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## Solvent-induced dual nucleophiles and the $\alpha$ -effect in the $S_N2$ versus E2 competition<sup>†</sup>

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We have quantum chemically investigated how microsolvation affects the various E2 and  $S_N2$  pathways, their mutual competition, and the  $\alpha$ -effect of the model reaction system  $HOO^-(H_2O)_n + CH_3CH_2Cl$ , at the CCSD(T) level. Interestingly, we identify the dual nature of the  $\alpha$ -nucleophile  $HOO^-$  which, upon solvation, is in equilibrium with  $HO^-$ . This solvent-induced dual appearance gives rise to a rich network of competing reaction channels. Among both nucleophiles,  $S_N2$  is always favored over E2, and this preference increases upon increasing microsolvation. Furthermore, we found a pronounced  $\alpha$ -effect, not only for  $S_N2$  substitution but also for E2 elimination, *i.e.*,  $HOO^-$  is more reactive than  $HO^-$  in both cases. Our activation strain and quantitative molecular orbital analyses reveal the physical mechanisms behind the various computed trends. In particular, we demonstrate that two recently proposed criteria, required for solvent-free nucleophiles to display the  $\alpha$ -effect, must also be satisfied by microsolvated  $HOO^-(H_2O)_n$  nucleophiles.

## Introduction

Bimolecular nucleophilic substitution ( $S_N2$ ) and elimination (E2) are ubiquitous, often mutually competing, reactions. For nearly 50 years, the mechanism of the  $S_N2$  reaction has been extensively studied experimentally and theoretically, including its temperature dependence,<sup>1–3</sup> benchmark studies,<sup>4–11</sup> steric effects,<sup>12–21</sup> and solvent effects.<sup>14,15,22–29</sup> Solvation can have a tremendous effect on chemical reactions.<sup>14,15,22,30–32</sup> Rate constants, for example, can be reduced by up to 16 orders of magnitude from the gas phase to aqueous solution.<sup>33</sup> Microsolvation bridges the gap between the gas phase and bulk solution and constitutes a powerful approach for obtaining a detailed understanding of how exactly solvent molecules affect reactions.<sup>33–39</sup> Most research in this direction has been focused on microsolvated  $S_N2$  reactions of  $Y^-(H_2O)_n + CH_3X$ .<sup>1,23,26,30,32,34,40–50</sup> Significantly less attention has been paid to microsolvated E2 reactions<sup>51–53</sup> or solvent effects on the  $S_N2$  versus E2 competition.<sup>25,27,37,54–56</sup>

Previous theoretical studies have shown that the stepwise introduction of solvent molecules (*e.g.*, HF,<sup>25</sup>  $CH_3OH$ <sup>27</sup> and  $H_2O$ <sup>37,55</sup>) as well as the increase in their solvation power favors  $S_N2$  relative to E2 pathways, and can lead to a switch from overall E2 to  $S_N2$  reactivity. The mechanism behind this solvation favoring the  $S_N2$  pathway is that  $S_N2$  reactions have a significantly lower characteristic distortivity and thus lower activation strain than E2 reactions and therefore suffer less from the reduced interaction between the nucleophile and the substrate when the nucleophile's basicity is attenuated upon solvation, as shown by Bickelhaupt, Hamlin *et al.*<sup>55</sup>

The peroxide anion  $HOO^-$  is an interesting nucleophile, because the monohydrated peroxide anion displays dual nucleophile character, where both  $HOO^-(H_2O)$  and  $(HOOH)(OH^-)$  species are similarly stable.<sup>57</sup> Moreover, as  $HOO^-(H_2O)$  reacts with  $CH_3Cl$ ,<sup>58</sup> both  $HOO^-$  and  $HO^-$  anions are possible attacking nucleophiles, and this was observed in our recent direct dynamic simulation work.<sup>36</sup> The introduction of water into the  $HOO^-$  nucleophile enriches the reaction dynamics for it adds the proton-transfer induced  $HO^-S_N2$  pathway; one can expect that if the substrate was ethyl halides, the E2 pathway will emerge and make the dynamics more complicated and interesting.

Besides,  $HOO^-$  is a typical  $\alpha$ -nucleophile, possessing a lone pair of electrons adjacent to the nucleophilic atom. The term  $\alpha$ -effect<sup>59</sup> has been used to describe the enhanced reactivity of  $\alpha$ -nucleophiles compared to that of normal nucleophiles by deviating from the Brønsted-type correlations found for normal nucleophiles.<sup>60</sup> There has been extensive discussion on the origin of the  $\alpha$ -effect, as well as a controversy about whether the  $\alpha$ -effect is controlled by the intrinsic properties of the

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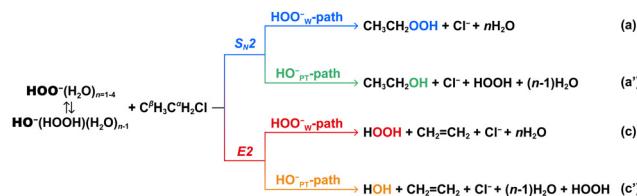
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**Scheme 1**  $S_N2$  and  $E2$  reaction pathways of  $HOO^-(H_2O)_{n=1-4} + CH_3CH_2Cl$  reactions. In pathway  $c'$ , one water is the conjugate acid of the nucleophile  $HO^-$ .

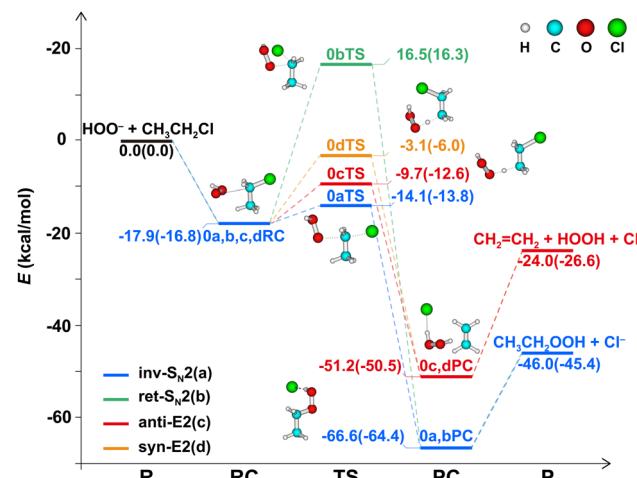
$\alpha$ -nucleophile or by external solvent effects.<sup>61–63</sup> In terms of the intrinsic properties, mechanisms such as ground state destabilization, transition state stabilization, and thermodynamic product stability were proposed to be the origin of the  $\alpha$ -effect.<sup>63–69</sup> The  $\alpha$ -effect has been observed in a variety of  $S_N2$  reactions,<sup>58,61,70–74</sup> yet fewer studies have addressed its relevance to  $E2$  reactions.<sup>75,76</sup> A recent study by Hamlin *et al.*<sup>77</sup> proposed two intrinsic criteria for the  $\alpha$ -nucleophile to display the  $\alpha$ -effect: (1) a higher energy HOMO and (2) a smaller HOMO lobe and overall amplitude of occupied orbitals on the nucleophilic center compared to the normal nucleophile. These criteria were proposed for solvent-free nucleophiles, and it is intriguing to examine whether they suit microsolvated nucleophiles.

In this study, we report a quantum chemical study on the  $HOO^-(H_2O)_n + CH_3CH_2Cl$  reaction (Scheme 1), where  $n = 0$  to 4 is the number of water molecules. We explore the full reaction pathways that, after formation of the initial  $E2$  or  $S_N2$  product complexes, lead to the separated products, as shown in Scheme 1. This is in agreement with earlier experimental and simulation studies on closely related microsolvated ion–molecule  $S_N2$  reactions, which have shown that the formation of the unsolvated ionic products strongly dominates the formation of the solvated ionic products because of dynamic bottlenecks which make solvent transfer from nucleophiles to leaving groups less likely.<sup>45,47,78</sup> The purpose of this study is three-fold, namely, to investigate the effect of solvation on (1) the competing  $S_N2$  and  $E2$  reaction pathways, (2) the competing normal  $HOO^-$ -pathways and the solvent-induced  $HO^-$ -pathways, and (3) the  $\alpha$ -effect on both, the  $S_N2$  and  $E2$  reactions, in terms of the nucleophiles' intrinsic properties.

## Results and discussion

### Potential energy surfaces

Fig. 1 depicts the potential energy surfaces (PESs) of the various  $S_N2$  and  $E2$  mechanistic pathways of  $HOO^-$  reacting with  $CH_3CH_2Cl$ , all of which display the typical “double-well” shape in which reactants (R) combine to a reactant complex (RC) that is connected *via* a transition state (TS) to product complexes (PC) and eventually separate products (P). The  $S_N2$  reaction is significantly more exothermic than the  $E2$  reaction, with respective reaction energies of  $-46.0$  and  $-24.0$   $kcal\ mol^{-1}$ . Depending on the site and spatial direction under which  $HOO^-$  attacks, the  $S_N2$  reaction proceeds either *via* backside



**Fig. 1** Potential energy profiles of the  $S_N2$  and  $E2$  reactions of  $HOO^- + CH_3CH_2Cl$  using the CCSD(T)/aug-cc-pVTZ//MP2/6-311++G(d,p) method. Energies in the normal text and parentheses are the electronic energy and enthalpy at 298.15 K, respectively.

substitution (inv- $S_N2$ , a) or *via* front side substitution that retains the stereostructure (ret- $S_N2$ , b),<sup>79</sup> whereas the  $E2$  reaction may take place either *via* anti-elimination (c) or *via* syn-elimination (d).<sup>80</sup> For the inv- $S_N2$  pathway, the nucleophile  $HOO^-$  attacks  $C^\alpha$  from the back side of the leaving group and causes the  $CH_3$ -moiety to undergo Walden inversion toward the product  $CH_3CH_2OOH$ . For the ret- $S_N2$  pathway,  $HOO^-$  attacks  $C^\alpha$  from the front side of the leaving group and the  $CH_3$ -moiety retains its geometry in product  $CH_3CH_2OOH$ . For anti- $E2$  or syn- $E2$  pathways,  $HOO^-$  attacks a  $\beta$  proton from the opposite or same side of the leaving group, respectively, leading to the abstraction of this proton and the formation of the  $E2$  products  $CH_2 = CH_2 + Cl^- + HOOH$ . The energy barriers decrease along ret- $S_N2$  ( $16.5$   $kcal\ mol^{-1}$ )  $>$  syn- $E2$  ( $-3.1$ )  $>$  anti- $E2$  ( $-9.7$ )  $>$  inv- $S_N2$  ( $-14.1$ ). The enthalpy and free energy values show the same trend (Table 1). The inv- $S_N2$  and anti- $E2$  pathways have the lowest barriers and therefore outperform other reaction channels. Therefore, we now focus on these two reaction mechanisms and label them, for simplicity, as  $S_N2$  and  $E2$ .

It has been shown that hydrated peroxide anions,  $HOO^-(H_2O)_n$ , tend to abstract a proton from  $H_2O$  and form more stable species  $HO^-(HOO)(H_2O)_{n-1}$  (Fig. S1, ESI<sup>†</sup>).<sup>36,38,57,58</sup> Consequently, the solvent-induced proton-transfer  $HO^-$ -moiety is a potential nucleophile to compete with the original nucleophile  $HOO^-$ . Hence, when  $HOO^-(H_2O)_n$  reacts with  $CH_3CH_2Cl$ , four pathways are possible: (a)  $HOO^-wS_N2$ , “W” indicates that the nucleophile is bound with water molecules; (a')  $HO^-PTS_N2$ , “PT” indicates that the  $HO^-$  nucleophile is induced by proton transfer from water molecules; (c)  $HOO^-wE2$ ; and (c')  $HO^-PTE2$  (Scheme 1). Note that in the hydrated system, the  $HOO^-wE2$  and  $HO^-PTE2$  pathways generate the same products.

To show the effect of individual solvent molecules, we plotted the potential energy profiles of the  $HOO^-(H_2O)_{n=0,1,2} + CH_3CH_2Cl$  reactions in Fig. 2 for both  $HOO^-w$ -paths (right panel) and  $HO^-PT$ -paths (left panel). We used 0, 1 and 2 as



**Table 1** Calculated overall barriers (in kcal mol<sup>-1</sup>) for the S<sub>N</sub>2 and E2 paths of HOO<sup>-</sup>(H<sub>2</sub>O)<sub>n</sub> + CH<sub>3</sub>CH<sub>2</sub>Cl reactions with the electronic energy (E), enthalpy (H, 298.15 K) and Gibbs free energy (G, 298.15 K) values reported<sup>a</sup>

	n	HOO <sup>-</sup> <sub>w</sub> -path			HOO <sup>-</sup> <sub>PT</sub> -path		
		ΔE <sup>‡</sup>	ΔH <sup>‡</sup>	ΔG <sup>‡</sup>	ΔE <sup>‡</sup>	ΔH <sup>‡</sup>	ΔG <sup>‡</sup>
inv-S <sub>N</sub> 2	0	-14.1	-13.8	-3.9	-14.0	-14.0	-5.4
	1	-4.6	-3.8	5.2	-1.4	-0.4	9.5
	2	0.6	1.7	13.2	2.7	3.8	16.3
	3	3.1	4.1	14.8	8.0	9.3	19.4
	4	8.0	8.3	18.2	10.0	11.0	22.4
anti-E2	0	-9.7	-12.6	-3.5	-12.7	-15.8	-8.0
	1	2.4	0.1	7.9	3.8	1.6	9.6
	2	7.9	6.9	19.2	8.8	6.9	17.4
	3	14.0	12.5	21.6	15.2	13.7	22.0
	4	17.9	16.4	25.1	18.2	16.2	26.9
ret-S <sub>N</sub> 2	0	16.5	16.3	27.5	20.9	20.5	29.4
syn-E2	0	-3.1	-6.0	4.1	-5.2	-8.6	-0.5

<sup>a</sup> Computed at CCSD(T)/aug-cc-pVTZ//MP2/6-311++G(d,p).

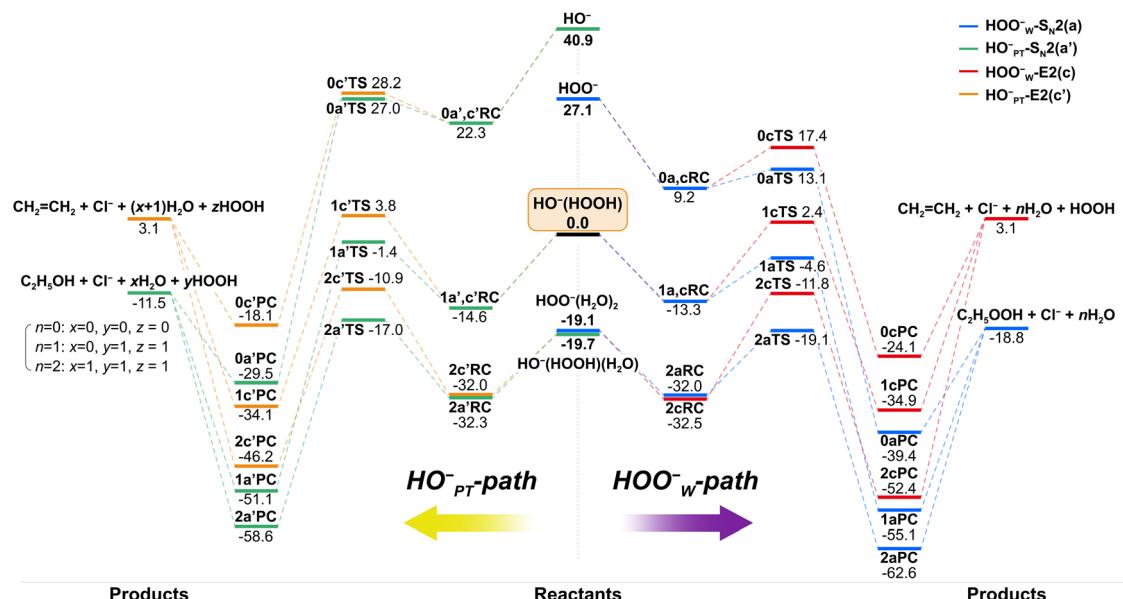
prefixes to denote the number of solvent molecules when naming the species. The corresponding transition state structures are shown in Fig. 3. The involvement of multiple H<sub>2</sub>O molecules complicates the structures, so the most stable structures of each species and corresponding energetics were used in the discussion. Information on higher-energy conformational isomers is provided in Fig. S2–S6 (ESI†) for interested readers. Using HO<sup>-</sup>(HOOH) + CH<sub>3</sub>CH<sub>2</sub>Cl as the reference point, as observed, the HOO<sup>-</sup><sub>w</sub>-S<sub>N</sub>2 reactions (−18.8 kcal mol<sup>-1</sup>) are more exothermic than HOO<sup>-</sup><sub>PT</sub>-S<sub>N</sub>2 reactions (−11.5 kcal mol<sup>-1</sup>), where both are more exothermic than E2 reactions (3.1 kcal mol<sup>-1</sup>). The addition of water molecules to the ion–molecule system stabilizes each species. For the singly- and doublyhydrated systems, the potential energy profiles of both S<sub>N</sub>2

and E2 reactions remain double-well shaped. However, due to the differential stabilization effect of water molecules on the reactants and transition states, the barrier heights changed differently for the S<sub>N</sub>2 and E2 reactions. The energetic values are presented in Table 1 and details will be discussed in the next section.

### Competition S<sub>N</sub>2 *versus* E2 and HOO<sup>-</sup> *versus* HO<sup>-</sup>

The reaction barrier,  $\Delta E^{\ddagger} = E(\text{TS}) - E(\text{R})$ , of all four paths increases systematically as the extent of microsolvation (*i.e.*, the number of solvent molecules) increases, as shown in Fig. 4. This trend of an increasing barrier agrees well with the decreasing reaction rate observed previously in experimental studies of microsolvated ion–molecule reactions.<sup>1,30,33,34,40,52,54,81,82</sup> Focusing on the competition between S<sub>N</sub>2 and E2 pathways (Fig. 4a and b), the barrier of E2 pathways increases slightly faster upon microsolvation than that of S<sub>N</sub>2 pathways. Thus, the extent to which the E2 barrier exceeds that of the S<sub>N</sub>2 barrier, *i.e.*  $\Delta \Delta E_1^{\ddagger} = \Delta E^{\ddagger}(\text{E2}) - \Delta E^{\ddagger}(\text{S}_N\text{2})$ , increases upon microsolvation. For instance, as *n* increases from 0 to 4, the  $\Delta \Delta E_1^{\ddagger}$  value increases from 4.3 to 7.0, 7.4, 10.8, and 9.9 kcal mol<sup>-1</sup> for the HOO<sup>-</sup><sub>w</sub>-path. Consequently, the S<sub>N</sub>2 paths always dominate and they do so even more at higher degrees of solvation. This enlargement of the E2-S<sub>N</sub>2 barrier difference ( $\Delta \Delta E_1^{\ddagger}$ ) with a stepwise increase of microsolvation is also computed for the HOO<sup>-</sup><sub>PT</sub>-pathways, with corresponding values of 1.3, 5.2, 6.2, 7.2 and 8.3 kcal mol<sup>-1</sup>, respectively. Our finding consolidates earlier work that focused on the E2 and S<sub>N</sub>2 reactions of microsolvated model systems HO<sup>-</sup>(H<sub>2</sub>O)<sub>n</sub> + CH<sub>3</sub>CH<sub>2</sub>X (X = Cl, Br, I) and F<sup>-</sup> + CH<sub>3</sub>CH<sub>2</sub>F + nHF.<sup>25,37,80</sup>

Regarding the competition between the  $\alpha$ -nucleophile HOO<sup>-</sup><sub>w</sub> and the normal nucleophile HO<sup>-</sup><sub>PT</sub> (Fig. 4c and d), the barrier height difference, defined as  $\Delta \Delta E_2^{\ddagger} = \Delta E^{\ddagger}(\text{HO}^-_{\text{PT}}) -$



**Fig. 2** Potential energy profiles of HOO<sup>-</sup>(H<sub>2</sub>O)<sub>n=0-2</sub> + CH<sub>3</sub>CH<sub>2</sub>Cl reactions for both, the HOO<sup>-</sup><sub>w</sub>-path (right) and the HOO<sup>-</sup><sub>PT</sub>-path (left). Energies (kcal mol<sup>-1</sup>) are relative to HO<sup>-</sup>(HOOH) + CH<sub>3</sub>CH<sub>2</sub>Cl at the level of CCSD(T)/aug-cc-pVTZ//MP2/6-311++G(d,p) without the ZPE correction.



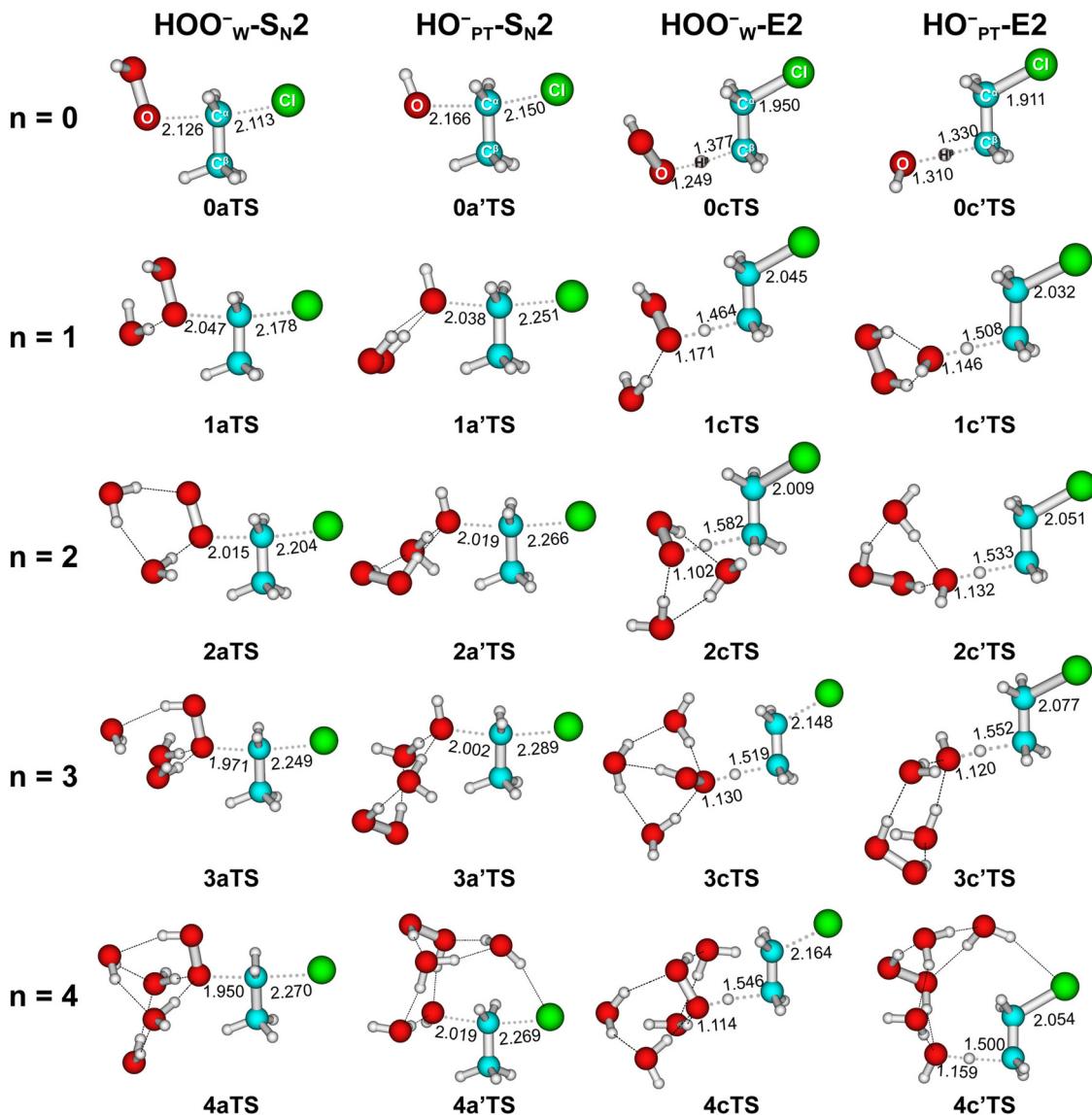


Fig. 3 Structures of transition states of  $\text{HOO}^-(\text{H}_2\text{O})_{n=0-4} + \text{CH}_3\text{CH}_2\text{Cl}$  optimized at MP2/6-311++G(d,p). Bond lengths are in angstrom.

$\Delta E^\ddagger(\text{HOO}^-\text{w})$ , is greater than zero for almost all cases, that is, the  $\alpha$ -nucleophile reacts faster in almost all cases. The main trend is that introducing microsolvation, *i.e.*, going from  $n = 0$  to  $n \geq 1$ , enhances the  $\alpha$ -effect, *i.e.*,  $\Delta\Delta E_2^\ddagger$ . However, the dependence of the  $\alpha$ -effect as a function of introducing more solvent molecules (*i.e.*, along  $n = 1, 2, 3, 4$ ) is less uniform. For the  $\text{S}_{\text{N}}2$  paths, the barrier difference  $\Delta\Delta E_2^\ddagger(\text{S}_{\text{N}}2)$  ranges from 2.0 to 4.9 kcal mol<sup>-1</sup>, whereas there is a smaller difference between the  $\text{E}2$  pathways (Fig. 4d), and  $\Delta\Delta E_2^\ddagger(\text{E}2)$  ranges from 0.4 to 1.4 kcal mol<sup>-1</sup>, indicating that the  $\text{HO}^-\text{-E}2$  path can be strongly competitive to the  $\text{HOO}^-\text{-E}2$  pathway, provided sufficient energy is available to pass the  $\text{E}2$  barriers.

In brief, among the four competing pathways, the  $\text{HOO}^-\text{w-S}_{\text{N}}2$  path dominates with incremental solvation. In what follows, we seek the reason for barrier height difference upon solvation by analyzing the properties of nucleophiles and transition states.

**Nucleophiles and HOMO-LUMO interactions.** Besides the hydrated peroxide anion nucleophiles  $\text{HOO}^-(\text{H}_2\text{O})_{n=1-4}$  and the associated PT-induced  $\text{HO}^-(\text{HOOH})(\text{H}_2\text{O})_{n=1-3}$  nucleophiles, we also considered the hydrated hydroxide anion nucleophiles  $\text{HO}^-(\text{H}_2\text{O})_{0-4}$ <sup>37</sup> for comparison. The latter was labeled as the  $\text{HO}^-\text{w}$ -path. The properties calculated include the energy level of the HOMO, the proton affinity (PA), and the ethyl cation affinity (ECA) of these nucleophiles.

In line with our previous studies,<sup>25,38,55,80</sup> we found herein that microsolvation lowers the energy of the HOMO of both  $\text{HOO}^-$  and  $\text{HO}^-$  systematically upon adding an additional solvent molecule, either  $\text{H}_2\text{O}$  or  $\text{HOOH}$ , by the HOMO-LUMO interaction with the  $\sigma_{\text{O}-\text{H}^*}$  (solvent) LUMO (Table S4, ESI<sup>†</sup>). The HOMO of the microsolvated peroxide anion in  $\text{HOO}^-(\text{H}_2\text{O})_{2,3,4}$  (-5.6, -6.0, and -6.6 eV) always remains higher in energy than that of the equivalently microsolvated hydroxide anion in  $\text{HO}^-(\text{HOOH})(\text{H}_2\text{O})_{1,2,3}$  (-6.5, -6.9, and -7.3 eV) and

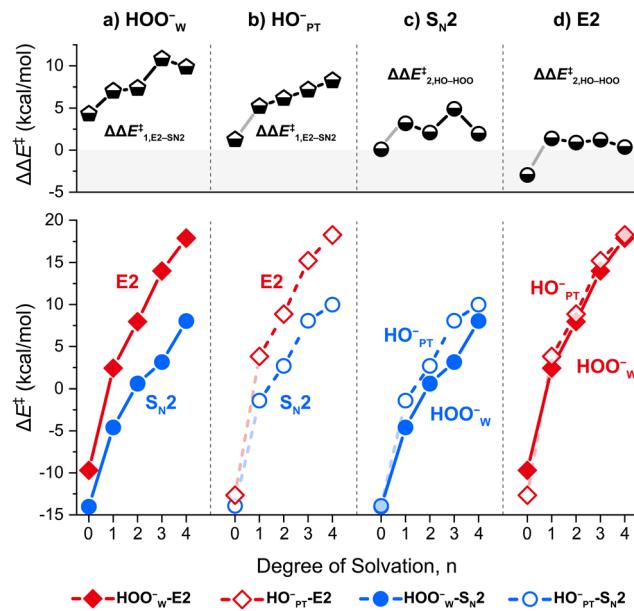


Fig. 4 Barrier heights  $\Delta E^\ddagger$  (bottom) and differences in barrier heights  $\Delta\Delta E^\ddagger$  (top) of the four mechanistic pathways of the model reaction system  $\text{HOO}^-(\text{H}_2\text{O})_{n=0-4} + \text{CH}_3\text{CH}_2\text{Cl}$ , computed at CCSD(T)/aug-cc-pVTZ//MP2/6-311+G(d,p) without the ZPE correction. Note: for the  $\text{HO}^-_{\text{PT}}$ -path,  $n = 0$  corresponds to the  $\text{HO}^- + \text{CH}_3\text{CH}_2\text{Cl}$  reaction.

$\text{HO}^-(\text{H}_2\text{O})_{2,3,4}$  clusters ( $-5.7$ ,  $-6.6$ , and  $-7.2$  eV; see Table S4, ESI<sup>†</sup>). As shown previously by Bickelhaupt *et al.*,<sup>25,55</sup> this situation gives rise to a smaller HOMO-LUMO energy gap and a more stabilizing HOMO-LUMO interaction with the LUMO (substrate) (in this work, the substrate is  $\text{CH}_3\text{CH}_2\text{Cl}$ ) and, therefore, to a lower barrier for the reactions of microsolvated  $\text{HOO}^-$  than for the corresponding reactions of microsolvated  $\text{HO}^-$ . Indeed, we found a strong correlation of the barrier heights of the  $\text{S}_{\text{N}2}$  and E2 reactions with the HOMO energy level of the microsolvated peroxide and hydroxide nucleophiles (Fig. 5 and Fig. S8, ESI<sup>†</sup>). The same holds true for the related quantities of proton affinity (PA) and ethyl cation affinity (ECA), which are measures of the nucleophiles' ability to bind with a proton and an ethyl cation. The PA and ECA values are defined as the enthalpy change of  $\text{NuH} \rightarrow \text{Nu}^- + \text{H}^+$  and  $\text{CH}_3\text{CH}_2\text{Nu} \rightarrow \text{Nu}^- + \text{CH}_3\text{CH}_2^+$ , respectively (see Table S5 for all computed PA and ECA values, ESI<sup>†</sup>).

With the introduction of solvent molecules, the PA or ECA values of nucleophiles decrease continuously, that is, both types of affinities become weaker. Gratifyingly, the overall barrier heights ( $\Delta E^\ddagger$ ) of the  $\text{S}_{\text{N}2}$  and E2 reactions have, again, a strong linear correlation with the PA or ECA of the various microsolvated peroxide and hydroxide nucleophiles ( $R^2 \approx 1.00$ ; see Fig. 5). In line with this, there is also a good to reasonable correlation between the height of the barrier and the amount of charge transferred from the nucleophile to the substrate, as computed with natural population analysis (NPA) atomic charges ( $R^2 = 0.85$ – $0.97$ ; see Fig. 5 and Table S7, ESI<sup>†</sup>).<sup>37,83</sup> Thus, a stronger charge transfer goes with a lower barrier.

**$\text{S}_{\text{N}2}$  versus E2 characteristic distortivity.** In this section, we analyze how the barrier heights correlate with the geometrical

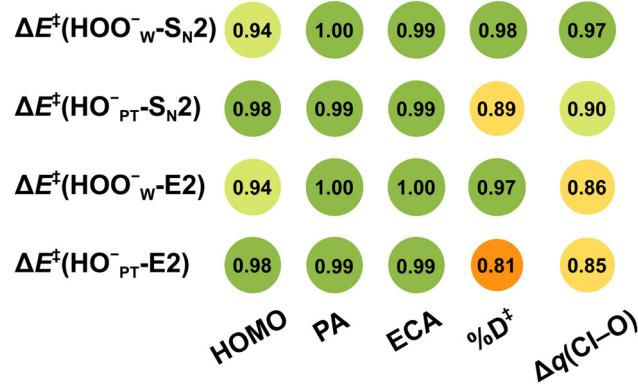


Fig. 5 Correlation coefficients ( $R^2$ ) for the linear correlation between the E2 and  $\text{S}_{\text{N}2}$  barrier heights and the physical properties of the  $\text{HOO}^-(\text{H}_2\text{O})_{n=0-4} + \text{CH}_3\text{CH}_2\text{Cl}$  reaction systems: nucleophile HOMO energy, PA and ECA, distortion parameter  $\%D^\ddagger$ , and charge asymmetry  $\Delta q(\text{Cl}-\text{O})$ .

characteristics of the transition states for the  $\text{S}_{\text{N}2}$  and E2 mechanistic pathways. As before, we focus on the pathway involving the geometrical configuration of solvent molecules that yields the most stable TS conformation for each reaction mechanism. In the first place, we recall that E2 elimination goes inherently with a higher characteristic distortivity than  $\text{S}_{\text{N}2}$  substitution which is a factor that contributes to a higher activation strain and thus a higher barrier. To put this on a more quantitative basis for the model reactions investigated herein, we have defined the distortion parameter  $\%D^\ddagger$ , which depends linearly on the stretch of  $\text{C}^\alpha\text{-Cl}$  and  $\text{H}^\beta\text{-C}^\beta$  bonds in the TS, as defined below:

$$\%D^\ddagger = \%C^\alpha Cl^\ddagger + \%H^\beta C^\beta \ddagger$$

$$\%C^\alpha Cl^\ddagger = 100 \times \left( \frac{r_{C^\alpha-Cl}^\ddagger - r_{C^\alpha-Cl}^{\text{Reactant}}}{r_{C^\alpha-Cl}^{\text{Reactant}}} \right) / r_{C^\alpha-Cl}^{\text{Reactant}}$$

$$\%H^\beta C^\beta \ddagger = 100 \times \left( \frac{r_{H^\beta-C^\beta}^\ddagger - r_{H^\beta-C^\beta}^{\text{Reactant}}}{r_{H^\beta-C^\beta}^{\text{Reactant}}} \right) / r_{H^\beta-C^\beta}^{\text{Reactant}}$$

Indeed, we found a reasonably strong linear correlation of the  $\text{S}_{\text{N}2}$  and E2 reaction barriers with  $\%D^\ddagger$  ( $R^2 = 0.81$  to  $0.98$ ; see Fig. 5). The trend is that the more geometrically distorted the TS, the higher the barrier.

Furthermore, we found that incremental microsolvation makes the transition structures in general more product-like, as elaborated upon in the following. The  $\text{S}_{\text{N}2}$  pathway and the transition state involve the breaking of the  $\text{C}^\alpha\text{-Cl}$  bond and the formation of the  $\text{O-C}^\alpha$  bond, whereas the E2 pathway and transition state involve the breaking of both the  $\text{C}^\alpha\text{-Cl}$  bond and the  $\text{H}^\beta\text{-C}^\beta$  bond and the formation of the peroxide or hydroxide  $\text{O-H}^\beta$  bond (Fig. 3). Table 2 shows that, as the degree of microsolvation increases from 0 to 3, for transition structures of inv- $\text{S}_{\text{N}2}$  reactions, the  $\text{O-C}^\alpha$  bond shortens and the  $\text{C}^\alpha\text{-Cl}$  bond lengthens systematically. In the case of the transition structures of *anti*-E2 reactions, there are a few irregularities but, by and large, the  $\text{O-H}^\beta$  bond shortens, and both the  $\text{C}^\alpha\text{-Cl}$  bond and the  $\text{H}^\beta\text{-C}^\beta$  bond lengthen upon going from the unsolvated



to the microsolvated situation. In line with these structural characteristics, the Cl leaving group becomes increasingly negatively charged in the transition states of both  $S_N2$  and E2 reactions when the degree of microsolvation increases, as reflected by the computed NPA charges (see Table S7, ESI<sup>†</sup>). For example, the negative charge of the leaving group  $q(\text{Cl})$  of the  $\text{HOO}^-_{\text{w}} S_N2$  path increases from  $-0.539$  to  $-0.675$  a.u., and the value of the  $\text{HOO}^-_{\text{w}}$  E2 path increases from  $-0.355$  to  $-0.580$  a.u. These findings are all consistent with the fact that the transition states become more product-like. Thus, as-stated above, microsolvation shifts the TS to a later, more product-like point along the reaction coordinate.

**Activation strain analysis.** The above analyses show that a higher distortion of the substrate is connected with a higher barrier (E2 higher than  $S_N2$ ) and also that a poorer donor-acceptor interaction capability of the nucleophile results in higher barriers (higher barriers upon adding solvent molecules). To gain more quantitative insight into this, we performed activation strain analyses that decompose the reaction energy barrier ( $\Delta E^\ddagger$ ) relative to separate reactants into the activation strain ( $\Delta E_{\text{strain}}$ ) and the TS interaction ( $\Delta E_{\text{int}}$ ), as shown below (for details, see the ESI<sup>†</sup>):<sup>80,84,85</sup>

$$\Delta E^\ddagger = \Delta E_{\text{strain}} + \Delta E_{\text{int}}$$

As shown in Fig. S9a and S9b (ESI<sup>†</sup>), the destabilizing strain energy of the E2 path is significantly larger than that of the  $S_N2$  path. The reason is the aforementioned larger characteristic distortion associated with the E2 path in which two bonds are breaking ( $\text{C}^\alpha\text{--Cl}$  and  $\text{C}^\beta\text{--H}$ ) in the substrate as compared to the lesser characteristic distortion associated with the  $S_N2$  path along which only one bond ( $\text{C}^\alpha\text{--Cl}$ ) is breaking in the substrate.<sup>86</sup> The higher activation strain is what makes the E2 barrier higher than the  $S_N2$  barrier, and this can only be inverted if the stabilizing nucleophile–substrate interaction is strong enough. As pointed out by Bickelhaupt *et al.*,<sup>86</sup> the E2 pathway goes with a higher TS acidity, *i.e.*, the lower LUMO in the TS, than the  $S_N2$  pathway. However, neither  $\text{HOO}^-$  nor  $\text{HO}^-$  are strong enough bases to cause an inversion of barrier heights

as determined by the unfavorably high activation strain for E2 reactions involving the  $\text{CH}_3\text{CH}_2\text{Cl}$  substrate. In view of the fact that introducing solvent molecules makes the nucleophile an even poorer electron donor, the E2 barrier becomes higher relative to the  $S_N2$  barrier as the degree of solvation increases.

### $\alpha$ -Effect

The  $\text{HOO}^-$  anion is an  $\alpha$ -nucleophile, featuring a lone-pair bearing heteroatom adjacent to the nucleophilic center. Its parent normal nucleophile is  $\text{HO}^-$ . The  $\alpha$ -effect refers to the dramatically enhanced reactivity of  $\alpha$ -nucleophiles compared to their parent normal nucleophiles by deviating downward from the Brønsted-type correlation (reaction barrier *versus* proton affinity) found for normal nucleophiles.<sup>87</sup> To evaluate whether microhydrated  $\text{HOO}^-$  anions display an  $\alpha$ -effect, two different Brønsted-type correlations between the barrier height and the basicity can be plotted by choosing different normal nucleophiles.

In the Brønsted-type I correlation, one plots reaction barriers ( $\Delta H^\ddagger$ ) against the proton affinity (PA) of  $\text{HOO}^-(\text{H}_2\text{O})_n$  and  $\text{HO}^-(\text{H}_2\text{O})_n$  for  $n = 0$  to 4. We found that there is a good correlation between  $\Delta H^\ddagger$  and the PA of the normal nucleophiles  $\text{HO}^-(\text{H}_2\text{O})_n$  for both  $S_N2$  (Fig. 6a) and E2 (Fig. 6b) reactions. Thus, nucleophiles with larger PA values have lower barriers and, therefore, a higher reactivity. All the points of  $\text{HOO}^-(\text{H}_2\text{O})_n$  deviate downward from this correlation line, with deviation values  $\Delta\Delta H^\ddagger = \Delta H^\ddagger(\text{HO}^- \text{ Brønsted path PA}) - \Delta H^\ddagger(\text{HOO}^-)$  of more than 5 kcal mol<sup>-1</sup>, revealing the existence of the  $\alpha$ -effect for microhydrated  $\text{HOO}^-$  anions. The barriers of E2 reactions are more sensitive to the PA, as reflected by the larger slope of the correlation lines.

In the Brønsted-type II correlation, at each degree of hydration  $n$ , one plots the  $\Delta H^\ddagger$  of a set of other reference normal nucleophiles, including  $\text{HO}^-$ ,  $\text{H}_2\text{N}^-$  and  $\text{HS}^-$ , against the PA, and compares this with the barrier for the  $\text{HOO}^-$  anion and its PA. As shown in Fig. 7, when  $n = 0$  to 3, the barriers of  $\text{HOO}^-(\text{H}_2\text{O})_n$  show a downward shift from the Brønsted-type II  $\Delta H^\ddagger$  *versus* PA correlation line of the normal nucleophiles for

**Table 2** Selected bond lengths (in Å) of the transition structures

HOO <sup>-</sup> <sub>w</sub> -path				HO <sup>-</sup> <sub>PF</sub> -path			
inv-S <sub>N2</sub> -TS				inv-S <sub>N2</sub> -TS			
<i>n</i>	<i>r</i> (O-C <sup>α</sup> )	<i>r</i> (C <sup>α</sup> -Cl)	<i>r</i> (H <sup>B</sup> -C <sup>B</sup> )	<i>r</i> (O-C <sup>α</sup> )	<i>r</i> (C <sup>α</sup> -Cl)	<i>r</i> (H <sup>B</sup> -C <sup>B</sup> )	
0	2.126	2.113	1.094	2.166	2.150	1.091	
1	2.047	2.178	1.094	2.038	2.251	1.091	
2	2.015	2.204	1.094	2.019	2.266	1.092	
3	1.971	2.249	1.094	2.002	2.289	1.092	
4	1.950	2.270	1.094	2.019	2.269	1.092	
anti-E2-TS				anti-E2-TS			
<i>n</i>	<i>r</i> (O-H <sup>B</sup> )	<i>r</i> (C <sup>α</sup> -Cl)	<i>r</i> (H <sup>B</sup> -C <sup>B</sup> )	<i>r</i> (C <sup>α</sup> -C <sup>B</sup> )	<i>r</i> (O-H <sup>B</sup> )	<i>r</i> (C <sup>α</sup> -Cl)	<i>r</i> (H <sup>B</sup> -C <sup>B</sup> )
0	1.249	1.950	1.377	1.458	1.310	1.911	1.330
1	1.171	2.045	1.464	1.433	1.146	2.032	1.508
2	1.102	2.009	1.582	1.442	1.132	2.051	1.533
3	1.130	2.148	1.519	1.412	1.120	2.077	1.552
4	1.114	2.164	1.546	1.411	1.159	2.054	1.500

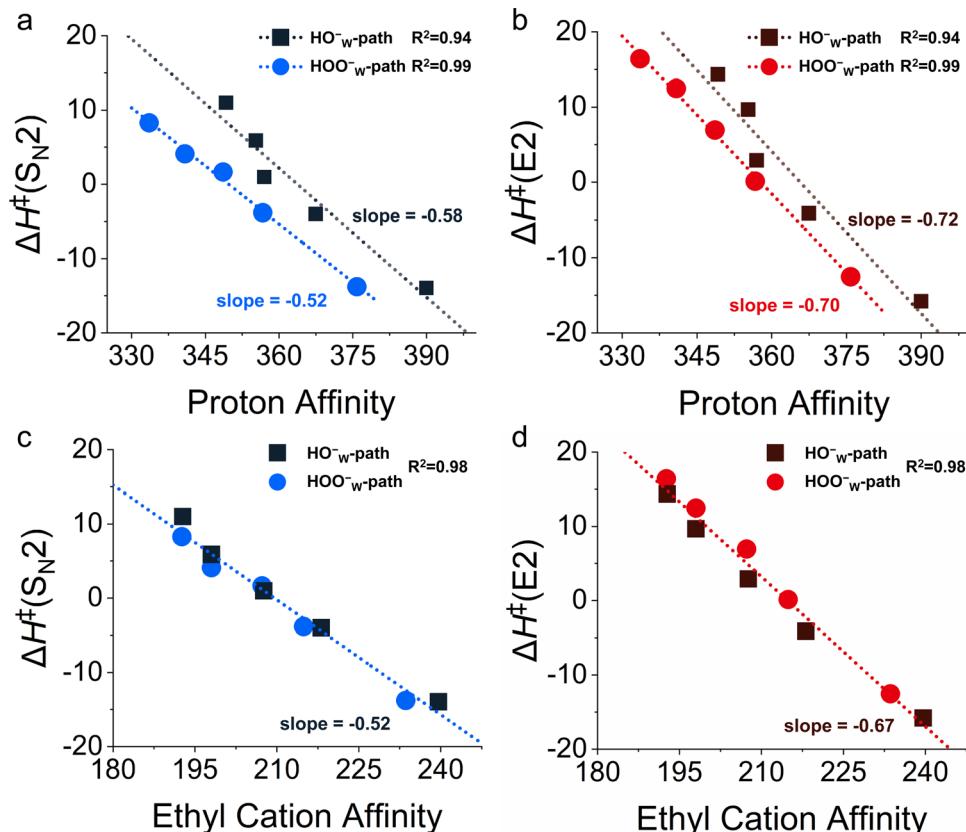


Fig. 6 The correlation between  $\text{HOO}^-(\text{H}_2\text{O})_n$  and  $\text{HO}^-(\text{H}_2\text{O})_n + \text{CH}_3\text{CH}_2\text{Cl}$  reaction barrier heights  $\Delta H^\ddagger$  and the (a) and (b) proton affinity or (c) and (d) ethyl cation affinity of nucleophiles. The barrier heights  $\Delta H^\ddagger$  (in kcal mol<sup>-1</sup>) have been computed at CCSD(T)/aug-cc-pVTZ//MP2/6-311+G(d,p). The PA and ECA of the nucleophiles have been computed at G3(MP2).

both  $\text{S}_{\text{N}}2$  and E2 reactions. This downward shift becomes less obvious for E2 when  $n = 1\text{--}3$ . This is consistent with the report of Hamlin *et al.*,<sup>77</sup> where the unsolvated nucleophiles were considered reacting with ethyl halides. To our knowledge, no such correlation has been plotted for microsolvated nucleophiles reacting with ethyl halides, nevertheless, examples with methyl halides exist.<sup>58,74,88-90</sup> Experimental studies by Bierbaum's group<sup>58,74</sup> and computational investigations by Ren's group<sup>88,90</sup> suggested that the  $\alpha$ -effect exists for  $\text{HOO}^-(\text{H}_2\text{O})$  and  $\text{HOO}^-(\text{CH}_3\text{OH})$  reacting with methyl chloride. Herein, we expand the exploration of this phenomenon to ethyl halides covering both  $\text{S}_{\text{N}}2$  and E2 reactions. In brief, the Brønsted-type II analyses also suggest that  $\text{HOO}^-(\text{H}_2\text{O})_n$  nucleophiles display the  $\alpha$ -effect as compared with their normal nucleophile counterparts.

The mechanism behind the  $\alpha$ -effect in  $\text{S}_{\text{N}}2$  reactions of unsolvated nucleophiles has been recently studied by Hamlin *et al.*<sup>77</sup> who identified two criteria an  $\alpha$ -nucleophile needs to fulfill in order to show the  $\alpha$ -effect: (1) a small HOMO lobe and overall reduction of occupied amplitude on the nucleophilic center, in order to reduce the repulsive occupied-occupied orbital overlap between the nucleophile and the substrate and (2) a sufficiently high HOMO energy level, in order to still engage in a strong HOMO<sub>nucleophile</sub>-LUMO<sub>substrate</sub> orbital interaction with the substrate. Herein, we examine whether the microsolvated  $\text{HOO}^-(\text{H}_2\text{O})_n$  nucleophiles satisfy the two

criteria. First, as shown in Fig. S10 (ESI<sup>†</sup>), the key occupied orbitals of nucleophiles, the HOMO lobes of  $\text{HOO}^-(\text{H}_2\text{O})_n$ , are smaller than those of  $\text{HO}^-(\text{H}_2\text{O})_n$  and  $\text{HO}^-(\text{HOOH})(\text{H}_2\text{O})_{n-1}$  for nearly all degrees of microsolvation. This is further confirmed by the fact that the total negative charge on the nucleophilic center is significantly lower in the  $\alpha$ -nucleophiles than in the corresponding normal nucleophile (see Fig. S1, ESI<sup>†</sup>). Second, if a DFT method is used, the HOMO levels of  $\text{HOO}^-(\text{H}_2\text{O})_n$  are consistently at higher energy than those of  $\text{HO}^-(\text{H}_2\text{O})_n$  and  $\text{HO}^-(\text{HOOH})(\text{H}_2\text{O})_{n-1}$ . So, both criteria<sup>77</sup> are indeed satisfied also in the case of microsolvated  $\alpha$ -nucleophiles  $\text{HOO}^-(\text{H}_2\text{O})_n$ . Accordingly, the microsolvated  $\text{HOO}^-(\text{H}_2\text{O})_n$  anion shall display an  $\alpha$ -effect, consistent with the result given by Brønsted-type correlations.

Our above analyses show that the rate-accelerating  $\alpha$ -effect in microsolvated E2 and microsolvated  $\text{S}_{\text{N}}2$  reactions goes hand-in-hand with a rise in orbital energy of the  $\pi$ -antibonding HOMO and the reduced amplitude of density on the nucleophilic center. The mechanism behind this is that the  $\alpha$ -nucleophile has, due to its reduced density on the nucleophilic center, less steric (Pauli) repulsion with the substrate than the normal nucleophile and, therefore, a more stabilizing overall interaction;<sup>77</sup> this difference in Pauli repulsion does not occur for the proton affinity because the proton has no occupied orbitals. However, if the ethyl cation affinity (ECA) instead



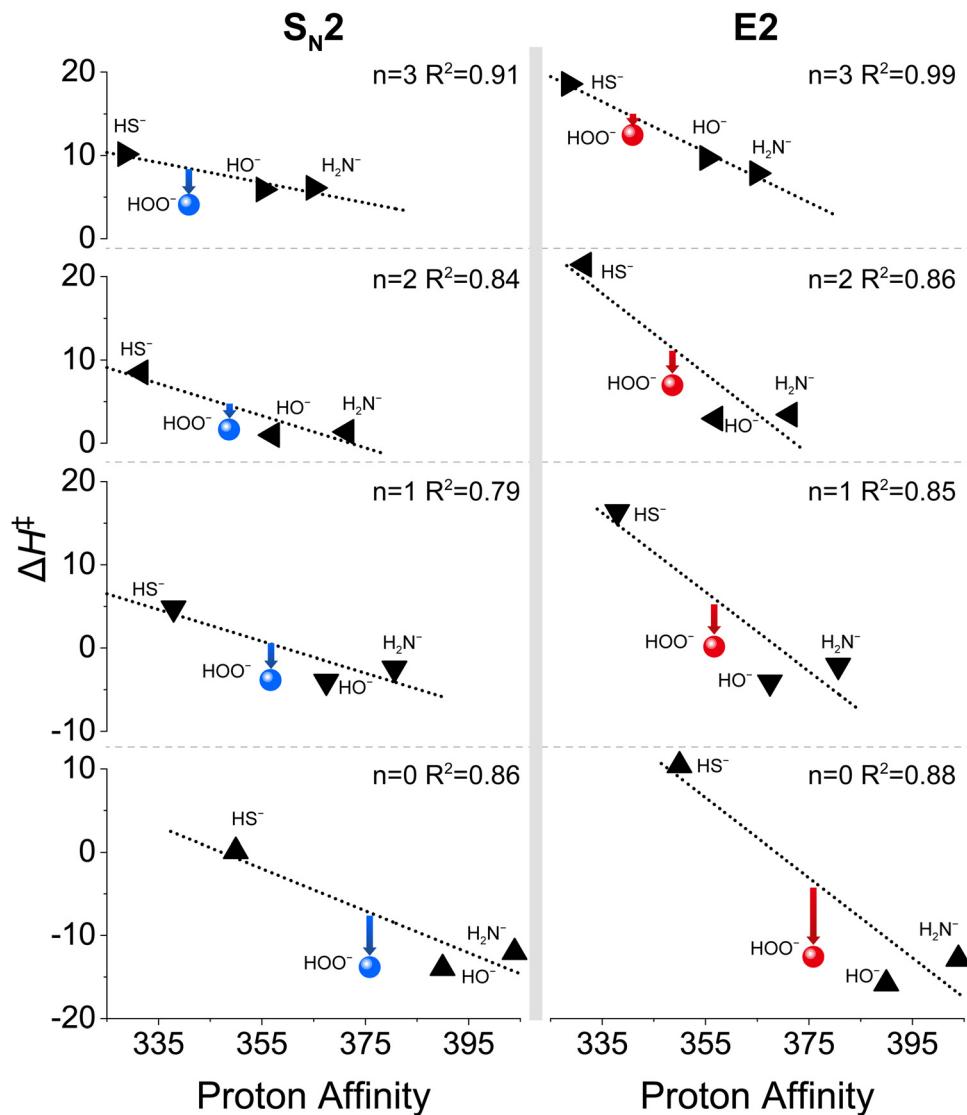


Fig. 7 The correlation between  $\text{HY}^-(\text{H}_2\text{O})_n$  ( $\text{Y} = \text{O, S, HN}$ ) +  $\text{CH}_3\text{CH}_2\text{Cl}$  reaction barrier heights  $\Delta H^\ddagger$  and the proton affinity of the nucleophiles. The barrier heights  $\Delta H^\ddagger$  (in  $\text{kcal mol}^{-1}$ ) are obtained with the CCSD(T)/aug-cc-pVTZ method.

of the PA is used as the Brønsted-correlation parameter, the  $\alpha$ -effect is diminished; that is, the downward deviation of the barriers for the  $\alpha$ -nucleophiles from the barrier *versus* the ECA Brønsted-type correlation is significantly reduced because the Pauli-reduction lowering in the case of  $\alpha$ -nucleophiles relative to normal nucleophiles now happens not only in the interaction with the substrate in the TS of the reaction but also in the interaction with the carbon acid  $\text{CH}_3\text{CH}_2^+$  which defines the ECA. We had found this previously for  $\text{S}_{\text{N}}2$  reactions of unsolvated  $\alpha$ -nucleophiles.<sup>77</sup>

Herein, we have been able to extend this finding to microsolvated nucleophiles and to E2 reactions. Thus, the barriers *versus* ECA correlations were constructed (Fig. S12d, ESI†), and the degree of downward deviation is indeed greatly reduced. In fact, as shown in Fig. 6c and d, the  $\Delta H^\ddagger$  values of both the  $\text{HOO}^-(\text{H}_2\text{O})_0$ -path ( $\alpha$ -nucleophile) and the  $\text{HO}^-(\text{H}_2\text{O})_0$ -path (normal nucleophile) have a good linear relationship with the ECA.

This phenomenon of reduced deviation is also observed when the  $\text{HY}^-(\text{H}_2\text{O})_{0-3}$  is used as the reference, where  $\text{Y} = \text{O, S, and HN}$  (Fig. S15, ESI†), *i.e.* type II correlation. Altogether, our computed Brønsted-type correlations reveal that the microhydrated  $\text{HOO}^-(\text{H}_2\text{O})_{0-4}$  nucleophiles exhibit the  $\alpha$ -effect in both  $\text{S}_{\text{N}}2$  and E2 reactions.

## Conclusions

We have computed highly accurate potential energy profiles for various E2 and  $\text{S}_{\text{N}}2$  pathways involved in the  $\text{HOO}^-(\text{H}_2\text{O})_{0-4} + \text{CH}_3\text{CH}_2\text{Cl}$  reaction system, involving both  $\text{HOO}^-$  and  $\text{HO}^-$  as attacking nucleophiles, based on a correlated CCSD(T)/aug-cc-pVTZ//MP2/6-311++G(d,p) approach. Our work provides both a benchmark description and a unified conceptual framework for a collection of interesting kinetic and structural phenomena

that occur in our chemically rich series of model reactions in which microsolvated  $\text{HOO}^-$  has a dual appearance due to the facile solvent-induced formation of microsolvated  $\text{HO}^-$ .

The  $\text{S}_{\text{N}}2$  path dominates the E2 path in our model systems. Adding water molecules further enhances the dominance of the  $\text{S}_{\text{N}}2$  reaction. This is so for both nucleophiles,  $\text{HOO}^-(\text{H}_2\text{O})_{0-4}$  and  $\text{HO}^-(\text{HOOH})(\text{H}_2\text{O})_{0-3}$ . Thus, the E2 barrier rises further above the  $\text{S}_{\text{N}}2$  barrier with each additional water molecule. This trend emerges from the combination of two factors: (i) the  $\text{S}_{\text{N}}2$  mechanism is associated with a smaller characteristic distortion and thus less activation strain  $\Delta E_{\text{strain}}$  than the E2 mechanism; (ii) therefore, as the nucleophile–substrate interaction  $\Delta E_{\text{int}}$  is weakened due to microsolvation, the barrier for the E2 path rises faster than that for  $\text{S}_{\text{N}}2$  and the latter pathway becomes more dominant.

In the  $\text{S}_{\text{N}}2$  substitution, the initial  $\text{HOO}^-$  nucleophile is clearly more reactive than the associated solvent-induced  $\text{HO}^-$  nucleophile. But, in the E2 elimination, the difference in reactivity is significantly smaller, with  $\text{HOO}^-$  still being somewhat more reactive. Thus, we found that the  $\text{HOO}^-(\text{H}_2\text{O})_{0-4}$  nucleophiles display the  $\alpha$ -effect in both the  $\text{S}_{\text{N}}2$  and, to a lesser extent, also in the E2 reaction. We show that the microsolvated  $\alpha$ -nucleophiles satisfy the earlier proposed criteria for the occurrence of the  $\alpha$ -effect, namely, a higher-energy HOMO and less occupied amplitude on the nucleophilic center, as compared to the corresponding normal nucleophile.

Our present work provides a unified description and rationalization of the reaction potential energy surface (PES) and kinetic and structural phenomena determined by this PES. A next leap forward that we envisage is the exploration of the complex dynamics taking place on this mechanistically rich multi-mechanistic PES of the  $\text{HOO}^-(\text{H}_2\text{O})_{0-4} + \text{CH}_3\text{CH}_2\text{Cl}$  reactions.

## Computational methods

All calculations were performed using the Gaussian 16 program.<sup>91</sup> To find an accurate method, the MP2,<sup>92</sup> B97-1,<sup>93</sup> B3LYP<sup>94</sup> and CAM-B3LYP<sup>95</sup> methods were tested on the reaction enthalpies of  $\text{HOO}^- + \text{CH}_3\text{CH}_2\text{Cl}$  to form  $\text{CH}_3\text{CH}_2\text{OOH} + \text{Cl}^-$  and  $\text{CH}_2 = \text{CH}_2 + \text{Cl}^- + \text{HOOH}$ . It turns out that the MP2/6-311++G(d,p) and MP2/aug-cc-pVTZ level of theories gave the best agreement with experimental values (Table S1, ESI†). We selected the MP2/6-311++G(d,p) method to perform the geometry optimization and frequency calculations throughout this work, for it is less time-consuming than the aug-cc-pVTZ basis set, and to be consistent with our previous work.<sup>38</sup> The nature of stationary points was confirmed by the frequencies under harmonic oscillator approximation, where energy minimum structures have no imaginary frequency and transition state structures have one imaginary frequency. The intrinsic reaction coordinate (IRC) calculations were performed for all transition states to ensure accuracy. On top of the geometries optimized with MP2/6-311++G(d,p) level of theory, single point calculations were performed using coupled cluster theory CCSD(T)<sup>96</sup>

with the aug-cc-pVTZ basis set.<sup>97,98</sup> If not specified, the energies reported in this work are at the CCSD(T)/aug-cc-pVTZ//MP2/6-311++G(d,p) level of theory.

## Data availability statement

The data that support the findings of this study are available in the ESI† of this article.

## Conflicts of interest

The authors declare no competing financial interests.

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## References

- 1 A. A. Viggiano, S. T. Arnold, R. A. Morris, A. F. Ahrens and P. M. Hierl, *J. Phys. Chem.*, 1996, **100**, 14397–14402.
- 2 J. V. Seeley, R. A. Morris, A. A. Viggiano, H. Wang and W. L. Hase, *J. Am. Chem. Soc.*, 1997, **119**, 577–584.
- 3 G. Caldwell, T. F. Magnera and P. Kebarle, *J. Am. Chem. Soc.*, 1984, **106**, 959–966.
- 4 A. P. Bento, M. Solà and F. M. Bickelhaupt, *J. Chem. Theory Comput.*, 2008, **4**, 929–940.
- 5 M. Swart, M. Solà and F. M. Bickelhaupt, *J. Chem. Theory Comput.*, 2010, **6**, 3145–3152.
- 6 J. Liang, Q. Su, Y. Wang and Z. Geng, *Bull. Chem. Soc. Jpn.*, 2015, **88**, 110–116.
- 7 I. Szabó and G. Czakó, *J. Phys. Chem. A*, 2017, **121**, 5748–5757.
- 8 V. Tajti and G. Czakó, *J. Phys. Chem. A*, 2017, **121**, 2847–2854.
- 9 D. A. Tasi, Z. Fábián and G. Czakó, *J. Phys. Chem. A*, 2018, **122**, 5773–5780.
- 10 D. A. Tasi, C. Tokaji and G. Czakó, *Phys. Chem. Chem. Phys.*, 2021, **23**, 13526–13534.
- 11 Z. Kerekes, D. A. Tasi and G. Czakó, *J. Phys. Chem. A*, 2022, **126**, 889–900.
- 12 A. A. Mohamed and F. Jensen, *J. Phys. Chem. A*, 2001, **105**, 3259–3268.
- 13 C. K. Regan, S. L. Craig and J. I. Brauman, *Science*, 2002, **295**, 2245–2247.
- 14 G. Vayner, K. N. Houk, W. L. Jorgensen and J. I. Brauman, *J. Am. Chem. Soc.*, 2004, **126**, 9054–9058.
- 15 X. Chen, C. K. Regan, S. L. Craig, E. H. Krenske, K. N. Houk, W. L. Jorgensen and J. I. Brauman, *J. Am. Chem. Soc.*, 2009, **131**, 16162–16170.



16 J. M. Garver, N. Eyet, S. M. Villano, Z. Yang and V. M. Bierbaum, *Int. J. Mass Spectrom.*, 2011, **301**, 151–158.

17 S. Nettey, C. A. Swift, R. Joviliano, D. O. Noin and S. Gronert, *J. Am. Chem. Soc.*, 2012, **134**, 9303–9310.

18 E. Carrascosa, J. Meyer, T. Michaelsen, M. Stei and R. Wester, *Chem. Sci.*, 2018, **9**, 693–701.

19 M. Gallegos, A. Costales and Á. Martín Pendás, *J. Comput. Chem.*, 2022, **43**, 785–795.

20 M. Gallegos, A. Costales and Á. Martín Pendás, *J. Phys. Chem. A*, 2022, **126**, 1871–1880.

21 X. Lu, C. Shang, L. Li, R. Chen, B. Fu, X. Xu and D. H. Zhang, *Nat. Commun.*, 2022, **13**, 4427.

22 J. Chandrasekhar, S. F. Smith and W. L. Jorgensen, *J. Am. Chem. Soc.*, 1984, **106**, 3049–3050.

23 P. M. Hierl, A. F. Ahrens, M. Henchman, A. A. Viggiano, J. F. Paulson and D. C. Clary, *J. Am. Chem. Soc.*, 1986, **108**, 3142–3143.

24 J. P. Bergsma, B. J. Gertner, K. R. Wilson and J. T. Hynes, *J. Chem. Phys.*, 1987, **86**, 1356–1376.

25 F. M. Bickelhaupt, E. J. Baerends and N. M. M. Nibbering, *Chem. – Eur. J.*, 1996, **2**, 196–207.

26 S. Kato, J. Hacaloglu, G. E. Davico, C. H. DePuy and V. M. Bierbaum, *J. Phys. Chem. A*, 2004, **108**, 9887–9891.

27 X. Liu, J. Zhang, L. Yang and W. L. Hase, *J. Am. Chem. Soc.*, 2018, **140**, 10995–11005.

28 F. Yu, *J. Phys. Chem. A*, 2022, **126**, 4342–4348.

29 Y. Li, C. Li, D. Gao and D. Wang, *J. Phys. Chem. A*, 2022, **126**, 5527–5533.

30 D. K. Bohme and A. B. Raksit, *Can. J. Chem.*, 1985, **63**, 3007–3011.

31 J. M. Garver, Y. Fang, N. Eyet, S. M. Villano, V. M. Bierbaum and K. C. Westaway, *J. Am. Chem. Soc.*, 2010, **132**, 3808–3814.

32 R. Otto, J. Brox, S. Trippel, M. Stei, T. Best and R. Wester, *Nat. Chem.*, 2012, **4**, 534–538.

33 D. K. Bohme and G. I. Mackay, *J. Am. Chem. Soc.*, 1981, **103**, 978–979.

34 J. Xie, X. Ma, J. Zhang, P. M. Hierl, A. A. Viggiano and W. L. Hase, *Int. J. Mass Spectrom.*, 2017, **418**, 122–129.

35 X. Ji, C. Zhao and J. Xie, *Phys. Chem. Chem. Phys.*, 2021, **23**, 6349–6360.

36 C. Zhao, X. Ma, X. Wu, D. L. Thomsen, V. M. Bierbaum and J. Xie, *J. Phys. Chem. Lett.*, 2021, **12**, 7134–7139.

37 X. Wu, S. Zhang and J. Xie, *Phys. Chem. Chem. Phys.*, 2022, **24**, 12993–13005.

38 X. Wu, C. Zhao and J. Xie, *ChemPhysChem*, 2022, **23**, e202200285.

39 Y. Hu, X. Wu and J. Xie, *Phys. Chem. Chem. Phys.*, 2023, **25**, 1947–1956.

40 D. K. Bohme and A. B. Raksit, *J. Am. Chem. Soc.*, 1984, **106**, 3447–3452.

41 X. G. Zhao, S. C. Tucker and D. G. Truhlar, *J. Am. Chem. Soc.*, 1991, **113**, 826–832.

42 R. A. J. O'Hair, G. E. Davico, J. Hacaloglu, T. T. Dang, C. H. DePuy and V. M. Bierbaum, *J. Am. Chem. Soc.*, 1994, **116**, 3609–3610.

43 S. Re and K. Morokuma, *J. Phys. Chem. A*, 2001, **105**, 7185–7197.

44 J. Xie, R. Otto, R. Wester and W. L. Hase, *J. Chem. Phys.*, 2015, **142**, 244308.

45 R. Otto, J. Brox, S. Trippel, M. Stei, T. Best and R. Wester, *J. Phys. Chem. A*, 2013, **117**, 8139–8144.

46 J. Xie, M. J. Scott, W. L. Hase, P. M. Hierl and A. A. Viggiano, *Z. Phys. Chem.*, 2015, **229**, 1747–1763.

47 J. Zhang, L. Yang, J. Xie and W. L. Hase, *J. Phys. Chem. Lett.*, 2016, **7**, 660–665.

48 B. Bastian, T. Michaelsen, L. Li, M. Ončák, J. Meyer, D. H. Zhang and R. Wester, *J. Phys. Chem. A*, 2020, **124**, 1929–1939.

49 X. Lu, L. Li, X. Zhang, B. Fu, X. Xu and D. H. Zhang, *J. Phys. Chem. Lett.*, 2022, **13**, 5253–5259.

50 R. Otto, J. Xie, J. Brox, S. Trippel, M. Stei, T. Best, M. R. Siebert, W. L. Hase and R. Wester, *Faraday Discuss.*, 2012, **157**, 41–57.

51 Y.-R. Wu and W.-P. Hu, *J. Am. Chem. Soc.*, 1999, **121**, 10168–10177.

52 N. Eyet, S. M. Villano, S. Kato and V. M. Bierbaum, *J. Am. Soc. Mass Spectrom.*, 2007, **18**, 1046–1051.

53 N. Eyet, S. M. Villano and V. M. Bierbaum, *J. Am. Soc. Mass Spectrom.*, 2008, **19**, 1296–1302.

54 N. Eyet, J. J. Melko, S. G. Ard and A. A. Viggiano, *Int. J. Mass Spectrom.*, 2015, **378**, 54–58.

55 T. Hansen, J. C. Roozee, F. M. Bickelhaupt and T. A. Hamlin, *J. Org. Chem.*, 2022, **87**, 1805–1813.

56 F. M. Lisboa and J. R. Pliego, *J. Mol. Model.*, 2022, **28**, 159.

57 D. J. Anick, *J. Phys. Chem. A*, 2011, **115**, 6327–6338.

58 D. L. Thomsen, J. N. Reece, C. M. Nichols, S. Hammerum and V. M. Bierbaum, *J. Am. Chem. Soc.*, 2013, **135**, 15508–15514.

59 J. O. Edwards and R. G. Pearson, *J. Am. Chem. Soc.*, 1962, **84**, 16–24.

60 W. P. Jencks and J. Carriuolo, *J. Am. Chem. Soc.*, 1960, **82**, 1778–1786.

61 E. Buncel, H. Wilson and C. Chuaqui, *J. Am. Chem. Soc.*, 1982, **104**, 4896–4900.

62 E. Buncel and I.-H. Um, *J. Chem. Soc., Chem. Commun.*, 1986, 595, DOI: [10.1039/C39860000595](https://doi.org/10.1039/C39860000595).

63 I.-H. Um and E. Buncel, *J. Org. Chem.*, 2000, **65**, 577–582.

64 M. Laloi-Diard, J.-F. Verchere, P. Gosselin and F. Terrier, *Tetrahedron Lett.*, 1984, **25**, 1267–1268.

65 I.-H. Um, E.-J. Lee and E. Buncel, *J. Org. Chem.*, 2001, **66**, 4859–4864.

66 J. F. Liebman and R. M. Pollack, *J. Org. Chem.*, 1973, **38**, 3444–3445.

67 A. Pross and S. S. Shaik, *J. Am. Chem. Soc.*, 1981, **103**, 3702–3709.

68 S. Hoz, *J. Org. Chem.*, 1982, **47**, 3545–3547.

69 M. M. Heaton, *J. Am. Chem. Soc.*, 1978, **100**, 2004–2008.

70 Y. Ren and H. Yamataka, *J. Org. Chem.*, 2007, **72**, 5660–5667.

71 Y. Ren and H. Yamataka, *Chem. – Eur. J.*, 2007, **13**, 677–682.

72 J. M. Garver, S. Gronert and V. M. Bierbaum, *J. Am. Chem. Soc.*, 2011, **133**, 13894–13897.



73 Y. Ren, X.-G. Wei, S.-J. Ren, K.-C. Lau, N.-B. Wong and W.-K. Li, *J. Comput. Chem.*, 2013, **34**, 1997–2005.

74 D. L. Thomsen, J. N. Reece, C. M. Nichols, S. Hammerum and V. M. Bierbaum, *J. Phys. Chem. A*, 2014, **118**, 8060–8066.

75 Y. Ren and H. Yamataka, *J. Comput. Chem.*, 2009, **30**, 358–365.

76 X.-G. Wei, X.-M. Sun, X.-P. Wu, Y. Ren, N.-B. Wong and W.-K. Li, *J. Org. Chem.*, 2010, **75**, 4212–4217.

77 T. Hansen, P. Vermeeren, F. M. Bickelhaupt and T. A. Hamlin, *Angew. Chem., Int. Ed.*, 2021, **60**, 20840–20848.

78 X. Liu, J. Xie, J. Zhang, L. Yang and W. L. Hase, *J. Phys. Chem. Lett.*, 2017, **8**, 1885–1892.

79 A. P. Bento and F. M. Bickelhaupt, *J. Org. Chem.*, 2008, **73**, 7290–7299.

80 F. M. Bickelhaupt, *J. Comput. Chem.*, 1999, **20**, 114–128.

81 J. V. Seeley, R. A. Morris and A. A. Viggiano, *J. Phys. Chem. A*, 1997, **101**, 4598–4601.

82 C. H. DePuy, S. Gronert, A. Mullin and V. M. Bierbaum, *J. Am. Chem. Soc.*, 1990, **112**, 8650–8655.

83 M. N. Glukhovtsev, A. Pross and L. Radom, *J. Am. Chem. Soc.*, 1996, **118**, 6273–6284.

84 F. M. Bickelhaupt and K. N. Houk, *Angew. Chem., Int. Ed.*, 2017, **56**, 10070–10086.

85 P. Vermeeren, S. C. C. van der Lubbe, C. Fonseca Guerra, F. M. Bickelhaupt and T. A. Hamlin, *Nat. Protoc.*, 2020, **15**, 649–667.

86 P. Vermeeren, T. Hansen, P. Jansen, M. Swart, T. A. Hamlin and F. M. Bickelhaupt, *Chem. – Eur. J.*, 2020, **26**, 15538–15548.

87 S. Hoz and E. Buncel, *Isr. J. Chem.*, 1985, **26**, 313–319.

88 W.-Y. Zhao, J. Yu, S.-J. Ren, X.-G. Wei, F.-Z. Qiu, P.-H. Li, H. Li, Y.-P. Zhou, C.-Z. Yin, A.-P. Chen, H. Li, L. Zhang, J. Zhu, Y. Ren and K.-C. Lau, *J. Comput. Chem.*, 2015, **36**, 844–852.

89 D. L. Thomsen, C. M. Nichols, J. N. Reece, S. Hammerum and V. M. Bierbaum, *J. Am. Soc. Mass Spectrom.*, 2014, **25**, 159–168.

90 L. Yun-Yun, Q. Fang-Zhou, Z. Jun, R. Yi and L. Kai-Chung, *J. Mol. Model.*, 2017, **23**, 192.

91 M. J. Frisch, G. W. Trucks, H. B. Schlegel, G. E. Scuseria, M. A. Robb, J. R. Cheeseman, G. Scalmani, V. Barone, G. A. Petersson, H. Nakatsuji, X. Li, M. Caricato, A. V. Marenich, J. Bloino, B. G. Janesko, R. Gomperts, B. Mennucci, H. P. Hratchian, J. V. Ortiz, A. F. Izmaylov, J. L. Sonnenberg, D. Williams-Young, F. Ding, F. Lipparini, F. Egidi, J. Goings, B. Peng, A. Petrone, T. Henderson, D. Ranasinghe, V. G. Zakrzewski, J. Gao, N. Rega, G. Zheng, W. Liang, M. Hada, M. Ehara, K. Toyota, R. Fukuda, J. Hasegawa, M. Ishida, T. Nakajima, Y. Honda, O. Kitao, H. Nakai, T. Vreven, K. Throssell, J. A. Montgomery Jr., J. E. Peralta, F. Ogliaro, M. J. Bearpark, J. J. Heyd, E. N. Brothers, K. N. Kudin, V. N. Staroverov, T. A. Keith, R. Kobayashi, J. Normand, K. Raghavachari, A. P. Rendell, J. C. Burant, S. S. Iyengar, J. Tomasi, M. Cossi, J. M. Millam, M. Klene, C. Adamo, R. Cammi, J. W. Ochterski, R. L. Martin, K. Morokuma, O. Farkas, J. B. Foresman and D. J. Fox, *Gaussian 16 Revision A.03*, Gaussian Inc., Wallingford CT, 2016.

92 M. J. Frisch, M. Head-Gordon and J. A. Pople, *Chem. Phys. Lett.*, 1990, **166**, 275–280.

93 F. A. Hamprecht, A. J. Cohen, D. J. Tozer and N. C. Handy, *J. Chem. Phys.*, 1998, **109**, 6264–6271.

94 P. J. Stephens, F. J. Devlin, C. F. Chabalowski and M. J. Frisch, *J. Phys. Chem.*, 1994, **98**, 11623–11627.

95 T. Yanai, D. P. Tew and N. C. Handy, *Chem. Phys. Lett.*, 2004, **393**, 51–57.

96 G. D. Purvis and R. J. Bartlett, *J. Chem. Phys.*, 1982, **76**, 1910–1918.

97 T. H. Dunning, *J. Chem. Phys.*, 1989, **90**, 1007–1023.

98 D. E. Woon and T. H. Dunning, *J. Chem. Phys.*, 1993, **98**, 1358–1371.

