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Intra-host π – π interactions in crown ether complexes revealed by cryogenic ion mobility-mass spectrometry†

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Cryogenic ion mobility-mass spectrometry was performed to investigate the relative abundance of conformers of dinaphtho-24-crown-8 (DN24C8) complexes with alkali metal cations M^+ ($M = \text{Li, Na, K, Rb, and Cs}$). The “closed” conformers of $M^+(\text{DN24C8})$ with short distances between two naphthalene rings in the crown ethers were predominantly observed for all complexes at 86 K. The two noncovalent interactions, host–guest and intra-host interactions, were analyzed separately by density functional theory calculations to reveal the origin of the stability of the closed conformers. As a result, it was revealed that the intra-host π – π interactions have a more critical role in determining the stability of the conformers than the host–guest interactions. The closed conformers of $M^+(\text{DN24C8})$ also have wider regions of the π – π interactions than those of the $M^+(\text{dibenzo-24-crown-8})$ complexes.

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1. Introduction

Noncovalent interactions (NCIs) influence the conformation of a variety of molecules and the formation of supramolecular systems.^{1–4} For example, π – π interactions play an important role in stabilizing the higher-order structure of DNA double strands in addition to hydrogen bonding between nucleobases. Biomolecules are designed to exhibit a variety of functionalities through a skillful combination of weak interactions. The in-depth understanding of NCIs such as π – π interactions could therefore provide guidelines for designing molecules with biomolecular-like functionalities. The key to the functional expression of soft molecules composed of weak NCIs is their flexibility. The structures of the soft molecules fluctuate among several conformers using thermal energy. This flexibility in structures makes detailed analysis of their conformation difficult. Therefore, supramolecular structures are mainly studied by X-ray crystallography in the solid phase without structural fluctuations.⁵

In recent years, gas-phase laser photodissociation spectroscopy has been actively performed to study the structures of cold supramolecular ions.⁶ In particular, the vibrational temperature of an isolated ion is lowered to ~ 10 K by employing a cryogenic ion trap. This cryogenic technique allows us to

observe vibrationally resolved sharp absorption spectra of supramolecules. However, complicated spectra are often obtained also at low temperatures due to the coexistence of multiple conformers of the flexible supramolecular ions. Therefore, it is still difficult to determine the structures and the abundance ratio of the conformers even from the cryogenic spectra in which the bands of the conformers are overlapped.

Host–guest complex is one of the simple models for a flexible supramolecule. Ion mobility-mass spectrometry (IM–MS) has been used to study conformers of host–guest complexes in the gas phase at room temperature.^{7–9} IM–MS is a powerful technique to separate conformers in the gas phase.^{7,10,11} Size- and conformer-selected absorption spectra of supramolecular ions can be obtained by the coupling of IM–MS with photodissociation spectroscopy.^{12–14} In the conventional drift-tube IM–MS, we measured the time distribution of the ions to pass through the drift tube (several milliseconds) and determined collision cross sections (CCSSs) of the ions with a buffer gas. Conformations of the complex ions can be determined from the comparison between experimental and theoretical CCSSs of the ions. Cryogenic IM–MS has also been performed using a drift tube cooled by liquid nitrogen.^{15–32}

Crown ether complexes with metal ions are one of the most common host–guest compounds. The encapsulation structures of crown ether complexes in crystal forms were studied by X-ray diffraction.^{33–36} Nuclear magnetic resonance studies of structures of the conformations of the dibenzo-30-crown-10 (DB30C10) complexes with alkali metal ions, M ($M = \text{Na and K}$), $M^+(\text{DB30C10})$, in solution also showed the same configuration of DB30C10 as that reported for the crystal.^{37,38}

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† Electronic supplementary information (ESI) available: Reduced density gradient isosurfaces and geometrical coordinates of conformers of $M^+(\text{dinaphtho-24-crown-8})$ ($M = \text{Li, Na, K, Rb, and Cs}$) in Table 1. See DOI: <https://doi.org/10.1039/d4cp00835a>



The conformation of the crown ether complexes in the gas phase was studied by ultraviolet photodissociation (UVPD)^{26,39–46} and IM–MS.^{8–10,26–30} Previously, we studied the conformations of $M^+(\text{DB24C8})$ ($M = \text{Na}, \text{K}, \text{Rb}, \text{and Cs}$) using cryogenic IM–MS. The coexistence of two conformers, **open** and **closed** conformers, were confirmed in the study, which have long and short distances between the two benzene rings in a molecule, respectively.³⁰ These conformers were also assigned by a UVPD spectrum of $M^+(\text{DB24C8})$ at ~ 10 K.⁴⁶ For the **closed** conformer of $K^+(\text{DB24C8})$, the relative orientation of the two benzene rings in the complex was found to be quite close to that of the benzene dimer. Therefore, this **closed** conformer was concluded to be highly stabilized by the strong π – π interaction between the benzene rings due to its special benzene dimer-like conformation.³⁰

In the present study, we investigated the effect of enhancement of π – π interaction by expanding the aromatic rings in the complexes on the conformation and the stability of the complexes. The conformation of the complexes of alkali metal cations ($\text{Li}^+, \text{Na}^+, \text{K}^+, \text{Rb}^+, \text{and Cs}^+$) with dinaphtho-24-crown-8 (DN24C8) shown in Fig. 1 was determined by combining the results of cryogenic IM–MS experiments at 86 K in the gas phase and quantum chemical calculations. The relative abundance ratio of the **closed** conformer of the $M^+(\text{DN24C8})$ complexes was larger than that of the $M^+(\text{DB24C8})$ complexes. The extraordinary stability of the **closed** conformers of the DN24C8 complexes was discussed based on a quantitative evaluation of the strength of intra-host π – π interactions between the two aromatic rings in the complexes.

2. Experimental and theoretical methods

2.1. Experimental methods

Details of the experimental setup were given elsewhere.^{30,31} $M^+(\text{DN24C8})$ complex ions were generated by electrospray ionization. 0.1 mM alkali metal chlorides ($\text{LiCl}, \text{NaCl}, \text{KCl}, \text{RbCl}, \text{or CsCl}$) and DN24C8 were dissolved in methanol. The solution was delivered to a metal capillary at $2.0 \mu\text{L min}^{-1}$. High voltage (+2.6 kV) was applied between the metal capillary and a heated desolvation capillary (~ 350 K). The entrance of an ion funnel was positioned close to the exit of the heated capillary. After focusing the ions by the funnel, ions were introduced into a quadrupole ion trap (QIT). Ions were accumulated for 40 ms in the QIT. To trap ions efficiently, helium gas was introduced to the QIT with 16 sccm. After accumulation, ions

were injected into the cryogenic ion drift tube by a pulsed electric field at a given time ($t = t_0$). The temperature and pressure of helium buffer gas in the drift tube were 86 K and 0.600 torr, respectively. A static electric field ($E = 6.1 \text{ V cm}^{-1}$) was applied in the drift tube for ion drifting. The E/N value was 9.0 Td under the above conditions (N is the number density of the buffer gas, $1 \text{ Td} = 10^{-17} \text{ V cm}^2$). Thus, the low-field condition (*ca.* $E/N < 10$ Td) was achieved in our experiments. After passing through the drift tube, the ions were transported to the acceleration region of the time-of-flight mass spectrometer (TOF-MS). The ions were accelerated to ~ 4 keV by pulsed electric fields at a given time later from the pulse for the injection into the drift tube: $t = t_0 + \Delta t$. The ions were detected by a microchannel plate of the TOF-MS. The delay time between the two pulses, Δt , was defined as “arrival time”. The drift velocity of the ions in the drift tube, v_d , was calculated numerically from the measured arrival time. It is known that v_d is proportional to E , and the proportional constant (K) is called ion mobility.^{47,48} The mobility K depends on the number density of the buffer gas N . To compare the mobility under different experimental conditions, the reduced mobility K_0 is defined as $K_0 = K(N/N_0)$, where N_0 is Loschmidt's number. From the Mason–Schamp equation, the reduced mobility K_0 in the drift tube was given as

$$K_0 = \frac{3q}{16N_0} \left(\frac{2\pi}{k_B \mu T_{\text{eff}}} \right)^{\frac{1}{2}} \frac{1}{\Omega} \quad (1)$$

where q is the charge of the ion, k_B is the Boltzmann constant, μ is the reduced mass of the ion and the buffer gas atom, T_{eff} is the effective temperature of the ions, and Ω is a CCS.^{47,48} The effective temperature is given by $T_{\text{eff}} = T_{\text{BG}} + m_B v_d^2 / 3k_B$, where T_{BG} is the buffer gas temperature, and m_B is the mass of buffer gas. Therefore, the CCS of the ion is calculated from the measured arrival time of the ion, Δt . In the IM–MS, we obtained a set of TOF mass spectra sequentially by scanning the arrival time. We obtained a plot of arrival time distribution (ATD), in which the total ion intensity of a certain TOF peak was shown as a function of the arrival time. We converted the ATDs to CCS distributions and analyzed them by fitting Gaussian functions with a width consistent with the CCS resolution of the apparatus.

2.2. Computational methods

Density functional theory (DFT) calculations were performed to determine the structures of the $M^+(\text{DN24C8})$ complex ions. The initial conformational search was performed by the CONFLEX 8 program with the MMFF94s force field.^{49,50} In total 20 conformers with relatively low steric energy on the CONFLEX were further optimized at the M05-2X/6-311++G(d,p) level with the Stuttgart RLC ECP for Rb and Cs atoms^{51–53} by the Gaussian 16 program.⁵⁴ The vibrational analysis, zero-point energy correction, and natural population analysis were also performed at the same level. The theoretical CCSs of candidate structures of $M^+(\text{DN24C8})$ obtained by DFT calculations were calculated by the trajectory method in the MOBCAL program.⁵⁵ In the

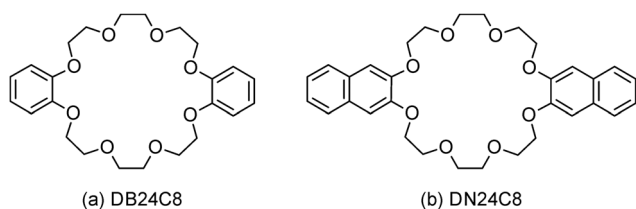


Fig. 1 Skeletal formulas of (a) DB24C8 and (b) DN24C8.



trajectory method, the Lennard-Jones parameters between constituent atoms of a complex ion and a He atom ($\epsilon_{X-\text{He}}$ and $\sigma_{X-\text{He}}$) are needed. In the present study, we used the previously reported $\epsilon_{X-\text{He}}$ and $\sigma_{X-\text{He}}$ for X = C, H, and O.⁵⁶ The values for X = Li, Na, K, Rb, and Cs were determined from calculations of X-He diatomic systems at the CCSD(T)/6-311+G(d,p) level with the Stuttgart RLC ECP for Rb and Cs atoms ($\epsilon_{\text{Li-He}} = 0.80165$ meV, $\sigma_{\text{Li-He}} = 1.7610$ Å, $\epsilon_{\text{Na-He}} = 0.7511$ meV, $\sigma_{\text{Na-He}} = 2.1925$ Å, $\epsilon_{\text{K-He}} = 0.5151$ meV, $\sigma_{\text{K-He}} = 2.6368$ Å, $\epsilon_{\text{Rb-He}} = 0.3224$ meV, $\sigma_{\text{Rb-He}} = 2.9985$ Å, $\epsilon_{\text{Cs-He}} = 0.2982$ meV, $\sigma_{\text{Cs-He}} = 3.1956$ Å).³⁰ Standard deviations of theoretical CCSs were less than $\pm 0.7\%$.

π - π interactions between the two aromatic rings in the $\text{M}^+(\text{DB24C8})$ and $\text{M}^+(\text{DN24C8})$ complexes were visualized by the NCI method⁵⁷ in the Multiwfn program.⁵⁸ The NCI method can identify the region where noncovalent interactions are working based on the values of the reduced density gradient (RDG) and the product of the sign of the second largest eigenvalue of electron density Hessian matrix and electron density ($\text{sign}(\lambda_2) \cdot \rho$). The extent of the π - π interactions can be distinguished from the fact that the RDG value is small (< 0.5 in atomic units) and the $\text{sign}(\lambda_2) \cdot \rho$ value is about zero. The strength of the π - π interactions was estimated by the domain analysis implemented in the Multiwfn program.

3. Results and discussion

3.1. Structure assignments of $\text{M}^+(\text{DN24C8})$ from IM-MS measurements

Fig. 2 shows CCS distributions of $\text{M}^+(\text{DN24C8})$ (M = Li, Na, K, Rb, and Cs) complexes measured with He buffer gas at 86 K in the drift tube. For $\text{Na}^+(\text{DN24C8})$, three Gaussian functions were necessary for the fittings (Fig. 2b). Therefore, at least three conformers with different CCSs coexist for this complex. For $\text{M}^+(\text{DN24C8})$ (M = Li, K, Rb, and Cs), CCS distributions were fitted well with one Gaussian function (Fig. 2a and c-e). Structures of the most stable conformers are shown in Fig. 3. Experimental CCSs ($^{\text{DT}}\text{CCS}_{\text{He},86}$), relative intensities determined from the CCS distributions, theoretical CCSs, and relative Gibbs free energies (ΔG_{86}) of the stable conformers are summarized in Table 1. In the most stable conformers of $\text{M}^+(\text{DN24C8})$, the distance between the centers of the two naphthalene rings in the complexes (r_c) was short. These conformers have the “closed” form of the DN24C8. For $\text{Na}^+(\text{DN24C8})$, conformers with the second and third lowest relative Gibbs free energies have longer r_c . Based on the conformation of DN24C8, we hereafter call the second and third stable conformers “closed’” and “open”, respectively.

The theoretical CCS of the **closed** conformer of $\text{Na}^+(\text{DN24C8})$ (185.0 Å²) was close to the experimental CCS of the strong band (186.4 ± 0.3 Å²). The experimental CCS of the shoulder on the strong band (191.7 ± 0.3 Å²) was in good agreement with the theoretical CCS of the **closed’** conformer (191.1 Å²). The larger CCS band had a relatively weak peak at 205.4 Å². The theoretical CCS of the **open** conformer (203.6 Å²) was close to this peak. The **closed’** and **open** conformers were less stable in Gibbs free

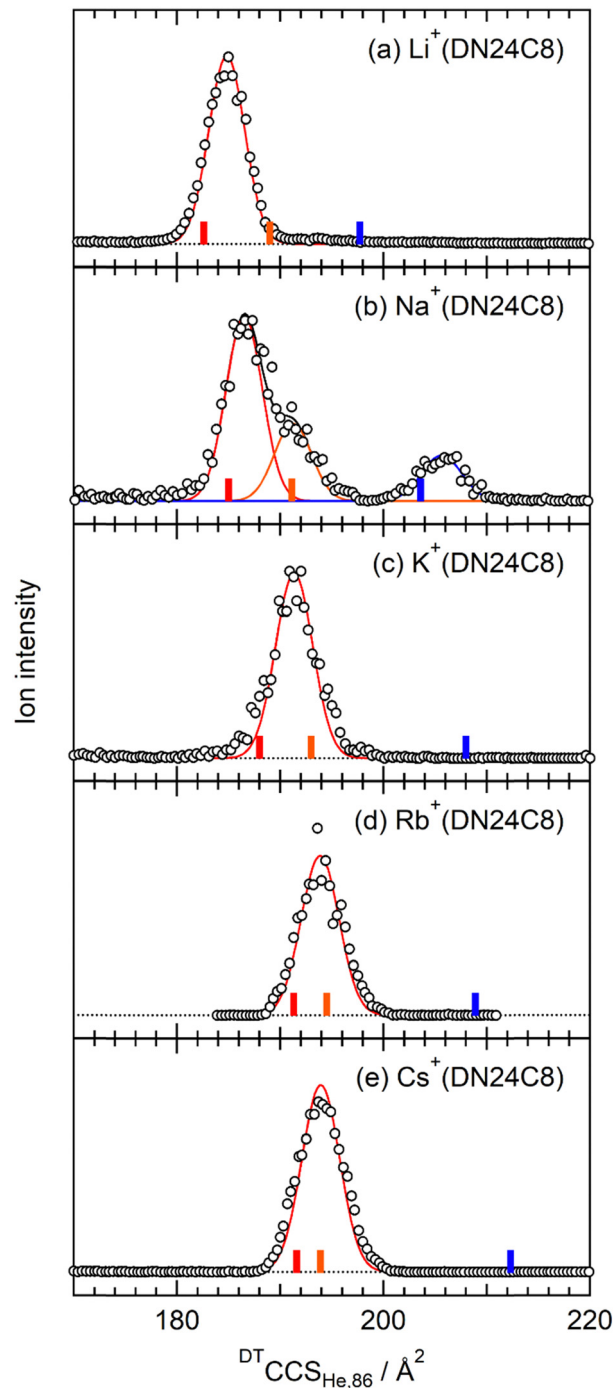


Fig. 2 CCS distributions of the $\text{M}^+(\text{DN24C8})$ (M = (a) Li, (b) Na, (c) K, (d) Rb, and (e) Cs) complexes obtained at 86 K. Curves are Gaussian functions for fitting the experimental data points (circles). Vertical red, orange, and blue bars indicate the theoretical CCSs of the **closed**, **closed’**, and **open** conformers.

energy than the **closed** conformer by 2.5 and 4.0 kJ mol^{−1}, respectively. It is worth noting that the Gibbs free energies are sensitive to the DFT functional. For example, the relative free energy with respect to the most stable conformer at 86 K, ΔG_{86} , is 4.0 kJ mol^{−1} for the **open** conformer of $\text{Na}^+(\text{DN24C8})$ in the DFT calculations at the M05-2X/6-311++G(d,p) level as shown in



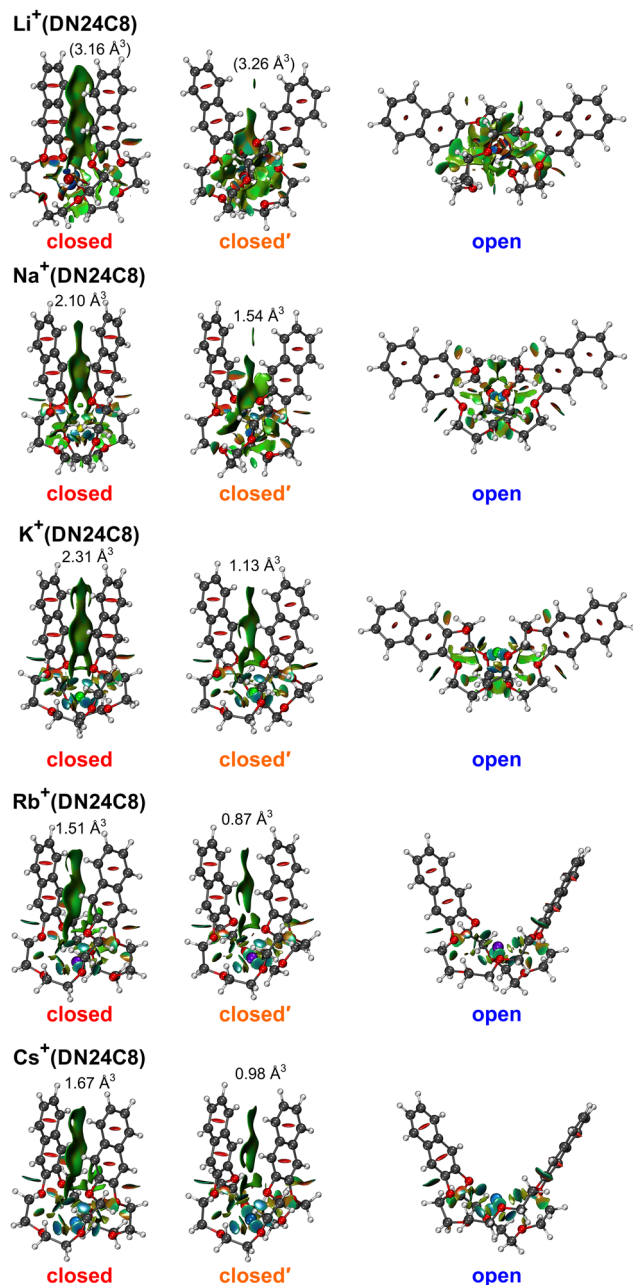


Fig. 3 The most stable **closed**, **closed'**, and **open** conformers of $M^+(\text{DN24C8})$ ($M = \text{Li, Na, K, Rb, and Cs}$) calculated at the M05-2X/6-311++G(d,p) level with the Stuttgart RLC ECP for Rb and Cs atoms.

Table 1, whereas it amounted to 15.3 kJ mol^{-1} at $\omega\text{B97X-D}/6\text{-311++G(d,p)}$. However, these energy differences were consistent with the order of relative ion intensity in the IM–MS regardless of the DFT functional.

For the $M^+(\text{DN24C8})$ ($M = \text{Li, K, Rb, and Cs}$) complexes, only one band was observed in each CCS distribution (Fig. 2a and c–e). In the calculations, the most stable conformers of these complexes were the **closed** ones similar to the most stable structure of the $\text{Na}^+(\text{DN24C8})$ complex. However, there are small differences in the

relative orientations of the two naphthalene rings in the **closed** conformers of the complexes: the rings are slightly tilted outward in $\text{Li}^+(\text{DN24C8})$ and $\text{Na}^+(\text{DN24C8})$, nearly parallel in $\text{K}^+(\text{DN24C8})$, and tilted inward in $\text{Rb}^+(\text{DN24C8})$ and $\text{Cs}^+(\text{DN24C8})$. The theoretical CCSs of the **closed** conformers were in good agreement with the experimental CCSs ($\sim 190 \text{ \AA}^2$). The theoretical CCSs of the **open** conformers ($\sim 210 \text{ \AA}^2$) were found to be much larger than the experimental CCS peaks. In addition, the **open** conformers were less stable in Gibbs free energy than the **closed** conformers by more than 6.2 kJ mol^{-1} at 86 K. This energy difference corresponded to the thermal abundance ratio of the **open** conformer being less than 0.02%. Since the **closed'** conformers were as unstable as the **open** conformers in Gibbs free energy at 86 K, probably due to steric hindrance, the thermal abundance ratio of the **closed'** conformers was lower than that of the **open** conformers. Thus, only the **closed** conformers were observed for $M^+(\text{DN24C8})$ ($M = \text{Li, K, Rb, and Cs}$) unlike $\text{Na}^+(\text{DN24C8})$.

3.2. Relative conformer stability analyzed by host–guest interactions

In our previous study on the DB24C8 complexes, the **open** conformers were strongly observed in $M^+(\text{DB24C8})$ ($M = \text{Na, K, Rb, and Cs}$) in addition to the **closed** conformers.³⁰ In contrast, the **closed** conformers were predominant for all five guest ions of the DN24C8 complexes in the present study. The $M^+(\text{DB24C8})$ and $M^+(\text{DN24C8})$ complexes are stabilized by strong host–guest interactions and relatively weak intra-host interactions. To show the differences in conformer stability between the $M^+(\text{DB24C8})$ and the $M^+(\text{DN24C8})$ complexes, host–guest and intra-host interactions were analyzed separately. In this section, we first analyze the host–guest interactions, which are governed by the positional relationships between the eight oxygen atoms and the guest M^+ ions in the complexes. The lengths of the $M^+\text{--O}$ bonds, $R(M^+\text{--O})$, in the most stable **closed**, **closed'**, and **open** conformers are plotted in Fig. 4. The $R(M^+\text{--O})$ lengths showed surprisingly close trends in the DB24C8 and DN24C8 complexes in the figure. The lengths were hardly dependent on both the conformation (**closed**, **closed'**, or **open**) and the host molecule (DB24C8 or DN24C8) except for the Li^+ complexes. For example, the average $R(M^+\text{--O})$ for the **closed** conformer of $\text{Na}^+(\text{DB24C8})$ was 2.48 \AA . This value is very close to that of the **open** conformer of $\text{Na}^+(\text{DB24C8})$ (2.47 \AA) and the **closed** (2.50 \AA), **closed'** (2.48 \AA), and **open** (2.47 \AA) conformers of $\text{Na}^+(\text{DN24C8})$. In contrast, the average $R(M^+\text{--O})$ depended sensitively on the guest ions. The average $R(M^+\text{--O})$ of the complexes have an order of $\text{Na}^+ (\sim 2.5 \text{ \AA}) < \text{K}^+ (\sim 2.8 \text{ \AA}) < \text{Rb}^+ (\sim 3.0 \text{ \AA}) < \text{Cs}^+ (\sim 3.1 \text{ \AA})$, which corresponds to the order of the ionic radii of the guest ions.⁵⁹ In summary, the host–guest interactions are exerted to almost the same extent for the complexes with the same guest ions. Therefore, the differences in conformer stability between the DB24C8 and the DN24C8 complexes could not be explained from the host–guest interactions. In the next section, we analyze the intra-host interactions, especially the π – π interactions between the two aromatic rings.



Table 1 Experimental CCSs ($^{DT}CCS_{He,86}$) and relative intensities of bands in CCS distributions of the $M^+(DN24C8)$ complexes measured with He buffer gas at 86 K, and CCSs and relative Gibbs free energies (ΔG_{86}) of stable conformers at 86 K obtained by DFT calculations. Standard deviations of $^{DT}CCS_{He,86}$ were estimated from five independent measurements. Standard deviations of calculated CCSs were less than $\pm 0.7\%$

Complexes	Experiments		Theory		
	$^{DT}CCS_{He,86}/\text{\AA}^2$	Relative intensity	CCS/ \AA^2	Conformers	$\Delta G_{86}/\text{kJ mol}^{-1}$
$Li^+(DN24C8)$	184.3 ± 0.4	1	182.6	closed	0
			189.0	closed'	13.9
			197.7	open	12.3
$Na^+(DN24C8)$	186.4 ± 0.3	1	185.0	closed	0
			191.1	closed'	2.5
			203.6	open	4.0
$K^+(DN24C8)$	191.2 ± 0.3	1	188.0	closed	0
			193.0	closed'	13.8
			208.0	open	12.8
$Rb^+(DN24C8)$	194.0 ± 0.6	1	191.3	closed	0
			194.5	closed'	9.9
			208.9	open	8.0
$Cs^+(DN24C8)$	193.7 ± 0.3	1	191.6	closed	0
			193.9	closed'	9.2
			212.3	open	6.2

For $Li^+(DB24C8)$ and $Li^+(DN24C8)$, the $R(M^+-O)$ lengths took a wide range of values because the Li^+ is too small to form an 8-coordination structure with the crown ethers. In the **closed** conformer of $Li^+(DN24C8)$, three oxygen atoms were located within 2.1 Å of the guest ion. However, in the **open** conformer there was only one oxygen atom within 2.1 Å of the guest ion, and there was no oxygen atom within 2.1 Å of the guest ion in the **closed'** conformer. Only the **closed** conformer was observed for $Li^+(DN24C8)$, in part because these oxygen positions highly stabilize the **closed** conformer. Although the $R(M^+-O)$ ($M = Na, K, Rb$, and Cs) lengths were similarly distributed for each complex in Fig. 4, two oxygen atoms (overlapped) and one oxygen atom were located more than 0.15 Å far from the guest ion in the **closed** and **closed'** conformers of $Na^+(DN24C8)$,

respectively. These oxygen positions should destabilize the most stable **closed** and **closed'** conformers, and consequently, the relative energies of the three (**closed**, **closed'**, and **open**) conformers become closer in this complex than in other complexes. At present, the reason why the **closed'** and the **open** conformers of the $Na^+(DN24C8)$ were observed can thus be explained from the consideration of the $R(M^+-O)$ distributions.

3.3. Quantitative evaluation of intra-host π - π interactions

The intra-host π - π interactions between the two aromatic rings of the most stable **closed** conformers of $M^+(DN24C8)$ are evaluated in this section by extracting the relative orientations and distances between the two naphthalene rings in the

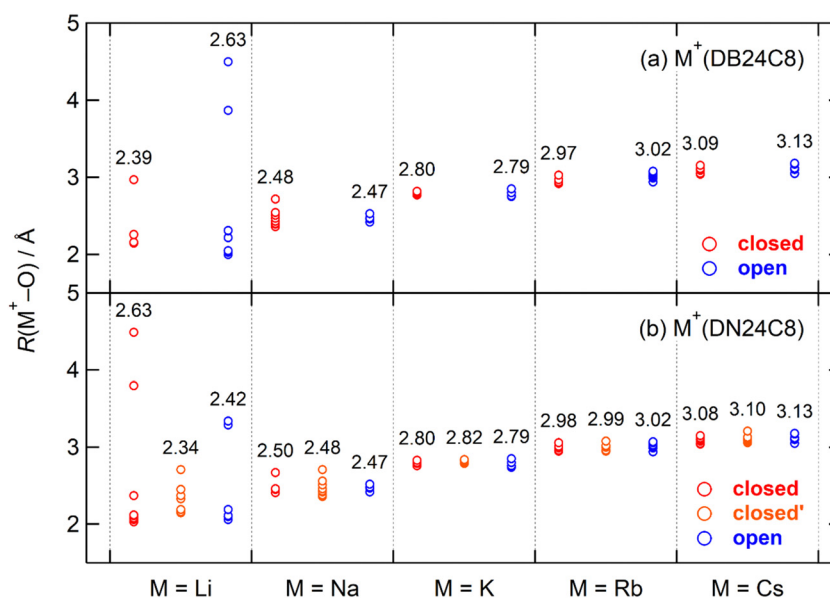


Fig. 4 Lengths of bonds between an alkali metal ion and oxygen atoms, $R(M^+-O)$, in (a) $M^+(DB24C8)$ and (b) $M^+(DN24C8)$ complexes obtained by DFT calculations. Red, orange, and blue markers are $R(M^+-O)$ in **closed**, **closed'**, and **open** conformers of the complexes, respectively. Some plots appear to overlap due to the high symmetry of the conformation. The average $R(M^+-O)$ for each conformer is shown near the plot.



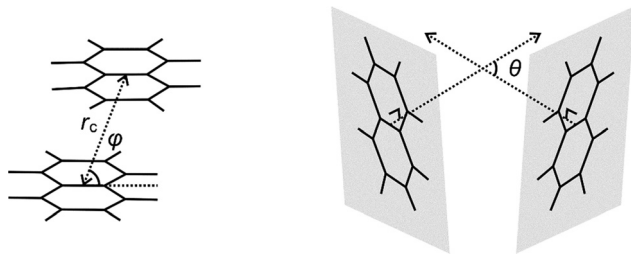


Fig. 5 Parameters of naphthalene ring positional relationship (r_c : distance between naphthalene ring centers, θ : angle formed by normal vectors of two naphthalene planes, ϕ : angle between the naphthalene plane and the central line).

Table 2 Parameter values of naphthalene ring positional relationship in stable **closed** conformers for $M^+(\text{DN24C8})$ and naphthalene dimer. The parameter values for naphthalene dimer were obtained from the parallel-displaced structure calculated at MP2/6-31+G* level⁶⁰

Complexes	$r_c/\text{\AA}$	$\theta/\text{deg.}$	$\phi/\text{deg.}$
$\text{Li}^+(\text{DN24C8})$	3.82	11.4	57.7, 67.6
$\text{Na}^+(\text{DN24C8})$	4.01	1.8	55.3
$\text{K}^+(\text{DN24C8})$	3.86	0.1	61.2
$\text{Rb}^+(\text{DN24C8})$	4.30	16.2	44.2, 60.2
$\text{Cs}^+(\text{DN24C8})$	4.25	16.3	45.3, 61.7
Naphthalene dimer	3.87	0	64.9

complexes from the calculated structures. Three positional relationship parameters (r_c , θ , ϕ) were defined similarly to the previous analysis of the DB24C8 complexes³⁰ as shown in Fig. 5. The r_c is the distance between the naphthalene ring centers, θ is the angle formed by normal vectors of the two naphthalene planes, and ϕ is the angle between the naphthalene plane and the central line. Table 2 summarizes the values of these parameters in the $M^+(\text{DN24C8})$ complexes. For comparison, the parameters of parallel-displaced conformer of a naphthalene dimer (C_{10}H_8)₂ obtained by *ab initio* calculation⁶⁰ are also shown in Table 2. As a result, the relative positions of

the naphthalene rings in $\text{K}^+(\text{DN24C8})$ were found to be quite close to that of the naphthalene dimer. Therefore, this **closed** conformer was concluded to show extraordinary stability because of the strong π - π interaction between the naphthalene rings due to its special naphthalene dimer-like conformation. These results were nearly identical to those of the **closed** conformer of $\text{K}^+(\text{DB24C8})$ becoming particularly stable due to the special orientation of benzene rings in the complex.³⁰

To visualize the π - π interaction between the two aromatic rings in the complexes, the NCI analysis was conducted. The RDG isosurfaces of the **closed** conformers of $\text{K}^+(\text{DB24C8})$ and $\text{K}^+(\text{DN24C8})$ are shown in Fig. 6. These surfaces were generated by evaluating the electron density calculated by DFT calculations and RDGs on cuboid grids with a 0.1 au step size. The green distribution indicates weak attractive interactions such as π - π interaction. In the **closed** conformers, π - π interactions work between the two aromatic rings as shown in Fig. 6. $\text{K}^+(\text{DN24C8})$ had wider regions of π - π interactions than $\text{K}^+(\text{DB24C8})$. Domain analysis was performed to confirm the difference in the size of the π - π interaction regions. The volume of the π - π interaction region between the two aromatic rings determined by the domain analysis in the Multiwfn program was 2.31 \AA^3 for $\text{K}^+(\text{DN24C8})$, which was larger than that of $\text{K}^+(\text{DB24C8})$ (1.44 \AA^3). Furthermore, the volume of the π - π interaction region for $\text{K}^+(\text{DN24C8})$ was equivalent to 89% of that for the parallel-displaced naphthalene dimer (2.59 \AA^3). This large interaction volume means that the π - π interaction between the two naphthalene rings of $\text{K}^+(\text{DN24C8})$ is comparable to the π - π interaction energy of a naphthalene dimer. These results suggest that the DN24C8 complexes are more stabilized than the DB24C8 complexes due to the large π - π interaction between the two aromatic rings in the complexes, supporting the experimental results that the **closed** conformers were predominantly observed in IM-MS.

4. Conclusions

We performed cryogenic ion mobility-mass spectrometry of the dinaphtho-24-crown-8 (DN24C8) complexes with alkali metal ions M^+ ($M = \text{Li}, \text{Na}, \text{K}, \text{Rb}, \text{and Cs}$) to investigate the effect of π - π interaction in the complex on the relative abundance of the conformers. The “**closed**” conformer with short distance between two naphthalene rings in the complex was predominant in the distributions of collision cross sections of $M^+(\text{DN24C8})$. For $\text{Na}^+(\text{DN24C8})$, the “**closed**” and the “**open**” conformers with longer naphthalene-naphthalene distance were also observed. The order of conformer stability is explained from the host-guest interactions and intra-host interactions. The attraction between oxygen atoms of the crown ether and the guest M^+ ion was strongly dependent on the guest M^+ ion and less dependent on the host molecule or the conformation. For the **closed** conformer of $\text{Na}^+(\text{DN24C8})$, two oxygen atoms were located far from the Na^+ ion, which destabilized the **closed** conformer and made the **closed'** and the

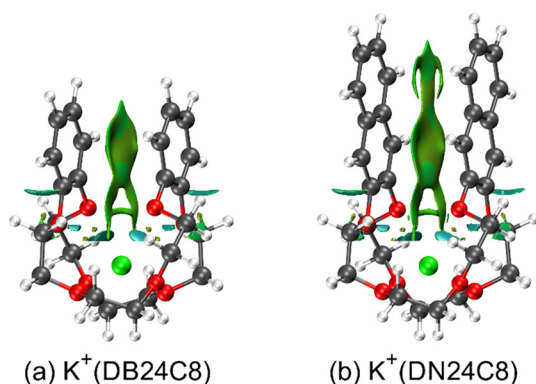


Fig. 6 RDG isosurfaces (RDG = 0.5 au) showing the π - π interaction between the two aromatic rings in **closed** conformers of (a) $\text{K}^+(\text{DB24C8})$ and (b) $\text{K}^+(\text{DN24C8})$. The surfaces are colored on a blue-green-red scale according to values of $\text{sign}(\lambda_2) \cdot \rho$, ranging from -0.025 to 0.02 au. The green distribution indicates weak attractive interactions such as π - π interaction.



open conformers relatively stable. The intra-host attractive π - π interaction region between the two aromatic rings of the **closed** conformer was wider in $M^+(\text{DN24C8})$ than in $M^+(\text{DB24C8})$. The stability of the crown ether complex was revealed to be controlled by a balance of noncovalent interactions.

Conflicts of interest

There are no conflicts to declare.

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References

- J. L. Atwood, G. A. Koutsantonis and C. L. Raston, *Nature*, 1994, **368**, 229.
- C. K. McLaughlin, G. D. Hamblin and H. F. Sleiman, *Chem. Soc. Rev.*, 2011, **40**, 5647.
- N. C. Seeman, *Annu. Rev. Biochem.*, 2010, **79**, 65.
- S. Zuluaga, P. Canepa, K. Tan, Y. J. Chabal and T. Thonhauser, *J. Phys.: Condens. Matter*, 2014, **26**, 133002.
- T. Akutagawa, K. Shitagami, S. Nishihara, S. Takeda, T. Hasegawa, T. Nakamura, Y. Hosokoshi, K. Inoue, S. Ikeuchi, Y. Miyazaki and K. Saito, *J. Am. Chem. Soc.*, 2005, **127**, 4397.
- T. Ebata and M. Fujii, *Physical chemistry of cold gas-phase functional molecules and clusters*, Springer, Singapore, 2019.
- S. C. Henderson, S. J. Valentine, A. E. Counterman and D. E. Clemmer, *Anal. Chem.*, 1999, **71**, 291.
- T. Wyttenbach, G. V. Helden, J. J. Batka, D. Carlat and M. T. Bowers, *J. Am. Soc. Mass Spectrom.*, 1997, **8**, 275.
- T. Wyttenbach, G. von Helden and M. T. Bowers, *Int. J. Mass Spectrom. Ion Processes*, 1997, **165–166**, 377.
- S. Lee, T. Wyttenbach, G. von Helden and M. T. Bowers, *J. Am. Chem. Soc.*, 1995, **117**, 10159.
- J. N. Bull, M. S. Scholz, E. Carrascosa, G. da Silva and E. J. Bieske, *Phys. Rev. Lett.*, 2018, **120**, 223002.
- J. Seo, S. Warnke, K. Pagel, M. T. Bowers and G. von Helden, *Nat. Chem.*, 2017, **9**, 1263.
- S. Warnke, A. Ben Faleh and T. R. Rizzo, *ACS Meas. Sci. Au*, 2021, **1**, 157.
- W. Zagorec-Marks, L. G. Dodson, P. Weis, E. K. Schneider, M. M. Kappes and J. M. Weber, *J. Am. Chem. Soc.*, 2021, **143**, 17778.
- H. Milloy and M. Elford, *Int. J. Mass Spectrom. Ion Phys.*, 1975, **18**, 21.
- R. Johnsen, A. Chen and M. A. Biondi, *J. Chem. Phys.*, 2008, **72**, 3085.
- T. Koizumi, N. Kobayashi and Y. Kaneko, *J. Phys. Soc. Jpn.*, 1980, **48**, 1678.
- F. Misaizu, N. Hori, H. Tanaka, K. Komatsu, A. Furuya and K. Ohno, *Eur. Phys. J. D*, 2009, **52**, 59.
- T. Wyttenbach, P. R. Kemper and M. T. Bowers, *Int. J. Mass Spectrom.*, 2001, **212**, 13.
- J. C. May and D. H. Russell, *J. Am. Soc. Mass Spectrom.*, 2011, **22**, 1134.
- K. A. Servage, J. A. Silveira, K. L. Fort and D. H. Russell, *Acc. Chem. Res.*, 2016, **49**, 1421.
- J. Ujma, K. Giles, M. Morris and P. E. Barran, *Anal. Chem.*, 2016, **88**, 9469.
- G. von Helden, T. Wyttenbach and M. T. Bowers, *Int. J. Mass Spectrom. Ion Processes*, 1995, **146–147**, 349.
- J. Gidden and M. Bowers, *Eur. Phys. J. D*, 2002, **20**, 409.
- P. Weis, T. Bierweiler, E. Vollmer and M. M. Kappes, *J. Chem. Phys.*, 2002, **117**, 9293.
- S. Tainaka, T. Ujihira, M. Kubo, M. Kida, D. Shimoyama, S. Muramatsu, M. Abe, T. Haino, T. Ebata, F. Misaizu, K. Ohshimo and Y. Inokuchi, *J. Phys. Chem. A*, 2020, **124**, 9980.
- K. Ohshimo, X. He, R. Ito and F. Misaizu, *J. Phys. Chem. A*, 2021, **125**, 3718.
- R. Ito, K. Ohshimo and F. Misaizu, *Chem. Phys. Lett.*, 2022, **794**, 139510.
- R. Ito, X. He, K. Ohshimo and F. Misaizu, *J. Phys. Chem. A*, 2022, **126**, 4359.
- K. Ohshimo, X. He, R. Ito, K. Tsunoda, S. Tainaka and F. Misaizu, *EPJ Tech. Instrum.*, 2023, **10**, 11.
- K. Ohshimo, R. Sato and F. Misaizu, *J. Phys. Chem. A*, 2020, **124**, 7999.
- K. Ohshimo, R. Sato, Y. Takasaki, K. Tsunoda, R. Ito, M. Kanno and F. Misaizu, *J. Phys. Chem. Lett.*, 2023, **14**, 8281.
- D. Bright and M. R. Truter, *J. Chem. Soc. B*, 1970, 1550.
- M. A. Bush and M. R. Truter, *J. Chem. Soc., Perkin Trans. 2*, 1972, 345.
- M. A. Bush and M. R. Truter, *J. Chem. Soc. D*, 1970, 1439.
- M. A. Bush and M. R. Truter, *J. Chem. Soc. B*, 1971, 1440.
- D. Live and S. I. Chan, *J. Am. Chem. Soc.*, 1976, **98**, 3769.
- M. Shamsipur and A. I. Popov, *J. Am. Chem. Soc.*, 1979, **101**, 4051.
- C. M. Choi, H. J. Kim, J. H. Lee, W. J. Shin, T. O. Yoon, N. J. Kim and J. Heo, *J. Phys. Chem. A*, 2009, **113**, 8343.
- C. M. Choi, J. H. Lee, Y. H. Choi, H. J. Kim, N. J. Kim and J. Heo, *J. Phys. Chem. A*, 2010, **114**, 11167.
- Y. Inokuchi, O. V. Boyarkin, R. Kusaka, T. Haino, T. Ebata and T. R. Rizzo, *J. Am. Chem. Soc.*, 2011, **133**, 12256.
- R. Kusaka, S. Kokubu, Y. Inokuchi, T. Haino and T. Ebata, *Phys. Chem. Chem. Phys.*, 2011, **13**, 6827.
- Y. Inokuchi, T. Ebata and T. R. Rizzo, *J. Phys. Chem. A*, 2018, **122**, 3754.



- 44 M. Kida, M. Kubo, T. Ujihira, T. Ebata, M. Abe and Y. Inokuchi, *ChemPhysChem*, 2018, **19**, 1331.
- 45 Y. Kitamura, S. Muramatsu, M. Kida, T. Ebata and Y. Inokuchi, *J. Phys. Chem. A*, 2019, **123**, 9185.
- 46 M. Kida, T. Ujihira, M. Kubo, S. Muramatsu, T. Ebata and Y. Inokuchi, *J. Phys. Chem. A*, 2023, **127**, 3210.
- 47 H. E. Revercomb and E. A. Mason, *Anal. Chem.*, 1975, **47**, 970.
- 48 V. Gabelica, A. A. Shvartsburg, C. Afonso, P. Barran, J. L. Benesch, C. Bleiholder, M. T. Bowers, A. Bilbao, M. F. Bush, J. L. Campbell, I. D. Campuzano, T. Causon, B. H. Clowers, C. S. Creaser, E. De Pauw, J. Far, F. Fernandez-Lima, J. C. Fjeldsted, K. Giles, M. Groessl, C. J. Hogan Jr, S. Hann, H. I. Kim, R. T. Kurulugama, J. C. May, J. A. McLean, K. Pagel, K. Richardson, M. E. Ridgeway, F. Rosu, F. Sobott, K. Thalassinou, S. J. Valentine and T. Wyttenbach, *Mass Spectrom. Rev.*, 2019, **38**, 291.
- 49 H. Goto and E. Osawa, *J. Am. Chem. Soc.*, 1989, **111**, 8950.
- 50 H. Gotō and E. Ōsawa, *J. Chem. Soc., Perkin Trans. 2*, 1993, 187.
- 51 B. P. Pritchard, D. Altarawy, B. Didier, T. D. Gibson and T. L. Windus, *J. Chem. Inf. Model.*, 2019, **59**, 4814.
- 52 L. von Szentpály, P. Fuentealba, H. Preuss and H. Stoll, *Chem. Phys. Lett.*, 1982, **93**, 555.
- 53 P. Fuentealba, H. Stoll, L. von Szentpály, P. Schwerdtfeger and H. Preuss, *J. Phys. B*, 1983, **16**, L323.
- 54 M. J. Frisch, G. W. Trucks, H. B. Schlegel, G. E. Scuseria, M. A. Robb, J. R. Cheeseman, G. Scalmani, V. Barone, G. A. Petersson, H. Nakatsuji, X. Li, M. Caricato, A. V. Marenich, J. Bloino, B. G. Janesko, R. Gomperts, B. Mennucci, H. P. Hratchian, J. V. Ortiz, A. F. Izmaylov, J. L. Sonnenberg, D. Williams-Young, F. Ding, F. Lipparini, F. Egidi, J. Goings, B. Peng, A. Petrone, T. Henderson, D. Ranasinghe, V. G. Zakrzewski, J. Gao, N. Rega, G. Zheng, W. Liang, M. Hada, M. Ehara, K. Toyota, R. Fukuda, J. Hasegawa, M. Ishida, T. Nakajima, Y. Honda, O. Kitao, H. Nakai, T. Vreven, K. Throssell, J. A. Montgomery, Jr., J. E. Peralta, F. Ogliaro, M. J. Bearpark, J. J. Heyd, E. N. Brothers, K. N. Kudin, V. N. Staroverov, T. A. Keith, R. Kobayashi, J. Normand, K. Raghavachari, A. P. Rendell, J. C. Burant, S. S. Iyengar, J. Tomasi, M. Cossi, J. M. Millam, M. Klene, C. Adamo, R. Cammi, J. W. Ochterski, R. L. Martin, K. Morokuma, O. Farkas, J. B. Foresman and D. J. Fox, *Gaussian 16 Revision C.01*, Gaussian Inc., Wallingford CT, 2016.
- 55 M. F. Mesleh, J. M. Hunter, A. A. Shvartsburg, G. C. Schatz and M. F. Jarrold, *J. Phys. Chem.*, 1996, **100**, 16082.
- 56 I. Campuzano, M. F. Bush, C. V. Robinson, C. Beaumont, K. Richardson, H. Kim and H. I. Kim, *Anal. Chem.*, 2012, **84**, 1026.
- 57 E. R. Johnson, S. Keinan, P. Mori-Sánchez, J. Contreras-García, A. J. Cohen and W. Yang, *J. Am. Chem. Soc.*, 2010, **132**, 6498.
- 58 T. Lu and F. Chen, *J. Comput. Chem.*, 2012, **33**, 580.
- 59 R. Shannon, *Acta Crystallogr., Sect. A: Cryst. Phys., Diffraction, Theor. Gen. Crystallogr.*, 1976, **32**, 751.
- 60 N. K. Lee, S. Park and S. K. Kim, *J. Chem. Phys.*, 2002, **116**, 7910.

