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Synergistic coupling of interface ohmic contact and LSPR effects over Au/Bi₂₄O₃₁Br₁₀ nanosheets for visible-light-driven photocatalytic CO₂ reduction to CO†

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The challenge of synergistically optimizing different mechanisms limits the further improvement of plasmon-mediated photocatalytic activities. In this work, an Au/Bi₂₄O₃₁Br₁₀ composite, combining an interface ohmic contact and localized surface plasmon resonance (LSPR), is prepared by a thermal reduction method. The LSPR effect induces the local resonance energy transfer effect and the local electric field enhancement effect, while the interface ohmic contact forms a stronger interface electric field. The novel synergistic interaction between the interface ohmic contact and LSPR drives effective charge separation and provides more active sites for the adsorption and activation of CO₂ with improved photocatalytic efficiency. The optimized 0.6 wt% Au (5.7 nm) over Bi₂₄O₃₁Br₁₀ nanosheets showed an apparently improved photocatalytic activity without any sacrificial reagents, specifically CO and O₂ yields of 44.92 and 17.83 μmol g⁻¹ h⁻¹, and demonstrated superior stability (only lost 6%) after continuous reaction for 48 h, nearly 5-fold enhanced compared to Bi₂₄O₃₁Br₁₀ and a great advantage compared with other bismuth-based photocatalysts.

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1. Introduction

Since the second industrial revolution, human beings have been in a state of over-use of fossil fuels, resulting in an energy crisis and severe environmental pollution problems.¹ So far, a variety of solutions have been proposed. Among them, the photocatalytic CO₂ reduction/conversion process can be regarded as the most promising one, and the driving force of this reaction is inexhaustible light energy, so it has attracted the interest of a large number of researchers.² To date, many kinds of photocatalysts for CO₂ reduction have been reported, such as transition metal oxides, transition metal sulfides, and metal-organic frameworks. The factors that affect the photocatalytic performance include: the spectral response range,³ the separation and migration of photogenerated charges,⁴ surface chemical reaction.⁵ Therefore, the construction of photocatalysts with high photoenergy efficiency, strong photogenerated charge transfer ability and a large specific surface area is of great significance for the photoreduction of CO₂.

Among the reported photocatalysts, two-dimensional (2D) ultrathin nanosheets have attracted much attention because of

their excellent photoinduced charge separation and transfer efficiency, large specific surface area, lots of surface-active sites and better substrate interaction. Therefore, it is of great research value to construct a highly efficient composite based on a 2D photocatalyst. Among many photocatalysts, BiOX (X = Cl, Br and I) easily form a nanosheet structure due to their unique “sandwich configuration” (X⁻-[Bi₂O₂]²⁺-X⁻).⁶ And their band gap is relatively narrow, and the utilization rate of sunlight is relatively high, which has aroused widespread concern. However, bismuth-based photocatalytic materials have a high photogenerated charge recombination rate and easy photo-corrosion, which limits their applications in practical applications. Cui *et al.* summarized three key strategies for improving photocatalytic efficiency through extensive experiments, such as composition adjustment, vacancy engineering and preparation of heterostructures.⁷ As a commonly used method to improve photocatalytic activity, the bismuth-rich strategy has constructed a series of semiconductors,⁸ such as Bi₅O₇Br,⁹ Bi₂₄O₃₁Cl₁₀ (ref. 10) and Bi₂₄O₃₁Br₁₀ (ref. 11), which show satisfactory photocatalytic performance under visible light. Among them, Bi₂₄O₃₁Br₁₀ ($E_g = 2.51$ eV), as a new bismuth halide compound, has a smaller band gap and better performance than BiOBr ($E_g = 2.75$ eV). Its crystal structure is composed of BiO₅ polyhedrons and Br⁻ ions.¹² Recently, Bi₂₄O₃₁Br₁₀ has been used in some photocatalytic systems.¹³⁻¹⁵ For example, Xiao *et al.* used Bi₂₄O₃₁Br₁₀ to realize the high

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selective conversion of benzyl alcohol to benzaldehyde.¹¹ Wang *et al.* used Bi₂₄O₃₁Br₁₀ as a catalyst to realize the efficient degradation of tetracycline hydrochloride.⁴ However, the weak photocatalytic reduction ability and high photogenerated charge recombination rate hinder the further development of Bi₂₄O₃₁Br₁₀ in the field of photocatalysis. To overcome these shortcomings, it is necessary to further develop Bi₂₄O₃₁Br₁₀.

One of the most attractive characteristics of noble metals (*e.g.*, Au, Ag, *etc.*) is the nanoscale local surface plasmon resonance (LSPR) effect, which is produced by conducting electrons in response to the collective oscillation of incident light under resonance conditions. Due to the LSPR effect, plasma metal is a kind of highly sensitive material for photothermal catalysis. Once exposed to light, they promote the activity of chemical reactions in the following ways: (1) local electric field enhancement (LEFE),¹⁶ (2) local resonance energy transfer (LRET),¹⁷ and (3) photothermal effect (PE).¹⁸ The LEFE can observably improve the band transition rate of the adjacent semiconductor, if the resonance energy of the plasma metal surpasses the band gap of the semiconductor.^{16,19} Furthermore, the plasmonic oscillation energy can be delivered *via* LRET to the semiconductor photocatalyst by local electric field or non-radiative dipole–dipole coupling, which requires spectral overlap between the LSPR peak and semiconductor light absorption range.^{17,20} The surface plasmon can improve the generation and separation of e⁻/h⁺ pairs through these two mechanisms. Ma *et al.* prepared core–shell (Au/AgAu)@CdS hybrid nanostructures with strong multi-interface plasmon coupling. Strong plasmon absorption and near-field enhancement caused by multi-interface plasmon coupling provide more plasmon energy, which significantly promotes the generation and separation of e⁻/h⁺ pairs in catalysts.²¹ Wang *et al.* have developed a simple surfactant-free nano-processing method, which limits Au nanoparticles to N-doped TiO₂ nano-bowl arrays. N-doping extends the optical absorption of TiO₂ to the visible region, which better overlaps the plasmon bands of Au nanoparticles. This design cooperates well with the Schottky junction and the plasmon LRET process, thus significantly improving the performance of photocatalytic hydrogen production in the full spectrum range.²² In addition, due to different work functions, the noble metal and the base photocatalyst will form Schottky contacts^{22,23} and ohmic contacts^{24,25} at the coupling interface, which will lead to energy band bending, which is beneficial to the migration of photogenerated electrons, thus improving the separation efficiency of charges. As far as we know, few people have studied the synergistic effect between the LSPR effect and interface contact on photocatalytic performance, and the internal mechanism is not clear.

In the work, 0D/2D Au/Bi₂₄O₃₁Br₁₀ (Au/BOB) was synthesized by a thermal reduction method, and the impact of the synergy between the LSPR of Au NPs and ohmic contact between interfaces on the photocatalytic reduction of CO₂ was investigated. Surprisingly, the CO yield of the composite can reach 44.92 μmol g⁻¹ h⁻¹, and the yield of CO only decreased 6% after continuous reaction for 48 h in the absence of sacrificial agents and photosensitizers. Emphatically, since the plasma band of Au NPs was mainly located in the visible band, the coupling between Au and

BOB can effectively improve the utilization of solar energy through the LRET mechanism. Therefore, the optical absorption of BOB was no longer restricted by the width of the bandgap. Multi-physical field coupling analysis (COMSOL) proved that the LEFE effect can significantly promote the generation and separation of photogenerated charges. In addition, the interface ohmic contact induces a strong interface electric field, which can directly improve the photogenerated charge transfer process and drive the photoelectron transfer from BOB to Au NPs, which is confirmed by DFT calculation. The synergistic effect of the interface ohmic contact and LSPR effect makes it possible to further optimize the photocatalyst, prolong the light absorption wavelength and improve the charge separation.

2. Experimental section

2.1. Chemicals and reagents

All reagents were used directly without further purification. Bismuth nitrate pentahydrate (Bi(NO₃)₃·5H₂O, AR), cetyl trimethyl ammonium bromide (CTAB, AR), sodium hydroxide (NaOH, AR), and methanol (99.5%) were purchased from Xi long Chemical Co., Ltd. Chlorauric acid tetrahydrate (HAuCl₄·4H₂O, AR) and silicon dioxide (SiO₂) were purchased from Aladdin Reagent (Shanghai) Co., Ltd. Sodium borohydride (NaBH₄, 96%) and anhydrous ethanol (AR) were purchased from Sinopec Chemical Reagent Co., Ltd. Deionized water was used throughout the whole experiment.

2.2. Synthesis of 2D Bi₂₄O₃₁Br₁₀ (BOB) ultrathin nanosheets

Firstly, 0.243 g of bismuth nitrate pentahydrate and 0.182 g of cetyl trimethyl ammonium bromide (CTAB) were added into 20 mL of deionized water. The suspension was ultrasonically treated for 10 min, and then stirred for 15 min until it was dissolved. Then NaOH aqueous solution (6 M) was added to adjust the pH value of the obtained solution to about 9. Then the uniform mixture was transferred to a Teflon lined stainless steel autoclave, which was held in a heated oven at 160 °C for 16 h. After naturally cooling down to room temperature, the solid was filtered, and washed several times with ethanol, and vacuum dried at 60 °C overnight to obtain (BOB) ultrathin nanosheets.

2.3. Synthesis of the Au NPs/Bi₂₄O₃₁Br₁₀ (Au/BOB) composite

100 mg of BOB prepared in the above experiment was added to 30 mL of ethanol solution (v(H₂O) : v(ethanol) = 1 : 1) and stirred at room temperature for 10 min to obtain a uniform suspension. Then, 12 mL of HAuCl₄ (50 mg L⁻¹) and NaOH (1 M) were added in turn and stirred for 30 min in the dark to ensure uniform dispersion. Then, 0.05 mL of NaBH₄ (112 mM) was dropwise added under vigorous stirring, and the reaction was carried out at room temperature for 3 h. The product was washed with ethanol and ultrapure water three times and then heated in a vacuum drying oven at 60 °C for 12 hours. Various mass ratios of Au NPs were investigated for optimization experiments, including 0.3, 0.6 and 1.2 wt%, and the samples are abbreviated as XAu/BOB, where X = 0.3, 0.6 and 1.2.



2.4. Photocatalytic CO₂ production experiments

The photocatalytic reduction CO₂ experiment was carried out in a 500 mL sealed quartz reactor, which can be irradiated vertically. The light source was a 300 W xenon lamp equipped with a 420 nm cut-off filter (CEL-HXF300-T3). Briefly, 50 mg photocatalyst powder was added to 100 mL aqueous solution. Prior to illumination, the reactor was degassed to completely remove air and dissolved oxygen by continuously blowing high-concentration CO₂ gas. During illumination, the system temperature was maintained at about 5 °C by circulating condensed water. The photocatalytic reduction CO₂ reaction lasted for 4 h, and the gas sample was collected every 1 h. Finally, the product was quantified using a gas chromatograph (Agilent 8890 GC) equipped with a FID detector, TCD detector and column (TDX-01). The product ¹³CO obtained from ¹³CO₂ reduction was confirmed by using a dual instrument of gas chromatography and mass spectrometry (GC-MS, Agilent 7890A GC and Agilent 5975C MS).

2.5. Characterization

The surface morphology of the samples was examined by using a field emission scanning electron microscope (FESEM) (Nova-NaNoSEM450, operation voltage 10 kV). The microstructure and elemental distribution were characterized by field emission transmission electron microscopy (FETEM, FEI Talos F200X), high-resolution TEM (HRTEM), high-angle annular dark field scanning TEM (HAADF-STEM) and selected area energy dispersive X-ray spectroscopy (EDX) mapping. The thickness and surface potential difference of the nanosheets were measured by atomic force microscopy (AFM, NT-MDT, NTEGRA). The operating conditions were an acceleration voltage 200 kV, a LaB6 filament, and a resolution of 0.235 nm. The optical properties of the samples were studied by using a UV-vis DIFFUSE reflectance spectrometer (UV-vis DRS, Shimadzu UV-2450) with BaSO₄ as a reference. The surface chemical composition was analyzed by X-ray photoelectron spectroscopy (XPS, Al Ka, 150 W, and C 1s 284.8 eV). The separation and migration rate of photogenerated charges were analyzed by using photoluminescence (PL) and the fluorescence lifetime spectrum (TRPL) obtained with a fluorescence spectrometer (Hitachi-F7000, 420 nm excitation). The Brunauer–Emmett–Teller (BET) method was employed to measure the surface area and pore size distribution of the samples using a NDVA2000e instrument with N₂ as the adsorbate at 77 K. The CO₂ adsorption capacity of the sample was evaluated by using a NDVA2000e instrument. The temperature programmed desorption data of CO₂ were recorded by using a TCD detector on an AutoChem II 2920 analyzer (measured flow rate: 1.339 mmol min⁻¹). *In situ* infrared spectroscopy (*In situ* FTIR, Perkin Elmer-3100) was used to determine the change trend of intermediates on the sample surface with time during photocatalytic CO₂ reduction. Electron paramagnetic resonance (EPR) data were recorded by using a Bruker A300 spectrometer (state frequency 9.54 GHz, modulation amplitude 2 G, locking frequency 100 kHz, and attenuation 30 dB). Multi-physical coupling analysis software (COMSOL Multiphysics) was used to analyze the local electric field of the sample surface.

2.6. Photoelectrochemical and electrochemical measurements

To study the photochemical properties of all samples, photocurrent was measured in a three-electrode system, where a platinum wire, Ag/AgCl (saturated KCl) and FTO substrate coated with the photocatalyst served as the counter electrode, reference electrode, and working electrode, respectively. The specific experiment was carried out using an electrochemical working station (CH2660C). The working electrode area was 2 cm² (1 cm × 2 cm) and the electrolyte was 0.1 M Na₂SO₄ solution. A 300 W xenon lamp (CEL-HXF300-T3) was the light source. The lamp was switched on and off to measure photocurrent under light irradiation and in the dark. Electrochemical impedance spectroscopy (EIS) was performed using the same working station. The Mott–Schottky curve (M–S) was scanned in 0.1 M Na₂SO₄ solution with a frequency range of 0.1 Hz to 100 kHz. The test amplitude of AC voltage was set at 1 V. In addition, the photoelectrochemical properties of the samples were studied in a three-electrode system, in which carbon rods, calomel electrodes and carbon cloth substrates coated with photocatalysts were used as counter electrodes, reference electrodes and working electrodes, respectively. The specific experiments were carried out on an electrochemical workbench (CH2660C). The working electrode area was 1 cm² (1 cm × 1 cm) and the electrolyte was 1 M KOH solution.

2.7. DFT calculation

In the density functional theory calculations, we built a planar model of Bi₂₄O₃₁Br₁₀ (002) and Bi₂₄O₃₁Br₁₀ (002) loading Au NPs. The atom number of Au NPs is 20. The vacuum layer along the Z direction was set to be 15 Å to avoid the interactions between images. Structural optimization was performed by using the Vienna *Ab initio* Simulation Package (VASP)²⁶ with the projector augmented wave (PAW) method.²⁷ The Perdew–Burke–Ernzerhof (PBE)²⁸ exchange–correlation functional with the generalized gradient approximation (GGA)²⁹ was adopted to describe the electronic exchange–correlation energy. The cut-off energy of the plane-wave basis was set at 450 eV. For the optimization of both geometry and lattice size, the Brillouin zone integration is performed with 2 × 2 × 1 Gamma *k*-point sampling.³⁰ The self-consistent calculations applied a convergence energy threshold of 10⁻⁵ eV. The equilibrium geometries and lattice constants were optimized with maximum stress on each atom within 0.01 eV Å⁻¹. Dipole correction was adopted along the Z direction in the calculation of the work function to describe the unsymmetrical surface. The isosurface level of the charge density difference was set to be 5 × 10⁻⁴ eÅ⁻³.

The Gibbs free energy of intermediates was calculated as $\Delta G = \Delta E + \Delta E_{zpe} - T\Delta S$, where ΔE is the electronic energy difference directly obtained by DFT calculation, and ΔE_{zpe} and ΔS indicate the zero-point energy and entropy of the surface adsorbing H atom, respectively. The entropies of the free molecules of H₂ and H₂O were referenced from the NIST database.³¹ Besides, the Kelvin temperature *T* was set to be 298.15 K. The reaction energy in photocatalytic CO₂ reduction is the reaction free energy of the potential limiting step. The ΔE_{zpe} and ($-T\Delta S$) of all

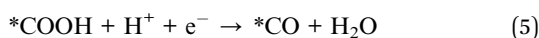


intermediates in the CO₂ reduction on Au/BOB are calculated by the following equation:

$$\Delta E_{zpc} = \frac{1}{2} \sum_i h\nu_i \quad (1)$$

$$-T\Delta S = K_B T \sum_i \ln \left(1 - e^{-\frac{h\nu_i}{K_B T}} \right) - \sum_i h\nu_i \left(\frac{1}{\frac{h\nu_i}{e^{K_B T}} - 1} \right) \quad (2)$$

where h , ν and K_B represent Planck's constant, vibrational frequencies, and Boltzmann's constant, respectively. The reduction pathway from CO₂ to CO was calculated according to:



3. Results and discussion

3.1. Structure characterization of the photocatalysts

Fig. 1a shows that Au NPs were *in situ* grown on the surface of 2D BOB nanosheets by impregnation reduction, and the details and particular characterization of the modified 0D/2D Au/BOB composite are displayed in the ESI.† The prepared samples are proved to be monoclinic BOB by XRD (JCPDS No. 75-0888),³² while no obvious diffraction peaks belonging to Au NPs are observed on 0.6Au/BOB, proving the high degree of fine dispersion of Au NPs (Fig. S1†). The content of Au NPs in 0.6Au/BOB was measured using an inductively coupled plasma analyzer. The results are consistent with expectations, indicating that Au NPs were successfully recombined in Au/BOB (Table S1†). Besides, the chemical structure information of BOB and 0.6Au/BOB is compared by FT-IR (Fig. S2†). The results show that the 0.6Au/BOB composite shows a new characteristic peak at 661 cm⁻¹, corresponding to the Au–Bi–O vibrational dynamics, which further proves the successful loading of Au NPs.³³

In order to intuitively observe the loading of Au NPs on the surface of BOB nanosheets, the microstructure of BOB and 0.6Au/BOB was obtained by SEM and TEM. As shown in Fig. S3a and b,† both BOB and 0.6Au/BOB exist in the form of irregular nanosheets. In addition, atomic force microscopy (AFM) further confirmed the nanosheet structure of BOB and 0.6Au/BOB (Fig. S4a and b†). The results show that BOB and 0.6Au/BOB are both ultrathin nanosheet structures with thickness close to 6.6 nm, which indicate that the loading of Au NPs will not change the morphology and size of the substrate. TEM further confirmed the results of SEM and AFM. As shown in Fig. 1b and c, BOB exists in the form of irregular nanosheets with no excess longitudinal clutter, which is conducive to the exposure of the active surface and the loading of Au NPs. There are a large

number of Au NPs in the TEM image of 0.6Au/BOB (Fig. 1d and e). In addition, through the statistical distribution, the diameter of Au NPs is calculated to be about 5.7 nm (Fig. S3c†). From the HRTEM spectra of 0.6Au/BOB (Fig. 1e), it can be clearly seen that the crystal plane spacings of 0.287 nm and 0.232 nm correspond to the (117) and (111) crystal planes of BOB and Au NPs, respectively. HRTEM indirectly shows that Au NPs were successfully loaded on BOB nanosheets. In addition, as shown in Fig. 1f and g, HAADF-STEM and EDX mappings on the selected areas indicate that the distributions of Bi, O and Br elements are uniform. And also, the uniform distribution of Au elements on the nanosheets can also be observed.³⁴ The above results show that Au NPs are very finely dispersed on BOB, which explains why the specific surface area of the composites increases with the loading of Au NPs (Fig. S5a and b, Table S2†). As a contrast, the same content of Au NPs (0.6Au/SiO₂) is loaded on the surface of commercial SiO₂, and the results show that the size of Au NPs has not changed significantly (Fig. S6†), which provides more reference data for the proof of the synergy effect and reduces the error of external experiments.

To further verify the chemical interaction between Au NPs and BOB, 0.6Au/BOB was tested by XPS and compared with BOB. As shown in Fig. S7a,† the 0.6Au/BOB composite determines the presence of Bi 4f, Br 3d, O 1s and Au 4f elements, which proves the successful loading of Au NPs. Fig. S7b† shows that BOB has two symmetry peaks at 159.20 and 164.45 eV, which are related to the characteristics of Bi 4f_{7/2} and Bi 4f_{5/2} of Bi³⁺. The characteristic peak of the 0.6Au/BOB composite shifts 0.1 eV to a lower value. As shown in Fig. S7c and d,† compared with BOB, the characteristic peaks of O 1s and Br 3d in the 0.6Au/BOB composite shift negatively and the binding energy decreases. The result further confirms that there is a strong interfacial interaction between Au NPs and BOB. In addition, the XPS spectrum of Au 4f, as shown in Fig. S7e,† is located at the peaks of 87.45 and 83.74 eV, corresponding to Au 4f_{5/2} and Au 4f_{7/2}, respectively, which is consistent with the results in the literature, and confirms that the Au element exists in the composite in the form of Au⁰.³⁵

3.2. Photophysical and electrochemical measurements

The optical properties of BOB and Au/BOB were determined using UV-vis DRS. As shown in Fig. 2a, pristine BOB exhibits a visible light response, and the edge is about 470 nm. The corresponding band gap is 2.54 eV (Fig. S8†), which is consistent with a previous report.³⁶ Interestingly, Au/BOB has an obvious wide peak near 570 nm, which may be caused by the LSPR effect of Au NPs,^{17,37} assigned to quadrupole and dipole plasmon resonance signals. In addition, Fig. 2a shows a broadened and enhanced light absorption edge of Au/BOB compared with BOB, indicating that the Au NPs form impurity energy levels within the band gap, effectively reduce the band gap, and improve the absorption of visible light. The DOS curve proves the correctness of the results (Fig. S9†).³⁸ Moreover, if the visible absorption spectrum of BOB partially overlaps with the LSPR peak of Au NPs, this may lead to the formation of a local resonance energy transfer (LRET) mechanism, similar to that of hot



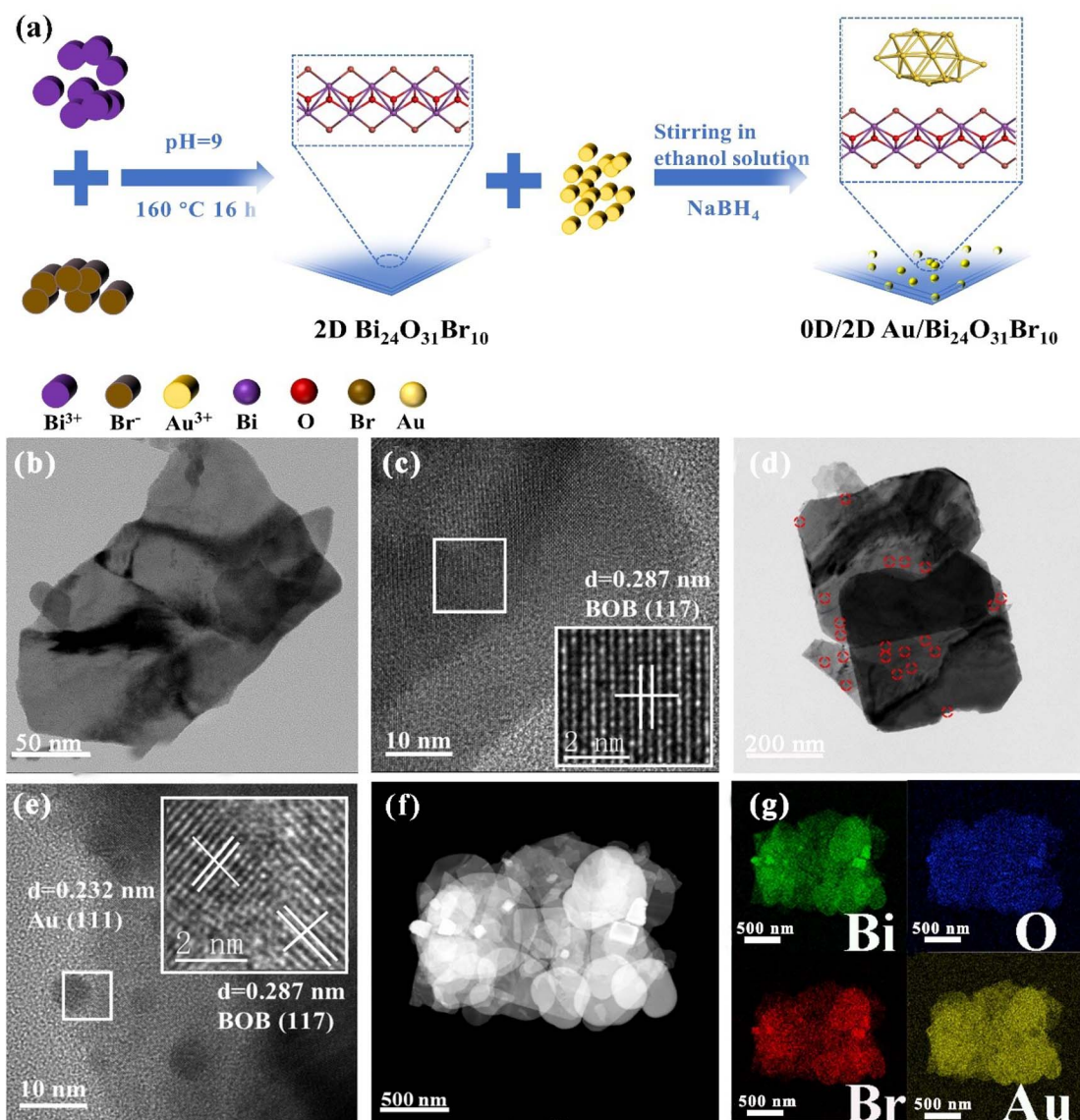


Fig. 1 (a) Synthesis schematic diagram of 0D/2D Au/BOB; TEM and HRTEM images of (b and c) BOB and (d and e) 0.6Au/BOB; HAADF-STEM image (f) and corresponding elemental mappings (g) of 0.6Au/BOB.

electrons at the interface.^{17,39} Therefore, we investigated the dependence of CO yield on light intensity at $\lambda > 570$ nm, and the results are shown in Fig. S10.† As observed, the rate of CO generation was linearly correlated with light intensity, indicating the presence of the LRET mechanism between the interfaces. Thus, the energy acquisition of the BOB CB can be extended beyond the spectral response range, which enhances visible light absorption, photogenerated charge generation and effective charge separation.

The photocatalytic properties of BOB and 0.6Au/BOB were studied by transient photocurrent response and electrochemical impedance spectroscopy (EIS). As shown in Fig. 2b, the photocurrent response of pure BOB is very weak, which is due to its fast charge recombination. The optical current intensity of the Au/BOB composite has been greatly improved. Among them, the average optical current intensity of 0.6Au/BOB

is $1.12 \mu\text{A cm}^{-2}$, which is about 11 times that of pure BOB. Therefore, the composite structure of 0.6Au/BOB can effectively separate the photogenerated charge, thus improving the photocatalytic performance. In addition, the near-field electromagnetic coupling of Au NPs also plays an important role, which shows that the yield of photogenerated charges in semiconductor BOB is improved in the vicinity of Au NPs with the LSPR effect. At the same time, electrochemical impedance spectroscopy (EIS) measurement shows that the radius of 0.6Au/BOB is smaller than that of BOB (Fig. 2c), implying less interface resistance for carrier transfer due to the improvement of electron mobility, which is consistent with the photocurrent results. All these data together determine the fast electron transfer in 0.6Au/BOB, which means that the formation of composites promotes charge separation. To confirm this improved charge separation, photoluminescence spectra (PL) and time-resolved



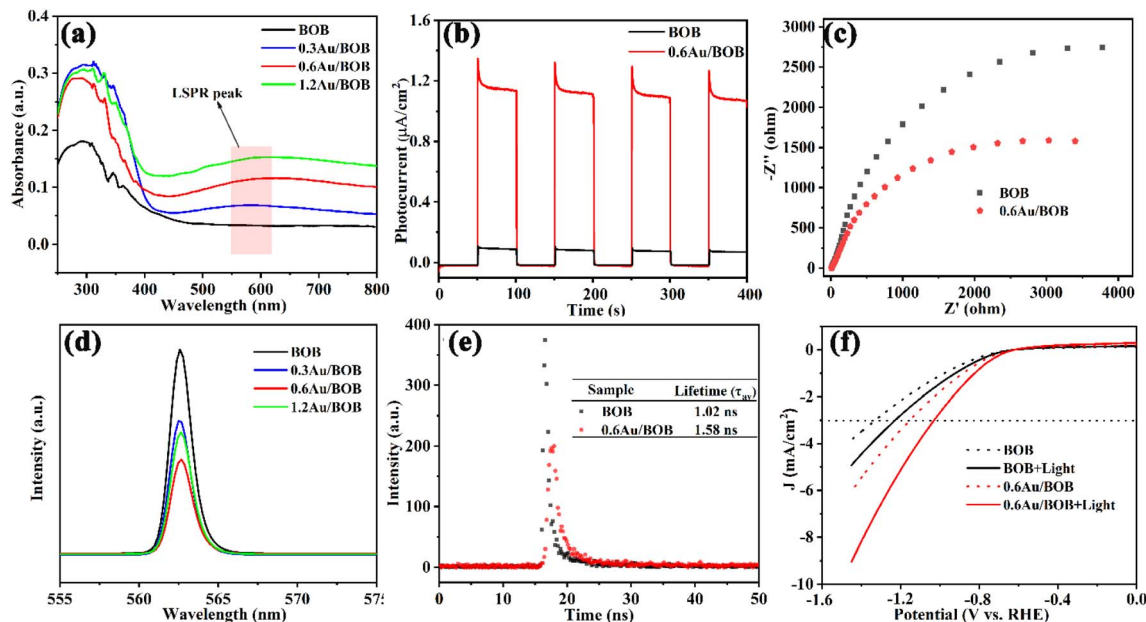


Fig. 2 (a) UV-vis diffuse reflection spectra of samples; (b) transient photocurrent responses ($\lambda > 420$ nm) and (c) EIS of BOB and 0.6Au/BOB; (d) PL spectra; (e) time-resolved PL (TRPL) spectra; (f) LSV of BOB and 0.6Au/BOB under visible light or not.

PL decay (TRPL) spectra of BOB and 0.6Au/BOB were recorded.⁴⁰ As shown in Fig. 2d and e, 0.6Au/BOB has the lowest PL intensity and longer lifetime (1.58 ns, Table S3[†]), proving effective suppression of charge recombination, which is consistent with the results of photocurrent response and EIS. The LSV results in Fig. 2f further proved that to reach the same cathodic current density (e.g. -3 mA cm⁻²), the lowest potential (-0.95 V vs. RHE) was needed for 0.6Au/BOB, while the highest potential (-1.25 V vs. RHE) was required for BOB.⁴¹ Therefore, the composite structure and LRET mechanism can effectively promote the separation and transport of photogenerated charges and inhibit their recombination.

3.3. DFT calculation evidence for the formation of the ohmic contact

To clarify the charge separation and transfer behaviors of the Au/BOB composites, Density Function Theory (DFT) calculations were employed to investigate the work function and spatial charge density difference of the as-prepared materials. As shown in Fig. 3a and b, the crystal structures of BOB and Au/BOB are simulated by DFT, and the work functions (ϕ) of BOB and Au/BOB are 7.04 and 4.99 eV, respectively. Obviously, the Au/BOB composite has a smaller work function, indicating easier overflowing electrons and higher separation efficiency of photogenerated charges. The work function of Au is 5.21 eV,^{42,43} which is much smaller than that of BOB (7.04 eV), implying that a unique ohmic contact will be formed between Au NPs and BOB (as mentioned by Tang *et al.*).⁴⁴ When BOB and Au NPs are in contact (Fig. 3c), the energy band of BOB will bend downward and form an electron accumulation region. Due to electron diffusion, Au NPs will form a hole accumulation area at the

interface. Theoretically, this charge distribution at the interface of the composite will form an internal electric field (IEF), and the direction of the IEF is from Au NPs to BOB. The differential charge distribution of Au/BOB confirms this conclusion. As shown in Fig. S11,[†] the electron consumption occurs on the Au NP side, while charge accumulation occurs on the BOB side.^{45,46} This result also confirms the electron enrichment trend of BOB in the dark, which is consistent with the results of XPS. Different from Schott contact, the downward bending of the band caused by ohm contact can reduce the transfer barrier of BOB conduction band electrons to Au NPs and greatly increases the separation efficiency of photogenerated charge under the action of the IEF.

Generally, the interface IEF takes on the dynamic driving force for the separation and migration of photogenerated charges. And the intensity of the IEF determines the efficiency of charge separation and migration. The IEF intensities of BOB and Au/BOB composites were calculated. Firstly, the intensity of the IEF is directly proportional to the surface potential of the catalyst, and a Kelvin probe force microscopy (KPFM) test was conducted.^{47,48} As shown in Fig. 3d and e, the surface potential of BOB is only 16.2 mV, and while the surface potential of the Au/BOB composite reaches 33.5 mV. Secondly, an experimental test was conducted based on the model proposed by Kanata *et al.*⁴⁹ (Fig. S12[†]). According to this model, the square roots of the products of surface voltage (V_s) and surface charge density (ρ) are proportional to IEF intensities. As shown in Fig. 3f, the IEF intensity of Au/BOB increases by 6.12 times compared to that of BOB. The above results show that the ohmic contact between Au NPs and BOB can effectively improve the IEF intensity and stimulate the separation and migration of photogenerated charges.



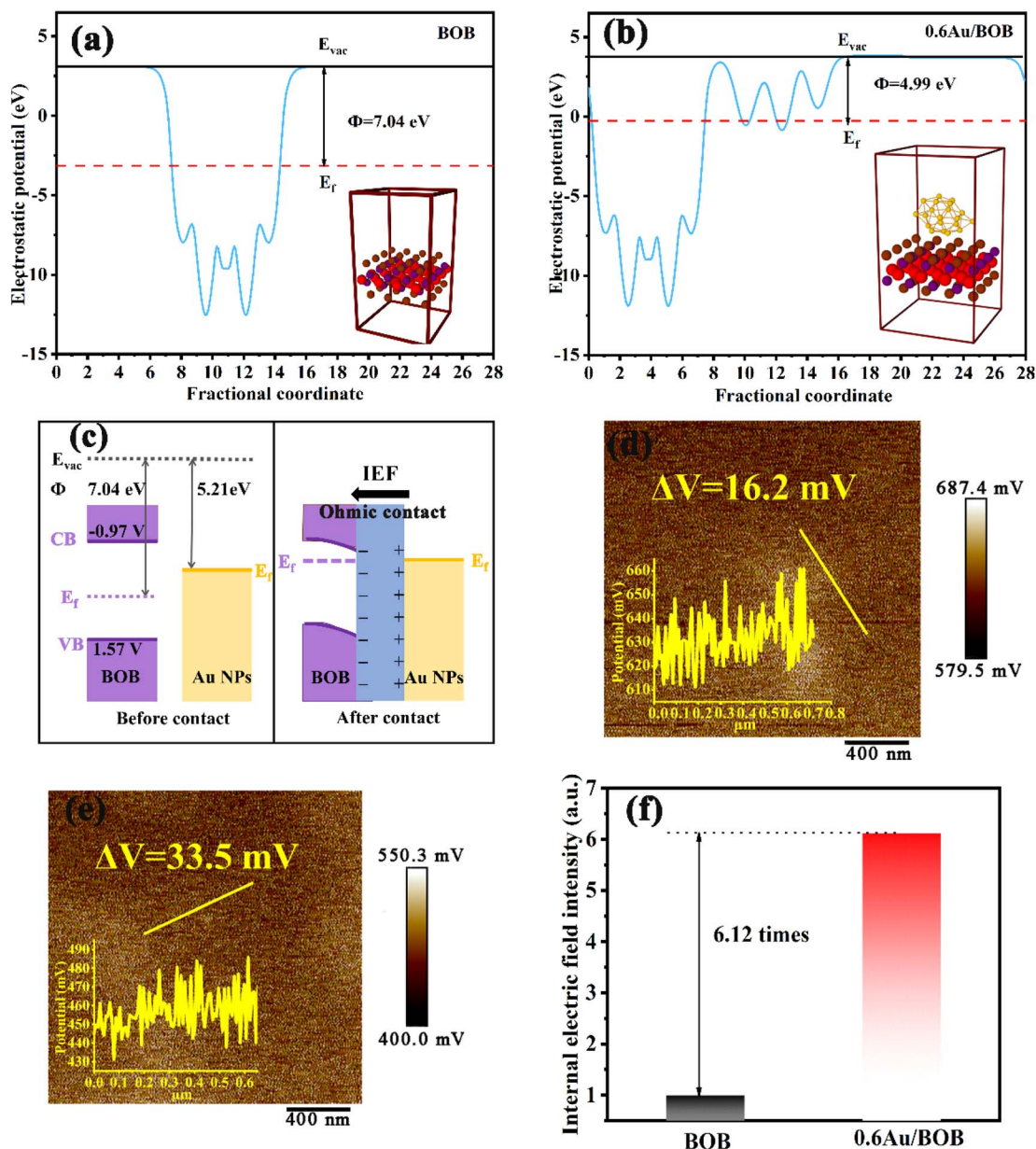


Fig. 3 Work functions of (a) BOB and (b) Au/BOB; (c) schematic diagram of the energy band structure of the ohmic contact; KPFM images of (d) BOB and (e) 0.6Au/BOB; internal electrical field intensities (f) of BOB and 0.6Au/BOB.

3.4. CO₂ photoreduction performance measurements

Photocatalytic performance is evaluated to reveal the advantages of the catalytic system. Here, the photocatalytic performance of BOB and different Au/BOB composites was studied. The photocatalytic reduction of CO₂ with H₂O (no sacrificial agent) was evaluated under simulated sunlight ($\lambda > 420$ nm). CO was detected as the major CO₂ reduction product with a small amount of CH₄ but no other C₂–C₄ hydrocarbon was detected. O₂ was also generated from H₂O oxidation. As shown in Fig. 4a, the CO yield of BOB is only 9.36 $\mu\text{mol g}^{-1} \text{h}^{-1}$. In contrast, the CO yield of 0.6Au/BOB is the highest, reaching 44.92 $\mu\text{mol g}^{-1} \text{h}^{-1}$. With the loading of Au NPs, the changing trend of properties is 0.6Au/BOB (44.92 $\mu\text{mol g}^{-1} \text{h}^{-1}$) > 1.2Au/BOB (34.68 $\mu\text{mol g}^{-1} \text{h}^{-1}$)

> 0.3Au/BOB (27.62 $\mu\text{mol g}^{-1} \text{h}^{-1}$). Obviously, the performance of O₂ also has the same trend, and the performance of 0.6Au/BOB is the highest, which can reach 17.83 $\mu\text{mol g}^{-1} \text{h}^{-1}$. Although the performance of the Au/BOB composite did not increase linearly with the increase of the Au NP content, it was improved. In addition, the performance of 1.2Au/BOB does not increase but decreases compared with that of 0.6Au/BOB, which can be attributed to the fact that superfluous Au NPs will shield the incident light, thus hindering the generation of photo-generated carriers and becoming the recombination center.

The excellent CO yield of 0.6Au/BOB can be attributed to the following aspects. Firstly, the IEF at the interface between Au NPs and BOB greatly improves the migration rate of



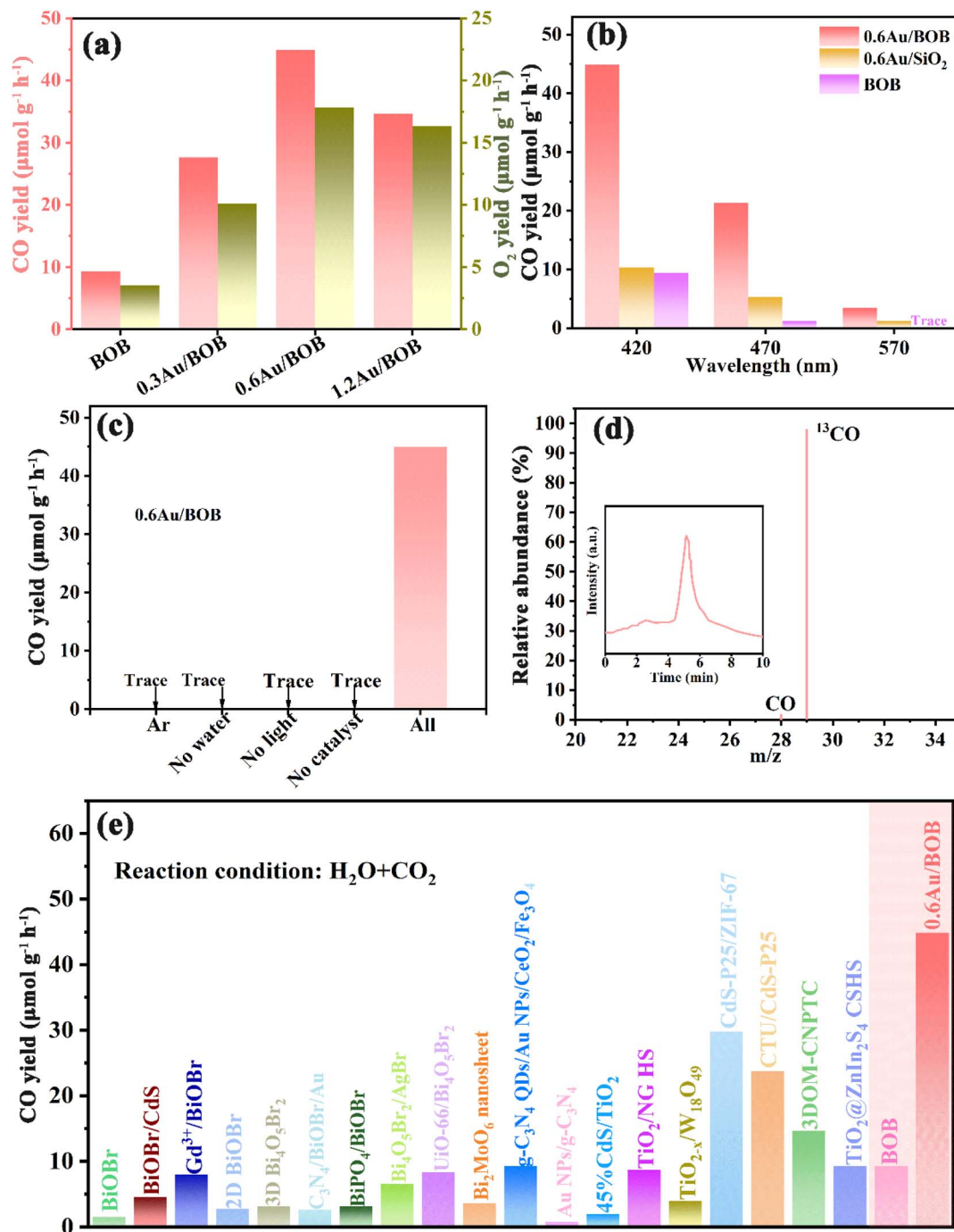


Fig. 4 (a) Photocatalytic CO yield and O₂ yield; (b) CO yield of BOB and 0.6Au/BOB under different wavelengths of light irradiation; (c) CO yield on 0.6Au/BOB under different reaction conditions; (d) mass spectrometry spectra of ¹³CO₂ isotope experiments in the presence of 0.6Au/BOB (inset: gas chromatography spectra); (e) comparison of the CO yield of 0.6Au/BOB and some representative photocatalysts under similar conditions.

photogenerated charges. Secondly, Au NPs formed a strong LSPR effect. Among them, the LEFE mechanism further enhances the charge transfer of the IEF. The LRET mechanism enhances the generation of photogenerated charges. The photocatalytic performance of 0.6Au/BOB composites at different wavelengths was investigated to confirm the positive influence of the LSPR effect of Au NPs on the photocatalytic performance

(Fig. 4b). The best photocatalytic performance was obtained under simulated sunlight ($\lambda > 420$ nm) because of more photon energy excited by its wide spectral range, which was further proved by the significant reduction of photocatalytic CO yield under the irradiation of incident light wavelengths $\lambda > 470$ nm and $\lambda > 570$ nm (ref. 42). Of course, this also proved that the LSPR effect of Au NPs can indeed broaden the spectral response



range of the photocatalyst, and then increase the yield of photogenerated charges. In addition, to verify the synergy effect of the IEF and LSPR effect on the improvement of photocatalytic performance, the activities of BOB and 0.6Au/SiO₂ at different wavelengths were studied. It can be found that when the incident wavelength was greater than 420 nm, the photocatalytic performance of 0.6Au/BOB was 5 times and 6 times that of BOB and 0.6Au/SiO₂, respectively. When the incident wavelength was greater than 570 nm, the photocatalytic CO yields of 0.6Au/BOB and 0.6 Au/SiO₂ are 3.44 $\mu\text{mol g}^{-1} \text{h}^{-1}$ and 0.98 $\mu\text{mol g}^{-1} \text{h}^{-1}$, respectively, and BOB has almost no light response. These data prove that it was the synergy effect of the IEF and LSPR effect that improved the photocatalytic CO yield of the 0.6Au/BOB composite, and both of them were indispensable.

Fig. 4c shows the effect of different experimental conditions on the photocatalytic reduction of CO₂ with 0.6Au/BOB. No CO is detected in the Ar atmosphere. Similarly, CO is not detected under no water, no light or no catalyst conditions, indicating that the integrity of reaction conditions is a prerequisite for the success of an experiment. These results also confirm that the added CO₂ is the main source of CO, which is further confirmed by the isotope labeling experiment of ¹³CO₂ (Fig. 4d). Using ¹³CO₂ as the isotope-labeled raw material, the *m/z* values of ¹³CO and CO were observed to be 29 and 28, respectively, in the

GC-MS dual instrument, and the abundance ratio is close to 49 : 1, which confirms the accuracy of the conjecture.⁵⁰ Furthermore, Table S4† shows the wavelength dependent CO yield of 0.6Au/BOB, except that a bandpass filter is equipped to obtain monochromatic incident light ($\lambda = 380, 420, 470$ and 570 nm). The optical power of the different monochromatic lights is given in Table S4.† Based on this, the apparent quantum yield (AQY) of the 0.6Au/BOB photocatalyst from CO₂ to CO can be determined (see the ESI† for the detailed calculation process). The action spectrum is presented in Fig. S13a.† The UV-visible absorption spectrum and action spectrum of the 0.6Au/BOB photocatalyst were found to be congruent, evincing its commendable light absorption and utilization capacities.⁵¹ In addition, the positive correlation trend between the obtained intensity and AQY at $\lambda = 420$ nm further supports this result (Fig. S13b†).⁵² Considering that the 0.6Au/BOB photocatalyst was prepared by thermal reduction, it is reasonable to suspect that crystal defects were formed in 0.6Au/BOB. As shown in Fig. S14,† unsurprisingly, strong oxygen vacancies did exist in 0.6Au/BOB and positively enhanced the photocatalytic CO₂ reduction performance.⁵³ Moreover, recycling stability is a key factor for the practical application of photocatalysts, in addition to excellent photocatalytic performance. As shown in Fig. S15,† the photocatalytic activity of 0.6Au/BOB only decreased by 6%

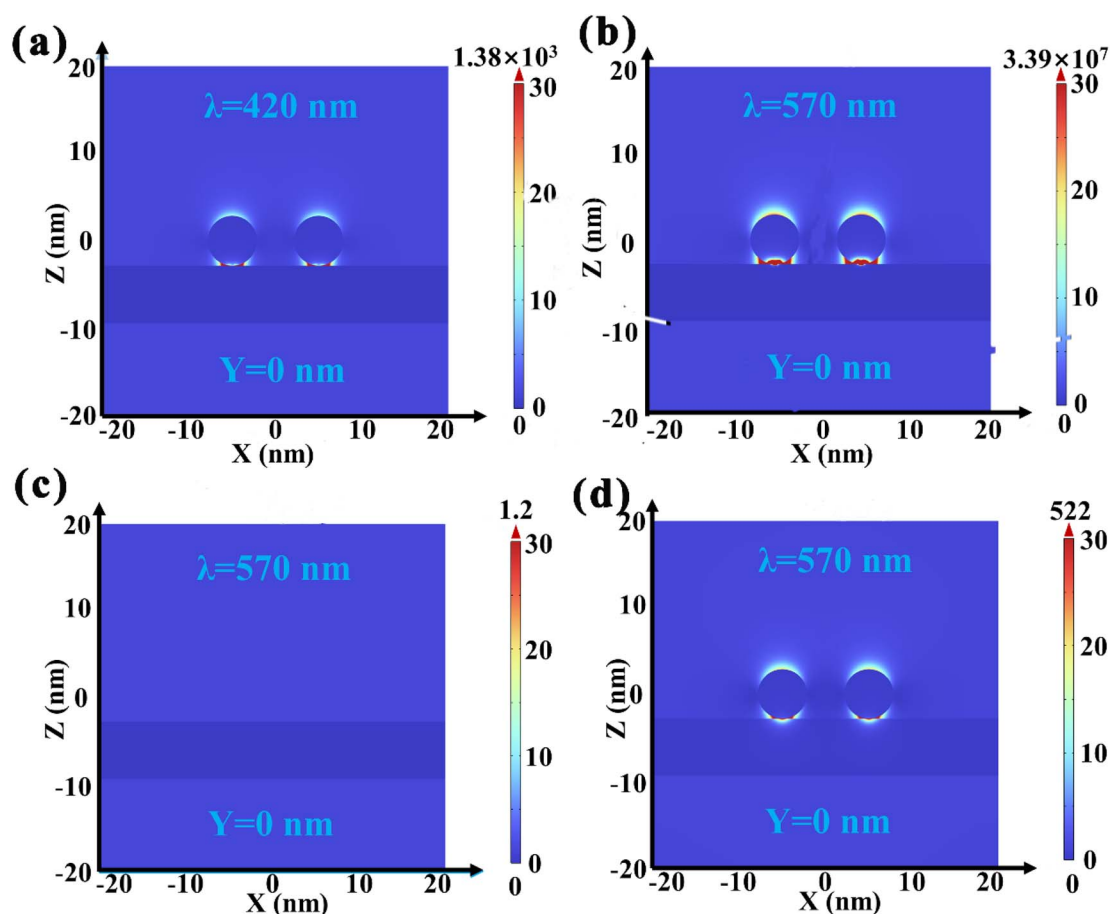


Fig. 5 COMSOL simulated the local electric field distribution of (a) 0.6Au/BOB at 420 nm and (b) 0.6Au/BOB, (c) BOB and (d) 0.6Au/SiO₂ under 570 nm monochromatic light excitation, respectively.



after 48 h of continuous reaction, and the activity showed a stable trend in the later stage of the reaction, indicating that loading Au NPs could not only improve the photocatalytic activity of BOB but also greatly improve its stability. In addition, the XRD, TEM, XPS and FTIR spectra of 0.6Au/BOB remained almost unchanged before and after the reaction, which also proved its excellent structural stability (Fig. S16 and S17[†]). For comparison, Fig. 4e and Table S5[†] summarize some important experiments in recent years, showing that loading Au NPs based on a bismuth-rich strategy has an advantage in photocatalytic CO₂ reduction.

3.5. Local electric field enhancement effect analysis

The inherent mechanism of the effect of Au NP LSPR on the photocatalytic performance of Au/BOB is surveyed by the multi-physical field coupling analysis (COMSOL) simulation method. Under the excitation of polarized 420 nm light along the *Y* axis (Fig. 5a), a hot spot representing the focused electromagnetic field is found at the interface.³⁵ The electromagnetic intensity of the local hot spot is much higher than that of the far-field excited light, and the electromagnetic field intensity of the far-field excited light is at least one order of magnitude higher than that of the far-field excited light. The appearance of a local hot spot means that the photogenerated electrons will migrate from BOB to Au NPs driven by the IEF. As shown in Fig. 5b, with the increase of the incident light wavelength, the electromagnetic intensity corresponding to the local hot spot increases by four orders of magnitude, which means that the local electric field enhancement effect (LEFE) can promote the migration of photogenerated charge.^{54,55} At the same time, the enhancement of the local electric field of 0.6Au/BOB can further explain the synergy effect of the LEFE and interfacial ohmic contact on the effective separation and migration of photogenerated charges, which can be confirmed by the lower local electric field of Au NPs and BOB (Fig. 5c and d).

3.6. Photocatalytic mechanism

Based on the above analysis, we proposed a feasible mechanism of Au/BOB to enhance photocatalytic CO₂ reduction (Fig. 6). Due to the different work functions (Φ), the contact between BOB and Au NPs leads to a charge redistribution with the BOB side negatively charged and the Au NPs positively charged. This charge distribution caused the energy band near the BOB side to bend downward, forming an internal electric field (IEF) from Au NPs to BOB. An ohmic contact was formed at the heterogeneous interface to realize the directional migration of carriers. Under the irradiation of visible light, the photogenerated electrons and photogenerated holes were separated. Driven by the IEF, the photogenerated electrons transferred to the BOB CB were quickly guided and migrated to the Au NP surface in one direction. At the same time, the LSPR effect of Au NPs promotes this process. On the one hand, the local resonance energy transfer (LRET) effect was induced by the LSPR effect, which broadened the spectral response range of BOB and increased the photocharge yield efficiency. On the other hand, the LSPR effect induced the local electric field enhancement effect (LEFE), which led to a strong local electric field on the Au NPs to further accelerate the transfer of photogenerated charge. Under the double synergistic action of the interface ohmic contact and LSPR effect, more photogenerated electrons migrated to the Au NP surface along the bending direction of the BOB energy band and participated in the reduction reaction of CO₂ to CO. Surely, the significantly increased CO₂ adsorption capacity and the appearance of double CO₂-TPD desorption peaks in the 0.6Au/BOB composite further proved that the synergistic effect promoted the adsorption and activation of CO₂ (Fig. S18[†]), which was consistent with its performance results. In addition, we also characterized the process of photocatalytic reduction of CO₂ to CO by *in situ* FTIR and Gibbs free energy, and the specific reaction process is shown in Fig. S19.[†] Noticeably, due to no sacrificial agent in the reaction system, the holes left by the BOB VB would react directly with H₂O to form O₂, as shown in Fig. 6.

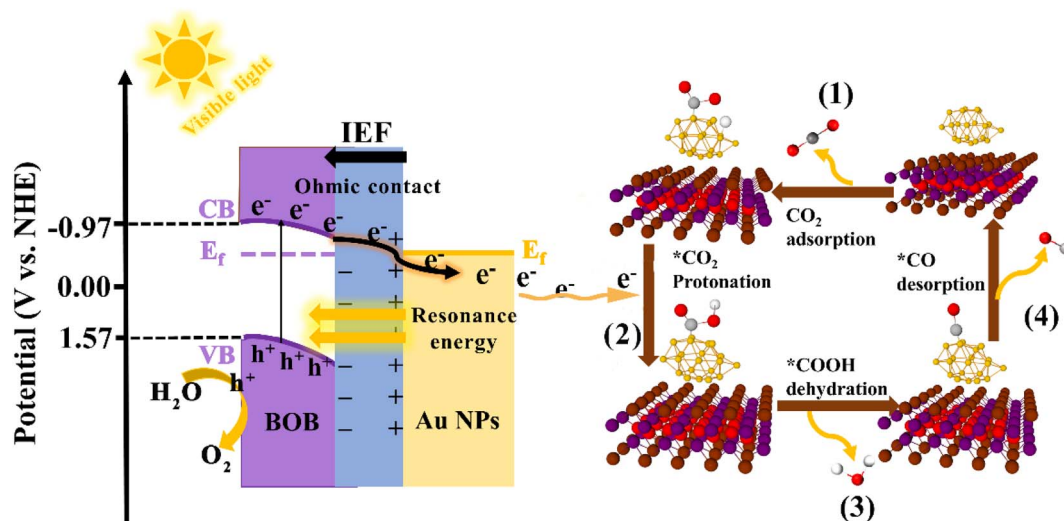


Fig. 6 The photocatalytic mechanism of Au/BOB and the detailed steps of reducing CO₂ to CO.



4. Conclusion

In summary, Au NPs with a dual function emphasized in the Au/Bi₂₄O₃₁Br₁₀ (Au/BOB) composite, and the ultrafast transfer of photogenerated charges from BOB to Au NPs was realized through the synergistic effect of the interfacial ohmic contact and LSPR effect. Here, the introduction of Au NPs could not only collect more electrons for CO₂ reduction as an electron trap, but also induced the local electric field enhancement effect (LEFE) to enhance the driving force of the internal electric field (IEF) for photogenerated charge transfer. In addition, the ultrahigh CO yield and selectivity were significantly attributed to the reduced energy barrier and strong CO₂ adsorption and activation. Under visible light irradiation, the optimal 0.6Au/BOB exhibited nearly 5 times higher activity than BOB (9.36 μmol g⁻¹ h⁻¹) without any sacrificial reagents. More importantly, 0.6Au/BOB demonstrated superior stability (only lost 6%) after continuous reaction for 48 h without obvious deactivation.

Data availability

The data is available in ESI.†

Author contributions

Jie Liu: conceptualization, methodology, software, investigation; Yu Xie: supervision, project administration, funding acquisition; Yiqiao Wang: writing – review & editing; Kai Yang, Shuping Su, Yun Ling and Pinghua Chen: writing – review & editing, investigation.

Conflicts of interest

The authors declare no competing interests.

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References

- 1 K. Maeda, *Adv. Mater.*, 2019, **31**, e1808205.
- 2 L. Liu, H. Huang, Z. Chen, H. Yu, K. Wang, J. Huang, H. Yu and Y. Zhang, *Angew. Chem., Int. Ed.*, 2021, **60**, 18303–18308.
- 3 D. Zhang, Y. Guo and Z. Zhao, *Appl. Catal., B*, 2018, **226**, 1–9.
- 4 C. Y. Wang, X. Zhang, H. Bin Qiu, G. X. Huang and H. Q. Yu, *Appl. Catal., B*, 2017, **205**, 615–623.
- 5 J. Wan, Y. Zhang, R. Wang, L. Liu, E. Liu, J. Fan and F. Fu, *J. Hazard. Mater.*, 2018, **384**, 121484.
- 6 K. Li, B. Peng and T. Peng, *ACS Catal.*, 2016, **6**, 7485–7527.
- 7 D. Cui, L. Wang, Y. Du, W. C. Hao and J. Chen, *ACS Sustainable Chem. Eng.*, 2018, **6**, 15936–15953.
- 8 M. Ji, Y. Liu, J. Di, R. Chen, Z. Chen, J. Xia and H. Li, *Appl. Catal., B*, 2018, **237**, 1033–1043.
- 9 X. Chen, M. Y. Qi, Y. H. Li, Z. R. Tang and Y. J. Xu, *Chin. J. Catal.*, 2021, **42**, 2020–2026.
- 10 X. Jin, C. Lv, X. Zhou, C. Zhang, B. Zhang, H. Su and G. Chen, *J. Mater. Chem. A*, 2018, **6**, 24350–24357.
- 11 X. Xiao, C. Zheng, M. Lu, L. Zhang, F. Liu, X. Zuo and J. Nan, *Appl. Catal., B*, 2018, **228**, 142–151.
- 12 R. L. Tang, C. L. Hu, B. L. Wu, Z. Fang, Y. Chen and J. G. Mao, *Angew. Chem., Int. Ed.*, 2019, **58**, 15358–15361.
- 13 Y. Liu, C. Miao, P. Yang, Y. He, J. Feng and D. Li, *Appl. Catal., B*, 2019, **244**, 919–930.
- 14 Y. Dai, C. Li, Y. Shen, S. Zhu, M. S. Hvid, L. C. Wu, J. Skibsted, Y. Li, J. W. H. Niemantsverdriet, F. Besenbacher, N. Lock and R. Su, *J. Am. Chem. Soc.*, 2018, **140**, 16711–16719.
- 15 C. Zhao, Z. Wang, X. Li, X. Yi, H. Chu, X. Chen and C. C. Wang, *Chem. Eng. J.*, 2020, **389**, 123431.
- 16 R. Verma, R. Belgamwar and V. Polshettiwar, *ACS Mater. Lett.*, 2021, **3**, 574–598.
- 17 X. Li, H. Jiang, C. Ma, Z. Zhu, X. Song, H. Wang, P. Huo and X. Li, *Appl. Catal., B*, 2021, **283**, 119638.
- 18 P. Zhang, L. Jun Wu, W. Guo Pan, S. Cheng Bai and R. Tang Guo, *Appl. Catal., B*, 2021, **289**, 120040.
- 19 L. Y. Lin, S. Kavadiya, X. He, W. N. Wang, B. B. Karakocak, Y. C. Lin, M. Y. Berezin and P. Biswas, *Chem. Eng. J.*, 2020, **389**, 123450.
- 20 W. Hou and S. B. Cronin, *Adv. Funct. Mater.*, 2013, **23**, 1612–1619.
- 21 L. Ma, Y. L. Chen, D. J. Yang, H. X. Li, S. J. Ding, L. Xiong, P. L. Qin and X. B. Chen, *Nanoscale*, 2020, **12**, 4383–4392.
- 22 X. Wang, R. Long, D. Liu, D. Yang, C. Wang and Y. Xiong, *Nano Energy*, 2016, **24**, 87–93.
- 23 C. Cheng, J. Zhang, R. Zeng, F. Xing and C. Huang, *Appl. Catal., B*, 2022, **310**, 121321.
- 24 H. Ren, T. Dittrich, H. Ma, J. N. Hart, S. Fengler, S. Chen, Y. Li, Y. Wang, F. Cao, M. Schieda, Y. H. Ng, Z. Xie, X. Bo, P. Koshy, L. R. Sheppard, C. Zhao and C. C. Sorrell, *Adv. Mater.*, 2019, **31**, 1–9.
- 25 C. Feng, L. Tang, Y. Deng, J. Wang, Y. Liu, X. Ouyang, Z. Chen, H. Yang, J. Yu and J. Wang, *Appl. Catal., B*, 2020, **276**, 119167.
- 26 P. Mazierski, W. Lisowski, T. Grzyb, M. J. Winiarski, T. Klimczuk, A. Mikołajczyk, J. Flisikowski, A. Hirsch, A. Kolakowska, T. Puzyn, A. Zaleska-Medynska and J. Nadolna, *Appl. Catal., B*, 2017, **205**, 376–385.
- 27 P. E. Blöchl, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1994, **50**, 17953–17979.
- 28 J. P. Perdew, K. Burke and M. Ernzerhof, *Phys. Rev. Lett.*, 1996, **77**, 3865–3868.
- 29 A. S. Rosen, V. Fung, P. Huck, C. T. O'Donnell, M. K. Horton, D. G. Truhlar, K. A. Persson, J. M. Notestein and R. Q. Snurr, *npj Comput. Mater.*, 2022, **8**, 1–10.
- 30 K. Hu, M. Wu, S. Hinokuma, T. Ohto, M. Wakisaka, J. I. Fujita and Y. Ito, *J. Mater. Chem. A*, 2019, **7**, 2156–2164.



- 31 J. Haslmayr and T. Renger, *J. Chem. Phys.*, 2013, **139**, 044103.
- 32 Z. W. Wang, Q. Wan, Y. Z. Shi, H. Wang, Y. Y. Kang, S. Y. Zhu, S. Lin and L. Wu, *Appl. Catal., B*, 2021, **288**, 120000.
- 33 S. G. Fard, M. Haghghi and M. Shabani, *Appl. Catal., B*, 2019, **248**, 320–331.
- 34 Y. Zhou, C. Zhang, D. Huang, W. Wang, Y. Zhai, Q. Liang, Y. Yang, S. Tian, H. Luo and D. Qin, *Appl. Catal., B*, 2022, **301**, 120749.
- 35 G. Yu, J. Qian, P. Zhang, B. Zhang, W. Zhang, W. Yan and G. Liu, *Nat. Commun.*, 2019, **10**, 1–8.
- 36 J. Guo, X. Liao, M. H. Lee, G. Hyett, C. C. Huang, D. W. Hewak, S. Mailis, W. Zhou and Z. Jiang, *Appl. Catal., B*, 2019, **243**, 502–512.
- 37 Y. Tong, L. Chen, S. Ning, N. Tong, Z. Zhang, H. Lin, F. Li and X. Wang, *Appl. Catal., B*, 2017, **203**, 725–730.
- 38 X. Cai, L. Mao, S. Yang, K. Han and J. Zhang, *ACS Energy Lett.*, 2018, **3**, 932–939.
- 39 R. Verma, R. Belgamwar, P. Chatterjee, R. Bericat-Vadell, J. Sa and V. Polshettiwar, *ACS Nano*, 2023, **17**, 4526–4538.
- 40 Y. B. Kim, J. S. Kim, S. K. Baek, Y. D. Yun, S. W. Cho, S. H. Jung and H. K. Cho, *Energy Environ. Sci.*, 2018, **11**, 2540–2549.
- 41 Y. Yoon, B. Yan and Y. Surendranath, *J. Am. Chem. Soc.*, 2018, **140**, 2397–2400.
- 42 H. Zhu, X. Yuan, Q. Yao and J. Xie, *Nano Energy*, 2021, **88**, 106306.
- 43 T. Han, X. Cao, K. Sun, Q. Peng, C. Ye, A. Huang, W. C. Cheong, Z. Chen, R. Lin, D. Zhao, X. Tan, Z. Zhuang, C. Chen, D. Wang and Y. Li, *Nat. Commun.*, 2021, **12**, 1–11.
- 44 M. Tang, Y. Ao, C. Wang and P. Wang, *Appl. Catal., B*, 2020, **270**, 118918.
- 45 Z. Miao, Q. Wang, Y. Zhang, L. Meng and X. Wang, *Appl. Catal., B*, 2022, **301**, 120802.
- 46 Z. Wu, J. Jing, K. Zhang, W. Li, J. Yang, J. Shen, S. Zhang, K. Xu, S. Zhang and Y. Zhu, *Appl. Catal., B*, 2022, **307**, 121153.
- 47 J. Xu, Z. Ju, W. Zhang, Y. Pan, J. Zhu, J. Mao, X. Zheng, H. Fu, M. Yuan, H. Chen and R. Li, *Angew. Chem., Int. Ed.*, 2021, **60**, 8705–8709.
- 48 J. Li, L. Cai, J. Shang, Y. Yu and L. Zhang, *Adv. Mater.*, 2016, **28**, 4059–4064.
- 49 T. Kanata, M. Matsunaga, H. Takakura, Y. Hamakawa and T. Nishino, *J. Appl. Phys.*, 1990, **68**, 5309–5313.
- 50 Z. Yu, K. Yang, C. Yu, K. Lu, W. Huang, L. Xu, L. Zou, S. Wang, Z. Chen, J. Hu, Y. Hou and Y. Zhu, *Adv. Funct. Mater.*, 2022, 1–10.
- 51 X. Wang, X. Wang, J. Huang, S. Li, A. Meng and Z. Li, *Nat. Commun.*, 2021, **12**, 1–11.
- 52 S. Singh, R. Verma, N. Kaul, J. Sa, A. Punjal, S. Prabhu and V. Polshettiwar, *Nat. Commun.*, 2023, **14**, 2551.
- 53 Z. Miao, Q. Wang, Y. Zhang, L. Meng and X. Wang, *Appl. Catal., B*, 2022, **301**, 120802.
- 54 M. Liu, Q. Kang, Z. Xie, L. Lu, K. Dai and G. Dawson, *J. Phys. D Appl. Phys.*, 2022, **55**, 043002.
- 55 M. Fang, X. Tan, Z. Liu, B. Hu and X. Wang, *Research*, 2021, **2021**, 9794329.

