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# Optimizing oxygen vacancies and electrochemical performance of $\text{CeO}_{2-\delta}$ nanosheets through the combination of di- and tri-valent doping

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Ceramic fuel cells presently hold an important position in the future of sustainable energy. However, new concepts and designs are vital for each individual cell's component materials to improve the overall power output and stability. The limited ionic conductivity of the electrolyte component is one major challenge among these. In the present work, we developed nanosheets with a cubic fluoride structure of  $\text{CeO}_2$  and introduced the di- and tri-valent doping of La and Sr to study their impact on oxygen vacancies and its ionic transport, keeping in mind the fact that  $\text{CeO}_2$  is reduced when exposed to a reducing atmosphere. The attained La- and Sr-doped fluorite structures of  $\text{CeO}_2$  exhibited good ionic conductivity of  $>0.05 \text{ S cm}^{-1}$  at low temperature, and their use in a fuel cell resulted in achieving a power output of  $>900 \text{ mW cm}^{-2}$  while operating at  $550^\circ\text{C}$ . Therefore, we have found that laterally combining di- and tri-valent doping could be textured to give a highly oxygen-deficient  $\text{CeO}_2$  structure with high ionic transport. Furthermore, various microscopic and spectroscopic analyses, such as HR-TEM, XPS, Raman, UV-visible, EIS, and density functional theory, were applied to investigate the change in structural properties and mechanism of the ionic transport of the synthesized La and Sr co-doped  $\text{CeO}_2$  electrolyte. This work provides some new insights for designing high-ionic-conductivity electrolytes from low-cost semiconductor oxides for energy storage and conversion devices.

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## 1 Introduction

The term solid oxide fuel cells (SOFCs) was coined for the feasible and attractive devices for converting chemical energy directly into electrical energy with negligible pollution. This is one of the compatible and competitive candidates for the efficient production of sustainable energy to meet the present and ever-increasing high energy demand. However, there is a problem associated with the required high ionic conduction and low-temperature operation, which has limited the

advancement and full commercialization of SOFCs.<sup>1</sup> Therefore, researchers are devoting great efforts to overcome this problem to enhance ionic conduction in the electrolyte and pursuing options for low-temperature operation.<sup>1–3</sup> In this regard, various technical efforts have been made in terms of micro-thin film technology to reduce the size of the induced electrolyte layer to overcome the associated problems,<sup>4–6</sup> which could help to overcome polarization losses under a reduced operating temperature; however, micro-thin film technology involves an expensive and complex method.<sup>7</sup> Therefore, Goodenough *et al.* proposed investigating and designing new oxide-ion conductors as alternative materials for use with the possible low temperatures and to be technically useful.<sup>8</sup> However, despite the great efforts of many researchers to design and investigate new oxide materials for high ionic conductivity, still the materials are constrained below the level of  $0.1 \text{ S cm}^{-1}$  ionic conductivity.<sup>9</sup>

The addition of ceria or ceria-based materials has drawn significant attention due to their highly active and demonstrated multifunctionality in many fields, such as catalytic and electronic applications, solar cells and photoelectrochemistry, lithium batteries, sensors and biosensor, fuel cells, and a variety of other energy-related applications.<sup>10–21</sup> One of the attractive characteristics of ceria is the easily changeable oxidation states in the  $\text{Ce}^{4+}/\text{Ce}^{3+}$  redox cycle that can be used to release and store oxygen, which would be entirely dependent on the

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concentration and type of oxygen vacancies in the lattice, as well as the surface structures and states. There are unique physical properties linked with  $\text{Ce}^{3+}$  ions and oxygen vacancies, but issues also exist, especially for doped ceria oxides, which concern the redox instability induced by the reduction from  $\text{Ce}^{4+}$  to  $\text{Ce}^{3+}$ , resulting in electronic conduction and an undesirable structural change.<sup>22,23</sup> However, ceria is considered a competent candidate as an electrolyte for SOFCs and has been reported in heterostructures as well as in dopants. In this regard, the use of  $7\text{CeO}_2\text{-}3\text{CuO}$  composite as an electrolyte was reported, where Ce played a key role in achieving a high ionic conductivity of  $0.137\text{ S cm}^{-1}$  at  $550\text{ }^\circ\text{C}$ .<sup>24</sup> Furthermore, Qiao Z. *et al.* reported doping ceria with ZnO in a composite weight ratio of  $7\text{LCP-}3\text{ZnO}$ , where the primary part was doped ceria and contributed significantly to the increase in the ionic conductivity.<sup>25</sup>

Moreover, Wachsmann and Lee reported co-doped bismuth-oxide-based materials and demonstrated the known highest conductivity as a solid-state oxide electrolyte for SOFCs.<sup>26</sup> Furthermore, the ceria was further tuned by introducing  $\delta\text{-Bi}_2\text{O}_3$  via the formation of coherent and high interfaces of  $\text{Er}_2\text{O}_3$ -stabilized  $\delta\text{-Bi}_2\text{O}_3$  and  $\text{Gd}_2\text{O}_3$ -doped  $\text{CeO}_2$ .<sup>27</sup> Interestingly, the designed layered structured materials with highly coherent interfaces demonstrated exceptionally high chemical stability even in a reducing environment and repeated redox cycles at elevated temperature. In addition, the introduction of thulium doping into  $\text{SrCeO}_2$  led to a significantly enhanced ionic conductivity  $0.13\text{ S cm}^{-1}$  at  $550\text{ }^\circ\text{C}$  resulting from the bulk doping; however, there remains a need to modify ceria with new dopants to achieve stable and high ionic conductivity at low operating temperature with high power density.

Inspired from the above discussion, we introduced dopants of La and Sr into ceria to form  $\text{La}_{0.15}\text{Sr}_{0.05}\text{Ce}_{0.80}\text{O}_2$  (LSC) and applied this as an electrolyte in CFC for the enhancement of ionic conductivity and power density. The as-prepared materials were physically characterized to analyze the phase pattern, microstructure and surface morphology, chemical and oxidation states, weight loss, and formation of oxygen vacancies via X-ray diffraction, high-resolution transmission electron microscopy (HR-TEM), field emission scanning electron microscopy (FE-SEM), XPS, and TGA. In addition, the as-prepared  $\text{LaSrCeO}_{2-\delta}$  was used as an electrolyte membrane, which generated a high content of oxygen vacancies, resulting in an enhanced ionic conductivity, and demonstrated high electrochemical performance compared to pure  $\text{CeO}_2$ , respectively.

## 2 Experimental methods

### 2.1 Materials and methods

**2.1.1 Preparation of  $\text{CeO}_2$ .** First,  $0.05\text{ mol Ce}(\text{NO}_3)_3\cdot 6\text{H}_2\text{O}$  was dissolved in  $100\text{ ml}$  deionized water to prepare  $0.5\text{ mol L}^{-1}$  solution, which was drop by drop added to  $0.5\text{ mol L}^{-1}$  of  $\text{Na}_2\text{CO}_3$  with stirring until a milky precipitate was formed. The precipitates were shifted into an autoclave for hydrothermal processing at  $170\text{ }^\circ\text{C}$  for  $6\text{ h}$  in a vacuum oven. The precipitate was then washed several times with distilled water and filtered, in addition to washing several times with ethanol to remove the

water molecules absorbed at the surface. The resultant precursor was dried in an oven at  $120\text{ }^\circ\text{C}$ , followed by calcination at  $700\text{ }^\circ\text{C}$  for  $1\text{ h}$  in air to obtain  $\text{CeO}_2$  nanoparticles.

**2.1.2 Preparation of  $\text{La}_{0.15}\text{Sr}_{0.05}\text{Ce}_{0.80}\text{O}_2$ .** An appropriate amount of  $\text{La}(\text{NO}_3)_3\cdot 6\text{H}_2\text{O}$  and  $\text{Sr}(\text{NO}_3)_3\cdot 6\text{H}_2\text{O}$  were dissolved in  $200\text{ ml}$  of  $\text{Ce}(\text{NO}_3)_3\cdot 6\text{H}_2\text{O}$ -deionized water solution. Another  $0.1\text{ mole}$  of  $\text{Na}_2\text{CO}_3$  was dissolved in  $100\text{ ml}$  deionized water, and then drop by drop added to  $0.5\text{ mol L}^{-1}$  solution of LSC with stirring until a milky precipitate was formed. Afterward, the precipitates were shifted into an autoclave for hydrothermal processing at  $170\text{ }^\circ\text{C}$  for  $6\text{ h}$  in a vacuum oven. The resulting product was collected by centrifugation and washed with deionized water several times. After undergoing a one-step direct pyrolysis process at  $700\text{ }^\circ\text{C}$  for  $4\text{ h}$ , the corresponding LSC nanosheets were obtained.

### 2.2 Material characterizations

The phase structures of the above-synthesized samples were determined using a Bruker D8 advanced X-ray diffractometer (Germany, Bruker Corporation) with  $\text{Cu K}\alpha$  radiation ( $\lambda = 1.5418\text{ \AA}$ ). The scan range of  $2\theta$  was set from  $20^\circ$ – $80^\circ$  with a scanning step of  $0.02^\circ$ . X-Ray diffractometer refinements were done using MJAD 6.5 software. High-resolution transmission electron microscopy was performed using the FEI Tecnai G1 F30 system along with energy-dispersive spectroscopy and element mapping of the LSC samples. The surface morphological features of the synthesized samples were studied by FESEM (JEOL JSM7100F). A Physical Electronics Quantum 2000 system with a source of  $\text{Al K}\alpha$  X-rays was hired for the surface and chemical analysis. CASA XPS software was used to analyze the XPS data. An NT-MDT Raman spectrometer at  $532\text{ nm}$  and with a  $20\text{ mW}$  solid-state laser source was used to measure the Raman spectra of pure  $\text{CeO}_2$  and gradiently doped  $\text{CeO}_2$ . Kelvin probe force microscopy (KPFM, Bruker Multimode 8) was employed to detect the surface potential height images for the LSC samples with the tapping mode in comparison with  $\text{CeO}_2$ .

### 2.3 Cell fabrication and electrochemical measurements

The preparation of dense electrolyte membranes on porous electrodes is an important step for the fabrication of high-performance CFCs. In our study, we used a simple dry-pressing method to prepare a ceria-based and doped  $\text{CeO}_2$  electrolyte on a  $\text{Ni}_{0.8}\text{Co}_{0.15}\text{Al}_{0.05}\text{LiO}_{2-\delta}$  symmetrical electrode. In details, a  $\text{Ni}_{0.8}\text{Co}_{0.15}\text{Al}_{0.05}\text{LiO}_{2-\delta}$  (NCAL) solid powder was obtained from Bamo Sci. & Tech. Joint Stock Ltd, China. A slurry of NCAL was made by mixing it with terpinol + ethanol ( $90\% + 10\%$  volume ratio). The obtained slurry of NCAL was fixed on porous Ni-foam of  $2.5\text{ mm}$  thickness and laterally followed by drying at  $130\text{ }^\circ\text{C}$  for  $4\text{ h}$  to remove organic compounds to serve as the electrode. However, the high-density filled  $\text{CeO}_2$  and LSC electrolyte powders were sandwiched in-between the Ni-foam painted-NCAL electrode and simultaneously pressed at  $200\text{ MPa}$  to obtain complete tri-layer fuel cells. The thickness of the fuel cells after pressing was about  $150\text{ }\mu\text{m}$  with an active-area of  $0.64\text{ cm}^2$ , respectively. An ITECH8511 system (ITECH Electrical Co., Ltd) was used to measure the electrochemical



performance of the fabricated fuel cell using  $\text{H}_2$  fuel and air with a flow rate of  $120\text{--}130\text{ mL min}^{-1}$  and  $100\text{ mL min}^{-1}$ , respectively. The current density–voltage ( $I$ – $V$ ) and current density–power density ( $I$ – $P$ ) were recorded for presentation of the  $I$ – $V$  and  $I$ – $P$  curves. A Gamry ref. 3000 system was used to measure the electrochemical impedance spectroscopy in open-circuit voltage (OCV) conditions. The impedance spectra were measured in the frequency range of  $10^6$  to  $0.01\text{ Hz}$  along with a  $10\text{ mV}$  dc signal. ZSIMPWIN 3.3 software was used to model the EIS data.

#### 2.4 Density functional theory (DFT) calculations

DFT calculations of the La and Sr co-doped  $\text{CeO}_2$  structure were performed with the Vienna *ab initio* simulation package with the projector augmented wave (PAW) method for the undoped and doped  $\text{CeO}_2$ . The ion exchange–correlation with electrons was described and controlled using the GGA +  $U$  (generalized gradient approximation) and PBE (Perdew–Burke–Ernzerhof) method with a plane wave cut-off energy of  $520\text{ eV}$ . The values of  $U_{\text{eff}} = U$  were set to  $6.30$  for Ce 4-f, respectively. A  $2 \times 2 \times 2$   $k$ -point grid Brillouin zone of each structure was used for the

calculations. The geometry of each supercell was relaxed until the forces at each atom became smaller than  $0.01\text{ eV \AA}^{-1}$  and the energy change was less than  $1 \times 10^{-4}\text{ eV}$ . The reaction pathways for the protons were accurately examined using the climbing-image nudged elastic band method (CI-NEB).<sup>28</sup> The vacancy formation energy ( $E_{\text{V}_\text{O}}$ ) was calculated as,

$$E_{\text{V}_\text{O}} = E_{\text{tot}}(v_\text{O}^q) - E_{\text{tot}}(\text{ideal}) + \mu + q(E_\text{F} + E_{\text{valan}} + \Delta V)$$

where  $E_{\text{V}_\text{O}}$  is the vacancy formation energy,  $E_{\text{tot}}(v_\text{O}^q)$  is the superlattice's total energy in the  $q$  charge state,  $E_{\text{tot}}(\text{ideal})$  is total energy of the ideal structure,  $\mu$  is the chemical potential,  $E_\text{F}$  is the orientation of the Fermi level,  $E_{\text{valan}}$  refers to the valence band maximum, and  $\Delta V$  is the change in the electrostatic potential with the defective lattice.

### 3 Results and discussion

#### 3.1 Structural and compositional study

To prove the incorporation of interstitial La and Sr-doping into the  $\text{CeO}_2$  crystal structure, various characterization techniques

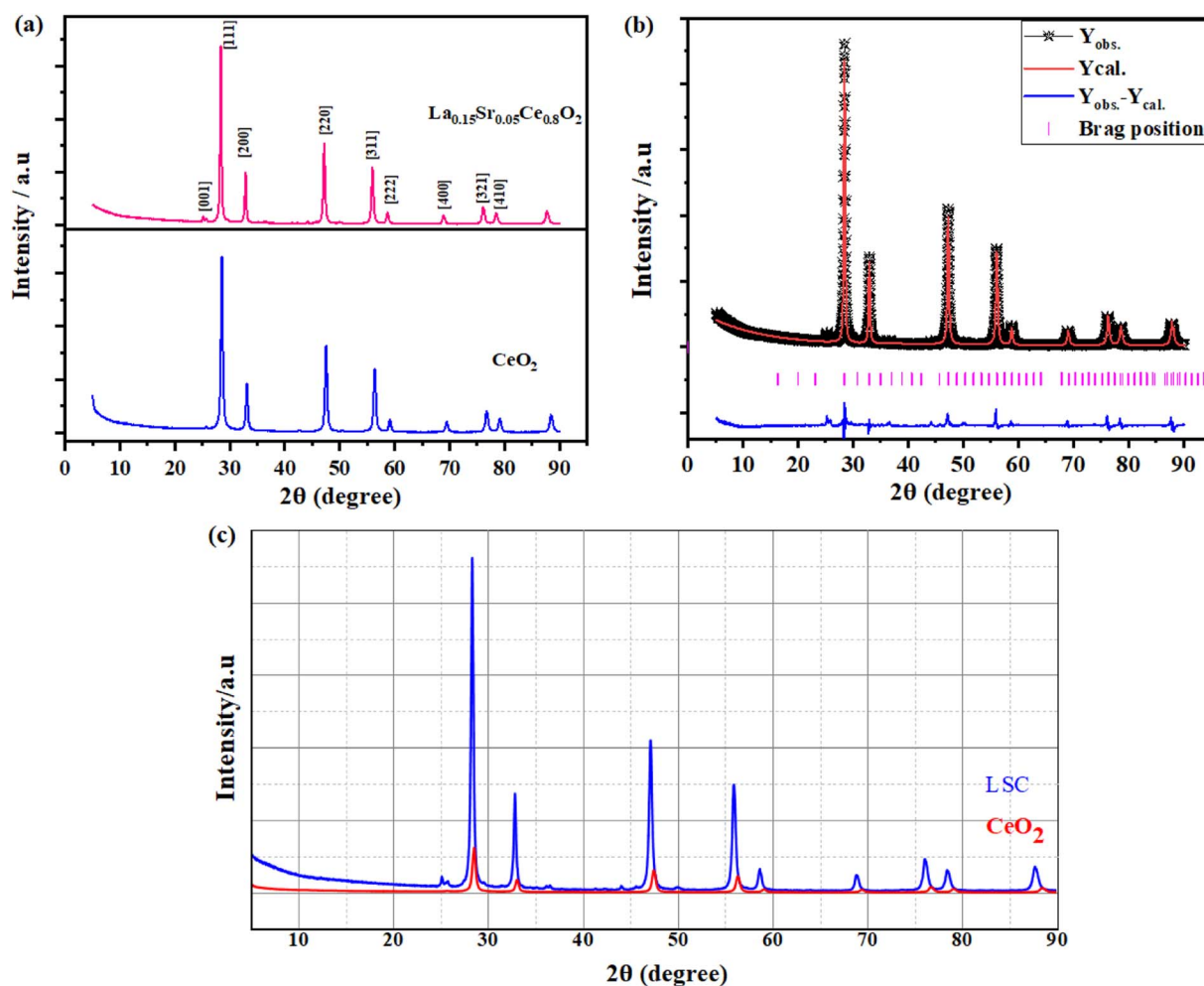


Fig. 1 (a and b) Comparison of the measured XRD data of pristine  $\text{CeO}_2$  and La, Sr-doped  $\text{CeO}_2$ , (b) Rietveld refinement using the Prof-Suit software for the individual La, Sr-doped  $\text{CeO}_2$  and (c) comparison of the peak shifting of the doped and undoped  $\text{CeO}_2$  XRD patterns, respectively.



were employed, and a detailed comparison between the doped and undoped materials is discussed below. The XRD patterns of  $\text{CeO}_2$ , and La-Sr co-doped  $\text{CeO}_2$  (LSC) are shown in Fig. 1a, where it can be observed that the pure  $\text{CeO}_2$  crystal and LSC structure had identical XRD patterns without any secondary phase with the  $Fm\bar{3}221$  space group.<sup>28</sup> However, Rietveld refinement of the XRD pattern for pristine La-Sr co-doped  $\text{CeO}_2$  (LSC) materials was done with the FullProf Suite software. Fig. 1b shows the Rietveld refinement pattern of LSC, where the calculated patterns are shown by solid red curves, and the black empty bullets show the observed intensities, while the differences between the observed and calculated intensities are presented by solid blue curves with a violet bar space group. Apparently, the diffraction peaks of the as-prepared target could be indexed as a typical fluoride structure with the space group  $Fm\bar{3}$ , except for the increase in the lattice constant, which can clearly be seen by the apparent shift of the XRD peaks to lower angles for the La, Sr-doped  $\text{CeO}_2$  sample. The increase in the lattice parameter for La-Sr Co-doping  $\text{CeO}_2$  could possibly be due to the larger ionic radii of La and Sr compared to  $\text{Ce}^{4+/3+}$  ions.<sup>29</sup>

Moreover, the HR-TEM images confirmed the nanosheets-type morphology of the synthesized LSC powders. The LSC nanosheets had diameters of 20–30 nm as shown in Fig. 2a–c. Afterward, HR-TEM was used to study the crystal planes and the lattice spacing of 0.26 nm corresponding to the (111) plane of cubic LSC fluorite structure, indicating that when the (111)

planes were exposed to the crystal face some structural defects/disorder were generated at the exposed surfaces, as shown in Fig. 2d and e. Furthermore, selected area electron diffraction (SAED) was performed for the multiplanes of the LSC pattern, as shown in Fig. 2f, where the main diffraction patterns of the (111), (002), and (220) planes can be observed.<sup>28,29</sup>

Furthermore, Fig. 3a shows the high annular dark-field image of  $\text{La}_{0.15}\text{Sr}_{0.05}\text{Ce}_{0.8}\text{O}_{2-\delta}$ , where the grain level interactions of each grain can be clearly seen with sophisticated and fine nano-sized sheets with diameters of 20–30 nm.<sup>30</sup> To confirm the elemental distribution at the particle level, energy-dispersive spectrometry (EDS) mapping was performed using HAADF-STEM, as shown in the image in Fig. 3a. Fig. 3b–f show the elemental mapping of  $\text{La}_{0.15}\text{Sr}_{0.05}\text{Ce}_{0.8}\text{O}_{2-\delta}$ , where a homogenous chemical spreading of each individual element, such as La, Sr, and Ce, could be seen, confirming the well-controlled nanosheets formation with a homogenous chemical concentration.<sup>31</sup>

### 3.2 Spectroscopic analysis of the LSC electrolyte

Fig. 4a displays the typical SEM image of the LSC electrolyte. The SEM image also shows that the LSC electrolyte seemed to be fully nanostructured. Further, XPS spectroscopy was used to deeply understand the characteristics of La and Sr doping on the structural properties.<sup>21,32</sup> Moreover, Fig. 4b–f show the XPS plots of the high-resolution XPS spectra and XPS survey of each

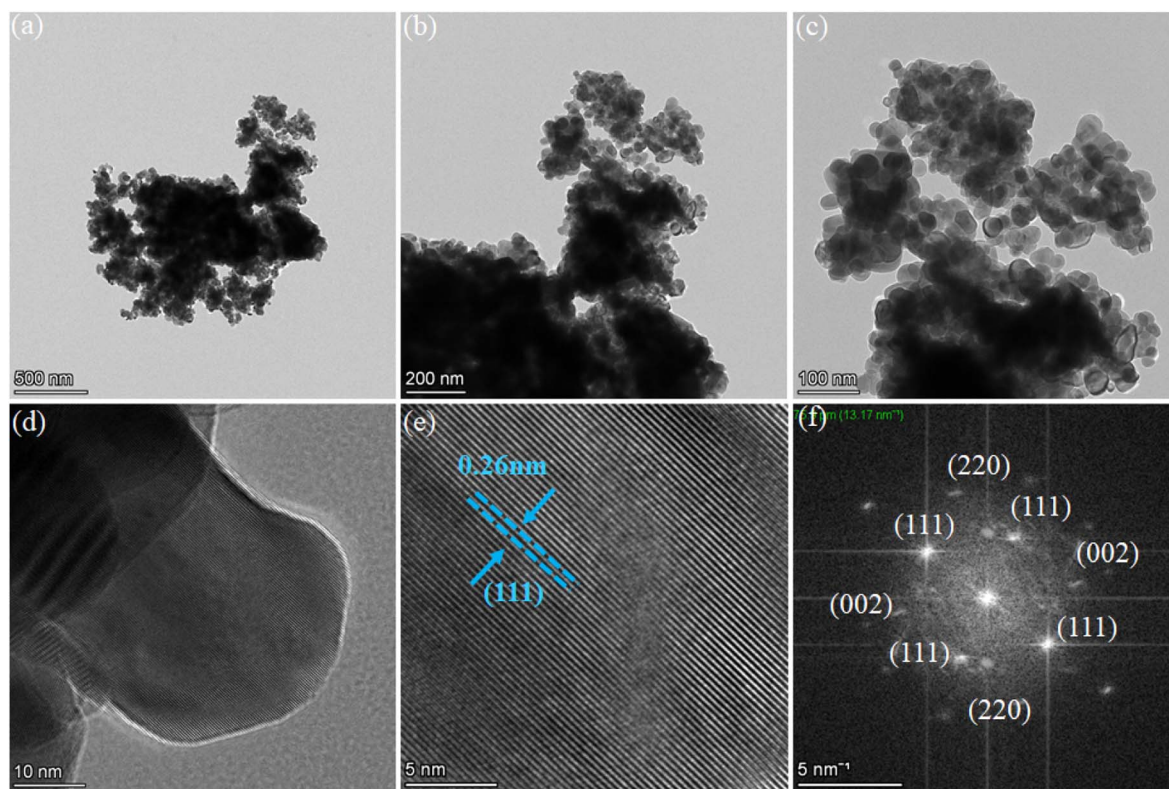


Fig. 2 (a–f) HR-TEM images of  $\text{La}_{0.15}\text{Sr}_{0.05}\text{Ce}_{0.8}\text{O}_{2-\delta}$ : (a–c) morphology display of LSC at different scales: 500, 200, and 100 nm, (d and e), crystal planes of the  $\text{La}_{0.15}\text{Sr}_{0.05}\text{Ce}_{0.8}\text{O}_{2-\delta}$  sample, where the (111) *d*-spacing plane can be observed, and (f) selected area electron diffraction pattern of  $\text{La}_{0.15}\text{Sr}_{0.05}\text{Ce}_{0.8}\text{O}_{2-\delta}$ , respectively.



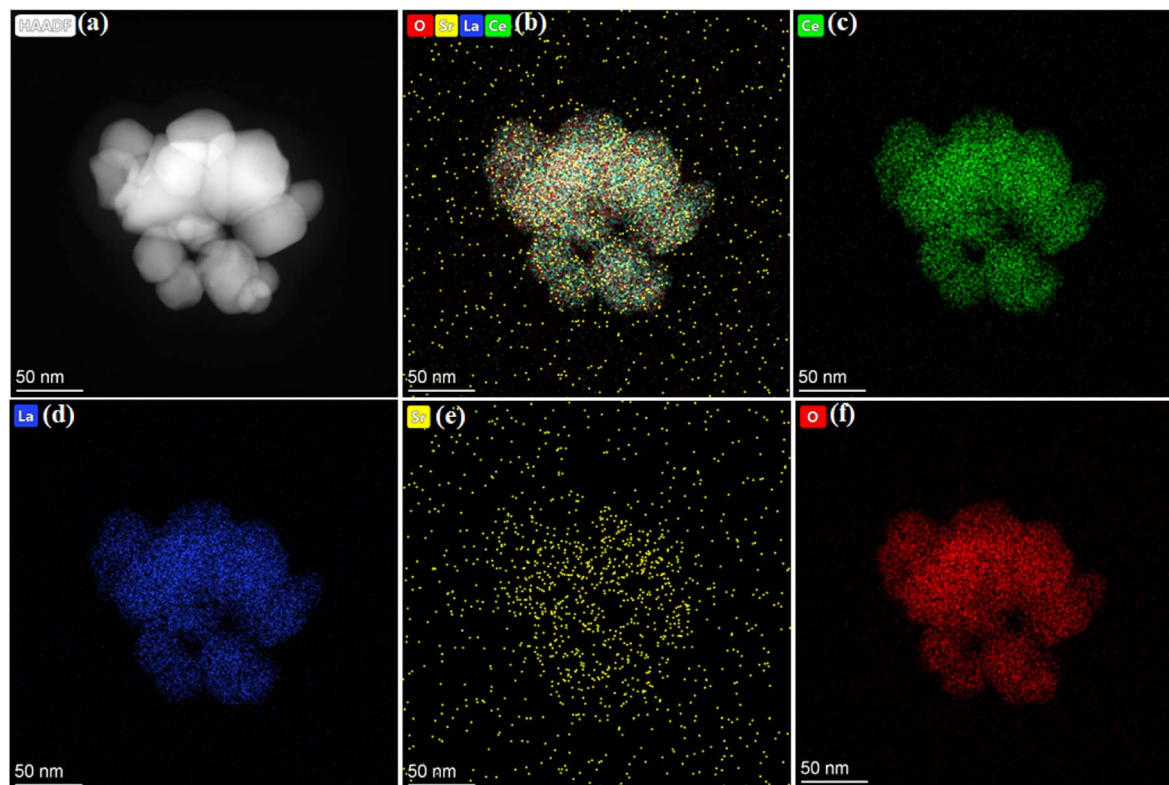


Fig. 3 (a) High annular dark-field image of  $\text{La}_{0.15}\text{Sr}_{0.05}\text{Ce}_{0.8}\text{O}_{2-\delta}$ . (b–f) combined and individual element mapping of each element present in  $\text{La}_{0.15}\text{Sr}_{0.05}\text{Ce}_{0.8}\text{O}_{2-\delta}$ , such as La, Sr, Ce, and O, respectively.

element in the pristine and doped sample for comparison. Fig. 4b shows the XPS spectra of the Ce-3d of pristine  $\text{CeO}_2$  and of the La- and Sr-doped sample. The  $\text{Ce}^{4+}$  and  $\text{Ce}^{3+}$  peaks of the LSC sample appeared at 881.8/901.2 eV and 898.2/916.3 eV, whereas for pure  $\text{CeO}_2$ , the  $\text{Ce}^{4+}$  and  $\text{Ce}^{3+}$  peaks were at 882.212/901.65 and 898.52/916.78 eV,<sup>16,28,29</sup> respectively, showing a significant change in the binding energy as a shift down to about  $0.5 \pm 0.05$  eV for LSC, related to the Ce-3d of  $\text{CeO}_2$ . However, the  $\text{Ce}^{3+}$  concentration in LSC was higher than that of the pure  $\text{CeO}_2$ , indicating an obvious signal for the presence of oxygen vacancies and maintaining the overall charge neutrality. The O 1s XPS spectra could further provide information about its ionic conductivity. However, the O 1s spectra of LSC comprised peaks for lattice  $\text{O}_2$ , highly oxidant  $\text{O}_2$ , surface-active oxygen, and surface adsorbed water, wherein the representative peaks were located at 528.8, 529.5, 530.1, and 531.4 eV, respectively. The LSCC sample showed a relatively high surface-active to lattice  $\text{O}_2$  ratio compared to that of the individual  $\text{CeO}_2$  (as shown in Fig. 4c). This indicator can be used for the assessment of LSC as having high adsorption abilities. The high surface-active to lattice  $\text{O}_2$  ratio also provides information that La- and Sr-doping in the  $\text{CeO}_2$  lattice provide better surface coverage for absorbing hydroxide species, which is mandatory for  $\text{O}^{2-}/\text{H}^+$  transport.<sup>22,23</sup> However, the XPS results confirmed that the gradiently doped  $\text{CeO}_2$  sample had a high surface-active oxygen species density, which could be described as providing an important role in proton conduction. The effects of gradient

doping on the proton transport and electrochemical performance are further explored in the following parts.

Furthermore, TGA/DSC analysis was carried out for  $\text{CeO}_2$  and LSC in the air at a temperature range of 25–800, and plots of the obtained data are shown in Fig. 5a and b, respectively. The initial weight loss at a low operating temperature of 25–180 could be due to the physisorbed water content.<sup>31,32</sup> The mass change (loss) progressively increased with temperature up to a high value above 600, when the powders started to be thermally reduced and released the lattice oxygen to form oxygen vacancies. The LSC sample showed the highest mass loss at this temperature, indicating that the LSC sample could exhibit more ionic conductivity compared to the pure  $\text{CeO}_2$  sample.<sup>33</sup> The differential scanning calorimetry (DSC) curve showed obvious thermal effects, indicating that there were 1st-order phase transitions in both samples. Furthermore, Raman spectroscopy was employed for structural characterization of the LSC in comparison to pure  $\text{CeO}_2$ . Fig. 5c displays the specific Raman shifts in the bands centered at  $463\text{ cm}^{-1}$  corresponding to the stretching vibrations of F<sub>2g</sub>, indicating the symmetric vibration of oxygen ions around  $\text{Ce}^{4+}$  ions in  $\text{CeO}_2$  (octahedra sites).<sup>29,34–36</sup> The significant shift to lower wavenumber values was due to the cell expansion resulting from the substitution of dopants ions in the  $\text{CeO}_2$  lattice and the subsequent loss of oxygen around the cation sites. A singlet Raman peak observed around  $\sim 1083\text{ cm}^{-1}$  in the calcined and sintered sample of LSC arose from the symmetric stretching vibration of  $\text{Na}_2\text{CO}_3$ , which



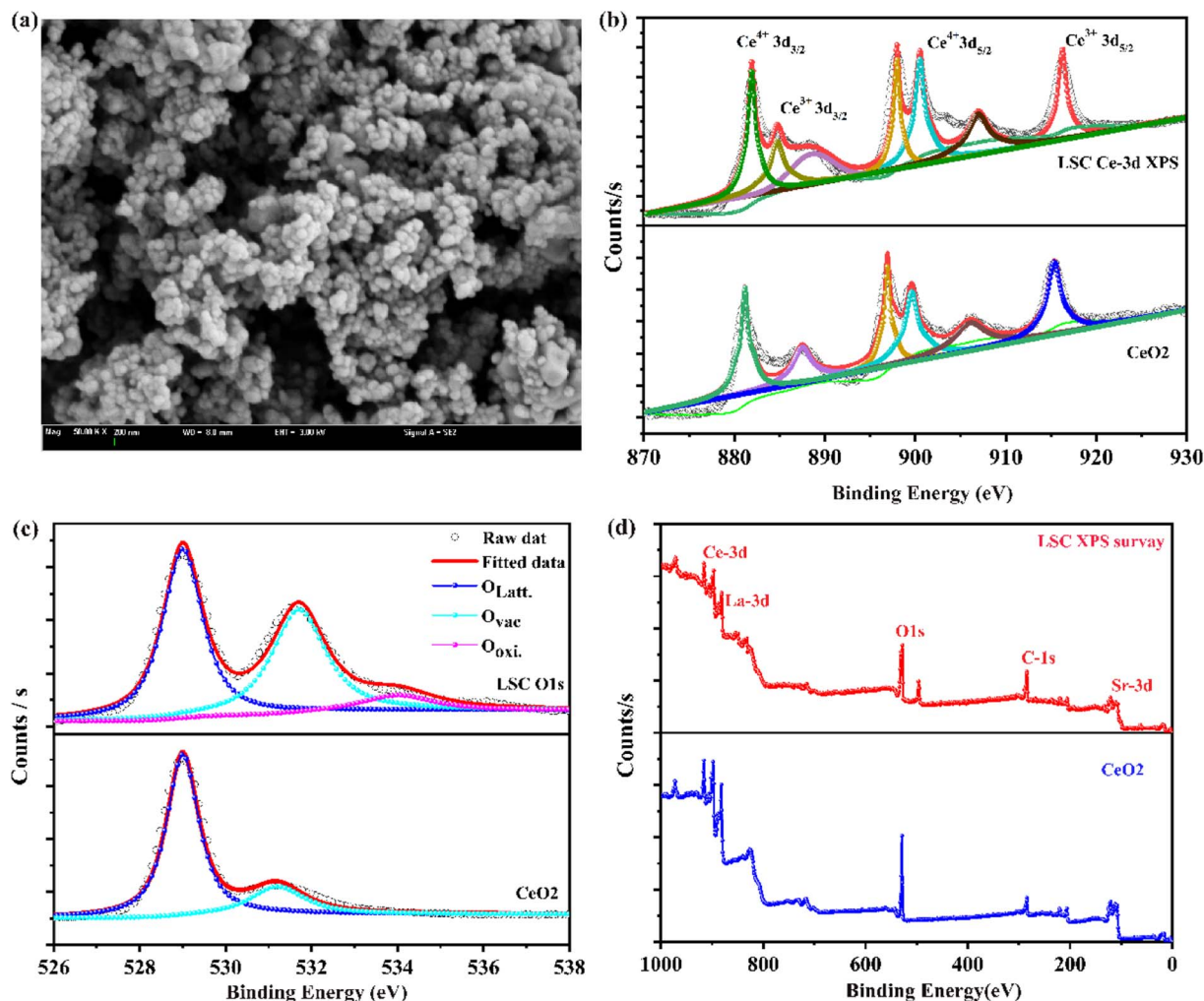


Fig. 4 (a) SEM image of  $\text{La}_{0.15}\text{Sr}_{0.05}\text{Ce}_{0.8}\text{O}_2$  powders, (b–d) X-ray photoelectron spectra of LSC and  $\text{CeO}_2$  corresponding high-resolution XPS spectra of Ce-3d, O 1s, and the XPS survey, respectively.

could be because of the precipitating agent. Further, the effect of the co-doping of La and Sr into  $\text{CeO}_2$  was investigated through the energy band gap by UV-vis spectroscopy, as shown in Fig. 5d. Initially, UV-vis was employed to record the absorbance spectra in the range from 200–800 nm. There was a sharp absorption peak, and the spectra of  $\text{CeO}_2$  showed slightly higher absorbance than LSC. It could be noticed that there was a red-shift at the absorption edge and this red-shift signified the shallow levels in the energy bandgap due to the increase in doping concentration. The optical bandgaps of both the intermediate and lower concentrations of doping could be calculated via the following relation:<sup>33</sup>  $(\alpha h\nu) = \beta_o (h\nu - E_g)^n$ , where  $\alpha$ ,  $h\nu$ ,  $E_g$ ,  $\beta_o$ , and  $n$  denote the absorbance coefficient, energy of photons, energy band gap, constant, and the type of bandgap, respectively. The linear part of the graph where  $h\nu$  and  $(\alpha h\nu)^2$  meet the  $x$ -axis gives the value of the bandgap.<sup>34</sup> The obtained energy bandgaps via that equation were 3.2 and 2.91 eV for  $\text{CeO}_2$  and LSC, respectively, which shows that co-doping induced a lower bandgap value than for pure  $\text{CeO}_2$ . This lowering of the energy bandgap value suggested the formation of an

intermediate energy level in-between the conduction and valence band levels. Additionally, the decrease in the bandgap is also helpful for reducing the activation energy for the transport of ionic conduction. Therefore, analysis of the energy bandgaps based on the co-doping of La and Sr into  $\text{CeO}_2$  provided us with evidence to assist in the formation of oxygen vacancies during the transport of charge carriers from the valence band to the conduction band, which can lead to the easy and fast transport of ion conduction.<sup>35–37</sup>

### 3.3 Electrochemical performance measurements

Fig. 6a–e show the  $I$ - $V$  and  $I$ - $P$  characteristic plots of the fuel cells with  $\text{CeO}_2$ , Sr-doped  $\text{CeO}_2$ , La-doped  $\text{CeO}_2$ , 10% Sr, 10% La-doped  $\text{CeO}_2$ , and 5% Sr- and 15% La-doped  $\text{CeO}_2$  electrolytes-based fuel cells operating at 550 °C, respectively. The highest OCV of above 1.1 V and power density of 910  $\text{mW cm}^{-2}$  were achieved for the 15% La and 5% Sr- $\text{CeO}_2$ -based electrolyte fuel cell, while the pure  $\text{CeO}_2$  electrolyte-based fuel cell exhibited a power densities of 520  $\text{mW cm}^{-2}$ ; while the values were 660 and 750  $\text{mW cm}^{-2}$  for the individual Sr and La-



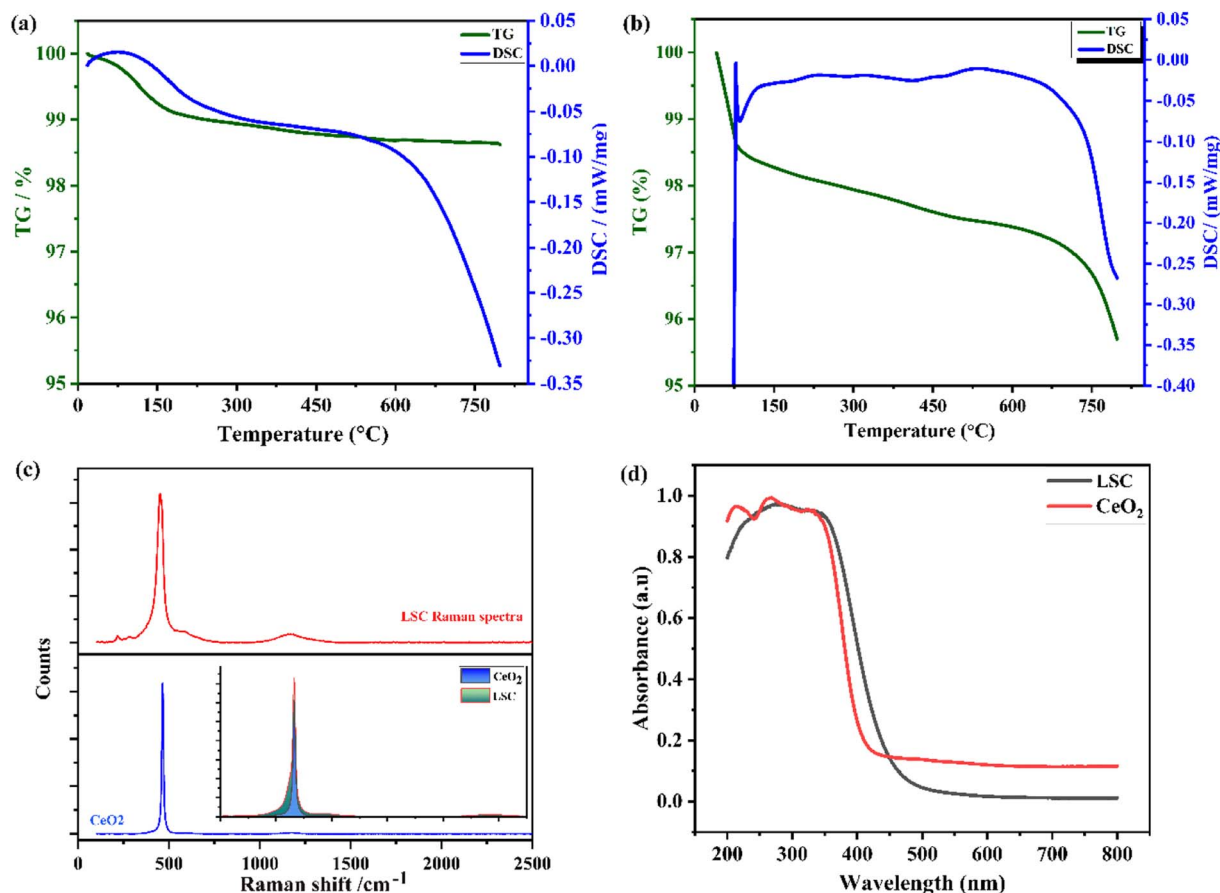


Fig. 5 (a and b) Thermogravimetric and differential scanning calorimetry analysis (TGA/DSC) for evaluating the thermal properties of  $\text{CeO}_2$  and LSC, respectively. (c and d) Raman and UV-visible absorbance and spectra of the  $\text{CeO}_2$  and LSC electrolytes, respectively.

doped  $\text{CeO}_2$  electrolyte-based fuel cells, respectively. The OCV value of 1.12 V at 550 °C was in good agreement with the Nernst theoretical potential, expressing the ionic transfer number in the LSC electrolyte was near to unity.<sup>18,38</sup>

Moreover, the 15%La and 5%Sr-doped  $\text{CeO}_2$ -based electrolyte fuel cell also displayed an admirable electrochemical performance at lower operating temperatures. Fig. 6e shows the  $I$ - $V$  and  $I$ - $P$  curves of the LSC electrolyte operating at different temperatures, and power densities of 640, 420, and 228  $\text{mW cm}^{-2}$  were achieved at 500 °C, 450 °C, and 400 °C, respectively. Therefore, the higher performance of the LSC electrolyte fuel cell over the  $\text{CeO}_2$  and individual La- and Sr-doped  $\text{CeO}_2$  cells further show the vital role of the co-doping of tri-valent La and di-valent Sr to drive faster ion transport with a lower energy barrier. Fig. 6f displays a typical-three-layered cross-sectional SEM image of the LSC electrolyte-based fuel cell attained after online sintering and a fuel test. The SEM image shows that the LSC electrolyte layer was fully dense, without obvious connected pores, and was well-adhered on the porous anode layer.

### 3.4 Electrochemical impedance and electrical properties

Moreover, EIS was used to understand the electrochemical process of the LSC electrolyte-based fuel cell compared with the pure  $\text{CeO}_2$  electrolyte-based cell. Fig. 6a–d show the Nyquist and

Bode plots of the EIS spectra in OCV conditions. Furthermore, model circuiting was used to model the EIS spectra, *e.g.*,  $R_0 - (R_1 - \text{CPE}_1) - (R_2 - \text{CPE}_2)$ , where  $R_0$  corresponds to the ohmic contribution to the total area specific resistance (ASR) that come from the electrolyte layer, and  $R_1$  and  $R_2$  refer to the electrode polarization resistances (charge- and mass-transport resistances).<sup>39</sup> The model circuiting results showed that both the ohmic-ASR and polarization resistance were dramatically reduced using the LSC electrolyte-based fuel cell compared with the pure  $\text{CeO}_2$  electrolyte cell. The low ohmic-ASR of the gradiently doped LSC strongly supported the electrochemical performance, whereby a highly improved ionic transport was observed. Moreover, the decrease in electrode polarization also showed that the electrocatalytic dissociation of  $\text{H}_2$  and oxygen reduction reaction with the LSC electrolyte were improved along with the ionic conductivity.<sup>41</sup> The fuel cell with the LSC electrolyte exhibited an ohmic-ASR of only  $0.08 \Omega \text{ cm}^2$ , while the charge and mass transport ASR ( $R_1$  &  $R_2$ ) were  $0.07$  and  $0.38 \Omega \text{ cm}^2$ , respectively at 550 °C. Whereas, the fuel cell with pure  $\text{CeO}_2$  showed an ohmic-ASR of  $0.26 \Omega \text{ cm}^2$  and  $R_2$  values of  $0.14$  and  $0.42 \Omega \text{ cm}^2$ , respectively.<sup>17,18</sup> The small ohmic-ASR could only be due to  $R_1$  the high ionic conductivity in the electrolyte and unfortunately is common when using conventional doping methods, like just trivalent dopants  $\text{Sm}^{3+}$ ,  $\text{Gd}^{3+}$ , and  $\text{Y}^{3+}$ , where the vacancies remain immobile when the sample is reduced or



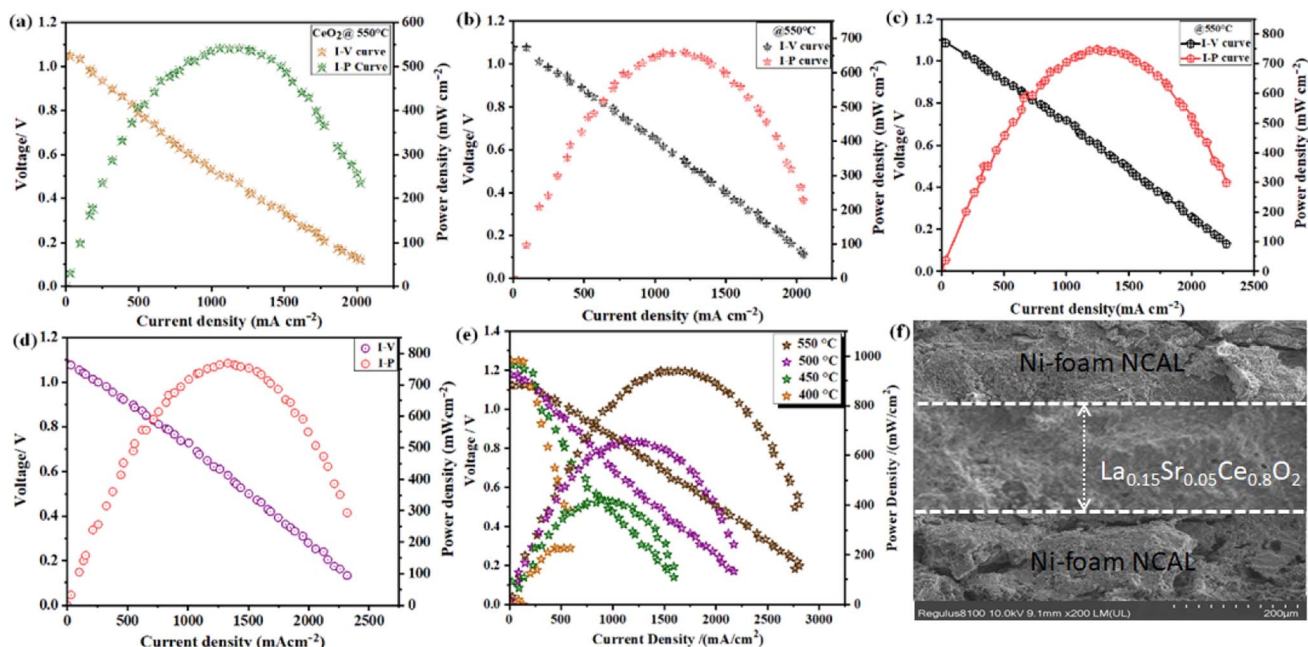


Fig. 6 (a–d) Typical  $I$ - $V$  characteristics compared among the  $\text{CeO}_2$  electrolyte, 20%Sr-doped  $\text{CeO}_2$ , 20%La-doped  $\text{CeO}_2$ , and 10%Sr- and 10% La-doped  $\text{CeO}_2$  electrolytes, respectively at 550 °C, and (e)  $I$ - $V$  characteristics for the LSCC cell at different temperatures from 400–550 °C with 5%Sr and 15% La-doped  $\text{CeO}_2$ . (f) Cross-sectional SEM images of a tri-layer fuel cell with an LSC electrolyte layer after online sintering and a fuel cell test, respectively.

contribute to electronic conduction when used as an electrolyte.<sup>19,20</sup> However, the ionic conductivity of the LSC and all the other prepared electrolytes were calculated using model circuiting of the measured impedance spectra. The ionic conductivity of  $0.05 \text{ S cm}^{-1}$  was achieved for LSC, as shown in Fig. 7e. These achieved results show that the ionic conductivity of the LSC fluorite structure was significantly higher than that of

most the used state-of-the-art oxygen- and proton-conducting electrolytes, *e.g.*, YSZ, doped ceria, and barium zirconate-based perovskite materials.<sup>17</sup>

### 3.5 Density functional theory calculations

In order to deepen the understanding of gradient doping on the electronic structure, the density of states (DOS) for each

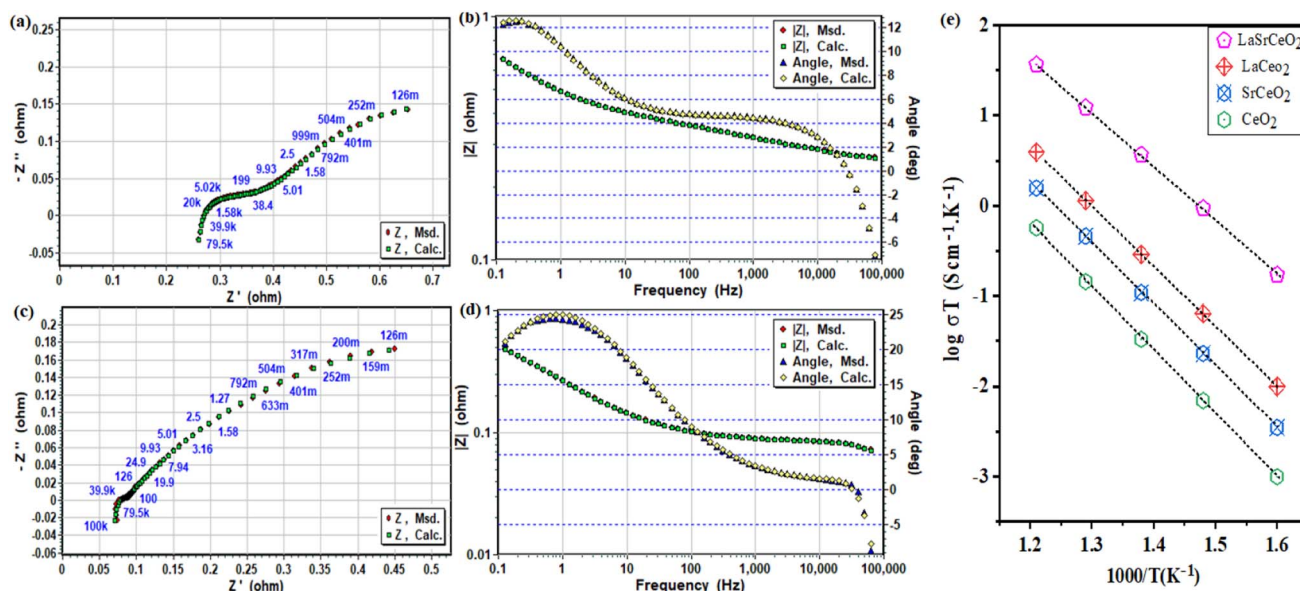


Fig. 7 (a and b) Nyquist and Bode plots of the impedance spectra measured in  $\text{H}_2/\text{air}$  at 550 °C with the corresponding fitting data for fuel cells with  $\text{CeO}_2$  electrolyte. (c and d) Nyquist and Bode plots of the impedance spectra measured in  $\text{H}_2/\text{air}$  at 550 °C with the corresponding fitting data for fuel cells with LSC electrolyte. (e) Comparison of the ionic conductivity of LSC with other developed electrolytes.



individual orbital was calculated by the DFT<sup>24</sup> method to investigate the accountable mechanism for the enhanced ionic conductivity. Fig. 8a–c display the optimized structures of CeO<sub>2</sub>, La-doped CeO<sub>2</sub>, and the La, Sr-doped CeO<sub>2</sub> structure.

However, using the conventional bulk doping, such as the Sm<sup>3+</sup>/Gd<sup>3+</sup>-doped case, oxygen vacancies can only be connected at long range. Therefore, O<sub>1</sub><sup>2−</sup> and O<sub>2</sub><sup>2−</sup> are partially occupied by different mixed coordination with d-orbital

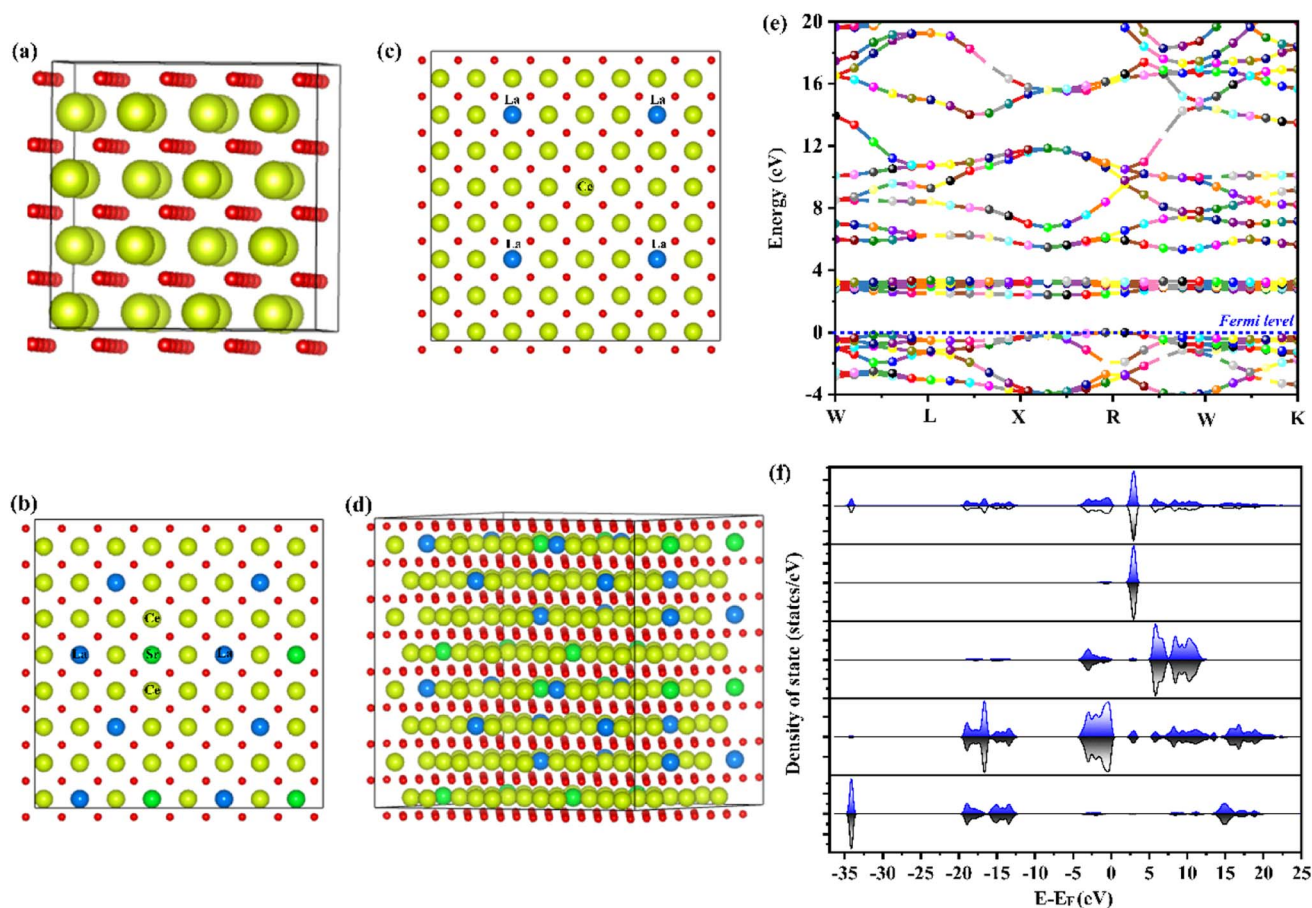


Fig. 8 (a–d) Optimized structure of pure CeO<sub>2</sub>, La-doped CeO<sub>2</sub>, and La and Sr-doped CeO<sub>2</sub>, respectively. (e and f) Calculated energy band structure and DOS results performed for the different orbitals of each species, such as s, p, f, and total DOS for the LSC structure.

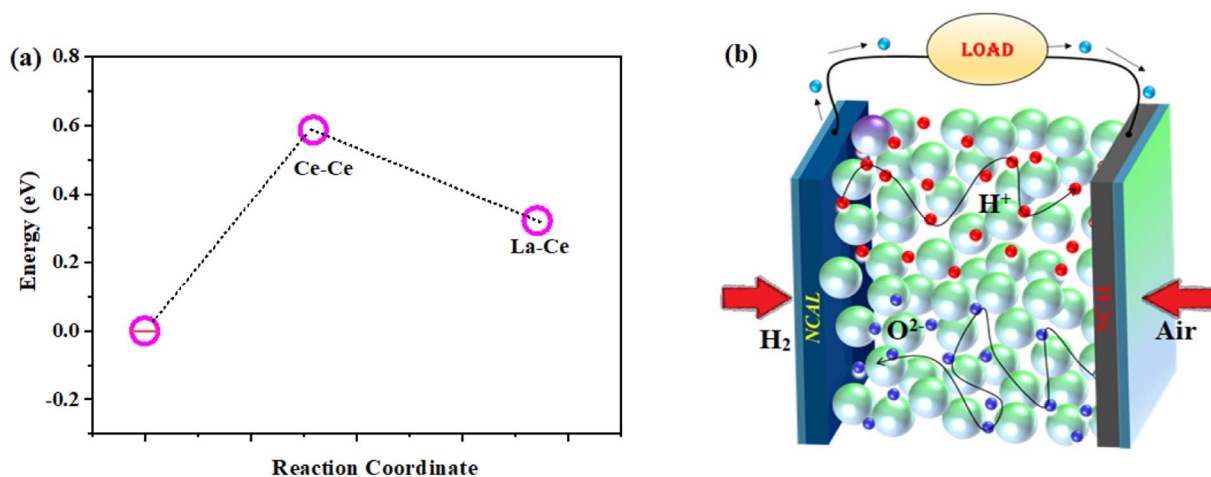


Fig. 9 (a) Calculated oxygen vacancy formation energy in the bulk structures of LSC for the possible migration of oxygen ions. (b) Schematic diagram of the system and power output mechanism of our prepared fuel cell using the LSC electrolyte.



metal ions. However, the arrangement of vacancies needs a high activation energy for a random distribution. However, we introduced a combination of di- and tri-valent doping with bulk doping to consider mixed valence doping to connect the small-range vacancies with low activation energies. This positional disorder of vacancies allowed forming easy transport units with variable local geometries. Meanwhile, the energy band structure and partial DOS for pure LSC structures were calculated as shown in Fig. 8e and f, respectively. The DOS showed a gradually increased DOS value near the Fermi level, correspond to a gradually enhanced chemical activity for the absorption of ions.<sup>25,26</sup> Also, the DOS plots of LSC revealed that La, Sr doping could narrow the bandgaps between the valence and conduction bands (Fig. 8e), corresponding to a gradually increased electron concentration.

Afterward, we calculated the vacancy formation energy along different sites of the LSC structure to support our experimental findings. Fig. 9a shows the obtained results, where the minimum vacancy formation energy was exhibited by the La- $\bar{V}_{\text{O}}$ -Ce site.<sup>40–45</sup> A schematic diagram of the fuel cell and the mechanism of the power output are shown in Fig. 9b.

## 4 Conclusion

The work performed herein demonstrated the excellent ionic conducting properties in nanosheets of CeO<sub>2</sub> with a suitable combination of di- and tri-valent doping of La and Sr. The dexterous design showed a drastic improvement in power output and ionic conductivity. Further, we used different microscopic and spectroscopic analyses to study the mechanism behind the drastic increase in ionic conductivity of the LSC nanosheets in a fuel cell operation. We found that the synthesized nanostructured LSC structure could easily be modified to form an oxygen-deficient structure that facilitates ionic transport effectively when a suitable amount of di- and tri-valent elements is doped. The *ex situ* spectroscopies, including EIS, UV-visible, XPS, and TGA, clearly described and confirmed our findings and the structural change properties of the LSC electrolyte. In conclusion, this method could form the basis of interest to develop new oxide and proton-conducting electrolytes, which could be useful for all energy devices and material systems.

## Conflicts of interest

There are no conflicts to declare.

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