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Structural diversity in the membrane-bound hIAPP dimer correlated with distinct membrane disruption mechanisms†

Qin Qiao,^a Guanghong Wei^b and Zhijian Song^{ab}

Amyloid deposits of the human islet amyloid polypeptide (hIAPP) have been identified in 90% of patients with type II diabetes. Cellular membranes accelerate the hIAPP fibrillation, and the integrity of membranes is also disrupted at the same time, leading to the apoptosis of β cells in pancreas. The molecular mechanism of hIAPP-induced membrane disruption, especially during the initial membrane disruption stage, has not been well understood yet. Herein, we carried out extensive all-atom molecular dynamics simulations investigating the hIAPP dimerization process in the anionic POPG membrane, to provide the detailed molecular mechanisms during the initial hIAPP aggregation stage in the membrane environment. Compared to the hIAPP monomer on the membrane, we observed not only an increase of α -helical structures, but also a substantial increase of β -sheet structures upon spontaneous dimerization. Moreover, the random coiled and α -helical dimer structures insert deep into the membrane interior with a few inter-chain contacts at the C-terminal region, while the β -sheet-rich structures reside on the membrane surface accompanied by strong inter-chain hydrophobic interactions. The coexistence of α and β structures constitutes a diverse structural ensemble of the membrane-bound hIAPP dimer. From α -helical to β -sheet structures, the degree of membrane disruption decreases gradually, and thus the membrane damage induced by random coiled and α -helical structures precedes that induced by β -sheet structures. We speculate that insertion of random coiled and α -helical structures contributes to the initial stage of membrane damage, while β -sheet structures on the membrane surface are more involved in the later stage of fibril-induced membrane disruption.

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1. Introduction

The human islet amyloid polypeptide (hIAPP), a 37-residue glucose-regulating hormone co-secreted with insulin, is associated with the development of type II diabetes (T2D), as it deposits as amyloid fibrils in more than 90% of T2D patients.¹ Experiments have indicated that the hIAPP exerts cytotoxicity by disrupting the cellular membrane during its amyloid aggregation process.^{2–4} Therefore, understanding the hIAPP-membrane interactions is essential to reveal the hIAPP pathogenic mechanisms, and thus provide guidance for the therapeutic strategies of T2D.

Various hypotheses have been proposed to explain the molecular mechanism of hIAPP-induced membrane damage, including: (1) the “pore formation” mechanism: the formation of hIAPP non-selective ion channels across the membrane breaking the cell osmosis;^{5–8} (2) the “detergent-like” mechanism: the extraction of lipids by the hIAPP N-terminal region leading to the membrane leakage;^{9–11} and (3) the “fibril elongation” mechanism: the elongation of hIAPP fibrils distorting the cell by inducing the mechanical pressure on the membrane surface.^{12,13} Several experimental studies indicated that the hIAPP-induced membrane damage is bi-phasic, with the 1st stage caused by prefibrillar oligomers, and the 2nd stage caused by fibril elongation,^{13,14} and thus different disruption mechanisms may be involved in different aggregation stages. Meanwhile, the particular mechanism of hIAPP-induced membrane damage also depends on experimental conditions, such as lipid compositions, hIAPP-to-lipid ratios, pH conditions, etc.^{15–21} Therefore, understanding the detailed molecular interactions during the hIAPP aggregation in a membrane environment is important for unravelling the molecular mechanism of hIAPP-induced membrane disruption.

^a Digital Medical Research Center, School of Basic Medical Sciences, Fudan University, Shanghai 200032, China. E-mail: qinqiao@fudan.edu.cn

^b Shanghai Key Laboratory of Medical Imaging Computing and Computer Assisted Intervention, Shanghai 200032, China

^c Department of Physics, State Key Laboratory of Surface Physics, Key Laboratory for Computational Physical Science (Ministry of Education), Fudan University, Shanghai 200438, China

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Understanding the conformational changes during hIAPP aggregation in a membrane is essential for identifying its cytotoxic species. However, due to the transient and heterogeneous nature of the hIAPP oligomerization process, it is still difficult to experimentally characterize the structural distribution of hIAPP intermediates in a membrane environment. On the one hand, although the presence of a membrane accelerates the hIAPP fibrillation,^{22–27} previous experiments also indicated the stabilization of non-amyloid α -helical structures as a key intermediate in membranes containing both anionic and zwitterionic lipids,^{28–30} and showed the coil-to- α -to- β structural transitions during the hIAPP fibrillation in the negatively charged DPPG membrane.³¹ On the other hand, sum frequency generation IR spectroscopy indicated a direct transition from coil to β structures of the hIAPP in the DPPG membrane without involving α -helical intermediates.³² Thus, the role of α -helical intermediates still remains controversial.^{33,34} Several experiments have utilized nano-discs to capture the hIAPP intermediates on membranes and obtained their residue-based secondary structures, *e.g.*, a NMR spectroscopy experiment revealed a three-stranded β -sheet structure of the hIAPP on the membrane of mixed zwitterionic DMPC and anionic DMPG lipids.³⁵ Recently, an isotope-labelled 2D-IR experiment indicated that the hIAPP oligomer contains α -helices in fragment L₁₂A₁₃ and β -sheets in fragment G₂₄A₂₅ on neutrally charged membranes.³⁶ However, the relations between the hIAPP oligomer structural properties and the corresponding membrane disruption effects still remain unstudied.

Computational studies have shed light on the structural properties of the membrane-bound hIAPP in different aggregation stages.^{21,37,38} As for monomers, in agreement with the EPR experiment of the hIAPP on anionic POPS vesicles,³⁹ several computational studies reported that the hIAPP monomer forms an amphiphilic α -helix parallel to the membrane surface at its N-terminal region, while its amyloid core and C-terminal regions remain flexible.^{40–42} During the aggregation stage, a previous simulation study utilized a highly mobile membrane mimetic model to accelerate the hIAPP inter-chain diffusion, and captured the hIAPP α -helix to β -sheet transitions on the membrane surface.⁴³ However, as the lipid mimetic model simplifies the lipid tail groups, this study may not be able to reflect the real hIAPP membrane insertion process.⁴³ The stability of membrane-bound α -helical and β -sheet structures has also been investigated *via* simulations. On the one hand, starting from the α -helical hIAPP structures on SDS micelles resolved by NMR,⁴⁴ a previous simulation study investigated the inter-chain interactions of the hIAPP dimer with the N-terminal region pre-inserted into the membrane,⁴⁵ and coarse-grained simulations revealed that the α -helical hIAPP monomers aggregate and form ion channels, starting from either pre-inserted or surface-bound conformations.^{46,47} On the other hand, several computational studies focused on revealing the different insertion orientations of β -sheet-rich protofibrils with the change of pH and membrane compositions,^{48–51} and the structural model of a β barrel for hIAPP ion channels has also been proposed.⁵² A recent

computational study created various structural models for hIAPP ion channels, and concluded that both helical-bundle and β -barrel ion channels are meta-stable after examining their structural stability *via* all-atom simulations.⁵³ In spite of the continuous efforts, there still lacks an understanding of molecular mechanisms during the hIAPP initial spontaneous aggregation process in the membrane environment.

In this study, we performed extensive (30- μ s long) all-atom molecular dynamics simulations to investigate the hIAPP dimerization process in the anionic POPG membrane. Dimer, as the smallest oligomer, provides the first and best opportunity to study the monomer–monomer interface that drives the initial aggregation process and experiments have indicated that the hIAPP dimer is the most probable oligomer at low concentration.^{54,55} Starting from the membrane-bound α -helical hIAPP monomer structures identified in our previous simulations,⁴¹ we observed a substantial increase of β structures and an increase of helical structures of the hIAPP upon its spontaneous dimerization in the POPG membrane, although the random coiled structures are still dominant. The coexistence of α and β structures constitutes a diverse structural ensemble of the membrane-bound hIAPP dimer. Meanwhile, the random coiled and α -helical dimer structures insert deep into the membrane interior, while the β -sheet-rich dimer structures reside on the membrane surface with strong inter-chain hydrophobic interactions. Moreover, from β -sheet to α -helical dimer structures, the degree of hIAPP-induced membrane disruption also increases gradually. In detail, the membrane deformation induced by β -sheet-rich structures on the membrane surface is almost negligible, while the deep insertion of random coiled and α -helical dimers induces global membrane disruption including membrane thinning and stretching. We speculate that the insertion of random coiled and α -helical structures contributes to the initial stage of membrane damage during hIAPP oligomerization, while β -sheet structures on the surface may be more involved in the later membrane disruption stage during hIAPP fibrillation. This study reveals the complexity of the membrane-bound hIAPP dimer structural ensemble, and sheds light on the molecular mechanisms of hIAPP-induced membrane disruption from a structural perspective.

2. Methods

2.1. Simulation details

We constructed a dimerization system from the α -helical structures identified in our previous simulations⁴¹ of hIAPP monomers in the POPG membrane. To shorten the free diffusion between monomers, we elevated the hIAPP concentration during the system construction, and then performed extensive simulations (accumulated to 30- μ s long) on its dimerization process in the membrane. The details of simulations are as follows.

2.1.1. Force-fields and the hIAPP sequence. Consistent with our previous study, we adopted the Amber99sb-ILDN force



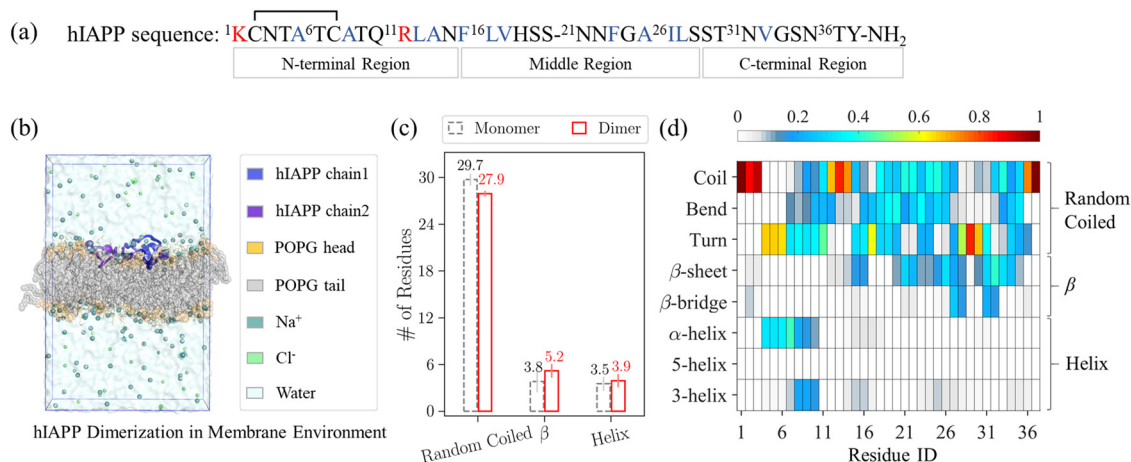


Fig. 1 (a) hIAPP amino-acid sequence, which is divided into the N-terminal region (K₁–N₁₄), the middle region (F₁₅–S₂₈), and the C-terminal region (S₂₉–Y₃₇). (b) Representative snapshot of the simulation system. (c) Secondary structural comparison between the membrane-bounded hIAPP monomer and dimer. The average number of residues adopting each category of secondary structures is plotted. Random coiled category includes coil, bend, and turn structures. β category includes β -sheet and β -bridge structures. Helix category includes α -helix, 5-helix, and 3-helix structures. The gray dashed bars are the monomer and red solid bars are the dimer, with the error bars indicating the standard deviations. (d) Probability of each residue to adopt each type of secondary structures of the hIAPP dimer in the membrane.

field⁵⁶ for the hIAPP, the Amber Slipid force field^{57,58} for the POPG bilayer, and the TIP3P model⁵⁹ for water. The N-terminus of the hIAPP was protonated and residue His₁₈ was deprotonated, as under the neutral pH condition. Consistent with its *in vivo* form, the cysteine residues C₂ and C₇ of the hIAPP are linked with a disulfide bond, and its C-terminus is amidated. The sequence of the hIAPP is shown in Fig. 1(a).

2.1.2. Construction and simulations of the dimer system.

We selected the last frames of the α -helical-rich trajectories of the hIAPP monomer in the POPG bilayer as the initial conformations of dimerization. For each conformation, we first decreased the simulation box-size in the *xy*-plane, by discarding the POPG lipids whose maximal value distances, *i.e.*, the Chebyshev distance,⁶⁰ to the hIAPP center of mass (COM) were larger than 2.4 nm in the *xy*-plane. To ensure the POPG bilayer is symmetric, we then kept 30 lipids closest to the COM of the hIAPP in each membrane leaflet, and the simulation box-size became around $4.7 \times 4.7 \times 14.5$ nm³. We solvated the simulation box, deleted the water molecules inside the membrane bilayer *via* the 'water_deletor.pl' (mdtutorials.com/gmx/membrane_protein/Files/water_deletor.pl) script, and added counter ions to keep the system charge neutral. We also added additional ~ 100 mM NaCl to mimic physiological conditions. The resulting reduced simulation box contains 1 hIAPP protein, 60 POPG lipids (30 lipids in each leaflet), 7138 water molecules, 72 Na⁺, and 15 Cl[−].

We then equilibrated the reduced simulation system. To avoid periodic crash, we first slightly increased both the *x* and *y* box lengths to 5.0 nm, and then performed energy minimization with the steepest method. The minimized system was then further equilibrated with a 200-ps long position restraint followed by a 200-ps long NPT equilibrium. In both the position restraint and NPT equilibrium, we utilized the Berendsen method⁶¹ for the thermodynamic coupling. The temperature

coupling at 310 K was applied separately on three groups, *i.e.*, hIAPP, POPG, and the rest of water and ions, with the coupling constant at 0.5 ps. The pressure coupling at 1 bar was semi-isotropic, in which the pressures of the *xy*-plane and the *z*-direction were coupled separately. The pressure coupling constant was 10 ps and the compressibility coefficient was 4.5×10^{-5} bar^{−1}. The electrostatic interactions were calculated *via* PME method,⁶² with real space cut-off at 1.6 nm. vDW interactions were switched to zero from 1.4 nm to 1.5 nm. hIAPP and POPG lipids were constrained with the LINCS algorithm,⁶³ and water molecules with the SETTLE algorithm.⁶⁴ After the equilibration, the size of the simulation box became around $4.7 \times 4.7 \times 13.4$ nm³.

We then duplicated the monomer conformations in either the *x* or *y* direction to create the dimer simulation system, as shown in Fig. S1 (ESI[†]). The dimer system contains 2 hIAPP, 120 POPG lipids (60 lipids in each leaflet), 14 276 water molecules, 144 Na⁺, and 30 Cl[−]. In the 6 initial conformations, the minimal distances between two monomers were larger than 1.8 nm. For each initial state, we performed one 5- μ s long NPT production run, and observed the spontaneous dimerization process as shown by the time evolution of minimal distances between two hIAPP chains in Fig. S2 (ESI[†]). In the production simulations, we utilized the Nose–Hoover temperature coupling⁶⁵ and the Parrinello–Rahman pressure coupling⁶⁶ for the thermodynamic couplings.

2.2. Analysis methods

2.2.1. Conformational ensemble of the dimerized hIAPP in the POPG membrane.

A hIAPP dimer complex is formed if there are inter-chain contacts, and an inter-chain contact is defined when the minimal distance between two chains is less than 0.45 nm. As shown in Fig. S2 (ESI[†]), the initially isolated monomers spontaneously aggregate to form a dimer, with the



increase of the number of inter-chain contacts. For the dimerized hIAPP, we calculated its secondary structures using the DSSP algorithm,⁶⁷ and further grouped the 8 types of secondary structures into 3 categories, *i.e.*, random coiled, β , and helix. In detail, the random coiled category includes coil, bend, and turn structures, the β category includes β -sheet and β -bridge structures, and the helix category includes α -helix, 5-helix, and 3-helix structures. To examine the contribution of different types of inter-chain interactions, we further calculated the residue-based inter-chain hydrogen bonds (H-bonds) and hydrophobic contacts. In detail, the H-bonds were calculated *via* the MDTraj package⁶⁸ with the 'baker_hubbard' criterion,⁶⁹ and the hydrophobic contact number was calculated as the number of contacts between the side-chains of hydrophobic residues. Besides, we also calculated the number of contacts between the hIAPP dimer and the POPG tail groups, to examine the insertion depth of the hIAPP dimer.

2.2.2. Clustering of the dimer complex. To illustrate the hIAPP conformational ensemble during the dimerization, we performed a two-step clustering to dissect the influence of single-chain conformations and inter-molecular interactions separately. The first step clustering is based on single-chain RMSD and the second step is based on protein-protein and protein-membrane contact patterns, as shown in the flowchart of Fig. S3 (ESI[†]). The detailed procedures of clustering are as follows.

First, we conducted the conformational clustering based on the single-chain RMSD *via* Daura's method,^{70,71} with an RMSD cut-off of 0.4 nm (step 1 in Fig. S3, ESI[†]), *i.e.*, the maximal RMSD to the cluster centre is less than the cut-off 0.4 nm within each cluster. Each frame is labelled by the cluster indices of its two hIAPP chains, and forms a pair of single-chain clusters. Depending on whether there are inter-chain contacts or not, the conformations of single-chain clusters are further divided into the dimerized part and the isolated part. The population distribution of single-chain clustering is shown in Fig. S4 (ESI[†]). We selected the top 12 clusters of the dimerized part with the population cut-off at 2.0%, as shown in the inset of Fig. S4 (ESI[†]), and their accumulating population reaches 85.9%. In comparison with the monomeric hIAPP in the membrane, we further calculated the single-chain structural properties in the dimerized part, such as the radius of gyration (R_g) and the number of contacts with POPG tail groups, to investigate the changes in the chain dimension and insertion depth upon the dimerization. Moreover, we also calculated the inter-chain H-bonds and the inter-chain hydrophobic contacts, to examine the contributions of different inter-chain interactions to the corresponding structural changes.

Second, we examined the pairs of single-chain clusters, and divided the dimerized pairs of single-chain clusters by performing additional *k*-means clustering based on the protein-protein and protein-membrane contact maps (step 2 in Fig. S3, ESI[†]). We selected the top 13 dimerized pairs with population cut-off at 2.0% for the clustering in the step 2. As shown in Fig. S5 (ESI[†]), the accumulated population of the top 13 dimerized pairs reaches 80.2%. The details of step 2 clustering are as

follows. For each frame, we first concatenated the residue-based inter-chain contact maps (37×37 elements) and residue-based contact maps with membrane tail groups (37×2 elements), and then performed the *k*-means clustering based on the concatenated vectors ($37 \times (37 + 2)$ elements) describing the inter-molecular contact patterns of the dimerized complex. The optimal number of clusters k^* was determined *via* the elbow method, *i.e.*, the *k* value at which the 2nd derivative of sum of squared distances (SSD) reaches the maximum. In the case that the SSDs were smaller than 10^4 and the 2nd derivatives of SSDs were smaller than 10^3 , we chose $k^* = 1$ as the optimal number of clusters, as shown in Fig. S6 (ESI[†]). This step 2 clustering of dimer pairs results in 22 dimer clusters, and the population distribution of these 22 clusters is shown in Fig. S7 (ESI[†]). In this way, each resulting dimer cluster corresponds to one specific type of single-chain conformations with one particular inter-molecular contact pattern. For each dimer cluster, we also calculated its structural properties, such as the average number of inter-chain contacts, the average number of tail contacts, the residue-based inter-chain H-bonds, the residue-based inter-chain hydrophobic contacts, and the residue-based tail contacts. To elucidate the hIAPP conformational changes upon dimerization, we tracked the time evolution of the 22 dimer clusters in each trajectory. The transitions among these clusters are shown in Fig. S8 (ESI[†]). We also labelled the inter-chain contacts of the initially collapsed dimer clusters in each trajectory, and plotted the initially collapsed contact map in Fig. S9 (ESI[†]).

2.2.3. Density distribution along the z-axis of the simulation box. We calculated the density distribution of the hIAPP, POPG head groups, POPG tail groups, and the water solvent along the z-axis of the simulation box, to illustrate the positions of the hIAPP dimer relative to the POPG membrane. Due to the size fluctuation of the simulation box under pressure coupling, we first shifted the z-coordinate of COM of the POPG bilayer to the centre of the simulation box before the density calculations.

2.2.4. Membrane deformation. We further calculated the membrane thickness, area per lipid, membrane curvature, tail order parameters S_{CD} s, and sizes of hydrophobic defects to characterize the membrane deformation induced by the hIAPP dimer. The thickness of the membrane was calculated according to ref. 72. In short, we discretized the *xy*-plane with the spacing of grids at around 1 nm. The z-coordinates of upper and lower surfaces were estimated from the z-coordinates of phosphorus atoms in the corresponding leaflet. Instead of using the nearest phosphorus atom,⁷² we estimated the z-coordinate of a grid point by averaging the phosphorus atoms whose Euclidean distances to the grid point are within 1.5 nm in the *xy*-plane, and thus smoothed the surface representation. The thickness of the bilayer was then estimated as the vertical distance between the grids of upper and lower surfaces. We estimated the thickness of each leaflet in the same way, where we used the phosphorus atoms and the last carbon atoms in the unsaturated acyl chains sn-2 to construct the z-coordinates of surfaces in each leaflet. The mean curvature of the membrane surface was also estimated *via* the derivatives of



the z-coordinates of grid points. The area per lipid was calculated by dividing the area of the xy-plane of the simulation box by the number of lipids in each leaflet. The order parameters S_{CD} s of the POPG tail groups were calculated according to ref. 73 for both the saturated acyl chain sn-1 and the unsaturated acyl chain sn-2. We calculated the packing defects of POPG, where the aliphatic carbons in tail groups are exposed, *via* the PackMem program,⁷⁴ and the xy-plane was discretized at $0.1 \times 0.1 \text{ nm}^2$ resolution. A defect is defined if there is no atom of the head group within the threshold 0.1 nm of the surface formed by the carbonyl C2 atoms (sp^3 carbon with one hydrogen) vertically. A defect is further divided into the deep part and the shallow part. The deep part corresponds to the area where there are neither head nor tail atoms within the threshold 0.1 nm, while the shallow part corresponds to the area where there are only tail atoms within the threshold 0.1 nm. When its size exceeds 99% of the defects in the pure POPG system, we define the defect as a large-size one. The large-size cutoff is 0.29 nm^2 for the deep defects, 0.32 nm^2 for the shallow defects, and 0.45 nm^2 for the overall (deep part plus shallow part) defects.

3. Results and discussion

3.1. Spontaneous hIAPP dimerization in the POPG membrane

We observed the spontaneous dimerization of the hIAPP in the POPG membrane bilayer in our simulations. As shown in Fig. S2 (ESI[†]), in all the six 5- μs long trajectories, the minimal distances between two hIAPP chains decrease with the increase of inter-chain contacts, and later remain at $\sim 0.2 \text{ nm}$. Fig. 1(b) shows the representative conformation of the last frame in the 1st trajectory.

3.2. Structural properties of the membrane-bound hIAPP dimer

3.2.1. hIAPP becomes more structured upon dimerization.

As shown in Fig. 1(c), the random coiled structures in the dimerized hIAPP decrease compared to those in the monomer, although they are still dominant, occupying 27.9 out of the 37 residues. Meanwhile, we observed a substantial increase of β structures and a slight increase of helical structures. From the residue-based secondary structure distribution shown in Fig. 1(d), the β structures are located at the middle and the C-terminal regions, while the helical structures are at the N-terminal region. In detail, the fragments $F_{15}L_{16}$, $S_{20}S_{28}$, and $T_{30}N_{35}$ have considerable probability to form β -sheets, and the N-terminal fragment T_4Q_{10} forms α -helices. The emergence of β -sheet structures is consistent with previous experiments that studied the hIAPP oligomerization process in the membrane environment, *e.g.*, previous NMR experiment of the hIAPP on nano-discs of mixed zwitterionic DMPC and anionic DMPG lipids reveals the three-stranded β -sheets at fragments A_8L_{12} , $F_{15}H_{18}$, and $I_{26}S_{29}$,³⁵ and a recent 2D IR experiment indicates that $G_{24}A_{25}$ forms β -sheets during oligomerization on zwitterionic DOPC vesicles.³⁶

3.2.2. Diverse structural ensemble of the dimerized hIAPP.

To illustrate its structural ensemble, we further projected the free energy landscape of the dimerized hIAPP on the number of inter-chain contacts (x-axis) and the number of contacts with tail groups (y-axis), as shown in Fig. 2(a). In addition, the distribution of dimer clusters (black circles in Fig. 2(a)) covers the free-energy minima, and the structural ensemble of the membrane-bound hIAPP dimer is heterogeneous. We further divided the 22 dimer clusters into 4 groups, based on the inter-molecular contact patterns. The 4 groups range from group 1 (G1, shaded in purple) of deep insertion and relatively few inter-chain contacts to group 4 (G4, shaded in orange) located on the membrane surface with a relatively high number of inter-chain contacts, while the G2 (shaded in blue) and G3 (shaded in green) lie in between. From the negative correlation between the number of inter-chain contacts and insertion depth, we believe that hIAPP-hIAPP interactions and hIAPP-membrane interactions are competing with each other. This observation is consistent with the previous experiments indicating that direct inter-chain interactions are not necessary for hIAPP poration.^{18,29,75} We further calculated the residue-based inter-chain H-bonds and inter-chain non-polar, *i.e.*, hydrophobic, contacts in each dimer group as shown in Fig. 2(b), and also the residue-based hIAPP contact number with POPG tail groups in each dimer group (Fig. 2(c)), to investigate the interactions between the hIAPP dimer and the POPG membrane. The secondary structure distribution and the centre conformations in each dimer group are also shown in Fig. 2(d-g). By examining the inter-molecular interactions and secondary structure distribution of each group, we outline the feature of each group in detail as follows.

G1: deep insertion with α -helical structures at T_4T_9 and $F_{15}V_{17}$ fragments. Among the 4 groups, G1 contains the fewest inter-chain contacts and inserts deepest. As shown in the purple circles of Fig. 2(b), some inter-chain contacts are formed in the C-terminal region of the hIAPP in G1, *i.e.*, the H-bond between S_{29} and S_{29} and the hydrophobic contacts between V_{32} and V_{32} . Meanwhile, we observed the insertion of the hIAPP throughout the whole sequence in G1. As shown in Fig. 2(c), not only the N-terminal amphiphilic fragment T_4T_9 and the hydrophobic fragment $L_{16}V_{17}$, but also the residue I_{26} at the amyloid core region and Y_{37} at the C-terminus, form a large number of contacts with POPG tail groups. As for the secondary structures (Fig. 2(d)), G1 contains the most abundant α -helical contents and the least β -sheet contents among the 4 groups. In detail, both T_4T_9 and $F_{15}V_{17}$ fragments have substantial α -helical contents, as shown in Fig. 2(d).

G2: α -helical fragment T_5R_{11} insertion and β -hairpin structures on the membrane surface. As for G2, its contact patterns with membrane tail groups are similar to those of G1, as its N-terminal α -helical-rich amphiphilic segment T_5R_{11} also inserts into membrane tail groups. However, the hydrophobic fragment $L_{16}V_{17}$ is on the membrane surface, as shown in Fig. 2(c). In contrast to G1, G2 contains substantial β -sheet structures,



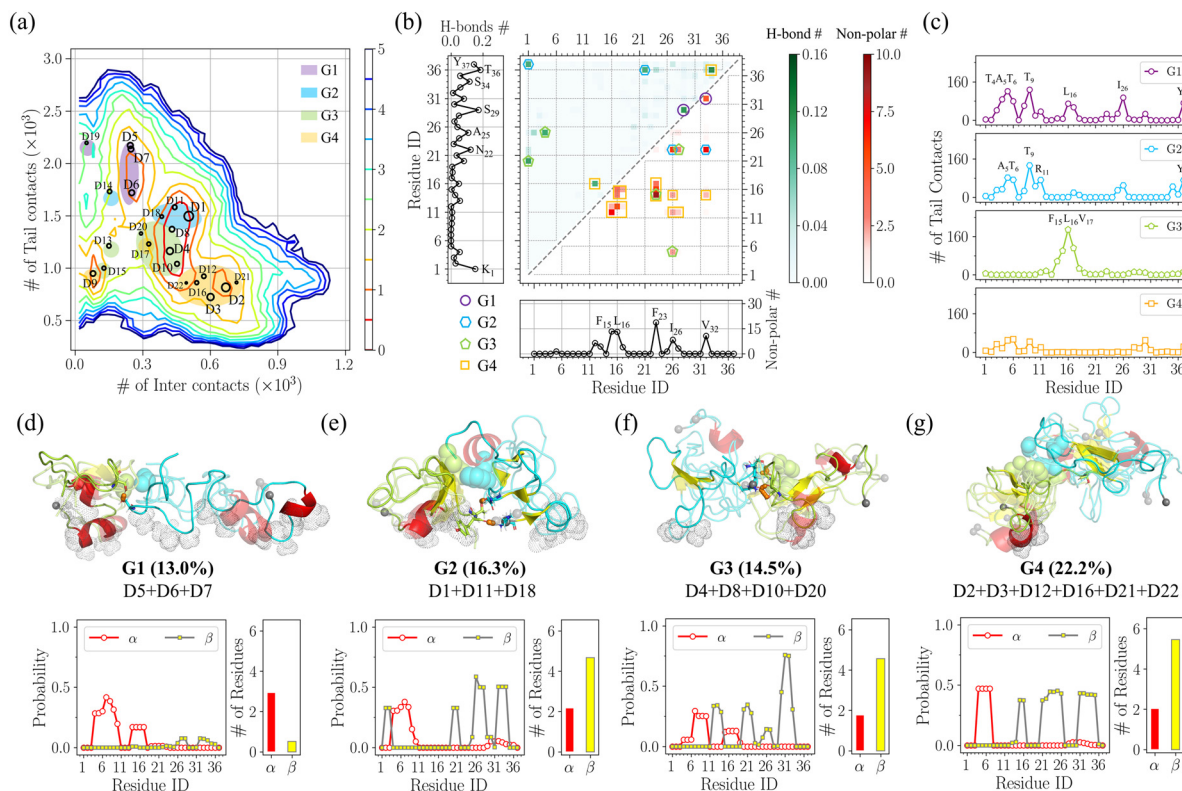


Fig. 2 Structural properties of the hIAPP dimer in the membrane. (a) 2-d free energy projection on the number of inter-chain contacts (x-axis) and the number of contacts with membrane tail groups (y-axis). The black circles correspond to the properties of the top 22 dimer clusters of the hIAPP in the membrane, and the size of each circle is proportional to the corresponding dimer cluster population. The dimer clusters are further grouped into 4 categories: group 1 (G1) in purple, G2 in blue, G3 in green, and G4 in orange. (b) Residue-based inter-chain interactions. The bottom right is the number of inter-chain non-polar, i.e., hydrophobic, contacts and the upper left is the number of inter-chain H-bonds, with the corresponding 1-d distribution shown by side. The contacts of different groups are also labelled with different markers: purple circles for G1, blue hexagons for G2, green pentagons for G3, and orange squares for G4. (c) Residue-based contact number of the hIAPP dimer with membrane tail groups in each group. (d)–(g) Representative conformations and residue-based secondary structural distribution of each group, from G1 in (d) to G4 in (g). In each group, the centre conformations of dimer clusters are aligned. The hIAPP dimers are shown in cartoon representation, with one chain in cyan, and the other chain in light-green. The secondary structure α -helices are in red and β -sheets are in yellow. N-termini (C α atoms of K₁ residues) of chains are in gray sphere. Sidechains of residues forming inter-chain hydrophobic contacts are shown in sphere, and residues forming inter-chain H-bonds are shown in stick with H-bonds shown as orange dashes. Sidechains of deeply inserted residues are shown in gray dots. The population of each group and its cluster IDs (excluding the initially collapsed clusters) are also labelled at the bottom. As for the secondary structure distribution in each group, red circles are the distribution of α -helices and yellow squares are that of β -sheets. Residue numbers adopting α -helices or β -sheets are shown by side as red or yellow bars.

and the β -sheet structures are located at N-terminal (C₂N₃), middle (S₂₀N₂₁, I₂₆–S₂₈), and C-terminal (V₃₂–S₃₄) regions, as shown in Fig. 2(e). In addition, these β -sheet fragments in G2 are on the surface of the membrane. Both H-bonds and hydrophobic contacts contribute to the inter-chain interactions in G2, as shown in the blue hexagons of Fig. 2(b). Moreover, these inter-chain contacts are adjacent to the β -sheet-rich regions, e.g., the H-bond between K₁ and Y₃₇ is close to the β -fragments C₂N₃ in one chain and V₃₂–S₃₄ in the other chain, while the hydrophobic contacts between F₂₃ and I₂₆ are close to the β -fragments S₂₀N₂₁ in one chain and I₂₆–S₂₈ in the other chain.

G3: hydrophobic fragment F₁₅–V₁₇ insertion and β -hairpin structures on the membrane surface. In G3, the hydrophobic fragment F₁₅–V₁₇ inserts deep with random coiled and α -helical structures. Different from G2, although it contains notable α -helical structures, the N-terminal fragment C₇–T₁₀ is located on the membrane surface without deep insertion.

Instead, the N-terminal residues (K₁, T₄, and A₅) form inter-chain interactions with the residues (S₂₀, A₂₅, and I₂₆) in the middle region, e.g., there are H-bonds formed by K₁/T₄ with S₂₀/A₂₅, and also hydrophobic contacts between A₅ and I₂₆. Moreover, both chains form intra-chain β -hairpin structures, and these β -hairpins are connecting fragments L₁₂–N₁₄ with T₃₀–V₃₂, S₂₀–N₂₂ with N₃₁–G₃₃, and A₂₅–L₂₇ with T₃₀–V₃₂. Consistent with G2, these β -sheet-rich structures in G3 are on the membrane surface, and also participate in the inter-chain interactions.

G4: β -sheet-rich structures on the membrane surface with strong inter-chain hydrophobic interactions. Opposite to G1, G4 contains the most abundant inter-chain contacts and has the shallowest insertion depth among the 4 groups. In G4, there are multiple inter-chain hydrophobic contacts among several fragments, as shown in the orange squares in Fig. 2(b). In detail, the hydrophobic contacts are between L₁₂A₁₃ and F₁₅–V₁₇, L₁₆V₁₇ and

L₁₆V₁₇, L₁₂ and A₂₆I₂₇, F₂₃ and F₁₅-V₁₇, F₁₅ and I₂₆, and F₁₅ and V₃₂, covering all the hydrophobic residues of the hIAPP, except A₅ and A₈ at the amphiphilic N-terminal region. In addition, there are also H-bonds between A₁₃/S₃₄ and L₁₆/T₃₆. As shown in Fig. 2(c), there are only a small number of tail contacts throughout the hIAPP sequence, indicating that hIAPP dimer structures in G4 are mainly located on the membrane surface. Moreover, G4 contains the highest β -sheet contents among the 4 groups, and its β -sheet structures are located at F₁₅L₁₆, N₂₁-I₂₆, and N₃₁-N₃₅ fragments, as shown in Fig. 2(g). These β -sheet-rich regions coincide with the inter-chain contacts, indicating that the β -sheet structures in the dimerized hIAPP correlate well with the formation of inter-chain interactions. Besides, the N-terminal T₄-C₇ segment contains notable α -helix contents with a few tail contacts.

Moreover, the β -sheet-rich structures on the membrane surface are more stable than the random coiled and α -helical structures deeply inserted into membrane interior. To compare their structural stability, we estimated the population of the 4 groups. For each group, we summed up the population of its dimer clusters, with the initially collapsed clusters excluded. As shown in Fig. 2(d-g), among the 4 groups, the deeply inserted G1 is the least populated (13.0%) while β -sheet-rich G4 on the membrane surface is the most populated (22.2%), with the G2 (16.3%) and G3 (14.5%) in between. The higher metastability of G4 is probably due to two factors: one is the strong inter-chain hydrophobic interactions in G4, and the other is the less severe membrane deformation it induces due to its shallow insertion. In comparison, the deep insertion of G1 leads to fewer inter-chain contacts and also more severe membrane deformation. The membrane deformation induced by different groups will be discussed in detail in the Section 3.3.

To elucidate the conformational rearrangements upon dimerization, we plotted the time evolution of the dimer clusters along each trajectory as shown in Fig. S8 (ESI†). In short, we observed the deepening of membrane insertion in G1 and also the elongation of β -hairpin structures on the membrane surface in G4. Besides, the initial inter-chain contacts contain both the hydrophobic and polar interactions, mostly involving the hydrophobic residues F₂₃, V₃₂, and the positively charged residues K₁, R₁₁, as shown in Fig. S9 (ESI†).

3.2.3. Emergence of single-chain structures distinct from the monomer ensemble. To demonstrate the structural changes induced by the hIAPP dimerization, we further examined the dimension and the insertion depth of the single chains in the dimerized hIAPP structural ensemble, by projecting its free energy landscape onto the radius of gyration (R_g) and the number of contacts with membrane tail groups, as shown in Fig. 3(a), in comparison with the properties of monomeric hIAPP in the membrane in our previous study.⁴¹

Dimerization induces more expanded single-chain conformations. As shown in the bottom right subplot of Fig. 3(a), compared to the hIAPP monomer (gray dotted line), the single-chain R_g distribution of the dimerized hIAPP shifts to larger values (red solid line). In detail, the major peak at

$R_g \sim 1.01$ nm in the monomer shifts to $R_g \sim 1.07$ nm in the dimerized hIAPP, while both the monomer and dimer contain the peaks at $R_g \sim 0.93$ nm. In particular, the R_g s of single-chain clusters C4 and C8 are significantly larger than the rest clusters, *i.e.*, the R_g of C4 in the β -sheet-rich G4 is ~ 1.15 nm, and the R_g of C8 in the deeply inserted G1 is ~ 1.30 nm. The dimension expansions of C4 and C8 are caused by different factors. In detail, the dimension expansion of C4 is due to the strong hydrophobic inter-chain interactions as shown in Fig. 3(b), while the dimension expansion of C8 is due to its deep insertion as reflected by its large number of tail contacts (Fig. 3(a)).

Dimerization induces wider distribution of the single-chain insertion depth. For the insertion depth, the dimerized hIAPP adopts a wider distribution of the number of tail contacts, as shown in the upper left subplot of Fig. 3(a). In contrast to the single high peak at $\sim 0.62 \times 10^3$ in the monomer ensemble, there are several peaks in the dimerized hIAPP, *i.e.*, one peak at $\sim 0.05 \times 10^3$ corresponding to the hIAPP on the membrane surface (single-chain cluster C4 in G4), and two flat peaks at $\sim 0.54 \times 10^3$ and $\sim 0.90 \times 10^3$, containing the majority of the single-chain clusters in the dimer. Moreover, the single-chain cluster C8 in G1 contains an exceptionally large number of tail contacts ($\sim 1.43 \times 10^3$), deeply inserted into the membrane interior.

Last but not least, the interactions between single-chain clusters also reflect the feature of each group of dimer clusters. In detail, the single-chain clusters in G1 (C1, C5, and C8, shaded in purple) contain the fewest inter-chain contacts among the 4 groups, while the single-chain clusters of G4 (C2, C4, and C9, shaded in orange) contain a considerable large number of hydrophobic contacts. Besides, the single-chain clusters in both G2 and G3 contain a mixture of inter-chain H-bonds and hydrophobic interactions.

From the analyses above, we conclude that the inter-chain hydrophobic interactions in dimerized hIAPP facilitate the β -sheet formation on the membrane surface, while coiled and α -helical structures of hIAPP are responsible for the membrane deep insertion. Moreover, the inter-chain interactions are mainly formed at the middle and C-terminal regions, while the insertion of α -helical and coiled structures is located at the N-terminal and middle regions. Previous NMR experiment observed that the hIAPP₁₋₁₉ fragment alone causes membrane disruption to the identical extent as the full-length hIAPP, but does not form amyloid fibrils, and concluded that hIAPP membrane disruption and fibrillation are two separate processes.⁷⁶ Meanwhile, the rat IAPP (rIAPP), which differs from the hIAPP by 6 residues at middle and C-terminal regions, is able to bind to the membrane, but is non-cytotoxic and does not form fibrils.^{77,78} As the 3 out of the 6 rIAPP residues different from hIAPP are the secondary structure breaker proline (A₂₅P, S₂₈P, and S₂₉P), the non-fibrillation of rIAPP implies the importance of hIAPP middle and C-terminal regions in forming inter-chain interactions in the membrane environment. Besides, the non-toxic IAPP analogues from cattle and



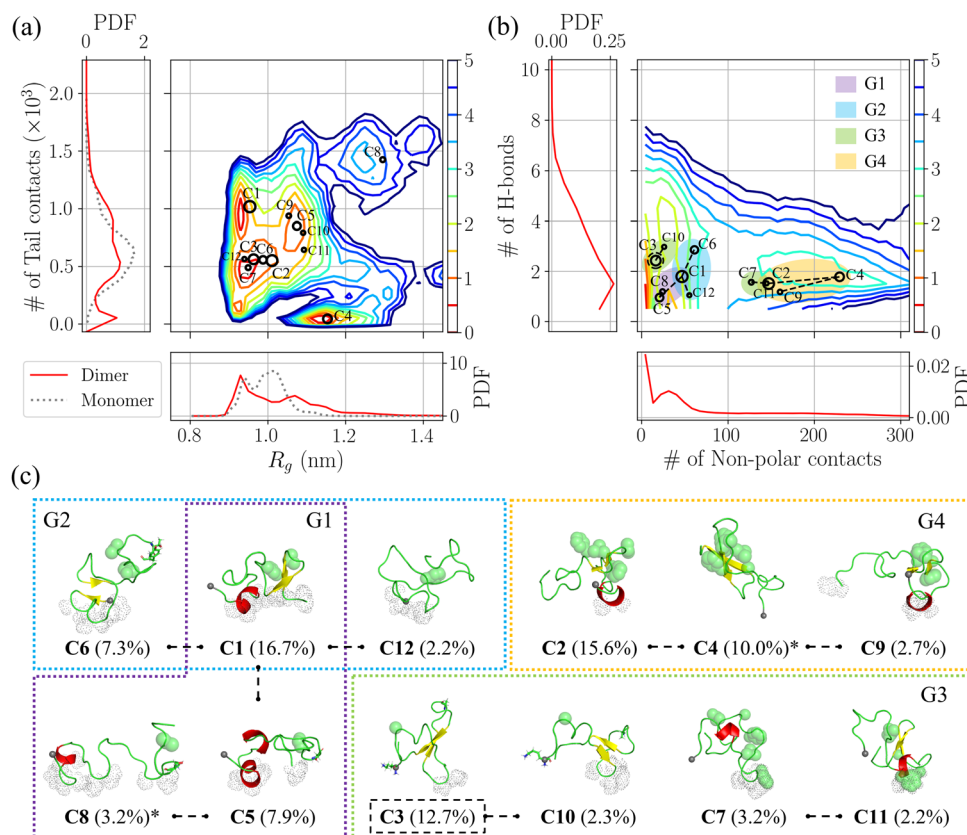


Fig. 3 Single-chain structural properties of the hIAPP dimer in the membrane. (a) 2-d free energy projection on the radius of gyration (R_g) (x-axis) and the number of tail contacts (y-axis) of each chain in the dimer ensemble, with the corresponding 1-d probability distribution shown along each axis (red solid lines), in comparison with the properties of the hIAPP monomer in the membrane (gray dotted lines). (b) 2-d free energy projection on the non-polar contacts (x-axis) and the inter-chain H-bonds (y-axis) formed by each single chain in the dimer ensemble. Similarly, the corresponding 1-d probability distribution is shown along each axis. In both (a) and (b), we labelled the average values of the top 12 single-chain clusters as the black circles in the contour plot. In (b), the groups of clusters are shaded with different colors, i.e., purple for group 1 (G1), blue for G2, green for G3, and orange for G4. The single-chain clusters forming a dimer complex with each other are connected with dashed black lines. The cluster C3 forming a dimer with itself is surrounded by dashed black circles. (c) Center conformations of the top 12 single-chain clusters arranged according to their corresponding structural groups. The population of each cluster is labelled in the bracket. The proteins are shown as cartoon representation, with α -helices in red and β -sheets in yellow. N-termini ($C\alpha$ atoms of K_1 residues) of hIAPP chains are in gray sphere. Sidechains forming inter-chain non-polar interactions are shown in sphere and residues forming inter-chain H-bonds are shown in stick. Sidechains of residues deeply inserted into membrane are shown in gray dot.

bears also contain at least one or two mutations to proline in the middle and the C-terminal regions and do not form amyloid fibrils either.^{79,80} Taken together, consistent with our simulation results, these experimental observations not only indicate the N-terminal to middle regions is responsible for hIAPP membrane binding and insertion, but also highlight the important role of middle and C-terminal regions during the hIAPP fibrillation in the membrane environment. Moreover, the correlation between cytotoxicity and amyloidogenicity among different IAPP isoforms supports the membrane disruption mechanism of the “fibril elongation” hypothesis, and suggests that hIAPP-induced membrane disruption may be a multi-step process.

The coexistence of α -helical and β -sheet structures in the membrane-bound hIAPP dimer indicates the complexity of hIAPP aggregation pathways in the membrane. Previous experiments have also indicated the structural diversity during the hIAPP oligomerization process, with some species toxic and

some species non-toxic.^{7,8} To explore the membrane deformation induced by various hIAPP dimer structures, we further calculated the changes of membrane properties, such as the membrane thickness, area per lipid, membrane curvature, tail ordering parameter, and hydrophobic packing defects, induced by each structural group of the hIAPP dimer in the following sections.

3.3. Different degrees of membrane deformation induced by different hIAPP dimer structures

3.3.1. Strongest membrane deformation in G1 vs. weakest membrane deformation in G4. We further examined the membrane disruption induced by different groups. As shown in Fig. 4(a), the dimer insertion depth ranges from the deepest in G1 to the shallowest in G4. Even in the deepest insertion of G1, the hIAPP dimer inserts into only one leaflet without permeabilizing into the other leaflet, and there is no water channel across the membrane. The asymmetric hIAPP insertion



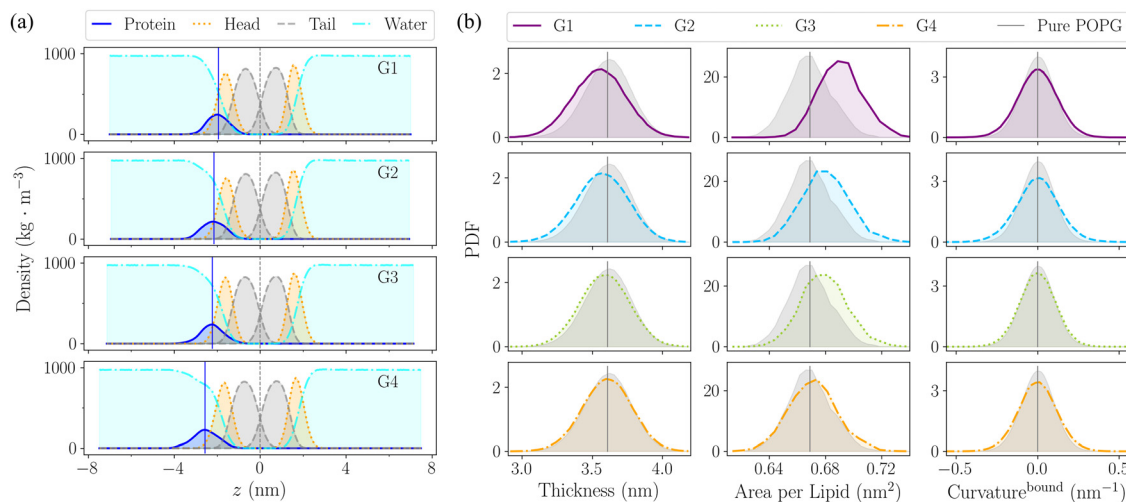


Fig. 4 Comparison of (a) hIAPP position and (b) membrane deformation among different groups of the hIAPP dimer in the membrane environment. (a) Density distribution of each component along the z -axis of the simulation box. Subplots from top to bottom correspond to the distributions from group 1 (G1) to G4. In each subplot, the blue solid curve is the hIAPP, the orange dotted curve is POPG head groups, the dashed gray curve is POPG tail groups, and the cyan dash-dotted curve is water solvent. Vertical blue solid lines indicate the density peaks of the hIAPP, and vertical gray dashed lines indicate the COMs of POPG tail groups. (b) Distribution of the membrane thickness (left), area per lipid in the xy -plane (middle), and the mean curvature (right) of the hIAPP-bound membrane leaflet. In each subplot, the distribution in group 1 (G1) is shown in a purple solid line, G2 in a blue dashed line, G3 in a green dotted line, and G4 in an orange dash-dotted line. The vertical gray lines indicate the average values of the pure POPG membrane, with the gray shaded areas indicating its distributions.

also induces unique membrane deformation. As shown in Fig. 4(b), correlated with the dimer insertion depth, the overall degrees of membrane deformation range from the most severe in G1 to the least significant in G4, with the corresponding details discussed in the following sections.

3.3.2. Membrane thinning and stretching in G1 vs. membrane bending in G2. Among the 4 groups, G1 induces the most severe membrane thinning and stretching, while G2 induces the strongest membrane bending in the hIAPP-bound leaflet. As shown in Fig. 4(b), among all the 4 groups, the distributions of membrane thickness shift to the smaller values, indicating the dimer-induced membrane thinning, while the distributions of area per lipid shift to the larger values, indicating the dimer-induced membrane stretching. From G1 to G4, the extent of membrane thinning and stretching decreases gradually. As for the membrane surface bending, we do not observe the overall bending induced by the hIAPP, and the average values of the mean curvature are nearly zero ($\sim 10^{-5} \text{ nm}^{-1}$). Nonetheless, the mean curvature distribution in the hIAPP-bound leaflet is broader than that of the pure POPG bilayer (right panel in Fig. 4(b)), while the mean curvature distribution in the hIAPP-free leaflet is identical to that of the pure POPG bilayer (Fig. S10, ESI†). Among the 4 groups, the curvature distribution in G2 exhibits the widest variance, indicating that G2 induces the most significant local disturbance of the hIAPP-bound membrane leaflet.

3.3.3. Global membrane disruption in G1 vs. local membrane disruption in G2. Although it only inserts into one of the membrane leaflets, the hIAPP dimer still causes notable disturbing effects on the other hIAPP-free leaflet, especially in G1, as reflected by the tail ordering and the thickness of each

membrane leaflet. As shown in Fig. 5(a), the tail ordering parameters S_{CD} s decrease in not only the hIAPP-bound (hollow markers) leaflet but also the hIAPP-free (solid markers) leaflet. In the hIAPP-free leaflet, G1 reduces the tail ordering and thickness most significantly, indicating that the deep insertion in G1 causes more global membrane disturbance. In comparison, in the hIAPP-bound leaflet, G2 reduces the tail ordering most significantly, consistent with its strong local membrane disturbance effects as reflected by its mean curvature in the hIAPP-bound leaflet. Meanwhile, the tail disordering correlates well with the thickness thinning in the corresponding leaflet, *i.e.*, the greater the disturbance of tail ordering in the leaflet, the smaller the corresponding thickness, as shown in Fig. 5(b).

The local membrane disturbance in G2 is also reflected by its large-size shallow defects in the hIAPP-bound leaflet. As shown in Fig. 5(c), in the hIAPP-bound leaflet, G2 contains significantly more large-size shallow defects ($> 0.32 \text{ nm}^2$) than the other 3 groups. Meanwhile, G1 contains the most large-size deep defects ($> 0.29 \text{ nm}^2$) among the 4 groups (Fig. 5(d)), which is consistent with its deepest insertion. Besides, as for the size distribution of overall defects (shallow defects plus deep defects) shown in Fig. S11(a) (ESI†), although G1 contains the most large-size overall defects ($> 0.45 \text{ nm}^2$), there is a substantial increase of overall defects larger than 1 nm^2 in G2. This increase is mainly due to the large-size shallow defects in G2. Meanwhile, as for the hIAPP-free leaflet, we observed a slight increase of large-size overall defects, and the degrees of size enlargement decrease gradually from G1 to G4, which is consistent with the degree of membrane stretching in each structural group, as shown in Fig. S11(b) (ESI†). In short, both the tail disordering and the size distribution of hydrophobic



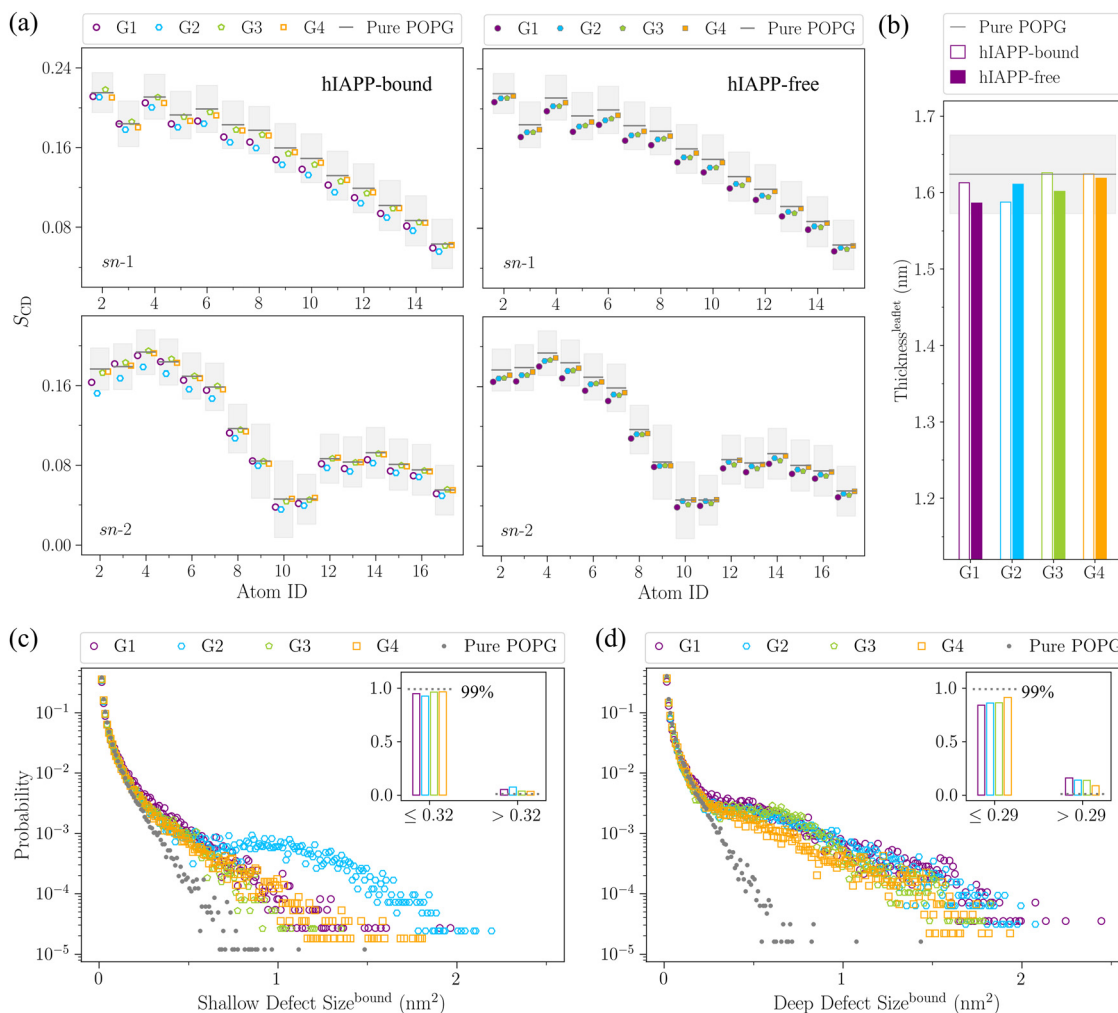


Fig. 5 Comparison of (a) order parameters S_{CD} s of membrane tail groups, (b) thickness of each membrane leaflet, and (c) and (d) distribution of sizes of packing defects in the hIAPP-bound membrane leaflet, among different groups of hIAPP dimer in membrane environment. In (a) and (b), the hollow markers/bars are for the hIAPP-bound membrane leaflet, and the solid markers/bars are for the hIAPP-free leaflet. (c) is of the shallow defects and (d) is of the deep defects. In each plot, the properties of group 1 (G1) are in purple, G2 in blue, G3 in green, and G4 in orange, in comparison with the properties of pure POPG membrane in gray. In (a) and (b), the gray shaded areas indicate the corresponding standard deviations of the pure POPG membrane. In (c) and (d), the inset plots are the distribution of defect sizes in two ranges, one for the normal size (within 99% of the pure POPG), and the other for the large size (exceeding 99% of the pure POPG).

packing defects indicate that the hIAPP deep insertion in G1 causes more global membrane disruption in the hIAPP-free leaflet, while the shallow insertion in G2 causes more local membrane disruption in the hIAPP-bound leaflet.

From the analyses above, we observe that different hIAPP dimer structures cause different degrees of membrane disruption, and we believe that the structural diversity of the membrane-bound hIAPP dimer may also imply the diversity of hIAPP toxic species. Previous experiments have indicated the structural diversity in the oligomerization stage, although the detailed structures of membrane-bound hIAPP oligomers have still been elusive.^{7,8} With computational modelling, both helical-bundle and β -barrel structures of oligomeric hIAPP ion channels have been proposed.^{46,52,53} Here, based on the structural and corresponding membrane disruption properties of the hIAPP dimer ensemble, we speculate that the random

coiled and α -helical structures in G1 may be the precursor of helical-bundle ion channels, while the mixed α - and β -sheet structures in G2 may be the precursor for the β -barrel ion channels.

To sum up, the degrees of membrane disruption decrease gradually from the random coiled and α -helical structures in G1 to the β -sheet-rich structures in G4, well correlated with the hIAPP insertion depth in each structural group, as shown in Fig. 6. In particular, the deep insertion of G1 causes the global membrane disruption, inducing the most significant membrane thinning and stretching, and also creating the large-size deep defects. In comparison, the shallow insertion of G2 with a mixture of α -helical and β -sheet structures causes the local membrane disruption and creates large-size shallow defects, resulting in the most severe membrane bending in the hIAPP-bound leaflet. Experiments have indicated that the hIAPP-induced membrane disruption process is bi-phasic,



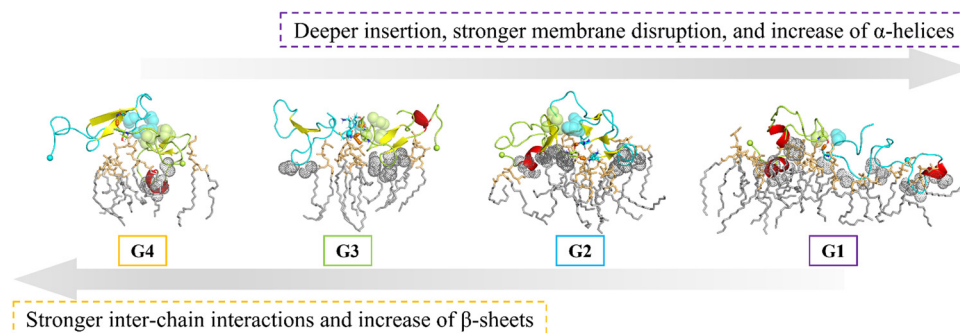


Fig. 6 Representative centre conformation of each group with the nearby POPG lipids, whose contact probabilities with the hIAPP dimer are higher than 95%. hIAPP dimers are shown in cartoon representation, with the α -helix in red and β -sheets in yellow. In each conformation, one chain is in cyan and the other chain is in light green. N-termini (the C α atoms of K1) are shown in a solid sphere. Sidechains forming inter-chain hydrophobic contacts are shown in a sphere, and residues forming inter-chain H-bonds are shown in a stick. Sidechains of deep inserting residues are shown in gray dots. POPG lipids nearby are shown in a stick, with head groups in orange and tail groups in gray.

including the 1st phase leakage caused by the prefibrillar oligomer, and the 2nd phase leakage caused by fibril elongation,^{13,14} and thus both the oligomers in the early nucleation stage and fibrils in the later elongation stage contribute to the hIAPP-induced membrane disruption. Besides, even during the initial oligomer poration stage, previous experiments have observed that the growth of oligomers in the membrane is in a step-wise manner, and there exist two types of oligomers, only one of which is toxic.^{7,8} Though the detailed structures have still been elusive experimentally, both helical bundles and β barrels of hIAPP ion-channels have been proposed in the computational modelling.^{46,52,53} In line with these experimental and computational results, we observed a diverse structural ensemble of the membrane-bound hIAPP dimer in our simulations, and different dimer structures induce the membrane disruption to different extents. Based on our simulations, we speculate that the random coiled and α -helical structures of G1 initialize the membrane disruption and are responsible for the membrane poration, while the β -sheet-rich structures in G4 may be the seed for hIAPP fibrillation in the membrane environment and more involved in the later stage of hIAPP-induced membrane disruption.

4. Conclusions

In this study, we performed extensive all-atom molecular dynamics simulations to investigate the hIAPP dimerization process in the anionic POPG membrane environment. From our simulations, we observed a diverse conformational ensemble of the hIAPP dimer in the POPG membrane environment, and different structural groups induce different degrees of membrane deformation. From G4 to G1, the hIAPP dimer shifts from the β -sheet-rich structures on the membrane surface to the random coiled and α -helical structures inserted into the membrane interior. Strong inter-chain hydrophobic interactions facilitate the β -sheet formation in G4, while only a few inter-chain interactions at the C-terminal region are involved in the random coiled and α -helical structures deeply inserted into the membrane in G1. The extent of hIAPP-induced membrane

disruption also increases from G4 to G1 gradually. In detail, the deep insertion of random coiled and α -helical structures in G1 induces the strongest membrane deformation among the 4 groups, including membrane thinning, membrane stretching, and global tail disordering. Meanwhile, the mixture of α -helical and β -sheet structures in G2 mainly causes membrane bending and local tail disordering. In contrast, the β -sheet-rich structures on the membrane surface in G4 induce negligible membrane deformation. We believe that the deeply inserted random coiled and α -helical dimer structures identified here contribute to the initial stage of membrane disruption, while the β -sheet-rich dimer structures on the membrane surface may be involved in the later stage during the hIAPP fibril elongation. This study reveals the complexity of the membrane-bound hIAPP dimer structural ensemble, providing a detailed structural explanation on the seemingly conflicting mechanisms of hIAPP-induced membrane disruption.

Author contributions

Q. Q. conceived the project, performed the simulations, analyzed the simulation data, and drafted the manuscript. The manuscript was written through contributions of all authors, and all authors have given approval to the final version of the manuscript.

Conflicts of interest

There are no conflicts of interest to declare.

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References

- 1 J. W. M. Höppener, B. Ahrén and C. J. M. Lips, *N. Engl. J. Med.*, 2000, **343**, 411–419.
- 2 S. M. M. Butterfield and H. A. A. Lashuel, *Angew. Chem., Int. Ed.*, 2010, **49**, 5628–5654.
- 3 S. H. Back and R. J. Kaufman, *Annu. Rev. Biochem.*, 2012, **81**, 767.
- 4 L. Haataja, T. Gurlo, C. J. Huang and P. C. Butler, *Endocr. Rev.*, 2008, **29**, 303–316.
- 5 T. A. Mirzabekov, M. Lin and B. L. Kagan, *J. Biol. Chem.*, 1996, **271**, 1988–1992.
- 6 A. Quist, I. Doudevski, H. Lin, R. Azimova, D. Ng, B. Frangione, B. Kagan, J. Ghiso and R. Lal, *Proc. Natl. Acad. Sci. U. S. A.*, 2005, **102**, 10427–10432.
- 7 N. B. Last, E. Rhoades and A. D. Miranker, *Proc. Natl. Acad. Sci. U. S. A.*, 2011, **108**, 9460–9465.
- 8 M. Birol, S. Kumar, E. Rhoades and A. D. Miranker, *Nat. Commun.*, 2018, **9**, 1312.
- 9 Y. A. Domanov and P. K. J. Kinnunen, *J. Mol. Biol.*, 2008, **376**, 42–54.
- 10 C. C. Lee, Y. Sun and H. W. Huang, *Biophys. J.*, 2012, **102**, 1059–1068.
- 11 A. Martel, L. Antony, Y. Gerelli, L. Porcar, A. Fluitt, K. Q. Hoffmann, I. Kiesel, M. Vivaudou, G. Fragneto and J. J. de Pablo, *J. Am. Chem. Soc.*, 2017, **139**, 137–148.
- 12 M. F. M. Engel, L. Khemtémourian, C. C. Kleijer, H. J. D. Meeldijk, J. Jacobs, A. J. Verkleij, B. de Kruijff, J. A. Killian and J. W. M. Hoppener, *Proc. Natl. Acad. Sci. U. S. A.*, 2008, **105**, 6033–6038.
- 13 B. O. W. Elenbaas, S. M. Kremsreiter, L. Khemtémourian, J. A. Killian and T. Sinnige, *BBA Adv.*, 2023, **3**, 100083.
- 14 J. R. Brender, E. L. Lee, K. Hartman, P. T. Wong, A. Ramamoorthy, D. G. Steel and A. Gafni, *Biophys. J.*, 2011, **100**, 685–692.
- 15 P. E. S. Smith, J. R. Brender and A. Ramamoorthy, *J. Am. Chem. Soc.*, 2009, **131**, 4470–4478.
- 16 M. F. M. Sciacca, J. R. Brender, D. K. Lee and A. Ramamoorthy, *Biochemistry*, 2012, **51**, 7676–7684.
- 17 J. R. Brender, S. Salamekh and A. Ramamoorthy, *Acc. Chem. Res.*, 2012, **45**, 454–462.
- 18 N. Bag, A. Ali, V. S. Chauhan, T. Wohland and A. Mishra, *Chem. Commun.*, 2013, **49**, 9155–9157.
- 19 M. F. M. Sciacca, F. Lolicato, G. Di Mauro, D. Milardi, L. D'Urso, C. Satriano, A. Ramamoorthy and C. La Rosa, *Biophys. J.*, 2016, **111**, 140–151.
- 20 M. F. M. Sciacca, C. Tempra, F. Scollo, D. Milardi and C. La Rosa, *Biochim. Biophys. Acta, Biomembr.*, 2018, **1860**, 1625–1638.
- 21 D. Milardi, E. Gazit, S. E. Radford, Y. Xu, R. U. Gallardo, A. Caflisch, G. T. Westermark, P. Westermark, C. La Rosa and A. Ramamoorthy, *Chem. Rev.*, 2021, **121**, 1845–1893.
- 22 A. Relini, O. Cavalleri, R. Rolandi and A. Gliozzi, *Chem. Phys. Lipids*, 2009, **158**, 1–9.
- 23 J. D. Knight and A. D. Miranker, *J. Mol. Biol.*, 2004, **341**, 1175–1187.
- 24 J. D. Knight, J. A. Hebda and A. D. Miranker, *Biochemistry*, 2006, **45**, 9496–9508.
- 25 K. Sasahara, D. Hall and D. Hamada, *Biochemistry*, 2010, **49**, 3040–3048.
- 26 L. Caillon, O. Lequin and L. Khemtémourian, *Biochim. Biophys. Acta, Biomembr.*, 2013, **1828**, 2091–2098.
- 27 X. Zhang, J. R. S. Clair, E. London and D. P. Raleigh, *Biochemistry*, 2017, **56**, 376–390.
- 28 J. A. Williamson, J. P. Loria and A. D. Miranker, *J. Mol. Biol.*, 2009, **393**, 383–396.
- 29 N. B. Last and A. D. Miranker, *Proc. Natl. Acad. Sci. U. S. A.*, 2013, **110**, 6382–6387.
- 30 A. Rawat, B. K. Maity, B. Chandra and S. Maiti, *Biochim. Biophys. Acta, Biomembr.*, 2018, **1860**, 1734–1740.
- 31 L. Fu, G. Ma and E. C. Y. Yan, *J. Am. Chem. Soc.*, 2010, **132**, 5405–5412.
- 32 J. Tan, J. Zhang, Y. Luo and S. Ye, *J. Am. Chem. Soc.*, 2019, **141**, 1941–1948.
- 33 C. A. De Carufel, N. Quittot, P. T. Nguyen and S. Bourgault, *Angew. Chem., Int. Ed.*, 2015, **54**, 14383–14387.
- 34 Y. Xing, E. H. Pilkington, M. Wang, C. J. Nowell, A. Kakinien, Y. Sun, B. Wang, T. P. Davis, F. Ding and P. C. Ke, *Phys. Chem. Chem. Phys.*, 2017, **19**, 30627–30635.
- 35 D. C. Rodriguez Camargo, K. J. Korshavn, A. Jussupow, K. Raltchev, D. Goricanec, M. Fleisch, R. Sarkar, K. Xue, M. Aichler, G. Mettenleiter, A. K. Walch, C. Camilloni, F. Hagn, B. Reif and A. Ramamoorthy, *eLife*, 2017, **6**, e31226.
- 36 S. S. Dicke, M. Maj, C. R. Fields and M. T. Zanni, *RSC Chem. Biol.*, 2022, **3**, 931–940.
- 37 X. Dong, Q. Qiao, Z. Qian and G. Wei, *Biochim. Biophys. Acta, Biomembr.*, 2018, **1860**, 1826–1839.
- 38 P. H. Nguyen, A. Ramamoorthy, B. R. Sahoo, J. Zheng, P. Faller, J. E. Straub, L. Dominguez, J. E. Shea, N. V. Dokholyan, A. de Simone, B. Ma, R. Nussinov, S. Najafi, S. T. Ngo, A. Loquet, M. Chiricotto, P. Ganguly, J. McCarty, M. S. Li, C. Hall, Y. Wang, Y. Miller, S. Melchionna, B. Habenstein, S. Timr, J. Chen, B. Hnath, B. Strodel, R. Kaye, S. Lesné, G. Wei, F. Sterpone, A. J. Doig and P. Derreumaux, *Chem. Rev.*, 2021, **121**, 2545–2647.
- 39 M. Apostolidou, S. A. Jayasinghe and R. Langen, *J. Biol. Chem.*, 2008, **283**, 17205–17210.
- 40 G. L. Dignon, G. H. Zerze and J. Mittal, *J. Phys. Chem. B*, 2017, **121**, 8661–8668.
- 41 Q. Qiao, G. Wei, D. Yao and Z. Song, *Phys. Chem. Chem. Phys.*, 2019, **21**, 20239–20251.
- 42 L. Khemtémourian, H. Fatafta, B. Davion, S. Lecomte, S. Castano and B. Strodel, *Front. Mol. Biosci.*, 2022, **9**, 849979.
- 43 M. Christensen, K. K. Skeby and B. Schiøtt, *Biochemistry*, 2017, **56**, 4884–4894.
- 44 S. M. Patil, S. Xu, S. R. Sheftic and A. T. Alexandrescu, *J. Biol. Chem.*, 2009, **284**, 11982.
- 45 Y. Zhang, Y. Luo, Y. Deng, Y. Mu and G. Wei, *PLoS One*, 2012, **7**, e38191.
- 46 M. Pannuzzo, A. Raudino, D. Milardi, C. La Rosa and M. Karttunen, *Sci. Rep.*, 2013, **3**, 2781.



- 47 X. Li, M. Wan, L. Gao and W. Fang, *Sci. Rep.*, 2016, **6**, 21614.
- 48 C. Poojari, D. Xiao, V. S. Batista and B. Strodel, *Biophys. J.*, 2013, **105**, 2323–2332.
- 49 N. Liu, M. Duan and M. Yang, *Sci. Rep.*, 2017, **7**, 7915.
- 50 Z. Qian, Y. Zou, Q. Zhang, P. Chen, B. Ma, G. Wei and R. Nussinov, *Biochim. Biophys. Acta, Biomembr.*, 2018, **1860**, 1818–1825.
- 51 Q. Tan, H. Liu, M. Duan and S. Huo, *Biochim. Biophys. Acta, Biomembr.*, 2021, **1863**, 183691.
- 52 J. Zhao, Y. Luo, H. Jang, X. Yu, G. Wei, R. Nussinov and J. Zheng, *Biochim. Biophys. Acta, Biomembr.*, 2012, **1818**, 3121–3130.
- 53 A. Sepehri, B. Nepal and T. Lazaridis, *J. Chem. Inf. Model.*, 2021, **61**, 4645–4655.
- 54 N. F. Dupuis, C. Wu, J.-E. Shea and M. T. Bowers, *J. Am. Chem. Soc.*, 2011, **133**, 7240–7243.
- 55 Y. Bram, A. Lampel, R. Shaltiel-Karyo, A. Ezer, R. Scherzer-Attali, D. Segal and E. Gazit, *Angew. Chem., Int. Ed.*, 2015, **54**, 2062–2067.
- 56 K. Lindorff-Larsen, S. Piana, K. Palmo, P. Maragakis, J. L. Klepeis, R. O. Dror and D. E. Shaw, *Proteins: Struct., Funct., Bioinf.*, 2010, **78**, 1950–1958.
- 57 J. P. M. Jämbbeck and A. P. Lyubartsev, *J. Phys. Chem. B*, 2012, **116**, 3164–3179.
- 58 J. P. M. Jämbbeck and A. P. Lyubartsev, *J. Chem. Theory Comput.*, 2013, **9**, 774–784.
- 59 W. L. Jorgensen, J. Chandrasekhar, J. D. Madura, R. W. Impey and M. L. Klein, *J. Chem. Phys.*, 1983, **79**, 926.
- 60 A. Bellet, A. Habrard and M. Sebban, *Metric Learning*, Springer, Cham, 2015.
- 61 H. J. C. Berendsen, J. P. M. Postma, W. F. van Gunsteren, A. DiNola and J. R. Haak, *J. Chem. Phys.*, 1984, **81**, 3684–3690.
- 62 T. Darden, D. York and L. Pedersen, *J. Chem. Phys.*, 1993, **98**, 10089–10092.
- 63 B. Hess, H. Bekker, H. J. C. Berendsen and J. G. E. M. Fraaije, *J. Comput. Chem.*, 1997, **18**, 1463–1472.
- 64 S. Miyamoto and P. A. Kollman, *J. Comput. Chem.*, 1992, **13**, 952–962.
- 65 S. Nose, *J. Chem. Phys.*, 1984, **81**, 511–519.
- 66 M. Parrinello, A. Rahman and R. A. Parrinello, *Phys. Rev. Lett.*, 1980, **45**, 1196–1199.
- 67 W. Kabsch and C. Sander, *Biopolymers*, 1983, **22**, 2577–2637.
- 68 R. T. McGibbon, K. A. Beauchamp, M. P. Harrigan, C. Klein, J. M. Swails, C. X. Hernández, C. R. Schwantes, L.-P. Wang, T. J. Lane and V. S. Pande, *Biophys. J.*, 2015, **109**, 1528–1532.
- 69 E. N. Baker and R. E. Hubbard, *Prog. Biophys. Mol. Biol.*, 1984, **44**, 97–179.
- 70 X. Daura, K. Gademann, B. Jaun, D. Seebach, W. F. Van Gunsteren and A. E. Mark, *Angew. Chem., Int. Ed.*, 1999, **38**, 236–240.
- 71 R. González-Alemán, D. Hernández-Castillo, A. Rodríguez-Serradet, J. Caballero, E. W. Hernández-Rodríguez and L. Montero-Cabrera, *J. Chem. Inf. Model.*, 2020, **60**, 444–448.
- 72 V. Gapsys, B. L. De Groot and R. Briones, *J. Comput. Aided. Mol. Des.*, 2013, **27**, 845–858.
- 73 J. P. Douliez, A. Ferrarini and E. J. Dufourc, *J. Chem. Phys.*, 1998, **109**, 2513–2518.
- 74 R. Gautier, A. Bacle, M. L. Tiberti, P. F. Fuchs, S. Vanni and B. Antonny, *Biophys. J.*, 2018, **1**, 436–444.
- 75 M. F. M. Engel, H. A. Yigittop, R. C. Elgersma, D. T. S. Rijkers, R. M. J. Liskamp, B. De Kruijff, J. W. M. Höppener and J. Antoinette Killian, *J. Mol. Biol.*, 2006, **356**, 783–789.
- 76 J. R. Brender, E. L. Lee, M. A. Cavitt, A. Gafni, D. G. Steel and A. Ramamoorthy, *J. Am. Chem. Soc.*, 2008, **130**, 6424–6429.
- 77 R. Soong, J. R. Brender, P. M. Macdonald and A. Ramamoorthy, *J. Am. Chem. Soc.*, 2009, **131**, 7079–7085.
- 78 H. M. Sanders, F. Chalyavi, C. R. Fields, M. M. Kostelic, M. H. Li, D. P. Raleigh, M. T. Zanni and M. T. Marty, *J. Am. Soc. Mass Spectrom.*, 2023, **34**, 986–990.
- 79 R. Akter, A. Abedini, Z. Ridgway, X. Zhang, J. Kleinberg, A. M. Schmidt and D. P. Raleigh, *Isr. J. Chem.*, 2017, **57**, 750–761.
- 80 R. Akter, R. L. Bower, A. Abedini, A. M. Schmidt, D. L. Hay and D. P. Raleigh, *ACS Chem. Biol.*, 2018, **13**, 2747–2757.

