same compound⁶⁶ except that coherent domain size are smaller in our case, probably due to the synthetic procedure.

Magnetic behaviour of the starting powder

Magnetic properties of **1** were recorded on a Vibrating Sample Magnetometer (VSM). At 300 K the $\chi_M T$ product (where χ_M is the molar magnetic susceptibility and T the temperature) is around $0.3 \text{ cm}^3 \text{ K}^{-1} \text{ mol}^{-1}$ corresponding to a Fe^{2+} in the Low Spin (LS) state with a slight High Spin (HS) residual fraction. On warming, $\chi_M T$ product remains close to 0 up to 380 K and increases abruptly above 380 K to reach a value of $3.25 \text{ cm}^3 \text{ K}^{-1} \text{ mol}^{-1}$ at 400 K as expected for a Fe^{2+} ion in the HS state. This value remains stable till 420 K. On cooling, the $\chi_M T$ product remains around the full HS value before rapidly falling down to *ca.* $0.3 \text{ cm}^3 \text{ K}^{-1} \text{ mol}^{-1}$ at 340 K, a signature of an almost complete HS to LS conversion (Fig. 2). It remains close to $0.3 \text{ cm}^3 \text{ K}^{-1} \text{ mol}^{-1}$ on cooling to 250 K. These features are characteristic of a cooperative SCO with a 40 K wide hysteresis ($T_{1/2\uparrow} = 380 \text{ K}$ and $T_{1/2\downarrow} = 340 \text{ K}$) as expected for the polymorph I of this compound. Several cycles (one cycle corresponding to crossovers from LS to HS in warming mode followed by HS to LS in cooling mode) were then performed between 300 and 420 K to evaluate the evolution of the hysteresis loop upon cycling. The first ten cycles showed a strong modification of the transition temperatures that decreased by almost 10 K in the warming mode ($T_{1/2\uparrow}$) leading to a smaller hysteresis loop. The next ten cycles showed a smaller evolution of both $T_{1/2\downarrow}$ and $T_{1/2\uparrow}$. It is worth noting that at least ten of cycles are thus required to reach a “stabilized” hysteresis loop for which $T_{1/2\downarrow}$ and $T_{1/2\uparrow}$



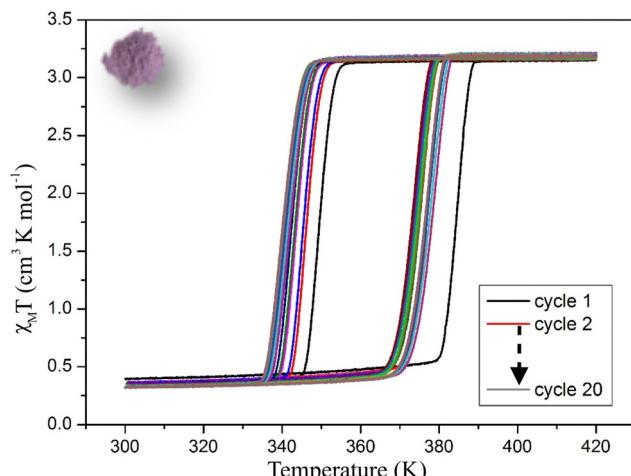


Fig. 2 $\chi_M T$ product versus T recorded on a powder sample of **1** showing the evolution of the magnetic properties upon 20 cycles. Scan rate = 5 K min⁻¹.

converge to 341 K and 374 K respectively leading to a 33 K stable hysteresis loop for this compound (Fig. S2a, ESI[†]). Therefore, upon thermal cycling, the width of the hysteresis decreases as a reflection of a weakening of the cooperative character of the compound. This might correspond to a fatigability of the sample, as already reported.⁶⁶ It is also important to mention that cooperativity is related to either the abruptness of the switching (the temperature range of the LS \leftrightarrow HS switching) or the width of the thermal hysteresis. Several thermodynamic models were developed to account for both characteristics and will be discussed in the following sections.

DSC measurements on the starting powder

The deviation from normal heat capacity (ΔC_p) has been recorded upon twenty cycles and is reported for the starting

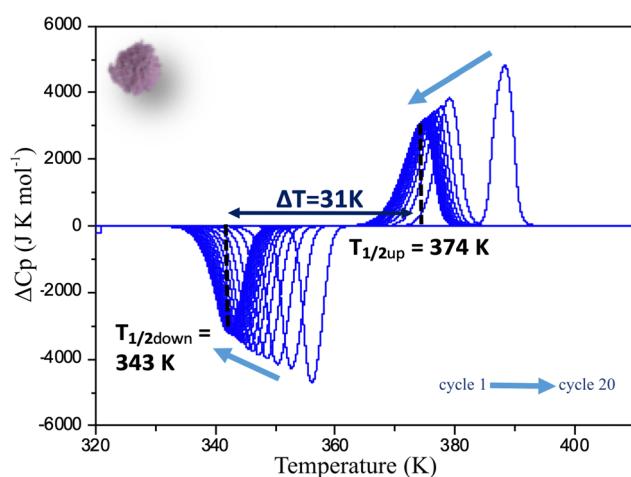


Fig. 3 Deviation from normal molar heat capacity (ΔC_p) vs. temperature for **1** recorded by DSC on a powder sample at 5 K min⁻¹. Arrows are showing the evolution from the first to the 20th cycle. ΔC_p peaks in the cooling mode are artificially represented negatives, that way, the endothermic process corresponds to a positive peak in the chart.

powder sample on Fig. 3. Two peaks corresponding to the LS \rightarrow HS, and HS \rightarrow LS SCO in the warming and cooling modes, respectively, are clearly detected. As already observed with magnetic measurements, a strong evolution of the two $T_{1/2}$ is observed upon cycling, which is more pronounced concerning $T_{1/2\downarrow}$. It is worth noting that at least 10 cycles are needed before reaching a stabilized hysteresis loop with $T_{1/2\uparrow}$ and $T_{1/2\downarrow} = 374$ K and 343 K respectively, perfectly matching the magnetic behaviour (Fig. 2 and Fig. S2a, ESI[†]). Additionally, a significant decrease of the absolute value peak maxima or minima accompanied by an enlargement of the peak is clearly visible in Fig. 3. The LS \leftrightarrow HS switching thus takes place on a wider temperature range upon cycles on either ascending or descending branches indicating a decrease of the abruptness of this switching. This effect may be linked to the structural fatigability that is evidenced by the decrease of the coherent domain sizes and the slight modification of the cell parameters after cycling (Fig. S3, ESI[†]). Two effects, thus, go with the thermal cycling: a lowering of the transition temperatures and a decrease of the cooperativity of the LS \rightarrow HS, and HS \rightarrow LS SCO as shown by the increase of DSC peak widths.

Elaboration of ceramics

Even if this compound is not purely molecular due to its 1-D polymeric nature, ceramics of **1** fall into the definition of “molecular ceramics” given in ref. 10. Such molecular ceramics of **1**, were then prepared using Cool-SPS. Different sintering parameters and sequences were explored and optimized (pressure, temperature, ramps, *etc.*), to design an efficient Cool-SPS process. Sintering sequences for which the temperature was applied before pressure were found adequate and extensively explored with respect to Cool-SPS accessible conditions (50 \leq T \leq 600 °C, 40 \leq P \leq 600 MPa). The sintering map obtained (see Fig. S4, ESI[†]) revealed that combining high temperatures and pressures ($T \geq 300$ °C, $P > 400$ MPa) is detrimental to the stability of **1**, while operating at either low temperature ($T \leq 200$ °C) or pressure ($P < 200$ MPa) would not yield dense ceramics. The highest densifications (95 % of relative density, where the relative density is the ratio between the geometrically measured density of the pellets and the density of the corresponding single crystal) were obtained for temperatures ranging between 250 and 300 °C, and pressures between 300 and 400 MPa.

Mechanical testing of molecular ceramics

Application of a significant pressure during a low temperature sintering process always questions the nature of the obtained samples, which may be dense but poorly cohesive compact lacking intergranular bonding, or ceramics showing improved strength thanks to optimized intergranular bonding through sintering. To address this point, a series of twelve ceramics (3 ceramics per processing condition from which nine could be used in the diametral compression test) was specifically elaborated and brought under a breaking test to evaluate their cohesiveness. Except for samples of the lowest density, ceramics obtained using Cool-SPS typically have a larger tensile

strength, demonstrating their higher cohesiveness, as compared to cold pressed powders (see Fig. S5a, ESI[†]) of similar relative density. This trend is confirmed when differentiating the elastic part of the mechanical response (see Fig. S5b, ESI[†]), where Cool-SPS ceramics display slopes ranging from 26 to 42 MPa mm⁻¹, *versus* 24 to 27 MPa mm⁻¹ for cold-pressed samples, *i.e.* an increase of almost 50 % in the best case. Since all samples are constituted of the same compound, displaying the same crystal structure, these results confirm a significant level of intergranular bonding in Cool-SPS processed molecular ceramics, responsible for their increased strength. Thus, providing they kept their SCO behaviour, the most cohesive and dense samples potentially form the first series of functional molecular ceramics. To establish this point, a ceramic **1** of 95 % relative density denoted as **C1** underwent a complete set of characterization studies identical to those performed on powder in order (i) to check if the functional feature (SCO) was preserved and (ii) to compare the behaviour of the ceramics *versus* the powders.

X-ray diffraction of the ceramics

X-ray diffraction patterns recorded for **C1** is shown in Fig. 4. It is clearly very similar to the pattern of the corresponding powder. The Rietveld refinement confirms that the polymorph I is preserved after sintering ($R_{wp} = 0.09$, $R_p = 0.07$, $R_{bragg} = 0.05$). Nevertheless some slight structural and/or microstructural details differ. The calculated profile reveals a different shape of coherent domain size and a slightly smaller a parameter. Even if the mean coherent domain size is of the same order (but slightly smaller: 35 nm), the anisotropic shapes differ: while the apparent sizes in a and c directions remain similar to the powder (30 nm *vs.* 34 nm for the powder), the size in the b direction is significantly smaller (140 nm *vs.* 260 for the powder). Such an effect of the reduction of the coherent domain size has been previously observed on powder after more than 10 cycles and attributed to fatigability,⁶⁶ which

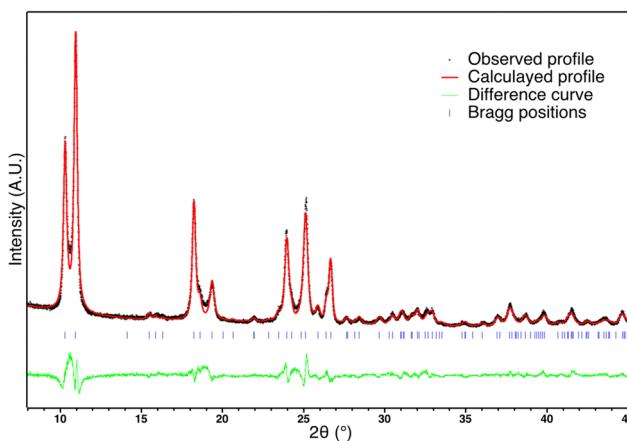


Fig. 4 Rietveld refinement results on powder X-ray diffraction pattern of molecular ceramic of **1** (**C1**) at 293 K. Experimental pattern (black cross), calculated pattern (red line), Difference curve (green line), Bragg positions (vertical lines).

was also confirmed for powder in this study (Fig. S3, ESI[†]). It was suggested that this microstructural fatigability was accompanied by a subtle structural variation corresponding to a slight rotation of the chain linked to the decrease of the a parameter.⁶⁷ Since a similar evolution of the a parameter is also detected on ceramics in addition to the reduction of the coherent domain size in the b direction, it seems that the sintering process leads to a very similar microstructural and structural evolution, but without cycling.

Magnetism and DSC of ceramics

Magnetic properties were recorded on a vibrating sample magnetometer (VSM). Several cycles of warming and cooling were performed on **C1** (Fig. 5). For the first cycle, upon warming from 300 K the $\chi_M T$ product remains at around $0.3 \text{ cm}^3 \text{ K}^{-1} \text{ mol}^{-1}$ until *ca.* 370 K and increases sharply to reach *ca.* $3.25 \text{ cm}^3 \text{ K}^{-1} \text{ mol}^{-1}$ at 380 K and higher, which is a signature of an almost complete LS to HS transition. On cooling the $\chi_M T$ product remains stable until *ca.* 345 K and then rapidly decreases to reach $0.3 \text{ cm}^3 \text{ K}^{-1} \text{ mol}^{-1}$, corresponding to the HS to LS SCO with a small amount of remaining HS complexes below the transition temperature. This feature proves that after sintering the molecular ceramics have conserved their SCO functionality, corresponding to a cooperative spin transition with a 39 K hysteresis width ($T_{1/2\uparrow} = 379 \text{ K}$ and $T_{1/2\downarrow} = 340 \text{ K}$). Upon cycling between 300 and 400 K an evolution of the transition temperatures is clearly observed. Interestingly, this phenomenon is much less pronounced than for the powder especially in the cooling mode. $T_{1/2\uparrow}$ decreases to 374 K whereas $T_{1/2\downarrow}$ remains almost unchanged at around 340 K, leading to a hysteresis loop of 34 K wide equivalent to the powder behaviour after 20 cycles (Fig. S2b, ESI[†]). The stabilized hysteresis loops of powder and ceramic are thus very similar even if the ceramic seems to present a slightly more abrupt transition ($\Delta T_{60} = 6.1 \text{ K}$ for ceramic, *vs.* 7.4 K for powder in warming mode, where ΔT_{60} is the temperature range in which 60% of the iron ions (from 80% LS to 80% HS) undergo the SCO).⁶⁸

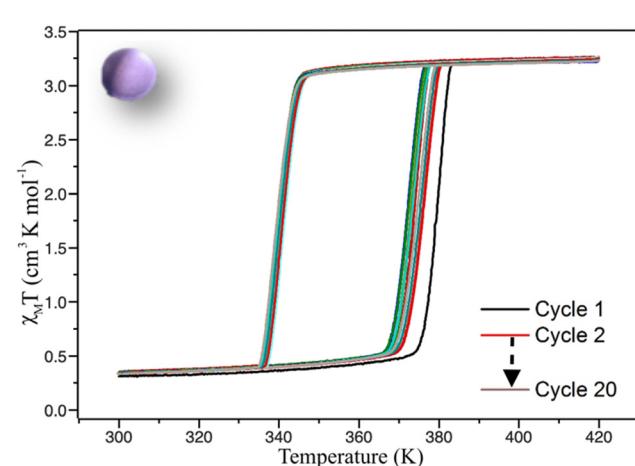


Fig. 5 $\chi_M T$ product *versus* T recorded for molecular ceramic of **1** (**C1**) showing the evolution of the magnetic properties in 20 cycles. Scan rate = 5 K min^{-1} .



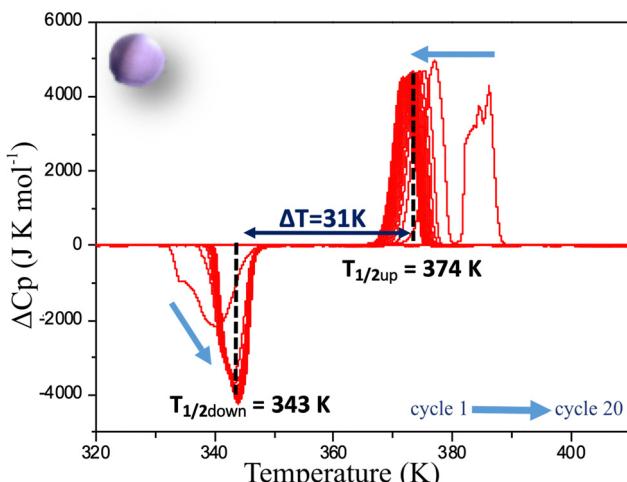


Fig. 6 Deviation from normal molar heat capacity (ΔC_p) vs. temperature for a molecular ceramic of **1** recorded by DSC at 5 K min^{-1} . Arrows show the evolution from the first to the 20th cycle.

As for the powder, a full DSC study has been performed on **C1** following the same procedure. The deviation from the normal heat capacity (ΔC_p) has been recorded upon twenty cycles and the two peaks corresponding to the LS \rightarrow HS, and HS \rightarrow LS transitions in the warming and cooling mode respectively are shown after different cycles in Fig. 6. The evolution of the positions of the peaks differs significantly from that observed for the powder. Although in this case, conversely to magnetic behaviour, a very singular first cycle is observed, the subsequent cycles behave differently. Indeed, $T_{1/2\uparrow}$, as for powder, slightly decreases to reach 374 K, but $T_{1/2\downarrow}$ increases instead, to stabilize much more rapidly than for the powder after a few cycles at 343 K. This behaviour is in line with the magnetic measurements confirming that the ceramic is already close to the stabilized state after sintering. This can be related to X-ray diffraction results that showed that the structural and microstructural properties after sintering are closer to those observed after several cycles on the powder. Sintering process would thus induce structural and microstructural modifications in the polycrystalline sample. As a consequence, only 5 cycles are needed to get a stabilized hysteresis loop from the molecular ceramics *versus* 12 for the powder (Fig. S6, ESI[†]). It is worth noting that, conversely to the powder, the absolute value of the DSC peak maxima or minima do not change, leading to the same peak widths and heights upon cycling. Thus, cycling the ceramics do not modify the abruptness/cooperativity of the transition demonstrating limited fatigability. Additionally, comparing the absolute values of DSC peak maxima or minima for the ceramic and the powder shows that they are systematically higher in the molecular ceramics with respect to the powder, and correspond to narrower peaks and then more steep transitions.

Discussion

Powder of the SCO material $[\text{Fe}(\text{Htrz})_2(\text{trz})](\text{BF}_4)$ has been successfully used to produce molecular ceramics with a relative

density as high as 95%. Either from magnetic or DSC measurements, the conservation of SCO properties in the molecular ceramics is undisputable, with the same global behaviour as the starting powder. X-ray diffraction study showed that Cool-SPS processing did not induce any irreversible phase transition as the structural analysis confirmed that the ceramics correspond to the same polymorph than the starting powder (recognized as polymorph I⁶³). Nevertheless, Rietveld refinement revealed that the sintering process induced structural and microstructural changes similar to those observed upon cycling with the powder.^{66,69} It is worth mentioning that one cycle of LS–HS–LS transition is occurring during the SPS process due to pressure/temperature modification, which can be partially detected upon SPS monitoring (Fig. S7, ESI[†]). This structural and microstructural relaxations lead to quicker stabilization of the hysteresis loop upon cycling for the ceramic as compared to the powder, which is well illustrated by the magnetic behaviour of both samples and even more clearly revealed by DSC. The main difference remaining after cycling between powder and ceramic concerns the sharpness of the transition which is significantly more pronounced for ceramics than for powder according to DSC results. To go further on that point, DSC measurements have been used to obtain the thermodynamic features of both samples and implemented in well-known thermodynamic models such as the Slusher & Drickamer (S&D)⁷⁰ or Sorai & Seki (S&S)⁷¹ ones. In order to emphasize only the differences between stabilized SCO behaviours, it is mandatory to work on cycles greater than the 10th one, we chose then to work on the 12th cycle. Integration of the ΔC_p peaks and $\Delta C_p/T$ peaks over T gave respectively the variation of enthalpy and entropy of the samples at the SCO. The results are tabulated in Table 1. The obtained values for the variation of enthalpy at the transition in the warming and cooling modes $\Delta H\uparrow$ and $\Delta H\downarrow$ are very similar, demonstrating that the transition is fully reversible for both samples (powder and ceramic). Then, the mean value has been chosen for ΔH in both cases. Conversely, as this compound is associated with an hysteresis behaviour, two distinct values of ΔS corresponding to $T_{1/2\uparrow}$ and $T_{1/2\downarrow}$ are obtained, the real one (used in the S&D model) should lie in between these two values but is not necessarily the average. The thermodynamic values obtained that way are very similar for powder and molecular ceramics and close to those found in the literature for this compound.⁷² It confirms that the SCO phase transition involves the same polymorph (polymorph I) in both cases. Nevertheless, variations of enthalpy of ceramic are slightly lower than powder probably due to a somewhat less complete transition for the ceramic as suggested

Table 1 Thermodynamic parameters extracted from DSC measurements and S&D model

Sample	$\langle \Delta H \rangle$ kJ mol ⁻¹	$\Delta S\uparrow$ J K ⁻¹ mol ⁻¹	$\Delta S\downarrow$ J K ⁻¹ mol ⁻¹	ΔS (S&D) J mol ⁻¹	Γ J mol ⁻¹
Powder	26.3	68.6	77.7	73.2	8500
Ceramic	24.9	65.2	73.3	69.3	8500



by the slightly lower $\chi_M T$ product obtained for the ceramic above 380 K.

These thermodynamic parameters have been then introduced into the S&D model to evaluate the so-called cooperativity parameter Γ . The fit of the magnetic curves that is illustrated in Fig. S8 (ESI†) leads to the same value of Γ for both samples (Table 1). It is not surprising since (i) the stabilized SCO cycles correspond to the same hysteresis width and (ii) Γ is representative of intermolecular interactions that should be the same between powder and ceramics as both correspond to the same polymorph. Because the S&D model cannot explain the difference in abruptness between both samples we also used the S&S model which introduces the notion of cooperative domains. This model uses the number n that represents the number of metal centres undergoing the SCO simultaneously.

Thus, the higher the n number and the more cooperative the transition, the more abrupt the SCO. Since this model cannot account for hysteresis behaviours, we used only the LS to HS transition and fit the corresponding DSC peak with the S&S model to extract n . Fig. S9 (ESI†) shows the result of the fit for both samples on the 12th cycle leading to $n_p = 17$ for the powder and $n_c = 35$ for the ceramic. This agrees with a steeper and more cooperative transition observed in ceramics.

In order to check how this parameter depends on thermal conduction capability of the samples, we also explore the effect of the temperature scan rate. Fig. 7 shows the ΔC_p over T curves for different scan rates for powder and ceramic between 0.5 K min⁻¹ to 30 K min⁻¹.

Three observations can be made: (i) at 1 K min⁻¹, the scan rate is high enough to artificially increase the width of the transition; (ii) the $\Delta C_{p\max}$ values are systematically higher for the ceramic than for the powder even at very low scan rates and (iii) ΔC_p peak widths and maxima did not converge between ceramics and powders even at very low scan rates. As $\Delta C_{p\max}$ directly influences the value of n in the S&S model, the value of n tends to artificially decrease when the scan rate increases. Nevertheless, as $\Delta C_{p\max}$ are systematically higher for the ceramic than for the powder, the value of the n S&S parameter may be effectively higher for the ceramics. In order to overcome this scan rate effect we calculated the n parameter by fitting the DSC

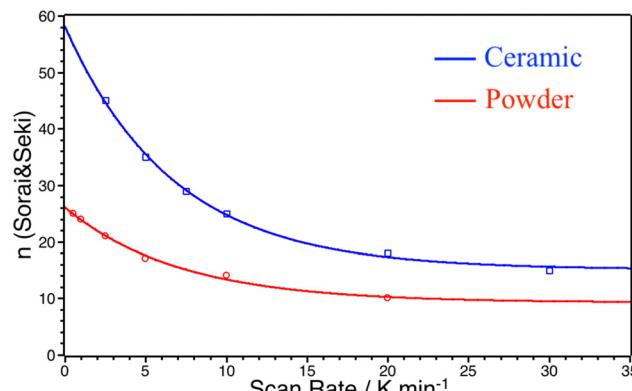


Fig. 8 Evolution of the cooperative domain size parameter n obtained from the S&S model for different temperature scan rates for powder (red circles) and ceramic (blue squares).

peaks of the powder (n_p) and ceramic (n_c) for all scan rates using the S&S model. The results are presented in Fig. 8 and show that n_c is indeed systematically higher than n_p at any scan rate and the effect is hindered at high rates. By extrapolating this evolution of n to scan rate = 0 K min⁻¹, one can evaluate the limit of the n value for the powder and the ceramic, ridding off the scan rate effect. The limit cooperative domain sizes are found, that way, to be $n_p = 26$ and $n_c = 84$ for the powder and the ceramic, respectively. The sintering of the powder into a molecular ceramic tends thus to increase the cooperative domain size, which becomes triple that of the corresponding powder. When the scan rate is increasing, the cooperative domain size rapidly decreases for ceramics, while it is much less affected for powder, which can be due to better thermal conductivity of ceramics.

Thereby, while the grinding process was proven to affect the SCO by decreasing the cooperativity directly linked to smaller crystal coherent domain sizes, sintering by Cool-SPS, even if the influence is smaller, leads to the opposite effect probably due to the improvement of thermal conductivity of the sintered samples. It could be then an interesting tool for understanding the possible different kinds of cooperativity in SCO compounds for example: microscopic intrinsic cooperativity and macroscopic microstructurally dependent cooperativity.

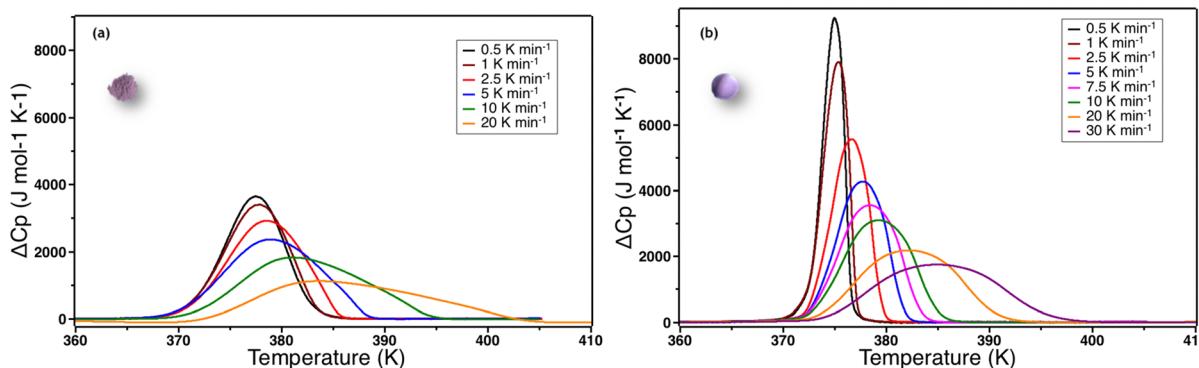


Fig. 7 Effect of the scan rate on the ΔC_p peaks for (a) powder and (b) molecular ceramics of 1.



Conclusions

Shaping of SCO materials has been widely studied during the last decade, from thick to thin films, bulk to nanomaterials, composites or alloys but the molecular ceramic form has not been investigated so far. Thanks to Cool-SPS, we have demonstrated that high density and cohesive functional molecular ceramics are accessible. Starting from the model $[\text{Fe}(\text{trz})_2(\text{Htrz})]\text{BF}_4$ SCO cooperative compound, we have obtained a new molecular ceramic of 95% relative density with the complete conservation of the SCO functionality. The global behaviour (transition temperature, hysteresis width) was preserved. This new way of shaping the SCO compounds is thus very promising for the insertion of SCO materials into specific devices using traditional ceramics, but also enables a more easy and reliable access to some specific measurements like dielectric, polarization or conductivity properties as well as to the conception of composites.

Additionally, from an academic point of view, deep investigation using simple thermodynamic models showed that SCO molecular ceramics present larger “cooperative domains” than the corresponding powders. Sintering would then improve the interdomain pathways to increase cooperativity at the macroscopic level probably due to much better thermal conduction capabilities. These specificities make opportunities to better understand both the notion of cooperativity in SCO materials and the notion of molecular ceramic itself, in which, beyond the contribution of intermolecular forces to cohesiveness at the coherent domain scale, functional properties like SCO are expressed beyond the boundaries of single particles or coherent domains.

Author contributions

The manuscript was written through contributions from all authors. All authors have given approval to the final version of the manuscript. MM wrote the manuscript and supervised all the work, LEK performed the synthesis of all samples and supervised the characterization, ND synthesized the starting powder, MJ carried out and supervised the sintering of the samples, DD and LE performed the DSC analysis, PR supervised the magnetic measurements and chemical analysis CG tested the thermodynamic models and carefully read the whole manuscript, and VM performed the diametral compression of pellets.

Conflicts of interest

The authors declare no conflict of interest.

Acknowledgements

This work was supported by the ANR-19-CE07-0026 (MEMORI). The authors acknowledge Centre National de la Recherche Scientifique (CNRS), the University of Bordeaux and Bordeaux INP for technical support and the Lebanon State for part of PhD. grant of Liza El-Khoury. Michaël Josse thanks CNRS and the University of Bordeaux for 6 months of sabbatical leave

during 2017-2018 and 2019–2020, for the development of molecular ceramics and Cool-SPS. The X-ray diffraction facility of the ICMCB is thanked for the powder diffraction experiments and more particularly Eric Lebraud for his useful technical support. The authors also thank U-Chan Chung for his support for the SPS. The UAR 3626 PLACAMAT is acknowledged for the TEM facilities and more specifically Marion Gayot for her help.

Notes and references

- 1 G. Barreto, P. Canhoto and M. Collares-Pereira, Parametric analysis and optimisation of porous volumetric solar receivers made of open-cell SiC ceramic foam, *Energy*, 2020, **200**, 117476.
- 2 X. Kong, L. Yang, Z. Cheng and S. Zhang, Bi-modified SrTiO_3 -based ceramics for high-temperature energy storage applications, *J. Am. Ceram. Soc.*, 2020, **103**, 1722–1731.
- 3 E. A. Mikhaleva, I. N. Flerov, E. V. Bogdanov, V. S. Bondarev, M. V. Gorev and E. Rysiakiewicz-Pasek, Size effect on sensitivity to external pressure and caloric effects in TGS: Ceramics and nanocomposites, *Mater. Today Commun.*, 2020, **25**, 101463.
- 4 K. Zou, Y. Dan, H. Xu, Q. Zhang, Y. Lu, H. Huang and Y. He, Recent advances in lead-free dielectric materials for energy storage, *Mater. Res. Bull.*, 2019, **113**, 190–201.
- 5 J. Guo, R. Floyd, S. Lowum, J.-P. Maria, T. Herisson de Beauvoir, J.-H. Seo and C. A. Randall, Cold Sintering: Progress, Challenges, and Future Opportunities, *Annu. Rev. Mater. Res.*, 2019, **49**, 275–295.
- 6 J.-P. Maria, X. Kang, R. D. Floyd, E. C. Dickey, H. Guo, J. Guo, A. Baker, S. Funihashi and C. A. Randall, Cold sintering: Current status and prospects, *J. Mater. Res.*, 2017, **32**, 3205–3218.
- 7 M. T. Sebastian, H. Wang and H. Jantunen, Low temperature co-fired ceramics with ultra-low sintering temperature: A review, *Curr. Opin. Solid State Mater. Sci.*, 2016, **20**, 151–170.
- 8 S. Li, X. Wang, J. Ren, C. Liu, Y. Hu and Y. Yang, Microstructure, mechanical property and corrosion behavior of biomedical Zn-Fe alloy prepared by low-temperature sintering, *J. Alloys Compd.*, 2023, **934**, 167812.
- 9 T. Herisson de Beauvoir, A. Sanggregorio, A. Bertrand, C. Payen, D. Michau, U.-C. Chung and M. Josse, Cool-SPS stabilization and sintering of thermally fragile, potentially magnetoelectric, $\text{NH}_4\text{FeP}_2\text{O}_7$, *Ceram. Int.*, 2019, **45**, 9674–9678.
- 10 T. Herisson de Beauvoir, V. Villemot and M. Josse, Cool-Spark plasma sintering: An opportunity for the development of molecular ceramics, *Solid State Sci.*, 2020, **102**, 106171.
- 11 T. Herisson de Beauvoir, A. Sanggregorio, I. Cornu, C. Elissalde and M. Josse, Cool-SPS: an opportunity for low temperature sintering of thermodynamically fragile materials, *J. Mater. Chem. C*, 2018, **6**, 2229–2233.
- 12 E. Coronado, Molecular magnetism: from chemical design to spin control in molecules, materials and devices, *Nat. Rev. Mater.*, 2020, **5**, 87–104.



13 B. L. Feringa and W. R. Browne, *Molecular Switches*, Wiley-VCH Verlag GmbH & Co. KGaA, Weinheim, Germany, 2011, vol. 1.

14 M. D. Manrique-Juárez, S. Rat, L. Salmon, G. Molnár, C. M. Quintero, L. Nicu, H. J. Shepherd and A. Bousseksou, Switchable molecule-based materials for micro- and nanoscale actuating applications: Achievements and prospects, *Coord. Chem. Rev.*, 2015, **308**, 395–408.

15 Y.-S. Meng, O. Sato and T. Liu, Manipulating Metal-to-Metal Charge Transfer for Materials with Switchable Functionality, *Angew. Chem., Int. Ed.*, 2018, **57**, 12216–12226.

16 *Molecular Machines and Motors*, ed. J.-P. Sauvage, V. Amendola, R. Ballardini, V. Balzani, A. Credi, L. Fabbrizzi, M. T. Gandolfi, J. K. Gimzewski, M. Gómez-Kaifer, C. Joachim, A. E. Kaifer, E. Katz, T. R. Kelly, J. Liu, C. Mangano, P. Pallavicini, A. R. Pease, L. Raehm, M. Sano, J.-P. Sauvage, J. P. Sestelo, A. N. Shipway, J.-F. Stoddart, M. Venturi and I. Willner, Springer Berlin Heidelberg, Berlin, Heidelberg, 2001, vol. 99.

17 M. A. Halcrow, *Spin-Crossover Materials*, John Wiley & Sons Ltd, Oxford, UK, 2013.

18 H. Li and H. Peng, Recent advances in self-assembly of spin crossover materials and their applications, *Curr. Opin. Colloid Interface Sci.*, 2018, **35**, 9–16.

19 G. Molnár, S. Rat, L. Salmon, W. Nicolazzi and A. Bousseksou, Spin Crossover Nanomaterials: From Fundamental Concepts to Devices, *Adv. Mater.*, 2018, **30**, 1703862.

20 K. Ridier, A.-C. Bas, Y. Zhang, L. Routaboul, L. Salmon, G. Molnár, C. Bergaud and A. Bousseksou, Unprecedented switching endurance affords for high-resolution surface temperature mapping using a spin-crossover film, *Nat. Commun.*, 2020, **11**, 3611.

21 K. S. Kumar and M. Ruben, Sublimable Spin-Crossover Complexes: From Spin-State Switching to Molecular Devices, *Angew. Chem., Int. Ed.*, 2021, **60**, 7502–7521.

22 K. Senthil Kumar and M. Ruben, Emerging trends in spin crossover (SCO) based functional materials and devices, *Coord. Chem. Rev.*, 2017, **346**, 176–205.

23 X. Zhang, T. Palamarcu, J.-F. Létard, P. Rosa, E. V. Lozada, F. Torres, L. G. Rosa, B. Doudin and P. A. Dowben, The spin state of a molecular adsorbate driven by the ferroelectric substrate polarization, *Chem. Commun.*, 2014, **50**, 2255.

24 A. Bousseksou, N. Negre, M. Goiran, L. Salmon, J. P. Tuchagues, M. L. Boillot, K. Boukheddaden and F. Varret, Dynamic triggering of a spin-transition by a pulsed magnetic field, *Eur. Phys. J. B*, 2000, **13**, 451–456.

25 G. Chastanet, M. Lorenc, R. Bertoni and C. Desplanches, Light-induced spin crossover—Solution and solid-state processes, *Comptes Rendus Chim.*, 2018, **21**, 1075–1094.

26 S. Decurtins, P. Gütlich, C. P. Köhler, H. Spiering and A. Hauser, Light-induced excited spin state trapping in a transition-metal complex: The hexa-1-propyltetrazole-iron (II) tetrafluoroborate spin-crossover system, *Chem. Phys. Lett.*, 1984, **105**, 1–4.

27 C. Lefter, R. Tan, J. Dugay, S. Tricard, G. Molnár, L. Salmon, J. Carrey, W. Nicolazzi, A. Rotaru and A. Bousseksou, Unidirectional electric field-induced spin-state switching in spin crossover based microelectronic devices, *Chem. Phys. Lett.*, 2016, **644**, 138–141.

28 Y. Qi, E. W. Müller, H. Spiering and P. Gütlich, The effect of a magnetic field on the high-spin \rightleftharpoons low-spin transition in $[\text{Fe}(\text{phen})_2(\text{NCS})_2]$, *Chem. Phys. Lett.*, 1983, **101**, 503–505.

29 M. Enamullah and W. Linert, Spin-crossover of $[\text{Fe}(\text{Cl}-\text{bzimpy})_2](\text{ClO}_4)_2$ induced by deprotonation, *J. Coord. Chem.*, 1995, **35**, 325–335.

30 F. Renz, P. A. De Souza, G. Klingelhöfer and H. A. Goodwin, *Molecular Sensors for Moisture Detection by Mössbauer Spectroscopy*, in *Industrial Applications of the Mössbauer Effect*, Springer Netherlands, Dordrecht, 2002, pp. 699–704.

31 E. Resines-Urien, E. Fernandez-Bartolome, A. Martinez-Martinez, A. Gamonal, L. Piñeiro-López and J. S. Costa, Vapochromic effect in switchable molecular-based spin crossover compounds, *Chem. Soc. Rev.*, 2022, **52**, 705–727.

32 P. Guionneau, Crystallography and spin-crossover. A view of breathing materials, *Dalt. Trans.*, 2014, **43**, 382–393.

33 W. Nicolazzi, Thermodynamical aspects of the spin crossover phenomenon, *C. R. Chim.*, 2018, **21**, 1060–1074.

34 T. Guillon, S. Bonhommeau, J. S. Costa, A. Zwick, J.-F. Létard, P. Demont, G. Molnár and A. Bousseksou, On the dielectric properties of the spin crossover complex $[\text{Fe}(\text{bpp})_2][\text{BF}_4]_2$, *Phys. Status Solidi*, 2006, **203**, 2974–2980.

35 C. Lefter, V. Davesne, L. Salmon, G. Molnár, P. Demont, A. Rotaru and A. Bousseksou, Charge Transport and Electrical Properties of Spin Crossover Materials: Towards Nanoelectronic and Spintronic Devices, *Magnetochemistry*, 2016, **2**, 18.

36 C. Lefter, I. A. Gural'skiy, H. Peng, G. Molnár, L. Salmon, A. Rotaru, A. Bousseksou and P. Demont, Dielectric and charge transport properties of the spin crossover complex $[\text{Fe}(\text{Htrz})_2(\text{trz})](\text{BF}_4)$, *Phys. Status Solidi RRL*, 2014, **8**, 191–193.

37 P. Guionneau, M. Marchivie, G. Bravic, J.-F. Létard and D. Chasseau, Structural Aspects of Spin Crossover. Example of the $[\text{Fe}^{\text{II}}\text{L}_n(\text{NCS})_2]$ Complexes, *Top. Curr. Chem.*, 2004, **234**, 97–128.

38 M. A. Halcrow, Structure:Function Relationships in Molecular Spin-Crossover Materials, in *Spin-Crossover Materials*, John Wiley & Sons Ltd, Oxford, UK, 2013, pp. 147–169.

39 G. Dupouy, M. Marchivie, S. Triki, J. Sala-Pala, J.-Y. Salaün, C. J. Gómez-García and P. Guionneau, The Key Role of the Intermolecular π – π Interactions in the Presence of Spin Crossover in Neutral $[\text{Fe}(\text{abpt})_2\text{A}_2]$ Complexes (A = Terminal Monoanion N Ligand), *Inorg. Chem.*, 2008, **47**, 8921–8931.

40 E. Tailleur, M. Marchivie, N. Daro, G. Chastanet and P. Guionneau, Thermal spin-crossover with a large hysteresis spanning room temperature in a mononuclear complex, *Chem. Commun.*, 2017, **53**, 4763–4766.

41 M. G. Reeves, E. Tailleur, P. A. Wood, M. Marchivie, G. Chastanet, P. Guionneau and S. Parsons, Mapping the cooperativity pathways in spin crossover complexes, *Chem. Sci.*, 2021, **12**, 1007–1015.

42 T. Delgado and M. Villard, Spin Crossover Nanoparticles, *J. Chem. Educ.*, 2022, **99**, 1026–1035.



43 K. S. Kumar and M. Ruben, Sublimable Spin-Crossover Complexes: From Spin-State Switching to Molecular Devices, *Angew. Chem., Int. Ed.*, 2021, **60**, 7502–7521.

44 H. J. Shepherd, G. Molnár, W. Nicolazzi, L. Salmon and A. Bousseksou, Spin Crossover at the Nanometre Scale, *Eur. J. Inorg. Chem.*, 2013, 653–661.

45 L. Kipgen, M. Bernien, F. Tuczek and W. Kuch, Spin-Crossover Molecules on Surfaces: From Isolated Molecules to Ultrathin Films, *Adv. Mater.*, 2021, **33**, 2008141.

46 S. Yazdani, J. Phillips, T. K. Ekanayaka, R. Cheng and P. A. Dowben, The Influence of the Substrate on the Functionality of Spin Crossover Molecular Materials, *Molecules*, 2023, **28**, 3735.

47 A. Enriquez-Cabrera, A. Rapakousiou, M. Piedrahita Bello, G. Molnár, L. Salmon and A. Bousseksou, Spin crossover polymer composites, polymers and related soft materials, *Coord. Chem. Rev.*, 2020, **419**, 213396.

48 P. Guionneau, M. Marchivie and G. Chastanet, Multiscale Approach of Spin Crossover Materials: A Concept Mixing Russian Dolls and Domino Effects, *Chem. – A Eur. J.*, 2021, **27**, 1483–1486.

49 M. Romanini, Y. Wang, K. Gürpinar, G. Ornelas, P. Lloveras, Y. Zhang, W. Zheng, M. Barrio, A. Aznar, A. Gràcia-Condal, B. Emre, O. Atakol, C. Popescu, H. Zhang, Y. Long, L. Balicas, J. Lluís Tamarit, A. Planes, M. Shatruk, L. Mañosa, M. Romanini, A. Gràcia-Condal, A. Planes, L. Mañosa, Y. Wang, G. Ornelas, Y. Zhang, M. Shatruk, H. Zhang, Y. Long, K. Gürpinar, O. Atakol, G. G. Ornelas, W. Brackenridge, P. Lloveras, M. Barrio, A. Aznar, J. L. Tamarit, W. Zheng, L. Balicas and B. Emre, Giant and Reversible Barocaloric Effect in Trinuclear Spin-Crossover Complex $\text{Fe}_3(\text{bntrz})_6(\text{tcnset})_6$, *Adv. Mater.*, 2021, **33**, 2008076.

50 P. Lloveras and J.-L. Tamarit, Advances and obstacles in pressure-driven solid-state cooling: A review of barocaloric materials, *MRS Energy Sustain.*, 2021, **8**, 3–15.

51 C. Bartual-Murgui, E. Natividad and O. Roubeau, Critical assessment of the nature and properties of Fe(II) triazole-based spin-crossover nanoparticles, *J. Mater. Chem. C*, 2015, **3**, 7916–7924.

52 T. Delgado, C. Enachescu, A. Tissot, A. Hauser, L. Guénée and C. Besnard, Evidencing size-dependent cooperative effects on spin crossover nanoparticles following their HS \rightarrow LS relaxation, *J. Mater. Chem. C*, 2018, **6**, 12698–12706.

53 M. Giménez-Marqués, M. L. García-Sanz De Larrea and E. Coronado, Unravelling the chemical design of spin-crossover nanoparticles based on iron(II)-triazole coordination polymers: Towards a control of the spin transition, *J. Mater. Chem. C*, 2015, **3**, 7946–7953.

54 L. Moulet, N. Daro, C. Etrillard, J. F. Létard, A. Grosjean and P. Guionneau, Rational control of spin-crossover particle sizes: From nano-to micro-rods of $[\text{Fe}(\text{Htrz})_2(\text{trz})](\text{BF}_4)$, *Magnetochemistry*, 2016, **2**, 10.

55 C. Faulmann, J. Chahine, I. Malfant, D. De Caro, B. Cormary and L. Valade, A facile route for the preparation of nanoparticles of the spin-crossover complex $[\text{Fe}(\text{Htrz})_2(\text{trz})](\text{BF}_4)$ in xerogel transparent composite films, *Dalt. Trans.*, 2011, **40**, 2480–2485.

56 A. Suzuki, M. Fujiwara and M. Nishijima, High spin/low spin phase transitions of a spin-crossover complex in the emulsion polymerization of trifluoroethylmethacrylate (TFEMA) using PVA as a protective colloid, *Colloid Polym. Sci.*, 2008, **286**, 525–534.

57 A. B. Gaspar, G. Molnár, A. Rotaru and H. J. Shepherd, Pressure effect investigations on spin-crossover coordination compounds, *Comptes Rendus Chim.*, 2018, **21**, 1095–1120.

58 J. Linares, E. Codjovi and Y. Garcia, Pressure and temperature spin crossover sensors with optical detection, *Sensors*, 2012, **12**, 4479–4492.

59 J. H. Askew, D. M. Pickup, G. O. Lloyd, A. V. Chadwick and H. J. Shepherd, Exploring the Effects of Synthetic and Postsynthetic Grinding on the Properties of the Spin Crossover Material $[\text{Fe}(\text{atrz})_3](\text{BF}_4)_2$ (atrz = 4-Amino-4H-1,2,4-Triazole), *Magnetochemistry*, 2020, **6**, 44.

60 D. Nieto-Castro, F. A. Garcés-Pineda, A. Moneo-Corcuera, B. Pato-Doldan, F. Gispert-Guirado, J. Benet-Buchholz and J. R. Galán-Mascarós, Effect of Mechanochemical Recrystallization on the Thermal Hysteresis of 1D Fe^{II} -triazole Spin Crossover Polymers, *Inorg. Chem.*, 2020, **59**, 7953–7959.

61 J. Kroeber, J.-P. Audiere, R. Claude, E. Codjovi, O. Kahn, J. G. Haasnoot, F. Groliere, C. Jay and A. Bousseksou, Spin Transitions and Thermal Hysteresis in the Molecular-Based Materials $[\text{Fe}(\text{Htrz})_2(\text{trz})](\text{BF}_4)$ and $[\text{Fe}(\text{Htrz})_3](\text{BF}_4)_2 \cdot \text{cnctdot.H}_2\text{O}$ (Htrz = 1,2,4-4H-triazole; trz = 1,2,4-triazolato), *Chem. Mater.*, 1994, **6**, 1404–1412.

62 A. Grosjean, P. Négrier, P. Bordet, C. Etrillard, D. Mondieig, S. Pechev, E. Lebraud, J. Létard and P. Guionneau, Crystal Structures and Spin Crossover in the Polymeric Material $[\text{Fe}(\text{Htrz})_2(\text{trz})](\text{BF}_4)$ Including Coherent-Domain Size Reduction Effects, *Eur. J. Inorg. Chem.*, 2013, 796–802.

63 M. Palluel, L. El Khoury, N. Daro, S. Buffière, M. Josse, M. Marchivie and G. Chastanet, Rational direct synthesis of $[\text{Fe}(\text{Htrz})_2(\text{trz})](\text{BF}_4)$ polymorphs: temperature and concentration effects, *Inorg. Chem. Front.*, 2021, **8**, 3697–3706.

64 V. Petříček, M. Dušek and L. Palatinus, Crystallographic Computing System JANA2006: General features, *Zeitschrift für Krist. - Cryst. Mater.*, 2014, **229**, 345–352.

65 J. Rodriguez-Carvajal, Recent advances in magnetic structure determination by neutron powder diffraction, *Phys. B Condens. Matter*, 1993, **192**, 55–69.

66 A. Grosjean, N. Daro, S. Pechev, L. Moulet, C. Etrillard, G. Chastanet and P. Guionneau, The Spin-Crossover Phenomenon at the Coherent-Domains Scale in 1D Polymeric Powders: Evidence for Structural Fatigability, *Eur. J. Inorg. Chem.*, 2016, 1961–1966.

67 A. Grosjean, *Matériaux polymériques 1D à transition de spin: investigations structurales multi-échelles*, PhD. thesis, University of Bordeaux, 2013.

68 M. Marchivie, P. Guionneau, J.-F. Létard and D. Chasseau, Towards direct correlation between spin crossover and structural properties in iron II complexes, *Acta Crystallogr. B*, 2003, **59**, 479–486.



69 A. Grosjean, N. Daro, S. Pechev, C. Etrillard, G. Chastanet and P. Guionneau, Crystallinity and Microstructural Versatility in the Spin-Crossover Polymeric Material $[\text{Fe}(\text{Htrz})_2(\text{trz})](\text{BF}_4)$, *Eur. J. Inorg. Chem.*, 2018, 429–434.

70 C. P. Slichter and H. G. Drickamer, Pressure-Induced Electronic Changes in Compounds of Iron, *J. Chem. Phys.*, 1972, 56, 2142–2160.

71 M. Sorai and S. Seki, Phonon coupled cooperative low-spin ${}^1\text{A}_1$ high-spin ${}^5\text{T}_2$ transition in $[\text{Fe}(\text{phen})_2(\text{NCS})_2]$ and $[\text{Fe}(\text{phen})_2(\text{NCSe})_2]$ crystals, *J. Phys. Chem. Solids*, 1974, 35, 555–570.

72 O. Roubeau, M. Castro, R. Burriel, J. G. Haasnoot and J. Reedijk, Calorimetric investigation of triazole-bridged Fe(II) spin-crossover one-dimensional materials: Measuring the cooperativity, *J. Phys. Chem. B*, 2011, 115, 3003–3012.

