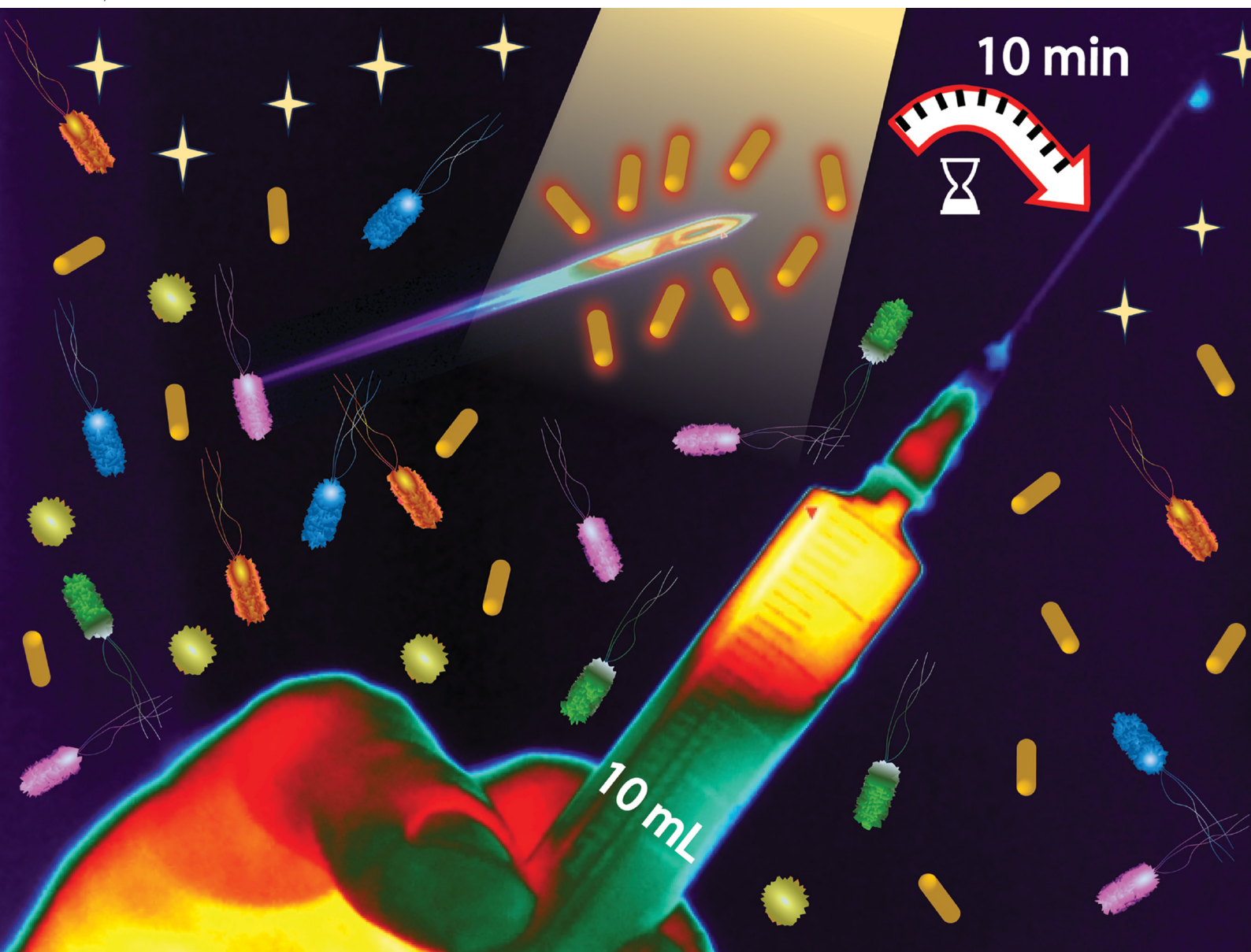


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White light thermoplasmonic activated gold nanorod arrays enable the photo-thermal disinfection of medical tools from bacterial contamination†

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The outspread of bacterial pathogens causing severe infections and spreading rapidly, especially among hospitalized patients, is worrying and represents a global public health issue. Current disinfection techniques are becoming insufficient to counteract the spread of these pathogens because they carry multiple antibiotic-resistance genes. For this reason, a constant need exists for new technological solutions that rely on physical methods rather than chemicals. Nanotechnology support provides novel and unexplored opportunities to boost groundbreaking, next-gen solutions. With the help of plasmonic-assisted nanomaterials, we present and discuss our findings in innovative bacterial disinfection techniques. Gold nanorods (AuNRs) immobilized on rigid substrates are utilized as efficient white light-to-heat transducers (thermoplasmonic effect) for photo-thermal (PT) disinfection. The resulting AuNRs array shows a high sensitivity change in refractive index and an extraordinary capability in converting white light to heat, producing a temperature change greater than 50 °C in a few minute interval illumination time. Results were validated using a theoretical approach based on a diffusive heat transfer model. Experiments performed with a strain of *Escherichia coli* as a model microorganism confirm the excellent capability of the AuNRs array to reduce the bacteria viability upon white light illumination. Conversely, the *E. coli* cells remain viable without white light illumination, which also confirms the lack of intrinsic toxicity of the AuNRs array. The PT transduction capability of the AuNRs array is utilized to produce white light heating of medical tools used during surgical treatments, generating a temperature increase that can be controlled and is suitable for disinfection. Our findings are pioneering a new opportunity for healthcare facilities since the reported methodology allows non-hazardous disinfection of medical devices by simply employing a conventional white light lamp.

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Introduction

Disinfection of harmful microorganisms is crucial in safeguarding household, public, and healthcare premises¹ and, until recent years, has been an effective tool for preventing and controlling infectious diseases. The Covid-19 pandemic highlighted the issues related to contact transmission and raised awareness of the importance of environmental hygiene for human safety. The 2022 Healthcare Cleaning Forum has underlined the crucial role of innovations like digitalization, tracking, automated disinfection, and antimicrobial surfaces, focusing on the sustainability of the newly adopted innovative solutions.¹ Routine cleaning and disinfection of healthcare structures and instruments could lower the viruses and bacteria

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transmission among patients and through objects, tools, or surfaces. Consequently, implementing antimicrobial mechanisms for frequently reused surfaces is of increasing interest. Antimicrobial surfaces can be realized, for instance, by adding coatings of different materials and producing biocidal agents with other techniques: the most common technologies are anti-adhesive surfaces, contact-active surfaces, or light-activated molecules.²

Disinfection methods can be classified as chemical and physical. The first category employs chemical agents, called biocides or germicides, that kill microorganisms interfering with their essential enzymes. Alternatively, physical disinfection methods involve using heat or radiation to eliminate microorganisms contaminating the surfaces. At sufficiently high temperatures, heat can lead to sterilization, thus removing all forms of life, including microorganisms. Chemical disinfection techniques may encounter practical difficulties since the antimicrobial compounds can cause skin and mucosal surfaces irritation and/or have an unpleasant smell; furthermore, they can be potentially flammable or cause deterioration or corrosion of metals. The usage of safer products, such as quaternary ammonium compounds (also known as quats or QACs), is not adequate for all kinds of bacteria and viruses, and the development of microbial persistence mechanisms has been reported to occur following exposure to sublethal levels of these compounds.^{3,4} These methods typically cause cell surface alteration, and changes in cell membrane permeability, causing damage to the intracellular constituents. Despite the high efficiency, reactions triggered by chemical disinfectants may produce toxic disinfection by-products, which can develop resistance traits such as generalised resistance (persistence or ability to form biofilms) or specific resistance (specific antibiotic resistance), as well as increasing human cancer risk.⁵ Commonly used physical-chemical disinfection methods include (but are not limited to) the use of ozone, chlorine dioxide, free chlorine, UV irradiation,⁶ Fenton and photo-Fenton reactions,⁷ semiconductor-assisted photocatalysis⁸ and hydrogen peroxide plus UV processes.⁹ These techniques come under the umbrella term of advanced oxidation processes (AOP) that make use of a physical trigger such as electromagnetic radiation to generate reactive oxygen species (ROS) (e.g., hydroxyl radicals $\cdot\text{OH}$) and superoxide radical anions ($\cdot\text{O}_2^-$). ROS can provoke the degradation of organic molecules by giving rise to a sequence of oxidation reactions that can potentially cause the mineralization of the target organic compound, such as lipidic components of the cell membranes.⁸ Although highly effective toward bacteria and virus inactivation, AOP processes have drawbacks, including the need for costly energy sources and chemical compounds to produce ROS. Therefore, the scientific community is very active in developing and augmenting new strategies and materials that can trigger AOP disinfection processes by solar or visible light irradiation.^{8,10,11}

Since the XIX century, Luis Pasteur introduced a paradigm shift by demonstrating the effectiveness of heating to disinfect and sterilize surgical instruments, dressing, objects, and liquids.¹² Today, the progress in nanomaterials synthesis and characterization has paved the way for using light as a cost-effective and

sustainable energy source to generate highly localized heat through PT agents. PT agents are light absorbers that convert light energy, of a suitable frequency, to heat. They often cause controlled and localized hyperthermia-producing protein denaturation to promote microorganism PT disinfection.¹³ PT disinfection is a physical disinfection method that, unlike photocatalysis, does not involve the trigger of (photo)chemical reactions but is only based on the heat produced by the PT agents.

Plasmonic nanomaterials can offer outstanding opportunities for achieving PT disinfection if used as PT agents. In particular, metal nanoparticles (NPs) exhibit excellent PT properties, providing a high light-to-heat conversion efficiency.¹⁴ Such a property arises from the Localized Surface Plasmon Resonance (LSPR) phenomenon that occurs when an electromagnetic radiation impinges on a metal NP with a particular frequency that is in resonance with one of the collective electronic oscillations of the nanomaterial. The light absorption is enhanced and confined close to the NP surface. Conduction electron oscillations increase the frequency of collisions with the lattice atoms resulting in Joule heating (thermoplasmonic effect). Therefore, the optical energy of metal NPs is converted into thermal energy with high efficiency and then released to the surrounding environment.¹⁵

Noble metal NPs such as gold (Au),¹³ silver (Ag),¹⁴ and copper (Cu)¹⁵ are extensively used as PT agents for disinfection because their resonance frequency lies in the visible range,¹⁶ making the process sustainable and cost-effective, unlike organic compounds that require UV light to generate heat.¹⁰ Among noble metal NPs, AuNPs benefit from several properties, including high PT efficiency, excellent biocompatibility, and chemical stability. Moreover, available synthesis methods currently allow tuning AuNP morphology and surface chemistry according to the desired chemical-physical properties and applications.¹⁷ Loeb *et al.* performed PT disinfection experiments using AuNPs having similar sizes but different shapes. Their work demonstrated a higher PT disinfection efficiency for Au nanorods (AuNRs) compared to Au nanocubes, highlighting AuNRs' excellent performance as nano-heaters and for broadband light sources.¹⁸ AuNRs are elongated NPs that, due to their peculiar morphology, are excellent candidates in several application fields, including biosensing,^{19,20} drug delivery,²¹ photocatalysis,²² tumour ablation by PT therapy,^{23,24} and PT disinfection.²⁵

AuNRs were also employed as building blocks in nanocomposites to realize devices for PT disinfection. AuNRs embedded in polydimethylsiloxane, an optically transparent and non-toxic silicone polymer, were used for fabricating microfluidic channels with a high PT response and water disinfection capabilities.²⁶ Loeb *et al.* coated the glass coverslip of a solar reactor with PT nanomaterials (AuNRs and carbon black) to enhance water disinfection capabilities. Their results highlighted the significant contribution of AuNRs that display their PT properties in the Near Infrared (NIR) range, leading to a temperature increase 78% higher than the temperature increase achieved with only carbon black under solar irradiation.²⁷



However, the reactor structure is designed for bulk water disinfection. Thus, the contributions of carbon black and high-density AuNRs to the generation of heat are more relevant despite the absence of an optical signal caused by the carbon black double layer opacity and to the spectrum broadening due to the AuNR aggregates.

In this work, we intend to take a step forward in the field of PT disinfection by tackling the challenge of achieving bacteria killing under white light irradiation. To this end, we exploit the extraordinary capability of an array of AuNRs firmly deposited on a glass substrate using a controlled electrostatic Layer-by-Layer (eLbL) assembly method. The resulting highly dense and well-monodispersed AuNRs array possesses a large absorption cross-section, making them promising candidates for broadband energy conversion through a PT process. Our experimental results, supported by a solid theoretical analysis, demonstrate that, under the investigated experimental conditions, white light irradiation of the AuNRs array produced a temperature increase suitable for achieving white light-triggered disinfection. We show that our highly efficient plasmonic platforms can be effectively utilized in healthcare facilities, allowing the white light-assisted thermal disinfection of different surgical instruments.

Experimental

Materials

Citrate-capped AuNRs having size $55 \text{ nm} \times 15 \text{ nm}$ (AuNRs) were purchased from Nanocomposix. Acetone, isopropanol, methanol, sodium hydroxide (NaOH), poly(sodium 4-styrenesulfonate)(PSS, $M_w \sim 70 \text{ kDa}$), and poly(allylamine hydrochloride)(PAH, $M_w \sim 50 \text{ kDa}$) were purchased from Merck. Deionized water was used to prepare the colloidal dispersions of AuNRs and the polyelectrolyte solutions.

Escherichia coli (*E. coli*) K12 MG1655 CGSC#7740 was obtained from the Coli Genetic Stock Centre (CGSC) collection. The minimal medium E supplemented with 0.40% glucose was used for bacterial growth. The chemicals required for preparing the bacterial growth medium were purchased from Merck or VWR International. LIVE/DEAD™ BacLight™ Bacterial Viability Kit for microscopy was purchased from Thermo Fisher Scientific.

Immobilization of AuNRs on glass substrates

The incorporation of AuNRs on glass substrates was achieved following the procedure reported in ref. 28. Briefly, $1 \text{ cm} \times 1 \text{ cm}$ sized glass substrates were thoroughly washed in an ultrasonic bath for 10 min by using, sequentially, methanol and acetone. An intermediate rinsing step in isopropanol was carried out between the two washing steps. Finally, the glass substrates were stored in isopropanol, rinsed, and dried under a stream of nitrogen before use. Before fabricating the polyelectrolyte (PE) multilayer, the substrates were immersed for 30 min in a 5 M NaOH solution to impart a negative charge triggering the electrostatic incorporation of the first positively charged PE.

PEs were deposited using eLbL by following the sequence: PAH-PSS-PAH. Accordingly, the building of the PE multilayer was performed by sequentially immersing the glass substrate in PAH, PSS, and PAH solutions for 10 min. In particular, the concentration of the PE solution was 1.6 mg mL^{-1} , and the pH was 2 for the PAH solution and 8 for the PSS, respectively. In-between two consecutive immersion steps, an intermediate washing step was performed (2 min immersion in water) to remove the excess of PE molecules. Finally, a final step of a 2 min immersion in water was carried out to avoid the counterion effect.²⁹ The substrates were then dried under a stream of nitrogen and stored in a refrigerator ($+4 \text{ }^\circ\text{C}$). Subsequently, the PE functionalized glass substrates were immersed for 16 h in a colloidal dispersion of AuNRs (suitably diluted to obtain an optical density of 1 at 790 nm). After this step, the AuNRs-modified substrates were washed with water and gently dried under a stream of nitrogen before their characterization.

Sensitivity to the refractive index variation and photo-thermal characterization

A double-face glass cell was used to investigate the optical response of the AuNR-modified substrates to the alteration of the refractive index (n) of the infiltrating medium. This configuration was suitably selected to study the optical and the PT behavior of the AuNR-modified substrates in two different conditions: (i) when only one side of the substrate underwent a n variation, and (ii) when both sides of the AuNRs substrate experienced an n variation.

To this end, a cell was fabricated, such that the AuNR substrate was sandwiched between two glass slides ($1.2 \text{ cm} \times 1.2 \text{ cm}$). First, to ensure a uniform $10 \text{ }\mu\text{m}$ gap between the plasmonic substrate and the two glass slides, the glue NOA-61 with $10 \text{ }\mu\text{m}$ glass microbeads was deposited on the corners of the AuNRs substrate on both sides. Then the AuNRs substrate was placed in between the two glass slides. After that, the cell was sealed by exposing it to UV light radiation for 1 min. At this stage, the first side of the resulting cell was infiltrated with the NOA-61 as a representative medium with a known n . Subsequently, the absorption spectrum was collected, and the PT measurements were performed. The same procedure was performed after infiltrating the second side of the double-face cell.

White light disinfection

The ability of the AuNR modified substrates to achieve surface disinfection, induced by white light irradiation, was investigated utilizing *E. coli* cells as a model bacterium.

AuNR substrates were immersed in $500 \text{ }\mu\text{L}$ of a 10^4 CFU mL^{-1} *E. coli* cells dispersion in minimal *E* (with no glucose). After 30 min, the substrate, contaminated with *E. coli* cells, was dried under a stream of nitrogen. The substrate was then characterized by absorption spectroscopy and irradiated with white light (14.7 W cm^{-2} for 10 min) by using the experimental setup described in section 2.5.4. After irradiation, the substrates were inspected by fluorescence microscopy using propidium iodide as a staining agent. The staining was realized by immersing the sample



for 10 min in 500 μL minimal E with 2 μL fluorescent dye propidium iodide. The substrate was dried under a stream of nitrogen before microscopy analysis. Control experiments used SYTO 9TM as a staining agent for identifying viable cells.

Characterizations

UV-Visible absorption spectrophotometry. A Lambda 365 spectrophotometer from PerkinElmer was used to collect the absorption spectra of the AuNR modified substrates and the double-face AuNRs cell. A diode array spectrophotometer HP8453 (Agilent Technologies) was used to measure the OD600 of *E. coli* cell cultures. An Ocean Optics USB spectrophotometer was used for measuring the white light source spectrum.

Scanning electron microscopy. A scanning electron microscopy (SEM) was used to study the morphology of AuNR substrates. The measurements were obtained by a field emission scanning electron microscopy (FE-SEM, Carl Zeiss, SUPRA 40VP) with an accelerating voltage of 2 kV.

Atomic force microscopy. The topography of the AuNR substrates was analyzed by atomic force microscopy (AFM, Nanoscope Multimode system, Veeco Instruments). The measurements were performed in tapping mode with a vertical resolution of 0.1 \AA and lateral resolution of 2 \AA .

White light photo-thermal measurements. A broadband light source with an operating wavelength range from 400 to 1000 nm was used for the PT characterization of AuNR substrates. The power density was tunable and an optical fiber was used to irradiate the whole sample area uniformly. A high-resolution thermal camera (FLIR, A655sc) was used to record the temperature variation during the irradiation and map the heating distribution. The thermal images are 640×480 pixels and have an accuracy of ± 0.20 $^{\circ}\text{C}$. The software FLIR ResearchIR Max was used to acquire and process data from the thermal camera.

Contrast phase and fluorescence microscopy. A ZEISS Axiolab 5 fluorescent microscope, equipped with contrast phase objectives and fluorescence modules, was used to collect micrographs of the bioactive AuNR substrates.

Results and discussion

AuNRs substrate characterization and white light photo-thermal response

AuNPs exhibit a strong LSPR effect and are tunable according to size, shape, and surface chemistry. Au is a noble metal with good biocompatibility that enables interaction with biological entities such as human cells and bacteria without affecting their viability.³⁰ AuNRs are an optimal model for studying white light thermoplasmonic properties from a theoretical and experimental perspective. AuNRs exhibit an absorption cross-section higher than Au Nanospheres (AuNSs).³¹ Indeed, in the absorption spectrum of AuNRs, the transverse plasmon band can be assimilated to the absorption peak of AuNSs having a diameter equal to the thickness of AuNRs. Such absorption signal is much less intense than the longitudinal plasmon

band, thus accounting for the lower absorption cross-section of AuNSs consequently, the less efficient thermoplasmonic conversion efficiency. Therefore, among AuNPs, AuNRs possess an excellent light-to-heat conversion ability, having a PT efficiency of 100% when irradiated at their resonance frequency (monochromatic light source), usually in the NIR range.³² The NIR light is a significant component of the white light source and solar spectrum; thus, exploiting this spectral range can open an opportunity to realize solar-light thermoplasmonic disinfection. Such an outstanding property inspired the goal of the present work. Here, we aim to push the borders of PT conversion efficiency to visible light by investigating AuNR-based platforms able to generate thermal energy upon white light irradiation. In addition, we also exploit the white light PT heating to achieve PT disinfection by using *E. coli* as a model pathogen to simulate a form of surface contamination.

The immersive eLbL method was used for the fabrication of the AuNR-based plasmonic platforms. The procedure relies on the electrostatic interactions involving a negatively charged glass substrate, PEs, and AuNRs. In particular, the AuNR substrates were obtained by performing 3 steps: (i) activation of the glass substrate, (ii) construction of a PE multilayer by eLbL, and (iii) incorporation of AuNRs. The glass surface activation was achieved by immersing it in NaOH. The PE multilayer was realized by alternating a weak positively charged PE (PAH), and a strong negatively charged PE (PSS), giving rise to the sequence PAH-PSS-PAH. A multilayer PE architecture, instead of a monolayer, is suitable for obtaining a high density of AuNRs. Indeed, it was already reported that the amount of deposited NPs increases with the number of PE layers underlying the NPs array.^{33,34} In addition, the number of PE layers increases the roughness of the substrate,³⁵ providing a high number of charged sites available for anchoring AuNRs. Finally, beyond the electrostatic attraction, the PEM multilayer assembly is also driven by an overcompensation effect, namely an excess of surface charges not paired with those in the underlying layer. Therefore, a multilayer, instead a monolayer, provides a higher surface charge number,³⁶ resulting in a higher AuNRs density.

After that, 16 hours of immersion in an AuNRs colloidal dispersion promoted the incorporation of the plasmonic NPs on the glass surface. In this step, a colloidal AuNRs dispersion with O.D. 1 was used.

Our preliminary experiments demonstrated that the O.D. 1 is suitable for obtaining AuNR substrates with optimal optical properties under the investigated experimental conditions. This value was carefully selected after performing several preliminary experiments consisting of immersion of the PEM-modified glass substrates in AuNRs at different O.Ds. Using a colloidal dispersion with an O.D. less than 1 resulted in substrates with a low amount of AuNRs. Conversely, by exceeding the O.D. 1, no significative increase in the AuNRs density on the substrates was detected. These results suggested that under the investigated experimental conditions, the underlying PAH layer can host only a limited number of AuNRs.

Fig. 1a reports the optical characterization performed by absorption spectroscopy of the resulting AuNRs substrate.



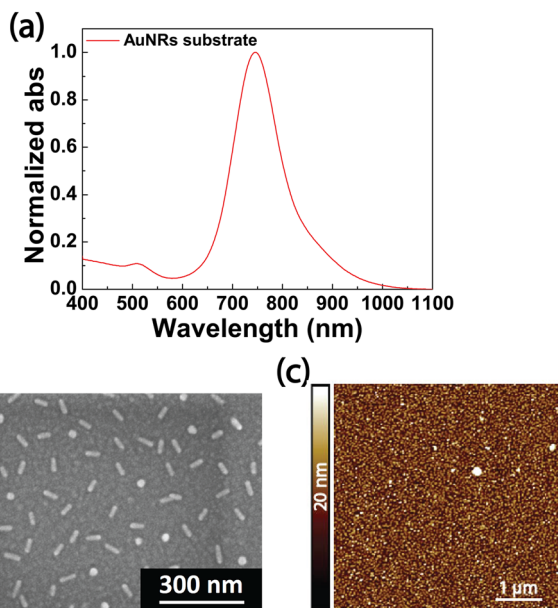


Fig. 1 Characterization of the AuNRs functionalized glass substrate fabricated through the immersive eLbL assembly method. Optical characterization of AuNR substrate performed by absorption spectrophotometry (a). Morphological characterization of the plasmonic substrate performed by SEM (b). Topographic image obtained by AFM (c).

It reveals the presence of two peaks corresponding to the transverse (LSP_T) and longitudinal (LSP_L) plasmon modes due to the LSPR phenomenon. The LSP_T peak is centered at 515 nm and much less intense than the LSP_L peak, located at lower energy, namely at 764 nm. The absorption spectrum (Fig. 1a) suggests that the optical properties of the AuNRs array on a glass substrate resemble the ones of AuNR colloidal dispersions where AuNRs are monodisperse due to their peculiar physiochemical features. Remarkably, the spectral response of the sample was acquired by using a bare glass substrate as a baseline. Considering, instead, the glass functionalized with PEM as the baseline, no variation is introduced in the absorption spectrum of the AuNRs substrate. Indeed, the PE multilayer shows an absorption peak in a wavelength range lower

than 300 nm. In this spectral range, the glass substrate exhibits a very intense absorption (see Fig. S1, ESI[†]), thus overlapping the contribution of the PE multilayer. The morphology of the AuNRs substrate (Fig. 1b), analyzed by SEM, pointed out an even AuNRs distribution, where AuNRs are well separated from each other, with no formation of aggregates. Such a result is ascribable to the PE multilayer that confers a uniform charge distribution on the glass substrate. The PE solutions were prepared in the absence of salts. Accordingly, we expect the PAH layer to display a stretched structure that fosters the incorporation of AuNRs occurring through electrostatic attractions involving the positively charged amine groups of the PAH and the negatively charged citrate molecules on the AuNRs surface.

Moreover, from the SEM micrographs analysis, it is computed that AuNR substrates have a fill fraction of $5.9\% \pm 0.3\%$ and an interparticle distance of (120 ± 22) nm. The topographic investigation performed with AFM is consistent with this hypothesis. Indeed, as displayed in Fig. 1c, the 2D AFM image of the AuNRs substrate shows a uniform distribution of asperities, having an average height profile of ~ 10 nm corresponding to the short axis of the NRs, in agreement with the generation of AuNRs monolayer.

Remarkably, the absorption spectrum of the AuNRs substrate shows the most intense LSP_L band between 600 nm and 900 nm. This absorption signal overlaps the white light source spectrum, as shown in Fig. 2a, which displays an intense signal in the wavelength range from 400 nm to 900 nm. Accordingly, the PT performance of AuNR substrates under white light was investigated using the optical setup described in Fig. 2b; it uses a white light beam that uniformly illuminates the entire sample area at normal incidence (Fig. 2c). During irradiation, a high-resolution thermal camera recorded the temperature profile and the spatial heating distribution on the sample surface in a defined region of interest (ROI). The short distance (1 cm) between the AuNRs substrate and the optical fiber maintains the light beam's convergence.

The experimental results from the PT investigation of the AuNR substrates under white light irradiation are reported in Fig. 3. In this set of experiments, the AuNR substrates were

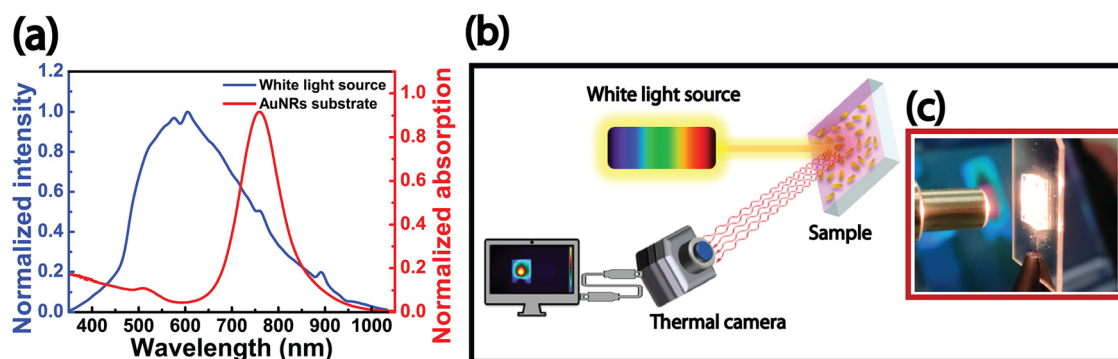


Fig. 2 Emission spectrum of the white light source intensity (blue curve), overlapped with the absorption spectrum of a representative AuNRs substrate (red curve) (a). Sketch of the PT optical setup used to perform the white light thermoplasmonic characterization of the AuNR substrates and PT disinfection experiments (b). Representative image of the white light source impinging the sample area (c).



illuminated for 5 min. The subsequent shutdown of the white light source allowed the sample to cool down for 2 min. The PT response of AuNR substrates was investigated by varying the power density of the white light source in a range from 0.949 W cm^{-2} to 28.6 W cm^{-2} . The analysis of the resulting thermographic images gave rise to the plot in Fig. 3a, reporting the maximum temperature increase (ΔT_{max}) as a function of irradiation time. Independently from the power density, Fig. 3a shows a progressive (exponential) temperature increase during the first 2 min of the 5 min irradiation, followed by a gradual decrease when the light source was turned off. Such a result is consistent with experiments carried out under laser light irradiation for colloidal dispersion of AuNRs³⁷ and for AuNRs immobilized on glass substrates²⁸ or embedded in polymeric matrices.³⁸

Fig. 3b reports the ΔT_{max} values as a function of the light source power density. Experimental results highlight that 5 min of white light irradiation can produce an appreciable temperature increase when the power density is higher than 5 W cm^{-2} . Indeed, when the minimum light source power density was set at 0.949 W cm^{-2} , a ΔT_{max} of $4.44 \text{ }^\circ\text{C}$ was measured, while the maximum ΔT_{max} value rose to $49.3 \text{ }^\circ\text{C}$ when the power density increased to 28.6 W cm^{-2} . Selecting a power density of 14.7 W cm^{-2} , the ΔT_{max} achieved was $44.7 \text{ }^\circ\text{C}$. This temperature increase value corresponds to a T_{max} of $69 \text{ }^\circ\text{C}$, suitable for performing the white light disinfection process in a proper time interval of 10 min. Such a temperature value is above the one identified in the paper by Annesi *et al.* ($65 \text{ }^\circ\text{C}$) for promoting a reduction of $2 \log \text{ CFU}$ of *E. coli* population (more than 90% in viability reduction) after 7.5 min of AuNRs NIR laser illumination.³⁹ Accordingly, in our experimental conditions,

we selected a white light irradiation time of 10 min as a suitable timeframe to induce PT disinfection. For this reason, the power density of 14.7 W cm^{-2} was chosen as a reference value for further PT experiments of characterization and disinfection.

The fitting of the experimental ΔT_{max} values as a function of the power density, reported in Fig. 3b, indicated that the power density increase of the white light source produces an exponential rise of ΔT_{max} , according to the equation:

$$\Delta T = 49.3 - 52.8 \times \exp\left(-\frac{\text{power density}}{5.8}\right) \quad (1)$$

Moreover, cycling experiments were carried out in three consecutive irradiation cycles on the same substrate at a power density of 14.7 W cm^{-2} . Results reported in Fig. 3c proved the PT stability of the samples. Indeed, the same value of the ΔT_{max} was attained in each cycle. A representative thermographic image of the AuNRs substrate under white light irradiation, shown in Fig. 3d, provides evidence of the uniform heating distribution on the sample's surface. A control experiment was performed by irradiating a bare glass substrate with the white light source, applying a power density of 14.7 W cm^{-2} for 10 min. The resulting time vs. ΔT_{max} profile, in Fig. 3e, yields a ΔT_{max} of $5 \text{ }^\circ\text{C}$, reaching a maximum temperature (T_{max}) of $30 \text{ }^\circ\text{C}$ after 10 min of irradiation. When the same experiment was carried out by irradiating the AuNRs substrate, a T_{max} of $69 \text{ }^\circ\text{C}$ and a ΔT_{max} of $44.7 \text{ }^\circ\text{C}$ were achieved (Fig. 3f), which is in complete agreement with the values reported in Fig. 3a. Such a ΔT_{max} value is 87.3% higher than the ΔT_{max} of the bare glass slide, thus demonstrating the ability of AuNRs substrate to convert white light into thermal energy efficiently. AuNRs are excellent PT transducers because they can entirely absorb the

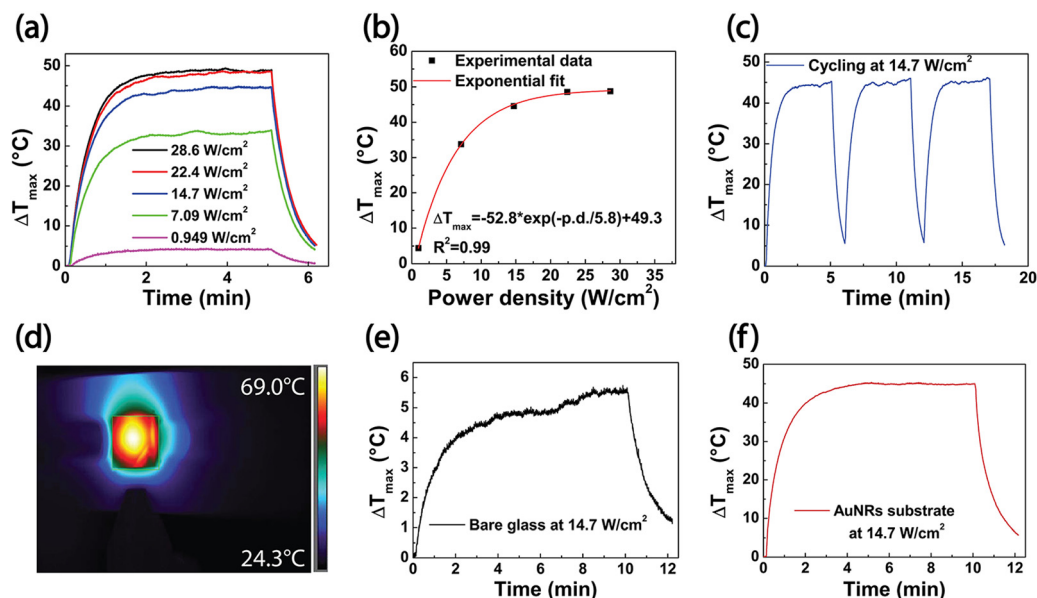


Fig. 3 PT characterization of AuNR substrates under white light irradiation. Time-temperature dependence as a function of the power density of the light source (a). Plot of the maximum temperature increase obtained for different power density values (b). Maximum temperature increase as a function of time for three consecutive cycling experiments performed at the power density of 14.7 W cm^{-2} . (c). Representative thermographic image of an AuNRs substrate acquired after 10 min of irradiation at a power density of 14.7 W cm^{-2} (d). Maximum temperature increase as a function of time achieved by irradiating a bare glass substrate (e) and an AuNRs substrate (f) with white light at a power density of 14.7 W cm^{-2} for 10 min.



impinging EM radiation due to the LSPR phenomenon. An important figure of merit to evaluate the PT properties of the AuNR substrates is the PT conversion efficiency, η_{heat} , defined as the ratio of the increased internal energy to the total incident radiation. Assuming that the PT experiments were performed at constant room temperature and with uniform white light irradiation, the following expression can be calculated, as typically reported for nanofluids:⁴⁰

$$\eta_{\text{heat}} = \frac{(c_{\text{glass}}m_{\text{glass}} + c_{\text{Au}}m_{\text{Au}})\Delta T}{IA\Delta t} \approx \frac{c_{\text{glass}}m_{\text{glass}}\Delta T}{IA\Delta t}, \quad (2)$$

where C_{glass} , C_{Au} , m_{glass} , and m_{Au} are the specific heat capacity and mass of glass and gold, ΔT is the temperature increase, Δt is the time interval, A is the illumination area, and I is the power density. Since the AuNR mass is negligible, the expression in eqn (2) can be simplified.^{41,42} In this formula, the critical role of the glass substrate is evident. For the same volume of AuNRs in the sample, the glass substrate confers a lower value for the specific heat capacity to the water value that plays the same role in colloidal solutions, thus facilitating the temperature increase.

Assuming the specific heat capacity of borosilicate glass of $779.7 \text{ J kg}^{-1} \text{ }^{\circ}\text{C}^{-1}$, the PT efficiency for bare glass and AuNRs substrate are 5.50% and 43.5%, respectively; they are calculated according to the data shown in Fig. 3e and f, for a broadband illumination of 10 min at a power density of 14.7 W cm^{-2} , and estimated at a wavelength of 600 nm (see Fig. 2a) on a glass substrate of $1 \text{ cm} \times 1 \text{ cm}$ with a mass of 0.212 g.

AuNRs substrate sensitivity to refractive index change

The AuNRs immobilization protocol discussed in section 2.2 was implemented to deliberately deposit AuNRs on both sides of the glass substrate in a one-shot process. However, a possible strategy for further improving the PT properties of the resulting thermoplasmonic substrates can involve the variation of the n of the infiltrating medium. It is well known that the photo-induced temperature increase is also related to the overall n of the chemical environment experimented with by the plasmonic NPs.⁴³ In the present section, we report and discuss the result of an experiment specifically designed to investigate the effect of the n variation on the AuNRs substrate optical response in two different conditions. In the first, a non-uniform n variation is achieved by varying the n of only one face of the AuNRs substrate. In the second, a uniform n variation is obtained by varying the n of both faces of the AuNRs substrate.

Hence, the sensitivity of the AuNR samples was studied by separating the upper and lower sides of the AuNR substrates. To this end, a double face glass cell reported in Fig. 4a was fabricated and infiltrated either with air (Case 1) or with a high n medium, on one side as reported in Fig. 4b (Case 2) and on both sides 4c (Case 3). The glue, NOA-61, was selected as a high n infiltrating medium ($n = 1.56$).

The optical and PT characterization was performed on these three different cases (Fig. 4).

The absorption spectra in Fig. 5a demonstrate that the LSPR position, initially at 766 nm for AuNRs substrate (Case 1), varied

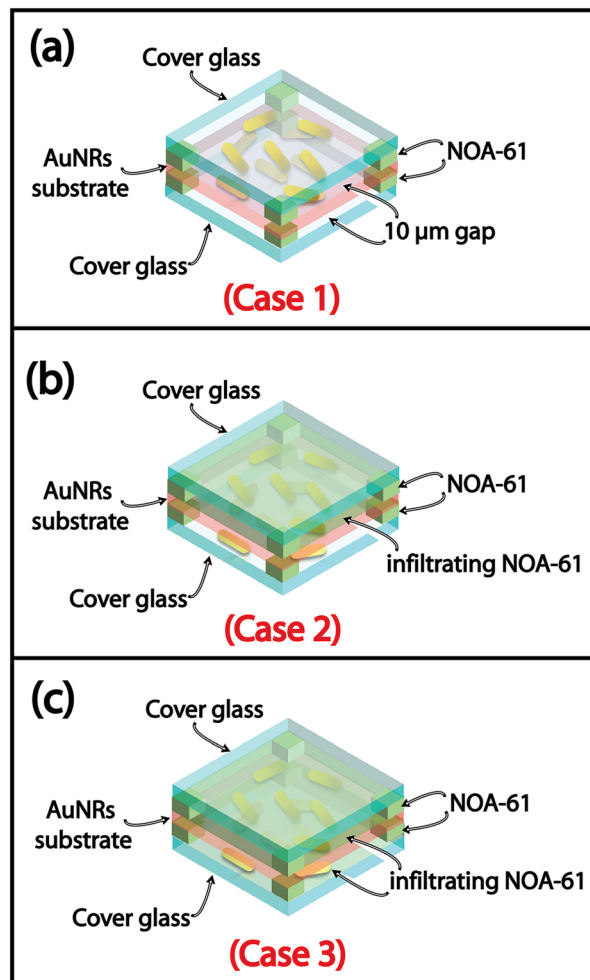


Fig. 4 Scheme of the double face glass cell used for investigating the optical and PT behavior of the substrate as a function of n . The double face glass cell (Case 1) was prepared by following the procedure reported in section 2.3 (a). The first side of the cell was infiltrated with NOA 61 (Case 2), and then optical and PT characterization was performed (b). Subsequently, the second side of the cell was infiltrated with NOA 61 (Case 3) so that both glass sides functionalized with AuNRs experienced the same n variation (c), before investigating the optical and the PT behavior.

only 14 nm in the non-uniform infiltration of NOA-61 (Case 2). Still, a shoulder is evident at 882 nm. When both the faces were infiltrated with NOA-61, the LSPR band peaked at 882 nm resulting in a total red-shift of 118 nm (Case 3). This peculiar spectroscopic behavior indicates that the cell absorption spectrum results from the contributions of the sample's two sides. The condition identified as Case 2 of the experiment is characterized by different n_s for each side, namely an uneven n variation. In other words, the AuNRs array experienced a n variation on one side and a n preservation on the other, resulting in a 14 nm low red shift of the LSPR position, along with the shoulder at 882 nm. Conversely, in Case 3, both sides of the substrate underwent the same n variation, resulting in a well-defined absorption peak at 882 nm and, therefore, in a greater LSPR shift of 118 nm. It is possible to infer that the shoulder accounts for the AuNRs, deposited on the first face that



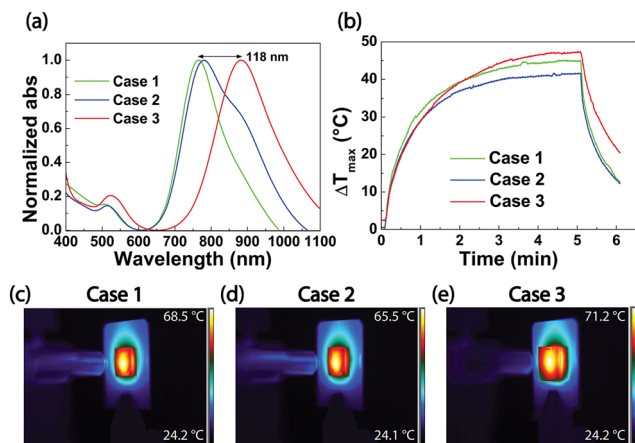


Fig. 5 Optical (a) and white light PT (b) characterization of AuNRs substrate investigated by varying, sequentially, the n of one side and two sides of the AuNRs substrate. Thermographic images were taken after 5 min of irradiation of the empty double-face cell (Case 1, c), the one-side infiltrated cell (Case 2, d), and the two sides infiltrated cell (Case 3, e).

experiences the n variation caused by the NOA-61. Meanwhile, the uniform distribution of NOA-61 on both faces of the AuNRs substrate gave rise to a symmetric LSPRI signal.

Moreover, white light PT experiments were performed by irradiating the sample in Cases 1, 2, and 3. Time-temperature plots are reported in the green (Case 1), blue (Case 2), and red (Case 3) traces of Fig. 5b. The thermal camera recorded 5 min of illumination at a power density of 14.7 W cm^{-2} plus 1 min with the beam turned off. The temperature increase is greater for Case 3 than for Case 1; therefore, a temperature increase was observed for a higher n of the surrounding medium. The ΔT_{max} measured for Case 3 is elevated by about $2.3 \text{ }^\circ\text{C}$ with respect to Case 1, as expected from theoretical investigations.⁴⁴ Meanwhile, ΔT_{max} is lower ($-3.4 \text{ }^\circ\text{C}$) in Case 2 with respect to Case 1.

The increase in temperature is dependent on the light source for its intensity, I , and on the AuNRs for the absorption cross-section of the sample σ_{abs} , the thermal conductivity k_{m} , the shape-correction factor β , and the equivalent radius of a sphere having the same volume of the nanorod R_{eq} ,⁴⁵ as described the eqn (3):

$$\Delta T_{\text{np}}^{\text{CW}} = \frac{I}{4\pi k_{\text{m}}} \frac{\sigma_{\text{abs}}}{R_{\text{eq}} \beta} \quad (3)$$

Moreover, the absorption cross-section of the AuNRs substrate depends on the n of the surrounding medium through its dielectric constant:

$$\sigma_{\text{ext}} \approx \sigma_{\text{abs}} = \frac{\omega}{3c} \varepsilon_{\text{m}}^{3/2} V \sum_j \frac{\left(\frac{1}{P_j^2}\right) \varepsilon_2}{\left[\varepsilon_1 + \left[\frac{(1-P_j)}{P_j}\right] \varepsilon_{\text{m}}\right]^2 + \varepsilon_2^2}, \quad (4)$$

where P_j is the depolarization factor dependent on the shape of NRs.¹⁹ As a result, the increased temperature becomes higher for larger dielectric constant values and, thus, for larger n_s of

the medium surrounding the AuNRs. However, the increased n entails a significant red-shift of the LSPRI position. The initial LSPRI wavelength at 766 nm shifts by more than 100 nm (882 nm). The corresponding intensity of the white light source spectrum at 882 nm , shown in Fig. 2c, is lower, reducing the PT temperature rise. Therefore, the experimental results shown in Fig. 5b suggest the occurrence of a competitive effect. The expected ΔT_{max} increase due to the rise of the n is counterbalanced by the decrease of the light source intensity at the resonance wavelength, which leads to a ΔT_{max} reduction. The thermal camera images in Fig. 5c–e show that the heating distribution is uniform on the sample surface, despite the increasing thickness resulting from embedding the AuNRs substrate in the glass cells.

Theoretical modelling of the optical and photo-thermal properties

The theoretical simulations were performed following two approaches to confront the experimental results. The first concerns the analytical solutions of the optical cross sections specified by Rayleigh–Drude approximation,^{46,47} whereas the second part considers the temperature distributions obtained *via* CFD simulations. The substrate effect is included following Yamaguchi's⁴⁸ and Rocher's⁴⁹ approaches, assuming that symmetrically arranged nanorods are $\sim 98.5 \text{ nm}$ for x - and $\sim 78.5 \text{ nm}$ for Z -axes away from one another. The complete calculations and all the theoretical details are provided in the ESI.†

Optical properties

From a theoretical point of view, absorption spectra are strongly dependent on the absorption coefficient of each material. Different formulas govern the absorption coefficient for continuous media and nanostructures that will be investigated. Hence, including the Gaussian size and shape distribution of AuNRs, the general equation for absorbance, Abs, assumes:

$$\text{Abs} = \log\left(\frac{I_0}{I_{\text{absM}} + I_{\text{absi}}}\right), \quad (5)$$

$$I_{\text{absM}} = I_0 \cdot (1 - R_{\text{M}}(\lambda)) \cdot \left(1 - \exp\left(-\left(\frac{4\pi \cdot \text{im}(n_{\text{M}}(\lambda))}{\lambda}\right) \cdot d_{\text{M}}\right)\right), \quad (6)$$

$$I_{\text{absi}} = I_0 \cdot (1 - R_{\text{coating}}) \cdot (1 - R_{\text{NR}_i}) \cdot \left(1 - \exp\left(\left(-\xi \cdot \sum_i (w_i \cdot \sigma_{\text{abs}_i}) \cdot l_{\text{p-h}}\right)\right)\right), \quad (7)$$

where I_0 is the intensity of the incident beam, I_{absM} is the intensity absorbed by the surrounding material; I_{absi} is the intensity absorbed by the AuNRs layer, $R_{\text{M}}(\lambda)$ is the wavelength function of the reflection coefficient of the surrounding medium; λ is the wavelength of the incident radiation; $n_{\text{M}}(\lambda)$ is the wavelength function of the n of the surrounding materials; d_{M} is the thickness of the surrounding materials; R_{coating} is the capping agent of AuNRs; R_{NR_i} is the reflection coefficient of AuNRs;



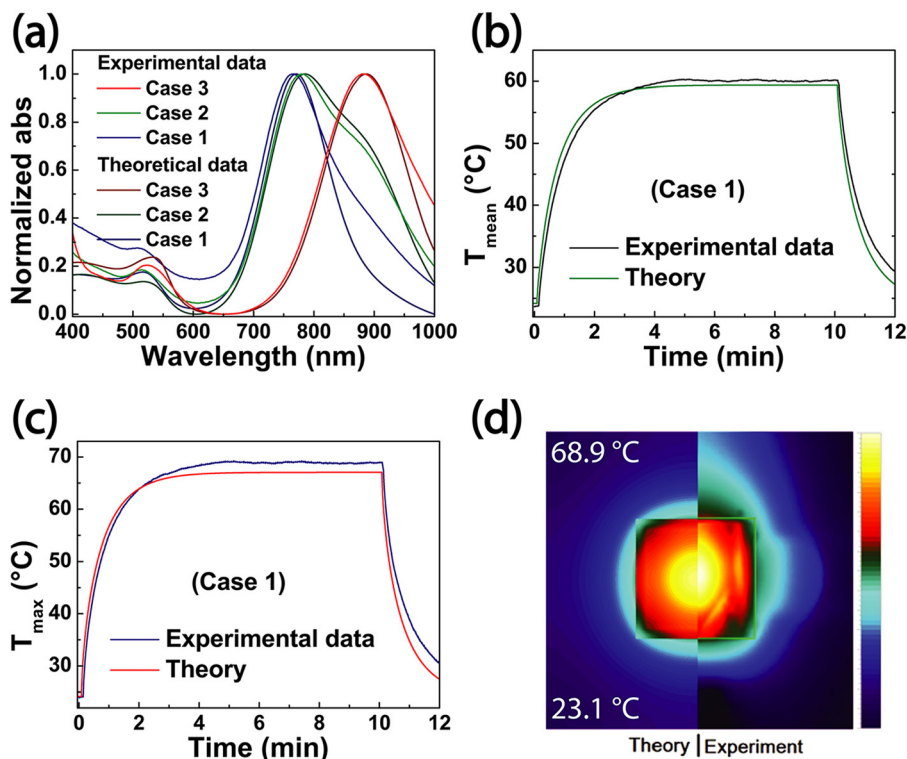


Fig. 6 Experimental (bright blue, green and red lines) and theoretical (deep colored lines) comparison of the absorption spectra for the three different cases (1, 2 and 3 previously discussed) (a). Experimental (black and blue curves) and theoretical (red and green curves) comparison of maximum (b) and average (c) temperature changes. Theoretical and experimental visual comparison of temperature distribution (Case 1) after 10 min of irradiation (d).

ξ is the AuNRs density; N is the number of AuNRs; w_i is the AuNRs polydispersity; σ_{abs_i} is the absorption cross section; $l_{\text{p-h}}$ is the AuNRs layer thickness (26.0 nm); i is assigned to AuNRs' size and shape Gaussian distributions; and im refers to the imaginary part. By construction, AuNRs were considered randomly oriented on both sides of the glass substrate, with an even density.⁴⁷ Fig. 6a reports the experimental (bright blue, green, and red lines) and the corresponding theoretical results (dark blue, dark green, and dark red colored lines, respectively) for the three different cases (1, 2, and 3) previously discussed in Fig. 5a. It is worth noting that the agreement between the experimental (Fig. 6a, bright colored lines) and theoretical (Fig. 6a, deep colored lines) results is excellent in the visible range of the electromagnetic spectrum (450–800 nm), while there is a slight deviation for shorter (below 450 nm) and longer (above 800 nm) wavelengths for all three cases.

The disagreement below 450 nm can be attributed to Rayleigh approximation that works under condition $\lambda/2\pi > d_{\text{NP}}$. For lower λ values, Mie theory is necessary to describe the multipoles interactions in each AuNRs. Above 800 nm, a substrate effect might play a crucial role. Indeed, the utilized equation⁴⁸ assumes a flat substrate. However, in the experimental cases, the different areas of the substrates do not fulfill the assumption. This inhomogeneity cannot be easily included in the theoretical calculation and could introduce the disagreement between the experimental and theoretical curves for longer wavelengths (Fig. 6a).

Photo-thermal properties

The heat transfer simulations were obtained only for Case 1, using the experimental parameters (white light irradiation, 14.7 W cm^{-2} for 10 min). This work introduces the approach where \dot{S}_e^{TOT} , \dot{S}_e^{NP} , and \dot{S}_e^{M} are treated as the converted heat from incident electromagnetic energy,⁵⁰ including Gaussian size and shape distributions, as follows:

$$\dot{S}_e^{\text{TOT}} = \dot{S}_e^{\text{NRs}} + \dot{S}_e^{\text{M}} = \frac{\int_{\lambda_0}^{\lambda_k} \left(\xi \cdot \sum_i^N (w_i \cdot \sigma_{\text{abs}_i} \cdot I_{\text{abs}_i}) \right) d\lambda}{\int_{\lambda_0}^{\lambda_k} d\lambda} + \frac{\int_{\lambda_0}^{\lambda_k} (A_{\text{absM}} \cdot I_{\text{absM}}) d\lambda}{\int_{\lambda_0}^{\lambda_k} d\lambda}, \quad (8)$$

where: \dot{S}_e^{TOT} is the source of energy for continuous materials; \dot{S}_e^{M} is the source of energy for the surrounding materials; A_{absM} is the absorption coefficient of a continuous material; \dot{S}_e^{NRs} is the source of energy for nanoparticles; λ_0 , λ_k are the lower and upper limit of the wavelength interval, here: 250 nm and 1100 nm, respectively.

Boundary conditions are based on adiabatic conditions and mixed (radiation and convection) conditions on the surfaces where the lamp impinges. Due to the lack of some parameters about the NOA61 glue, epoxy resin's properties have been assumed, which can be found in the ESL.† Transient simulations follow the SIMPLE algorithm and the second-order



computational scheme for the (Y.1)–(Y.3) equations. The calculations have been performed *via Ansys.Fluent* software (version 22.1) using the *Tryton* supercomputer possessing 72 cores (Intel[®]Xeon[®]Processor E5 v3@2.3 GHz). The control parameters are established at 0.5 except for pressure and density, whose values equal 0.3. Spatial and time independence are verified and obtained *via* Richardson and Roache extrapolations.⁵¹

Fig. 6b and c report the time-temperature profiles obtained for average (Fig. 6b) and maximum (Fig. 6c) for the theory and the experiments. Fig. 6d shows a corresponding visual comparison of the temperature distributions.

A good agreement between theory and experiments is worth observing, except for a slight mismatch. This difference can be ascribed to the fact that the experimental thermal characterization utilizes a thermal camera that also measures possible irregularities not present in the theoretical model. In addition, another potential difference can be ascribed to the difference in the boundary conditions since a perfect symmetrical assumption was considered in the theoretical part.

White light photo-thermal disinfection

Thermoplasmonic white light disinfection is expected to be a valuable alternative to conventional disinfection induced by irradiation with UV light. Moreover, white light sources are safer for human health and the environment than UV light sources. The germicidal UV lamps show emissions between 200 and 280 nm (UV-C). Although effective for disinfection, as UV-C light sources emit at 254 nm, they can induce several injuries, including sunburn, skin cancer, photokeratitis, retinal damages,^{52,53} and corneal damages,⁵⁴ as this wavelength can damage DNA. UV-C radiation is also used in domestic air purifiers for air disinfection. Such wavelengths produce ozone, a molecule harmful to the environment but can also adversely affect the respiratory system.⁵⁵ Moreover, the UV-light sources are more expansive with respect to white light sources. Our work, conversely, aims at demonstrating for the first time that it is possible achieving disinfection by using a white light source that is safer for health and the environment and more affordable with respect to UV lamps.

To demonstrate the extraordinary capability of AuNR substrates to perform white light-assisted PT bacteria disinfection, experiments were carried out as follows:

AuNR substrates were immersed in an *E. coli* solution containing 10^4 CFU mL⁻¹ concentration for 30 min.

The substrates were then irradiated for 10 min under white light to induce the PT killing of *E. coli* cells.

The samples were immersed for 10 min in minimal E medium with the fluorescent dye Propidium Iodide to detect dead cells by fluorescence microscopy.

The absorption spectroscopy analysis of the substrate (Fig. 7a, black curve) performed after step 1 revealed a 10 nm red shift of the LSPRI (Fig. 7a, green curve), accounting for the presence of *E. coli* cells on the substrate. After steps 2 and 3, the absorption spectrum highlighted a 35 nm red shift (Fig. 7a, red curve) associated with the presence of propidium iodide dye,

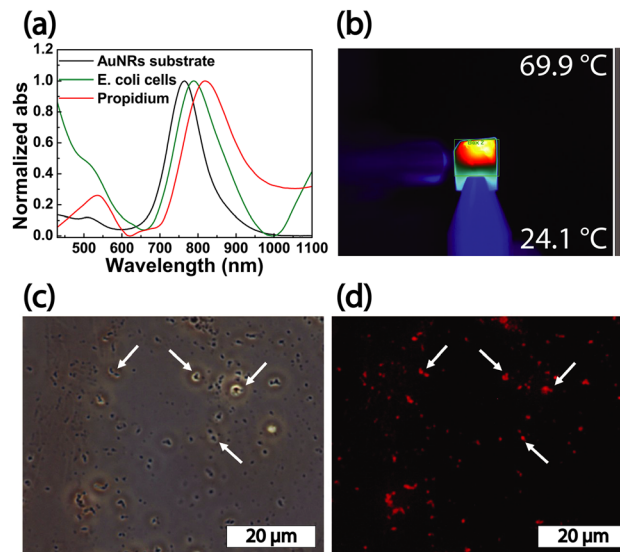


Fig. 7 Optical characterization of the AuNRs substrate, performed in all the steps of the PT disinfection experiment (a): after the AuNRs substrate preparation (black track), after the immersion in the *E. coli* dispersion (green track), and after the staining with the propidium iodide fluorescent dye (red track). Before the irradiation with a white light source at a power density of 14.7 W cm^{-2} , the substrate was immersed for 30 min in the *E. coli* dispersion 10^4 CFU mL⁻¹ to simulate a form of bacterial contamination. Thermographic image of the contaminated AuNRs substrate acquired at the end of the irradiation time (b). Contrast phase (c) and fluorescence (d) microscopy images of the AuNRs substrate, after the white light PT disinfection experiment.

that rises the n of the medium, responsible for an additional peak at 500 nm. The white light PT disinfection experiments were performed at a power density of 14.7 W cm^{-2} for 10 min providing a maximum temperature value of 69.9 °C (Fig. 7b) sufficient to kill *E. coli* cells.³⁹ After the white light disinfection experiment, the substrate was investigated by contrast phase microscopy. The high magnification image (Fig. 7c) showed spots with high contrast that can be associated with the presence of *E. coli* cells characterized by their typical elongated morphology. By investigating the same area by fluorescence microscopy, spots with red fluorescence appeared (Fig. 7d). These features were detected in the same position occupied by the higher contrast spots. They, therefore, can be safely associated with dead *E. coli* cells stained with propidium iodide dye. Accordingly, we can safely point out that under the investigated experimental conditions, the PT properties of AuNRs substrate are suited to achieve the white light disinfection of the substrate contaminated by *E. coli* cells.

Control experiments were performed without white light-assisted PT disinfection in verifying the AuNR substrate's biocompatibility and to further support the results reported in Fig. 7. The AuNRs substrate (a fresh sample) was immersed in 10^4 CFU mL⁻¹ *E. coli* cells solution for 30 min and, subsequently, for 10 min in Minimal E with fluorescent dye SYTO 9 that only binds to live bacteria cells. The absorption spectroscopy analysis reported in Fig. 8a (AuNRs sample, black curve; AuNRs sample + *E. coli* cells, red curve) shows an overall



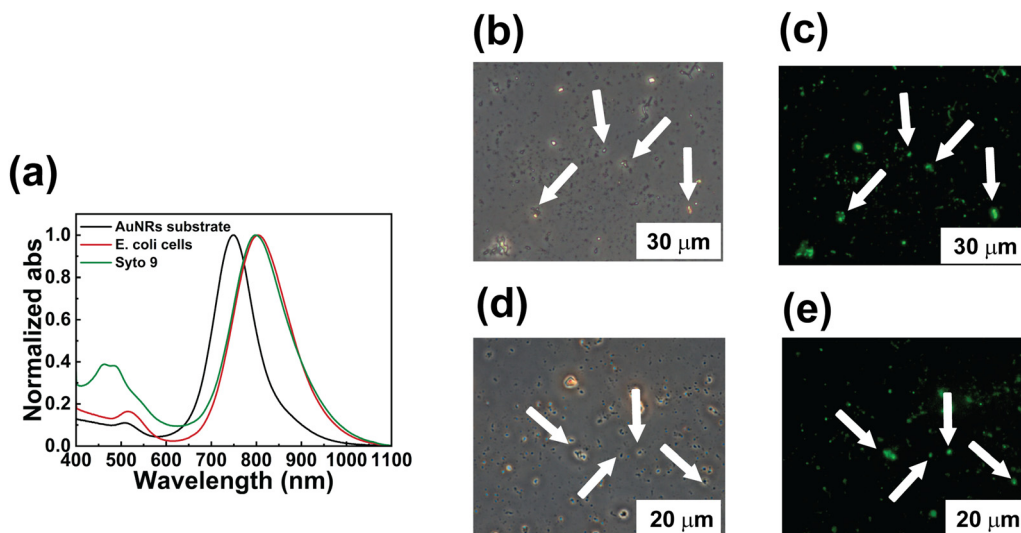


Fig. 8 Control experiment performed to assess the *E. coli* viability after a 30 min contact with AuNR substrate without irradiation. Absorption spectroscopy characterization of the AuNR substrate performed in each step of the control experiment (a). Contrast phase (b and d) and fluorescent microscopy images (c and e) at low (b and c) and high (d and e) magnification.

red-shift of about 55 nm (Fig. 8a, green curve) after the experimental steps, plus the peak of the absorption of the green fluorescent dye. Contrast phase (Fig. 8b and d) and fluorescent micrographs (Fig. 8c and e) acquired at different magnifications show a distribution of high contrast and green spots, respectively, corresponding to the live *E. coli* cells, thus demonstrating that the bacteria killing can be achieved only under proper irradiation conditions and cannot be ascribed to the mere contact with the AuNRs substrate.

Disinfection of medical tools from bacterial contamination

Nanotechnology plays a pivotal role in improving the performance of disinfection applications. AuNRs can convert the absorbed visible light into heat. As a result, the plasmonic platform constituted by AuNRs array deposited on a glass substrate acquires self-disinfectant abilities, under irradiation. The PT properties of AuNR substrates can be exploited to disinfect close surfaces through heat conduction, using the samples as PT transducers.

Biomedical tools used in surgery must be frequently sterilized, and the PT properties of AuNR substrates can be used for this purpose. This work developed a thermo-optical setup for testing AuNRs substrates' ability to achieve white light photo-thermal disinfection of biomedical tools. The optical setup (Fig. 9) was realized so that the AuNRs substrate is brought in close contact with the specific tool, and the illumination (heating generation) is performed through the AuNRs substrate. In this way, it is possible to avoid the direct deposition of AuNRs on the biomedical tool that has to be disinfected. The selected metallic biomedical devices are a bistoury, scissors, and a spatula (Fig. 10a–c, respectively). The disinfection ability of the platform is applied to the tools' surface through heat conduction that requires few minutes of illumination from a white light source to provide a uniform temperature distribution

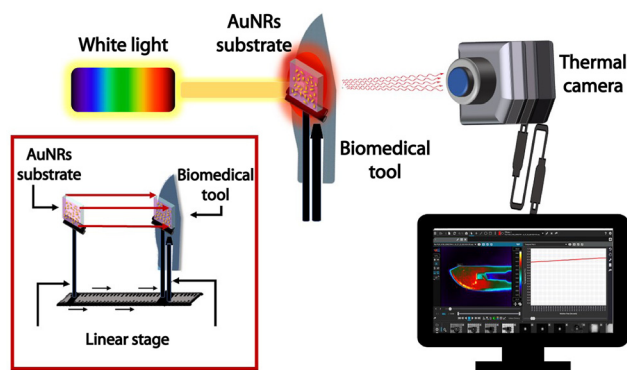


Fig. 9 Schematic of the thermo-optical setup implemented to disinfect the biomedical tools using the AuNR substrates as optical transducers. The sample holders that hold the biomedical tool and the AuNRs substrate are brought in close contact (within a few μm distances) using a high-precision linear stage.

according to the thickness and shape of the instruments. The PT experiments were performed by irradiating the samples for 10 min using the white light source (14.7 W cm^{-2}). The thermal camera measurements were collected on the back of the tools to verify the heat distribution through the thickness of the devices. High-temperature values of $67.1 \text{ }^\circ\text{C}$, $73.1 \text{ }^\circ\text{C}$, and $76.1 \text{ }^\circ\text{C}$ are respectively achieved in the areas of interest: the bistoury, the scissors, and the spatula in Fig. 10d–f, respectively. The corresponding temporal plots in Fig. 10g–i show how the different thicknesses impact the temperature rise: the maximum value is obtained after 2 min of irradiation of the bistoury and the spatula, while it requires almost 5 min for the thicker scissors. The temperature values and the heating distribution point out that the plasmonic substrate's disinfection properties can be effectively transferred to other metallic tools. The achieved temperature values are suitable for producing bacterial and viral



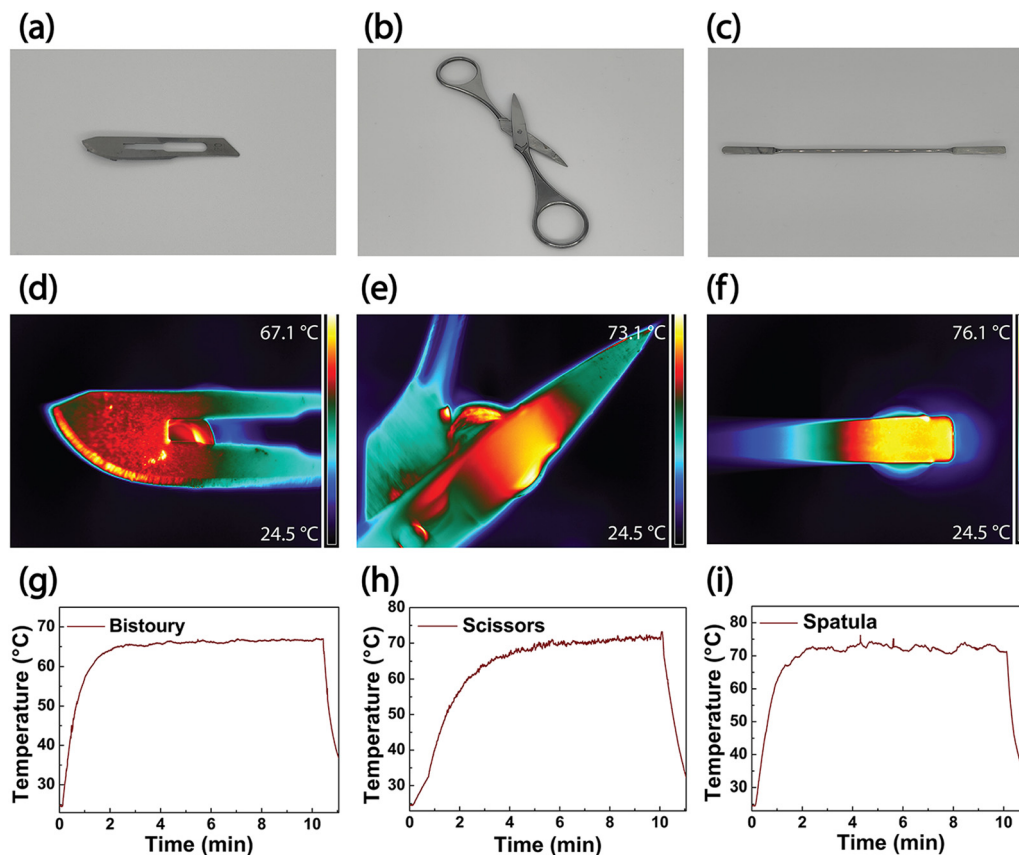


Fig. 10 Biomedical tools disinfection using the AuNR substrates as PT transducers. Bistoury (a), scissors (b), and a spatula (c) were adhered on the AuNRs substrate and irradiated by white light to achieve high-temperature values of 69.9, 65.4, and 68.3 °C, respectively. The thermal camera images show the uniform temperature distribution in the areas of interest for bistoury (d), scissors (e) and spatula (f) respectively. The corresponding time-temperature profiles are plotted in (g)–(i).

disinfection, as previously demonstrated in Fig. 7. Although the power density used in the present work is relatively high, our results pave the way to a valuable alternative to conventional disinfection methods. However, ongoing and future studies are devoted to developing a novel hybrid NPs generation with higher photothermal efficiency⁵⁶ in the visible spectrum.

Conclusions

In this work, we have reported a low-cost, reproducible white light-triggered PT transducer (thermo-plasmonic based) for energy-saving and facile bacteria disinfection. It exploits the properties of an AuNRs array adsorbed on a glass substrate through the immersive eLbL method enabling the PT disinfection of *E. coli* cells. Refractive index sensitivity experiments have demonstrated that both sides of the substrate are functionalized with AuNRs as a result of the immersive method. Thus, the absorption spectrum is given by both sides becoming highly sensitive to the AuNRs environment. The viability tests on *E. coli* cells have demonstrated that the plasmonic platform is biocompatible and that bacteria cell death is only due to white light irradiation. Moreover, the PT efficiency of 43.5% obtained under white light illumination confirms the effective

contribution of AuNRs array PT properties. Theoretical and experimental comparisons reveal good agreement, thus strengthening the experimental output in terms of *n* sensitivity and PT capabilities. Using the AuNR samples as white light-activated PT transducers, several medical tools such as a bistoury, scissors, and spatula have been illuminated and then thermally disinfected. Our findings are pioneering a new opportunity for healthcare facilities since the reported methodology allows non-hazardous disinfection of medical devices by simply using a conventional white light lamp. Our results pave the way for exploiting the PT properties of plasmonic NPs to achieve bacterial disinfection under solar light irradiation, thus pioneering a new solution for more environmentally friendly and affordable disinfection procedures, also improving the life quality in less developed countries.

Author contributions

FZ performed the photo-thermal experiments and analysed the data; PR, PZ, and DM implemented the theoretical studies; MLS prepared the AuNR substrates; SL, KJ realized the SEM and AFM characterization; NG, DE, JS, and MM provided support and feedback on samples realization and characterization; DD,



provided and prepared the bacteria culture; FP designed the methodology, performed the experiments, and supervised the work; LD conceived and formulated the idea, supervised the project, and acquired the funds. FZ, FP, and LD wrote the manuscript.

Conflicts of interest

There are no conflicts to declare.

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References

- 1 S. Mohapatra, *Essentials of Neuroanesthesia*, 2017, pp. 929–944, DOI: [10.1016/B978-0-12-805299-0.00059-2](https://doi.org/10.1016/B978-0-12-805299-0.00059-2).
- 2 W. Bäumlner, D. Eckl, T. Holzmann and W. Schneider-Brachert, *Crit. Rev. Microbiol.*, 2022, **48**, 531–564.
- 3 B. Z. Lakicevic, H. M. W. Den Besten and D. De Biase, *Front. Microbiol.*, 2021, **12**, 738470.
- 4 Environmental Health & Safety, (accessed 27/03/2023).
- 5 M. Cho, J. Kim, J. Y. Kim, J. Yoon and J. H. Kim, *Water Res.*, 2010, **44**, 3410–3418.
- 6 Z. T. Hu, Y. Chen, Y. F. Fei, S. L. Loo, G. Chen, M. Hu, Y. Song, J. Zhao, Y. Zhang and J. Wang, *Sci. Total Environ.*, 2022, **837**, 155720.
- 7 S. Giannakis, T. M. Le, J. M. Entenza and C. Pulgarin, *Water Res.*, 2018, **143**, 334–345.
- 8 I. De Pasquale, C. Lo Porto, M. Dell’Edera, F. Petronella, A. Agostiano, M. L. Curri and R. Comparelli, *Catalysts*, 2020, **10**, 1382.
- 9 D. J. Weber, W. A. Rutala, D. J. Anderson, L. F. Chen, E. E. Sickbert-Bennett and J. M. Boyce, *Am. J. Infect. Control*, 2016, **44**, e77–84.
- 10 F. Petronella, A. Truppi, T. Sibillano, C. Giannini, M. Striccoli, R. Comparelli and M. L. Curri, *Catal. Today*, 2017, **284**, 100–106.
- 11 N. López-Vinent, A. Cruz-Alcalde, G. Moussavi, I. del Castillo Gonzalez, A. Hernandez Lehmann, J. Giménez and S. Giannakis, *Chem. Eng. J.*, 2022, **450**, 137904.
- 12 G. Birgand, R. Ahmad, A. N. H. Bulabula, S. Singh, G. Bearman, E. C. Sánchez and A. Holmes, *Lancet*, 2022, **400**, 2250–2260.
- 13 G. M. Santos, F. I. d S. Ferrara, F. Zhao, D. F. Rodrigues and W.-C. Shih, *Opt. Mater. Express*, 2016, **6**, 1217–1229.
- 14 S. Maity, L. N. Downen, J. R. Bochinski and L. I. Clarke, *Polymer*, 2011, **52**, 1674–1685.
- 15 L. Jauffred, A. Samadi, H. Klingberg, P. M. Bendix and L. B. Oddershede, *Chem. Rev.*, 2019, **119**, 8087–8130.
- 16 Z. Tang, D. Ma, Q. Chen, Y. Wang, M. Sun, Q. Lian, J. Shang, P. K. Wong, C. He, D. Xia and T. Wang, *J. Hazard. Mater.*, 2022, **437**, 129373.
- 17 I. Hammami, N. M. Alabdallah, A. A. Jomaa and M. Kamoun, *J. King Saud Univ., Sci.*, 2021, **33**, 101560.
- 18 S. Loeb, C. Li and J.-H. Kim, *Environ. Sci. Technol.*, 2018, **52**, 205–213.
- 19 J. Cao, T. Sun and K. T. V. Grattan, *Sens. Actuators, B*, 2014, **195**, 332–351.
- 20 Y. Ziai, F. Petronella, C. Rinoldi, P. Nakielski, A. Zakrzewska, T. A. Kowalewski, W. Augustyniak, X. Li, A. Calogero, I. Sabała, B. Ding, L. De Sio and F. Pierini, *NPG Asia Mater.*, 2022, **14**, 18.
- 21 A. Jahangiri-Manesh, M. Mousazadeh, S. Taji, A. Bahmani, A. Zarepour, A. Zarrabi, E. Sharifi and M. Azimzadeh, *Pharmaceutics*, 2022, **14**, 664.
- 22 J. Zhao, S. C. Nguyen, R. Ye, B. Ye, H. Weller, G. A. Somorjai, A. P. Alivisatos and F. D. Toste, *ACS Cent. Sci.*, 2017, **3**, 482–488.
- 23 S. Liao, W. Yue, S. Cai, Q. Tang, W. Lu, L. Huang, T. Qi and J. Liao, *Front. Pharmacol.*, 2021, **12**, 664123.
- 24 V. Frantellizzi, V. Verrina, C. Raso, M. Pontico, F. Petronella, V. Bertana, A. Ballesio, S. L. Marasso, S. Miglietta, P. Rosa, S. Scibetta, V. Petrozza, M. S. De Feo, G. De Vincentis, A. Calogero, R. Pani, G. Perotto and L. De Sio, *Mater. Today Adv.*, 2022, **16**, 100286.
- 25 A. G. Al-Bakri and N. N. Mahmoud, *Molecules*, 2019, **24**, 2661.
- 26 C. Fang, L. Shao, Y. Zhao, J. Wang and H. Wu, *Adv. Mater.*, 2012, **24**, 94–98.
- 27 S. K. Loeb, J. Kim, C. Jiang, L. S. Early, H. Wei, Q. Li and J. H. Kim, *Environ. Sci. Technol.*, 2019, **53**, 7621–7631.
- 28 F. Petronella, D. De Biase, F. Zaccagnini, V. Verrina, S.-I. Lim, K.-U. Jeong, S. Miglietta, V. Petrozza, V. Scognamiglio, N. P. Godman, D. R. Evans, M. McConney and L. De Sio, *Environ. Sci.: Nano*, 2022, **9**, 3343–3360.
- 29 G. A. Evtugyn and T. Hianik, *Curr. Anal. Chem.*, 2011, **7**, 8–34.
- 30 X. Hu, Y. Zhang, T. Ding, J. Liu and H. Zhao, *Front. Bioeng. Biotechnol.*, 2020, **8**, 990.
- 31 K. M. Mayer and J. H. Hafner, *Chem. Rev.*, 2011, **111**, 3828.
- 32 L. Martinez Maestro, P. Haro-González, J. Coello and D. Jaque, *Appl. Phys. Lett.*, 2012, **100**, 201110.
- 33 X. Zan and Z. Su, *Langmuir*, 2009, **25**, 12355.



- 34 M. E. Martínez-Hernández, X. Sandúa, P. J. Rivero, J. Goicoechea and F. J. Arregui, *Chemosensors*, 2022, **10**(2), 77.
- 35 R. A. Ghostine, R. M. Jisr, A. Leahaf and J. B. Schlenoff, *Langmuir*, 2013, **29**, 11742.
- 36 E. Guzmán, R. G. Rubio and F. Ortega, *Adv. Colloid Interface Sci.*, 2020, **282**, 102197.
- 37 X. Huang and M. A. El-Sayed, *J. Adv. Res.*, 2010, **1**, 13–28.
- 38 C. C. Chuang, C. C. Cheng, P. Y. Chen, C. Lo, Y. N. Chen, M. H. Shih and C. W. Chang, *Int. J. Nanomed.*, 2019, **14**, 181–193.
- 39 F. Annesi, A. Pane, M. A. Losso, A. Guglielmelli, F. Lucente, F. Petronella, T. Placido, R. Comparelli, M. G. Guzzo, M. L. Curri, R. Bartolino and L. De Sio, *Materials*, 2019, **12**, 1530.
- 40 H. Zhang, H.-J. Chen, X. Du and D. Wen, *Sol. Energy*, 2014, **100**, 141–147.
- 41 J. Lin and D. Shi, *Appl. Phys. Rev.*, 2021, **8**, 011302.
- 42 J. Lin, Y. Zhao and D. Shi, *MRS Commun.*, 2020, **10**, 155–163.
- 43 A. Guglielmelli, F. Pierini, N. Tabiryan, C. Umeton, T. J. Bunning and L. De Sio, *Adv. Photonics Res.*, 2021, **2**, 2000198.
- 44 L. M. Maestro, Q. Zhang, X. Li, D. Jaque and M. Gu, *Appl. Phys. Lett.*, 2014, **105**, 181110.
- 45 T. d L. Pedrosa, S. Farooq and R. E. de Araujo, *Nanomaterials*, 2022, **12**, 4188.
- 46 J. W. Strutt, *Philos. Mag.*, 1871, **41**, 447–454.
- 47 C. F. Bohren and D. R. Huffman, *Absorption and scattering of light by small particles*, 1983.
- 48 T. Yamaguchi, S. Yoshida and A. Kinbara, *Thin Solid Films*, 1974, **21**, 173–187.
- 49 P. Royer, J. L. Bijeon, J. P. Goudonnet, T. Inagaki and E. T. Arakawa, *Surf. Sci.*, 1989, **217**, 384–402.
- 50 P. e. a. Radomski, Polish Academy of Science (PAN), Bialystok 2022.
- 51 P. J. Roache, *Annu. Rev. Fluid Mech.*, 1997, **29**, 123–160.
- 52 M. O. P. Alvarenga, S. R. M. Veloso, A. Bezerra, B. P. Trindade, A. S. L. Gomes and G. Q. M. Monteiro, *J. Photochem. Photobiol.*, 2022, **9**, 100096.
- 53 T. Dai, M. S. Vrahas, C. K. Murray and M. R. Hamblin, *Expert Rev. Anti-Infect. Ther.*, 2012, **10**, 185.
- 54 S. Kaidzu, K. Sugihara, M. Sasaki, A. Nishiaki, T. Igarashi and M. Tanito, *Free Radical Res.*, 2019, **53**, 611.
- 55 California Air source board, Hazardous Ozone-Generating Air Purifiers <https://ww2.arb.ca.gov/our-work/programs/air-cleaners-ozone-products/hazardous-ozone-generating-air-purifiers> (accessed 8 June 2023).
- 56 Z. Zhang, J. Wang and C. Chen, *Theranostic*, 2013, **3**(3), 223–238.

