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Investigating the behaviour of NaCl brines and  
hydrocarbons in porous alumina using low-field  
NMR relaxation and diffusion methods

The main focus in our research group is the utilization of low-field nuclear magnetic resonance (NMR) to study the behaviour of fluids in porous materials. In this particular study, we use proton NMR relaxation and diffusion experiments to monitor two different fluid phases during a spontaneous displacement process in order to assess the effect of NaCl concentration on the wettability of alumina towards brine solutions.

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## Investigating the behaviour of NaCl brines and hydrocarbons in porous alumina using low-field NMR relaxation and diffusion methods

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The behaviour of multiple fluid phases within a porous medium is hard to predict. NMR measurements offer an excellent tool to probe such systems in a fast and non-invasive way. Such systems can be relevant to hydrocarbon recovery, catalysis, and CO<sub>2</sub> and H<sub>2</sub> geo-storage, among others. Since electrolyte solutions are always present in subsurface reservoirs, understanding their behaviour within porous media is highly important. In this study, we use NMR relaxation and diffusion methods to investigate the diffusion coefficients and strength of interactions between alumina surfaces and brines at various NaCl concentrations, focusing on the effect of salt concentration on transport and interactions within the porous structure. Furthermore, we study the spontaneous displacement of dodecane, a model hydrocarbon, from the same alumina pellets using the same brine solutions. Results show that brines of lower salinity consistently displace more dodecane in total, after soaking dodecane-saturated pellets in a brine solution for several days. This indicates that increased salt concentrations can reduce wettability towards the aqueous phase in simple metal oxide surfaces and highlights the capabilities of NMR to efficiently study such systems.

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### Introduction

Investigating the behaviour of multiple fluids in porous media is of relevance in many areas of science and technology, such as hydrocarbon recovery, CO<sub>2</sub> storage, catalysis, and separation. One of the most important parameters to consider is the wettability of the pore surfaces towards each of those fluids. It controls the distribution of fluid phases within the pore spaces and determines the flow behaviour for both the wetting and non-wetting phases.

One particularly important aspect is the interactions between electrolyte solutions and mineral surfaces. Understanding such interactions is of significance as these systems are relevant to areas such as hydrocarbon recovery, carbon dioxide and hydrogen geo-storage, and geothermal applications. Low salinity waterflooding is one example of an area where interfacial wettability in porous rock plays an important role. The technique has been used to displace hydrocarbons

from porous rocks for secondary and tertiary oil recovery. The effectiveness of such techniques may be explained through expansion of electrical double layer and multicomponent ion exchange mechanisms;<sup>1,2</sup> however, uncertainties remain when it comes to implementations on reservoir scales.<sup>3,4</sup>

Nuclear magnetic resonance (NMR) relaxation has been used in numerous studies in recent years to probe the behaviour of fluids within restricted environments, such as biological cells, catalyst supports and rock materials.<sup>5–7</sup> Several aspects of porous structures can be probed, such as pore size distribution<sup>8,9</sup> and exchange of fluid molecules between different pore environments,<sup>10</sup> under the condition of sufficiently low diffusion coupling between these pores.<sup>11</sup> NMR also offers the possibility to investigate wettability of fluids with a fast, non-invasive and non-destructive approach,<sup>12,13</sup> which can offer advantages over traditional, more complicated methods to assess wettability, such as the USBM and Amott-Harvey index,<sup>14</sup> or contact angle methods<sup>15</sup> which can sometimes yield inconsistent results<sup>16</sup> and cannot fully represent the effects of complex geometries found within porous structures.

In this study, we investigate the potential of NMR to study the wetting properties of NaCl brines, as well their ability to spontaneously displace dodecane, a model hydrocarbon compound, from the pores of alumina pellets through an imbibition process at room temperature and atmospheric pressure.

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Since most natural rocks consist predominantly of metal oxide-based minerals, it would be reasonable to study porous extrudates of pure metal oxides, such as alumina, as an analogue for actual rock samples. The advantage of this approach for systematic studies is in the more controlled and reproducible pore structures offered by these materials, as opposed to the intrinsic heterogeneous nature of rocks in terms of structure and chemical composition, which makes experimental results difficult to interpret. NaCl has been chosen since it is by far the most common electrolyte found in natural brines.

## Theory

### NMR relaxation

When a sample containing spin-bearing particles is placed under a strong magnetic field, these spins will precess around an axis parallel to that external magnetic field, which is labelled as the  $z$ -axis, with an angular velocity ( $\omega_0$ ) given by:

$$\omega_0 = -\gamma B_0 \quad (1)$$

where  $\gamma$  is the gyromagnetic ratio of the particle and  $B_0$  is the strength of the applied magnetic field. A slight majority of the spins will precess pointing towards the direction of that external magnetic field, which is their lower energy state, generating a macroscopic net magnetisation ( $\bar{M}$ ). Electromagnetic pulses can be used to manipulate this magnetization, to obtain useful information about various physical and chemical properties of the sample. More details on the physics of magnetic resonance can be found in various textbooks.<sup>17–19</sup>

After the sample's net magnetisation is disturbed by this pulse, it relaxes back to its equilibrium state, parallel to the  $z$ -axis. The restoration of the  $z$ -axis component back to the equilibrium level is referred to as "longitudinal relaxation", and the diminishing of the  $x$  and  $y$  axes' components back to zero is called "transverse relaxation". For spins in liquid

molecules, both longitudinal and transverse relaxation occur exponentially, and their rates can provide information about the molecular dynamics existing within the sample.<sup>20</sup> Longitudinal relaxation is characterised by the " $T_1$ " time constant, while transverse relaxation is characterised by the " $T_2$ " time constant.

In the case of fluids confined within some restrictive geometry, such as in a porous medium, relaxation times are usually much shorter than those of the same liquids in their bulk state. This is because the surfaces of the confining solid act as relaxation sinks, whether these involve spin bearing particles, paramagnetic species or surface adsorption sites.<sup>21</sup>

For porous media with limited pore connectivity,<sup>22</sup> both  $T_1$  and  $T_2$  relaxation time distributions can usually be correlated to the pore size distribution of a porous medium.<sup>9</sup> The relationship between  $T_1$  and  $T_2$  values can be used to draw information about the wettability of the porous medium towards a fluid, since it is related to the surface residency time of the fluid's molecules.<sup>23</sup> The  $T_1/T_2$  ratio can be linked to an energy of adsorption for a fluid molecule onto a surface. This relationship comes from the different dependencies of  $T_1$  and  $T_2$  on molecular motions, and can be mathematically proven for surface adsorption on both paramagnetic<sup>24,25</sup> and non-paramagnetic<sup>26,27</sup> sites. Several experimental studies have shown the validity of such measurements.<sup>26,28</sup>

### NMR diffusion

The precession frequency of a spin-bearing nucleus depends on the strength of the external magnetic field it experiences. Modern NMR instruments are capable of generating magnetic field gradients (using short pulses of constant current applied to suitable gradient coils), which effectively label different positions within the sample, according to their different precession frequencies. In the case of a one-dimensional gradient

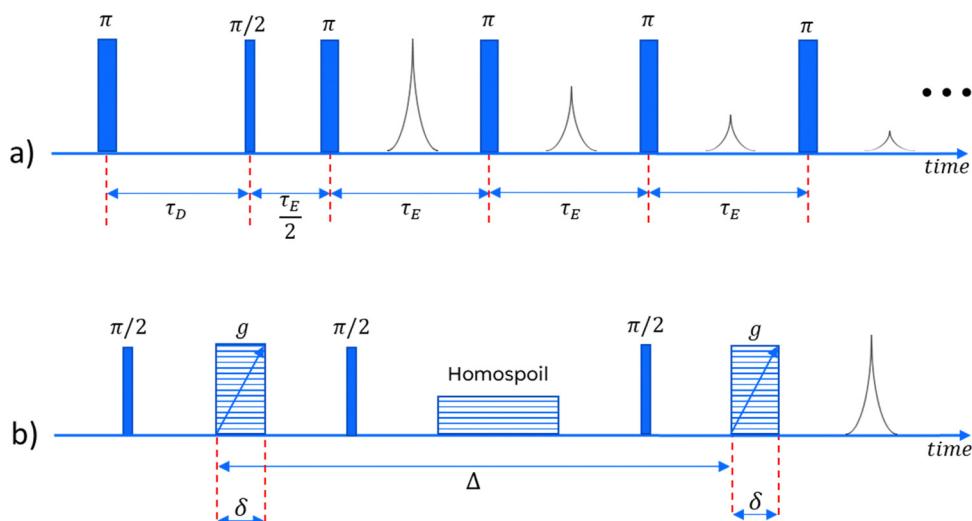


Fig. 1 (a) The IR-CPMG pulse sequence, which is repeated for several different values of delay time ( $\tau_D$ ) while the echo time ( $\tau_E$ ) remains constant. (b) The PGSTE pulse sequence, repeated for several different magnitudes of magnetic field gradient ( $g$ ). The duration of the gradient pulses is  $\delta$ , and the time interval between those pulses is  $\Delta$ . Striped boxes represent gradient pulses, along the direction of the external field.



along the  $z$ -axis, the precession frequency at each point will be:<sup>29</sup>

$$\omega(z) = \omega_0 - \gamma g z \quad (2)$$

where  $z$  is the position of the spin along the  $z$ -axis, and  $g$  is the magnetic field gradient, or the slope of the magnetic field's variation along the same axis.

After a radiofrequency pulse creates a net magnetisation, the movement of molecules between regions of different magnetic fields introduces additional dephasing of the spin ensemble. The extent of this dephasing results in a signal attenuation, which can be exploited by carefully designed pulse sequences in order to accurately quantify the diffusion coefficients of the molecules within the sample. Further information about diffusion measurements and the theory behind them can be found in the literature.<sup>30–32</sup>

### Pulse sequences

In order to obtain meaningful information about the NMR behaviour of any sample, we must use a series of well-timed radiofrequency pulses. In this study, for the relaxation measurements, we employ an IR-CPMG sequence<sup>33</sup> (a combination of the Inversion Recovery and Carr-Purcell-Meiboom-Gill pulse sequences), capable of obtaining information about both longitudinal and transverse relaxation in one single experiment. For the diffusivity measurements a pulsed field gradient stimulated echo (PGSTE) sequence is used.<sup>34</sup> Both of these sequences are illustrated in Fig. 1.

### Numerical inversion

In the presence of multiple relaxation behaviours within the sample, these different components are hidden behind a single NMR signal. To separate the contributions of each component, numerical inversion methods are commonly used. Effectively, it is a process where the experimental results are fitted with a sum of a large number of relaxation decays. Mathematically, this multi exponential decay can be expressed as a Fredholm integral of the first kind:<sup>35</sup>

$$\frac{M(t)}{M_0} = \int_0^\infty k(T, t) \cdot f(T) dT + \varepsilon(t), \quad (3)$$

where  $f$  is the desired distribution of relaxation times (roughly akin to a probability density function),  $t$  is time,  $T$  is the relaxation time constant ( $T_1$  or  $T_2$ ),  $\varepsilon$  is the noise contained within the signal, and  $k$  is referred to as the “kernel”, which depends on the type of the experiment. For example:

$$k(T_1, t) = 1 - 2 \exp\left(-\frac{t}{T_1}\right), \quad (4)$$

$$k(T_2, t) = \exp\left(-\frac{t}{T_2}\right), \quad (5)$$

are the kernels used for IR and CPMG experiments, respectively, in order to obtain  $T_1$  and  $T_2$  relaxation time distributions.

Conventional algorithms used for curve fitting cannot handle such complex forms, and thus a regularisation process

must be used.<sup>36</sup> The inverse problem of finding a realistic set of  $f(T)$  values can be solved using methods such as non-linear least squares, maximum entropy, or equivalent. These methods can be used for both 1D and 2D NMR experiments, sometimes involving both relaxation and diffusion measurements.<sup>37</sup> Further information about the inversion process can be found in the literature.<sup>35,38–40</sup>

## Experimental

### Sample preparation

For the experiments performed in this study, cylindrical pellets of pure alumina (approximately 2 mm in diameter and 5 mm in length) were used. The brine samples have been prepared using NaCl of purity above 99%, dissolved in deionised water. The concentration range used is from 1% w/v up to 25% w/v, which is close to the solubility limit of NaCl at room temperature. The alumina pellets were dried overnight in a vacuum oven at 70 °C, in order to remove any residual moisture and other potential impurities. Afterwards, they were soaked in the prepared brine samples in order to saturate the pores, for at least 48 hours before an experiment. Before the measurements, the pellets were removed from the liquid, and carefully dried on a filter paper so that the outer surface would be dry, without removing any significant amount of liquid from within the pores of the pellet. This is so that the obtained signal will come exclusively from the pore-bound fluid, and not from an external bulk liquid layer.

### Methodology

NMR relaxation and diffusion experiments were performed on the alumina samples, saturated with brines of various NaCl concentrations.

Imbibition experiments were also performed, whereby brines of different salinities were used to spontaneously displace dodecane from within the alumina pellets. The pellets were dried and soaked in dodecane for at least 48 h before the experiments. They were then placed in a container full of brine, and removed after a designated time interval, after which a  $T_1-T_2$  measurement was performed. The process was repeated for ten different time steps, for all ten of the brine solutions used, with the longest soaking interval lasting about one week. The pellets inherently have a higher water wetting tendency and dodecane is not miscible with the aqueous phase, thus the brine solutions would imbibe into the pores of the pellet, displacing some of the dodecane.

This displacement process can be studied by looking at the volume fractions of water and dodecane through time. These two can be easily quantified using NMR since their relaxation times within the pores of the alumina pellets differ significantly. Their respective volume fractions can be estimated by calculating the integrals under the peaks representing each of the fluids in a 2D  $T_1-T_2$  map,<sup>41</sup> obtained by inverting the experimental data by methods discussed in the previous section, utilising the Tikhonov regularisation method.<sup>42</sup>



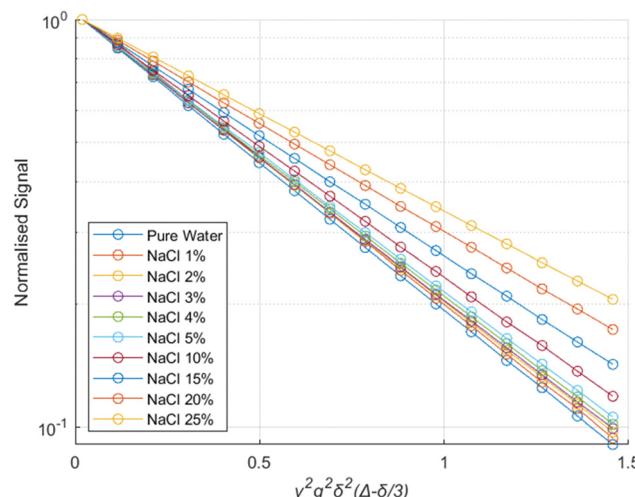


Fig. 2 Stejskal-Tanner plot for brines of different salinities within the pores of alumina pellets. The solid lines represent exponential fits using eqn (6).

Diffusion coefficients were estimated by fitting the experimental results with an exponential equation of the form:

$$\frac{S}{S_0} = \exp \left[ -D \gamma^2 g^2 \delta^2 \left( \Delta - \frac{\delta}{3} \right) \right], \quad (6)$$

where  $S/S_0$  is the normalised signal,  $D$  is the diffusion coefficient,  $\gamma$  is the gyromagnetic ratio of the observed nucleus,  $g$  is the magnetic field gradient,  $\Delta$  is the time interval between the gradient pulses and  $\delta$  is the duration of the gradient pulses (as illustrated in Fig. 1b).

## NMR hardware

A Magritek Spinsolve™ spectrometer was used, with a magnetic field of approximately 1 Tesla, corresponding to a proton Larmor frequency of around 43 MHz. A gradient coil is included in the instrument, which can be used to apply one-dimensional gradients up to  $163 \text{ mT m}^{-1}$ . For the IR-CPMG sequence, 16 inversion recovery cycles have been used for each experiment, with delay times logarithmically spaced between 1 and 3000 ms. The echo time for the CPMG part of the sequence was 250  $\mu\text{s}$ . For the PGSTE experiments,  $\delta$  was fixed at 2 ms and  $\Delta$  at 300 ms, while  $g$  was varied between 14 and  $130 \text{ mT m}^{-1}$ , in 16 linearly spaced intervals.

## Results and discussion

### Diffusion measurements

The experimental results from the PGSTE measurements are presented in a Stejskal-Tanner plot<sup>43</sup> in Fig. 2. The data for all different NaCl concentrations follow clear mono-exponential trends, implying that there is a single diffusion regime within the samples over the time scale probed by PGSTE measurements. This behaviour is typical of porous materials showing a macroscopically homogeneous pore structure and is often referred to as *quasi-homogeneous* behaviour.<sup>44,45</sup> A similar behaviour is observed for the same brine samples in their bulk state.

The results show that the diffusion coefficients of the water molecules consistently decrease with increasing salinity, both in bulk and within the alumina samples, as shown in Fig. 3.

The reduction in the diffusion coefficients for the brines within the pores are mostly due to the increase in the viscosity

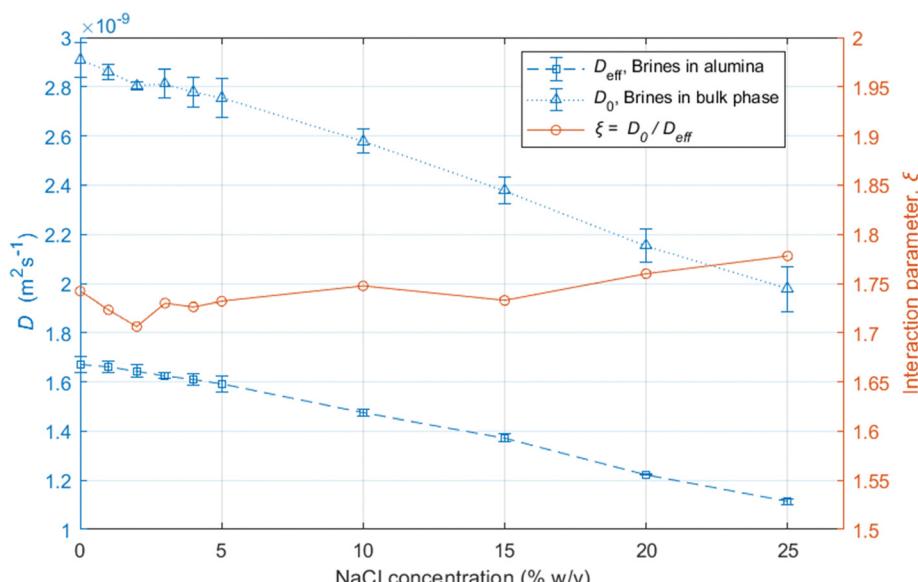
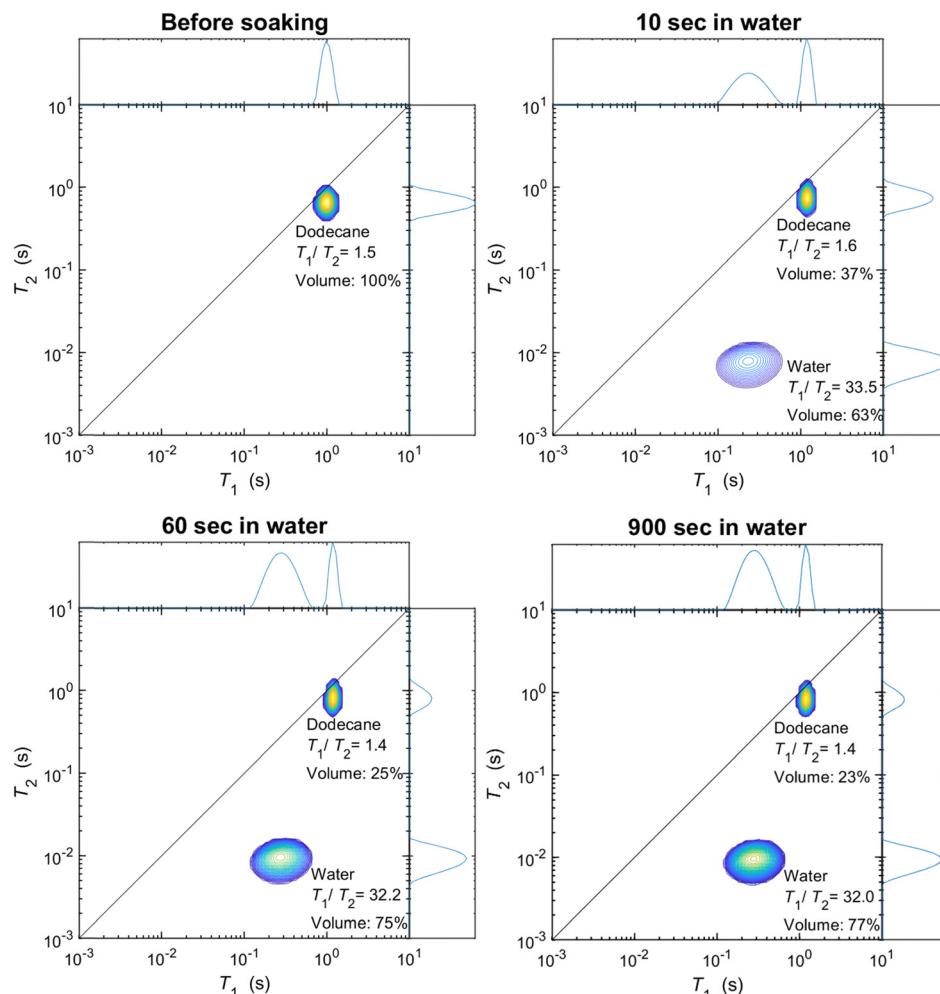


Fig. 3 Diffusion coefficients and  $\xi$  values versus salinity, for NaCl brines in mixed-phase alumina pellets. Error bars were calculated from the standard deviations between three iterations of the same measurements, using different pellets. All the lines between data points are a guide to the eye, not fittings.





**Fig. 4**  $T_1$ – $T_2$  maps of the alumina pellet at four different time steps during an imbibition experiment. Initially, only one peak is present, representing the hydrocarbon phase. After soaking the pellet in water, a second peak appears, whose volume increases at longer soaking times.

of the solutions due to the addition of NaCl,<sup>46</sup> also reflected by the bulk brine measurements.

The ratio between the diffusivity of a liquid in its bulk state and the same liquid confined within pores can provide information related to the tortuosity of the porous medium.<sup>44,47</sup> However, this is only valid in the case of molecules which do not interact strongly with the surfaces, as shown elsewhere.<sup>48</sup> Otherwise, a PFG interaction parameter is defined, as:

$$\xi = \frac{D_0}{D_{\text{eff}}} \quad (7)$$

where  $D_0$  is the bulk diffusion coefficient and  $D_{\text{eff}}$  is the diffusion coefficient within some restricted environment. This parameter can be used to characterise the reduction of diffusivity within a porous medium, which effectively includes contributions of both tortuosity and surface interactions, and thus it can also be referred to as “apparent tortuosity”.<sup>49,50</sup>

The  $\xi$  value for the alumina samples is slightly lower for the low salinity brines, indicating that the decrease in the movement of water molecules within the alumina samples is slightly greater for lower salinities. This could be partly due to stronger

interaction of water molecules with the pore walls for the low-salinity brines, which is supported by the relaxation measurements presented in the next section. Since that effect would only have an impact on the surface layers and not the bulk liquid in the middle of the pores, the variation in  $\xi$  is not very significant.

#### Relaxation measurements – dodecane displacement

As discussed in the experimental section, we look at the volume fractions of brine and dodecane at several different time steps during the soaking process (10 s, 20 s, 1 min, 10 min, 15 min, 1 h, 24 h, 1 week). This information is obtained by integrating the peaks representing each of the fluid phases in the  $T_1$ – $T_2$  maps, such as the examples presented in Fig. 4. The position of the same peaks provides the value of the  $T_1/T_2$  ratio for each of the fluids.

$T_1$ – $T_2$  experiments have also been performed in alumina pellets fully saturated with each of the different brine concentrations tested. The results of such measurements are presented in Fig. 5, along with the volume fractions of the same



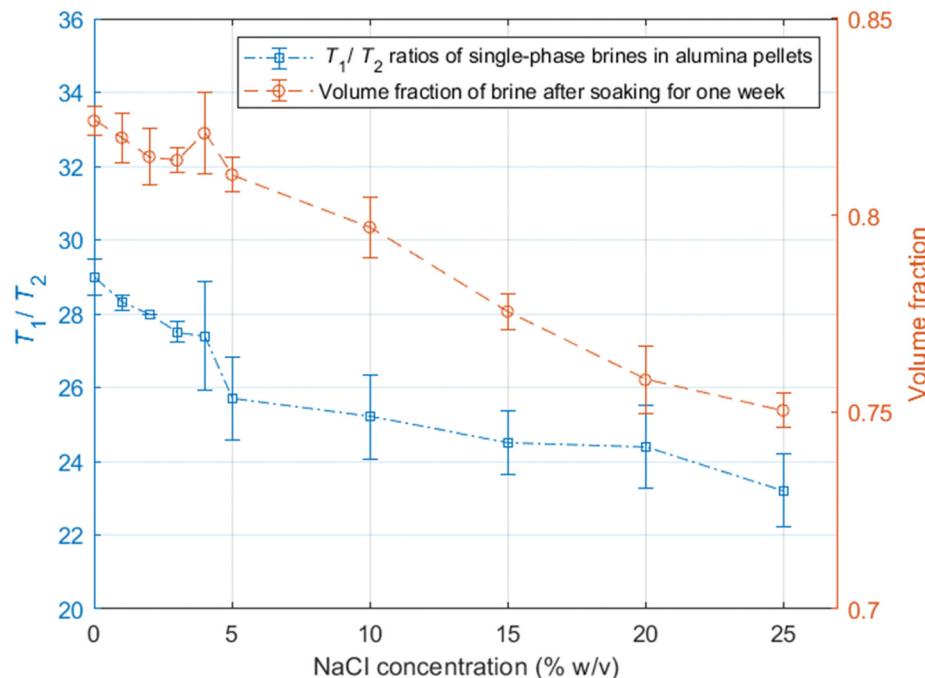


Fig. 5  $T_1/T_2$  ratios of brines in alumina pellets, and volume fractions of brines after a dodecane-saturated pellet was soaked in them for a period of one week, versus NaCl concentration. Error bars were calculated from the standard deviations between three iterations of the same measurements, using different alumina pellets. Variation in error bars can be attributed to differences in pore structure between different alumina pellets used in the repeated experiments. All the lines between data points are a guide to the eye, not fittings.

brine samples, after a dodecane-saturated pellet was soaked in them for a period of one week.

It is evident from the data in Fig. 5 that there is a strong correlation between the  $T_1/T_2$  ratio of the brines and the amount of dodecane they are able to displace. During the early time steps, no significant trends were observed between different salinities, and the rate of displacement was uncorrelated with salinity. This is possibly attributed to slight differences in pore structure and permeability between individual pellets, which would probably affect the initial displacement rates. However, looking at the final time steps, after a week of soaking, there is a clear trend, showing that brines of lower

salinities are consistently displacing more dodecane. This is also illustrated in Fig. 6.

The results show that a fraction of dodecane will always remain trapped in the pores of the alumina pellets, regardless of the soaking time in the brine solutions. This is because the non-wetting phase forms discontinuous ganglia surrounded by the mobile wetting phase,<sup>51,52</sup> which are very difficult to displace spontaneously by the wetting phase.

By looking at the  $T_1/T_2$  ratios of the brines at each step, we can see that brines of lower salinities generally tend to have higher  $T_1/T_2$  values, the same as the single-phase experiments. This confirms that surface affinity is higher for the brines at

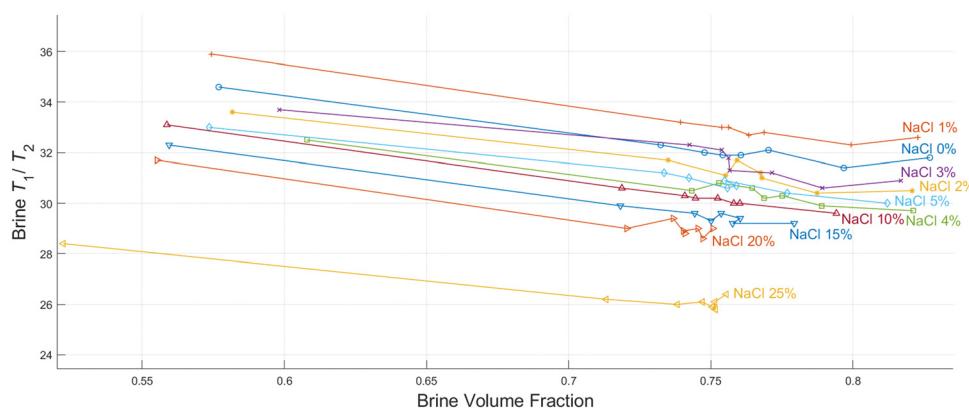


Fig. 6 Plot showing the  $T_1/T_2$  ratio of the ten different brine samples, versus the volume fraction of the same brines during one set of the imbibition experiments. Each data point represents a different time step, with the leftmost points representing 10 seconds of soaking time and the rightmost point representing a week of soaking, for each salinity. The solid lines connecting the data points act as a guide to the eye.



lower NaCl concentrations, even in multi-phase environments, which explains their ability to displace more dodecane. These results agree with the findings of Lee *et al.*,<sup>2</sup> who observed that in low salinities the water layers on silica surfaces are thicker compared to the water layers for higher salinities.

Furthermore, the  $T_1/T_2$  ratio for each of the brine samples is consistently reduced when the brine volume fraction is increased, as shown in Fig. 4 and 6. This could be attributed to the distribution of fluids within the pores. Since brine is the wetting phase, it mostly occupies the limited space between the pore surface and the non-wetting phase, which occupies the middle of the pores. When the volume fraction of the brine is lower, a bigger proportion of the water hydrogens experience enhanced surface relaxation and additional dephasing of the spin ensemble due to the magnetic field distortion caused by the magnetic susceptibility of the solid, leading to reduced  $T_2$  values. Nevertheless, these dephasing effects would be the same regardless of the salinity, and thus differences in  $T_1/T_2$  between the different brines can be interpreted as a clear indication of surface affinity.

## Conclusions

In this study, the effects of NaCl concentration in the wettability and the capacity for hydrocarbon displacement through spontaneous imbibition were investigated for several different brine samples in porous alumina pellets.

Both NMR relaxation and diffusion results show that interactions between the aqueous phase and the alumina surfaces are stronger for lower NaCl concentrations, indicating higher levels of wettability for lower salinity brines. Imbibition experiments show that low salinity brines consistently displace more dodecane. A strong correlation between the  $T_1/T_2$  ratio of the brines and their ability to displace hydrocarbons from the pores of the alumina pellets was also observed.

Such results show that the  $T_1/T_2$  ratio of fluids confined within the pores of a solid medium can be a valid indication of wettability in multi-phase systems and can help unravel the effect of salt concentration on transport and dynamics of fluids inside the pores, which has important implications for subsurface applications, where multi-phase fluid behaviour is of critical importance, whether the goal is to displace a fluid phase, or keep it trapped within the pores.

It is worth noting that NMR might be able to help us evaluate the ways fluids are distributed within pores, in systems where multiple fluid phases are present. Future studies could further investigate these effects, by comparing NMR results with imaging techniques for example.

## Author contributions

Aristarchos Mavridis: investigation, writing – original draft, data curation, formal analysis, validation, methodology. Carmine D'Agostino: conceptualisation, funding acquisition, supervision, project administration, writing – review & editing.

Mark Sankey, Kuhan Chellappah: supervision, project administration, writing – review & editing.

## Conflicts of interest

There are no conflicts to declare.

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## References

- 1 K. S. Sorbie and I. R. Collins, *SPE - DOE Improved Oil Recovery Symposium. Proceedings*, 2010, **1**, 760–777.
- 2 S. Y. Lee, K. J. Webb, I. R. Collins, A. Lager, S. M. Clarke, M. O'Sullivan, A. F. Routh and X. Wang, *Proceedings - SPE Symposium Improved Oil Recovery*, 2010, **1**, 534–544.
- 3 P. I. Sagbana, K. Sarkodie and W. A. Nkrumah, *Petroleum*, 2023, **9**, 317–330.
- 4 F. Liu and M. Wang, *Fuel*, 2020, **267**, 117112.
- 5 Y.-Q. Song, *J. Magn. Reson.*, 2013, **229**, 12–24.
- 6 L. F. Gladden and J. Mitchell, *New J. Phys.*, 2011, **13**, 035001.
- 7 K. T. O'Neill, E. O. Fridjonsson, D. Smeed, T. A. J. Hopper and M. L. Johns, *Icarus*, 2023, **399**, 115544.
- 8 S. Davies, M. Z. Kalam, K. J. Packer and F. O. Zelaya, *J. Appl. Phys.*, 1990, **67**, 3171–3176.
- 9 R. L. Kleinberg, *Magn. Reson. Imaging*, 1994, **12**, 271–274.
- 10 K. E. Washburn and P. T. Callaghan, *Phys. Rev. Lett.*, 2006, **97**, 175502.
- 11 M. Fleury, T. Gimmi and M. Mazurek, *Clays Clay Miner.*, 2022, **70**, 417–437.
- 12 A. Valori and B. Nicot, *PETROPHYSICS The SPWLA Journal of Formation Evaluation and Reservoir Description.*, 2019, **60**, 255–263.
- 13 A. Valori, F. Ali and W. Abdallah, in Day 2 Tue, March 27, 2018, SPE, 2018, vol. 2018, pp. 26–28.
- 14 W. Anderson, *J. Pet. Technol.*, 1986, **38**, 1246–1262.
- 15 J. Bachmann, S. K. Woche, M.-O. Goebel, M. B. Kirkham and R. Horton, *Water Resour. Res.*, 2003, **39**, 1353.
- 16 S. Iglauer, C. H. Pentland and A. Busch, *Water Resour. Res.*, 2015, **51**, 729–774.
- 17 M. H. Levitt, *Spin Dynamics: Basics of Nuclear Magnetic Resonance*, John Wiley & Sons Ltd, 2008.
- 18 J. Keeler, *Understanding NMR Spectroscopy*, John Wiley & Sons, Ltd, 2010.
- 19 R. J. Abraham, J. Fisher and P. Loftus, *Introduction to NMR spectroscopy*, Wiley, 1988.
- 20 J. Kowalewski, *Nucl. Magn. Reson.*, 2012, **41**, 196–243.
- 21 J.-P. Korb, *New J. Phys.*, 2011, **13**, 035016.



22 K. R. McCall, D. L. Johnson and R. A. Guyer, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1991, **44**, 7344–7355.

23 J. Mitchell, M. D. Hürlimann and E. J. Fordham, *J. Magn. Reson.*, 2009, **200**, 198–206.

24 R. L. Kleinberg, W. E. Kenyon and P. P. Mitra, *J. Magn. Reson., Ser. A*, 1994, **108**, 206–214.

25 P. J. McDonald, J.-P. Korb, J. Mitchell and L. Montelihet, *Phys. Rev. E: Stat., Nonlinear, Soft Matter Phys.*, 2005, **72**, 011409.

26 J. Mitchell, L. M. Broche, T. C. Chandrasekera, D. J. Lurie and L. F. Gladden, *J. Phys. Chem. C*, 2013, **117**, 17699–17706.

27 C. D'Agostino, J. Mitchell, M. D. Mantle and L. F. Gladden, *Chem. – A Eur. J.*, 2014, **20**, 13009–13015.

28 N. Robinson, C. Robertson, L. F. Gladden, S. J. Jenkins and C. D'Agostino, *ChemPhysChem*, 2018, **19**, 2472–2479.

29 D. Sinnaeve, *Concepts Magn. Reson. Part A Bridg. Educ. Res.*, 2012, **40**, 39–65.

30 W. S. Price, *NMR Studies of Translational Motion*, Cambridge University Press, 2009.

31 W. S. Price, *Concepts Magn. Reson.*, 1997, **9**, 299–336.

32 D. S. Grebenkov, *Concepts Magn. Reson., Part A*, 2008, **32A**, 277–301.

33 S. Vashaei, M. Li, B. Newling, B. MacMillan, F. Marica, H. T. T. Kwak, J. Gao, A. M. M. Al-harbi and B. J. J. Balcom, *J. Magn. Reson.*, 2018, **287**, 113–122.

34 S. A. Willis, T. Stait-Gardner, A. M. Torres and W. S. Price, *Diffusion NMR of Confined Systems: Fluid Transport in Porous Solids and Heterogeneous Materials*, 2016, pp. 16–51.

35 J. Mitchell and T. C. Chandrasekera, *J. Chem. Phys.*, 2014, **141**, 224201.

36 P. C. Hansen, *Inverse Probl.*, 1992, **8**, 849–872.

37 P. T. Callaghan, C. H. Arns, P. Galvosas, M. W. Hunter, Y. Qiao and K. E. Washburn, *Magn. Reson. Imaging*, 2007, **25**, 441–444.

38 Y.-Q. Song, L. Venkataraman, M. D. Hürlimann, M. Flaum, P. Frulla and C. Straley, *J. Magn. Reson.*, 2002, **154**, 261–268.

39 K. Dunn, D. Bergman and G. LaTorraca, *Nuclear magnetic resonance: Petrophysical and logging applications*, Pergamon, 2002.

40 V. Bortolotti, R. J. S. Brown, P. Fantazzini, G. Landi and F. Zama, *Inverse Probl.*, 2016, **33**, 015003.

41 B. Nicot, N. Vorapalawut, B. Rousseau, L. F. Madariaga, G. Hamon and J.-P. Korb, *Petrophysics*, 2016, **57**, 19–29.

42 J. Mitchell, T. C. Chandrasekera and L. F. Gladden, *Prog. Nucl. Magn. Reson. Spectrosc.*, 2012, **62**, 34–50.

43 E. O. Stejskal and J. E. Tanner, *J. Chem. Phys.*, 1965, **42**, 288–292.

44 C. D'Agostino, J. Mitchell, L. F. Gladden and M. D. Mantle, *J. Phys. Chem. C*, 2012, **116**, 8975–8982.

45 M. Dvoryashkin, R. Valiullin and J. Kärger, *Phys. Rev. E: Stat., Nonlinear, Soft Matter Phys.*, 2007, **75**, 041202.

46 D. E. Goldsack and R. Franchetto, *Can. J. Chem.*, 1977, **55**, 1062–1072.

47 M. D. Mantle, D. I. Enache, E. Nowicka, S. P. Davies, J. K. Edwards, C. D'Agostino, D. P. Mascarenhas, L. Durham, M. Sankar, D. W. Knight, L. F. Gladden, S. H. Taylor and G. J. Hutchings, *J. Phys. Chem. C*, 2011, **115**, 1073–1079.

48 N. Robinson and C. D'Agostino, *Top. Catal.*, 2020, **63**, 319–327.

49 C. D'Agostino, R. Liuzzi, L. F. Gladden and S. Guido, *Soft Matter*, 2017, **13**, 2952–2961.

50 J. A. Ward-Williams, V. Karsten, C. M. Guédon, T. A. Baart, P. Munnik, A. J. Sederman, M. D. Mantle, Q. Zheng and L. F. Gladden, *Chem.: Methods*, 2022, **2**, e202200025.

51 C. Wang, Y. Mehmani and K. Xu, *Proc. Natl. Acad. Sci. U. S. A.*, 2021, **118**, e2024069118.

52 I. Zarikos, A. Terzis, S. M. Hassanizadeh and B. Weigand, *Sci. Rep.*, 2018, **8**, 13228.

