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# Composition dependence of X-ray stability and degradation mechanisms at lead halide perovskite single crystal surfaces†

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The multiple applications of lead halide perovskite materials and the extensive use of X-ray based techniques to characterize them highlight a need to understand their stability under X-ray irradiation. Here, we present a study where the X-ray stability of five different lead halide perovskite compositions (MAPbI<sub>3</sub>, MAPbCl<sub>3</sub>, MAPbBr<sub>3</sub>, FAPbBr<sub>3</sub>, CsPbBr<sub>3</sub>) was investigated using photoelectron spectroscopy. To exclude effects of thin film formation on the observed degradation behaviors, we studied clean surfaces of single crystals. Different X-ray resistance and degradation mechanisms were observed depending on the crystal composition. Overall, perovskites based on the MA<sup>+</sup> cation were found to be less stable than those based on FA<sup>+</sup> or Cs<sup>+</sup>. Metallic lead formed most easily in the chloride perovskite, followed by bromide, and only very little metallic lead formation was observed for MAPbI<sub>3</sub>. MAPbI<sub>3</sub> showed one main degradation process, which was the radiolysis of MAI. Multiple simultaneous degradation processes were identified for the bromide compositions. These processes include ion migration towards the perovskite surface and the formation of volatile and solid products in addition to metallic lead. Lastly, CsBr formed as a solid degradation product on the surface of CsPbBr<sub>3</sub>.

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## Introduction

Lead halide perovskites have gained great interest in the scientific community due to their useful properties in multiple applications, including photovoltaics, light emitting diodes (LEDs), photodetectors, and more recently, their use in batteries and hydrogen production.<sup>1–4</sup> The general chemical formula for these materials is ABX<sub>3</sub>, where A represents a monovalent cation (*e.g.*, methylammonium (MA<sup>+</sup>), CH<sub>3</sub>NH<sub>3</sub><sup>+</sup>; formamidinium (FA<sup>+</sup>), CH(NH<sub>2</sub>)<sub>2</sub><sup>+</sup>; Cs<sup>+</sup>), B is a divalent cation (*e.g.*, Pb<sup>2+</sup>) and X is typically a halide (*e.g.*, Br<sup>−</sup>, I<sup>−</sup>, Cl<sup>−</sup>).

Containing heavy elements, such as lead, iodide, or bromide, enables lead halide perovskites to detect high energy radiation (*e.g.*, X-rays or  $\gamma$ -rays) with high sensitivity and low detection limits.<sup>5</sup> Therefore, perovskites have been suggested

as the active material in photodetection applications, such as medical imaging, where studies indicate comparable performance to currently used materials.<sup>5</sup> Additionally, they have been tested as extra-terrestrial solar cells for spacecraft power systems.<sup>6</sup> Their performance under these conditions, such as absence of atmosphere (lack of moisture or oxygen) and extreme temperature fluctuations, has been tested on lead halide perovskite solar cells for a period of 10 months and shown potential.<sup>6</sup>

Despite their potential, their stability remains a challenge to address. Exposure to humidity, oxygen, illumination, and heat can all contribute to material deterioration.<sup>3</sup> Particularly in their applications as photodetector or outer space solar cell, commercial implementation requires stability against high energy radiation.<sup>7,8</sup> For this purpose, the X-ray stability of different perovskite compositions has to be assessed.

In addition to the previously mentioned applications, various characterization techniques which are extensively used for the study of lead halide perovskites rely on X-rays, such as powder and single crystalline X-ray diffraction (XRD), scanning X-ray fluorescence (XRF), energy dispersive X-ray analysis (EDX) and photoelectron spectroscopy (PES). To obtain reliable results, it is crucial to identify and minimize any degradation from the measurement process itself. For example, PES is extensively used to study surfaces and interfaces of lead halide

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perovskites<sup>9,10</sup> and to investigate causes for their degradation when exposed to environmental stressors like humidity or heat.<sup>3,11,12</sup> However, the interpretation of the results becomes more challenging if X-ray induced initial shifts in binding energy position,<sup>13</sup> or the formation of metallic lead (Pb<sup>0</sup>) are present.<sup>8,14</sup> Therefore, understanding the degradation mechanisms behind high energy radiation is necessary to prevent radiation-induced damage and to ensure a correct interpretation of measurement results.

Despite the high importance of understanding the degradation mechanisms caused by X-ray radiation, limited research has been done so far and most studies have focused on thin films. Considering the improving performance of perovskite single crystals in devices (*e.g.* solar cells or photodetectors),<sup>15–17</sup> there is an interest to understand the stability behavior of single crystals both for single crystals devices and as a model system for thin films.

Some examples are a comparative X-ray stability study on CsPbBr<sub>3</sub>, CsMAPbBr<sub>3</sub>, and MAPbBr<sub>3</sub> thin films on top of a half solar cell structure (FTO/TiO<sub>2</sub>/Perovskite), where a stability dependence on the monovalent cation was found by PES characterization.<sup>18</sup> Specifically, the purely inorganic CsPbBr<sub>3</sub> was found to be rather stable with no detection of metallic lead (Pb<sup>0</sup>), followed by CsMAPbBr<sub>3</sub> which contained a steady amount of Pb<sup>0</sup> throughout the whole experiment. In contrast, metallic lead was detected after two hours of X-ray irradiation on the MAPbBr<sub>3</sub> thin film with the Pb<sup>0</sup> concentration increasing over time. Density functional theory (DFT) calculations suggested that the metallic lead initiates MABr vacancies which then migrate towards the surface. Other studies in literature compared MAPbI<sub>3</sub>, FAPbI<sub>3</sub>, Cs<sub>x</sub>MA<sub>1-x</sub>PbI<sub>3</sub> and Cs<sub>x</sub>MA<sub>y</sub>FA<sub>1-x-y</sub>PbI<sub>3</sub> thin films after exposure to a controlled amount of  $\gamma$  radiation and observed a formation of Pb<sup>0</sup> for all compositions except FAPbI<sub>3</sub>.<sup>19</sup> MAPbI<sub>3</sub> thin films were investigated *via* PES under ultrahigh vacuum conditions and it was reported that soft X-rays caused a rapid decomposition of the organic cation and, after long exposure, also the formation of metallic lead.<sup>20</sup> Furthermore, a comparative stability study on MAPbI<sub>3</sub>, MAPbBr<sub>3</sub>, Cs<sub>0.15</sub>FA<sub>0.85</sub>PbI<sub>3</sub>, Cs<sub>0.1</sub>MA<sub>0.15</sub>FA<sub>0.75</sub>PbI<sub>3</sub>, CsPbI<sub>3</sub>, and CsPbBr<sub>3</sub> thin films using PES and XRD has shown no  $\gamma$ -ray radiation damage on most compositions, except MAPbI<sub>3</sub> and CsPbI<sub>3</sub>.<sup>21</sup> MAPbI<sub>3</sub> demonstrated degradation to PbI<sub>2</sub> and CsPbI<sub>3</sub> degraded to Pb<sup>0</sup> at the surface. However, experiments on full solar cell stacks of these compositions showed that the performance of the MAPbI<sub>3</sub> solar cell was only slightly affected by  $\gamma$ -rays. These results together with DFT calculations suggested that the radiation-induced degradation was reversible in a full solar cell stack for MAPbI<sub>3</sub>. Additionally, Milotti *et al.* investigated the X-ray stability of FAPbBr<sub>3</sub> thin films using PES and micro-photoluminescence recently.<sup>22</sup> They observed the formation of metallic lead and the generation of Br<sub>2</sub> gas. However, prolonged irradiation led to re-oxidation of metallic lead and FA<sup>+</sup> and Br<sup>-</sup> ion migration, which was reported as self-healing of the material. In our previous study, we compared the X-ray stability between CsPbBr<sub>3</sub> and CsFAPbI<sub>3</sub> thin films, where we reported a higher stability of CsPbBr<sub>3</sub><sup>23</sup> and identified a

distinct degradation path for each composition. In CsPbBr<sub>3</sub>, X-ray exposure resulted in defect accumulation, eventually leading to the formation of Pb<sup>0</sup>. In contrast, we found degradation of the organic cation FA<sup>+</sup> without the formation of Pb<sup>0</sup> in CsFAPbI<sub>3</sub>.

Polycrystalline thin films typically degrade initially at the grain boundaries under environmental stressors. In other words, variations in the film formation (and therefore grain structure) due to the preparation methods may significantly impact degradation processes.<sup>3</sup> To avoid these external factors, single crystal surfaces instead of thin film surfaces can be studied. We have shown that clean surfaces of perovskite single crystals can be prepared in vacuum through cleaving,<sup>10,11,24</sup> which enables us to study intrinsic behaviors of the material. Therefore, to take our previous study further and to gain insight into the behavior of lead halide perovskite surfaces with X-rays without the effects of grain boundaries, we decided to use perovskite single crystals in the work presented here. There are few studies on the X-ray stability of perovskite single crystals reported in literature, despite the implementation of single crystals in practical applications increasing considerably.<sup>25</sup> For instance, Motoki *et al.* studied the X-ray stability of MAPbI<sub>3</sub> single crystals using in-house PES.<sup>26</sup> They compared core levels of the perovskite after exposure to X-rays for durations from 0 to 200 hours and found minor oxygen contributions. The core level spectra N 1s, I 3d and Pb 4f showed a decrease and remained constant after 40 hours, indicating X-ray induced degradation and eventually stabilization. Other studies focused on X-ray and e-beam induced degradation of MAPbBr<sub>3</sub> single crystals using various techniques, such as surface photovoltage spectroscopy, XRD and PES. Exposing the crystals to X-rays under ambient conditions lead to quenching of free excitons into a new bound species, with the material recovering after one week under dark conditions and low humidity.<sup>27</sup> On the other hand, both X-ray and e-beam irradiation under ultra-high vacuum (UHV) conditions induced the formation of metallic lead<sup>28,29</sup> and the release of MABr and Br<sub>2</sub>.<sup>28</sup>

In the study presented here, the X-ray stability of clean surfaces of different lead halide perovskite single crystals was investigated *via* synchrotron-based PES, as PES allows following the compositional (chemical) changes on the perovskite surface and therefore to identify degradation pathways. To observe general, composition-dependent trends and compare the effects of the halide and the A-site cation, five different perovskite single crystal compositions were analyzed, and the results were divided in two main groups. MAPbBr<sub>3</sub>, MAPbI<sub>3</sub> and MAPbCl<sub>3</sub> were measured to study the impact of the halide on the X-ray stability (MA<sup>+</sup>-series) and MAPbBr<sub>3</sub>, FAPbBr<sub>3</sub> and CsPbBr<sub>3</sub> (Br-series) were measured and compared to investigate the effect of the A-site cation.

## Experimental methods

### Synthesis of single crystals

Perovskite single crystals, including MAPbI<sub>3</sub>, MAPbBr<sub>3</sub>, MAPbCl<sub>3</sub>, FAPbBr<sub>3</sub>, and CsPbBr<sub>3</sub>, with diameters ranging from 0.5 to 1 cm, were synthesized using the inverse temperature crystallization



(ITC) method. The starting materials were used as bought without further treatment. Methylammonium chloride ( $\text{CH}_3\text{NH}_3\text{Cl}_3$ , MACl), methylammonium iodide ( $\text{CH}_3\text{NH}_3\text{I}$ , MAI), methylammonium bromide ( $\text{CH}_3\text{NH}_3\text{Br}$ , MABr), formamidinium bromide ( $\text{CH}(\text{NH}_2)_2\text{Br}$ , FABr) and cesium bromide (CsBr) were purchased at Sigma Aldrich. Lead chloride ( $\text{PbCl}_2$ ), lead iodide ( $\text{PbI}_2$ ) and lead bromide ( $\text{PbBr}_2$ ) were bought at TCI. To prepare the precursor solutions, 1 M solutions of 1:1 AX:PbX (A = MA/FA and X = I/Br/Cl) were made using different solvents such as  $\gamma$ -butyrolactone (GBL) for  $\text{MAPbI}_3$ , dimethylformamide (DMF) for  $\text{MAPbBr}_3$  and  $\text{FAPbBr}_3$ , dimethyl sulfoxide (DMSO) for  $\text{MAPbCl}_3$ . The solutions were stirred at room temperature until all the precursors were fully dissolved. For  $\text{CsPbBr}_3$ , CsBr and  $\text{PbBr}_2$  (in a molar ratio of 1:2) were dissolved on DMSO to form a saturated perovskite precursor solution. After filtration through a 0.45  $\mu\text{m}$  PTFE filter, each solution was transferred to an open glass vial and heated to specific temperatures: 100 °C for  $\text{MAPbI}_3$ , 80 °C for  $\text{MAPbBr}_3$  and  $\text{FAPbBr}_3$ , 90 °C for  $\text{MAPbCl}_3$ , and 120 °C for  $\text{CsPbBr}_3$ . Once the desired temperature was reached, selected seed crystals were introduced into a new solution and allowed to grow with the vial sealed until the desired crystal size was achieved.

The most promising crystals were chosen for characterization using photoelectron spectroscopy. The remaining crystals from the same batch were grounded using a mortar and pestle and characterized by powder X-ray diffraction analysis. Fig. S1 (ESI<sup>†</sup>) demonstrates that all the synthesized materials exhibit a single phase, free from any impurities, which is consistent with the results obtained from the single crystal XRD profiles from literature.

### Photoelectron spectroscopy and beamline set-up

The PES measurements were carried out at the SMS-branch of the FlexPES<sup>30</sup> beamline at the MAX IV synchrotron source in Lund. To measure the single crystal surfaces, the single crystals were mounted on sample plates using a two-component epoxy (EPO-TEK H20E) and cured at 100 °C for one hour under ambient conditions. The single crystals were then transferred into a preparation chamber under ultra-high vacuum (UHV) conditions and cleaved at around  $1 \times 10^{-8}$  mbar to obtain a clean crystal surface for the PES measurements. Subsequently, they were measured at  $10^{-9}$ – $10^{-10}$  mbar.

The X-rays for the measurements were generated using a linearly polarized undulator (period length = 54.4 mm) and monochromated *via* a plane grating (Zeiss SX700). To detect the ejected photoelectrons, a hemispherical spectrometer (Scienta DA30-L (W)) in normal emission mode was used. Relevant core level spectra (N 1s, C 1s, Cs 4d, Pb 4f, I 4d, Br 3d, Cl 2p, O 1s) were recorded in loops at 600 eV photon energy with a step size of 0.1 eV and a pass energy of 100 eV for each crystal composition.

All core levels were intensity normalized with respect to the Pb 4f intensity and energy calibrated against Pb 4f<sub>7/2</sub> ( $\text{Pb}^{2+}$ ) = 138.54 eV binding energy,<sup>24</sup> unless otherwise mentioned. The quantification was done using peak fitting with the pseudo-Voigt function with a linear, or Shirley background. All N 1s

core levels were fitted using a single peak, I 4d, Cl 2p and FA-based Br 3d core levels were fitted using a single spin-orbit doublet. On the other hand, all Pb 4f core levels and Cs 4d and Br 3d core levels from  $\text{CsPbBr}_3$  were fitted with 2 peaks to account for the different contributions. As single crystal surfaces were investigated, a homogeneous distribution of their surface was assumed. The inelastic mean free path (IMFP) was estimated using the TTP-2M method. The IMFP ranged between 0.71 and 1.56 nm depending on the core level and the crystal. Details of the calculation can be found in the ESI<sup>†</sup> in Tables S1 and S2.

The sample degradation varies depending on the X-ray intensity. To avoid instant degradation through beam damage and to facilitate the study of degradation over larger timescales the beam was defocused and had a fullwidth half maximum (FWHM) of  $1 \times 0.4$  (horizontal  $\times$  vertical) mm in normal incidence, and the incident beam intensity was regulated *via* the exit slit. Considering a Gaussian intensity distribution, the standard deviation was estimated to be 0.43 mm (horizontal)  $\times$  0.17 mm (vertical) from the FWHM. Based on these values and assuming an elliptical beam, we estimate that 68% of the incident beam intensity fall in an area of 0.23 mm<sup>2</sup>. The number of photons was calculated based on the time required to measure one loop of core levels and the number of photons per loop, considering the corresponding slit and beam size for each sample. The fluence (total number of photons per cm<sup>2</sup>) was calculated by multiplying the number of photons with 0.68/0.0023 (see Table S3 in the ESI<sup>†</sup> to obtain all values per loop).

The exit slit was adjusted to 10  $\mu\text{m}$  (low flux, LF =  $1.2 \times 10^{11}$  photons per second) for the  $\text{MAPbBr}_3$  and  $\text{MAPbCl}_3$  single crystals, and to 30  $\mu\text{m}$  (high flux, HF =  $3.6 \times 10^{11}$  photons per second) for the more X-ray-stable  $\text{MAPbI}_3$  and  $\text{CsPbBr}_3$  compositions.  $\text{FAPbBr}_3$  was measured under both LF and HF. The flux values were theoretically estimated based on the used beamline settings as is described in the FlexPES beamline paper.<sup>30</sup> These theoretical values are expected to be somewhat higher than the actual ones at the used photon energy, but give a good estimate, which is based on the beamline settings used in the experiment. The results presented for  $\text{FAPbBr}_3$  throughout the paper were measured under high flux, as this allows us to follow the degradation quicker and, therefore, improves comparability to the measurements of the other compositions. The core levels of  $\text{FAPbBr}_3$  measured using low flux can be found in the ESI<sup>†</sup> in Fig. S2. It should be noted that the exposure time for the measurement under high flux of photons was reduced to achieve the same photon count.

## Results and discussion

The work presented here studied five distinct lead halide perovskite single crystal compositions ( $\text{MAPbI}_3$ ,  $\text{MAPbBr}_3$ ,  $\text{MAPbCl}_3$ ,  $\text{FAPbBr}_3$ , and  $\text{CsPbBr}_3$ ), which can be categorized into two series. The first series consists of a constant cation (methylammonium) with varying halides (iodide, bromide, chloride) while the second series consists of a constant halide



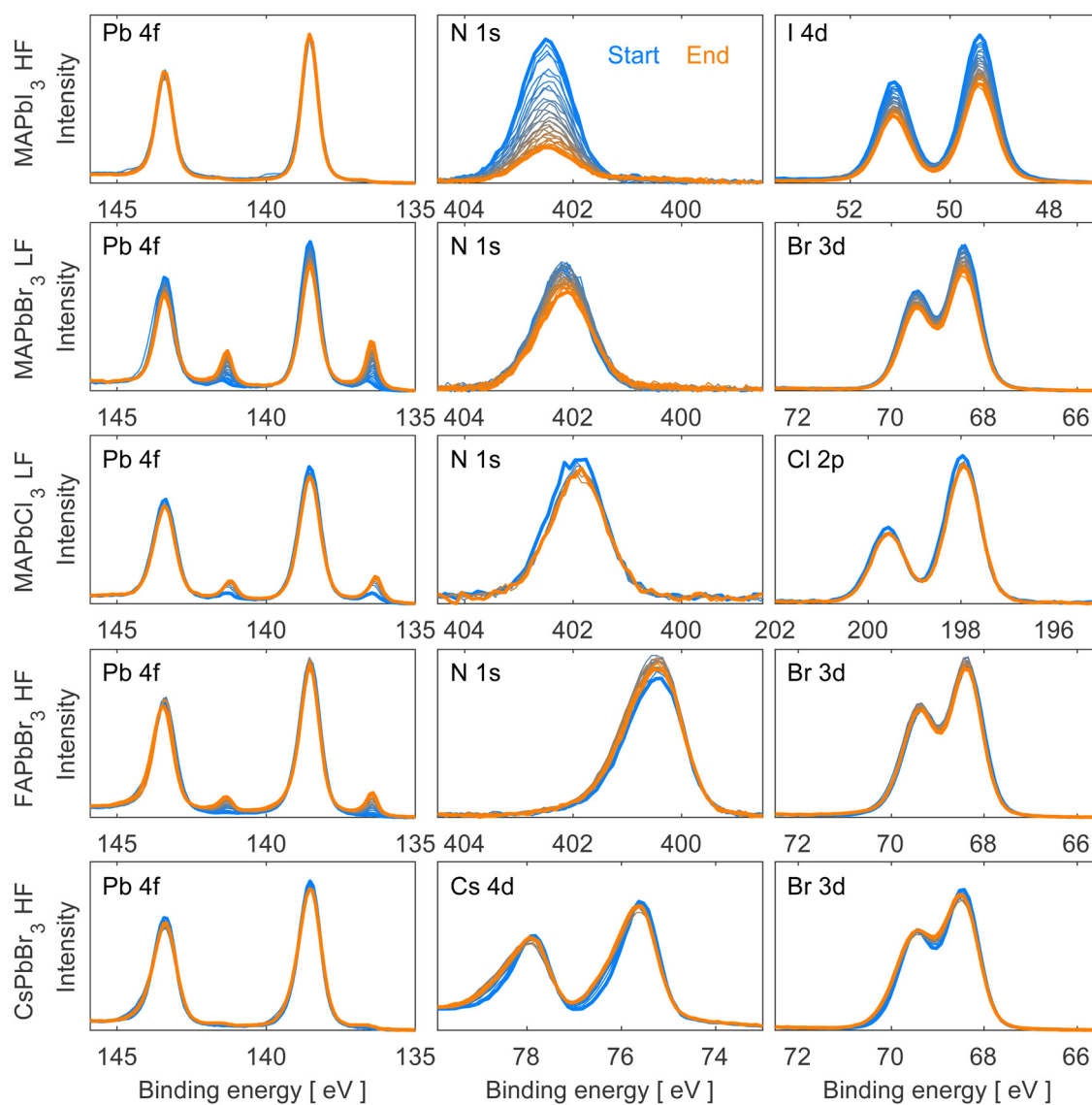
(bromide) with varying cations (methylammonium, formamminium, cesium). Depending on the composition of the material, different degradation patterns were observed, including the formation of metallic lead ( $\text{Pb}^0$ ) and/or the degradation of the organic cation.

Fig. 1 depicts the evolution of Pb 4f, N 1s, Cs 4d, and halide (I 4d, Br 3d, Cl 2p) core levels spectra for all the studied materials during irradiation. Fig. S3 in the ESI† provides additional core level spectra (C 1s and O 1s). All core level spectra presented in both figures were measured using a photon energy of 600 eV, intensity normalized to the total amount of Pb 4f ( $\text{Pb}^0 + \text{Pb}^{2+}$ ), and energy calibrated against Pb 4f<sub>7/2</sub> ( $\text{Pb}^{2+}$ ) at 138.54 eV binding energy. Fig. S4 in ESI† shows the same core level spectra as Fig. 1 but normalized to the perovskite contribution ( $\text{Pb}^{2+}$ ). All the measurements were

carried out in loops for each single crystal under constant irradiation with the initial spectrum shown in blue and the final in orange.

The beam flux was adjusted depending on the exact crystal composition to avoid immediate beam degradation and enable degradation processes to happen within a reasonable time-frame. Since  $\text{MAPbI}_3$ ,  $\text{FAPbBr}_3$  and  $\text{CsPbBr}_3$  exhibit greater stability towards X-rays, they were exposed to a higher beam flux ( $\text{HF} = 3.6 \times 10^{11}$  photons per second). Conversely, due to their lower stability towards X-rays,  $\text{MAPbCl}_3$  and  $\text{MAPbBr}_3$  were measured at a lower beam flux ( $\text{LF} = 1.2 \times 10^{11}$  photons per second) to slow down the degradation processes.

We observe no oxygen contamination and a clear C 1s peak corresponding to the perovskite and no adventitious carbon, which usually appears lower binding energies



**Fig. 1** Pb 4f, N 1s/Cs 4d and I 4d/Br 3d/Cl 2p core level spectra of clean surfaces of  $\text{MAPbI}_3$ ,  $\text{MAPbBr}_3$ ,  $\text{MAPbCl}_3$ ,  $\text{FAPbBr}_3$  and  $\text{CsPbBr}_3$  single crystals, recorded using the FlexPES beamline (MAX IV) with a photon energy of 600 eV. The spectra are intensity normalized to the total amount of Pb 4f and energy calibrated against Pb 4f<sub>7/2</sub> ( $\text{Pb}^{2+}$ ) = 138.54 eV binding energy. Low flux ( $\text{LF} = 1.2 \times 10^{11}$  photons per second), high flux ( $\text{HF} = 3.6 \times 10^{11}$  photons per second).



(around 284.8 eV)<sup>31,32</sup> for all measured single crystals except for FAPbBr<sub>3</sub> and CsPbBr<sub>3</sub> which both present a small contribution of adventitious carbon and oxygen (Fig. S3, ESI†). Despite not having a perfectly clean surface for these two crystals, the small amounts of oxygen and adventitious carbon do not change during the experiments, and therefore do not take part in the degradation process and should not interfere with our results.

As can be seen in Fig. 1 and Fig. S4 (ESI†), all measured single crystals present a similar shape of their Pb 4f core level corresponding to Pb<sup>2+</sup> with the Pb 4f<sub>7/2</sub> peak calibrated at 138.54 eV. A metallic lead (Pb<sup>0</sup>) contribution appears during X-ray irradiation at a binding energy of 137.0 eV<sup>33</sup> and will be discussed in detail in the following paragraphs. All Br-based compounds also present a similar shape and position of Br 3d doublet associated with Br<sup>-</sup> (68.4 and 69.5 eV). The N 1s core level presents the same shape for all four hybrid materials with a clear difference in position for FA<sup>+</sup>-based (400.5 eV (Br)) and MA<sup>+</sup>-based (401.8 eV (Cl), 402.1 eV (Br) and 402.5 eV (I)) perovskites due to the different chemical environments of the elements in the organic cations. For MAPbI<sub>3</sub> and MAPbCl<sub>3</sub>, the observation of a I 4d doublet with the I 4d<sub>5/2</sub> at 49.4 eV and Cl 2p doublet with the Cl 2p<sub>3/2</sub> at 197.9 eV confirms the presence of these anions. The detection of Cs 4d doublet with the Cs 4d<sub>5/2</sub> at 75.5 eV confirms the presence of Cs<sup>+</sup> in the CsPbBr<sub>3</sub> compound.

Analyzing Fig. 1, MAPbI<sub>3</sub> exhibited only small amounts of metallic lead during the experiment, while the organic cation (observed through nitrogen) and iodide intensities decreased significantly. In contrast, the other hybrid perovskite single crystals (MAPbBr<sub>3</sub>, MAPbCl<sub>3</sub>, and FAPbBr<sub>3</sub>) formed significant amounts of metallic lead under irradiation accompanied by a decrease of the organic cation and halide content in the case of MAPbBr<sub>3</sub> and MAPbCl<sub>3</sub>. On the other hand, FAPbBr<sub>3</sub> depicts only slight changes in the Br 3d spectra and an increase in the N 1s peak intensity, in agreement with literature.<sup>22</sup> Lastly, the fully inorganic CsPbBr<sub>3</sub> perovskite single crystal demonstrated minimal metallic lead formation and peak broadening of the Br 3d and the Cs 4d core level spectra. Similar results were previously reported by our group for CsPbBr<sub>3</sub> thin films.<sup>23</sup>

To gain quantitative information about these changes and therefore the mechanisms behind them, all core levels were curve fitted. Notably, the N 1s and halide core level spectra of all crystal compositions except CsPbBr<sub>3</sub> could be fitted well using the perovskite fit contributions. In other words, no additional non-perovskite contributions were found for those core levels for most compositions. Instead, we observed changes in peak intensities. For CsPbBr<sub>3</sub>, however, additional, non-perovskite peak contributions were found at higher binding energy for the Cs 4d and Br 3d core levels. Fits of the first and last measurement can be found in Fig. S5 and S6 in the ESI† and fitting details are described in the experimental section.

The trends of the metallic lead formation for the different studied perovskite crystal compositions and their dependence on the fluence (photons cm<sup>-2</sup>) were examined after careful curve fitting of all individual Pb 4f spectra. The contribution of metallic lead as a function of fluence (photons cm<sup>-2</sup>) is

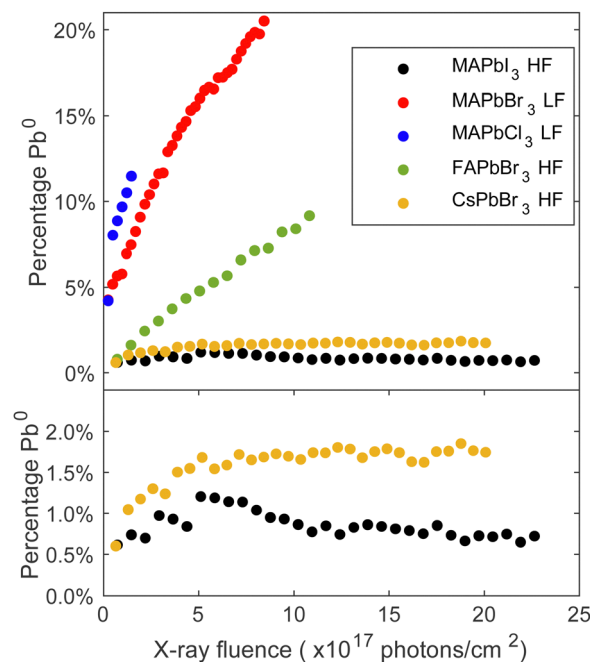


Fig. 2 Percentage Pb<sup>0</sup> (Pb<sup>0</sup> intensity/total Pb intensity) on the surface of each material (MAPbI<sub>3</sub>, MAPbBr<sub>3</sub>, MAPbCl<sub>3</sub>, FAPbBr<sub>3</sub>, CsPbBr<sub>3</sub>) as a function of fluence (photons cm<sup>-2</sup>). Top panel: All compositions. Bottom panel: Zoom on MAPbI<sub>3</sub> and CsPbBr<sub>3</sub>.

presented in Fig. 2, which clearly describes two trends. Firstly, focusing on the series with the same cation (MA<sup>+</sup>) at the A-site and varying X-site halides (I<sup>-</sup>, Br<sup>-</sup>, Cl<sup>-</sup>): a small amount of metallic lead was detected on the MAPbI<sub>3</sub> crystal surface, which is consistent with the results presented in Fig. 1 and previous literature observations.<sup>34</sup> About 1% of metallic lead was found after exposing the surface to a dose of  $5 \times 10^{17}$  photons cm<sup>-2</sup>. After this point the amount of metallic lead decreased somewhat to stabilize at below 1% despite using the higher flux (HF) and exposing the material to  $22.5 \times 10^{17}$  photons cm<sup>-2</sup>. However, the lower flux (LF) caused approximately 10% metallic lead on the MAPbBr<sub>3</sub> surface after receiving  $2.4 \times 10^{17}$  photons cm<sup>-2</sup> and 16% after receiving  $5 \times 10^{17}$  photons cm<sup>-2</sup> and increased further until the end of the experiment ( $8.3 \times 10^{17}$  photons cm<sup>-2</sup>). In comparison, MAPbCl<sub>3</sub> formed 10% metallic lead after exposure to only half the fluence ( $1.2 \times 10^{17}$  photons cm<sup>-2</sup>) compared to MAPbBr<sub>3</sub>.

In the series with a constant halide (Br<sup>-</sup>) and varying cations (Cs<sup>+</sup>, FA<sup>+</sup>, MA<sup>+</sup>), the MAPbBr<sub>3</sub> crystal was least stable. FAPbBr<sub>3</sub> formed 5% metallic lead after exposure to  $5 \times 10^{17}$  photons cm<sup>-2</sup> using the higher flux (HF) and the fully inorganic CsPbBr<sub>3</sub> demonstrated only a small amount of metallic lead formation. Approximately 1.6% metallic formed after receiving  $5 \times 10^{17}$  photons cm<sup>-2</sup> using the higher flux (HF), and this amount did not increase much more under continued irradiation. This is expected as CsPbBr<sub>3</sub> lacks the volatile organic cation.<sup>35</sup> Other studies have reported the presence metallic lead in all these compositions under environmental stressors such as light or X-rays.<sup>22,27,34,36,37</sup> However, most studies focus on polycrystalline thin films, where degradation is different and facilitated by the

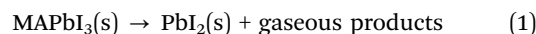


grain boundaries and vacancy formation. It is worth noting that exposing different spots on the same FAPbBr<sub>3</sub> single crystal to an equal number of photons (e.g.  $7.4 \times 10^{17}$  photons cm<sup>-2</sup>) but with a different flux leads to different amounts of metallic lead formation: the spot measured with the high flux (HF) formed twice as much metallic lead as the spot measured with low flux (LF), as shown in Fig. S7 in the ESI† Our previous study also presented a dependence on the flux and found that once Pb<sup>0</sup> is detected, we observe an accelerated formation of Pb<sup>0</sup>.<sup>23</sup> On the contrary, it is possible to prevent the formation of Pb<sup>0</sup> by using a sufficiently low flux. Therefore, the choice of flux has a direct impact on the formation rate of Pb<sup>0</sup> and, hence, on the degradation process of the single crystal surface. In the context of the FAPbBr<sub>3</sub> crystal examined in this experiment, the crystal has the capacity to absorb a greater quantity of photons when they are distributed over a longer period of time. In summary, both the organic cation and the halide anion impact the degradation mechanism and the kinetics of forming metallic lead.

We will now continue to discuss the degradation mechanisms for each compound in detail starting with the MA<sup>+</sup>-series. To be consistent with the results presented in Fig. 1, all core level intensities are divided by the total lead intensity (Pb<sup>0</sup> + Pb<sup>2+</sup>) and initial ratios of nitrogen and halide were set to those expected from the perovskite stoichiometry (1 MA<sup>+</sup> cation, represented by the nitrogen, and 3 halides). Fig. 3 depicts the evolution of core level intensities of the MA<sup>+</sup>-series crystals as a function of fluence, showing the lead and nitrogen on the left axis (black) and the halide on the right axis (blue). The axes are scaled in a 1 : 3 (left : right) ratio for MAPbBr<sub>3</sub> and MAPbCl<sub>3</sub>. This means that the intensities of the perovskite components are scaled in relation to their abundance in the pristine perovskite crystal. For MAPbI<sub>3</sub> the axes are scaled in a 1 : 1 (left : right) ratio to be able to directly follow a radiolysis process expected for methylammonium iodide and how this develops during exposure. Overall, all materials, except MAPbI<sub>3</sub>, show a notable decrease in their Pb<sup>2+</sup> intensities which stems from the formation of metallic lead and is consistent with the findings reported in Fig. 2. Furthermore, a general reduction in the halide and organic cation contents is observed for all compositions.

Analyzing Fig. 3 in depth, different reaction mechanisms can be proposed. Starting with MAPbI<sub>3</sub>, a significant and rapid decrease in the intensities of its N 1s core level, accompanied by a decrease in its I 4d core level, is observed. Previous studies have reported the formation of MAI and PbI<sub>2</sub>, where PbI<sub>2</sub> subsequently decays to metallic lead and iodine.<sup>20,38</sup> Here, the Pb<sup>2+</sup> signal remains almost stable throughout our measurements, indicating only a small amount of metallic lead formation. However, a possible degradation to MAI and PbI<sub>2</sub> still holds and requires the simultaneous removal of one nitrogen and one iodide. With the scaling of the y-axis used here, this could have been observed through overlapping data points for N and I in Fig. 3. Considering the difference in IMFP of nitrogen and iodide (0.76 nm and 1.50 nm, respectively, further details in Table S1 in the ESI†), a stronger decrease of the N intensity would be expected if MAI is mainly removed in the surface region. However, here we observe that the I intensity decreases relatively more. This suggests that in addition to MAI removal from the perovskite surface, also some additional iodide is removed relative to the remaining perovskite components (Pb and organic cation). It is important to note that the PES measurements are carried out under UHV conditions. Hence, any gaseous degradation products are being quickly removed from the surface and any reversible reactions, as others have reported for MAPbI<sub>3</sub>,<sup>21</sup> cannot be observed in this experiment.

Overall, the results indicate the simultaneous disappearance of (at least) one iodide per organic cation from the surface of the single crystal. Therefore, our results align best with the hypothesis of MAI radiolysis, accompanied by the generation of PbI<sub>2</sub>:



Moving to the analysis of MAPbBr<sub>3</sub>, the Pb<sup>2+</sup> intensity decreases throughout the experiment, and is accompanied by the formation of metallic lead, a result that has been reported frequently in literature.<sup>27–29</sup> Furthermore, in both nitrogen and bromide an overall intensity decrease is observed, suggesting a partial removal of the organic cation and the halide from the substrate. It is important to note, that the observed decrease of

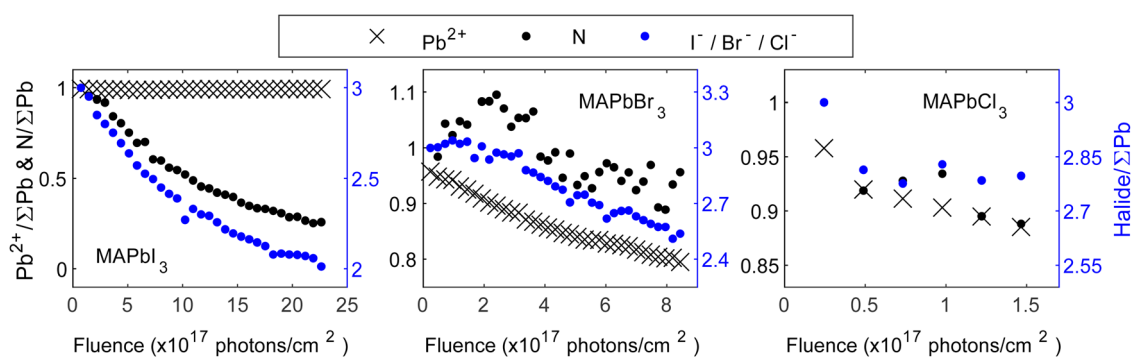
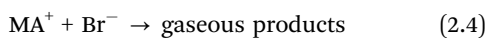
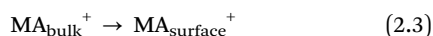
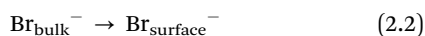


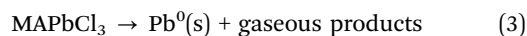
Fig. 3 Core level intensity evolution for the crystal surfaces of MAPbI<sub>3</sub> (left), MAPbBr<sub>3</sub> (middle) and MAPbCl<sub>3</sub> (right) as a function of the fluence (photons cm<sup>-2</sup>). Intensity normalized by the total amount of lead (Pb<sup>0</sup> + Pb<sup>2+</sup>) and initial ratios set to those expected from the perovskite stoichiometry (1 nitrogen and 3 halides). The left axes (black) refer to Pb<sup>2+</sup> and N, the right axes (blue) refer to the halide.



the bromide and nitrogen intensities is delayed relative to the decrease of  $\text{Pb}^{2+}$ , *i.e.*, the cause of the initial decrease in the signals of the different components are different and thus originate from different reactions and kinetics. Furthermore, the N intensity initially increases followed by a decrease, and the Br intensity decreases more than the N intensity. This indicates that multiple processes are involved in the degradation, in contrast to the behavior observed for  $\text{MAPbI}_3$  where the organic cation and the halide disappear together. Ion migration has been frequently reported for lead halide perovskites and could explain some of the observed behavior.<sup>39–42</sup> Specifically, the N intensity could increase initially due to ion migration and subsequently decrease due to radiolysis of  $\text{MABr}$ . The bromide intensity, on the other hand, may decrease faster than nitrogen due to a combination of  $\text{MABr}$  radiolysis and metallic lead formation. Overall, we propose that the degradation to  $\text{MABr}$  and  $\text{Br}_2^{28}$  is accompanied by ion migration of methylammonium ions from the bulk towards the surface as well as metallic lead formation:



In the case of  $\text{MAPbCl}_3$ , the intensity decreases of  $\text{Pb}^{2+}$ , nitrogen and chloride follow each other. Given the scaling of the y-axis, this means that we observe a simultaneous loss of all components in contrast to both  $\text{MAPbI}_3$  and  $\text{MAPbBr}_3$ . Specifically, the formation of metallic lead and the loss of the organic cation and halide happen simultaneously, suggesting that radiolysis of  $\text{MACl}$  occurs simultaneously with metallic lead formation:



In general, two different reactions have been discussed for removal of MAX from the perovskites in the scientific community:<sup>3,43</sup> MAX can form  $\text{NH}_3$  and  $\text{CH}_3\text{X}$  or  $\text{CH}_3\text{NH}_2$  and  $\text{HX}$ . As an example, it was found that  $\text{MAPbI}_3$  forms  $\text{CH}_3\text{NH}_2$  and  $\text{HI}$  at lower temperatures and  $\text{MAI}$  and  $\text{NH}_3$  at higher temperatures.<sup>44</sup> However, from our experiments it is not possible to distinguish which gaseous products are formed.

After discussing the trends for MA-series we will now focus on the trends for remaining perovskites containing the  $\text{Br}^-$  anion. Fig. 4 shows the core level intensities evolution of  $\text{FAPbBr}_3$  as a function of fluence (photons  $\text{cm}^{-2}$ ) where the intensity relation between the left:right axis is 1:3. Similar to Fig. 3, all core level intensities were normalized by the total amount of lead ( $\text{Pb}^0 + \text{Pb}^{2+}$ ) and initial ratios were set to those expected from the perovskite stoichiometry (1  $\text{FA}^+$  cation and 3  $\text{Br}^-$ ). While the lead content decreases continuously, the bromide and nitrogen intensities initially increase, followed by a subsequent decrease, similar to the observations made for  $\text{MAPbBr}_3$ . The bromide intensity returns close to its initial relative value of  $3 \text{ Br}^- / \sum \text{Pb}$ , while an excess of nitrogen

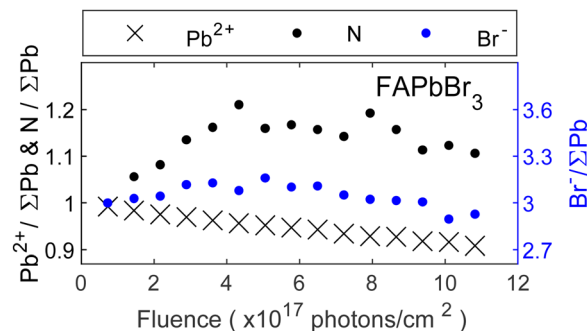
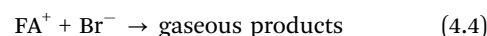
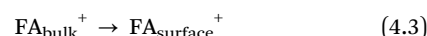
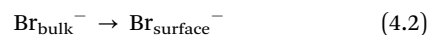
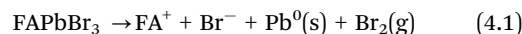


Fig. 4 Core level intensity evolution for the  $\text{FAPbBr}_3$  crystal surface as a function of fluence (photons  $\text{cm}^{-2}$ ). Intensity normalized by the total amount of lead ( $\text{Pb}^0 + \text{Pb}^{2+}$ ) and initial ratios set to those expected from the perovskite stoichiometry (1 cation represented by nitrogen and 3 bromide).

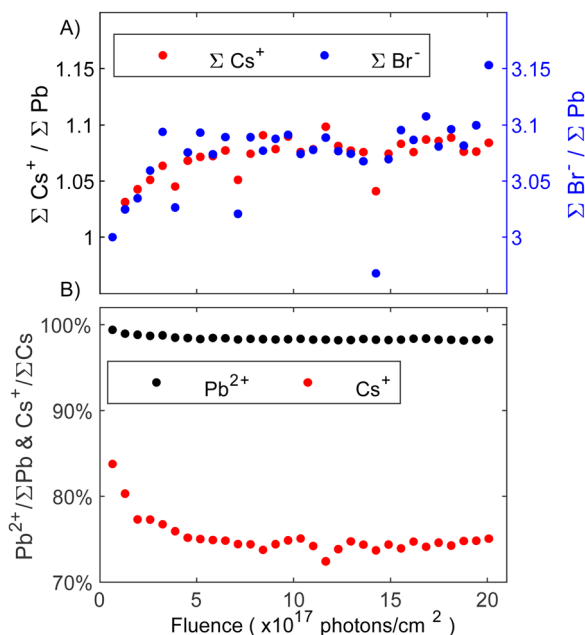
remains at the surface. This suggests that the metallic lead formation is also accompanied by other reactions. Previous studies have reported self-healing effects in  $\text{FAPbBr}_3$  thin films.<sup>22</sup> However, we observe a concentration increase beyond expected levels, followed by a decline and no recovery of metallic lead to  $\text{Pb}^{2+}$ . In contrast, self-healing implies an initial decrease followed by restoration of the original composition. Therefore, ion migration from the bulk to the surface seems more likely, as it would lead to higher concentrations of  $\text{FA}^+$  and  $\text{Br}^-$  at the surface but leave the bulk  $\text{FABr}$ -deficient. At the same time, the eventual decrease of nitrogen and bromide indicates a removal of the organic cation and bromide ions, most likely by radiolysis. Notably, the organic cation exhibits a stronger increase compared to the halide. Considering that the formation of metallic lead involves the loss of bromide ions, a decrease together with  $\text{Pb}^{2+}$  would be expected. However, this seems to be compensated by the bromide ion migration. Overall, similar reactions to  $\text{MAPbBr}_3$  seem likely:



However, the formation of metallic lead occurs at a slower rate compared to  $\text{MAPbBr}_3$  and ion migration leads to a temporary excess of  $\text{FABr}$  in the surface region.

Moving to the experiments on  $\text{CsPbBr}_3$ , an in-depth analysis of the core level spectra *via* curve fitting showed that the degradation process involves the formation of a secondary species at higher binding energy in the Cs 4d and Br 3d core level spectra (fits can be seen in Fig. S5 and S6 in the ESI†). In this context, it is important to note that minor C 1s and O 1s contributions were present (see Fig. S3, ESI†). However, given their consistent intensity levels throughout the experiment and considering their low overall intensities, no significant impact on the presented results is expected. Fig. 5A depicts the evolution of core level intensities of total Cs 4d and total Br 3d



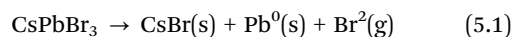


**Fig. 5** (A) Core level intensity evolution of total  $\text{Cs}^+$  (left, black axis) and total  $\text{Br}^-$  (right, blue axis) of the  $\text{CsPbBr}_3$  surface as a function of fluence ( $\text{photons cm}^{-2}$ ). Intensity normalized by the total amount of lead ( $\text{Pb}^0 + \text{Pb}^{2+}$ ) and initial ratios set to those expected from the perovskite stoichiometry (1  $\text{Cs}^+$  and 3  $\text{Br}^-$ ). The axes are scaled in a 1:1 ratio to one another. (B) Comparison of the percentage of decrease in perovskite  $\text{Pb}^{2+}$  and  $\text{Cs}^+$  contributions at the surface relative to total Pb 4f and total Cs 4d core level intensity, respectively.

normalized by total lead ( $\text{Pb}^0 + \text{Pb}^{2+}$ ) as a function of fluence. The initial values were set to the expected  $\text{CsPbBr}_3$  perovskite stoichiometry (1  $\text{Cs}^+$  and 3  $\text{Br}^-$ ). Cesium intensities are represented by the left axis (black), while bromide is represented by the right axis (blue) with a linear 1:1 scaling between the two axes. Both, the total amount of Cs and Br increase together in a 1:1 ratio relative to the total Pb 4f intensities, suggesting that the excess of Cs and Br forms in the same reaction. As the intensity increase is observed relative to the total lead intensity, we suggest that ion migration of  $\text{Br}^-$  and  $\text{Cs}^+$  is responsible.

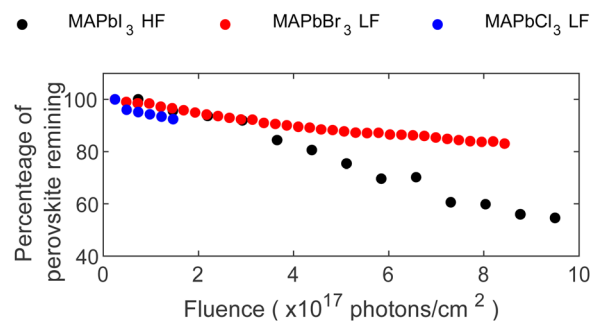
Fig. 5B shows the percentage of the perovskite  $\text{Pb}^{2+}$  and  $\text{Cs}^+$  core level intensities relative to the total amount of Pb and Cs, respectively. It can be seen in Fig. 5B that the perovskite Pb 4f contribution ( $\text{Pb}^{2+}$ ) presents an initial small decrease and then is almost constant throughout the entire experiment while the intensity of the perovskite  $\text{Cs}^+$  quickly decreases to 75% indicating a change in the surface composition. Considering that overall  $\text{Cs}^+$  and  $\text{Br}^-$  increase by the same amount (Fig. 5A), the formation of  $\text{CsBr}$  seems likely. This aligns with our previous study on thin films where  $\text{CsBr}$  was identified as a degradation product with the  $\text{CsBr}$  core levels appearing at higher binding energies than the perovskite Cs and Br.<sup>23</sup> Furthermore, PES experiments on  $(\text{Cs}_{0.06}\text{FA}_{0.79}\text{MA}_{0.15})\text{Pb}(\text{I}_{0.85}\text{Br}_{0.15})_3$  perovskite thin films with varying amounts of KBr as a passivator similarly showed a secondary Br 3d contribution at higher binding energy, which was assigned to KBr.<sup>45</sup> Despite different cations, one would expect similar

trends in the binding energy shift for KBr and  $\text{CsBr}$ . We therefore assign the new Cs and Br species to  $\text{CsBr}$ . However, higher quantities of  $\text{CsBr}$  are formed significantly quicker on the single crystal surface investigated here than on the thin film surface studied previously. In this context it is important to note that small, non-perovskite Cs 4d and Br 3d contributions were observed from the start, which contrasts our previous study. This initial  $\text{CsBr}$  presence could explain why the degradation is accelerated, as the initial  $\text{CsBr}$  may act as a catalyst for further degradation. The  $\text{CsBr}$  layer stabilizes after receiving around  $5 \times 10^{17}$  photons  $\text{cm}^{-2}$  (Fig. 5B). Alongside  $\text{CsBr}$ , traces of metallic lead resulting from the degradation to  $\text{PbBr}_2$  on the surface were detected. Overall, the following degradation reactions proposed:



Finally, our findings on the stability comparison of the  $\text{MA}^+$ -series are summarized in Fig. 6. Considering the different degradation mechanism of each perovskite composition, the predominant reaction was used to show the percentage of remaining perovskite as a function of X-ray fluence. For  $\text{MAPbI}_3$ , the percentage of N relative to total lead was used, while the percentage  $\text{Pb}^{2+}$  relative to total lead was used for  $\text{MAPbBr}_3$  and  $\text{MAPbCl}_3$ . Given the complexity of the  $\text{FAPbBr}_3$  and  $\text{CsPbBr}_3$  degradation, selecting one core level as the main indicator of the remaining perovskite was not feasible. Consequently, these two compositions were excluded from this figure.

Comparing all investigated compositions in the  $\text{MA}^+$ -series,  $\text{MAPbCl}_3$  shows the quickest degradation, and therefore lowest stability against X-rays, even under low flux conditions (LF). However, the experiment was stopped early in comparison due to the rapid accumulation of metallic lead on the surface, limiting the discussion of long-term behavior. The potential for longer experiments to study the evolution of  $\text{MAPbCl}_3$  remains open for exploration. Continuing with  $\text{MAPbBr}_3$ , a similar but slower behavior as for  $\text{MAPbCl}_3$  under LF conditions can be found with the formation of metallic lead and a continuous decline in  $\text{Pb}^{2+}$  content. In comparison,  $\text{MAPbBr}_3$



**Fig. 6** Percentage of perovskite remaining at the surface as function of X-ray fluence ( $\text{photons cm}^{-2}$ ). Determined by the percentage of  $\text{N}^+/\sum\text{Pb}$  in  $\text{MAPbI}_3$  and the percentage of  $\text{Pb}^{2+}/\sum\text{Pb}$  in  $\text{MAPbBr}_3$  and  $\text{MAPbCl}_3$ .





retained 93% of its perovskite composition after exposure to  $1.8 \times 10^{17}$  photons  $\text{cm}^{-2}$  while  $\text{MAPbCl}_3$  retained only 89% after the same X-ray exposure.

In contrast,  $\text{MAPbI}_3$  behaved differently, *i.e.* formation of only little metallic lead despite exposure to higher photon flux (HF). Furthermore, its degradation mechanism was characterized by a consistent loss of the organic cation together with iodide, suggesting MAI radiolysis as a main degradation reaction. In contrast,  $\text{MAPbBr}_3$  involved multiple processes, including significant formation of  $\text{Pb}^0$  and ion migration. Comparable amounts of N and  $\text{Pb}^{2+}$  disappear for  $\text{MAPbI}_3$  and  $\text{MAPbBr}_3$ , respectively, after receiving  $1.8 \times 10^{17}$  photons  $\text{cm}^{-2}$ . Therefore, both compositions show similarly strong X-ray degradation but with different degradation paths. Radiolysis of  $\text{MA}^+$  is slowed down in  $\text{MAPbBr}_3$  compared to  $\text{MAPbI}_3$ , while  $\text{MAPbBr}_3$  forms significantly larger amounts of  $\text{Pb}^0$ . The difference in radiolysis could be related to an enhanced stability of  $\text{MABr}$  compared to  $\text{MAI}$ .<sup>46,47</sup> Ion migration is in general observed in different perovskite compositions.<sup>39–41</sup> McGovern *et al.* compared the halide migration ability of  $\text{MAPbI}_3$  and  $\text{MAPbBr}_3$  thin films and found similar activation energies for bromide migration in  $\text{MAPbBr}_3$  as for iodide ions in  $\text{MAPbI}_3$ .<sup>47</sup> Furthermore, they found that  $\text{MA}^+$  ion migration is hindered in  $\text{MAPbBr}_3$  but not in  $\text{MAPbI}_3$ .<sup>47</sup> In contrast, our experiments give direct evidence of ion migration in  $\text{MAPbBr}_3$  induced by X-ray irradiation of both  $\text{Br}^-$  and  $\text{MA}^+$  through initial increase of both ion concentrations. On the other hand, ion migration is not directly observed for  $\text{MAPbI}_3$ , as the intensity decrease of the  $\text{MA}^+$  and  $\text{I}^-$  signals are most likely due to radiolysis. However, this does not exclude that ion migration also occurs in the  $\text{MAPbI}_3$ , but if so, it is a slower process than radiolysis. Finally, the degradation of  $\text{MAPbBr}_3$  slowed down compared to  $\text{MAPbI}_3$  after receiving  $3 \times 10^{17}$  photons  $\text{cm}^{-2}$  resulting in a reduction to 80%  $\text{Pb}^{2+}$  for  $\text{MAPbBr}_3$  while the intensity of N in  $\text{MAPbI}_3$  decreased to 60% after receiving  $8.2 \times 10^{17}$  photons  $\text{cm}^{-2}$ .

## Conclusions

We successfully synthesized, *in situ* cleaved and characterized several lead halide perovskite single crystal compositions ( $\text{MAPbI}_3$ ,  $\text{MAPbCl}_3$ ,  $\text{MAPbBr}_3$ ,  $\text{FAPbBr}_3$ ,  $\text{CsPbBr}_3$ ). The study focused on their X-ray stability and different degradation mechanisms have been found. Distinct degradation pathways were identified for each material. In this context,  $\text{MAPbI}_3$  and  $\text{CsPbBr}_3$  form only small amounts of metallic lead, which stabilized during further X-ray radiation.  $\text{MAPbCl}_3$ ,  $\text{MAPbBr}_3$ ,  $\text{FAPbBr}_3$  formed substantial quantities of metallic lead, which continued to increase during the experiment. Additionally, our study revealed that while  $\text{MAPbI}_3$  and  $\text{MAPbCl}_3$  showed one main degradation pathway (radiolysis resulting in  $\text{PbI}_2$  and  $\text{Pb}^0$  formation, respectively),  $\text{MAPbBr}_3$ ,  $\text{FAPbBr}_3$  and  $\text{CsPbBr}_3$  underwent multiple parallel degradation processes. These included the formation of metallic lead, the formation of volatile and solid degradation products and ion migration. Specifically, the degradation of  $\text{MAPbBr}_3$ ,  $\text{FAPbBr}_3$  and  $\text{CsPbBr}_3$

involved the migration of both, the monovalent cation and the halide. Furthermore, our in-depth analysis revealed peak broadening in the cesium and bromide core levels of  $\text{CsPbBr}_3$ , which was attributed to the formation of  $\text{CsBr}$ , which was potentially catalyzed by the presence of  $\text{CsBr}$  on the cleaved surface. Lastly, the  $\text{MA}^+$ -series was compared to one another regarding their remaining perovskite after X-ray exposure. Using our experimental settings and exposure time,  $\text{MAPbCl}_3$  exhibited lowest stability against X-rays in the same fluence range compared to the other two crystals followed by  $\text{MAPbI}_3$  and  $\text{MAPbBr}_3$ . Comparing  $\text{MAPbI}_3$  and  $\text{MAPbBr}_3$ , both are similarly unstable under X-rays, but the dominating reactions are different.  $\text{MAPbBr}_3$  is more resistant to  $\text{MA}^+$  radiolysis but presents an easier formation of  $\text{Pb}^0$ .  $\text{MAPbI}_3$ , on the other hand, shows only little  $\text{Pb}^0$  formation and its degradation is dominated by MAI radiolysis.

## Conflicts of interest

There are no conflicts to declare.

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## References

- 1 S. Porwal, D. Kumar, S. Ghosh, S. Kansal, S. Priya, A. Chandra and T. Singh, in *Low-Dimensional Halide Perovskites*, Elsevier, 2023, pp. 301–333.
- 2 X. Shen, K. Kang, Z. Yu, W. H. Jeong, H. Choi, S. H. Park, S. D. Stranks, H. J. Snaith, R. H. Friend and B. R. Lee, *Joule*, 2023, 7, 272–308.
- 3 J. Zhuang, J. Wang and F. Yan, *Nano-Micro Lett.*, 2023, 15, 84.
- 4 Z.-Y. Chen, N.-Y. Huang and Q. Xu, *Coord. Chem. Rev.*, 2023, 481, 215031.
- 5 S. Jia, Y. Xiao, M. Hu, X. He, N. Bu, N. Li, Y. Liu, Y. Zhang, J. Cui, X. Ren, K. Zhao, M. Liu, S. Wang, N. Yuan, J. Ding, Z. Yang and S. Liu, *Adv. Mater. Technol.*, 2022, 7, 2100908.



- 6 W. Delmas, S. Erickson, J. Arteaga, M. Woodall, M. Scheibner, T. S. Krause, K. Crowley, K. T. VanSant, J. M. Luther, J. N. Williams, J. McNatt, T. J. Peshek, L. McMillon-Brown and S. Ghosh, *Adv. Energy Mater.*, 2023, **13**, 2203920.
- 7 X. He, Y. Deng, D. Ouyang, N. Zhang, J. Wang, A. A. Murthy, I. Spanopoulos, S. M. Islam, Q. Tu, G. Xing, Y. Li, V. P. Dravid and T. Zhai, *Chem. Rev.*, 2023, **123**, 1207–1261.
- 8 T. J. Jacobsson, S. Svanström, V. Andrei, J. P. H. Rivett, N. Kornienko, B. Philippe, U. B. Cappel, H. Rensmo, F. Deschler and G. Boschloo, *J. Phys. Chem. C*, 2018, **122**, 13548–13557.
- 9 P. Schulz, D. Cahen and A. Kahn, *Chem. Rev.*, 2019, **119**, 3349–3417.
- 10 A. García-Fernández, B. Kammlander, S. Riva, D. Kühn, S. Svanström, H. Rensmo and U. B. Cappel, *Inorg. Chem.*, 2023, **62**, 15412–15420.
- 11 B. Kammlander, S. Svanström, D. Kühn, F. O. L. Johansson, S. Sinha, H. Rensmo, A. G. Fernández and U. B. Cappel, *Chem. Commun.*, 2022, **58**, 13523–13526.
- 12 S. Svanström, A. García-Fernández, T. J. Jacobsson, I. Bidermane, T. Leitner, T. Sloboda, G. J. Man, G. Boschloo, E. M. J. Johansson, H. Rensmo and U. B. Cappel, *ACS Mater. Au*, 2022, **2**, 301–312.
- 13 R. Larciprete, A. Agresti, S. Pescetelli, H. Pazniak, A. Liedl, P. Lacovig, D. Lizzit, E. Tosi, S. Lizzit and A. Di Carlo, *Materials*, 2021, **14**, 3954.
- 14 J. Liang, X. Hu, C. Wang, C. Liang, C. Chen, M. Xiao, J. Li, C. Tao, G. Xing, R. Yu, W. Ke and G. Fang, *Joule*, 2022, **6**, 816–833.
- 15 Y. Guan, C. Zhang, Z. Liu, Y. Zhao, A. Ren, J. Liang, F. Hu and Y. S. Zhao, *Adv. Mater.*, 2022, **34**, 2203201.
- 16 L. De Marco, G. Nasti, A. Abate and A. Rizzo, *Solar RRL*, 2022, **6**, 2101085.
- 17 Z. Chen, B. Turedi, A. Y. Alsalloum, C. Yang, X. Zheng, I. Gereige, A. Alsaggaf, O. F. Mohammed and O. M. Bakr, *ACS Energy Lett.*, 2019, **4**, 1258–1259.
- 18 J. Hieulle, D. Son, A. Jamshaid, X. Meng, C. Stecker, R. Ohmann, Z. Liu, L. K. Ono and Y. Qi, *Adv. Funct. Mater.*, 2023, **33**, 2211097.
- 19 V. V. Ozerova, N. A. Emelianov, D. P. Kiryukhin, P. P. Kushch, G. V. Shilov, G. A. Kichigina, S. M. Aldoshin, L. A. Frolova and P. A. Troshin, *J. Phys. Chem. Lett.*, 2023, **14**, 743–749.
- 20 W.-C. Lin, W.-C. Lo, J.-X. Li, Y.-K. Wang, J.-F. Tang and Z.-Y. Fong, *npj Mater. Degrad.*, 2021, **5**, 13.
- 21 A. G. Boldyreva, L. A. Frolova, I. S. Zhidkov, L. G. Gutsev, E. Z. Kurmaev, B. R. Ramachandran, V. G. Petrov, K. J. Stevenson, S. M. Aldoshin and P. A. Troshin, *J. Phys. Chem. Lett.*, 2020, **11**, 2630–2636.
- 22 V. Milotti, S. Cacovich, D. R. Ceratti, D. Ory, J. Barichello, F. Matteocci, A. Di Carlo, P. M. Sheverdyeva, P. Schulz and P. Moras, *Small Methods*, 2023, **7**, 2300222.
- 23 S. Svanström, A. García Fernández, T. Sloboda, T. J. Jacobsson, H. Rensmo and U. B. Cappel, *Phys. Chem. Chem. Phys.*, 2021, **23**, 12479–12489.
- 24 A. García-Fernández, S. Svanström, C. M. Sterling, A. Gangan, A. Erbing, C. Kamal, T. Sloboda, B. Kammlander, G. J. Man, H. Rensmo, M. Odelius and U. B. Cappel, *Small*, 2022, **18**, 2106450.
- 25 M. Ghasemi, S. Yuan, J. Fan, B. Jia and X. Wen, *J. Mater. Chem. A*, 2023, **11**, 3822–3848.
- 26 K. Motoki, Y. Miyazawa, D. Kobayashi, M. Ikegami, T. Miyasaka, T. Yamamoto and K. Hirose, *J. Appl. Phys.*, 2017, **121**, 085501.
- 27 G. Armaroli, L. Ferlauto, F. Lédée, M. Lini, A. Ciavatti, A. Kovtun, F. Borgatti, G. Calabrese, S. Milita, B. Fraboni and D. Cavalcoli, *ACS Appl. Mater. Interfaces*, 2021, **13**, 58301–58308.
- 28 H. Syafutra, J.-H. Yun, Y. Yoshie, M. Lyu, S. N. Takeda, M. Nakamura, L. Wang and M.-C. Jung, *Nanomaterials*, 2020, **10**, 1253.
- 29 C. Wang, B. R. Ecker, H. Wei, J. Huang and Y. Gao, *J. Phys. Chem. C*, 2018, **122**, 3513–3522.
- 30 A. Preobrajenski, A. Generalov, G. Öhrwall, M. Tchapyguine, H. Tarawneh, S. Appelfeller, E. Frampton and N. Walsh, *J. Synchrotron Radiat.*, 2023, **30**, 831–840.
- 31 H. Xie, X. Liu, L. Lyu, D. Niu, Q. Wang, J. Huang and Y. Gao, *J. Phys. Chem. C*, 2016, **120**, 215–220.
- 32 B. Conings, J. Drijkoningen, N. Gauquelin, A. Babayigit, J. D'Haen, L. D'Olieslaeger, A. Ethirajan, J. Verbeeck, J. Manca, E. Mosconi, F. De Angelis and H.-G. Boyen, *Adv. Energy Mater.*, 2015, **5**, 1500477.
- 33 B. Philippe, B.-W. Park, R. Lindblad, J. Oscarsson, S. Ahmadi, E. M. J. Johansson and H. Rensmo, *Chem. Mater.*, 2015, **27**, 1720–1731.
- 34 J. Liang, X. Hu, C. Wang, C. Liang, C. Chen, M. Xiao, J. Li, C. Tao, G. Xing, R. Yu, W. Ke and G. Fang, *Joule*, 2022, **6**, 816–833.
- 35 S. Ullah, J. Wang, P. Yang, L. Liu, S.-E. Yang, T. Xia, H. Guo and Y. Chen, *Mater. Adv.*, 2021, **2**, 646–683.
- 36 J. D. McGettrick, K. Hooper, A. Pockett, J. Baker, J. Troughton, M. Carnie and T. Watson, *Mater. Lett.*, 2019, **251**, 98–101.
- 37 X. Dai, C. Fei, P. Kandlakunta, L. Zhao, Z. Ni, L. R. Cao and J. Huang, *IEEE Trans. Nucl. Sci.*, 2022, **69**, 1850–1856.
- 38 B. Philippe, B.-W. Park, R. Lindblad, J. Oscarsson, S. Ahmadi, E. M. J. Johansson and H. Rensmo, *Chem. Mater.*, 2015, **27**, 1720–1731.
- 39 E. J. Juarez-Perez, L. K. Ono, I. Uriarte, E. J. Cocinero and Y. Qi, *ACS Appl. Mater. Interfaces*, 2019, **11**, 12586–12593.
- 40 E. Bi, Z. Song, C. Li, Z. Wu and Y. Yan, *Trends Chem.*, 2021, **3**, 575–588.
- 41 L. Chen, H. Wang, W. Zhang, F. Li, Z. Wang, X. Wang, Y. Shao and J. Shao, *ACS Appl. Mater. Interfaces*, 2022, **14**, 10917–10926.
- 42 L. McGovern, I. Koschany, G. Grimaldi, L. A. Muscarella and B. Ehrler, *J. Phys. Chem. Lett.*, 2021, **12**, 2423–2428.
- 43 A. F. Akbulatov, L. A. Frolova, N. N. Dremova, I. Zhidkov, V. M. Martynenko, S. A. Tsarev, S. Yu. Luchkin, E. Z. Kurmaev, S. M. Aldoshin, K. J. Stevenson and P. A. Troshin, *J. Phys. Chem. Lett.*, 2020, **11**, 333–339.
- 44 J. A. McLeod and L. Liu, *J. Phys. Chem. Lett.*, 2018, **9**, 2411–2417.
- 45 Z. Andaji-Garmaroudi, M. Abdi-Jalebi, F. U. Kosasih, T. Doherty, S. Macpherson, A. R. Bowman, G. J. Man,



- U. B. Cappel, H. Rensmo, C. Ducati, R. H. Friend and S. D. Stranks, *Adv. Energy Mater.*, 2020, **10**, 2002676.
- 46 L. Shi, M. P. Bucknall, T. L. Young, M. Zhang, L. Hu, J. Bing, D. S. Lee, J. Kim, T. Wu, N. Takamure, D. R. McKenzie, S. Huang, M. A. Green and A. W. Y. Ho-Baillie, *Science*, 2020, **368**, eaba2412.
- 47 L. McGovern, M. H. Futscher, L. A. Muscarella and B. Ehrler, *J. Phys. Chem. Lett.*, 2020, **11**, 7127–7132.

