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Experimental verification of SO₂ and S desorption contributing to defect formation in MoS₂ by thermal desorption spectroscopy†

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The defect-free surface of MoS₂ is of high importance for applications in electronic devices. Theoretical calculations have predicted that oxidative etching could be responsible for sulfur vacancy formation. No direct experimental evidence, however, points out the role of adsorbed oxygen on sulfur vacancy formation for MoS₂, especially on an insulating SiO₂/Si substrate. Herein, by applying thermal desorption spectroscopy, we found that sulfur loss can be tightly coupled to adsorbed oxygen, as confirmed by observation of SO₂ desorption. With annealing MoS₂, even under ultrahigh vacuum, oxygen molecules adsorbed on MoS₂ assist the sulfur atom in dissociating from MoS₂, and thus, defects are formed as the result of SO₂ desorption from 200 °C to 600 °C. At higher temperatures (over 800 °C), on the other hand, direct sulfur desorption becomes dominant. This finding can be well explained by combining the morphology investigation enabled by atomic layer deposition at defective sites and optical transitions observed by photoluminescence measurements. Moreover, a preannealing treatment prior to exfoliation was found to be an effective method to remove the adsorbed oxygen, thus preventing defect formation.

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1. Introduction

MoS₂, a two-dimensional (2D) layered material, is promising as a channel material for next-generation field-effect transistors (FETs) because a natural thin body can overcome the scaling limit for the Si gate length.^{1,2} Although the dangling-bond-free surface of the layered MoS₂ channel is expected to ideally provide an electrically inert interface, sulfur vacancies in reality have been recognized as a dominant defect in MoS₂ due to the lowest formation energy of ~1.3–1.5 eV under S-poor conditions.^{3,4} Moreover, sulfur vacancies are known to introduce defect states in the band gap, which degrades the FET performance.^{5–8} Therefore, controlling sulfur vacancies has still been an critical issue due to the limited understanding of its formation mechanism, even though many healing processes based on various adsorbates or S vapor annealing have been proposed.^{9–12}

The stability of MoS₂ has been intensively studied thus far.^{13,14} The conventional mechanical exfoliation process for device fabrication leads to exposure to ambient air, which introduces adsorbates on the MoS₂ surface and edge. These adsorbates, such as oxygen and water, greatly affect chemical

stability, since thermodynamic calculations suggest that most 2D materials show oxidation tendencies.¹⁵ The long-term exposure of MoS₂ flakes to ambient air has proven the gradual oxidation from the edges to the interior of MoS₂¹⁶ because the coordinatively unsaturated edge is energetically more favorable for oxidation than the basal plane.^{17–20} For intentional exposure to an oxygen environment,²¹ the basal plane of MoS₂ is oxidized, and MoO₃ is formed during oxidation above 400 °C. This is further supported by the aggressive oxidation using oxygen plasma, where layer-by-layer oxidation was observed since MoO₃ formed on the MoS₂ surface prevented further oxidation.²² Interestingly, for an oxygen/air environment at lower temperatures of 300–340 °C, however, layer-by-layer anisotropic etching of MoS₂ results in triangular pits, which are initiated *via* intrinsic defects on the basal plane of MoS₂.^{21,23} This oxidative etching is explained by the reaction of MoS₂ + O₂ → MoO₃↑ + SO₂↑,^{18,24} where both MoO₃ and SO₂ are volatile. Thus, oxidative etching is considered to be an important process when MoS₂ is exposed to oxygen or air.

Recently, atomic-resolution scanning tunneling microscopy (STM) suggested that the O₂ adsorbed on the basal plane of MoS₂ volatilized as SO₂ by removing S, leaving S vacancies with O saturation on the basal plane.¹⁰ This O₂-assisted S vacancy formation in ambient air at room temperature (RT) is estimated to be energetically spontaneous with −0.49 eV by first principles calculation. On the other hand, in a high vacuum environment, the creation of S vacancies has been confirmed not at 127 °C but at elevated temperatures >627 °C by atomic-resolution STM²⁵

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and transmission electron microscopy (TEM).²⁶ These results indicate that the formation of S vacancies at RT cannot be fully explained without the existence of oxygen adsorbed on the basal plane of MoS₂, suggesting the importance of O₂ adsorption and reaction mechanism. Therefore, further theoretical studies on the initial step of oxidative etching have been conducted in detail.²⁷ However, direct experimental evidence of SO₂ desorption as well as direct S desorption has not yet been proven.

Here, thermal desorption mass spectrometry (TDS) is used to study the surface reaction and desorption kinetics in Si^{28,29}, Ge,³⁰ and HOPG³¹ by detecting a small amount of desorption species using a quadrupole mass spectrometer (QMS). By using this TDS system, it may be possible to detect SO₂ and/or S desorption from MoS₂ flakes transferred onto SiO₂/Si substrates in ambient air. Although no measurement method to observe S vacancies with atomic resolution is available for MoS₂ on an insulating substrate, it is well known that Al₂O₃ during atomic layer deposition (ALD) is adsorbed only at the defect sites;^{32–35} this will allow us to visualize the defect evolution with the assistance of surface topology observed macroscopically by atomic force microscopy (AFM). In this research, a quantitative characterization of the defect evolution mechanism of MoS₂ on SiO₂ at elevated temperatures up to ~1000 °C is realized by combining TDS and ALD.

2. Experimental

2H-MoS₂ crystals purchased from SPI supplies were used in this study. First, a 110 nm SiO₂/n⁺-Si substrate was sonicated by acetone and isopropyl alcohol for 10 min in laboratory air. After a few exfoliations by metallic tweezers, the bulk MoS₂ flakes were directly transferred to the SiO₂/n⁺-Si substrate. This procedure provides a sufficiently large surface area of MoS₂ for the TDS measurements, while it prevents any tape residue on the substrate since the tape is not used. Other 2H-bulk transition metal dichalcogenides (WS₂, WSe₂, MoSe₂ and MoTe₂) were grown by chemical vapor transport method³⁶ and treated by the same method applied for MoS₂. Alternatively, MoS₂ flakes were

clashed into powders using a mortar to further increase the surface area of MoS₂. Then, the MoS₂ powders were placed on the SiO₂/Si substrate. Fig. 1a illustrates the TDS apparatus, which allows selective heating of only the Si substrate *via* infrared radiation from the underlying power-operated lamp under ultrahigh vacuum (UHV) conditions of $\sim 4 \times 10^{-8}$ Pa. Since this TDS is cold wall system, the desorption from the inner chamber wall can be avoided during the sample heating. The QMS functions by manipulating the electric field between four orthogonal rods to collect ionized desorbed molecules and atoms and thus enables distinguishing the desorbed species according to its atomic/molecular mass-to-charge ratio. The bulk MoS₂ flakes were heated from room temperature to ~1000 °C with a heating rate of 20 °C min⁻¹, and TDS spectra were detected in multi-ion detection mode with a sensitivity up to 10⁻¹⁵ A.

For ALD, Al₂O₃ was deposited on MoS₂ flakes in a hot wall chamber at 200 °C with trimethylaluminum (TMA), water, and nitrogen as the precursor, oxidant, and purge gas, respectively.³² The pulse time for TMA/water is 0.1 s/0.5 s, respectively. Ten cycles were performed to obtain ~2 nm-thick Al₂O₃. The surface morphology was measured with dynamic force mode by AFM. The Raman and photoluminescence (PL) spectra were measured at RT in ambient air using a 488 nm excitation laser with a power of 0.06 mW to avoid the degradation of MoS₂ owing to laser heating.

3. Results and discussion

First, the desorption of adsorbates on the SiO₂/Si substrate without MoS₂ flakes should be examined. As shown in Fig. S1a,† in addition to the small amount of H₂, CO₂ and N₂, H₂O physisorbed on the SiO₂/Si substrate was clearly detected in the temperature range of 100–300 °C because the SiO₂ surface is hydrophilic.³⁷ After transferring MoS₂ flakes onto the SiO₂/Si substrate, the desorption of H₂O and CO₂ was enhanced, as shown in Fig. S1b.† This indicates that many kinds of gas

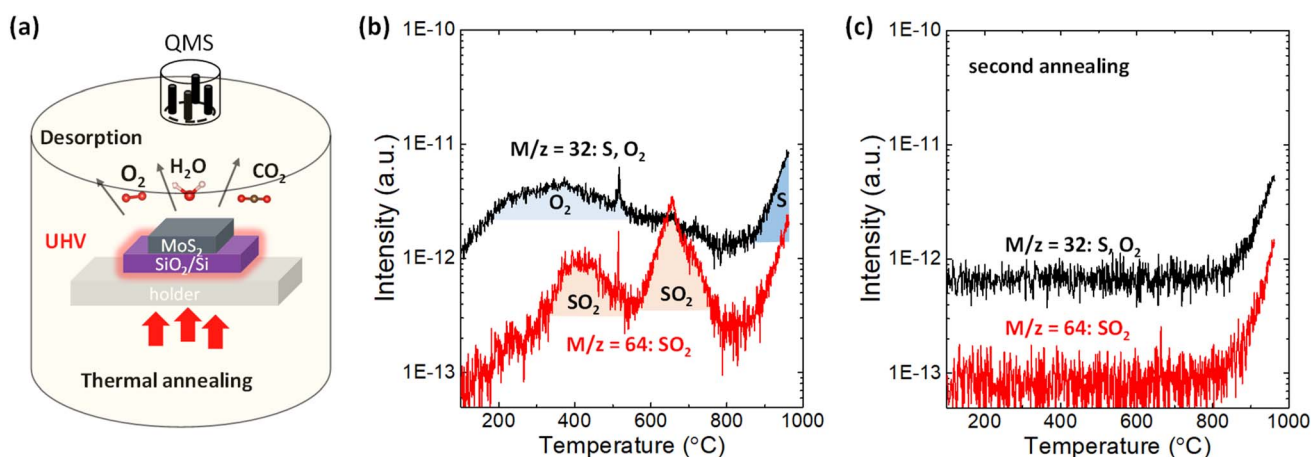


Fig. 1 (a) Illustration of thermal desorption spectroscopy applied in this experiment. (b) TDS spectra with $m/z = 32$ (O₂ or S) and $m/z = 64$ (SO₂) for bulk MoS₂ flakes on the SiO₂/Si substrate. (c) Comparison of the TDS spectra for $m/z = 32$ between the first annealing and second annealing.



species are adsorbed on MoS₂ flakes when MoS₂ flakes are mechanically transferred in ambient air.

The desorption of sulfur and sulfur-related compounds is of the highest interest. Fig. 1b shows the TDS spectra for mass-to-charge ratios of $m/z = 32$ and 64. S and O₂ can be assigned for the same mass-to-charge ratio of $m/z = 32$ because QMS can only filter substances by mass-to-charge ratio. To separate them, the second annealing experiments were successively carried out without removing the sample, as shown in Fig. 1c and S1c.† The broad peak from 200 to 600 °C for $m/z = 32$ is totally removed in the second annealing experiment, with the sharp intensity tail remaining at ~800 °C. It can be expected that direct S desorption is observable even in the second annealing, while O₂ is not detected once it has totally desorbed in the first annealing experiment. Therefore, the broad peak from 200 to 600 °C in Fig. 1b is assigned as the contribution from O₂. Here, it should be noted that the contribution from water adsorbed to the sample was excluded in this discussion, even though the water desorption was observed in Fig. S1.† According to the recent paper,³⁸ no degradation of monolayer MoS₂ was observed after two weeks exposure to the environment of both N₂ and N₂ with the 75% humidity, suggesting that water itself have no dominant contribution for oxidation. Moreover, the desorption tails observed for all species at a high temperature of ~800 °C in Fig. S1† basically resulted from the inner pressure enhancement in the TDS chamber, not from the real increase in desorption. Nevertheless, the TDS spectrum of $m/z = 32$ in the second annealing experiment is greatly enhanced when the surface area of MoS₂ increases by preparing the MoS₂ powder, which is reflected as a threshold temperature shift to the lower side in Fig. S2a.† This indicates that the desorption tail for $m/z = 32$ contains real desorption, that is, the contribution from S. Moreover, Fig. S2b† compares the TDS spectra of chalcogens (S, Se and Te) for various transition metal dichalcogenides. The direct Te ($m/z = 128$) desorption from MoTe₂ is stronger than others, which is consistent with the common understanding of the stability of 2D materials.

In addition to direct S desorption, the sulfur oxidation product, that is, SO₂, was also observed and expressed as $m/z = 64$ in Fig. 1b, where two peaks are observed at ~400 °C and ~650 °C. It should be noted that the $m/z = 64$ was assigned to be the SO₂ rather than S₂ based on the isotope method.³⁹ Sulfur mainly has stable four isotopes: ³²S (95.02%), ³³S (0.5%), ³⁴S (4.21%), and ³⁶S (0.02%). The second abundant ³⁴S ($m = 34$) was used to distinguish S₂ ($m/z = 68$) and SO₂ ($m/z = 66$), as shown in Fig. S3(b).† The desorption signal of $m/z = 66$ well follows $m/z = 64$, while no clear feature was observed for $m/z = 68$. Moreover, from the viewpoint of activation energy of S₂ desorption, the higher temperature is expected. Therefore, $m/z = 64$ was assigned to SO₂ rather than S₂. Furthermore, in the second annealing experiment in Fig. 1c, both peaks disappeared, suggesting that oxygen that constitutes SO₂ comes from adsorbates on the MoS₂ flakes not from inside of the MoS₂ flakes. Therefore, it is considered that adsorbed O₂ reacted with S in MoS₂ and desorbed as SO₂ at ~400 °C and ~650 °C. That is, it is suggested that sulfur vacancy formation at relatively low temperature (~400 °C & 650 °C) is caused by SO₂ desorption

assisted by oxygen adsorption on the MoS₂ basal plane, while it is caused by direct S desorption at high temperature (over 800 °C). Although the theoretical simulation¹⁰ has suggested SO₂ desorption instead of direct S desorption for S vacancy formation, this is the first experimental observation.

The origin of the two SO₂ peaks in Fig. 1b may be related to different intermediates for the final SO₂ formation, as suggested in the theoretical calculation.²⁷ The formation of the intermediate “OSOMo”, in which one oxygen atom is inserted into the Mo–S bond, leads to the desorption of SO₂ with an energy barrier of 1.49 eV. On the other hand, the formation of the intermediate “Mo–OSO”, which can be seen as the SO₂ molecule with one oxygen atom bonded to two adjacent Mo atoms, leads to the desorption of SO and SO₂ with energy barriers of 0.41 eV and 0.78 eV, respectively. Interestingly, in the reproduced TDS experiment in Fig. S3,† it is found that the SO signal highly overlaps with the SO₂ signal in the low-temperature region, where the first peak is observed in Fig. 1b. The overlapping feature suggests that the formation of SO and SO₂ at relatively low temperatures is simultaneous. The slight intense peak for SO at 650 °C can be attributed to the decomposition of SO₂ to SO and O due to the ionization in QMS.³⁹ Although the rigorous separation of SO and SO₂ is difficult at present, two SO₂ peaks may result from the different intermediates for the final SO₂ desorption.

To confirm whether defects form as a result of desorption during TDS annealing, 2 nm Al₂O₃ was deposited by ALD since Al₂O₃ is formed selectively on defect sites due to the dangling-bond free inert surface of the 2D material.^{32–35} For the present ALD-assisted morphology characterization, bulk MoS₂ were intentionally selected to avoid strain-enhanced ALD growth on monolayer MoS₂ due to the transfer process. To show the importance of ALD, the surface topography of MoS₂ annealed up to 835 °C in the TDS chamber without ALD was examined to assess the resolution of AFM. As shown in Fig. S4,† no clear defects could be identified. The situation was inverted when ALD was applied to the MoS₂ flake after TDS annealing. As shown in Fig. 2b, for the MoS₂ flake annealed up to 200 °C in the TDS chamber, the growth of Al₂O₃ is distributed in dots on the MoS₂ surface with a greater concentration at the step edge. It should be noted that 200 °C was selected as the starting temperature for defect characterization since the Al₂O₃ growth temperature in ALD was 200 °C. As the TDS annealing temperature increases, Al₂O₃ growth is enhanced, occupying more surface area of MoS₂ (Fig. 2b–e). For the MoS₂ flake annealed up to 630 °C, the growth of Al₂O₃ is no longer preferential at the grain boundaries and the step edges but uniformly covers the entire MoS₂ surface. The morphology evolution is schematically illustrated below the AFM images.

Next, the Al₂O₃ coverage was extracted from the AFM images and plotted as a function of annealing temperature in Fig. 2a. The coverage rate is relatively slow in the temperature range from 350 to 500 °C, where the first peak of SO₂ desorption is observed in Fig. 1b. In contrast, there was a significant increase in the coverage rate from 550 to 630 °C. This temperature interval matches the second peak of SO₂ desorption. These results suggest that the evolution of the Al₂O₃ coverage is correlated to SO₂ desorption. However, further quantitative



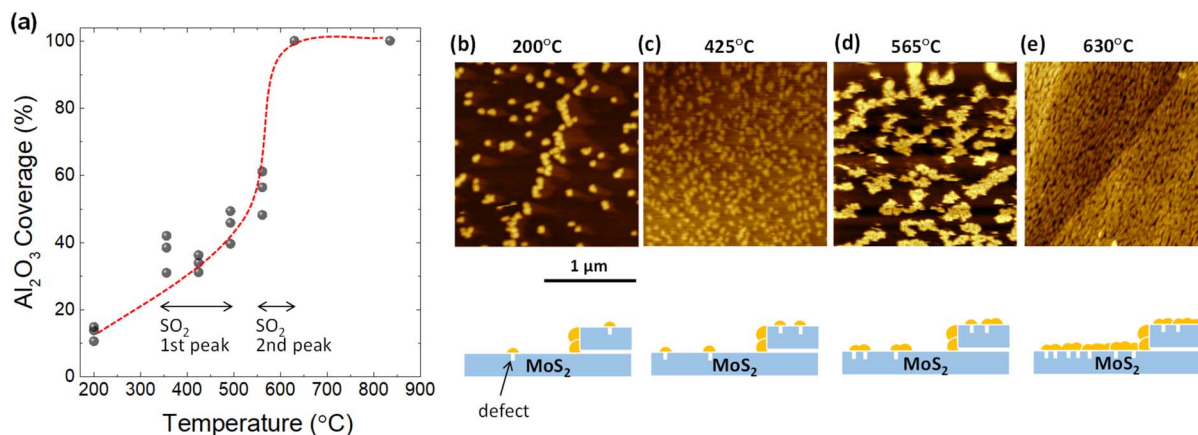


Fig. 2 (a) Coverage estimation of ALD- Al_2O_3 on TDS-annealed bulk MoS_2 flakes with varying annealing temperatures. The gray point represents the coverage extracted from different AFM images. The red dashed line is an eye guide to reflect the coverage evolution rate. (b)–(e) Representative AFM images for $\text{Al}_2\text{O}_3/\text{MoS}_2$ at annealing temperatures of 200 °C, 425 °C, 565 °C, and 630 °C, respectively. In addition, the Al_2O_3 coverage evolution with increasing annealing temperature is schematically illustrated. The yellow particles represent deposited Al_2O_3 .

estimation of the defect density is not possible here since the Al_2O_3 particle density does not form a one-to-one correspondence with the defect density. That is, new Al_2O_3 growth occurs not only at the defect site but also on the Al_2O_3 that has already been grown. However, the importance of ALD-assisted AFM characterization is to provide clear evidence that the formation of defects on the basal plane of MoS_2 flakes does arise when SO_2 desorption occurs.

To reveal the structural change in MoS_2 after TDS annealing, Raman measurements were conducted. Fig. 3a shows the Raman spectra of monolayer MoS_2 at different TDS annealing temperatures. The frequency difference of E_{2g}^1 and A_{1g} is $\sim 19.4 \text{ cm}^{-1}$, which matches well with the characteristics of

monolayer MoS_2 .⁴⁰ No detectable shifts of the E_{2g}^1 and A_{1g} peaks were found for any temperature range. On the other hand, the full width at half maximum (FWHM) of these two peaks shown in Fig. 3b is slightly broadened when the annealing temperature was elevated to 835 °C; at this temperature, direct S desorption becomes dominant in the TDS spectra of Fig. 1b. This indicates that some damage to the crystallinity is induced by considerable direct S desorption. The change in Raman spectra at temperatures above 800 °C has also been reported.⁴¹ Although Raman spectra are found to be insensitive to temperatures below 630 °C, this is not consistent with Fig. 2a, where a clear Al_2O_3 coverage evolution is observed. In Fig. S3,[†] the desorption of SO_2 and S was detected but not for MoO , MoO_2 , and MoO_3 ,

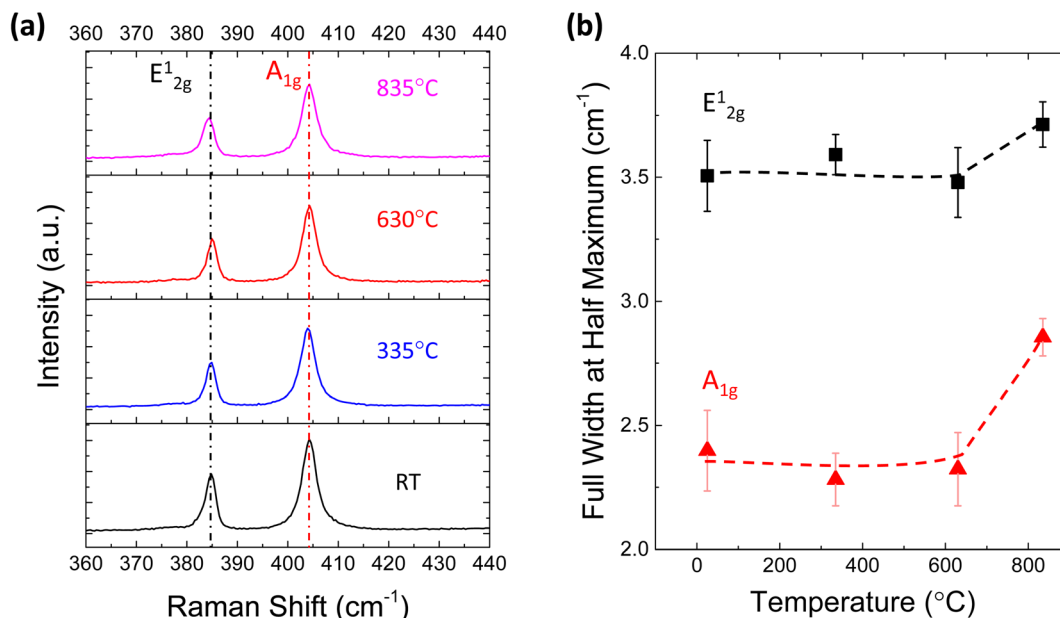


Fig. 3 (a) Raman spectra for monolayer MoS_2 annealed at different temperatures in the TDS chamber. (b) Full width at half maximum (FWHM) of the Raman peaks at different temperatures.



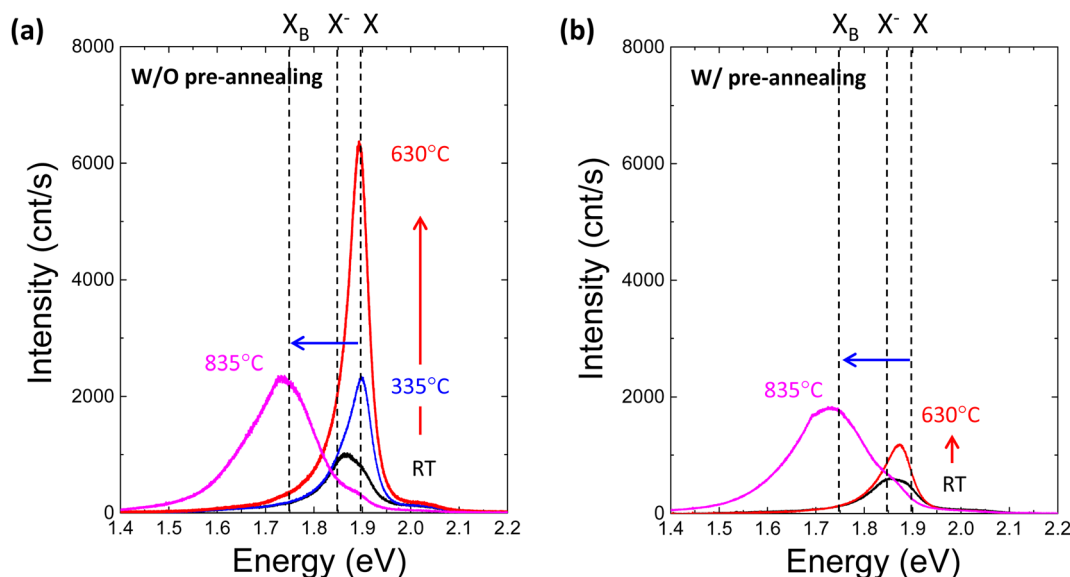


Fig. 4 (a) Photoluminescence spectra for monolayer MoS₂ annealed at different temperatures in the TDS chamber without preannealing in Ar. X, X⁻ and X_B represent exciton, trion and defect-induced bound exciton, respectively. The red arrow indicates the intensity enhancement with increasing annealing temperature from RT to 630 °C, while the blue arrow indicates the transition from exciton to defect-induced bound exciton due to high-temperature annealing at 835 °C. (b) Photoluminescence spectra for monolayer MoS₂ annealed at different temperatures in the TDS chamber with preannealing in Ar.

suggesting that the original MoS₂ crystal lattice can be retained since Mo is located at the center of the MoS₂ crystal structure. Therefore, Raman spectra could be insensitive to the relatively small change in crystallinity due to S vacancy formation accompanied by SO₂ desorption.

To further investigate the effect of defect formation more sensitively, PL measurements were performed for monolayer MoS₂ annealed at different temperatures in the TDS chamber. As shown in Fig. 4a, the conventional broad PL peak mixed with exciton (X, ~1.9 eV) and trion (X⁻, ~1.85 eV) was observed.^{42–44} As the TDS annealing temperature increased from RT to 630 °C, the intensity of the exciton peak drastically increased. Since the MoS₂ flakes were exposed to ambient air during the PL measurement after removing the sample from the TDS chamber, the exciton intensity enhancement is attributed to the oxygen chemical adsorption on S vacancies created by SO₂ desorption during TDS annealing.⁴³ As shown in Fig. S5,[†] the trion contribution was reduced at elevated temperature, while the exciton component became dominant. This indicates the p-type doping by oxygen passivation at the defect site.⁴⁵ Moreover, when the annealing temperature increased to 835 °C, the defect-induced bound exciton peak (X_B, ~1.75 eV) became dominant.^{46–49} Although SO₂ formation at relatively low temperatures is limited by the amount of oxygen adsorbed initially on MoS₂ flakes, direct S desorption at high temperatures depends only on the temperature. Therefore, a much broader X_B peak implies a significant amount of S vacancy formation and clustering due to the enhanced direct S desorption, which is also supported by the broadening of the FWHM in the Raman peaks (Fig. 3b).

Four different types of experiments, namely, of TDS, AFM, Raman and PL, surprisingly, can be well explained by the two

common kinds of defect formation mechanisms in the low and high annealing temperature regions. That is, at a low annealing temperature, the oxygen adsorbed on the basal plane of MoS₂ flakes is the key factor, and it volatilizes as SO₂ by removing S from MoS₂, leaving S vacancies. Then, at a high annealing temperature, direct S desorption becomes dominant. Here, it is reasonable to consider that defect formation at low annealing temperatures could be suppressed by removing oxygen adsorbed on MoS₂ flakes before TDS annealing. To verify this idea, the MoS₂ crystals were preannealed at 500 °C for one hour in an Ar gas flow to remove oxygen adsorbed on the MoS₂ surface. Then, MoS₂ flakes were transferred onto the SiO₂/Si substrate by mechanical exfoliation in ambient air, followed by TDS annealing. The PL spectra of preannealed monolayer MoS₂ are shown in Fig. 4b. The PL intensity enhancement at temperatures below 630 °C is greatly suppressed, which strongly indicates that S vacancy formation due to SO₂ desorption was hindered due to the great reduction in oxygen adsorbed on the MoS₂ surface. On the other hand, at a high temperature of 835 °C, the X_B peak was almost the same, which is also evidence for direct S desorption. These results clearly support the two kinds of defect formation mechanisms in the low and high annealing temperature regions. To date, S vacancies have been experimentally observed and recognized as dominant defects because they have the lowest formation energy of ~1.3–1.5 eV compared with other types of defects.^{3,4} However, this large formation energy could not explain the high S vacancy concentration of ~10¹³ cm⁻².⁵⁰ The present study experimentally proves that oxygen adsorbed on MoS₂ assists S vacancy formation by desorbing as SO₂ since this O₂-assisted S vacancy formation is energetically spontaneous at -0.49 eV.¹⁰



4. Conclusions

Through this experimental desorption study using TDS in combination with ALD, Raman, and PL, two kinds of S vacancy formation mechanisms are realized, that is, O₂-assisted SO₂ desorption at low annealing temperatures (400 °C to 650 °C) and direct S desorption at high annealing temperatures (over 800 °C). The key finding is that the initially adsorbed oxygen causes S vacancy formation through SO₂ desorption. As a solution, the removal of oxygen is quite effective for preventing the formation of defects and further oxidation.

Conflicts of interest

There are no conflicts to declare.

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References

- 1 S. B. Desai, S. R. Madhupathy, A. B. Sachid, J. P. Llinas, Q. Wang, G. H. Ahn, G. Pitner, M. J. Kim, J. Bokor, C. Hu, H.-S. P. Wong and A. Javey, *Science*, 2016, **354**, 99–102.
- 2 Y. Liu, X. Duan, H.-J. Shin, S. Park, Y. Huang and X. Duan, *Nature*, 2021, **591**, 43–53.
- 3 J. Hong, Z. Hu, M. Probert, K. Li, D. Lv, X. Yang, L. Gu, N. Mao, Q. Feng, L. Xie, J. Zhang, D. Wu, Z. Zhang, C. Jin, W. Ji, X. Zhang, J. Yuan and Z. Zhang, *Nat. Commun.*, 2015, **6**, 6293.
- 4 W. Zhou, X. Zou, S. Najmaei, Z. Liu, Y. Shi, J. Kong, J. Lou, P. M. Ajayan, B. I. Yakobson and J.-C. Idrobo, *Nano Lett.*, 2013, **13**, 2615–2622.
- 5 S. McDonnell, R. Addou, C. Buie, R. M. Wallace and C. L. Hinkle, *ACS Nano*, 2014, **8**, 2880–2888.
- 6 H. Qiu, T. Xu, Z. Wang, W. Ren, H. Nan, Z. Ni, Q. Chen, S. Yuan, F. Miao, F. Song, G. Long, Y. Shi, L. Sun, J. Wang and X. Wang, *Nat. Commun.*, 2013, **4**, 2642.
- 7 N. Fang and K. Nagashio, *J. Phys. D: Appl. Phys.*, 2018, **51**, 65110.
- 8 N. Fang, S. Toyoda, T. Taniguchi, K. Watanabe and K. Nagashio, *Adv. Funct. Mater.*, 2019, **29**, 1904465.
- 9 Z. Yu, Y. Pan, Y. Shen, Z. Wang, Z.-Y. Ong, T. Xu, R. Xin, L. Pan, B. Wang, L. Sun, J. Wang, G. Zhang, Y. W. Zhang, Y. Shi and X. Wang, *Nat. Commun.*, 2014, **5**, 5290.
- 10 J. Pető, T. Ollár, P. Vancsó, Z. I. Popov, G. Z. Magda, G. Dobrik, C. Hwang, P. B. Sorokin and L. Tapasztó, *Nat. Chem.*, 2018, **10**, 1246–1251.
- 11 T. Yanase, F. Uehara, I. Naito, T. Nagahama and T. Shimada, *ACS Appl. Nano Mater.*, 2020, **3**, 10462–10469.
- 12 M. Makarova, Y. Okawa and M. Aono, *J. Phys. Chem. C*, 2012, **116**, 22411–22416.
- 13 D. Rhodes, S. H. Chae, R. Ribeiro-Palau and J. Hone, *Nat. Mater.*, 2019, **18**, 541–549.
- 14 Q. Liang, Q. Zhang, X. Zhao, M. Liu and A. T. S. Wee, *ACS Nano*, 2021, **15**, 2165–2181.
- 15 Y.-R. Chang, T. Nishimura and K. Nagashio, *ACS Appl. Mater. Interfaces*, 2021, **13**, 43282–43289.
- 16 J. Gao, B. Li, J. Tan, P. Chow, T.-M. Lu and N. Koratkar, *ACS Nano*, 2016, **10**, 2628–2635.
- 17 D. Lv, H. Wang, D. Zhu, J. Lin, G. Yin, F. Lin, Z. Zhang and C. Jin, *Sci. Bull.*, 2017, **62**, 846–851.
- 18 R. C. Longo, R. Addou, S. KC, J.-Y. Noh, C. M. Smyth, D. Barrera, C. Zhang, J. W. P. Hsu, R. M. Wallace and K. Cho, *2D Mater.*, 2017, **4**, 025050.
- 19 J. Martincová, M. Otyepka and P. Lazar, *Chem.-Eur. J.*, 2017, **23**, 13233–13239.
- 20 Y. Li, Z. Zhou, S. Zhang and Z. Chen, *J. Am. Chem. Soc.*, 2008, **130**, 16739–16744.
- 21 M. Yamamoto, T. L. Einstein, M. S. Fuhrer and W. G. Cullen, *J. Phys. Chem. C*, 2013, **117**, 25643–25649.
- 22 H. Zhu, X. Qin, L. Cheng, A. Azcatl, J. Kim and R. M. Wallace, *ACS Appl. Mater. Interfaces*, 2016, **8**, 19119–19126.
- 23 J. Wu, H. Li, Z. Yin, H. Li, J. Liu, X. Cao, Q. Zhang and H. Zhang, *Small*, 2013, **9**, 3314–3319.
- 24 T. N. Walter, F. Kwok, H. Simchi, H. M. Aldosari and S. E. Mohny, *J. Vac. Sci. Technol., B*, 2017, **35**, 21203.
- 25 M. Liu, J. Shi, Y. Li, X. Zhou, D. Ma, Y. Qi, Y. Zhang and Z. Liu, *Small*, 2017, **13**, 1602967.
- 26 Q. Chen, H. Li, W. Xu, S. Wang, H. Sawada, C. S. Allen, A. I. Kirkland, J. C. Grossman and J. H. Warner, *Nano Lett.*, 2017, **17**, 5502–5507.
- 27 L. M. Farigliano, P. A. Paredes-Olivera and E. M. Patrito, *J. Phys. Chem. C*, 2020, **124**, 13177–13186.
- 28 N. Ikegami, N. Ozawa, Y. Miyakawa, N. Hirashita and J. Kanamori, *Jpn. J. Appl. Phys.*, 1992, **31**, 2020–2024.
- 29 W. Widdra, S. I. Yi, R. Maboudian, G. A. D. Briggs and W. H. Weinberg, *Phys. Rev. Lett.*, 1995, **74**, 2074–2077.
- 30 S. K. Wang, K. Kita, C. H. Lee, T. Tabata, T. Nishimura, K. Nagashio and A. Toriumi, *J. Appl. Phys.*, 2010, **108**, 54104.
- 31 Y.-R. Chang, N. Higashitarumizu, H. Kawamoto, F.-H. Chu, C.-J. Lee, T. Nishimura, R. Xiang, W.-H. Chang, S. Maruyama and K. Nagashio, *Chem. Mater.*, 2021, **33**, 186–194.
- 32 N. Takahashi, K. Watanabe, T. Taniguchi and K. Nagashio, *Nanotechnology*, 2015, **26**, 175708.
- 33 Y. Xuan, Y. Q. Wu, T. Shen, M. Qi, M. A. Capano, J. A. Cooper and P. D. Ye, *Appl. Phys. Lett.*, 2008, **92**, 13101.
- 34 H. G. Kim and H.-B.-R. Lee, *Chem. Mater.*, 2017, **29**, 3809–3826.
- 35 X. Wang, S. M. Tabakman and H. Dai, *J. Am. Chem. Soc.*, 2008, **130**, 8152–8153.
- 36 K. Ueno, *J. Phys. Soc. Jpn.*, 2015, **84**, 121015.
- 37 K. Nagashio, T. Yamashita, T. Nishimura, K. Kita and A. Toriumi, *J. Appl. Phys.*, 2011, **110**, 24513.
- 38 Y. Song, M. Park, J. Park, H. S. Ahn, T. K. Kim and S.-Y. Ju, *Nanomaterials*, 2022, **12**.



- 39 D. E. Pierce, R. P. Burns, H. M. Dauplaise and L. J. Mizerka, *Tribol. Trans.*, 1991, **34**, 205–214.
- 40 C. Lee, H. Yan, L. E. Brus, T. F. Heinz, J. Hone and S. Ryu, *ACS Nano*, 2010, **4**, 2695–2700.
- 41 L. Wang, X. Ji, F. Chen and Q. Zhang, *J. Mater. Chem. C*, 2017, **5**, 11138–11143.
- 42 K. F. Mak, C. Lee, J. Hone, J. Shan and T. F. Heinz, *Phys. Rev. Lett.*, 2010, **105**, 136805.
- 43 H. Nan, Z. Wang, W. Wang, Z. Liang, Y. Lu, Q. Chen, D. He, P. Tan, F. Miao, X. Wang, J. Wang and Z. Ni, *ACS Nano*, 2014, **8**, 5738–5745.
- 44 S. Mouri, Y. Miyauchi and K. Matsuda, *Nano Lett.*, 2013, **13**, 5944–5948.
- 45 D. Mastrippolito, S. Palleschi, G. D'Olimpio, A. Politano, M. Nardone, P. Benassi and L. Ottaviano, *Nanoscale*, 2020, **12**, 18899–18907.
- 46 S. Tongay, J. Suh, C. Ataca, W. Fan, A. Luce, J. S. Kang, J. Liu, C. Ko, R. Raghunathanan, J. Zhou, F. Ogletree, J. Li, J. C. Grossman and J. Wu, *Sci. Rep.*, 2013, **3**, 2657.
- 47 E. Mitterreiter, B. Schuler, A. Micevic, D. Hernangómez-Pérez, K. Barthelmi, K. A. Cochrane, J. Kiemle, F. Sigger, J. Klein, E. Wong, E. S. Barnard, K. Watanabe, T. Taniguchi, M. Lorke, F. Jahnke, J. J. Finley, A. M. Schwartzberg, D. Y. Qiu, S. Refaely-Abramson, A. W. Holleitner, A. Weber-Bargioni and C. Kastl, *Nat. Commun.*, 2021, **12**, 3822.
- 48 K. Greben, S. Arora, M. G. Harats and K. I. Bolotin, *Nano Lett.*, 2020, **20**, 2544–2550.
- 49 N. Fang, K. Nagashio and A. Toriumi, *2D Mater.*, 2016, **4**, 15035.
- 50 Z. G. Yu, Y.-W. Zhang and B. I. Yakobson, *Nano Lett.*, 2015, **15**, 6855–6861.

