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# Graphene nanoplatelets (GNPs): a source to bring change in the properties of Co–Ni–Gd-ferrite/GNP nanocomposites

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The nanocomposites of  $\text{Co}_{0.5}\text{Ni}_{0.5}\text{Gd}_{0.03}\text{Fe}_{1.97}\text{O}_4$ /graphene nanoplatelets (CNGF/GNPs) were synthesized by a cost-effective sol–gel auto combustion (SGAC) route. The X-ray diffraction analysis confirmed the cubic structure of the as-prepared nanocomposites, and a crystallite size of 32.28 nm was observed for the 7.5 wt% GNPs. Irregular and unique nanoparticles consisting of short stacks of graphene sheets having a platelet shape were confirmed by the morphological analysis of the as-prepared nanocomposites. Raman analysis revealed a spinel crystal structure along with a new vibrational mode due to the GNPs. The energy bandgap was 3.98 eV for the composite with 7.5 wt% GNP concentration. It was observed that the percentage temperature coefficient of resistance (TCR%) rapidly decreased with an increase in temperature both in low- and high-temperature ranges. Dielectric studies carried out in the frequency range  $10^4$ – $10^7$  Hz confirmed that the graphene-added composites had high values for both the real and imaginary parts of permittivity at low frequencies. A decrease in saturation magnetization with an increase in GNP concentration was observed compared with the pure CNGF samples. Hence, the as-prepared composites are useful for application in high-frequency devices as well as spintronics.

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## 1. Introduction

The incorporation of nano-sized particles into the matrix of a standard material results in a nanocomposite, which exhibits drastic improvement in several properties, such as mechanical strength and magnetic, optical, electrical, and thermal conductivities. An excellent example of the nanocomposites is graphene with soft ferrites, which show tremendous changes and enhancement in properties and hence are successfully used in various fields, such as wastewater treatment, improvement of broadband absorption in gigahertz frequencies and microwave absorption, and energy storage sensors.<sup>1–4</sup> The structure of graphene is capable of modifying the different properties of soft ferrites, such as non-conductivity to conductivity, thereby intensifying their magnetism and improving the structural and mechanical parameters. Moreover, it can also improve the thermal and dielectric properties of the nanocomposites. These

promising structural, electrical, optical, magnetic, mechanical, and thermal properties of graphene have opened new avenues for learning materials science.<sup>5,6</sup> Compared with other nanomaterials, soft ferrites are less expensive, reliably stable, and easy to prepare. Their high permeability at radio frequencies, high Curie temperature, high electrical resistivity, low eddy current loss, fine environmental stability, and mechanical hardness make them suitable for application in many devices, such as microwave apparatus, inductors, choke coils, transformers, and electromagnetic interface (EMI) suppressors. One class of soft ferrites is cubic spinel ferrites (SFs) with  $\text{MFe}_2\text{O}_4$ , where M is a divalent ion, such as  $\text{Mn}^{2+}$ ,  $\text{Ni}^{2+}$ ,  $\text{Fe}^{2+}$ ,  $\text{Co}^{2+}$ , or  $\text{Mg}^{2+}$ . A lot of studies and research have been done on cubic ferrites and many mixed ferrites. Saqib *et al.* studied the structural, electrical, and magnetic properties of Mg-doped  $\text{ZnFe}_2\text{O}_4$  ferrites.<sup>7</sup> Different ferrites are considered unique for different properties; cobalt ferrite has special magnetic properties<sup>8–10</sup> and outstanding chemical and structural stability.<sup>11–13</sup> The substitution of  $\text{Fe}^{3+}$  ions with rare earth metals with larger ionic radii than  $\text{Fe}^{3+}$  into the spinel structure alters the electromagnetic properties of ferrites. Consequently, the inconsistent behavior of doped nickel ferrites is observed.<sup>14–16</sup>

The preparation techniques used for cubic ferrites with different concentrations and the distribution of different cations at the tetrahedral (A-) and octahedral (B-) sites of the sample lattice play a vital role in tuning the properties. Thus,

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any addition and alteration in the distribution and ratio of these cations can change their magnetic behavior. Satalkar *et al.* prepared  $\text{Zn}_{0.7-x}\text{Ni}_x\text{Mg}_{0.2}\text{Cu}_{0.1}\text{Fe}_2\text{O}_4$  ( $x = 0.0\text{--}0.7$ ) ferrites by a sol-gel auto-combustion (SGAC) method. They found that  $\text{Ni}^{2+}$  addition plays a part in determining the magnetic properties and changes the cationic distribution.<sup>17</sup> Therefore, the usage of suitable additives plays a vital role in gaining high-performance ferrites.

Tremendous work has been done on the addition of rare earth metals to soft ferrites at varying concentrations by using a variety of mechanical and chemical procedures, such as sol-gel, hydrothermal, and traditional ceramic technology. Elements of the lanthanide series, such as gadolinium, neodymium, cerium, and yttrium, are added to different ferrites due to their special characteristics. These elements are frequently used with ferrite combinations at different concentrations in many applications as they enhance optical, structural, magnetic, electrical, and morphological properties. Pawar *et al.* reported that  $\text{Gd}^{3+}$  substitution in the Co-Zn ferrite increases the saturation magnetization and coercivity. Consequently, these  $\text{Gd}^{3+}$ -substituted materials are appropriate for use in permanent magnets and magnetic recording media.<sup>18</sup> The properties of cubic ferrites are enhanced when rare earth metals with higher ionic radii are substituted for  $\text{Fe}^{3+}$ . The concentration of rare earth metals plays a pivotal role in determining the properties of the cubic ferrites.<sup>19</sup> Poudel *et al.* studied that the concentration of  $\text{Gd}^{3+}$  in Ni ferrite enriched its properties.<sup>16</sup> Hossain, M. D., *et al.* studied the properties tailored by  $\text{Gd}^{3+}$  substitution at the B sites, replacing the  $\text{Fe}^{3+}$  ions of the Ni-Zn-Co ferrite. Due to substitution with gadolinium, they found changes in the structural properties, as well as electrical properties.<sup>20</sup> It is worth mentioning that the content of graphene plays an important, leading role in defining the dielectric behavior of carbon materials along with their magnetic anisotropy. This subsequently would impact the microwave weakening behavior.

In this work, we followed the sol-gel auto-combustion route to prepare Ni-doped cobalt ferrites with a stoichiometric formula of  $\text{Co}_{0.5}\text{Ni}_{0.5}\text{Gd}_{0.03}\text{Fe}_{1.97}\text{O}_4$  (CNGF) and loaded them on the surface of graphene nanoplatelets (GNP) using a sonochemical approach. Therefore, this work aims to investigate the impact of graphene nanoplatelets (GNPs) on the spectral, electrical, magnetic, and dielectric properties of the CNGF composites prepared *via* the sol-gel auto combustion (SGAC) process.

## 2. Materials and methods

### 2.1 Materials

The following chemicals were used for the synthesis of the  $\text{Co}_{0.5}\text{Ni}_{0.5}\text{Gd}_{0.03}\text{Fe}_{1.97}\text{O}_4$  (CNGF) sample and the CNGF/GNP composites. Iron nitrate nonahydrate  $[(\text{Fe}(\text{NO}_3)_2 \cdot 9\text{H}_2\text{O})]$ , nickel nitrate hexahydrate  $[(\text{Ni}(\text{NO}_3)_2 \cdot 6\text{H}_2\text{O})]$ , aqueous ammonia  $[(\text{NH}_4\text{OH})]$ , cobalt nitrate hexahydrate  $[(\text{Co}(\text{NO}_3)_2 \cdot 6\text{H}_2\text{O})]$ , gadolinium nitrate hexahydrate  $[\text{Gd}(\text{NO}_3)_2 \cdot 6\text{H}_2\text{O}]$ , citric acid ( $\text{C}_6\text{H}_8\text{O}_7$ ), and graphene nanoplatelets (GNP).

### 2.2 Synthesis of CNGF sample

The CHNF sample was prepared by the sol-gel auto-combustion (SGAC) method. All chemicals at proper stoichiometric proportions were sonicated individually for 10 min at 353 K to make their homogenous solutions in distilled water. All prepared solutions were mixed again by using a magnetic stirrer for up to 15 minutes until well combined, followed by the addition of the citric acid solution.  $\text{NH}_4\text{OH}$  was added dropwise to maintain the pH of the solution at 7 and stirring was further continued for 20 minutes. The samples were heated at 523 K until a gel was formed, after which the magnetic stirrer was turned off while heating was continued, which further led to self-combustion, and the solution turned into ash. It was mechanically ground by using a mortar and pestle to make a fine powder, which was calcined at 1123 K for 5 h. The calcined samples were again ground to obtain their powder form.

### 2.3 Synthesis of CNGF/GNP nanocomposites

The sonochemical approach was used to prepare the CNGF/GNP nanocomposites by keeping the stoichiometric quantities of GNP at 2.5 wt%, 5 wt%, and 7.5 wt% of the total metal nitrate salts. The required amount of GNP was dissolved in distilled water, and the mixture was sonicated for half an hour to get a homogeneous dispersion. All the metal nitrate salts dissolved in distilled water were separately added in dispersed GNP, and each mixture was again sonicated for 20 minutes. A calculated amount of citric acid dissolved in distilled water was added to the metal nitrate salts -GNP mixture and sonicated further for 10 minutes. This mixture was placed on a magnetic stirrer for thorough mixing, and this was followed by the SGAC method to get ash, which was ground into a fine powder and then calcined at 1123 K for 5 h. Fig. 1 illustrates the step-by-step fabrication process.

### 2.4 Characterization techniques used

The structural properties and phase confirmation of the prepared samples were studied by X-ray diffraction (XRD) analysis using a Bruker D8 Advanced X-ray diffractometer (XRD) using Cu-K $\alpha$  source with a wavelength of 1.54 (Å). To inspect the surface morphology, scanning electron microscopic (SEM) images were captured using a Nova Nano SEM series 450. Raman analysis was used to confirm the lattice structure of the ferrites along with  $\text{Gd}^{3+}$  substitution of and addition of GNP by using laser excitation at 532 and 780 nm in the range of 200–1800  $\text{cm}^{-1}$ . The energy bandgap ( $E_g$ ) of the as-prepared samples was determined through UV-visible spectroscopy. Current-voltage ( $I$ - $V$ ) measurements and hence resistivity were calculated by using the two-probe technique on an electric meter KEITHLEY Model 2401. The dielectric behavior was studied using an LCR Meter (IM3533 series) in the frequency range of 4 Hz to 8 MHz at room temperature. The magnetic properties were investigated by studying the  $M$ - $H$  loop at 303 K by using a vibrating sample magnetometer (VSM-175).



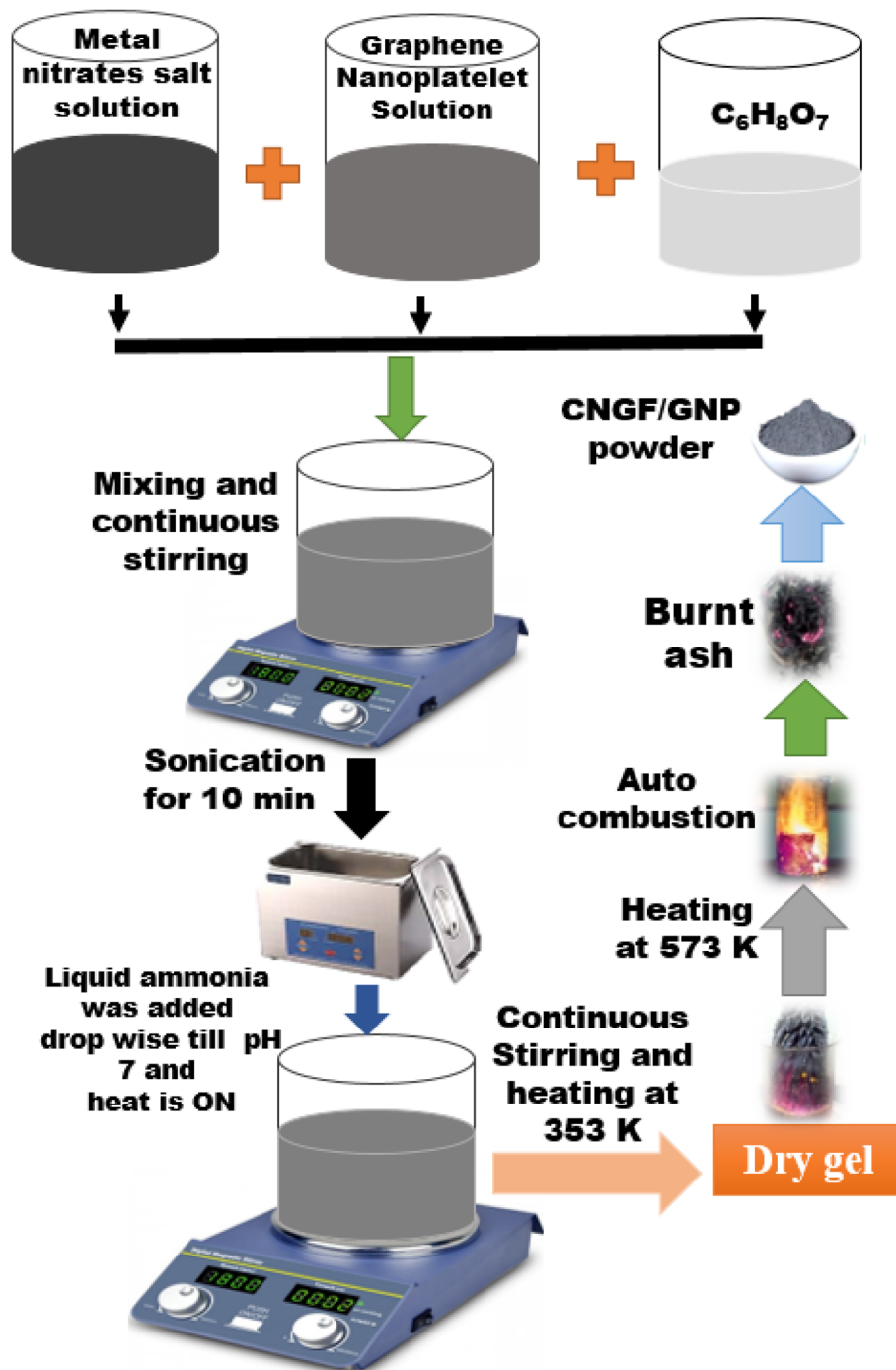


Fig. 1 A schematic representation of the synthesis of NCGF/GNP nanocomposites.

### 3. Results and discussion

#### 3.1 Structural properties

To discuss the effect of GNP on the structure of the CNGF nanoparticles, XRD analysis and Raman spectroscopy were used, while the effect of GNP on the surface morphology was studied using the micrographs obtained from SEM.

**3.1.1 X-Ray diffraction analysis.** XRD spectra were used to understand the structural changes in the CNGF crystals after the addition of GNP, as shown in Fig. 2. GNP showed diffraction

peaks (as shown in Fig. 2) at a  $2\theta$  value of  $26.38^\circ$ , which denotes the (002) crystal plane.<sup>21</sup> In addition to that, the CNGF/GNP nanocomposites revealed diffraction peaks at  $2\theta$  of  $30.2^\circ$ ,  $35.6^\circ$ ,  $37.2^\circ$ ,  $43.3^\circ$ ,  $53.3^\circ$  and  $57.3^\circ$ , corresponding to the (220), (311), (222), (400), (422), and (511) planes, respectively.<sup>22</sup> Further, using Braggs equation:  $2d \sin \theta = n\lambda$ , the  $d$ -spacing for pure GNP calculated based on  $2\theta = 26.8^\circ$  corresponding to the (002) plane was  $3.38 \text{ \AA}$ , which is larger than the values of the CNGF nanoparticles and CNGF/GNP nanocomposites, as shown in Table 1. Being a cubic crystal structure, the three lattice



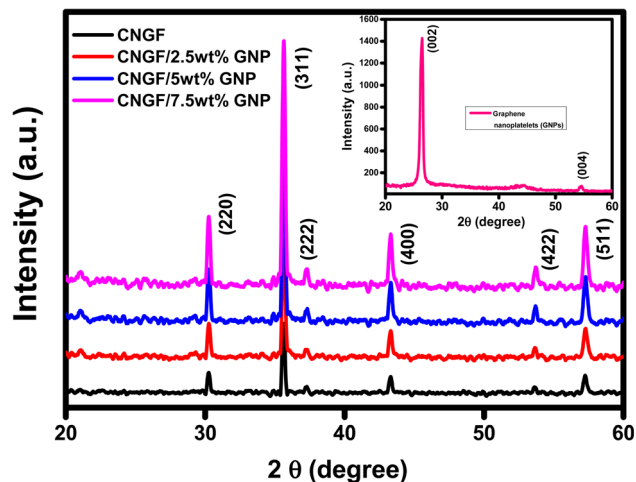


Fig. 2 The XRD spectra of the CNGF nanoparticles and CNGF/GNP nanocomposites.

constants are the same  $a = b = c$ , and their value was calculated using the formula;  $a = d_{hkl}[h^2 + k^2 + l^2]^{1/2}$ , where  $h$ ,  $k$ , and  $l$  Miller indices of the high-intensity peak (311) were used to calculate the value of the lattice constants as listed in Table 1. The values of unit cell volume ( $V = a^3$ ) were calculated by using the value of the lattice constant. The lattice constant and unit cell volume of the CNGF/GNP nanocomposites did not change significantly. This may be due to the incorporation of GNP in the CNGF nanoparticle spinel lattice. The crystalline size ( $D$ ) of the as-synthesized nanocomposites was calculated using the Scherrer equation;  $D = 0.94\lambda/\beta \cos \theta$ , where  $\beta$  is the full width at half maxima (FWHM) of the highest intensity peak of the samples,  $\lambda = 1.5406 \text{ \AA}$  is the wavelength of the used  $\text{Cu K}\alpha$  radiation source, and  $\theta$  represents the Bragg's reflection angle. The smaller crystallite size of the CNGF/GNP nanocomposites compared with the pure CNGF sample may be due to the incorporation of GNP in the CNGF lattice (as shown in Table 1). The dislocation density ( $\delta = 1/D^2$ ) is a measure of the number of dislocations per unit volume of a crystalline material and follows the opposite trend to crystallite size *i.e.* when the crystallite size of the material is the maximum, the dislocation line density is minimum and *vice versa*. The crystallite size was minimum for the CNGF/2.5wt% GNP nanocomposite and it had the maximum dislocation line density, as shown in Table 1. It is clear from Table 1 that the packing factor ( $p$ ) also was smaller for the CNGF/GNP nanocomposites compared with the CNGF nanoparticles. The increase in the values of microstrain ( $\epsilon$ ) of

the CNGF/GNP nanocomposites to that of the CNGF nanoparticles, along with a deteriorating trend of crystalline size, suggests their enhanced electrocatalytic activity.<sup>23</sup>

**3.1.2 Scanning electron microscopy analysis.** SEM images were used to analyze the morphology of the CNGF nanoparticles and CNGF/GNP nanocomposites, as shown in Fig. 3(a–d). Irregular particle morphology was observed in the nanoparticles, as seen in Fig. 3(a). The reported morphology of GNP is flake-like; moreover, graphene nanoplatelets are unique nanoparticles consisting of short stacks of graphene sheets having a platelet shape.<sup>24</sup> The SEM images of the as-prepared CNGF/GNP nanocomposites in Fig. 3(a–d) indicate that the GNP was densely covered with the CNGF nanoparticles. This indicates that the mechanism of electrostatic attraction between the nanoparticles and graphene<sup>25</sup> gives rise to the agglomeration of nanoparticles in the flake-like structure of GNP due to weak van der Waals, electrostatic, and capillary forces, which are responsible for producing magnetic interactions between the particles. The SEM images showed that samples with graphene were thicker than CNGF, which is due to the presence of different functional groups like carboxyl and carbonyl and some traces of oxygen, while, in the absence of graphene, the nanoparticles showed scattered behavior (as shown in Fig. 3(a)).<sup>24</sup>

**3.1.3 Raman spectroscopy analysis.** Fig. 4 reveals the Raman spectra of pure GNP, the CNGF sample, and the CNGF/GNP composites. Pure GNP showed two vibrational bands around  $1343 \text{ cm}^{-1}$  and  $1621 \text{ cm}^{-1}$  corresponding to the D-band and G-band, respectively. The characteristic peaks at  $689.1 \text{ cm}^{-1}$ ,  $545.6 \text{ cm}^{-1}$ ,  $471.6 \text{ cm}^{-1}$ ,  $316.8 \text{ cm}^{-1}$ , and  $370.1 \text{ cm}^{-1}$  of the SFs corresponded with five active Raman modes  $A_{1g}(R)$ ,  $3 T_{2g}(R)$  and  $E_g$  respectively, where A, E, and T denote one-, two- and three-dimensional display of the SFs, respectively, and g denotes symmetry at the center of inversion.<sup>26</sup> Raman spectra of the samples containing GNP at concentrations 2.5%, 5%, and 7.5% are shown in Fig. 4, showing slight changes in the first five active modes: the  $A_{1g}(1)$  mode of vibration at around  $689.1 \text{ cm}^{-1}$  changed to  $686.6 \text{ cm}^{-1}$  due to GNP addition due to the symmetric stretching of the oxygen atoms between Fe–O and metal ions at the tetrahedral (A-) site.  $T_{2g}(3)$ , which is assigned to the asymmetric bending of oxygen at the octahedral (B-) site relative to Fe, changed from  $545.6$  to  $543.3 \text{ cm}^{-1}$  due to the insertion of GNP.  $T_{2g}(2)$ , which is responsible for the asymmetric stretching of the Fe–O bond at the B-site, was originally found at  $471.6 \text{ cm}^{-1}$  but shifted to  $467.7 \text{ cm}^{-1}$  due to the addition of GNP. A low-intensity peak at  $370.2 \text{ cm}^{-1}$  representing the  $E_g$  mode, which is responsible for

Table 1 The structural parameters of the CNGF nanoparticles and CNGF/GNP nanocomposites

Samples	$a$ (Å)	$D$ (nm)	$V$ (Å) <sup>3</sup>	$d$ (Å)	$\delta \times 10^{-4}$ (nm <sup>-2</sup> )	$\epsilon$	$p$
CNGF	8.345	40.18	581.21	2.518	6.193	0.00294	159.69
CNGF/2.5wt%GNP	8.335	28.76	579.18	2.515	12.131	0.00411	114.44
CNGF/5wt%GNP	8.343	36.38	580.88	2.459	7.552	0.00325	144.64
CNGF/7.5wt%GNP	8.344	32.80	581.05	2.518	9.292	0.00361	130.38





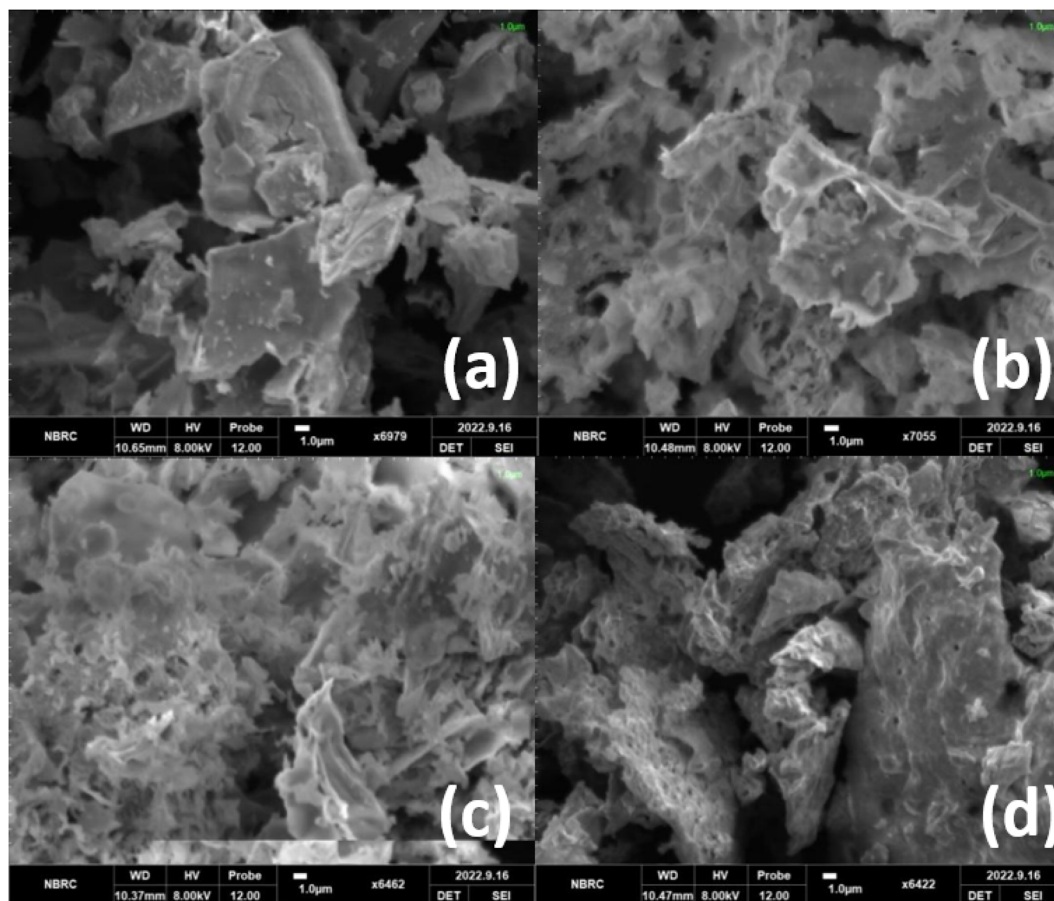


Fig. 3 The SEM images of the (a) CNGF nanoparticles, (b) CNGF/2.5wt%GNP nanocomposite, (c) CNGF/5wt%GNP nanocomposite, and (d) CNGF/7.5wt%GNP nanocomposite.

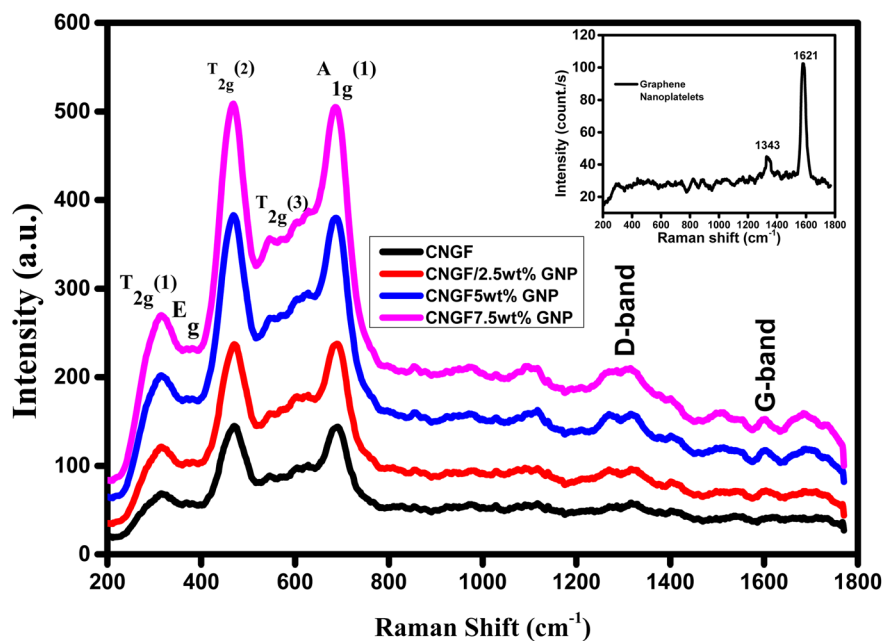


Fig. 4 The Raman spectra of GNP (inset), CNGF, and CNGF/GNP composites.



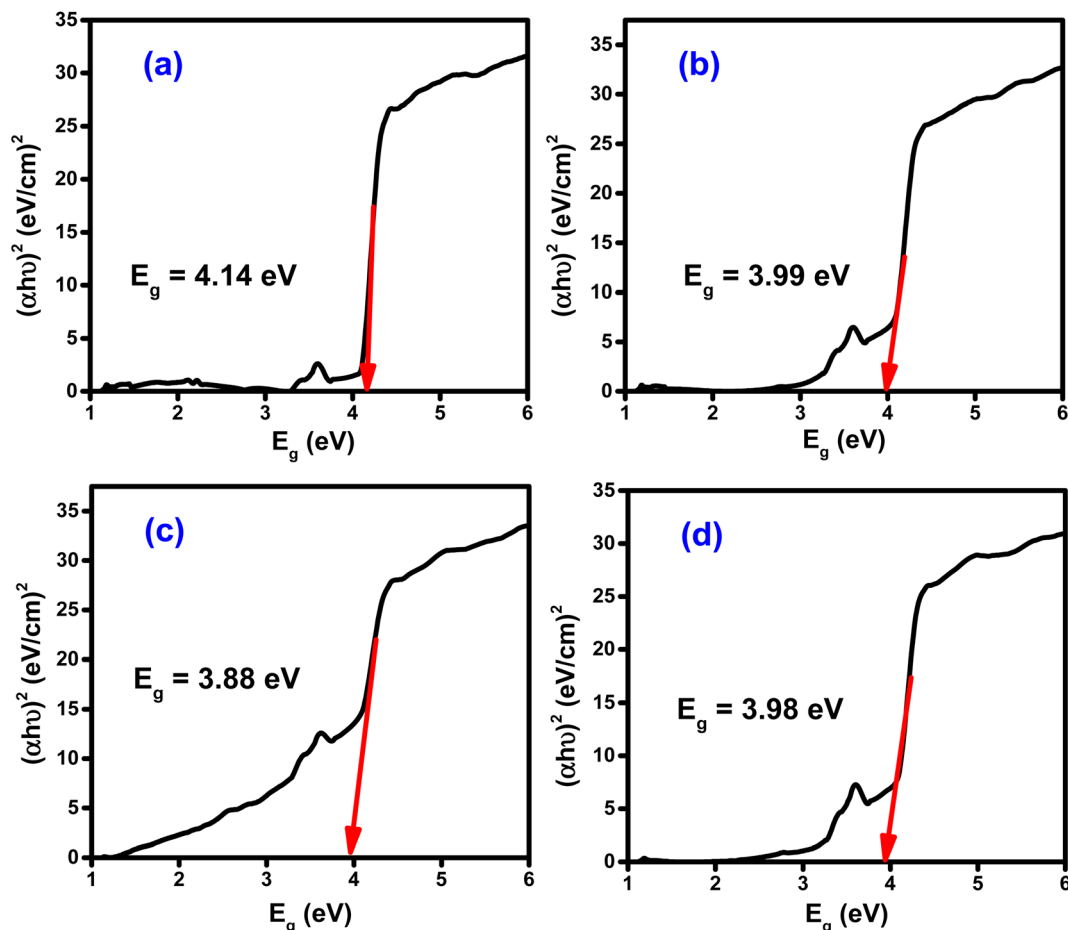


Fig. 5 (a–d) The Tauc plots for the CNFG sample and CNFG/GNP composites.

the symmetric bending of the oxygen ions with respect to iron, moved to  $365.8\text{ cm}^{-1}$  in the CNFG/GNP composites. The  $T_{2g}$  (1) peak assigned to the translational motion of the complete tetrahedron ( $\text{Fe-O}_4$ ) was present at  $316.8\text{ cm}^{-1}$  and moved to  $314.5\text{ cm}^{-1}$ .<sup>27,28</sup>

### 3.2 Optical studies by UV-vis analysis

The optical transition nature and energy bandgap of the as-prepared CNFG sample and CNFG/GNP composites were measured by using the Tauc plot method. SFs and their composites have a direct energy bandgap.<sup>29</sup> The Tauc plot equation is given as:<sup>26</sup>  $(\alpha h\nu)^2 = K(h\nu - E_g)$ , where “2” is for direct allowed transitions, the absorption coefficient ( $\alpha$ ) was determined by employing the following equation:<sup>26</sup>  $\alpha = 2.303 \log(A)/l$ , where  $A$  indicates the absorbance, and  $l$  is the path length of light along which absorbance takes place, and the energy bandgap ( $E_g$ ) was calculated by extrapolating the linear part of

the curves in Fig. 5(a–d). The  $E_g$  values were 4.14 eV, 3.99 eV, 3.88 eV, and 3.94 eV for CNFG, CNFG/2.5wt% GNP, CNFG/5wt% GNP, and CNFG/7.5wt% GNP, respectively. The  $E_g$  of the CNFG sample is the widest due to the interaction between the  $\text{O}^{2-}$  p and Fe-3d energy states, resulting in orbital overlap, while the change in energy bandgap is due to the insertion of GNP in the host CNFG sample.

### 3.3 I–V measurements analysis

The two-probe method was used to determine the DC resistivity of the CNFG nanoparticles and CNFG-GNP (2.5%, 5%, and 7.5%) nanocomposites. The resultant measured resistivity values of each given sample are summarized in Table 2. The calculated resistivity of the integrated CNFG nanoparticle was  $2.44 \times 10^9\ \Omega\text{ cm}$  at room temperature. The general behavior of ferrite nanoparticles is more insulating, but they can conduct due to the temperature dependency of the electrons, which

Table 2 The DC electrical resistance and resistivity of the CNFG nanoparticles and CNFG-GNP nanocomposites

Samples	CNFG	CNFG-GNP (2.5%)	CNFG-GNP (5%)	CNFG-GNP (7.5%)
Resistance ( $\Omega$ )	$1.36 \times 10^9$	$5.85 \times 10^7$	$1.23 \times 10^8$	$4.41 \times 10^7$
Resistivity ( $\Omega\text{ cm}$ )	$2.44 \times 10^9$	$1.05 \times 10^8$	$2.21 \times 10^8$	$7.89 \times 10^7$



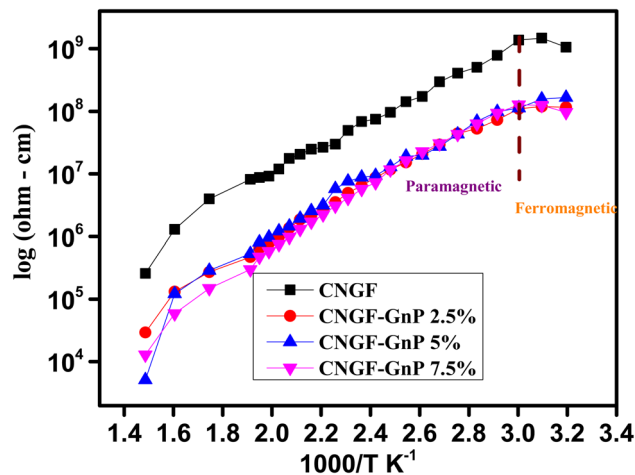


Fig. 6 Arrhenius plots of all the samples.

causes the transfer of the electrons between the ferrous ( $\text{Fe}^{2+}$ ) and ferric ( $\text{Fe}^{3+}$ ) ions at the octahedral  $B$ -lattice sites and transfer of holes between the  $\text{Ni}^{2+}$  and  $\text{Ni}^{3+}$  ions.<sup>30</sup> The  $\log \rho$  vs.  $1000/T$  (Arrhenius plots) graphs are shown in Fig. 6 at different temperatures ranging from 313 K to 673 K. Based on the variation in the slope of the curve, the resistivity curve can be split into two regions, namely the ferromagnetic region and paramagnetic region. The change in resistivity is in line with

a specific temperature at which the log values show a transformation; this is described as the Curie temperature or the transition temperature and is denoted as ( $T_C$ ). This process or mechanism also causes the transformation of magnetic collation from ferromagnetic to paramagnetic, thus forming two distinct regions in the plots. It is worth noting here that the region below  $T_C$  is called the ferromagnetic region, whilst the region above the  $T_C$  is referred to as the paramagnetic region.<sup>31</sup>

On the contrary, the resistivity of graphene is poor, and hence it is found to be a good conductor with a pristine surface containing encapsulated  $\text{sp}^2$ -hybridized carbon atoms inside the hexagonal lattice with certain constraints. This is because of the availability of functional groups like ethers, carboxyl, and hydroxyls, which can react with atmospheric humidity to form acids and other compounds. It is worth noting here that the oxidation level has been shown to directly resist the electrical properties of graphene nanoplatelets due to the quasi-ID active bonding between the carbon sheets with a conducting nature and the defects, according to the variable range hopping model.<sup>32</sup> The temperature dependency of the resistivity of the fabricated samples is shown in Fig. 7(a–c), which reveals that in the low-temperature range, the resistivity of CNGF increases with the increase in temperature, but in the medium- and high-temperature ranges, resistivity increases when temperature decreases, confirming the semiconducting nature of the gadolinium-doped ferrites. The same trend is found when GNPs

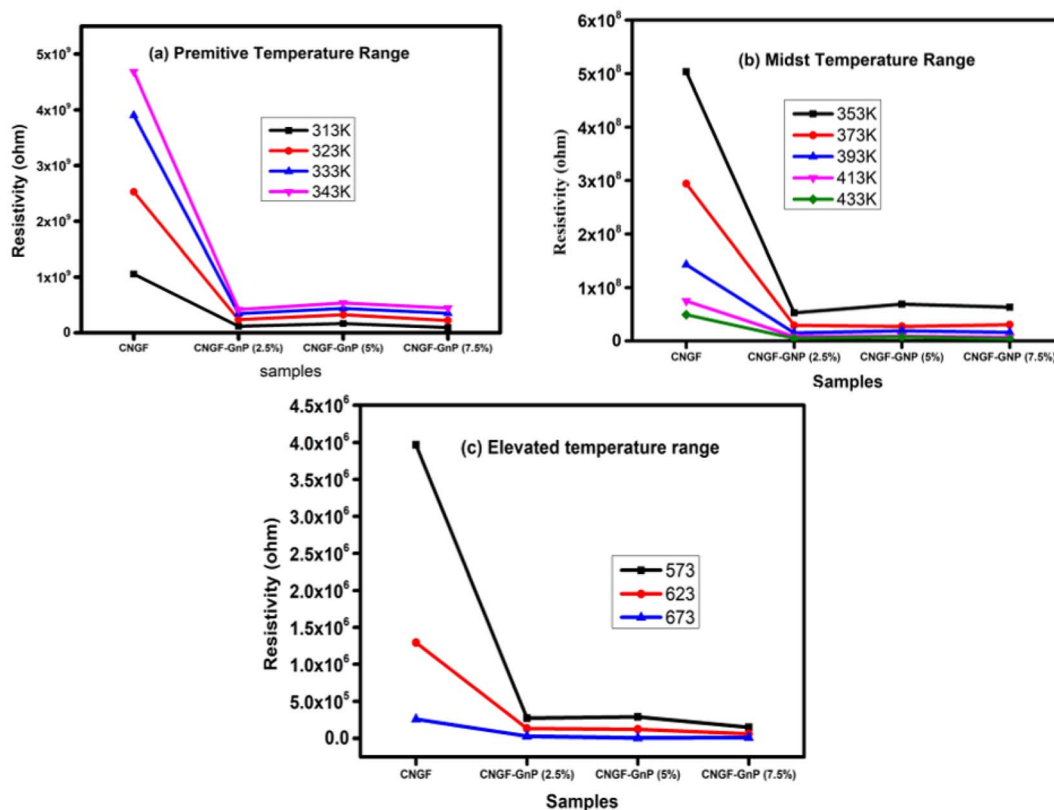


Fig. 7 The temperature-dependent DC resistivity ( $\rho_{DC}$ ) of the CNGF nano particles and CNGF-graphene nanoplatelet composites in the (a) lower temperature range (b) medium temperature range and (c) elevated temperature range.



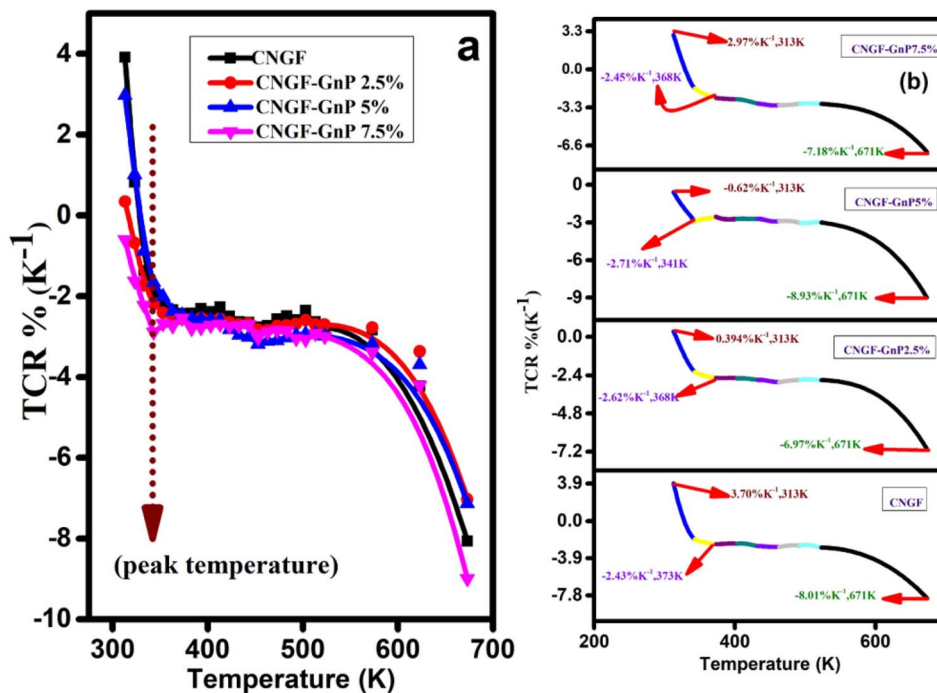


Fig. 8 (a) Temperature versus the temperature coefficient of resistance percentage for all the samples; (b)  $T$  versus TCR% for each sample.

are added to CNGF but with less diversity.<sup>20</sup> These reveal that lesser ferrite particles could get into the honeycomb or hexagonal structures of GNP, causing a rise in the density of functional groups that act as charge carriers. Therefore, the conductivity of graphene remains supreme, and the resistivity of the composites deteriorates with different concentrations of GNP, as shown in Fig. 7(a–c). The resistivity of the CNGF-GNP nanocomposites is less compared with the gadolinium-doped cobalt nickel ferrite nanoparticles. It is, therefore, comprehended that the CNGF-GNP composites have changed the resistivity of the ferrite nanoparticles, and they can be used in the preparation of supercapacitors.<sup>33</sup> TCR% is the percentage of the relative change in resistivity per degree temperature change. TCR is a valuable resistor performance stability indicator. TCR% is essential for estimating the relationship of resistivity with ambient temperature variations, both at low and high operating temperatures. TCR% is calculated from resistivity using the formula:<sup>42</sup>  $\text{TCR} (\%) = [(1/\rho) \times (d\rho/dT)] \times 100\%$ . The TCR% values were calculated for all samples from 313 K to 673 K and plotted against temperature, as shown in Fig. 8(a). The TCR% decreased rapidly with an increase in temperature both in the low-temperature range and high-temperature range, as shown in Fig. 8(b). The CNGF sample without graphene nanoplatelets showed the maximum value of  $3.70 \text{ K}^{-1}$  at a temperature of 313 K, while the composites with graphene nanoplatelets showed lesser values at 313 K due to the carbon in graphene. The peak values of TCR% were  $-2.43\% \text{ K}^{-1}$  at 373 K,  $-2.62\% \text{ K}^{-1}$  at 368 K,  $-2.71\% \text{ K}^{-1}$  at 341 K, and  $-2.45\% \text{ K}^{-1}$  at 368 K for CNGF, CNGF-GNP 2.5%, CNGF-GNP 5% and CNGF-GNP 7.5% respectively. Materials with temperature-dependent electrical resistance can be used as a source of radiant heat and are used

in many applications, such as bolometric devices for microwave measurements, astronomy, and thermal cameras.<sup>27</sup>

### 3.4 Dielectric studies

The term permittivity or dielectric constant refers to the measure of capacitance observed during the formation of an electric field.<sup>34</sup> The electronic distribution of the synthesized samples was predicted by measuring their polarizations through dielectric analysis.<sup>20</sup> The values of dielectric constant ( $\epsilon'$ ) and dielectric constant loss ( $\epsilon''$ ) were plotted as functions log of frequency ( $10^4 \text{ Hz}$  to  $10^7 \text{ Hz}$ ) at room temperature, as shown in Fig. 9(a) and (b), respectively. The dielectric constant and loss initially decreased due to the reduction in the space charge polarization effect. Then, it remained approximately constant. On the other hand, a high dielectric constant was observed at low frequencies. This is because a potential barrier is created at the grain boundaries by space charge polarization. Subsequently, there is an accumulation of charge at the grain boundary, resulting in an increase in the real and imaginary parts of the permittivity values. The dielectric dispersion can be explained by the dominance of the impact of grain boundaries rather than the grains. Koop's phenomenological theory attributes this to the Maxwell–Wagner kind of interfacial polarization. Both  $\epsilon'$  and  $\epsilon''$  show maximum values when the frequency is low because of interfacial and dipole polarization.<sup>35</sup> The CNGF nanoparticles showed lower values of dielectric constant and dielectric constant loss, while the CNGF/GNP nanocomposites showed higher values of both at low frequencies, proving the graphene and ferrite interactions and the presence of better polarization due to graphene. Moreover, the CNGF/5wt%GNP nanocomposite exhibited higher values of the real





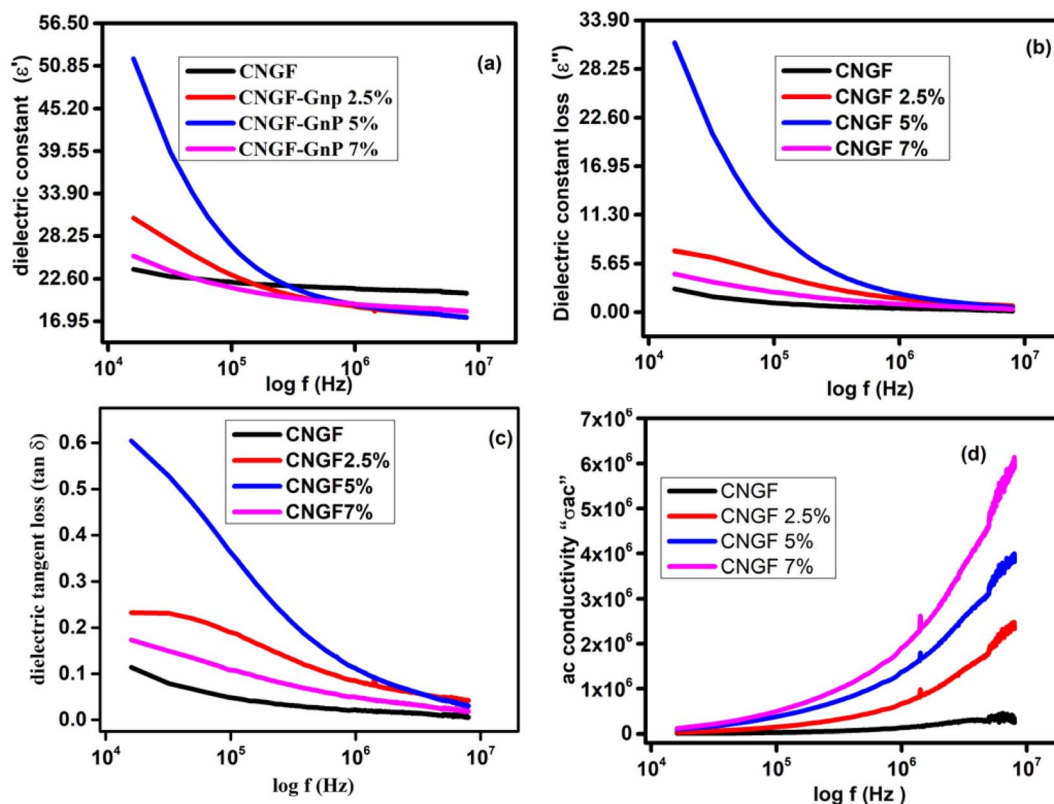


Fig. 9 The plots of (a)  $\epsilon'$  versus frequency, (b)  $\epsilon''$  versus frequency, (c)  $\tan \delta$  versus frequency, and (d)  $\sigma_{ac}$  versus frequency.

and imaginary parts of the dielectric constant than those of the CNGF nanoparticles. As the frequency was increased, the values of  $\epsilon'$  and  $\epsilon''$  remained constant for all the samples. At higher frequencies, due to fast variation in the applied field, the dipoles cannot follow the exact variation speedily, and the response of the permanent dipoles along with the contribution of charge carriers decreases, leading to a decrease in the dielectric constants. Graphene enhances the dielectric constant at lower frequencies, making it a significant factor in the development of high-temperature ceramic supercapacitors.<sup>36</sup> The values of dielectric constant and dielectric constant loss increased when graphene nanoparticles were added.<sup>37</sup> In CNGF, the net polarization is due to two types of charges: (1) due to electron transfer at the octahedral site between  $\text{Fe}^{2+}$  and  $\text{Fe}^{3+}$ , and (2) due to p-type charge carriers, namely  $\text{Ni}^{3+}/\text{Ni}^{2+}$  and  $\text{Co}^{3+}/\text{Co}^{2+}$ , at the B-site. At higher frequencies, the dielectric constant is not dependent on the frequency, so polarization remains unchanged at these frequencies; this shows that polarization depends on the variation in the frequency range. Moreover, due to the insertion of doped ferrites into graphene nanoplatelets, an interface is produced, and charges accumulate at this interface with different relaxation times when current is passed through the sample, which can be well-described by the Maxwell-Wagner (MW) effect.<sup>38</sup>

The dielectric loss tangent ( $\tan \delta = \epsilon''/\epsilon'$ ) illustrates its dependence on the energy dissipation frequency.<sup>39</sup> Energy dissipation in the samples due to conduction (electrical) and

resonance occurring at the domain wall is described by calculating the dielectric loss tangent. The plot of  $\tan \delta$  and frequency is shown in Fig. 9(c), indicating its dependency on the frequency *i.e.*, at a higher frequency, the dielectric loss tangent has the lowest value.<sup>32</sup> One of the characteristics of nanocomposites is that, with a change in frequency, the *ac* conductivity changes, as noticed in Fig. 9(d), and this behavior can be well-explained by Koop's theory, according to which conductivity shows variation/dispersion in the higher frequency range ( $>10^5$  Hz). At room temperature, transitions between  $\text{Fe}^{2+}$  to  $\text{Fe}^{3+}$  at the B-site increase (a hopping mechanism) and lessen the effect of oxygen ions, thus forming well-conducting grains and poor-conducting grain boundaries.<sup>40</sup> Firstly, the presence of multiple interfaces between the graphene nanoplatelets and the modified magnetic materials leads to increased interfacial polarization and the subsequent relaxation process. This phenomenon contributes to the enhancement of dielectric loss. Moreover, recent research findings indicate that charges can be transferred across the interface of a graphene-ferrite heterostructure due to the distinct physical properties of these materials. Hence, in this study, it is reasonable to expect a charge transfer process to take place at the interface between the GNPs and the CNGF sample. This would result in the introduction of free carriers from ferrite into graphene, leading to electric polarization in the as-prepared composites. This process significantly contributes to the increase in dielectric loss.<sup>41</sup>



The electrical performance of the synthesized sample was calculated by studying the  $Q$  factor as a function of frequency. As the  $Q$  factor reveals the association of stored energy with waste (thermal) energy, it is useful to describe the efficiency of the material in terms of storing energy. The  $Q$  factor *versus* frequency graphs of all samples are shown in Fig. 10, which clearly demonstrate that the samples with graphene nanoplatelets have lower values of  $Q$ -factor than CNGF. The  $Q$  factor values increase as the frequency increases due to the AC frequency and skin effect.<sup>42</sup>

The samples were also analyzed by studying complex impedance, which in turn constitutes a real  $Z'$  and an imaginary part  $Z''$  and measures the opposition caused by the material when current is passed through it. Fig. 11 shows the variation in  $Z'$  with frequency for all four samples: without (CNGF) and with graphene (CNGF-GNP 2.5%, 5%, and 7%). All the curves show high impedance at low frequencies; the values decrease when the frequency is increased up to  $10^7$  Hz, and then the impedance remains at a constant value even when the frequency is increased further. The high value of impedance at low frequencies is due to high dipole polarization and interfacial polarization, which in turn are due to the grain boundaries present in the crystal structure and electrode contribution, respectively. CNGF had the maximum impedance, while the CNGF-graphene nanocomposites showed lesser values, which reveals that the crystal grain boundary defects are lessened due to the addition of graphene. Further, the XRD results prove the reduction in crystal size due to graphene, and hence the CNGF-GNP nanocomposites have low impedance.<sup>43</sup>

When the frequency is increased, the effect of dipole polarization and space charge participation reduces due to their inability to respond at these high frequencies; thus, impedance decreases, leading to a constant value at which conduction becomes prominent due to the hopping of electrons. The imaginary part of the impedance is plotted against frequency for all the prepared samples in Fig. 12, which shows high values at low frequencies, a decrease in value as the frequency rises,

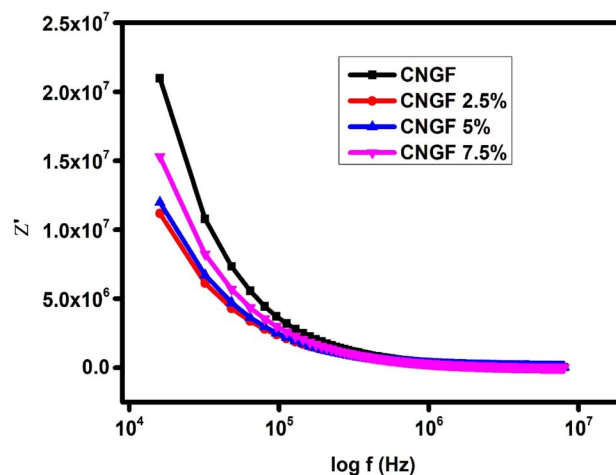


Fig. 11 The variation of the real part of complex impedance with frequency.

and then reaching stability at high values of frequency; this is in agreement with the fact that polarization is compatible at low frequencies but lags at high frequencies, and thus due to the transfer of electrons, conduction is possible at high frequencies.

The physical interpretation of the imaginary part is represented as the energy storage component, and the results revealed that graphene-added composites had higher values for the imaginary part of impedance than the bare spinel ferrite. The relationship between the imaginary parts and real parts of impedance was plotted based on the Nyquist theorem. All the curves showed incomplete semi-circles, thus demonstrating dependence on one dielectric relaxation (Fig. 13(a–d)).

### 3.5 Magnetic studies

Experimental evidence demonstrates that pure CNGF exhibits characteristic soft ferromagnetic behavior, accompanied by a small coercivity.<sup>44</sup> To investigate the effect of GNP on the magnetic properties of the CNGF sample, the magnetization

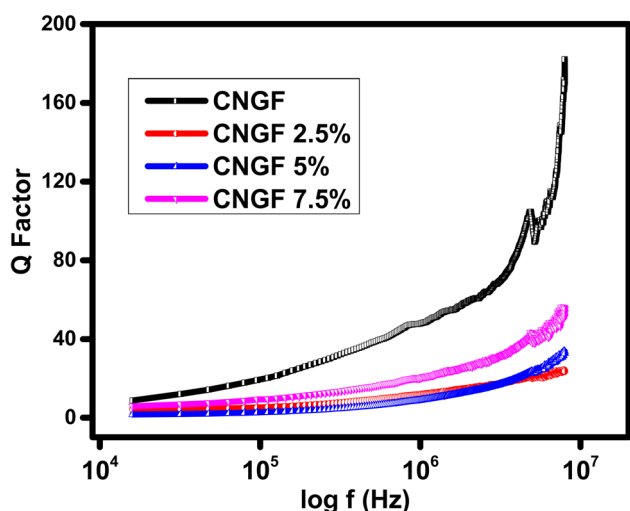


Fig. 10  $Q$  factor variation *versus* frequency.

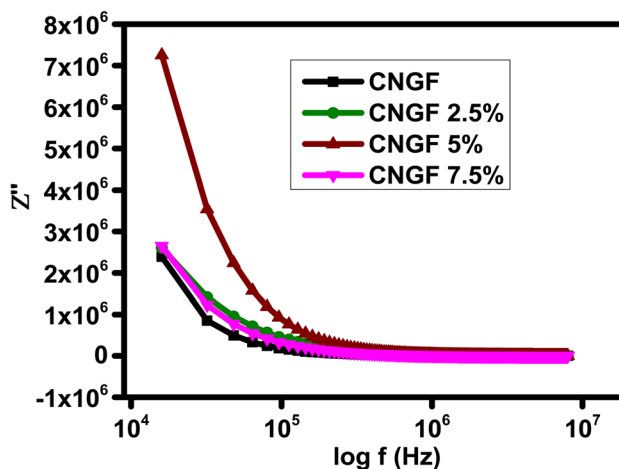


Fig. 12 The variation of the imaginary part of complex impedance with frequency.

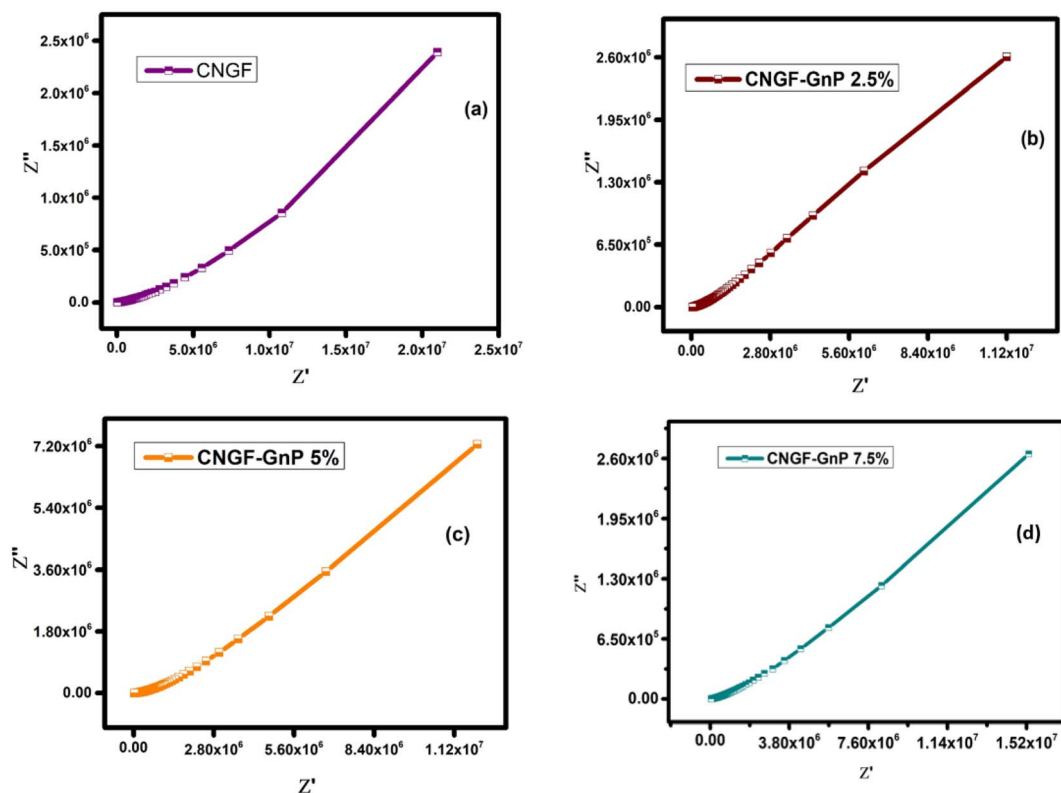


Fig. 13 Nyquist plots showing the variations between the imaginary and real parts of the impedance of the (a) CNGF nanoparticles, (b) CNGF-GnP 2.5% composites, (c) CNGF-GnP 5% composites, and (d) CNGF-GnP 7.5% composites.

versus applied field ( $M-H$ ) loops were recorded at room temperature, as shown in Fig. 14. It is clear from Fig. 14 that the saturation magnetization ( $M_s$ ) of the CNGF sample was the greatest, while it decreased with the insertion of GNP in the CNGF sample, indicating that GNP has poor ferromagnetic properties, while ferrites are good, and their combination results in weak magnetization due to the formation of a single domain rather than domain walls. Furthermore, it has also been observed from the XRD analysis that the sizes of CNGF-

GNPs nanocomposites are reduced due to the formation of a domain.<sup>45</sup> The combined magnetization graphs of the samples plotted against the applied field are shown in Fig. 14. The trends of  $M_s$  and  $M_r$  are shown in Fig. 15. The values of maximum saturation magnetization, remanence, squareness ratio, and coercivity are tabulated in Table 3. The  $H_c$  values are also shown in Table 3, which reveals the maximum value for graphene, while the minimum value was recorded for the CNGF nanoparticles. Coercivity values depend on anisotropy, and

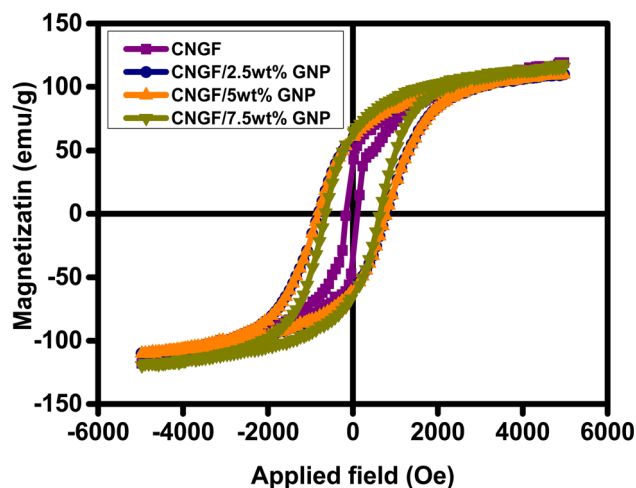


Fig. 14 The  $M-H$  loops for all the as-prepared composites.

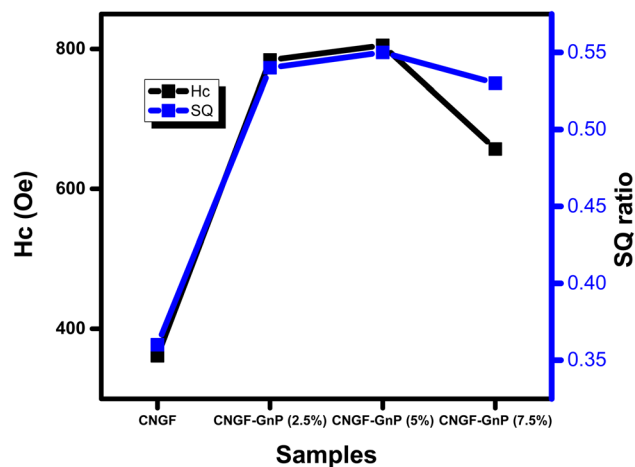


Fig. 15 Variation in the  $H_c$  and SQ ratios versus the samples.



Table 3 The magnetic parameters of the CNFG nanoparticles and CNFG-GNP nanocomposites

Parameters	CNFG	CNFG/2.5wt%GNP	CNFG/5wt%GNP	CNFG/7.5wt%GNP
$M_s$ (emu g <sup>-1</sup> )	119.8	110.5	110	116.3
$M_r$ (emu g <sup>-1</sup> )	43.3	60	60	61.3
$H_c$ (Oe)	361.1	784	805	657
$K$ (erg cm <sup>-3</sup> )	45 062	90 241.6	92 239.5	76 409.1
SQ	0.36	0.54	0.55	0.53

hence, the magneto-crystalline anisotropic parameters  $K$  were calculated by using  $M_s$  and  $H_c$  and are tabulated in Table 3. The CNFG nanoparticles had the minimum anisotropic constant value, while it was increased for the samples with GNPs. The change in coercivity value is due to the change in magneto-crystalline anisotropy as a consequence of the material added in the Nickel ferrites to replace the Ni ions at the octahedral sites.<sup>45</sup> The squareness ratio was calculated for each sample, as tabulated in Table 3. Moreover, the squareness ratio and coercivity values of the samples are plotted in Fig. 16.

The saturation magnetization decreased when GNPs were added to the spinel ferrites. It was also obvious that an increase in the GNP ratio led to an increase in the  $M_s$  value of the CNFG-GNP nanocomposites, while the anisotropic constant increased when GNPs were added. A further change in the weight ratio of GNPs led to a decrease in  $K$  values. Previous studies have revealed that graphene has a diamagnetic nature, so its momentum is in the opposite direction to the applied external magnetic field. However, our present findings show that the CNFG-GNP nanocomposites have different magnetic behavior from that of their constituents *i.e.*, spinel ferrites, which are ferromagnetic, and graphene, which is diamagnetic. Therefore, it can be hypothesized that in the presence of an applied external field, the exchange coupling phenomenon exists between them due to which the synthesized neoteric nanocomposites have wide and shifted hysteresis loops as the momentums of CNFG and GNPs interact.<sup>46</sup> Hence, it can be stated that the CNFG nanoparticles attached to the layers of GNPs act as functionalizing agents. Accordingly, it can be concluded that the magnetic properties of as-prepared Co<sub>0.5</sub>

Ni<sub>0.5</sub> Fe<sub>2-x</sub>Gd<sub>x</sub>O<sub>4</sub>-GNP can be tuned by adjusting the ratio of GNPs. Such graphene-based spinel ferrites can be used in semiconductor applications, such as spintronics.<sup>47</sup>

## 4. Conclusion

The structural, dielectric, electrical, optical, and magnetic properties of CNFG nanoparticles and changes in these properties due to graphene nanoplatelets are reported in this study by fabricating CNFG-GNP nanocomposites with varying GNP concentrations *via* a sol-gel auto-combustion route. It was observed that crystalline size  $D$  decreased from 40.1817 nm to 28.76267 nm due to the addition of the graphene nanoplatelets, besides providing the confirmation of GNP in the crystal structure. The SEM images revealed that samples with graphene were thicker than CNFG. The Raman spectra also showed the substitution of GNPs in the CNFG lattice. The TCR% decreased rapidly with an increase in temperature both in the low-temperature range and high-temperature range; moreover, the peak values of TCR% were  $-2.43\% \text{ K}^{-1}$  at 373 K,  $-2.62\% \text{ K}^{-1}$  at 368 K,  $-2.71\% \text{ K}^{-1}$  at 341 K and  $-2.45\% \text{ K}^{-1}$  at 368 K for CNFG, CNFG-GNP 2.5%, CNFG-GNP 5wt% and CNFG-GNP 7.5wt%, respectively. With the change in frequency, the variations in permittivity and conductivity were observed. Moreover, at higher frequencies, the external alternating field vibrated so rapidly that the dipole of the materials remained stationary giving no polarization, while the  $\varepsilon''$  intercept presented a higher value at low frequencies due to the addition of the graphene nanoplatelets, which provided more ions to jump in the direction of the field at low frequencies. CNFG had the maximum impedance, while the CNFG-GNPs nanocomposites showed lower values. The saturation magnetization ( $M_s$ ) value of the CNFG nanoparticles was higher, while it decreased when graphene nanoplatelets were added, reaching the minima when GNP was at 2.5% and 5%. Meanwhile, the squareness ratio increased from 0.36 for CNFG to 0.55 for the GNGF-GNP nanocomposites.

## Conflicts of interest

There are no conflicts to declare.

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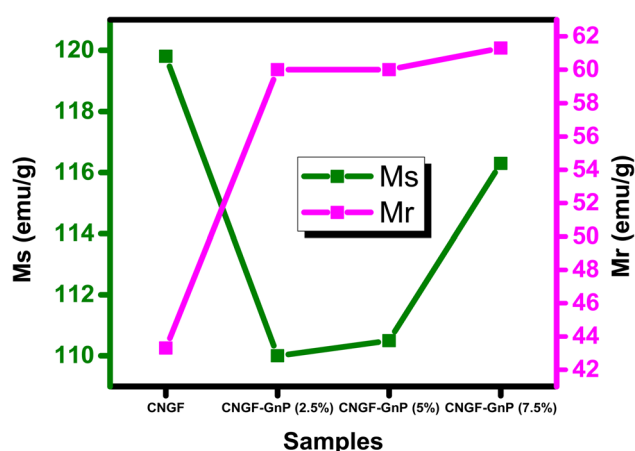


Fig. 16 The  $M_s$  and  $M_r$  trend of the samples.





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