


Cite this: *RSC Adv.*, 2023, 13, 5437

The influences of AlGa_N barrier epitaxy in multiple quantum wells on the optoelectrical properties of AlGa_N-based deep ultra-violet light-emitting diodes†

Tien-Yu Wang,^a Wei-Chih Lai,^{id} *^{ab} Qiao-Ju Xie,^a Shun-Hao Yang,^a Sheng-Po Chang,^{id} ^{abc} Cheng-Huang Kuo^d and Jinn-Kong Sheu^{id} ^{ab}

The growth conditions of the AlGa_N barrier in AlGa_N/AlGa_N deep ultra-violet (DUV) multiple quantum wells (MQWs) have crucial influences on the light output power of DUV light-emitting diodes (LEDs). The reduction of the AlGa_N barrier growth rate improved the qualities of AlGa_N/AlGa_N MQWs, such as surface roughness and defects. The light output power enhancement could reach 83% when the AlGa_N barrier growth rate was reduced from 900 nm h⁻¹ to 200 nm h⁻¹. In addition to the light output power enhancement, lowering the AlGa_N barrier growth rate altered the far-field emission patterns of the DUV LEDs and increased the degree of polarization in the DUV LEDs. The enhanced transverse electric polarized emission indicates that the strain in AlGa_N/AlGa_N MQWs was modified by lowering the AlGa_N barrier growth rate.

Received 20th November 2022
Accepted 6th February 2023

DOI: 10.1039/d2ra07368d

rsc.li/rsc-advances

1. Introduction

Deep ultraviolet (DUV) light has shown its importance in many fields, such as water purification, UV photolithography, high-density optical data storage, and portable chemical/biological agent detection/analysis systems.^{1–5} Furthermore, UV communication has been presented as a promising candidate for military wireless communication.^{6,7} A record-breaking data rate has been reported continuously in the past few years.^{8–11} Many applications require the light output power of the UV light source to be as high as possible. Especially, the output power of the UV light source essentially determines the transmission distance of UV communication, and its bandwidth has a major impact on the data rate.¹² The light output efficiency of AlGa_N-based DUV light-emitting diodes (LEDs) has been improved by the optimization of the epitaxy techniques for improving the material quality of AlN.^{13–17} Although the efficiency of DUV LEDs has been greatly improved, their performance is still hindered by high defect density, poor carrier injection, and low carrier

confinement. Many studies have reported strategies of defect reduction *via* substrate or template,^{18–20} methods of high carrier confinement, and hole injection.^{21–35} Low light extraction efficiency (LEE) is another challenge in achieving high efficiency for DUV LEDs. The LEE of the DUV LEDs has been improved by utilizing the flip-chip design,^{36,37} a patterned sapphire substrate,^{20,38} a patterned p-type layer,³⁹ and the nanowire structure.^{40,41}

AlGa_N/AlGa_N multiple quantum wells (MQWs) are the main active layers in DUV LEDs. The quality of AlGa_N/AlGa_N MQWs dominated the optoelectrical characteristics of DUV LEDs, such as leakage current and light output power. Furthermore, the defect reduction of AlGa_N/AlGa_N MQWs would improve the internal quantum efficiency (IQE) of DUV LEDs. The growth conditions of epitaxy directly affect the quality of AlGa_N/AlGa_N MQWs. Previously, Bryan *et al.*⁴² reported that the growth condition of the V/III ratio could impact the IQE of AlGa_N/AlGa_N MQWs. (The V/III ratio in metalorganic chemical vapor deposition (MOCVD) is normally defined as the ratio of the molar flow rate of ammonia (NH₃) to the total molar flow rates of trimethylgallium (TMGa) and trimethylaluminum (TMAI).) They found that AlGa_N/AlGa_N MQWs grown with high V/III ratio could achieve a high IQE of ~80%. The same group also reported that surface kinetics influence the growth of AlGa_N in terms of Al composition inhomogeneity and crystal quality.⁴³ Therefore, the growth conditions are crucial during the growth of AlGa_N/AlGa_N MQWs. During the growth of AlGa_N/AlGa_N MQWs, the condition of Al_{0.7}Ga_{0.3}N barrier might strongly affect the growth of the following Al_{0.4}Ga_{0.6}N wells. In this

^aDepartment of Photonics, National Cheng Kung University, Tainan 70101, Taiwan. E-mail: weilai@ncku.edu.tw

^bAdvanced Optoelectronic Technology Center, Research Center for Energy Technology and Strategy, National Cheng Kung University, Tainan 70101, Taiwan

^cInstitute of Microelectronics and Department of Electrical Engineering, National Cheng Kung University, Tainan 70101, Taiwan

^dInstitute of Lighting and Energy Photonics, College of Photonics, National Yang Ming Chiao Tung University, Tainan 71150, Taiwan

† Electronic supplementary information (ESI) available. See DOI: <https://doi.org/10.1039/d2ra07368d>



study, we varied the growth rates of the $\text{Al}_{0.7}\text{Ga}_{0.3}\text{N}$ barrier to investigate the influences on the electrical and optical properties of DUV LEDs.

2. Experimental

Epitaxial wafers were prepared on a 2-in (0001) sapphire substrate by MOCVD (Thomas Swan close-coupled showerhead 31×2 in). The source materials of gallium (Ga), aluminum (Al), and nitrogen (N) were TMGa, TMAI, and NH_3 , respectively. The n-type doping source was silane (SiH_4), and the p-type doping source was bis(cyclopentadienyl)magnesium(II) (Cp_2Mg). Initially a 20 nm-thick AlN nucleation layer was deposited on sapphire at 500 °C, and a 3 μm -thick undoped AlN layer was grown at 1250 °C. The screw and edge dislocation densities of the AlN epitaxy were $1.03 \times 10^8 \text{ cm}^{-2}$ and $1.64 \times 10^9 \text{ cm}^{-2}$, respectively. Following the AlN layer, a 20-pair AlN (2.5 nm)/ $\text{Al}_{0.8}\text{Ga}_{0.2}\text{N}$ (2.5 nm) superlattice and a 0.5 μm -thick $\text{Al}_{0.68}\text{Ga}_{0.32}\text{N}$ layer were grown as transition layers. A 1.5 μm -thick $\text{Al}_{0.6}\text{Ga}_{0.4}\text{N}:\text{Si}$ layer was grown as an n-type layer. The active layer was a five-pair $\text{Al}_{0.7}\text{Ga}_{0.3}\text{N}:\text{Si}$ (12 nm)/ $\text{Al}_{0.4}\text{Ga}_{0.6}\text{N}$ (2 nm) MQWs structure grown at 1050 °C, and a 2 nm-thick undoped $\text{Al}_{0.65}\text{Ga}_{0.35}\text{N}$ was grown as the last barrier. The $\text{Al}_{0.7}\text{Ga}_{0.3}\text{N}$ barriers grown at different growth rates were obtained by modifying the molar flows of TMGa and TMAI, and the molar flow of silane was also adjusted to fix the doping concentration. The growth rates of the $\text{Al}_{0.7}\text{Ga}_{0.3}\text{N}$ barriers were 900, 400, and 200 nm h^{-1} , respectively. After the MQW structure, an electron blocking layer, $\text{Al}_{0.8}\text{Ga}_{0.2}\text{N}$ (7 nm), and three-pair $\text{Al}_{0.8}\text{Ga}_{0.2}\text{N}:\text{Mg}$ (7.5 nm)/ $\text{Al}_{0.7}\text{Ga}_{0.3}\text{N}:\text{Mg}$ (3.5 nm) were grown. The p-type layer consisted of a 10-pair $\text{Al}_{0.8}\text{Ga}_{0.2}\text{N}:\text{Mg}$ (1 nm)/ $\text{Al}_{0.48}\text{Ga}_{0.52}\text{N}:\text{Mg}$ (1 nm) superlattice, a 20 nm-thick Al-grading AlGaIn:Mg layer (Al composition from 47% to 25%), and 10 nm-thick GaN:Mg layer.

The chip processes were performed to fabricate $500 \mu\text{m} \times 500 \mu\text{m}$ flip-chip DUV LEDs. The mesa area was revealed by the inductively coupled plasma (ICP) etching process, and Ti (20 nm)/Al (120 nm)/Ti (100 nm)/Au (50 nm) was subsequently deposited as the n-contact. Indium tin oxide (20 nm)/Al (100 nm) was deposited on the p-GaN as the p-contact. Through the thinning and dicing processes, the DUV LED chip was flipped and mounted on the AlN submount by AuSn eutectic bonding. The epitaxial structure and the schematic of the flip-chip DUV LED are shown in Fig. S1(a) and (b) in ESI,[†] respectively. A Keysight B1500 semiconductor parameter analyzer was employed to measure the current–voltage (I – V) characteristics of the flip-chip DUV LEDs for high- and low-current measurements. A calibrated integrating sphere and a spectrometer (Ocean Optics USB2000) were used to measure the output power and the emission spectra at room temperature. Moreover, the surface morphology of the AlGaIn/AlGaIn MQWs was measured by atomic force microscopy (AFM, Bruker Dimension Icon).

3. Results and discussion

The fabricated DUV LEDs with AlGaIn barrier grown at 900, 400, and 200 nm h^{-1} are denoted as BG900, BG400, and BG200,

respectively. Fig. 1 reveals the forward I – V characteristics and dynamic resistances of the AlGaIn-based DUV LEDs with different barrier growth rates. BG900, BG400, and BG200 have forward voltages (V_f) of 5.994, 5.915, and 5.948 V at 100 mA, respectively. The V_f of the DUV LEDs decreased slightly after the decrease of the AlGaIn barrier growth rate. In addition, the dynamic resistances at 350 mA of DUV LEDs with different AlGaIn barrier growth rates were in the range of 19.6–19.8 Ω . The inset of Fig. 1 shows the reverse I – V curves of the DUV LEDs, and the reverse leakage currents at -15 V of BG900, BG400, and BG200 were -0.2 , -1.15 , and $-2.49 \mu\text{A}$, respectively. The reverse leakage current at -15 V increased with the decreasing AlGaIn barrier growth rate. By reducing the AlGaIn barrier growth rate, the forward I – V changed slightly, but the reverse leakage current increased. The cause of the increased reverse currents of DUV LEDs with low AlGaIn barrier growth rates remains unclear.

The light output power–current (L – I) and external quantum efficiency–current (EQE– I) curves of the DUV LEDs with different AlGaIn barrier growth rates are presented in Fig. 2(a). The light output powers of all the DUV LEDs increased initially with the current and then reached the maximum output power at certain currents. The light output powers and EQEs of the DUV LEDs increased with decreasing AlGaIn barrier growth rate for all the driving currents. The 100 mA light output powers of BG900, BG400, and BG200 were 10.01, 14.01, and 15.51 mW, and the corresponding EQEs were 2.26%, 3.16%, and 3.49%, respectively. The 100 mA light output power enhancement was 40% by reducing the barrier growth rate from 900 nm h^{-1} to 400 nm h^{-1} , and the 100 mA light output enhancement could reach the largest value of 55% when the barrier growth rate was reduced to 200 nm h^{-1} . The enhancements of the maximum light output power were 55% and 83% for the barrier growth rates of 400 and 200 nm h^{-1} , respectively, compared with the barrier growth rate of 900 nm h^{-1} . The driving currents for the maximum light output power increased with the decrease of the barrier growth rate. The peak EQEs of BG900, BG400, and BG200 were 2.48%, 3.87%, and 4.27%, respectively. The corresponding improvements of EQE were 56% and 73% for the reductions of the barrier growth rates from 900 nm h^{-1} to 400

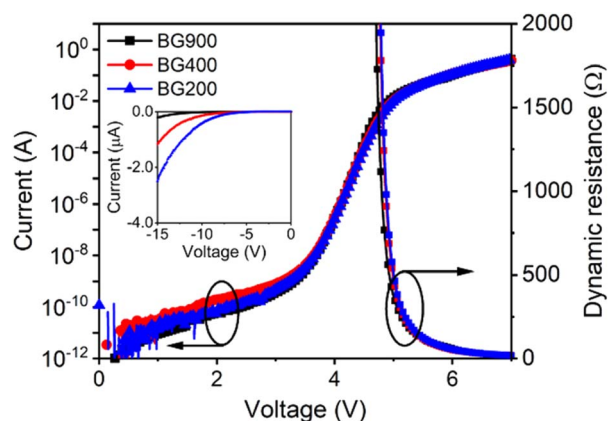


Fig. 1 Semi-log forward I – V curves of the DUV LEDs varying AlGaIn barrier growth rate. The inset of figure is linear reverse I – V curves.



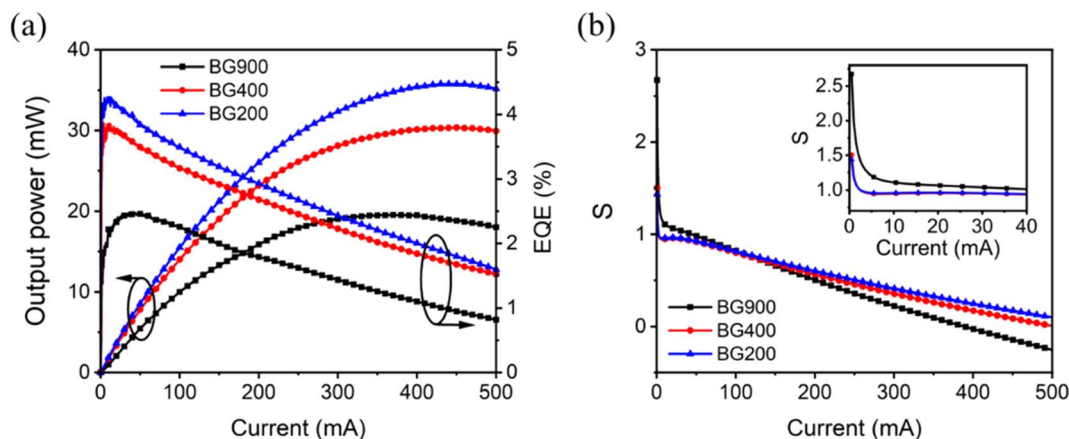


Fig. 2 (a) Light output power–current (L – I) and external quantum efficiency–current (EQE– I) curves of the DUV LEDs varying AlGaIn barrier growth rate. (b) S –current (S – I) curves of the DUV LEDs varying AlGaIn barrier growth rate. The inset of (b) is S – I curves at low injection current.

and 200 nm h^{-1} , respectively. Furthermore, the efficiency droops (defined as $1 - \text{EQE}_{450\text{mA}}/\text{EQE}_{\text{peak}}$) of the DUV LEDs were improved by decreasing the barrier growth rate. The efficiency droops of BG900, BG400, and BG200 were 62%, 57%, and 58%, respectively. The enhancements of the light output power and EQE of the DUV LEDs could be attributed to the quality improvement of the low AlGaIn barrier growth rate.

The L – I curves were further analysed by $\text{dlog } L/\text{dlog } I$ (ref. 44) (S is defined as the differential of the $\log L$ – $\log I$ curve) to obtain the S – I curves to further elucidate the dominant carrier recombination process in the DUV LEDs with different AlGaIn barrier growth rates. BG400 and BG200 presented lower S values than BG900 in the low injection current region (<40 mA), as presented in the inset of Fig. 2(b). The reduction of the S value in the low current region implies that the non-radiative recombination *via* defect was suppressed by lowering the barrier growth rate. In the mid and high injection current regions (>150 mA) in Fig. 2(b), the S values of BG400 and BG200 were larger than that of BG900, and all these S values were less than 1. The dominant carrier recombination process in this current range would be the carrier overflow, and a low S value implies severe carrier overflow. Therefore, a high AlGaIn barrier growth rate would have a large carrier overflow. A low AlGaIn barrier growth rate would have good carrier confinement in AlGaIn/AlGaIn MQWs, which would effectively reduce the carrier overflow process and improve the efficiency droop.

The surface morphology of AlGaIn can be changed by the growth conditions. The high growth rate induced 2D nucleation in the steps and the AlGaIn growth tended to be in 3D growth mode.⁴³ The high growth rate might increase the surface roughness of the AlGaIn barrier. Fig. 3 shows the AFM images of the MQWs with different AlGaIn barrier growth rates. The root mean square (RMS) surface roughness (R_q) of the AlGaIn/AlGaIn MQWs increased from 3.95 nm to 6.94 nm when the AlGaIn barrier growth rate was increased from 200 nm h^{-1} to 900 nm h^{-1} . The high growth rate increased the surface roughness of the AlGaIn barrier. This phenomenon might deteriorate the interface quality of the AlGaIn/AlGaIn MQWs and weaken the IQE. Moreover, the high growth rate of the AlGaIn barrier might

enhance the inhomogeneity of the Al composition of the barrier, which might also decrease the IQE. Given the fixed ammonia molar flow during the growth of MQWs, the reduction of the AlGaIn barrier growth rate also implies the increase of the V/III ratio (from 1900 to 6300) in AlGaIn growth. Bryan *et al.*⁴² reported that the high V/III ratio during the growth of AlGaIn/AlGaIn MQWs could suppress the midgap defect luminescence due to the reduction of the point defects. Therefore, the AlGaIn barrier grown at low growth rate could have a better surface quality and fewer defects than that grown at high growth rate. These phenomena could result in the enhancement of the 100 mA light output power and the reduction of the S value in the low current range of DUV LEDs. Moreover, the improved interfaces of AlGaIn/AlGaIn MQWs due to a low AlGaIn barrier growth rate would improve the MQWs' carrier confinements, which could suppress the carrier overflow and enhance the S value of DUV LEDs in the mid and high current regions.

To further understand the impacts of the AlGaIn barrier growth rate on the optical properties, far-field emission pattern and polarization measurements were taken. The far-field emission patterns of the DUV LEDs are presented in Fig. 4. The DUV LED with the highest growth rate, BG900, showed the lowest light intensity at each measured angle. The light intensities were augmented at all angles with the reduction of the barrier growth rate. The increases in light intensity in the far-field emission patterns are consistent with the enhancements of the light output power for the DUV LEDs with decreasing barrier growth rates. In addition, the enhancement of the light intensity at 90° was higher than those at 30° and 150° . The reduction of barrier growth rate primarily enhanced the normal emission and led to a shrink in the emission angles (the angle width at half maximum intensity) of the DUV LEDs. The normal light intensity to lateral light intensity ratio of the DUV LEDs might be related to the strain in the AlGaIn/AlGaIn MQWs due to the band structure change. By reducing the barrier growth rate, the surface roughness of the MQWs was diminished, and the strain in the MQWs might be altered at the same time. The large enhancement of the normal light intensity of the DUV LEDs



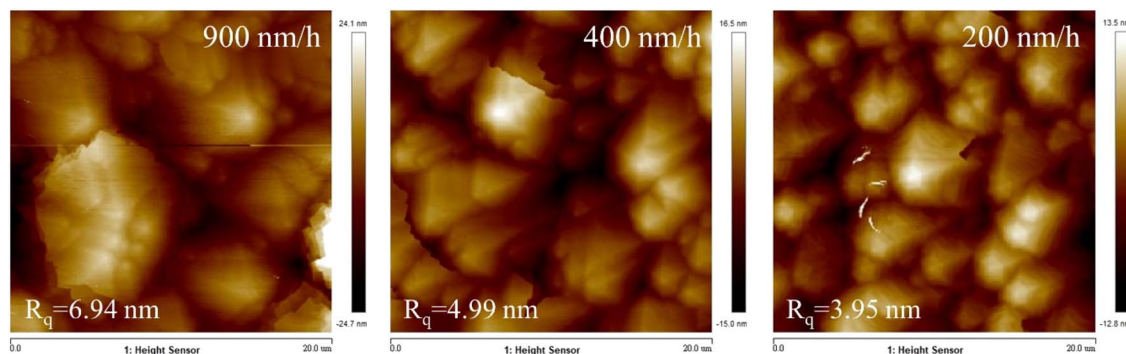


Fig. 3 AFM images of the DUV LEDs varying AlGaIn barrier growth rate.

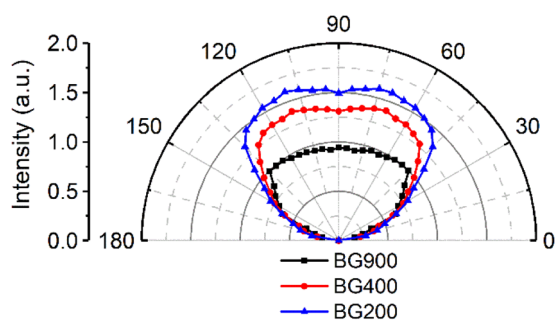


Fig. 4 Far-field emission patterns of the DUV LEDs varying AlGaIn barrier growth rate.

with low barrier growth rates might be attributed to the change of the strain in the MQWs.

The light emission of the AlGaIn-based DUV LEDs contained both TE- and TM-polarized lights, which originated from different band-to-band transitions. Many studies reported that the TE- and TM-polarized lights of AlGaIn-based LEDs mainly propagate in the normal (perpendicular to the c -axis, $E \perp c$) and lateral directions (parallel to the c -axis, $E \parallel c$), respectively.^{40,41} Given that the AlGaIn barrier growth rate affected the normal light intensity, polarization-dependent electroluminescence (EL) was performed to investigate the change of the TE- and TM-polarized lights of the DUV LEDs. A spectrometer was used to measure the EL spectrum of the sample at an angle of 10° to the bottom surface of the sample. The TE- and TM-polarized EL emissions could be distinguished by the polarizer (a Glan-Taylor prism) inserted between the lenses. The EL spectra of the TE- and TM-polarized lights of all the DUV LEDs are provided in Fig. S2 in the ESI.[†] The degree of polarization (DOP) is defined as $DOP = (I_{TE} - I_{TM}) / (I_{TE} + I_{TM})$, where I_{TE} and I_{TM} are the integrated intensities of the TE- and TM-polarized EL spectra, respectively. Fig. 5 shows that the DOP in the DUV LEDs varied with the AlGaIn barrier growth rates. The DOP increased with decreasing AlGaIn barrier growth rate. The result indicates that the low AlGaIn barrier growth rate augmented the TE-polarized emission in the DUV LEDs. Zhang *et al.*⁴⁵ reported that the compressive strain in the AlGaIn/AlGaIn MQWs could enhance the TE-polarized light emission. With the decrease of the AlGaIn barrier growth rate, the increase of DOP could be attributed to the increase of the compressive strain in the MQWs. Given that

the low AlGaIn barrier growth rate enhanced the TE-polarized emission of the DUV LEDs, the normal light intensity of the far-field pattern was significantly increased, and the emission angle was also shrunk, as shown in Fig. 4.

To understand how the growth rate affected the AlGaIn barrier, the AlN (104) X-ray diffraction reciprocal space mapping (RSM) was performed on the sample of a 400 nm-thick $Al_{0.7}Ga_{0.3}N$ barrier layer grown with different growth rates to analyse the strain. The detailed epitaxial structure is presented in Fig. S3.[†] The structure from AlN to n-AlGaIn layers was retained so that the situation of strain before growing MQWs was the same as that in the DUV LED structure. The AlN (104) RSMs of AlGaIn barrier layers grown with different growth rate are presented in Fig. 6, and it was found that the contour lines of the AlGaIn barrier layer shifted along the relaxation line, which meant that the strain was partially relaxed. All AlGaIn barrier layers grown with different growth rates showed partial relaxation. According to the RSM data, the calculated Al composition, in-plane strain (ϵ_a), and relaxation of in-plane strain (R_a) of AlGaIn barrier layers grown with different growth rates are summarized in Table 1. The ϵ_a and R_a are defined as

$$\epsilon_a = \frac{a_{\text{actual}} - a_{\text{relaxed}}}{a_{\text{relaxed}}}$$

and

$$R_a = \frac{a_{\text{actual}} - a_{\text{strained}}}{a_{\text{relaxed}} - a_{\text{strained}}}$$

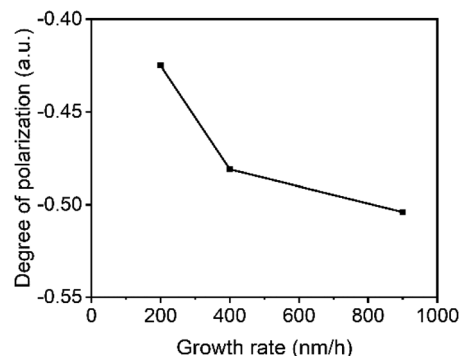


Fig. 5 Dependence of DOP at 100 mA current on AlGaIn barrier growth rate.



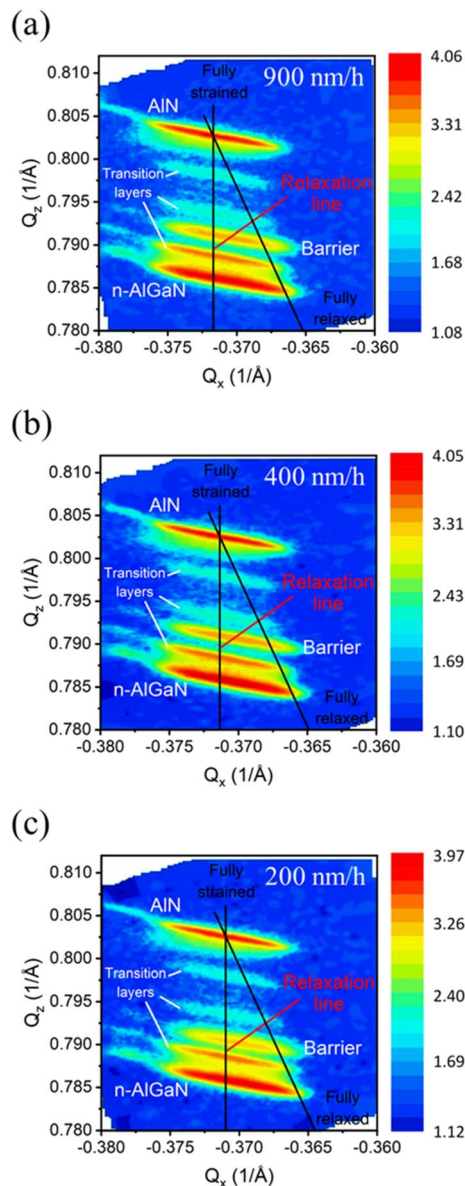


Fig. 6 The reciprocal space mappings of AlN (104) reflection plane of barrier layers with growth rate (a) 900, (b) 400, and (c) 200 nm h⁻¹.

where a_{actual} is the actual measured lattice constant, a_{relaxed} is the fully relaxed lattice constant, and a_{strained} is the fully strained lattice constant. The lattice constant a_{strained} is set to be a_{AlN} under the condition of pseudomorphic growth on AlN buffer layer. In almost the same Al content, the value of R_a

increased with increasing the AlGa_N barrier growth rate. In other words, the AlGa_N barrier grown at low growth rate tended to sustain the in-plane strain. Therefore, the AlGa_N barrier grown at high growth rate would be prone to relax its lattice constant compared to the AlGa_N barrier grown at low growth rate. Furthermore, the surface morphology of MQWs changed with varying barrier growth rates, as presented in Fig. 3. This phenomenon implied that there could be a change of strain even at the 60 nm-thick total barrier thickness in MQWs. Therefore, with lowering the barrier growth rate, the compressive strain in the barrier increased, and the lattice constant of the barrier shrank compared to the AlGa_N barrier grown at high growth rate. The well layer on the AlGa_N barrier grown at low growth rate was under larger compressive strain than the well layer on the AlGa_N barrier grown at high growth rate. The compressive strain altered the band structure of MQWs, so the TE-polarized emission was enhanced, and the DOP in the emission increased.

Finally, we compared our results to several previous reports.^{46–48} The flip-chip DUV LEDs with similar emission wavelengths are listed in Table 2. Zhang *et al.* proposed a dislocation filter (DF) layer to improve the unfavourable step-bunching morphology of AlN layer grown at a high growth rate.⁴⁶ The DF layer made a transition from step-bunching to step-flow morphology, which made it feasible to grow high quality AlN at high growth rate. Hu *et al.* proposed a superlattice electron deceleration layer (SEDL) to improve the performance of DUV LED.⁴⁷ The SEDL could decrease the electron injection current density and increase the hole injection current density. Hence, the radiative recombination rates in MQWs were raised by SEDL. Sun *et al.* reported the performance of DUV LED grown on nano-patterned sapphire substrates (NPSS) in combination with the optimization of V/III ratio and Si doping in MQWs.⁴⁸ They also found that the IQE of MQWs was improved from 56.2% to 70.3% when the V/III ratio was increased from 500 to 1000. However, further increasing V/III from 1000 to 2000 in MQWs slightly reduced the IQE from 70.3% to 69.2%. This phenomenon could be attributed to the degradation of crystal quality of MQWs under high V/III ratio. Even if the AlGa_N MQWs were grown at high V/III ratio, the crystal quality was also important. In our study, the V/III ratio was increased by lowering the growth rate of the AlGa_N barrier. According to the AlN (104) RSM results, the AlGa_N barrier grown at high growth rate tended to relax the strain, and the relaxation of strain usually caused defects. In other words, the barrier grown at low growth rate tended to sustain the strain, which caused fewer defects. Therefore, by lowering the growth rate, the V/III ratio

Table 1 The calculated Al compositions, in-plane strains, and relaxation of in-plane strain of barrier layer grown with different growth rates

Growth rate of barrier (nm h ⁻¹)	Al composition (%)	In-plane strain, ϵ_a (%)	Relaxation of in-plane strain, R_a (%)
900	69.9	−0.46	40.1
400	70.2	−0.48	37.6
200	69.5	−0.53	31.4



Table 2 The performances of the AlGaIn-based DUV LEDs

Ref.	Techniques	Structure of LED chip	Wavelength (nm)	MQWs pairs	Maximum EQE (%)	Maximum power enhancement (%)
46	Dislocation filtering (DF) layer	Flip-chip	276.6	5	2.28	40.4
47	Superlattice electron deceleration layer (SEDL)	Flip-chip	275	5	3.43	200
48	Optimization of V/III ratio and Si doping in MQWs and nano-patterned sapphire substrates (NPSS)	Flip-chip	276.1	10	4.01	n/a
This work	Reduction of barrier growth rate	Flip-chip	279	5	4.27	83

was increased without degrading the crystal quality of the AlGaIn barrier. Both effects were beneficial to the optical properties of MQWs. Our DUV LEDs with AlGaIn barrier grown at low growth rate presented the maximum EQE of 4.27% and the maximum power enhancement of 83%, which were comparable to the previous reports. Furthermore, the DOP of DUV LEDs was increased by reducing the growth rate of the AlGaIn barrier in MQWs.

4. Conclusions

In summary, the growth conditions of the AlGaIn barrier in the AlGaIn/AlGaIn MQWs presented crucial influences on the light output power of the DUV LEDs. The reduction of the AlGaIn barrier growth rate improved the surface roughness of the AlGaIn/AlGaIn MQWs, and the smooth surface could be favourable to enhancing the IQE. Furthermore, the accompanying high V/III ratio during the growth could lessen the AlGaIn barrier defects. The light output power enhancement could reach 83% when the AlGaIn barrier growth rate was reduced from 900 nm h⁻¹ to 200 nm h⁻¹. Furthermore, lowering the AlGaIn barrier growth rate enhanced the normal light intensity in the far-field emission patterns and increased the DOP in the DUV LEDs. The enhanced TE-polarized emission implies that the strain in the MQWs was compressive through the decrease of the AlGaIn barrier growth rate.

Conflicts of interest

There are no conflicts to declare.

Acknowledgements

We extend our gratitude to the National Science and Technology Council, Taiwan for their financial support under Grant Numbers: MOST 111-2221-E-006-033-MY3, MOST 107-2221-E-006-186-MY3, MOST 110-2221-E-006-153 and MOST 110-2218-E-006-025-MBK. We would also like to thank the core facility centre of the National Cheng Kung University for the support of the SEM observation and ICP etcher.

References

- W. L. Morison, *Phototherapy and photochemotherapy for skin disease*, CRC Press, 2005.
- III-Nitride Ultraviolet Emitters: Technology and Applications*, ed. M. Kneissl and J. Rass, Springer, Cham, 2016.
- H. Hirayama, T. Yatabe, N. Noguchi and N. Kamata, *Electron. Commun. Jpn.*, 2010, **93**, 24–33.
- M. Martens, F. Mehnke, C. Kuhn, C. Reich, V. Kueller, A. Knauer, C. Netzel, C. Hartmann, J. Wollweber, J. Rass, T. Wernicke, M. Bickermann, M. Weyers and M. Kneissl, *IEEE Photonics Technol. Lett.*, 2014, **26**, 342–345.
- F. Mehnke, M. Guttmann, J. Enslin, C. Kuhn, C. Reich, J. Jordan, S. Kapanke, A. Knauer, M. Lapeyrade, U. Zeimer, H. Krüger, M. Rabe, S. Einfeldt, T. Wernicke, H. Ewald, M. Weyers and M. Kneissl, *IEEE J. Sel. Top. Quantum Electron.*, 2017, **23**, 29–36.
- Z. Xu and B. M. Sadler, *IEEE Commun. Mag.*, 2008, **46**, 67–73.
- G. A. Shaw, A. M. Siegel, J. Model and D. Greisokh, in *Unattended Ground Sensor Technologies and Applications VII*, SPIE, 2005, vol. 5796, pp. 214–225.
- A. Vavoulas, H. G. Sandalidis, N. D. Chatzidiamantis, Z. Xu and G. K. Karagiannidis, *IEEE Commun. Surv. Tutor.*, 2019, **21**, 2111–2133.
- O. Alkhazragi, F. Hu, P. Zou, Y. Ha, Y. Mao, T. K. Ng, N. Chi and B. S. Ooi, *Optical Fiber Communication Conference (OFC) 2020, OSA*, 2020, p. M3I.5.
- K. Kojima, Y. Yoshida, M. Shiraiwa, Y. Awaji, A. Kanno, N. Yamamoto and S. Chichibu, *2018 European Conference on Optical Communication (ECOC)*, 2018, pp. 1–3.
- O. Alkhazragi, F. Hu, P. Zou, Y. Ha, C. H. Kang, Y. Mao, T. K. Ng, N. Chi and B. S. Ooi, *Opt. Express*, 2020, **28**, 9111–9122.
- L. Guo, Y. Guo, J. Wang and T. Wei, *J. Semicond.*, 2021, **42**, 081801.
- H. Hirayama, S. Fujikawa, N. Noguchi, J. Norimatsu, T. Takano, K. Tsubaki and N. Kamata, *Phys. Status Solidi*, 2009, **206**, 1176–1182.
- H. Hirayama, S. Fujikawa, J. Norimatsu, T. Takano, K. Tsubaki and N. Kamata, *Phys. Status Solidi C*, 2009, **6**, S356–S359.
- S. M. Islam, V. Protasenko, K. Lee, S. Rouvimov, J. Verma, H. Xing and D. Jena, *Appl. Phys. Lett.*, 2017, **111**, 091104.



- 16 M. Shatalov, W. Sun, R. Jain, A. Lunev, X. Hu, A. Dobrinsky, Y. Bilenko, J. Yang, G. A. Garrett, L. E. Rodak, M. Wraback, M. Shur and R. Gaska, *Semicond. Sci. Technol.*, 2014, **29**, 084007.
- 17 C. Pernot, S. Fukahori, T. Inazu, T. Fujita, M. Kim, Y. Nagasawa, A. Hirano, M. Ippommatsu, M. Iwaya, S. Kamiyama, *et al.*, *Phys. Status Solidi*, 2011, **208**, 1594–1596.
- 18 J. R. Grandusky, S. R. Gibb, M. C. Mendrick, C. Moe, M. Wraback and L. J. Schowalter, *Appl. Phys. Express*, 2011, **4**, 082101.
- 19 N. Susilo, S. Hagedorn, D. Jaeger, H. Miyake, U. Zeimer, C. Reich, B. Neuschulz, L. Sulmoni, M. Guttmann, F. Mehnke, C. Kuhn, T. Wernicke, M. Weyers and M. Kneissl, *Appl. Phys. Lett.*, 2018, **112**, 041110.
- 20 T. Takano, T. Mino, J. Sakai, N. Noguchi, K. Tsubaki and H. Hirayama, *Appl. Phys. Express*, 2017, **10**, 031002.
- 21 H. Hirayama, Y. Tsukada, T. Maeda and N. Kamata, *Appl. Phys. Express*, 2010, **3**, 031002.
- 22 R. K. Mondal, V. Chatterjee and S. Pal, *Opt. Mater.*, 2020, **104**, 109846.
- 23 A. Pandey, W. J. Shin, X. Liu and Z. Mi, *Opt. Express*, 2019, **27**, A738–A745.
- 24 X. Wang, H.-Q. Sun and Z.-Y. Guo, *Opt. Mater.*, 2018, **86**, 133–137.
- 25 S. Wang, Y. A. Yin, H. Gu, N. Wang and L. Liu, *J. Disp. Technol.*, 2016, **12**, 1112–1116.
- 26 Q. Si, H. Chen, S. Li, S. Lu and J. Kang, *IEEE Photonics J.*, 2017, **9**, 1–7.
- 27 X. Yang, H. Sun, X. Fan, Z. Zhang, J. Sun, X. Yi and Z. Guo, *Superlattices Microstruct.*, 2017, **101**, 293–298.
- 28 S. Sumiya, Y. Zhu, J. Zhang, K. Kosaka, M. Miyoshi, T. Shibata, M. Tanaka and T. Egawa, *Jpn. J. Appl. Phys.*, 2008, **47**, 43.
- 29 F. Mehnke, C. Kuhn, M. Guttmann, C. Reich, T. Kolbe, V. Kueller, A. Knauer, M. Lapeyrade, S. Einfeldt, J. Rass, T. Wernicke, M. Weyers and M. Kneissl, *Appl. Phys. Lett.*, 2014, **105**, 051113.
- 30 Y. Hou and Z. Guo, *Opt. Commun.*, 2019, **433**, 236–241.
- 31 Y. Kuo, J. Chang, F. Chen, Y. Shih and H. Chang, *IEEE J. Quantum Electron.*, 2016, **52**, 1–5.
- 32 W. Hu, P. Qin, W. Song, C. Zhang, R. Wang, L. Zhao, C. Xia, S. Yuan, Y. Yin and S. Li, *Superlattices Microstruct.*, 2016, **97**, 353–357.
- 33 J. Simon, V. Protasenko, C. Lian, H. Xing and D. Jena, *Science*, 2010, **327**, 60–64.
- 34 J. Zhang, W. Tian, F. Wu, W. Yan, H. Xiong, J. Dai, Y. Fang, Z. Wu and C. Chen, *IEEE Photonics J.*, 2013, **5**, 1600310.
- 35 X. Bao, P. Sun, S. Liu, C. Ye, S. Li and J. Kang, *IEEE Photonics J.*, 2015, **7**, 1–10.
- 36 H.-Y. Ryu, I.-G. Choi, H.-S. Choi and J.-I. Shim, *Appl. Phys. Express*, 2013, **6**, 062101.
- 37 K. H. Lee, H. J. Park, S. H. Kim, M. Asadirad, Y.-T. Moon, J. S. Kwak and J.-H. Ryou, *Opt. Express*, 2015, **23**, 20340–20349.
- 38 Y. K. Ooi and J. Zhang, *IEEE Photonics J.*, 2018, **10**, 1–13.
- 39 P. Zhao, L. Han, M. R. McGoogan and H. Zhao, *Opt. Mater. Express*, 2012, **2**, 1397–1406.
- 40 Y. K. Ooi, C. Liu and J. Zhang, *IEEE Photonics J.*, 2017, **9**, 1–12.
- 41 M. Djavid and Z. Mi, *Appl. Phys. Lett.*, 2016, **108**, 051102.
- 42 Z. Bryan, I. Bryan, J. Xie, S. Mita, Z. Sitar and R. Collazo, *Appl. Phys. Lett.*, 2015, **106**, 142107.
- 43 I. Bryan, Z. Bryan, S. Mita, A. Rice, L. Hussey, C. Shelton, J. Tweedie, J.-P. Maria, R. Collazo and Z. Sitar, *J. Cryst. Growth*, 2016, **451**, 65–71.
- 44 K.-S. Kim, D.-P. Han, H.-S. Kim and J.-I. Shim, *Appl. Phys. Lett.*, 2014, **104**, 091110.
- 45 S. Zhang, Y. Zhang, N. Tang, W. Wang, X. Chen, L. Fu, C. He, Y. Lv, Z. Feng, F. Xu, T. Yu, W. Ge and B. Shen, *Superlattices Microstruct.*, 2021, **150**, 106749.
- 46 Y. Zhang, H. Long, J. Zhang, B. Tan, Q. Chen, S. Zhang, M. Shan, Z. Zheng, J. Dai and C. Chen, *CrystEngComm*, 2019, **21**, 4072–4078.
- 47 J. Hu, J. Zhang, Y. Zhang, H. Zhang, H. Long, Q. Chen, M. Shan, S. Du, J. Dai and C. Chen, *Nanoscale Res. Lett.*, 2019, **14**, 347.
- 48 Y. Sun, F. Xu, N. Zhang, J. Lang, J. Wang, B. Liu, L. Wang, N. Xie, X. Fang, X. Kang, Z. Qin, X. Yang, X. Wang, W. Ge and B. Shen, *CrystEngComm*, 2021, **23**, 1201–1206.

