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Elucidating the effect of Ag interlayer formation on the intrinsic mechanical properties of free-standing ITO/Ag/ITO thin films†

Seung Jin Oh,^a Sangmin Lee,^a Kyung Cheol Choi,^b Jeong Hyun Kwon^{b,*c}
and Taek-Soo Kim^{b,*a}

The lack of characterization of the mechanical behavior of brittle indium tin oxide (ITO)-based electrodes has been a core issue for the development of advanced transparent and flexible electronics. In this study, we report the intrinsic mechanical properties of ITO/Ag/ITO (IAI) thin films measured by a free-standing tensile testing method for robust transparent electrodes. To understand the effect of Ag interlayer formation on the mechanical properties, Ag thickness is controlled (8–50 nm) by considering the stage of the film formation. Based on the tensile test, IAI thin films with a film-likely formed Ag interlayer (12–14 nm) exhibited outstanding mechanical robustness compared with island-(8 nm) or film-type (>25 nm) Ag interlayers. Notably, the most superior elongation and tensile strength ($0.57 \pm 0.08\%$ and 514.3 ± 71.5 MPa) are demonstrated for 14 nm-thick Ag interlayer, which is two-fold higher than pristine ITO thin films (0.27% and 264.6 MPa). The study provides intrinsic mechanical property values and fundamental insights into the fracture mechanisms that govern the mechanical behavior of oxide/metal/oxide multilayer thin films, which is valuable for the development of flexible electronics.

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Introduction

Based on their favorable electrical ($<10^{-4}$ Ω cm), optical (>90%), and chemical stability properties, indium tin oxide (ITO) thin films have been used as transparent electrodes for electronic devices.^{1–5} However, the relatively high sheet resistance and brittleness of ITO thin films are critical drawbacks for their application in advanced flexible electronics.^{6–8} Recently, oxide/metal/oxide thin films have been extensively studied as alternative transparent electrodes owing to their outstanding performance, such as low sheet resistance, enhanced flexibility, and roll-to-roll process compatibility. In particular, ITO/Ag/ITO (IAI) multilayered films have attracted significant attention in the fields of flexible organic light-emitting diodes (OLEDs), flexible organic solar cells, thin-film transistors (TFTs), touch screens, semiconductors, and aeronautical engineering.^{9–14}

Besides, a low-temperature process of IAI thin films is critically required for optoelectronic applications, which is achievable using sputtering and evaporation methods.^{15,16} The IAI multilayered structure satisfies all the requirements for electronics with a reduced amount of ITO thin films combined with tens of nanometer-thick Ag interlayers, whereas the typical thickness of the ITO thin films ranges between 150 and 250 nm.^{17,18}

Several studies have focused on optimizing the electrical and optical properties of ITO/Ag/ITO thin films. Jeong *et al.* reported that a Ag thickness of 16 nm exhibited the best transparency (86.5%) versus conductivity ($4.4 \Omega \text{ sq}^{-1}$) with a figure of merit of $49.4 \times 10^{-3} \Omega^{-1}$.¹⁹ Recently, Wang *et al.* revealed that ~ 10 nm of Ag incorporation showed the optimal combination of electrical ($6.1 \Omega \text{ sq}^{-1}$) and optical properties (90.8%) with a higher figure of merit of $62.3 \times 10^{-3} \Omega^{-1}$.⁸ However, the tendency of catastrophic fracture of the brittle IAI thin films remains the cause of the sudden failure of devices. Studies on the intrinsic mechanical behaviors of IAI thin films have rarely been reported because direct tensile testing is restricted owing to the challenges in handling nano-thick and brittle IAI thin films. Therefore, previous studies have been limited to comparing the bending stability of ITO and IAI multilayered thin films on plastic substrates, which was demonstrated to be affected by residual stress.^{16,20–22} Additionally, to the best of our knowledge,

^a Department of Mechanical Engineering, Korea Advanced Institute of Science and Technology (KAIST), Daejeon 34141, Republic of Korea. E-mail: tskim1@kaist.ac.kr

^b School of Electrical Engineering, Korea Advanced Institute of Science and Technology (KAIST), Daejeon 34141, Republic of Korea

^c Department of Display and Semiconductor Engineering, SunMoon University, Asan 31460, Republic of Korea. E-mail: love6539@sunmoon.ac.kr

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the effect of the Ag interlayer on the mechanical properties of IAI thin films has not been elucidated thus far. Therefore, exploring their intrinsic mechanical properties and fracture mechanism is essential for expanding the applications of transparent conductive electrodes (TCEs).

In this study, we report the tensile and fracture behaviors of low-temperature-processed IAI multilayer thin films, as a function of the thickness of the Ag interlayer. The IAI thin films with Ag thicknesses of 8, 10, 11, 12, 13, 14, 25, and 50 nm are prepared. Each IAI structure is tested in the free-standing state to confirm whether the regulation of the Ag interlayer thickness can improve the mechanical stability and simultaneously enhance the optical and electrical properties of the thin films. Previous literature reported the aging effect of the IAI thin films when the film is exposed to oxygen, air, and light illumination.^{23–25} To exclude the aging effect, whole samples were stored in a vacuum desiccator (approximately -0.1 mPa, 0% relative humidity) and the test was performed within 2 weeks after the deposition. Therefore, this study aims to provide a foundation for a better understanding of the mechanical behavior of oxide/metal/oxide multilayers and new opportunities for the design of flexible electronics and mechanical nanocomposites.

Results and discussion

Preparation of free-standing IAI thin films

The preparation of a free-standing IAI thin film for tensile testing is shown in Fig. 1. As shown in Fig. 1a, the IAI thin film

was deposited onto a Cu/SiO₂/Si square wafer (35×35 mm). For ITO deposition, the thicknesses of the bottom and top ITO thin films were symmetrically defined as 45 nm to achieve effective antireflection effects, which exhibited the highest red/green/blue (RGB) average transmittance based on an optical simulation (Fig. S1, ESI†).^{19,26,27} To determine the effect of the Ag interlayer formation on the mechanical properties of the IAI thin films, the Ag thicknesses were varied as 8, 10, 11, 12, 13, 14, 25, and 50 nm. The IAI thin films are denoted as IAI(8), IAI(10), IAI(11), IAI(12), IAI(13), IAI(14), IAI(25), and IAI(50), respectively. The XPS depth profiles show that the IAI thin films were successfully fabricated as intended with symmetric ITO thin films at the top and bottom of the structure (Fig. S2, ESI†). A stable Ag interlayer was formed between the well-defined symmetrical ITO thin films with a proper stoichiometric composition without any detectable interfacial reaction or diffusion. Additionally, the as-deposited ITO thin films were amorphous, and the Ag interlayer was polycrystalline and dominantly oriented in the (111) direction (Fig. S3, ESI†).^{28,29} After the deposition of the IAI thin film, the as-deposited IAI/Cu layers were patterned into a dog-bone shape using a femtosecond laser, and subsequently free-standing tensile tests were performed (Fig. 1c).^{30,31} To obtain a free-standing IAI thin film, the IAI/Cu film was peeled off from the SiO₂/Si substrate with the assistance of the surface tension of water (Fig. 1d). A slight force was addressed using a commercial tape to initiate and propagate the delamination between the Cu/SiO₂ interface (Fig. S4, ESI†). The overall transfer process took less than



Fig. 1 Schematics for the preparation of free-standing IAI thin films. (a) Deposition of Cu (1 μm), ITO (45 nm), Ag (8–50 nm), and ITO (45 nm) thin films on a SiO₂/Si wafer. (b) Photos of the fabricated ITO and IAI thin films as a function of Ag thickness (scale bar = 10 mm). (c) Femtosecond laser patterning of the IAI/Cu layer to fabricate the dog-bone-shaped specimen. (d) Transfer of the IAI/Cu layer onto the water surface. (e) Etching of the sacrificial Cu layer. (f) Prepared free-standing IAI thin films for the tensile test.

1 minute. Owing to the weak interfacial adhesion energy of Cu/SiO₂ in water, the IAI/Cu film was carefully transferred onto the water surface without any damage.³² Subsequently, the sacrificial Cu layer was etched with a 0.035 M ammonium persulfate solution (Fig. 1e). Finally, the free-standing IAI thin film was obtained by substituting the Cu etchant with deionized water to prevent any possible chemical damage to the IAI thin film (Fig. 1f).

Effect of Ag interlayer thickness on mechanical properties

Photographs of the tensile test procedure are shown in Fig. 2a. By utilizing the high surface tension of water, free-standing IAI thin films are successfully tested without disrupting deformations, such as wrinkles and deflection. Additionally, the effect of residual stress was excluded because the IAI thin films on the water surface could freely release the substrate constraint.³³ Moreover, the effect of surface roughness was negligible because there was no significant difference in root mean square roughness (R_q) with respect to the Ag interlayer thickness (Fig. S5, ESI†). Gripping of the free-standing and dog-bone-shaped IAI thin films was accomplished by van der Waals adhesion between the polydimethylsiloxane (PDMS) and IAI thin films. After gripping, a tensile test was conducted by horizontal loading using a high-resolution linear actuator at a strain rate of approximately $5.8 \times 10^{-5} \text{ s}^{-1}$ (Fig. 2b and Fig. S6, ESI†). During the tensile test, the digital image correlation

(DIC) method was used to precisely measure the applied strain by tracking the designated patterns using a charge-coupled device (CCD) camera.^{34,35} At the last moment of the tensile test, the fracture of the IAI thin films occurred catastrophically according to the crack branching phenomenon (that is, multiple shear bands), which represents the typical failure of brittle materials.^{31,36} This tendency of brittle cracking of the IAI thin films may result in a sudden increase in their electrical resistivity, thereby causing unexpected failure of the devices. The characteristics of completely brittle failure were also observed in the stress–strain curves of the IAI thin films, as a function of the Ag thickness (Fig. 2c). Independent of the Ag interlayer thickness, the curves exhibited a linear increase in the applied stress according to the applied strain, and plastic deformation was not detected. Therefore, considering the precise intrinsic elongation, it is essential to conservatively prevent the breakage of brittle IAI thin films for the design of flexible electronics.

To date, the intrinsic Young's modulus and elongation of IAI thin films have not yet been reported. Based on the stress–strain curves, Table S1 (ESI†) lists the intrinsic Young's modulus, elongation, and resultant tensile strength of the IAI thin films depending on the Ag thickness. As shown in Fig. 2d, the Young's modulus of the IAI thin films with Ag thicknesses ranging from 8 to 13 nm increased from $74 \pm 6 \text{ GPa}$ to $95 \pm 4 \text{ GPa}$ and decreased for 14, 25, and 50 nm to $93 \pm 2 \text{ GPa}$, $78 \pm 4 \text{ GPa}$, and $63 \pm 10 \text{ GPa}$, respectively. Compared with the 90 nm-thick reference ITO thin

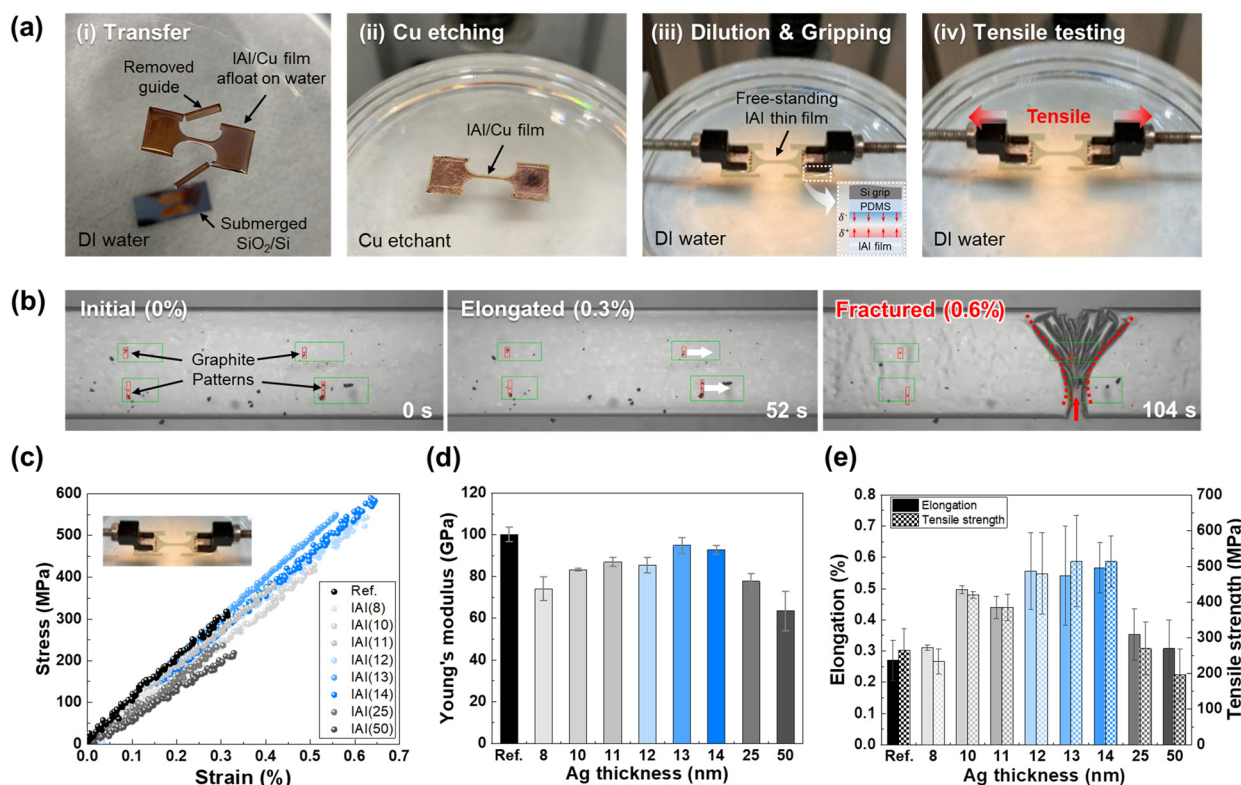


Fig. 2 Tensile testing of free-standing IAI thin films. (a) Photos of the tensile testing procedure for IAI thin films. (b) CCD camera images during the tensile loading. (c) Representative stress–strain curves of the ITO (Reference) and IAI thin films as a function of Ag thickness. (d) Intrinsic Young's modulus of the IAI thin films. (e) Intrinsic elongation and tensile strength of the IAI thin films.

film, the Young's modulus values of the fabricated IAI films were relatively low because of the premature formation of the Ag interlayer and its lower Young's modulus (< 60 GPa)³⁷ than that of the ITO thin films (103.2 GPa). Interestingly, the intrinsic Young's moduli of IAI thin films did not obey the rule of mixture, which support that our reporting values are worthy for the accurate mechanical analysis of flexible electronics (Fig. S7, ESI†). Additionally, the rule of mixture results indicates that the films' compactness maximized for the IAI(13) and IAI(14) thin films. As shown in Fig. 2e, the elongation of the IAI thin films with Ag thicknesses ranging from 8 to 14 nm increased from $0.31 \pm 0.01\%$ to $0.57 \pm 0.08\%$, and decreased for 25 and 50 nm to $0.35 \pm 0.08\%$ and $0.31 \pm 0.09\%$, respectively. The highest elongation was recorded for IAI(14) thin films ($0.57 \pm 0.08\%$). Thus, the elongation of the fabricated IAI thin films was superior, independent of the Ag thickness, compared with that of the reference ITO thin film ($0.27 \pm 0.06\%$). The tensile strength of the IAI thin films showed a similar trend to the elongation as a function of Ag thickness. These results suggest that IAI thin films with an optimized Ag thickness of 12–14 nm, as an alternative to conventional ITO thin films, can significantly improve the mechanical stretchability of flexible electronics. Additionally, the measured intrinsic mechanical properties enable the transition from IAI multilayer modeling to equivalent single-layer modeling for accurate finite element analysis (FEA) modeling to describe the mechanical behavior of IAI thin films.^{38,39}

To gain a better insight into the role of the Ag interlayer on the mechanical properties, structural characterization of the Ag/ITO (AI) and IAI thin films was conducted using SEM. Fig. 3a shows the surface morphologies of the AI thin films for the as-deposited Ag interlayer and the layer after exposure to the argon/oxygen (Ar/O₂) plasma atmosphere, as a function of the Ag thickness. For the as-deposited Ag layer, the film formation followed the multistep processes of metal growth on oxides: (1) the nucleation of discrete nano-sized clusters (islands), followed by coalescence into large and irregular clusters, (2) the formation of a premature film-like network at a specific percolation threshold, and (3) the transition to a continuous film with increasing metal thickness.⁴⁰ To understand the Ag film formation, a thermodynamic criterion can be introduced, and the growth mode of Ag is determined by the following energy difference (eqn (1)):

$$\Delta\gamma = -\gamma_s + \gamma_i + \gamma_f \quad (1)$$

where $\Delta\gamma$ is the energy difference between the substrate (ITO) and film (Ag), γ_s is the surface free energy of the substrate, γ_f is the free energy of the film, and γ_i is that of the interface between the substrate and film. Generally, in the case of metal films, the clusters are more bound to each other than to the substrates (that is, $\Delta\gamma > 0$). Therefore, the Volmer–Weber (island) growth mode is observed for the Ag thin films.^{41–44}



Fig. 3 SEM images of the surface morphologies of Ag interlayer and fractured surfaces of the IAI thin films as a function of Ag thickness. (a) Surface morphologies of the as-deposited and plasma-treated Ag layer on the bottom ITO surface. Plasma-treated images showing the effect of Ar/O₂ plasma during the top ITO deposition (scale bar = 500 nm). (b) Top, tilted, and cross-sectional images of the fractured surface of IAI thin films. Defective regions in the top ITO and Ag interlayer are exposed at the tilted images.

For the images of plasma-treated Ag interlayer, the diffusion of Ag was confirmed after exposure to the ITO sputtering atmosphere without direct collision of ITO clusters. Therefore, the change in the Ag morphology originates from the Ar/O₂ plasma, which enhances the wetting property of the exposed bottom ITO surface. The Ar/O₂ plasma in the partial O₂ environment increases the surface energy of the ITO surface, thereby contributing to the formation of a film-like state.^{45,46} With a decrease in γ_i , the Ag atoms interact more strongly with the ITO substrate than with each other. Additionally, the local thermal energy that occurred during the sputtering of the top ITO thin films may have induced additional diffusion and percolation of the Ag clusters. The low-temperature diffusion of the island-type Ag interlayer is accelerated by the small size of Ag nuclei.^{47,48} Based on our observations of the AI films, we can conclude that IAI(8) contains isolated Ag islands, IAI(12–14) contains a film-like Ag interlayer, and IAI(25–50) contains a continuous Ag film interlayer.

It is widely known that structural defects, such as voids, density-deficient boundaries, and pinholes, can significantly impact the fracture behavior of amorphous oxide thin films.^{49–51} Fig. 3b represents the structural characteristics of the IAI thin films as a function of the Ag interlayer thickness. For the IAI(8) thin film, the Ag islands functioned as a point nucleation site for the top ITO, thereby the top ITO formed a

distinct columnar microstructure. Our previous observation using high-angle annular dark field-scanning transmission electron microscope (HAADF-STEM) found that the boundaries of ITO particles are significantly density-deficient.³¹ Therefore, the IAI(8) thin films exhibiting poor mechanical properties due to the island Ag interlayer itself and the defective top ITO boundaries.⁵⁰ For the case of IAI(10–14), the film-like Ag interlayer contributed to the formation of a denser top ITO by functioning as a plane nucleation site. Therefore, the critical defects, which are density-deficient boundaries of the columnar ITO structure, were reduced. Thus, in combination with the densified Ag interlayer and top ITO, the Young's modulus, elongation, and tensile strength were significantly increased. For the IAI(25) and IAI(50) thin films, Young's modulus decreased owing to the increased portion of the Ag film layer. Additionally, elongation was also deteriorated owing to the extended and intensified grain boundaries around the crystallites as the crystallization of the Ag interlayer progressed (Fig. S3, ESI†). Therefore, film-like formation of the Ag interlayer is very important to enhance the mechanical properties of oxide/metal/oxide thin films.

Through the cross-sectional SEM observation, it was difficult to confirm the actual in-plane morphology of the Ag interlayer of IAI thin films. To thoroughly clarify the Ag interlayer formation, the interfacial morphologies of the Ag layer were

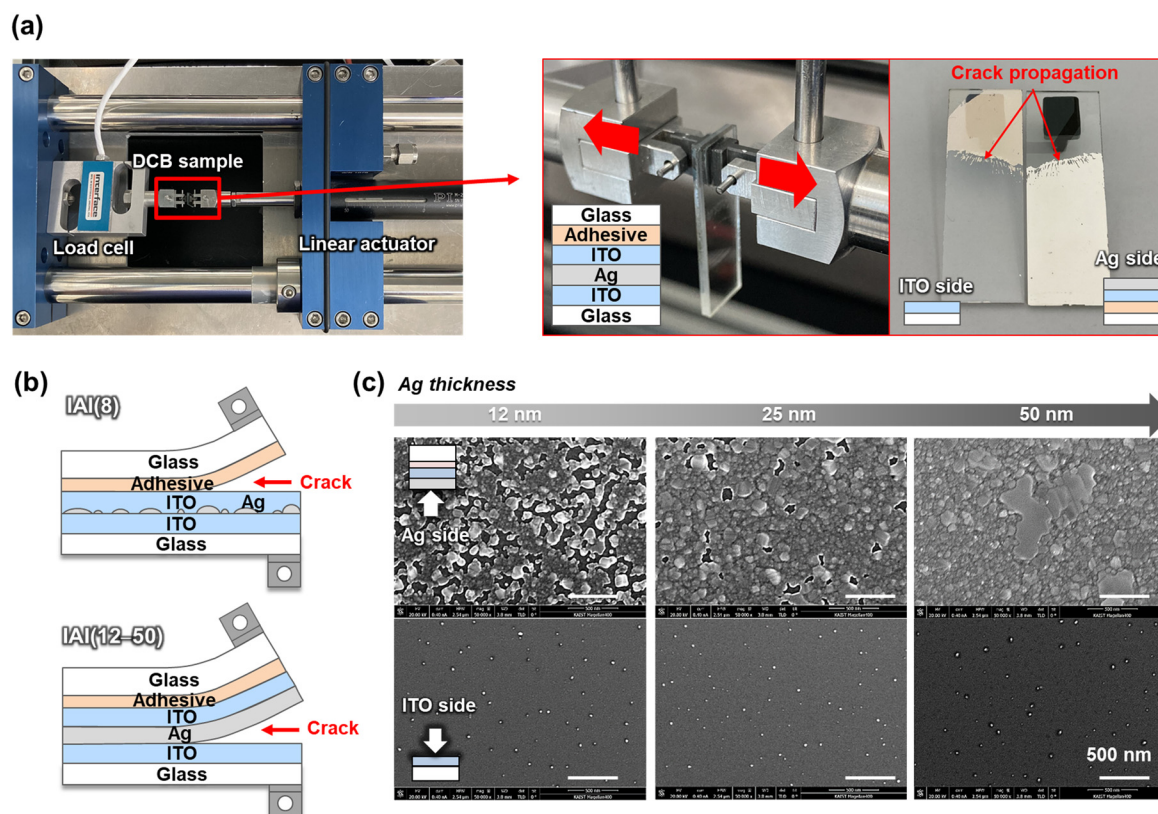


Fig. 4 DCB testing method for the IAI thin films and resultant delaminated surfaces. (a) Photo of the DCB equipment and a representative delaminated IAI(25) specimen. (b) Illustrations of the resultant crack pathway depending on Ag thickness. Interlayers of the IAI(8) specimens were not delaminated and interfacial fracture occurred for the IAI(12–50) specimens. (c) SEM images of the delaminated surface morphologies as a function of Ag thickness.

further characterized by delaminating their interfaces. For the careful separation of the IAI multilayer interfaces, a double cantilever beam (DCB) method was adopted as shown in Fig. 4a. The DCB method could apply mode I loading (opening mode) to the multilayer films, which is advantageous to induce delamination of the weakest interface of the multilayer film. As a result of the DCB test, the Ag interlayer and bottom ITO interface were clearly delaminated. Fig. 4b illustrates the resultant crack pathway for the IAI(8), (12), (25), and (50) specimens under Mode I loading. For the IAI(8) specimen, the crack did not penetrate into the interfaces of the IAI multilayer thin films, resulting in delamination between the epoxy and top ITO interface. Because of the isolated Ag islands, the top ITO particles could cover the exposed bottom ITO surfaces. This may have resulted in cohesive bonding between the bottom and top ITO layers, with good mechanical interlocking. For the IAI(12), (25), and (50) specimens, interfacial delamination was successfully achieved between the film-like Ag interlayer and the bottom ITO surface. The reason for the delamination at the Ag/bottom ITO interface is that the top ITO/Ag interface is strongly bonded by the sputtered high-energy clusters compared with the Ag/bottom ITO interface formed by e-beam evaporation.^{52,53} Fig. 4c shows the delaminated surface morphologies of the Ag interlayer and bottom ITO as a function of Ag thickness. The actual surface morphologies of the Ag interlayer were well coincided with the plasma-treated Ag layer (Fig. 3a), which exhibited better wetting characteristics than that of as-deposited Ag. Thus, we confirmed that the morphology of the as-deposited Ag interlayer was highly influenced by sputtering atmosphere of the top ITO thin film. The results also

implies that the surface treatment on the interlayer could significantly change the intrinsic mechanical properties of the oxide/metal/oxide thin films.

Based on the results of the tensile test, SEM analysis, and DCB test, the function of the Ag interlayer formation in the tensile and fracture behavior of the IAI thin films is proposed. As illustrated in Fig. 5, the Ag islands in IAI(8) thin films act as a point nucleation site for the top ITO, resulting in the formation of a distinct columnar microstructure. In this case, the density-deficient boundaries of the columnar structure of the top ITO serve as critical defect sites, leading to poor mechanical properties of the IAI thin films. For IAI(10–14), the film-like Ag interlayer functions as a plane nucleation site for the top ITO, contributing to the formation of a denser structure. This leads to a reduction in the voids at the boundary of the columnar ITO structure and thus, improves the mechanical properties of the IAI thin films. Additionally, the Ag interlayer covers the boundaries of the ITO particles, reducing stress concentration. Therefore, the densified Ag interlayer and reduced voids at ITO boundaries, in combination, significantly increase the Young's modulus and elongation. Additionally, IAI(25, 50) thin films with a film-type Ag interlayer show a decreased Young's modulus owing to the increased proportion of the Ag layer and deteriorated elongation caused by the formation of grain boundaries as a result of the crystallization of the Ag layer into larger agglomerates. In the future, exploring the mechanical properties of IAI thin films with controlled wetting characteristics of the Ag interlayer would be beneficial for gaining a better understanding of how to design the functional multilayer thin films.



Fig. 5 Schematic of the structural changes in the IAI thin films as a function of Ag interlayer formation. The effect of the Ag interlayer and the associated mechanical properties of the IAI thin films is proposed.

Mechanical, electrical, and optical optimization toward robust thin film electrodes

The optical and electrical properties of the as-prepared IAI thin films are shown in Fig. 6 and Table S2 (ESI†). Fig. 6a shows representative samples of the IAI(8), IAI(12), IAI(25), and IAI(50) thin films deposited onto soda-lime glass substrates. As shown in Fig. 6b and c, the optical transmittance was the highest (83.2%) for an Ag thickness of 12 nm. The sheet resistance of the as-deposited IAI thin films as a function of Ag thickness is plotted in Fig. 6c. The sheet resistance rapidly decreased from $110.5 \pm 4.6 \Omega \text{ sq}^{-1}$ to $1.63 \pm 0.04 \Omega \text{ sq}^{-1}$ as the Ag thickness increased, as previously reported in the literature.¹⁵ The thin films exhibited high optical transmittance ($> 80\%$ in the visible range) and low sheet resistance ($< 15 \Omega \text{ sq}^{-1}$), indicating their potential for use as transparent electrodes in flexible electronic devices. At the initial stage of the Ag film formation, discontinuous Ag islands are formed according to the Volmer–Weber growth mode, which is known to have a negative effect on conductivity owing to the comparable thickness to the electron mean free path (MFP).⁵⁴ As the deposition progresses, these islands grow and percolate to form a semicontinuous/conductive film, thereby the sheet resistance significantly decreases. To select an efficient TCE, the figure of merit is a key parameter that represents the optical and electrical performance of a film. By correlating the transmittance measured at a wavelength of 550 nm and the sheet resistance, the figure of merit of the IAI thin films with varying Ag thicknesses was calculated, as defined by Haacke:

$$\phi = T^{10}/R_s \quad (2)$$

The IAI(12) film exhibited the highest figure of merit ($32.2 \times 10^{-3} \Omega$), which is comparable with reported as-deposited IAI thin films.^{15,55–57} Thus, using the low-temperature process without any substrate heating or post-annealing steps, IAI thin films with a high optical transmittance and low sheet resistance were obtained. Overall, it has been discovered that the film-like formation of the Ag interlayer in IAI thin films plays a crucial role in achieving optimal mechanical, electrical, and optical performance (Fig. 6d).

Experimental section

Materials and sample preparation

Sputtering. For the preparation of the IAI thin films, a 1 μm -thick Cu film was pre-deposited onto a $35 \times 35 \text{ mm}$ SiO_2/Si wafer (300 nm/525 μm) by e-beam evaporation as a sacrificial template. Subsequently, a 45 nm-thick bottom indium tin oxide (ITO) thin film was deposited onto the Cu/ SiO_2/Si wafer using a radio frequency (RF) sputtering system. A 3 inch-diameter target composed of In_2O_3 (90 wt%) and SnO_2 (10 wt%) was used at a distance of 150 mm from the substrate. The RF power, base pressure, total pressure, and oxygen partial pressure were 100 W, $3\text{--}6 \times 10^{-6}$ torr, 5×10^{-3} torr, and 5% (argon flux of 20 sccm), respectively. Subsequently, Ag thin films with thicknesses of 8, 10, 11, 12, 13, 14, 25, and 50 nm were deposited by

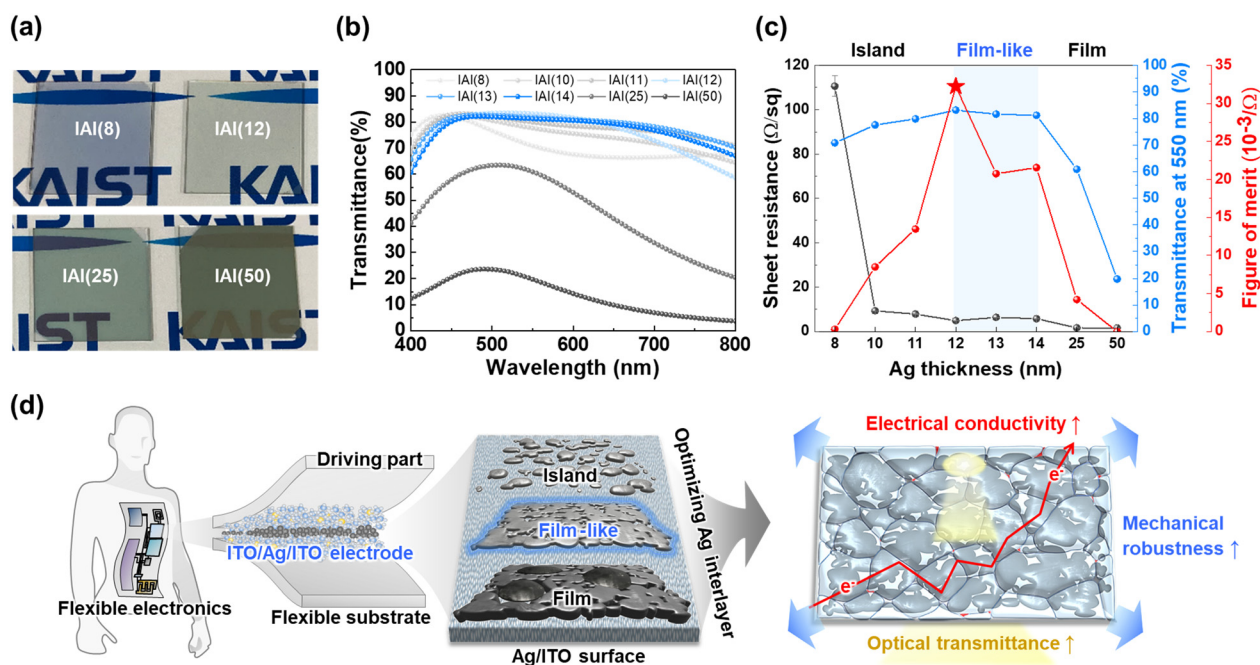


Fig. 6 Optical and electrical characterization for the IAI thin films as a function of Ag thickness. (a) Photograph showing the difference in the transmittance of the IAI(8), IAI(12), IAI(25), and IAI(50) thin films on the glass substrate. (b) UV-Vis spectra of the IAI thin films. The IAI(12–14) thin films exhibited higher optical transmittance among the films. (c) Sheet resistance, representative transmittance at 550 nm, and figure of merit of the IAI thin films with different Ag thicknesses. (d) Insight on optimizing the Ag interlayer formation for robust design of advanced flexible electronics.

thermal evaporation. The actual thicknesses of the 8, 12, 25, and 50 nm Ag thin films measured by SEM were 8.6 ± 0.4 , 12.1 ± 0.8 , 25.1 ± 0.2 , and 48.9 ± 0.5 nm, respectively (Fig. S8, ESI[†]). Finally, symmetrical IAI thin films were prepared by fabricating 45 nm-thick top ITO thin films on the Ag/ITO thin films. During the deposition process, the substrate was maintained at 25 °C. The as-deposited ITO thin films of the IAI thin films were amorphous and the Ag thin films were polycrystalline with a preferred orientation of (111), as characterized by XRD.

Femtosecond laser patterning

The sputtered ITO thin film on the Cu/SiO₂/Si substrate was fabricated in a dog-bone shape using a femtosecond laser (Pharos, L2K Co., Korea). The pulse power (10 mW) and frequency (100 kHz) of the laser were optimized to cut the ITO and sacrificial Cu layers, which was required for the clear separation of the dog-bone-shaped specimen. The total, grip section, and gage lengths of the fabricated dog-bone specimens were 17.2, 5, and 10 mm, respectively. Additionally, the widths of the grip and reduced sections were 5 and 0.8 mm, respectively (Fig. S9, ESI[†]).

Free-standing tensile test

A water surface with a high surface tension (73 mN m^{-1}) and low viscosity (1.002 mPa s^{-1}) was used as a platform to fabricate the IAI thin films in a free-standing state.³⁰ On the left and right sides of the free-standing IAI thin film, a high-resolution load cell (LTS-50GA, KYOWA, Japan) and linear actuator (M-111.1DG, PI, Germany) were placed on an anti-vibration table (Fig. S5, ESI[†]). For the tensile test, load values were acquired using the load cell, and *in situ* strain values were obtained using a DIC system with a charge-coupled device (CCD) camera. After gripping the IAI thin film with PDMS-coated Si grips, the tensile test was conducted at a strain rate of $\sim 5.8 \times 10^{-5} \text{ s}^{-1}$. The tensile test environment was uniformly maintained at a room temperature of 25 °C and relative humidity (RH) of approximately 20%. For the IAI thin films as a function of Ag thickness, 3–5 samples were evaluated among the fabricated 10–20 samples. The representative free-standing tensile test results show highly reproducible curves (Fig. S10, ESI[†]).

Characterization of the ITO thin films

After the tensile test, the IAI thin films were transferred onto a Si substrate to measure the thickness of each film using atomic force microscopy (Park Systems, XE-100, Korea). Before transfer, the IAI thin films were rinsed with deionized water immediately after the tensile test to dilute the residual ammonium persulfate solution. The scanned area ($20 \times 5 \mu\text{m}$) was measured at a scan frequency of 0.4 Hz, and the thickness was determined by measuring the height gap between the Si substrate and IAI thin film. The crystallographic state of the IAI thin films was characterized by X-ray diffractometer (Ultima IV, RIGAKU, Japan) with Cu radiation ($\lambda = 1.5406 \text{ \AA}$). A 2θ measurement with an initial inclination angle of 2° was conducted to amplify the reflection peaks of the IAI film. Additionally, the microstructures of the IAI thin films were characterized by

scanning electron microscopy (Magellan 400, FEI Co., USA) at an accelerating voltage of 10–20 kV. To quantitatively determine the optical and electrical performances of the IAI thin films, the transmittance was characterized by ultraviolet/visible/near infrared (UV-VIS/NIR) spectroscopy (Lambda 1050, PerkinElmer, Korea) and a four-point probe was used to measure the sheet resistance of the IAI thin films depending on the Ag thickness.

Double cantilever beam (DCB) test

The DCB test was performed to evaluate the fracture energy of the IAI interfaces and characterize the surface morphology of the Ag interlayer (Fig. S11, ESI[†]). The DCB specimens were fabricated into a sandwich structure by joining glass substrates on the IAI/glass substrates using a brittle-type epoxy (Epo-Tek 353ND, Epoxy Technology, USA). The epoxy was cured at room temperature ($\sim 25 \text{ }^\circ\text{C}$) for 72 h to prevent thermal damage to low-temperature-processed IAI thin films. A delaminator (Delaminator Adhesion Test System; DTS Company, USA) was used for the DCB test, which consisted of a load cell (SSM-FDH 200 N, Interface, USA), linear actuator (M-235.5DG, PI, Germany), and loading grips. The DCB tests were conducted at room temperature ($\sim 25 \text{ }^\circ\text{C}$) under controlled RH ($\sim 30\%$).

Conclusions

In summary, our findings demonstrate the importance of controlling the Ag interlayer formation in IAI thin films to achieve superior mechanical, structural, optical, and electrical properties. According to the tensile test results, Ag interlayer formation in a film-like structure (12–14 nm) exhibited optimum mechanical properties. The resulting films showed approximately a two-fold increase in both elongation (0.54–0.56%) and tensile strength (480.0–514.3 MPa) compared to the reference ITO thin film. The superior mechanical properties observed in the films can be attributed to the film-like formation of the Ag interlayer, which reduces defects and covers voids between the boundaries of the ITO particles. On the other hand, island-type Ag (8 nm) induces distinct columnar growth of the top ITO with density-deficient boundaries, and the film-type Ag (25, 50 nm) has grain boundaries along the Ag crystallites, which act as critical defects. Consequently, the use of IAI thin films with a film-like Ag interlayer that exceeds the island-type Ag layer is crucial to simultaneously enhance the mechanical robustness, electrical performance, and transparency for a variety of advanced flexible electronic applications.

Author contributions

S. J. O, J. H. K., and T.-S. K conceived and designed the experiments and prepared the manuscript. J. H. K fabricated the ITO/Ag/ITO (IAI) multilayer thin films used in the study. S. J. O conducted the IAI thin film free-standing tensile test and double cantilever beam test. S. J. O and S. L analyzed the

results. All authors discussed the results and commented on the manuscript.

Conflicts of interest

There are no conflicts to declare.

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References

- H. Kim, J. S. Horwitz, G. Kushto, A. Pique, Z. H. Kafafi, C. M. Gilmore and D. B. Chrisey, *J. Appl. Phys.*, 2000, **88**, 6021–6025.
- Y. C. Lin, J. Y. Li and W. T. Yen, *Appl. Surf. Sci.*, 2008, **254**, 3262–3268.
- V. Teixeira, H. N. Cui, L. J. Meng, E. Fortunato and R. Martins, *Thin Solid Films*, 2002, **420**, 70–75.
- C. H. Hsu, Z. X. Zhang, C. Y. Shi, P. H. Huang, W. Y. Wu, D. S. Wu, P. Gao, C. J. Huang, S. Y. Lien and W. Z. Zhu, *J. Mater. Chem. C*, 2022, **10**, 12350–12358.
- C. C. Lee, C. J. Shih, G. Kumar, S. Biring, S. Sen and S. W. Liu, *J. Mater. Chem. C*, 2017, **5**, 12050–12056.
- S. M. Lee, H. W. Koo, T. W. Kim and H. K. Kim, *Surf. Coat. Technol.*, 2018, **343**, 115–120.
- C. Peng, Z. Jia, D. Bianculli, T. Li and J. Lou, *J. Appl. Phys.*, 2011, **109**, 103530.
- H. L. Wang, C. M. Tang, Q. Shi, M. Y. Wei, Y. F. Su, S. S. Lin and M. J. Dai, *Ceram. Int.*, 2021, **47**, 7666–7673.
- K. J. Ko, S. R. Shin, H. B. Lee, E. Jeong, Y. J. Yoo, H. M. Kim, Y. M. Song, J. Yun and J. W. Kang, *Mater. Today Energy*, 2021, **20**, 100704.
- Y. L. Zhang, Z. Liu, C. G. Ji, X. L. Chen, G. F. Hou, Y. L. Li, X. Zhou, X. H. Cui, X. F. Yang, C. C. Ren, D. Liu, L. J. Guo, Y. Zhao and X. D. Zhang, *ACS Appl. Energy Mater.*, 2021, **4**, 6553–6561.
- C. C. Wu, *Rsc. Adv.*, 2018, **8**, 11862–11870.
- Y. Kim, G. J. Jeon, M. K. Lee, S. H. Lee and S. H. K. Park, *ECS Trans.*, 2016, **75**, 247–251.
- J. H. Kwon, J. Park, M. K. Lee, J. W. Park, Y. Jeon, J. Bin Shin, M. Nam, C. K. Kim, Y. K. Choi and K. C. Choi, *ACS Appl. Mater. Interfaces*, 2018, **10**, 15829–15840.
- W. Y. Ji, T. Wang, B. Y. Zhu, H. Zhang, R. Wang, D. D. Zhang, L. Z. Chen, Q. Y. Yang and H. Z. Zhang, *J. Mater. Chem. C*, 2017, **5**, 4543–4548.
- N. Ren, J. Zhu and S. Ban, *AIP Adv.*, 2017, **7**, 055009.
- S. H. Park, S. J. Lee, J. H. Lee, J. Kal, J. Hahn and H. K. Kim, *Org. Electron.*, 2016, **30**, 112–121.
- M. Girtan and B. Negulescu, *Opt. Mater.: X*, 2022, **13**, 100122.
- C. Guillen and J. Herrero, *Sol. Energy Mater. Sol. Cells*, 2008, **92**, 938–941.
- J. A. Jeong and H. K. Kim, *Sol. Energy Mater. Sol. Cells*, 2009, **93**, 1801–1809.
- W. Y. Yang, Y. Gong, C. Y. Yao, M. Shrestha, Y. Y. Jia, Z. Qiu, Q. H. Fan, A. Weber and W. Li, *Lab Chip*, 2021, **21**, 1096–1108.
- T. H. Kim, S. H. Park, D. H. Kim, Y. C. Nah and H. K. Kim, *Sol. Energy Mater. Sol. Cells*, 2017, **160**, 203–210.
- S. Park, J. Yoon, S. Kim and P. Song, *RSC Adv.*, 2021, **11**, 3439–3444.
- I. Baia, M. Quintela, L. Mendes, P. Nunes and R. Martins, *Thin Solid Films*, 1999, **337**, 171–175.
- B. Tyagi, H. B. Lee, N. Kumar, W. Y. Jin, K. J. Ko, M. M. Ovhal, R. Sahani, H. J. Chung, J. Seo and J. W. Kang, *Nano Energy*, 2022, **95**, 106978.
- J. H. Song, J. W. Jeon, Y. H. Kim, J. H. Oh and T. Y. Seong, *Superlattices Microstruct.*, 2013, **62**, 119–127.
- K. P. Sibin, N. Selvakumar, A. Kumar, A. Dey, N. Sridhara, H. D. Shashikala, A. K. Sharma and H. C. Barshilia, *Sol. Energy*, 2017, **141**, 118–126.
- E. H. Ko, H. J. Kim, S. J. Lee, J. H. Lee and H. K. Kim, *RSC Adv.*, 2016, **6**, 46634–46642.
- H. Ferhati, F. Djeflal and A. Benhaya, *Superlattices Microstruct.*, 2019, **129**, 176–184.
- M. Y. Wang, A. Barnabe, Y. Thimont, J. M. Wang, Y. C. He, Q. R. Liu, X. L. Zhong, G. B. Dong, J. M. Yang and X. G. Diao, *Electrochim. Acta*, 2019, **301**, 200–208.
- J. H. Kim, A. Nizami, Y. Hwangbo, B. Jang, H. J. Lee, C. S. Woo, S. Hyun and T. S. Kim, *Nat. Commun.*, 2013, **4**, 2520.
- S. J. Oh, J. H. Kwon, S. Lee, K. C. Choi and T. S. Kim, *ACS Appl. Mater. Interfaces*, 2021, **13**, 16650–16659.
- C. H. Lee, J. H. Kim, C. Y. Zou, I. S. Cho, J. M. Weisse, W. Nemeth, Q. Wang, A. C. T. Van Duin, T. S. Kim and X. L. Zheng, *Sci. Rep.*, 2013, **3**, 2917.
- S. J. Oh, B. S. Ma, C. H. Yang and T. S. Kim, *ACS Appl. Electron. Mater.*, 2022, **4**, 3980–3987.
- J. H. Kim, K. L. Jang, K. Ahn, T. Yoon, T. I. Lee and T. S. Kim, *Sci. Rep.*, 2019, **9**, 7071.
- T. I. Lee, W. Jo, W. Kim, J. H. Kim, K. W. Paik and T. S. Kim, *ACS Appl. Mater. Interfaces*, 2019, **11**, 13416–13422.
- J. Yoon, Y. Jang, K. Kim, J. Kim, S. Son and Z. Lee, *Carbon*, 2022, **196**, 236–242.
- H. Mizubayashi, J. Matsuno and H. Tanimoto, *Scr. Mater.*, 1999, **41**, 443–448.

- 38 R. Long and M. L. Dunn, *J. Appl. Phys.*, 2014, **115**, 233514.
- 39 M. Gu, W. J. Song, J. Hong, S. Y. Kim, T. J. Shin, N. A. Kotov, S. Park and B. S. Kim, *Sci. Adv.*, 2019, **5**, eaaw1879.
- 40 J. Yun, *Adv. Funct. Mater.*, 2017, **27**, 1606641.
- 41 M. Volmer, *Z. Phys. Chem.*, 1926, **119**, 277–301.
- 42 L. Vj, N. P. Kobayashi, M. S. Islam, W. Wu, P. Chaturvedi, N. X. Fang, S. Y. Wang and R. S. Williams, *Nano Lett.*, 2009, **9**, 178–182.
- 43 X. Yang, P. Gao, Z. Yang, J. Zhu, F. Huang and J. Ye, *Sci. Rep.*, 2017, **7**, 44576.
- 44 T. Lee, D. Kim, M. E. Suk, G. Bang, J. Choi, J. S. Bae, J. H. Yoon, W. J. Moon and D. Choi, *Adv. Funct. Mater.*, 2021, **31**, 2104372.
- 45 J. S. Kim, R. H. Friend and F. Cacialli, *J. Appl. Phys.*, 1999, **86**, 2774–2778.
- 46 H. S. Im, S. K. Kim, T. J. Lee and T. Y. Seong, *Vacuum*, 2019, **169**, 108871.
- 47 W. J. Dong, H. K. Yu and J. L. Lee, *Sci. Rep.*, 2020, **10**, 12819.
- 48 P. Buffat and J. P. Borel, *Phys. Rev. A: At., Mol., Opt. Phys.*, 1976, **13**, 2287–2298.
- 49 S. R. Foltyn, L. Civale, J. L. Macmanus-Driscoll, Q. X. Jia, B. Maiorov, H. Wang and M. Maley, *Nat. Mater.*, 2007, **6**, 631–642.
- 50 S. Tsukimoto, M. Moriyama and M. Murakami, *Thin Solid Films*, 2004, **460**, 222–226.
- 51 M. L. Chang, L. C. Wang, H. C. Lin, M. J. Chen and K. M. Lin, *Appl. Surf. Sci.*, 2015, **359**, 533–542.
- 52 A. Lazauskas, V. Grigaliunas, A. Guobiene, M. Andrulevicius and J. Baltrusaitis, *Thin Solid Films*, 2012, **520**, 6328–6333.
- 53 P. J. Martin, *Gold Bull.*, 1986, **19**, 102–116.
- 54 Y. L. Cha, J. H. Jo, D. J. Kim and S. H. Kim, *Nanomaterials*, 2022, **12**, 2060.
- 55 C. Guillen and J. Herrero, *Opt. Commun.*, 2009, **282**, 574–578.
- 56 Y. S. Park, K. H. Choi and H. K. Kim, *J. Phys. D: Appl. Phys.*, 2009, **42**, 235109.
- 57 J. H. Kim, T. W. Kang, S. N. Kwon, S. I. Na, Y. Z. Yoo, H. S. Im and T. Y. Seong, *J. Electron. Mater.*, 2017, **46**, 306–311.