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# Influence of the physical properties of the layered oxyselenides Bi<sub>2</sub>YO<sub>4</sub>Cu<sub>2</sub>Se<sub>2</sub> through Ni doping

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We have synthesized a series of Ni-doped layered oxyselenides  $Bi_2YO_4Cu_{2-x}Ni_xSe_2$  ( $0 \le x \le 0.4$ ). The crystal structure and physical properties were studied through X-ray diffraction, and electric and thermo transport measurements. We also performed DFT calculations to study the electric structure of the designed  $Bi_2YO_4Ni_2Se_2$ , which is similar to that of  $KNi_2Se_2$ .

#### Introduction

In recent decades, compounds with layered crystal structures have attracted people's increasing attention due to their interesting physical properties, such as superconductivity, <sup>1-3</sup> charge density wave, <sup>4-9</sup> topological insulation <sup>10</sup> *etc.* In particular, the discovery of iron-based superconductors in 2008 aroused wide concern all over the world. <sup>2</sup> Up to now, the mechanism of high-temperature superconductivity remains a mystery, which requires the joint efforts of theoretical and experimental researchers. According to previous research experience, high-temperature superconductors tend to adopt layered crystal structures. And it has been an important way to search for new superconductors through exploring new layered compounds.

Among the known iron-based superconductors, they all have the  $Fe_2Pn_2/Fe_2Ch_2$  (Pn = P, As; Ch = S, Se, Te) conducting layer, which is believed playing an important role in the superconductivity. In fact, all the iron-based superconductors have analogues in CuCh based compounds. For example, the 1111type LaOFeP/As and LaOCuS/Se adopt the same crystal structure. The 122-type KFe<sub>2</sub>Se<sub>2</sub> and KCu<sub>2</sub>Se<sub>2</sub>, the 32 522-type Sr<sub>3</sub>Sc<sub>2</sub>O<sub>5</sub>Fe<sub>2</sub>As<sub>2</sub> and Sr<sub>3</sub>Sc<sub>2</sub>O<sub>5</sub>Cu<sub>2</sub>S<sub>2</sub> etc. All have the same situations.13-15 Many iron-based superconductors were discovered by substituting Fe<sub>2</sub>Pn<sub>2</sub> layers for Cu<sub>2</sub>Ch<sub>2</sub> layers in the known CuCh-based compounds. In 2002, a new layered CuSebased compound Bi<sub>2</sub>YO<sub>4</sub>Cu<sub>2</sub>Se<sub>2</sub> was reported and we studied the physical and electric properties lately. 16-18 The Cu<sub>2</sub>Se<sub>2</sub> layer is defined as the conducting layer and the Bi<sub>2</sub>YO<sub>4</sub> layer is defined as the block layer. As every block layer provide one electron to conducting layer, which causes a mixed state of Cu<sup>2+</sup>/Cu<sup>+</sup>, it exhibits p-type metallic behaviour which is rare in

#### Results and discussion

The crystal structure of Bi<sub>2</sub>YO<sub>4</sub>Cu<sub>2</sub>Se<sub>2</sub> is shown in Fig. 1(a), which is alternatively stacked by Bi<sub>2</sub>YO<sub>4</sub> layer and Cu<sub>2</sub>Se<sub>2</sub> layer. Fig. 1(b) shows the crystal structure of KNi<sub>2</sub>Se<sub>2</sub>, which has the ThCr<sub>2</sub>Si<sub>2</sub>-type structure.<sup>20</sup> Comparing these two structures, it is not difficult to find that their crystal structure is the same except the block layer. Fig. 2 shows the X-ray diffraction patterns of Bi<sub>2</sub>YO<sub>4</sub>Cu<sub>2</sub>-<sub>x</sub>Ni<sub>x</sub>Se<sub>2</sub>. All the diffraction peaks of the doped samples match the Bragg peaks of Bi<sub>2</sub>YO<sub>4</sub>Cu<sub>2</sub>Se<sub>2</sub>, which indicates high quality samples. The fitting lattice parameters of Bi<sub>2</sub>YO<sub>4</sub>Cu<sub>2</sub>-<sub>x</sub>Ni<sub>x</sub>Se<sub>2</sub> are shown in Fig. 2(b). The lattice parameter *a* and *c* for Bi<sub>2</sub>YO<sub>4</sub>Cu<sub>2</sub>Se<sub>2</sub> are 3.862 and 24.31 Å respectively, which is corresponding to previous results.<sup>16,17</sup> The lattice parameter a is decreased with increasing Ni content, while the lattice parameter *c* is the opposite. It indicates that the in-plane

CuCh-based compounds. Similarly, p-type metallic ground state is also found in KCu<sub>2</sub>Se<sub>2</sub>, in which the K layer also provides one electron to Cu<sub>2</sub>Se<sub>2</sub> layer per unit cell. In 2012, the analogue KNi<sub>2</sub>Se<sub>2</sub> was reported superconducting below 0.8 K, which has aroused people's extensive research interest. 19 If the Cu<sub>2</sub>Se<sub>2</sub> layer in Bi<sub>2</sub>YO<sub>4</sub>Cu<sub>2</sub>Se<sub>2</sub> was replaced by Fe<sub>2</sub>Pn<sub>2</sub> or Ni<sub>2</sub>Pn<sub>2</sub> layers, we think new superconductors may be achieved. We try our best to synthesize Bi<sub>2</sub>YO<sub>4</sub>Ni<sub>2</sub>Se<sub>2</sub> through solid state reaction. Unfortunately, all the attempts ended in failure. On the other hand, it is also valuable to study the influence of the physical proper-ties of the layered oxyselenide Bi<sub>2</sub>YO<sub>4</sub>Cu<sub>2</sub>Se<sub>2</sub> through Ni doping at Cu site. In this paper, we successfully doped 20% Ni at Cu site in Bi<sub>2</sub>YO<sub>4</sub>Cu<sub>2</sub>Se<sub>2</sub>. The temperature coefficient of resistivity can be effectively adjusted through Ni doping. Furthermore, we studied the thermoelectric properties of the Ni-doped samples. Finally, we study the electric structure of the designed Bi<sub>2</sub>YO<sub>4</sub>Ni<sub>2</sub>Se<sub>2</sub> through the first principle calculation, which exhibit similar electronic structure with KNi2Se2. This will tempt people to explore the synthesis of Bi<sub>2</sub>YO<sub>4</sub>Ni<sub>2</sub>Se<sub>2</sub> through other methods.

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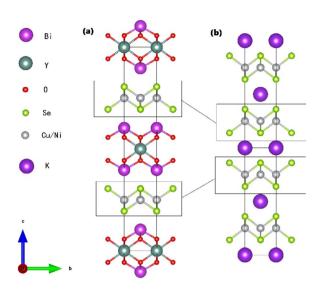


Fig. 1 The crystal structure of Bi<sub>2</sub>YO<sub>4</sub>(Cu,Ni)<sub>2</sub>Se<sub>2</sub> (a) and KNi<sub>2</sub>Se<sub>2</sub> (b)

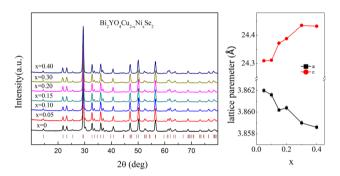


Fig. 2 The X-ray diffraction patterns and the fitted lattice parameters of  $Bi_2YO_4Cu_{2-x}Ni_xSe_2$  measured at 300 K.

interaction is strengthened and the interlayer interactions are weakened. The reported lattice parameters a and c of KCu<sub>2</sub>Se<sub>2</sub> are 4.008 and 13.643 Å respectively. Compared with KCu<sub>2</sub>Se<sub>2</sub>, the thicker block layer in Bi<sub>2</sub>YO<sub>4</sub>Cu<sub>2</sub>Se<sub>2</sub> makes the parameter a smaller. In other words, the in-plane interaction of Bi<sub>2</sub>YO<sub>4</sub>-Cu<sub>2</sub>Se<sub>2</sub> is stronger than that of KCu<sub>2</sub>Se<sub>2</sub>. For KNi<sub>2</sub>Se<sub>2</sub>, the fitted lattice parameter a and c according to experiment result are 3.89295 and 13.3158 Å respectively and the calculated lattice parameter a and c according to DFT calculation are 3.9680 and 13.0479 Å respectively. The lattice parameter a of KNi<sub>2</sub>Se<sub>2</sub> is larger than that of Bi<sub>2</sub>YO<sub>4</sub>Cu<sub>2</sub>Se<sub>2</sub>. According to the above results, the lattice parameter a of Bi<sub>2</sub>YO<sub>4</sub>Ni<sub>2</sub>Se<sub>2</sub>, if it exists, should be larger than that of KNi<sub>2</sub>Se<sub>2</sub>. This may be the reason we cannot doped much more Ni in Bi<sub>2</sub>YO<sub>4</sub>Cu<sub>2</sub>Se<sub>2</sub>.

Fig. 3 shows the temperature dependence of the electric resistivity of Bi<sub>2</sub>YO<sub>4</sub>Cu<sub>2-x</sub>Ni<sub>x</sub>Se<sub>2</sub>. For Bi<sub>2</sub>YO<sub>4</sub>Cu<sub>2</sub>Se<sub>2</sub>, the resistivity increases with increasing temperature from 2 to 300 K, which indicates a metallic transport behaviour and is similar with previous results.<sup>17</sup> The value of resistivity is increased with increasing Ni content. Compared with KNi<sub>2</sub>Se<sub>2</sub>, the resistivity is relatively higher. Unfortunately, no superconducting behaviour was found until 2 K for all the Ni doped Bi<sub>2</sub>YO<sub>4</sub>Cu<sub>2</sub>Se<sub>2</sub>.

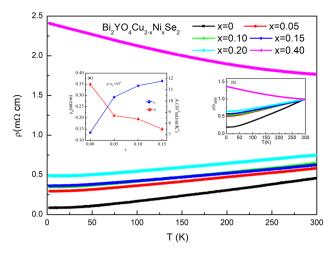


Fig. 3 The temperature dependence of the resistivity of  $Bi_2YO_4$ - $Cu_{2-x}Ni_xSe_2$  from 2 to 300 K. Inset (a) is the parameter  $\rho_0$  and A fitted through the Fermi liquid behavior equation with different Ni con-tent. Inset (b) is the reduced resistivity of  $Bi_2YO_4Cu_{2-x}Ni_xSe_2$ .

According to previous study,17 the metallic behaviour originates from the mixed valences of Cu ion in Cu<sub>2</sub>Se<sub>2</sub> layer. Every Bi<sub>2</sub>YO<sub>4</sub> unit provides one electron to Cu<sub>2</sub>Se<sub>2</sub> unit, so every two Cu atoms only pro-vide 3 electrons to keep balance. While when Ni is doped at Cu site, every Ni atom can provide two electrons, namely electronic doping. For undoped Bi<sub>2</sub>YO<sub>4</sub>Cu<sub>2</sub>Se<sub>2</sub>, the carriers are mainly p-type. Regardless of other factors, Ni doping at Cu site decreases the carrier con-centration and the resistivity would increase with increasing Ni content. Our experimental results are consistent with the viewpoint. As shown in Fig. 3, the reduced resistivity increases with increasing Ni content. Furthermore, the Bi<sub>2</sub>YO<sub>4</sub>Cu<sub>1.6</sub>Ni<sub>0.4</sub>Se<sub>2</sub> exhibits semiconducting/insulating transport behaviour from 2 to 300 K. The resistivity (x = 0–0.2) obey the Fermi liquid behaviour from 2 K to 50 K. We fitted the resistivity curves through the Fermi liquid behaviour equation  $\rho = \rho_0 + AT^2$ . Where  $\rho_0$  represents the residual resistivity. The results are shown in the inset (a) of Fig. 3. As can be seen, the residual resistivity in-creases with x. As known, the residual resistivity  $\rho_0$  is in direct proportion to  $m^*\hbar/n\tau_0$ , where  $m^*$ , n and  $\hbar/\tau_0$  represent the carrier effective mass, carrier concentration and the scattering rate related to the disordered potential. For Bi<sub>2</sub>YO<sub>4</sub>Cu<sub>1.6</sub>Ni<sub>0.4</sub>Se<sub>2</sub>, it exhibits semiconducting transport behaviour that the resistivity decreases with arising temperature. Further study is needed to clarify the ground state of Ni doped Bi<sub>2</sub>YO<sub>4</sub>Ni<sub>2</sub>Se<sub>2</sub>. According to previous study, the ground state of KNi<sub>2</sub>Se<sub>2</sub> is metallic.<sup>20</sup> Our following DFT calculation result also indicates the metallic ground state of Bi2YO4Ni2Se2. It should not turn it to semiconductor through Ni doping. The semiconducting-like phenomenon should not be intrinsic, which may be caused by defect or grain boundary introduced with Ni doping.

Fig. 4 shows the thermoelectric properties of  $\mathrm{Bi_2YO_4Cu_2}_{-x^2}\mathrm{Ni}_x\mathrm{Se_2}$  (x=0,0.2,0.3). Fig. 4(b) shows the Seebeck coefficient curves of  $\mathrm{Bi_2YO_4Cu_2}_{-x}\mathrm{Ni}_x\mathrm{Se_2}$ . The Seebeck coefficient also exhibits metallic behaviour from 100 to 300 K, increasing with temperature linearly according to the equation

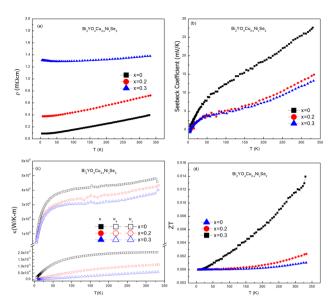


Fig. 4 The thermoelectric properties of Bi<sub>2</sub>YO<sub>4</sub>Cu<sub>2-x</sub>Ni<sub>x</sub>Se<sub>2</sub>

$$S = \frac{8\pi^2 k_{\rm B}^2 Tm^*}{3eh^2} \left(\frac{\pi}{3n}\right)^{2/3}$$

where  $k_{\rm B}$  is the Boltzmann constant,  $m^*$  is the effective mass of carrier, and n is the carrier concentration, which applies for metals or degenerate semiconductors under assumption of parabolic band and acoustic phonon scattering approximation. The Seebeck coefficients are positive for all the samples ranging from 5 to 330 K, indicating p-type carriers dominated in these materials, which is consistent with previous results. The Seebeck coefficient decreases with increasing Ni content. Fig. 3(c) shows the total thermal conductivity curves. The total thermal conductivity decreases with increasing Ni content. The total thermal conductivity can be divided into the electron component ( $\kappa_e$ ) and the lattice component ( $\kappa_l$ ):  $\kappa = \kappa_e + \kappa_l$ . The electron component is proportional to the electrical conductivity according to Wide-Franz law,  $\kappa_e = L\sigma T = LT/\rho$ . The decreasing resistivity indicates the electron component of the total thermal conductivity increase. Then the lattice thermal conductivity decreases with increasing Ni content, which originates from the increasing defect scattering induced by Ni. The dimensionless figure of merit is defined as  $ZT = S^2/\rho \kappa T$ , which is the measure of the efficiency of a thermoelectric material. We calculated the ZT of the synthesized samples, which is shown in Fig. 4(d). Although the induce of Ni de-crease the resistivity of the sample, the ZT decrease with increasing Ni content because of the great degree reduction of the Seebeck coefficient. The introduction of Ni is unfavourable to the improvement of thermoelectric properties in Bi<sub>2</sub>YO<sub>4</sub>Cu<sub>2</sub>Se<sub>2</sub>.

The plane-wave projector augmented method as implement in the Vienna *ab initio* simulation package (VASP) was used to calculate the density of states (DOS) and energy band. The Perdew-Bruke-Ernzerhof (PBE) form of generalized gradient approximation (GGA) was chosen as the exchange-correlation potential.<sup>23</sup> Due to the ex-change-correlation effects of the strongly localized Cu and Ni 3d electrons, we adopted the DFT +

U method. We tested different values of the screened Coulomb parameter U. The added Coulomb interactions do not change the ground state. An energy cut-off of 520 eV was used in all cases. The convergence criterion for energy was set to be 10<sup>-6</sup> eV per unit cell and the forces on all relaxed atoms were less than  $0.01 \text{ eV A}^{-1}$ . According to previous study, nonmagnetic states is adopted to calculate the electronic band structure. Fig. 5 shows the band structure of Bi<sub>2</sub>YO<sub>4</sub>Cu<sub>2</sub>Se<sub>2</sub> and Bi<sub>2</sub>YO<sub>4</sub>Ni<sub>2</sub>Se<sub>2</sub>. The calculated band structure of Bi<sub>2</sub>YO<sub>4</sub>Cu<sub>2</sub>Se<sub>2</sub> is consistent to previous result. The Fermi energy located at the valence band, indicating the p-type metallic ground state, which is consistent with the experiment results. Furthermore, as the replace of Cu by Ni, the Fermi energy move deeper in valence band, namely similar p type metallic ground state. The conduction band is attached to the valence band. The semiconducting behaviour in Ni-doped Bi<sub>2</sub>YO<sub>4</sub>Cu<sub>2</sub>Se<sub>2</sub> observed in our experiment may be non-intrinsic. For Bi<sub>2</sub>YO<sub>4</sub>Cu<sub>2</sub>Se<sub>2</sub>, the bands exhibit complicated character near the Fermi surface. Along the  $Z-\Gamma$  line, the energy dispersion is very weak. The energy dispersion is strong in the layer-parallel directions, which originates from the strong hybridizations between Cu and Se. Similar quasi-twodimensional properties are found in Bi<sub>2</sub>YO<sub>4</sub>Ni<sub>2</sub>Se<sub>2</sub>, which is consistent with KNi2Se2.20 There are also two bands crossing the Fermi energy, indicting the multi-band character. Compared the band structure of Bi<sub>2</sub>YO<sub>4</sub>Ni<sub>2</sub>Se<sub>2</sub> and KNi<sub>2</sub>Se<sub>2</sub>, the structure of the valence band is very similar. Moreover, the thick Bi<sub>2</sub>YO<sub>4</sub> layer make the 2D property more obvious. The ThCr2Si2-type super-conductor KNi<sub>2</sub>Se<sub>2</sub> was discovered with a superconducting transition temperature of 0.8 K. We believe the Bi<sub>2</sub>YO<sub>4</sub>Cu<sub>2</sub>Se<sub>2</sub> may be a superconducting material if it can be successful synthesized.

# Experimental

Bi<sub>2</sub>YO<sub>4</sub>Cu<sub>2-x</sub>Ni<sub>x</sub>Se<sub>2</sub> sample was prepared by reacting a stoichiometric mixture of Bi<sub>2</sub>O<sub>3</sub>, Y<sub>2</sub>O<sub>3</sub>, Y, Cu, Ni and Se. The raw materials weighed by stoichiometric ratio were mixed. After thoroughly grounding in an agate pestle and mortar, the mixtures were pressed into pellets under 12 MPa. The pellets

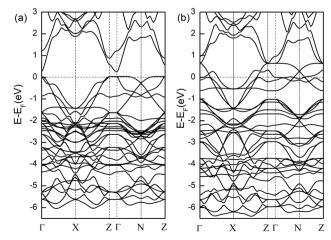


Fig. 5 The band structure of Bi<sub>2</sub>YO<sub>4</sub>Cu<sub>2</sub>Se<sub>2</sub> and Bi<sub>2</sub>YO<sub>4</sub>Ni<sub>2</sub>Se<sub>2</sub>.

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were then sealed under vacuum (<10<sup>-4</sup> Pa) in the silica tubes which had been baked in drybox for 1-2 h at 150 °C. The ampoules were heated to 830 °C and maintained at this temperature for 24 h. Finally, the furnace is shut down and cools to room temperature naturally. The obtained samples were reground, pelletized, and heated for another 24 hours at 830 °C followed by furnace cooling. The X-ray powder diffraction patterns were recorded at room temperature on a PANalytical diffractometer (X'Pert PRO MRD) with Cu Ka radiation (40 kV, 40 mA). Structural refinement of powder samples was carried out by using Rietica software. The electrical resistivity and thermoelectric properties were measured using a Quantum Design PPMS (Physical Properties Measurement System) from 300 to 2 K.

## Conclusions

In summary, we have synthesized a series of layered oxyselenides Bi<sub>2</sub>YO<sub>4</sub>Cu<sub>2-x</sub>Ni<sub>x</sub>Se<sub>2</sub> through solid state reaction. There crystalline structures and physical properties were thoroughly studied. Along with the increase of Ni content, the resistivity increase. It exhibits semiconducting transport behaviour in Bi<sub>2</sub>YO<sub>4</sub>Cu<sub>1.6</sub>Ni<sub>0.4</sub>Se<sub>2</sub>, which may be caused by the diffraction of impurity or defects. The crystal structure and band structure of Bi<sub>2</sub>YO<sub>4</sub>Ni<sub>2</sub>Se<sub>2</sub> are like those of KNi<sub>2</sub>Se<sub>2</sub>. We suspect that Bi<sub>2</sub>YO<sub>4</sub>Ni<sub>2</sub>Se<sub>2</sub> should also exhibit superconductivity if it can be successfully synthesized.

## Author contributions

Conceptualization, L. X. and J. W.; methodology, L. X.; validation, Z. L., W. K. and C. W.; formal analysis, Y. Z.; investigation, S. T.; resources, S. T.; data curation, L. X.; writing—original draft preparation, S. T.; writing—review and editing, L. X.; all authors have read and agreed to the published version of the manuscript.

### Conflicts of interest

There are no conflicts to declare.

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