



Cite this: *EES Catal.*, 2023,  
1, 392

## Strategies for reducing the overpotential of one-dimensional Si nanostructured photoelectrodes for solar hydrogen production

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Using solar energy to split water to produce hydrogen, the most promising green energy source, is an advanced way to solve the current energy crisis and environmental pollution problems. The key to this technology lies in the development of high-performance semiconductor photoelectrodes. Si has been recognized as a promising candidate for solar water splitting photoelectrodes due to its excellent light absorption properties (band gap of 1.1 eV), high carrier mobility, abundance on earth, and established manufacturing technology. To improve the performance of Si-based photoelectrodes, one-dimensional (1D) Si nanostructures have been constructed and investigated. 1D Si nanostructured photoelectrodes have many advantages over planar Si photoelectrodes, including a large specific surface area that provides more reaction sites, antireflection effect that greatly enhances the absorption of sunlight, and improved carrier utilization by decoupling the light absorption and carrier collection paths. However, the poor catalytic activity of Si itself for the water splitting reactions (oxygen evolution reaction and hydrogen evolution reaction) makes the photoelectrocatalytic (PEC) water splitting on Si-based photoelectrodes require a large overpotential. In addition, the narrow band gap of Si leads to a low photovoltage, and the valence band energy level of Si is insufficient to oxidize water, which also leads to an increased overpotential for the PEC water splitting by the Si-based photoelectrodes. In order to reduce the overpotential of 1D Si nanostructured photoelectrodes for PEC water splitting, several strategies have been proposed, such as loading co-catalysts and constructing hetero- and homo-junctions. After introducing PEC water splitting process on Si, and the preparation methods and advantages of 1D Si nanostructures, this review paper provides an overview of the strategies and recent advances in reducing the overpotential of 1D Si nanostructured photoelectrodes for PEC water splitting.

Received 4th April 2023,  
Accepted 16th May 2023

DOI: 10.1039/d3ey00073g

rsc.li/eescatalysis

### Broader context

Hydrogen energy is one of the most promising future clean energy sources. One-dimensional (1D) Si nanostructures can be used as photoelectrodes to produce hydrogen through the solar-driven water splitting process, providing a sustainable production pathway for hydrogen energy. However, the 1D Si nanostructured photoelectrodes still face the problem of high overpotential in photoelectrocatalytic water splitting due to some inherent factors of Si, such as poor catalytic activity, narrow bandgap, and unsuitable valence band position. This review focuses on the strategies and recent research progress in reducing the overpotential of 1D Si nanostructured photoelectrodes. The reducing of the overpotential is a common problem in the research of photoelectrocatalytic materials. In addition to 1D silicon nanostructures, the contents of this review will be equally applicable and instructive for other semiconductor photoelectrodes. Therefore, this paper will be of broad interest to a broader group of researchers in the field of photoelectrocatalysis.

## Introduction

Energy crisis and environmental pollution are two serious issues that hinder social progress and endanger human health.<sup>1</sup> With the development of human society, the energy demand is becoming increasingly high.<sup>2,3</sup> The world's energy supply is still mainly from non-renewable fossil energy sources such as oil, coal, and natural gas.<sup>4</sup> These energy sources have limited reserves on Earth

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and will continually diminish as we use them.<sup>5</sup> The uneven distribution of these energy sources across the planet and the struggle over their development, control, use, and pricing are the root causes of a series of international disputes today. Furthermore, these energy sources release energy into electrical, thermal, and chemical forms mainly through combustion. The inevitable chemical products such as dust, sulfide, nitrogen compounds, and carbon dioxide formed during the combustion process are the leading cause of environmental pollution, such as haze and the greenhouse effect, are also responsible for increasing human respiratory and skin diseases. For these reasons, using renewable energy sources, which are abundant or even inexhaustible, is considered an alternative to traditional fossil energy sources. How to improve the utilization efficiency of renewable energy and minimise environmental pollution in the process of energy utilisation has become a vital issue that must be addressed in the pursuit of sustainable development in human society.

Solar energy represents the most reasonable choice among the many renewable resources such as solar, wind, tidal, nuclear, and bioenergy.<sup>6</sup> There have been various approaches developed for the utilization of solar energy. One method is directly collecting solar radiation and converting it into thermal energy through interaction with matter.<sup>7</sup> Another common method is to convert solar energy into electrical energy using the photovoltaic effect, also known as the solar cell.<sup>8</sup> In addition, solar energy can also be converted into chemical energy, biological energy, and so on.<sup>9,10</sup> Of particular interest has been the use of solar energy to achieve PEC water splitting to produce hydrogen,<sup>11–14</sup> as hydrogen energy is considered to be one of the most promising future energy due to its high calorific value and non-polluting combustion products.<sup>15–17</sup> Semiconductor photocatalysts are the key to achieving PEC water splitting for hydrogen production.<sup>18</sup> In 1972, Fujishima *et al.* first achieved water splitting using semiconductor TiO<sub>2</sub> under illumination.<sup>19</sup> Due to the wide band gap of TiO<sub>2</sub>, it can only be excited by ultraviolet light, which accounts for just a tiny fraction of the entire solar spectrum. Therefore, narrow bandgap semiconductors have been used for PEC water splitting and hydrogen production, such as CdS,<sup>20,21</sup> Fe<sub>2</sub>O<sub>3</sub>,<sup>22,23</sup> BiVO<sub>4</sub>,<sup>24,25</sup> GaAs,<sup>26,27</sup> Si,<sup>28–31</sup> *etc.* Due to the advantages such as proper band structure and band gap ( $E_g = 1.1$  eV), high carrier mobility, and abundance on Earth, Si has emerged as a promising candidate for solar hydrogen production.<sup>30,32</sup> The design and construction of Si nanostructures can help greatly enhance the efficiency of solar water splitting for hydrogen production.<sup>33,34</sup> Among the various nanomaterials, one-dimensional (1D) nanomaterials have attracted tremendous attention. In particular, ordered arrays of 1D nanostructures, such as 1D nanowire arrays, nanorod arrays, and nanotube arrays, have shown great potential for PEC water splitting.<sup>35–38</sup> In this regard, 1D Si nanostructured photoelectrodes for solar water splitting to generate hydrogen have aroused much interest in recent years.<sup>39,40</sup>

Nevertheless, 1D Si nanostructures have some shortcomings in PEC water splitting applications, one of the important ones being the high overpotential for water splitting. First, the poor catalytic activity of Si for water splitting reactions (HER and

OER) results in sluggish kinetics on the surface of 1D Si nanostructures, and a large overpotential is required to achieve water splitting.<sup>41,42</sup> Secondly, the narrow bandgap of Si gives rise to a low photovoltage, thus requiring a large external voltage to split the water.<sup>43</sup> Several strategies have been developed to reduce the overpotential of PEC water splitting on 1D Si nanostructured photoelectrodes. In this review paper, we first briefly introduce the basic principles of PEC water splitting on Si surfaces, summarize the preparation methods of 1D Si nanostructures, elucidate the advantages of 1D Si nanostructured photoelectrodes compared with planar photoelectrodes, then review the research progress in reducing the overpotential of PEC hydrogen production on 1D Si nanostructured photoelectrodes by modifying co-catalysts and constructing hetero- or homo-junctions, and finally present some of our envisions for this field. It should be noted that some of the examples may be beyond the scope of 1D Si nanostructures, most typically the Si microwire arrays exemplified several times. Since Si microwires and Si nanowires have similar morphological characteristics, the difference being the diameter of the wires, the conclusions drawn with Si microwires would also apply to Si nanowires.

## PEC water splitting on Si

The complete water splitting reaction consists of the cathodic hydrogen evolution reaction (HER) and the anodic oxygen evolution reaction (OER).



The free energy change for splitting water to form H<sub>2</sub> and 1/2O<sub>2</sub> is + 237 kJ mol<sup>-1</sup>, corresponding to an electric potential change  $\Delta E = 1.23$  V.<sup>44</sup> This means that a voltage of at least 1.23 V is applied to split water. The standard reduction potentials of HER and OER are 0 V and 1.23 V vs. SHE, respectively. However, for some reasons, the actual potential of the electrode reaction exceeds the standard electrode potential. The difference between actual potential and reversible potential is called overpotential ( $\eta$ ), which produces an increased thermodynamic driving force for the electrochemical process.<sup>45</sup> Both the cathodic half-reaction and the anodic half-reaction have an overpotential. The overpotential ( $\eta$ ) of an electrochemical reaction is the sum of the cathodic overpotential ( $\eta_c$ ) and the anodic overpotential ( $\eta_a$ ). The electrode overpotential can be divided into its component overpotentials as in the following equation:<sup>45</sup>

$$\eta = \eta_y + \eta_k + \eta_e + \eta_t$$

where the component overpotentials are:  $\eta_y$ , electrolyte resistance;  $\eta_k$ , activation or reaction kinetics limitation;  $\eta_e$ , electronic resistance of electrodes; and  $\eta_t$ , concentration loss or mass transfer limitation. Taking into account the overpotential, the actual voltage required to split water is 1.6–1.8 V.<sup>30,46</sup>



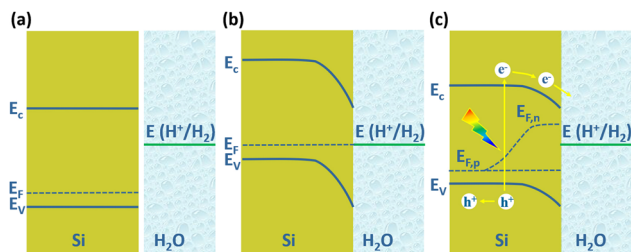


Fig. 1 Schematic diagram of Si/electrolyte interface. (a) Before contact, (b) equilibrium reached after contact in the dark, (c) under illumination.

The PEC water splitting reaction on the Si photoelectrode surface is associated with the energy band structure of the Si/electrolyte interface, as shown in Fig. 1 for the example of p-Si. The valence band edge energy ( $E_v$ ) and conduction band edge energy ( $E_c$ ) levels of Si are  $-5.2$  eV and  $-4.1$  eV, corresponding to electrochemical potentials of  $0.7$  V and  $-0.4$  V (vs. NHE), respectively.<sup>47</sup> In the case of p-Si, the Fermi level ( $E_F$ ) is near the top of the valence band (Fig. 1a). When p-Si is in contact with the electrolyte, since the  $E_F$  of p-Si is lower than that of the electrolyte, electrons will transfer from the electrolyte to p-Si until the Fermi levels of the two are equal. This process leads to the formation of net electron space-charge regions on the p-Si surface, causing the energy bands on the surface to bend downward (Fig. 1b). For n-Si, the process is just the opposite, with the energy band on the n-Si surface bending upward. Under illumination, electrons in the valence band of p-Si are excited to the conduction band, causing the equilibrium between p-Si and electrolyte to be disrupted. The departure from equilibrium is quantified in terms of the quasi-Fermi level, *i.e.*, electron quasi-Fermi level ( $E_{F,n}$ ) and hole quasi-Fermi level ( $E_{F,p}$ ) (Fig. 1c). The photogenerated electrons are transferred to the solution driven by the built-in electric field to reduce  $H^+$  to produce  $H_2$ .

When a beam of photons are incident on a semiconductor, in addition to reflection and transmission, the remaining photons are absorbed by the semiconductor and excite photogenerated carriers (electron-hole pairs). Some of the photogenerated carriers may be lost during transport to the semiconductor's surface as a result of recombination, and the remainder can be collected on the surface of the semiconductor and transferred to the external circuit or to the solution. Fig. 2 depicts the collection of photogenerated carriers within p-Si near the p-Si/electrolyte interface.<sup>48</sup> The light penetration depth (given by  $1/\alpha$ , where  $\alpha$  is the absorption coefficient) in a semiconductor is related to the wavelength of the incident light and the band structure of the semiconductor. The light penetration depth of direct bandgap semiconductors is relatively shallow, while that of indirect bandgap semiconductors is relatively deep. For the semiconductor Si with an indirect bandgap, visible light penetrates within it to a depth of  $100 \mu\text{m}$ .<sup>49</sup> The diffusion length of the minority carrier ( $L_n$ ) refers to the average length that the photogenerated hole travels from its creation to recombination with the electron. The  $L_n$  is related to the lattice defects and doping of the semiconductor.



Fig. 2 Schematic diagram of the carrier transportation process at the p-Si/electrolyte interface.

The more lattice defects and the higher the concentration of impurities, the higher the chance of recombination of photogenerated carriers, the shorter the lifetime and the shorter the diffusion length. The width of the space charge region (or the depletion layer) ( $W$ ) depends on the type and density of dopants in Si and the potential bias.

In the space-charge region (region I), photogenerated electrons and holes are spatially separated by the built-in electric field and transported by drift. Although some photogenerated carriers recombine, the collection efficiency of photogenerated carriers in this region is high. In region II, photogenerated carrier migration mainly relies on diffusion due to concentration differences. A part of the photogenerated carriers will recombine, and the remaining part can diffuse to the space charge region (region I), where the collection efficiency of photogenerated carriers is also relatively high. In region III beyond the  $L_n + W$ , most of the photogenerated minority carrier electrons will be recombined with holes and lost during transport to the semiconductor surface, and the collection efficiency of photogenerated carriers in this region is very low.  $L_n + W$  determines the effective collection length of photogenerated carriers. In region IV, the arriving light intensity is very weak, and very few photogenerated carriers are produced.

Another critical step in photogenerated carrier collection is the transfer of photogenerated carriers that reach the surface of the semiconductor to the solution. In addition to the type and concentration of redox pairs in solution and the ambient temperature, the surface states of the semiconductor can have an essential effect on photogenerated carrier collection here, the defect states of the semiconductor surface act as recombination centres for the photogenerated carriers. Even if the photogenerated carriers reach the semiconductor surface, some may recombine at the surface. In addition, regarding electrochemical reaction kinetics, the charge transfer from the electrode to the solution needs to overcome a certain energy potential barrier and requires an overpotential. The electrocatalytic activity of the semiconductor surface directly affects



the magnitude of the overpotential required for charge transfer at the semiconductor/electrolyte interface.

## Preparation of 1D Si nanostructures

Due to the unique physical and chemical properties of Si, 1D Si nanostructures have been extensively studied and a variety of preparation methods have been developed. For solar hydrogen production applications, ordered arrays of 1D Si nanostructures can be easily and directly fabricated into photoelectrodes and fully utilize the merits of 1D nanostructures. Therefore, we focus here on ordered 1D silicon nanostructure arrays and briefly review some typical preparation methods.

### Metal-assisted chemical etching

Metal-assisted chemical etching (MACE) has been widely used to prepare 1D Si nanostructure arrays due to its simplicity, low cost, and ease of scale-up preparation. Based on previous studies, Peng *et al.* developed a method for the preparation of Si nanowire arrays by MACE. The preparation of Si nanowires by MACE is using single-crystal Si wafer as the raw material. First, metal (Ag, Au, Pt, Cu, *etc.*) nanoparticles or thin films are deposited on the Si wafer by chemical or physical methods; then, Si is etched in an oxidizing HF solution containing oxidizing agents (nitric acid and nitrate,  $\text{H}_2\text{O}_2$ , *etc.*) with the metal as a catalyst.<sup>50,51</sup> The metal catalyst allows the etching of Si to proceed in a specific direction to form straight holes. When the adjacent holes penetrate each other, the remaining hole walls take on a 1D structural morphology. Ag nanoparticles are the most commonly used catalyst for the preparation of 1D Si nanostructures by the MACE method. Based on the experimental observations, Peng *et al.* proposed an

electrokinetic model to explain the microscopic dynamic origin of motility of Ag nanoparticles in Si during the chemical etching process, as shown in Fig. 3a.<sup>52</sup> The energy generated in the *in situ* bipolar electrochemical reaction on the surface of the Ag nanoparticles can be directly converted into mechanical work, which drives the movement of the Ag nanoparticles in Si with the help of HF.

A very simple method to prepare metal catalysts is to reduce metal ions in solution into metal nanoparticles using the reducing Si-H bonds on the surface of HF-treated Si wafer. Subsequently, the ordered Si nanowire arrays can be obtained by MACE with the catalytic effect of metal catalysts. The benefits of the Si-H bond reduction method for the preparation of metal catalysts are simplicity, low cost and short time consumption, but the method cannot precisely control the morphology of the metal catalysts, resulting in the low controllability of the prepared Si nanowires (Fig. 3b). Therefore, Peng *et al.* proposed a method to prepare regular Ag nanoholes on Si sheets by physical vapor deposition using silica nanosphere as a template for lithography.<sup>53</sup> The Si nanowire arrays prepared with such Ag nanoholes as catalysts were highly ordered, as shown in Fig. 3c. Using this method, the diameter, spacing, and alignment of Si nanowires can be easily and precisely tuned by the adjustment of the lithography pattern. In addition to Si nanowires perpendicular to the substrate, tilted Si nanowires,<sup>52</sup> zigzag Si nanowires,<sup>54,55</sup> Si nanoholes,<sup>56</sup> *etc.* have also been prepared by changing the MACE process conditions.

### Chemical vapor deposition

Chemical vapor deposition (CVD) is the condensation of one or more compounds from the gas phase onto a substrate where a



**Fig. 3** (a) Catalytic-induced electrokinetics for autonomous Ag particle movement in bulk Si. SEM images of 1D Si nanostructures prepared by MACE. (b) Perpendicular Si nanowire arrays. (c) Highly order Si nanowire arrays, (d) slanted Si nanowire arrays, reproduced with permission from ref. 52, Copyright © 2008 Wiley-VCH Verlag GmbH & Co. KGaA, Weinheim. and (e) highly ordered Si nanohole arrays. (a and d) Reproduced with permission from ref. 52, Copyright © 2008 Wiley-VCH Verlag GmbH & Co. KGaA, Weinheim. (b) Reproduced with permission from ref. 51, Copyright © 2013 Wiley-VCH Verlag GmbH & Co. KGaA, Weinheim. (c) Reproduced with permission from ref. 53, © 2007 American Institute of Physics. (e) Reproduced with permission from ref. 56, Copyright © 2010, American Chemical Society.





reaction occurs to produce a solid deposit.<sup>57</sup> In the CVD process, the deposition of solids can follow different mechanisms to form different morphologies. Among them, vapour-liquid-solid (VLS) is a mechanism widely used to prepare 1D nanostructures. 1D Si nanostructures can be synthesized “bottom-up” by employing metal nanoclusters as catalysts *via* a VLS mechanism.<sup>58,59</sup> Fig. 4a illustrates the growth mechanism of a Si nanowire with Au as the catalyst.<sup>60</sup> At high temperature, Si atoms decomposed from the Si precursor ( $\text{SiH}_4$ ,  $\text{SiCl}_4$ , etc.) form eutectic droplets together with Au catalyst. As the reaction proceeds, Si is continuously dissolved by the Au-Si eutectic droplets in increasingly high concentrations, leading to supersaturation and precipitation of crystals. The solid-liquid interface forms the growth interface, which acts as a sink for Si to continue to incorporate into the lattice, thus allowing the Si nanowires to grow in the presence of Au-Si alloy droplets. The diameter of the Si nanowire is determined by the size of the catalyst droplet. The nanowire growth process catalyzed by Au nanoclusters based on the VLS mechanism was observed by *in situ* TEM.<sup>61</sup> To prepare highly ordered arrays of 1D Si nanostructures, lithography patterning technique was utilized to prepared ordered catalyst arrays for the VLS process.<sup>62</sup> Fig. 4 b,c show a virtually defect-free Si wire array exhibiting extremely narrow diameter and length distributions, as well as

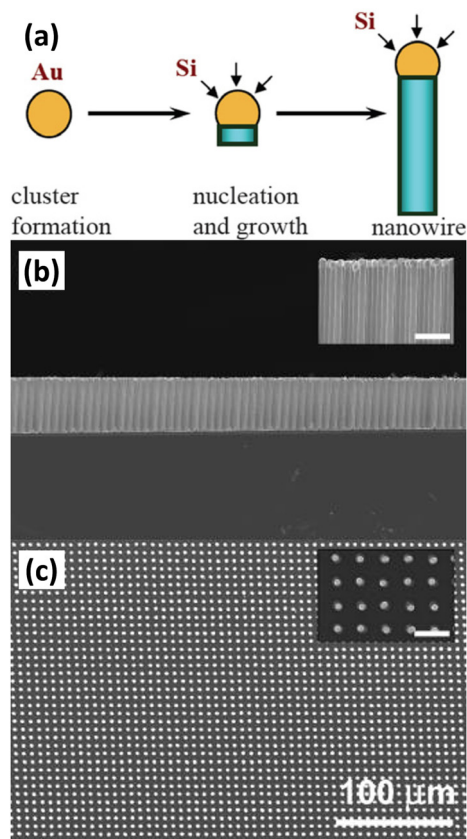


Fig. 4 (a) Schematic diagram of VLS growth of Si nanowires. Reproduced with permission from ref. 60, © 2006 IOP Publishing Ltd Side-view (b) and top-view (c) SEM images of a Au-catalyzed Si wire array. Reproduced with permission from ref. 62, © 2007 American Institute of Physics.

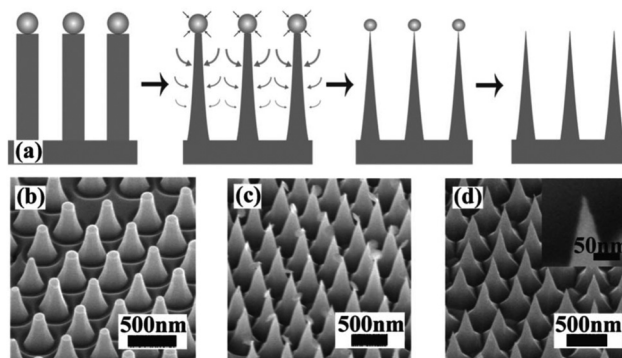


Fig. 5 (a) Schematic of sharpening process. (b)–(d) SEM images show the morphologies of each step. Reproduced with permission from ref. 67, © 2008 American Institute of Physics.

highly controlled wire placement. Highly ordered Si nanowire arrays have also been prepared by VLS growth of Si nanowires within nanopores of the anodic aluminum oxide (AAO) template.<sup>63</sup>

In addition to the VLS mechanism, several other mechanisms have been developed for the preparation of 1D Si nanostructures, such as oxide-assisted growth (OAG) mechanism,<sup>64</sup> solution-liquid-solid (SLS) mechanism,<sup>65</sup> etc.

### Other methods

MACE and CVD are the most commonly used methods for the preparation of 1D Si nanostructures that are used as photoelectrodes for solar hydrogen production. In addition to these, several other preparation methods have been explored. For example, using ordered arrays of gold particles produced by nanosphere lithography as catalysts, ordered arrays of vertically aligned Si nanowires were fabricated by molecular beam epitaxy (MBE) method.<sup>66</sup> Close-packed arrays of Si nanopillars over an entire 4 in. wafer were prepared by the synergy of Langmuir-Blodgett (LB) assembly and reactive ion etching (RIE).<sup>67</sup>  $\text{SiO}_2$  particles with diameters ranging from 50 to 800 nm were assembled on the Si wafer by LB method. Si nanopillars were formed by  $\text{Cl}_2$  based selective and anisotropic RIE of the Si substrate. Further, sharp Si nanocones could be prepared by the undercutting through control of the etching conditions, as shown in Fig. 5. Compared with MACE, RIE has better morphological controllability, but it requires more expensive equipment and a complex process, which substantially increases the cost of preparing 1D Si nanostructures.

## Advantages of 1D Si nanostructures

For solar hydrogen production applications, the 1D Si nanostructures provide several potential advantages over the traditional planar Si photoelectrode.<sup>36,68</sup>

### Large surface area

When a semiconductor photoelectrode is immersed in an electrolyte, the electrochemically active species in the solution



diffuse to the electrode surface and exchange electrons with the electrode to be oxidized or reduced. The larger the specific surface area of the electrode, the larger the electrochemical reaction site and the more opportunities for the electrochemically active species to contact the electrode surface and react. Therefore, increasing the specific surface area of the photoelectrode is beneficial for increasing the photocurrent density and reducing the overpotential required to achieve the desired current. 1D Si nanostructured photoelectrodes have a much larger specific surface area than planar Si photoelectrodes. For example, the surface area of a photoelectrode with an area of  $1 \mu\text{m} \times 1 \mu\text{m}$  is  $1 \times 10^{-12} \text{m}^2$ . When nanowires with a diameter of 100 nm and a length of  $10 \mu\text{m}$  are grown on the surface of the electrode, and the distance between each nanowire is 100 nm, the surface area of the photoelectrode is  $7.95 \times 10^{-11} \text{m}^2$ , which is about 80 times that of the planar electrode. In addition, the large specific surface area provides more sites that can be used to load catalysts, reducing the high activity requirement for individual sites by increasing the loading of catalysts per unit area, thus allowing the replacement of noble metal catalysts with inexpensive transition metal catalysts that have relatively poor catalytic activity.<sup>69</sup>

### Enhanced light absorption

Planar Si photoelectrodes have a strong reflective effect on light, resulting in a large portion of the light being lost. The nanostructured photoelectrode overcomes this drawback and dramatically improves the light absorption of the photoelectrode through the antireflection effect.<sup>34</sup> Especially, 1D nanostructured arrays of Si can generate significant “light trapping” effects, resulting in better antireflection effects and better light absorption.<sup>70</sup> Peng *et al.* investigated the light reflection properties of single-side polished Si wafers with Si nanowire arrays on the surface.<sup>71,72</sup> It was found that the reflectance of Si wafers was as high as 17–40% in the 300–1000 nm range, while the reflectance of Si nanowire arrays was only 2.1% or less. The morphology of the 1D Si nanostructures will affect their antireflectivity. The optical absorption properties of thin film and 1D nanostructures (nanowires and nanocones) of a-Si:H are shown in Fig. 6.<sup>73</sup> Photographs of the three samples are quite different. The Si thin film is mirror-like and highly reflective. The Si nanowires reflect less light while the Si nanocones look black, exhibiting enhanced absorption due to suppression of reflection from the front surface. The measured and calculated absorption spectra of the three samples over the wavelength range of 400–800 nm quantitatively demonstrate the visible absorbance of the three samples. 1D Si nanostructures have a significantly enhanced light absorption capacity. In particular, the Si nanocones maintain an absorption rate of more than 93% between 400 and 650 nm.

The Si nanowire arrays prepared by MACE have excellent anti-reflection effect.<sup>74–76</sup> However, the MACE process is carried out in an aqueous solution, and the surface tension of water causes the Si nanowires to agglomerate into bundles during the drying process, which may reduce the antireflection effect. It was found that post-treatment of Si nanowire arrays prepared

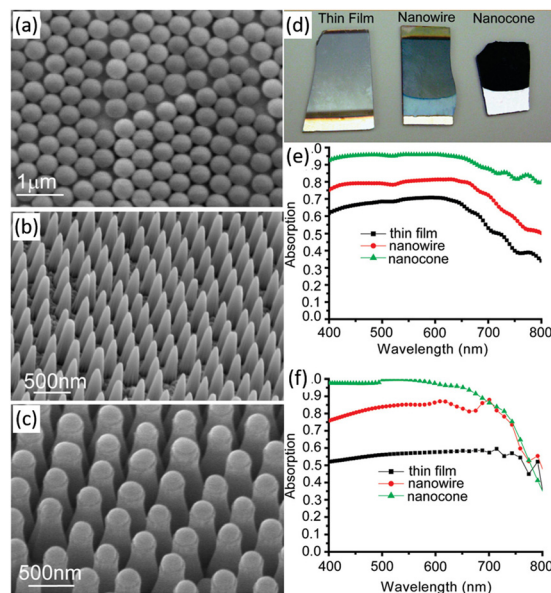


Fig. 6 SEM images (a–c), photograph (d) and the measured (e) and calculated (f) absorption spectra of Si thin film, nanowires and nanocones. Reproduced with permission from ref. 73, Copyright © 2009, American Chemical Society.

by MACE in KOH solution was able to separate Si nanowires from bundles by sharpening the top-ends of the nanowires.<sup>77</sup> As shown in Fig. 7a and b, the light reflectivity of Si nanowire arrays in the 300–1000 nm range was dramatically reduced after tapered Si nanowires formed by KOH treatment. The Si nanowires treated with 60 s in KOH showed the best antireflection performance with an average reflectance of less than 1%. This increase in reflectivity by changing the geometry of the

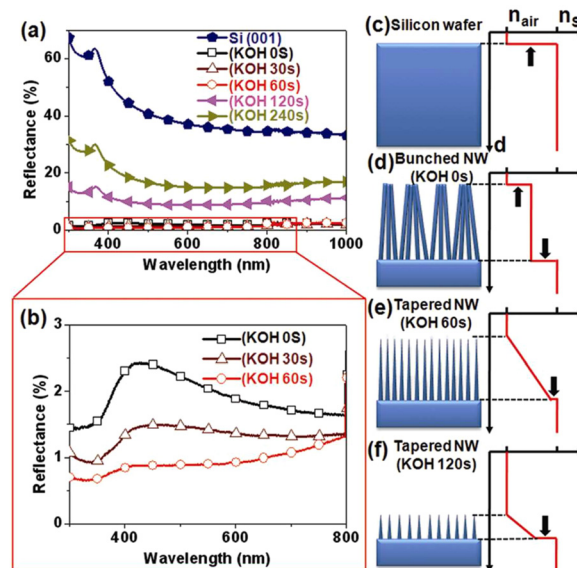


Fig. 7 (a and b) Optical reflective spectra recorded from different samples. (c–f) Schematic representation of the correlation between Si nanowire morphology and effective refractive index profile. Reproduced with permission from ref. 77, Copyright © 2011, The Optical Society.



Si nanowires is attributed to the multiple scattering and graded refractive index (GRI) effects, as shown in Fig. 7c–f.

### Decoupled light absorption and charge-carrier collection

For PEC hydrogen production, the semiconductor absorbs light to produce electron–hole pairs. The electrons or holes need to be transported to the surface of the semiconductor to react with the electrochemically active material in the solution. As mentioned earlier, the indirect bandgap semiconductor Si has a light penetration depth of up to 100  $\mu\text{m}$ . For conventional planar Si photoelectrodes, some photogenerated carriers need to travel distances of up to several tens of microns to reach the surface where the reaction takes place. In this process, many carriers will be recombined, resulting in low carrier utilization (Fig. 8a). The construction of photoelectrodes with 1D Si nanostructures is an excellent solution to this problem. As shown in Fig. 8b, after the 1D Si nanostructure is excited by light to generate photogenerated carriers, the minority carriers can be quickly transported to the surface for HER or OER reactions through the radial direction of the 1D nanostructure, which is in the nanometer scale in two dimensions. The short distance of this process dramatically degrades the chance of carrier recombination and improves carrier utilization. At the same time, the 1D Si nanostructure is continuous in one dimension, providing high-speed channels for the transport of the majority carriers to the substrate. This feature allows the use of low-purity Si with a short minority carrier diffusion length, significantly reducing the manufacturing cost of photoelectrodes.

Based on the advantages of 1D Si nanostructures, their superiority as photoelectrodes for PEC water splitting has been extensively reported. It is well established that 1D Si nanostructured photoelectrodes exhibit lower overpotential and higher photocurrent than planar Si photoelectrodes.<sup>78–80</sup> Oh *et al.* investigated the PEC hydrogen production performance of Si nanowire arrays fabricated by MACE.<sup>81</sup> As shown in

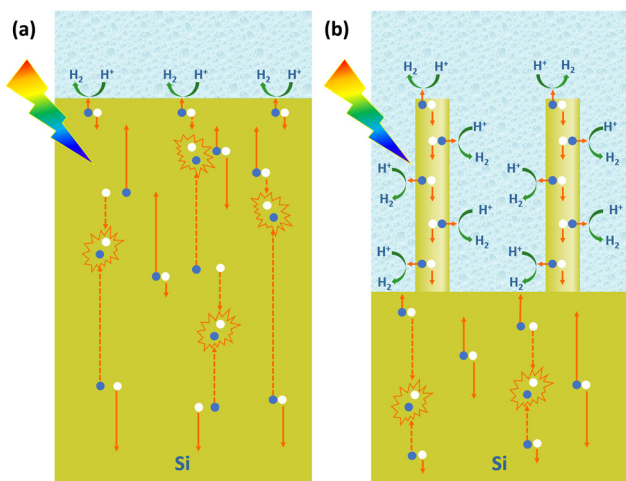


Fig. 8 Schematic diagram of the photogenerated carrier separation and transportation in the planar Si photoelectrode (a) and 1D Si nanostructured photoelectrode (b). The blue balls and white balls represent photogenerated electrons and photogenerated holes, respectively.

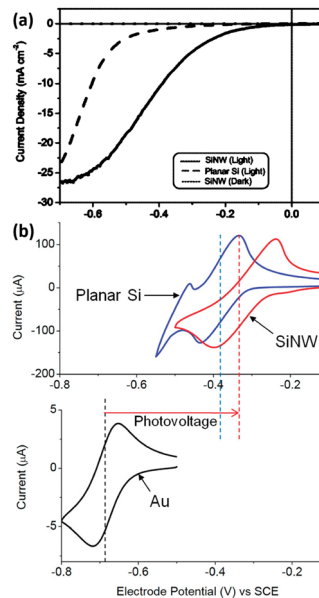


Fig. 9 (a) PEC HER performance comparison of the Si nanowire and planar Si. (b) Cyclic voltammety curves of methylviologen on the Si nanowire and planar Si photoelectrode under illumination and on the Au electrode in the dark. Reproduced with permission from ref. 81, Copyright © 2012, American Chemical Society.

Fig. 9a, the HER photocurrent for planar Si is very small until the potential reaches  $-0.5$  V. In contrast, the Si nanowire photocathode produces a much higher photocurrent, with an onset potential of  $0.2$  V, which is more positive than the planar Si. This positive onset potential comes from the higher photovoltage of the Si nanowire photoelectrodes, as shown in Fig. 9b for the cyclic voltammety test results of methylviologen. The increase in the photovoltage of the Si nanowire photocathode is attributed to the reduction of the overpotential, the favorable shift of the flat-band potential, and the light trapping or antireflection effect of the nanowires.

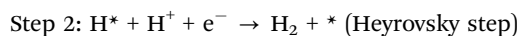
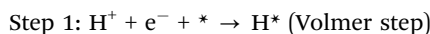
### Modification with co-catalysts

The kinetics of both HER and OER are sluggish on the bare Si surface, resulting in a large overpotential for water splitting. The most effective solution to this problem is to load a co-catalyst with high catalytic activity on the surface of Si. Then the next question is which material to choose as a co-catalyst. Purely from the point of view of reducing the overpotential, the catalytic activity of the co-catalyst is desired to be as high as possible. A direct criterion for the kinetic rate of a particular reaction at an electrode is the exchange current density ( $j_0$ ). When an electrode reaction is in equilibrium, the reaction rate in the anodic direction is equal to the reaction rate in the cathodic direction. The values of the anode-direction current ( $I_a$ ) and the cathode-direction current ( $I_c$ ) are equal, *i.e.*,  $I_a = I_c = I_0$ . The current density  $I_0/A$  ( $A$  is the electrode area)  $= j_0$ , *i.e.*, the exchange current density.<sup>82</sup> The value of  $j_0$ , which can be calculated from the Tafel curve, directly reflects the speed of the electrode

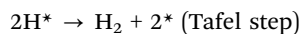




reaction. The larger  $j_0$  is, the faster the reaction rate is. The  $j_0$  of HER on Si is very low ( $\approx 10^{-8}$  A cm $^{-2}$ ), indicating the sluggish HER kinetics on Si. In addition to the evaluation factor  $j_0$ , which can only be obtained by measurement, catalytic activity descriptors based on DFT calculations have been proposed to predict the catalyst activity and guide catalyst design. For HER, the free energy change of H adsorption on the catalyst surface ( $\Delta G_H$ ) is the descriptor of the catalytic activity of a catalyst.<sup>83</sup> The HER occurring on the electrode surface undergoes a two-step process as follows:



or



where \* denotes a free site on the surface and H\* denotes the hydrogen atom adsorbed on the surface. In the first step, known as the Volmer step, the H<sup>+</sup> in solution receives an electron and adsorbs to the catalyst surface, forming adsorbed hydrogen H\*. In the second step, the H\* is converted to H<sub>2</sub> and detached from the catalyst surface by the Heyrovsky step or the Tafel step. From the whole HER reaction process, it can be seen that the intensity of H adsorption on the catalyst surface should be moderate. The rate of the whole reaction is limited by the Volmer step or the Heyrovsky/Tafel step, respectively, if the connection of H to the catalyst surface is too weak or too strong. Therefore, for a HER catalyst, it is implied that the closer the  $\Delta G_H$  is to zero, the higher the catalytic activity. Based on the above insight,  $j_0$  and  $\Delta G_H$  of various catalytic materials were measured and calculated and the volcano maps have been plotted.<sup>1</sup>

### Noble metals

In the volcano diagram, precious metals such as Pt are located at the top of the volcano in favor of HER. Among them, Pt with  $\Delta G_H$  close to 0 eV and the  $j_0$  up to  $10^{-3}$  A cm $^{-2}$  is recognized to be the catalyst with the highest HER catalytic activity available.<sup>83</sup> Therefore, Pt is an ideal co-catalyst for achieving high-performance HER on 1D Si nanostructures.<sup>81,84,85</sup> Wang *et al.* prepared Si nanowire arrays using MACE and loaded Pt nanoparticles on Si nanowires.<sup>86</sup> Electroless deposition (ELD) and atomic layer deposition (ALD) were utilized to synthesize Pt nanoparticles on Si nanowires. As shown in Fig. 10, the PEC HER of the Si nanowire photocathode was significantly improved after loading Pt nanoparticles. The Pt/Si nanowire photocathode prepared by the ALD method performed best. This may benefit from the uniform attachment of small-sized Pt nanoparticles on the Si surface and the effective charge transfer ensured by single-crystal Pt nanoparticles with single-crystal Si. It was also reported that the Pt electrocatalysts prepared by electron-beam evaporation showed much higher catalytic activity for HER than that prepared by electroless deposition.<sup>69</sup> The electroless Pt films produced a linear  $J$ - $E$  electrocatalytic response, indicating a resistive behavior. In

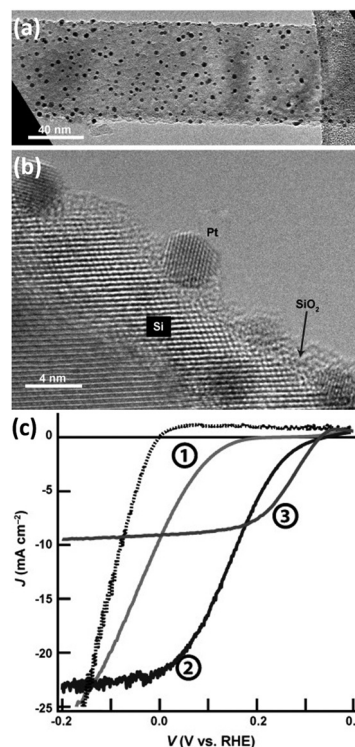


Fig. 10 (a) TEM and (b) HRTEM images of Pt nanoparticles modified Si nanowire, (c) PEC HER performance of naked Si nanowires (①), and Si nanowires modified with Pt nanoparticles prepared by ELD (②) and ALD (③). Reproduced with permission from ref. 86, Copyright © 2013 Wiley-VCH Verlag GmbH & Co. KGaA, Weinheim.

contrast, the e-beam deposited Pt films showed the expected exponential response. However, the e-beam-evaporated Pt films formed ohmic contact with p-Si, resulting in no photovoltage.

In addition to the preparation method, the catalytic activity of Pt nanoparticle-modified 1D Si nanostructures was greatly affected by the Pt loading. Yang *et al.* investigated the effect of ALD cycles on the PEC performance of Pt/Si nanowires.<sup>87</sup> As shown in Fig. 11a, the average diameter of Pt nanoparticles is approximately 0.8 nm after one ALD cycle. As the number of cycles increases, the particle size increases, and the size distribution is broader. Fig. 11b shows that the electrochemical performance also improves with the increase in ALD cycles, which may benefit from the increase in Pt loading. Photoelectrochemical tests showed that blank Si had almost no catalytic

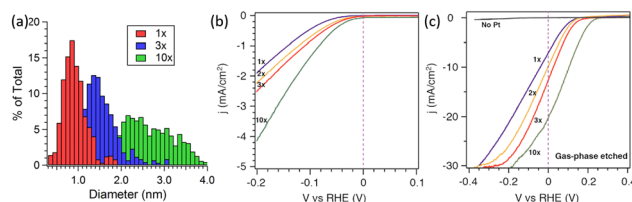


Fig. 11 (a) Histogram of Pt particle size for various numbers of ALD cycles, (b) electrochemical and (c) photoelectrochemical performance of Pt cocatalysts loaded via ALD in 0.5 M H<sub>2</sub>SO<sub>4</sub> electrolyte. Reproduced with permission from ref. 87, Copyright © 2013, American Chemical Society.





activity (Fig. 11c). After loading Pt nanoparticles, the PEC performance of the photoelectrode was greatly improved. As the number of cycles increased from 1 to 10, the photocurrent density at 0 V vs. RHE increased from 7 to 21 mA cm<sup>-2</sup>, and the slope of the *j*-*v* curve increased, indicating that the HER catalytic activity increased.

Other platinum group metals such as Ru and Ir have also been used as HER co-catalysts for Si nanowire photoelectrodes.<sup>88,89</sup> Shao *et al.* prepared Ru nanoparticles modified Si nanowire photoelectrode by reducing Ru ions by Si-H bonds.<sup>88</sup> The optimized Ru/Si nanowires exhibited low overpotential, good stability and low theoretical activation energy for HER. Ir-Au bimetal nanoparticles were also used as highly active catalyst to modify Si nanowires.<sup>89</sup> Besides, Ag was found to be a good cocatalyst for modifying 1D Si nanostructures to improve the HER kinetics.<sup>90,91</sup>

### Non-noble metals

In addition to catalytic activity, the cost of the catalyst should also be considered when selecting a co-catalyst. Due to the scarcity and high price of noble metals, there is a desire to replace them with non-noble metal catalysts. Metal Ni and its alloys are the most studied co-catalysts for modifying 1D Si nanostructured photoelectrodes thanks to their low price and relatively high catalytic activity. The performance of Ni and Ni-Mo as co-catalysts for the modification of 1D Si nanostructured photoelectrodes was evaluated and compared with Pt.<sup>69</sup> It was found that the catalytic activity for HER increased in the following order: Ni < Ni-Mo ≤ Pt, as shown in Fig. 12a. The electrocatalytic activity of Ni-Mo deposited for 180 s was even better than that of the electroless deposited Pt. However, the open circuit potential (*V*<sub>oc</sub>) and short circuit photocurrent densities for the photocathodes followed by the order Ni-Mo < Ni < Pt (Fig. 12b) due to the “pinch-off” effect.

To improve the performance of the Ni-Mo co-catalyst-modified Si photocathode, Lewis *et al.* prepared Si microwire arrays with buried radial n<sup>+</sup>p junction.<sup>92</sup> The Ni-Mo-coated n<sup>+</sup>p-Si microwire photoelectrode yielded a *V*<sub>oc</sub> as high as 0.46 V. Further, a light-scattering layer of TiO<sub>2</sub> nanoparticles were introduced to this system to eliminate the light blocking by the Ni-Mo catalyst layer (Fig. 13a and b).<sup>93</sup> The TiO<sub>2</sub> layer

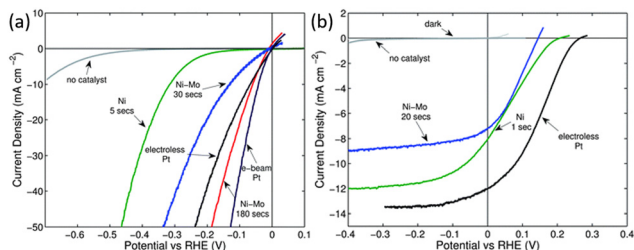


Fig. 12 (a) Polarization data of the HER activities of various catalysts on p<sup>+</sup>-Si microwire electrodes, (b) polarization data of the HER activities of various catalysts on illuminated p<sup>+</sup>-Si microwire array photoelectrodes. Reproduced with permission from ref. 69, Copyright © The Royal Society of Chemistry 2011.

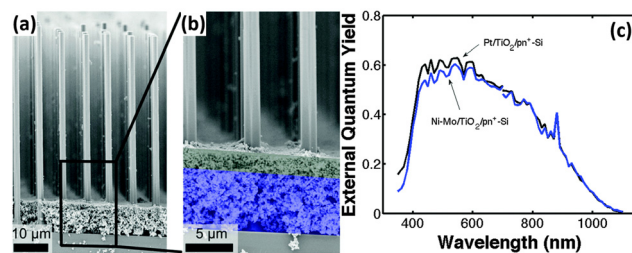


Fig. 13 (a and b) SEM images of the Ni-Mo/TiO<sub>2</sub>/n<sup>+</sup>p-Si microwire array, (c) spectral response data for the best-performing n<sup>+</sup>p-Si microwire-array photocathode loaded with Ni-Mo and TiO<sub>2</sub> scattering particles (blue curve) and for the best-performing n<sup>+</sup>p-Si microwire-array photocathode loaded with Pt and TiO<sub>2</sub> scattering particles (black curve). Reproduced with permission from ref. 93, Copyright © The Royal Society of Chemistry 2015.

scattered light back into the Si microwire array, while optically obscuring the underlying Ni-Mo catalyst film. In turn, the mass loading of Ni-Mo films is sufficient to produce high catalytic activity for HER on a geometric area basis. The best-performing Ni-Mo/TiO<sub>2</sub>/n<sup>+</sup>p-Si microwire array photoelectrode yielded a photovoltage of 420 mV. Fig. 13c shows that the external quantum efficiency (EQE) of Ni-Mo/TiO<sub>2</sub>/n<sup>+</sup>p-Si microwire array photoelectrode is close to that of Pt/TiO<sub>2</sub>/n<sup>+</sup>p-Si microwire array photoelectrode.

Tung *et al.* investigated the PEC HER performance of p-Si micropillar array photoelectrodes electrodeposited with Ni-M (M = Co, Fe, Mo) alloys as catalysts.<sup>94</sup> As shown in Fig. 14, the performance of p-Si micropillar array photoelectrodes modified with Ni-M co-catalyst showed the following trend Fe < Co < Mo. The *in situ* and *ex situ* tracking characterizations (XPS/XAS, GAXRD) revealed that the *in situ* formed surface Ni<sup>2+</sup> species during solar-driven hydrogen evolution is the primary factor for yielding the active site in Ni-M/p-Si micropillar arrays photocathodes. The presence of high-valence Mo species (Mo<sup>4+</sup> and Mo<sup>6+</sup>) could stabilize the high valence of Ni<sup>2+</sup> species near the catalyst surface, leading to the high catalytic activity of Ni-Mo.

In order to achieve the high catalytic activity, Ni-Mo catalysts with high loading on the Si surface are required. Since

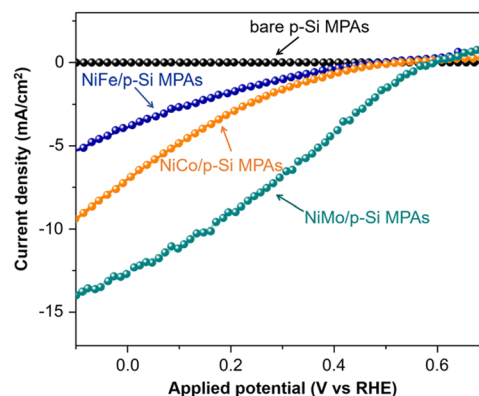


Fig. 14 PEC HER performance of p-Si micropillar arrays modified with Ni-Fe, Ni-Co and Ni-Mo as cocatalysts. Reproduced with permission from ref. 94, Copyright © 2020 Wiley-VCH Verlag GmbH & Co. KGaA, Weinheim.



Ni–Mo is impermeable to light, the resulting parasitic light absorption will affect the photoelectrode performance. To solve this problem, Vijselaar *et al.* designed and fabricated a Si microwire array photoelectrode by varying the fraction of Ni–Mo catalyst coverage over the microwires.<sup>95</sup> Three types of photoelectrodes were designed: completely exposed microwires (a), microwires passivated by SiO<sub>2</sub> with exposed tops (b), and microwires with a defined exposed area in the range of 2–36 μm from the top (c). Ni–Mo nanoparticles were electrodeposited on the exposed area of the microwires. The photoelectrical measurements indicated that the loading of Ni–Mo significantly decreases photocurrent densities due to blocking and/or scattering the light. In contrast, the photoelectrochemical HER test results show that photocurrent and photovoltage substantially increased after loading the Ni–Mo catalyst. For the c-type samples, the PEC HER performance was optimized by adjusting the distance between Si microwires and the spatioselective coverage of Ni–Mo on the microwires, resulting in a photovoltage of up to 495 mV.

Si is easily oxidized to form Si oxide on the surface, which will hinder charge transfer. Therefore, the direct use of 1D Si nanostructures as photoanodes for water oxidation is rarely investigated because of the great difficulties. Tung *et al.* demonstrated the decoration of isolated Ni nanoparticles as both protecting layer and OER electrocatalysts on the surfaces of n-Si microwires.<sup>96</sup> The stability of the photoelectrode was determined by the loading of the Ni catalyst. If the loading is too small, the stability is poor; and if the loading is too large, the light absorption is affected. The optimized Ni/n-Si microwire arrays photoelectrode showed an OER onset potential of 1.08 V *vs.* RHE and a saturation photocurrent of 35.5 mA cm<sup>-2</sup>.

### MoS<sub>x</sub>

The DFT-calculated  $\Delta G_{\text{H}}$  of the MoS<sub>2</sub> is slightly positive at +0.08 eV, suggesting that it might be a suitable catalyst for HER.<sup>97,98</sup> Tran *et al.* deposited MoS<sub>2</sub> co-catalyst on Si nanowires by photo-assisted electrodeposition, as shown in Fig. 15a and b.<sup>99</sup> The layer structured MoS<sub>2</sub> was assembled tightly on the surface of Si nanowires. The HER onset potential of MoS<sub>2</sub>/Si nanowire photoelectrode under illumination is 0.25 V *vs.* RHE, which is about 200 mV more positively shifted than that of Si nanowires without loading of MoS<sub>2</sub> co-catalyst (Fig. 15c). As shown in Fig. 15d, the MoS<sub>2</sub>/Si nanowire photoelectrode exhibited comparable PEC performance to the Si nanowire photoelectrode modified with Pt nanoparticles, indicating the great potential of MoS<sub>2</sub> as a candidate for Pt. Similar results were obtained on the thermolysis-prepared MoS<sub>2</sub>/Si photoelectrode.<sup>100</sup> MoS<sub>2</sub> acted not only as a catalyst but also as a light absorber to improve the PEC activity of the Si nanowire photoelectrode.<sup>101</sup> The MoS<sub>2</sub>/Si Schottky junction enhances the photo-generation of electron–hole pairs and suppresses carrier recombination. The low charge-transfer resistances at the semiconductor/catalyst and catalyst/electrolyte interfaces, slow carrier recombination dynamics, and consequently efficient charge carrier separation resulted in superior performance.<sup>102</sup> Zinc 1T-phase MoS<sub>2</sub> synthesized by the hydrothermal method was drop cast



Fig. 15 TEM (a) and HRTEM (b) images of the MoS<sub>2</sub>/Si nanowires. Comparison of the photoelectrochemical HER performance: (c) Si nanowires and MoS<sub>2</sub>/Si nanowires, (d) Si nanowires modified with MoS<sub>2</sub> or Pt nanoparticles in the dark or under illumination. Reproduced with permission from ref. 99, Copyright © 2012 Wiley-VCH Verlag GmbH & Co. KGaA, Weinheim.

onto Si nanowires.<sup>103</sup> While improving the catalytic activity of the photoelectrode, MoS<sub>2</sub> also prevents the growth rate of the natural oxide layer SiO<sub>x</sub> on the Si surface, significantly improving the stability of the photoelectrode. In addition, a Ag interface layer was introduced between Si nanowire and MoS<sub>2</sub> to improve carrier utilization.<sup>104</sup>

The multiple valence states of Mo lead to the presence of a wide range of molybdenum sulfides. In addition to MoS<sub>2</sub>, other molybdenum sulfides have been used as co-catalysts for 1D Si nanostructure photoelectrodes and have shown excellent performance.<sup>105–108</sup> Hou *et al.* modified 1D Si pillars with bioinspired incomplete cubane-like Mo<sub>3</sub>S<sub>4</sub> clusters to evolve hydrogen at a rate comparable to that of Pt.<sup>109</sup> After modified with Mo<sub>3</sub>S<sub>4</sub> clusters, the onset potential of the photoelectrode shifted from -0.4 V *vs.* RHE to 0.15 V *vs.* RHE. That is, the Mo<sub>3</sub>S<sub>4</sub> co-catalyst reduces the HER overpotential of the p-Si photoelectrode by 550 mV. Further, 1D Si pillars with a diameter of 3 μm and a length of 50 μm were prepared by photolithography to improve the performance of the p-Si photoelectrodes. Compared with planar photoelectrodes, p-Si photoelectrodes with pillar structure showed great superiority in onset potential, saturation current density, and incident photon-to-current conversion efficiency (IPCE). The limiting current density and IPCE of the photoelectrode are reduced after loading the co-catalyst, which can be due to the co-catalyst hindering the light absorption of the photoelectrode. Therefore, the balance between catalytic activity and light absorption performance should be considered when designing and



## EES Catalysis

constructing the photoelectrodes. Too little catalyst will not achieve a sufficient catalytic effect, while too much catalyst will affect light absorption.

## Metal compounds

Inspired by the excellent HER catalytic activity of MoS<sub>2</sub>, other metal sulfides have also been investigated as co-catalysts for 1D Si nanostructured photoelectrodes. P-Si nanowires were modified with tungsten sulfides (crystalline WS<sub>2</sub> and amorphous WS<sub>3</sub>) prepared by pyrolysis.<sup>110</sup> Compared with the crystalline WS<sub>2</sub>, the amorphous WS<sub>3</sub> had larger surface area and more active sites, resulting in higher catalytic activity in HER. CoS<sub>2</sub> was used as a co-catalyst to enhance the HER catalytic activity of Si microwire photoelectrodes while protecting Si from oxidation and reductive degradation.<sup>111</sup> It was suggested that more vacancies in the S-3p unoccupied states of the CoS<sub>2</sub>-Si electrode were presented with a lower negative charge of S<sub>2</sub><sup>2-</sup> to form weaker S-H bond strength, promoting water splitting efficiency. 2D layered VS<sub>2</sub> synthesized by CVD was wrapped on the surface of Si nanowires for highly efficient PEC water reduction.<sup>112</sup>

CoP has been reported to be one of the most active, acid-stable, earth-abundant HER electrocatalysts.<sup>113,114</sup> Therefore, it was utilized as a co-catalyst for 1D Si nanostructure photoelectrodes.<sup>115-117</sup> Hollow Co<sub>2</sub>P spheres were prepared on the Si nanowire surface by photo-assisted electrodeposition and phosphorization. This resulted in a remarkably positively shifted onset potential for PEC HER.<sup>115</sup> Conformal and continuous Co<sub>2</sub>P was coated on the highly ordered Si nanowire arrays by drop-casting method followed by low-temperature phosphorization.<sup>116</sup> The Co<sub>2</sub>P layer improved the HER catalytic activity of the photoelectrode while protecting the Si from photo-oxidation and improving the lifetime of the photoelectrode. FeP was also an efficient HER co-catalyst for the 1D Si nanostructured photoelectrodes.<sup>118</sup> Tetragonal phase Ni<sub>12</sub>P<sub>5</sub> nanoparticles was utilized as an efficient co-catalyst to modify p-Si nanowire arrays, an optimal overpotential of 143 mV was required for 20 mA cm<sup>-2</sup> current density.<sup>119</sup>

Uniform amorphous NiCoSe<sub>x</sub> shells were photo-assisted electrodeposited on the bamboo shoot-like Si nanopillars to fabricate the p-Si/NiCoSe<sub>x</sub> core/shell nanopillar array photoelectrode (Fig. 16a).<sup>120</sup> The good optical transparency and the excellent electrocatalytic activity boosted the generation and utilization of the photogenerated electrons. As shown in



Fig. 16 (a) SEM image of the p-Si/NiCoSe<sub>x</sub> core/shell nanopillar array, (b) PEC HER performance. Reproduced with permission from ref. 120, Copyright © The Royal Society of Chemistry 2016.

Fig. 16b, the onset potential ( $j = -0.1 \text{ mA cm}^{-2}$ ) of the p-Si/NiCoSe<sub>x</sub> core/shell nanopillar array photoelectrode is about 0.25 V vs. RHE, which is much more positive than that of the p-Si nanopillar photoelectrode. It is worth noting that the performance of the planar photoelectrodes in terms of onset potential and photocurrent is much worse compared to the 1D nanostructured photoelectrodes. The phase structure of a catalyst predominantly affects its catalytic activity. Three possible polymorphs of nickel selenide (orthorhombic phase NiSe<sub>2</sub>, cubic phase NiSe<sub>2</sub>, and hexagonal phase NiSe) were investigated as bifunctional electrocatalysts for both PEC water reduction and oxidation on 1D Si nanowire arrays.<sup>121</sup> It was found that the orthorhombic NiSe<sub>2</sub> nanocrystals exhibited the best PEC water splitting performance, owing to the more metallic nature and the formation of fewer surface oxides.

In addition to the metal sulfides, metal phosphides, and metal selenides mentioned above, metal oxides,<sup>122-125</sup> metal borides,<sup>126</sup> and metal carbides,<sup>127</sup> and other metal compounds<sup>128,129</sup> have also been used as HER or OER co-catalysts for 1D Si nanostructured photoelectrodes.

## Carbon materials

Although metals and metal compounds have achieved remarkable success in enhancing the performance of 1D Si nanostructured photoelectrodes as co-catalysts, the high cost (mainly referring to precious metals), light scattering, and stability issues have prompted the development of non-metallic catalysts. Sim *et al.* first demonstrated that N-doped graphene as a co-catalyst can significantly enhance the PEC HER performance of p-Si photocathodes.<sup>130</sup> Then they applied the idea to 1D Si nanowire array photoelectrodes.<sup>131</sup> As shown in Fig. 17a, the N-doped graphene quantum sheets (N-GQS) were obtained by replacing some C atoms with N atoms through N<sub>2</sub>-plasma treatment. The

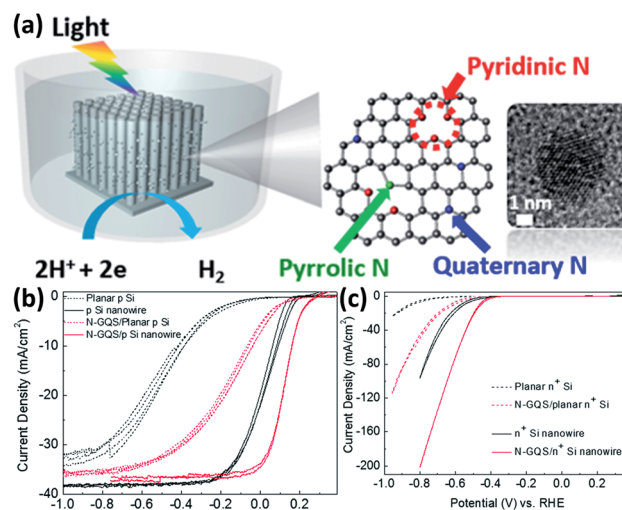


Fig. 17 (a) Schematic of N-doped graphene quantum sheets (N-GQSs) decorated on a Si nanowire photocathode. (b) Photoelectrochemical and (c) electrochemical HER performance of the planar and 1D nanostructured Si photoelectrodes with/without N-GQS catalysts. Reproduced with permission from ref. 131, Copyright © The Royal Society of Chemistry 2015.





N-GQS were modified on the surfaces of 1D Si nanowires by drop-casting. Photoelectrochemical and electrochemical tests showed that N-GQS has excellent HER catalytic activity (Fig. 17b). For the planar Si electrode, the PEC HER onset potential was positively shifted by 0.3 mV after modification with N-GQS. However, the difference is minimal for Si nanowire photoelectrodes, only 0.09 V. Due to the excellent light transmission of N-GQS, the photocurrent of the photoelectrode did not decrease significantly after loading the co-catalyst. Yu *et al.* fabricated reduced graphene oxide (rGO)-decorated Si nanowire array photocathode by a flexible and scalable electrochemical reduction method.<sup>132</sup> In addition to 2D nanostructured graphene, another characteristic carbon material, 1D nanostructured carbon nanowires, is also used as a co-catalyst for 1D Si nanostructured photoelectrodes. Si nanowire arrays were modified with hierarchically branched carbon nanowires that exhibited outstanding catalytic activity for HER and excellent stability against photo-oxidation and corrosion.<sup>133</sup>

For co-catalysts used to modify 1D Si nanostructures, the highest catalytic activity is still that of the noble metal Pt. However, considering issues such as cost and abundance, some non-precious metal catalysts are good choices, such as Ni-based alloys, sulfides of Mo, *etc.* Although the activity of these co-catalysts still is still somewhat inferior to that of Pt, it is expected to be further improved by several means, such as the high entropy alloy catalysts which have emerged in recent years. The strategy of single-atom catalysts, combining small amounts of noble metals with non-precious metals, can also be considered to obtain high performance at relatively low cost.

## 1D Si nanostructured heterojunction

The loading of co-catalysts with high electrocatalytic activity on the surface of 1D Si nanostructures solves the problem of sluggish reaction kinetics of water splitting on Si surfaces and dramatically reduces the water splitting overpotential. However, the narrow band gap of Si results in a low photovoltage, and its valence band potential is insufficient to oxidize water. These make it necessary to apply a large applied voltage to enable the PEC water splitting reaction. Combining Si with other semiconductors to form heterostructures effectively solve the high overpotential due to energy band structure.

### Type I heterojunction

Different heterojunctions can be formed between two semiconductors, depending on their conductivity types and energy band structures. A Type I heterojunction is formed between two semiconductors with a straddling bandgap. In this structure, both photogenerated electrons and photogenerated holes are transported from the semiconductor with wider band gap to the semiconductor with narrower band gap. The accumulation of photogenerated carriers within one semiconductor will greatly increase the chance of recombination of photogenerated electrons and holes. From this point of view, the Type I heterojunction is unfavorable for PEC hydrogen production

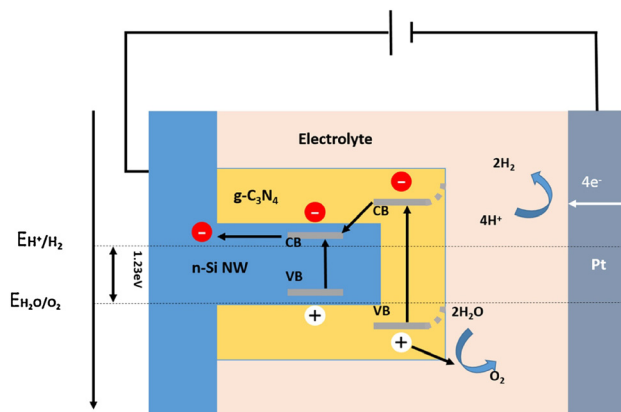


Fig. 18 Schematic diagram of energy band alignment of the Si/g-C<sub>3</sub>N<sub>4</sub> heterojunction. Reproduced with permission from ref. 135, Copyright © 2016 Published by Elsevier B.V.

applications. Nevertheless, there have still been some reports of the 1D Si nanostructures-based Type I heterojunctions for solar water splitting applications, including Si nanowires/CdSe quantum dots,<sup>134</sup> Si nanowires/g-C<sub>3</sub>N<sub>4</sub> shell arrays,<sup>135</sup> Si nanowires/g-C<sub>3</sub>N<sub>4</sub> nanosheets,<sup>136</sup> and Si nanowires/SrTiO<sub>3</sub> nanoparticle-coupled g-C<sub>3</sub>N<sub>4</sub> nanosheets.<sup>137</sup> Fig. 18 shows the schematic diagram of the energy band alignment and charge transporting process in the Si/g-C<sub>3</sub>N<sub>4</sub> heterojunction.<sup>135</sup> Under illumination, electrons in the valence bands of Si and g-C<sub>3</sub>N<sub>4</sub> jump to the conduction bands and generate holes in the valence bands. The holes are transferred to the surface of the photoelectrode to oxidize water and produce O<sub>2</sub>. The photogenerated electrons are transferred from the conduction band of g-C<sub>3</sub>N<sub>4</sub> to the conduction band of Si under the effect of the built-in electric field. They are transferred through an external circuit to the counter electrode for HER.

### Type II heterojunction

Another structure that is more favourable for the separation of the photogenerated carriers than the Type I heterojunction is the Type II heterojunction. Fig. 19a shows the energy band structure of a Type II heterojunction consisting of a p-type semiconductor (S1) and an n-type semiconductor (S2) with staggered band alignment.<sup>138</sup> Since the Fermi level of S2 is higher than S1; electrons are injected from S2 to S1 at the

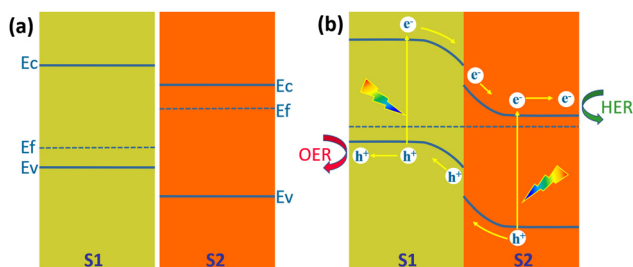


Fig. 19 Schematic diagram of the energy band diagram of a Type II heterojunction, (a) in the dark before contact, (b) under illumination after contact.



interface after contact between them. The built-in electric field formed on their respective surfaces leads to the bending of the energy bands, with S1 bending downward and S2 bending upward (Fig. 19b). Under illumination, electrons in the valence band of both semiconductors jump to the conduction band and create holes in the valence band. Driven by the built-in electric field, electrons in the S1 conduction band flow to the S2 conduction band and transport to the S2 surface for HER, and holes in the S2 valence band flow to the S1 valence band and transport to the S1 surface for OER. This energy band structure will significantly reduce the recombination rate of the photogenerated carriers. However, it can also be seen from Fig. 19b that the reduction ability of electrons and the oxidation ability of holes are weakened. Type II heterojunctions can also be composed of two semiconductors of the same conducting type, *i.e.*, n/n and p/p. In both cases, the minor difference in Fermi energy levels generates a smaller built-in electric field, which is not conducive to photogenerated carrier separation and transfer.<sup>138</sup>

1D Si nanostructures-based Type II heterostructure photoanodes were fabricated by using n-Si nanowires core and  $\alpha$ -Bi<sub>2</sub>O<sub>3</sub> shell and Sn<sub>3</sub>O<sub>4</sub> nanosheets.<sup>139,140</sup> Excellent PEC water splitting performance was achieved due to the enhanced light harvesting, efficient photogenerated electron-hole separation and rapid charge transfer at the photoanode. Although a Ag interlayer is introduced between the p-Si nanowire core and the semiconducting MoS<sub>2</sub> shell, the charge transfer between Si and MoS<sub>2</sub> follows the Type II mechanism.<sup>104</sup> Ta<sub>2</sub>O<sub>5</sub> and N-doped graphene quantum dots (NGQD) were sequentially decorated on the surface of the p-Si nanowires.<sup>141</sup> As shown in Fig. 20a, the p-SiNWs-Ta<sub>2</sub>O<sub>5</sub>-NGQD photoelectrodes exhibited more positive onset potential and much higher photocurrents than p-Si nanowires and p-Si nanowires-NGQD. The photogenerated carriers transfer between Si and Ta<sub>2</sub>O<sub>5</sub> and between Ta<sub>2</sub>O<sub>5</sub> and NGQD follows the Type II mechanism, that is, a double Type II heterojunction is formed (Fig. 20b).

### Z-Scheme heterojunction

Type II heterojunctions do promote the separation of photogenerated carriers, but at the expense of the reducing capacity of photogenerated electrons and the oxidizing capacity of photogenerated holes. In this regard, Z-shaped heterojunctions are very beneficial for solar water splitting for hydrogen production.

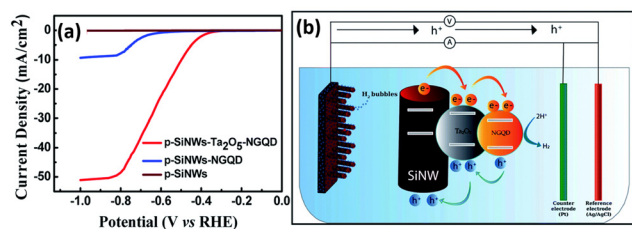


Fig. 20 (a) PEC HER performance of the p-SiNWs-Ta<sub>2</sub>O<sub>5</sub>-NGQD photoelectrodes. (b) Schematic diagram of the charge transfer process in the p-SiNWs-Ta<sub>2</sub>O<sub>5</sub>-NGQD photoelectrode. Reproduced with permission from ref. 141, Copyright © The Royal Society of Chemistry 2022.

The Z-scheme junction is a heterojunction structure favourable for photoelectrochemical water splitting. The Z-scheme junction works similarly to the natural photosynthesis system. In terms of energy band structure, the Z-scheme heterojunction is similar to the type II heterojunction in that both are composed of two semiconductors with a staggered band alignment. However, the migration process of photogenerated carriers within Z-scheme heterojunctions is completely different from that of Type II heterojunctions due to the internal electric field, interfacial defects, and crystal facets.<sup>142</sup> Fig. 21 illustrates the schematic band diagram of a typical Z-scheme junction.<sup>138</sup> To form a Z-scheme junction, the two semiconductors (S1 and S2) should have staggered alignment of band structures (Fig. 21a). When two semiconductors are in contact, electrons will flow from S1 to S2 because of the different Fermi levels until the Fermi levels of the two semiconductors are equal and equilibrium is reached. The energy bands of the two semiconductors are bent upward (S1) or downward (S2) at the interface. Under illumination, electrons jump from the valence band to the conduction band and create holes in the valence band. Driven by the built-in electric field generated by the energy band bending, the holes in the valence band of S1 and the electrons in the conduction band of S2 are transported to the interface and recombined (Fig. 21b). The electrons in the conduction band of S1 and holes in the valence band of S2 are transported to the surface of the semiconductor and undergo reduction and oxidation reactions, respectively. The recombination of electrons and holes at the interface of S1 and S2 significantly facilitates the separation and transfer of photogenerated carriers, thus improving their utilization. At the same time, the strong redox capacity of photogenerated electrons and photogenerated holes is maintained.

Yang *et al.* first realized the Z-scheme junction based on 1D Si nanostructures.<sup>143</sup> The Si nanowires were prepared by the metal-assisted electroless etching method and then coated with conformal TiO<sub>2</sub> film by atomic layer deposition (ALD). Fig. 22a shows that the heterojunction formed by n-Si nanowires and TiO<sub>2</sub> has the best performance, exhibiting the lowest onset potential for OER and the largest photocurrents at the same potentials. The performance of nanowire photoelectrodes is better than that of planar photoelectrodes for both n-Si and p-Si, which benefits from the low reflectivity and large surface

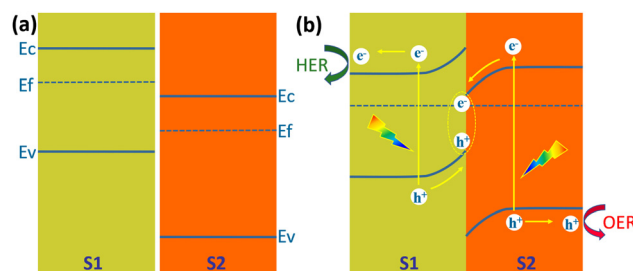


Fig. 21 Schematic diagram of the energy band diagram of a Z-scheme heterojunction, (a) in the dark before contact, (b) under illumination after contact.



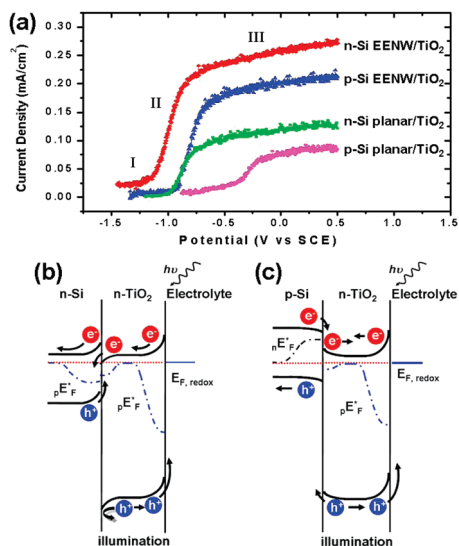


Fig. 22 (a) PEC OER performance on different Si/TiO<sub>2</sub> photoelectrodes. Band structures diagram of the n-Si/n-TiO<sub>2</sub> (b) and p-Si/n-TiO<sub>2</sub> heterostructures. Reproduced with permission from ref. 143, Copyright © 2009, American Chemical Society.

area. Compared with p-Si/TiO<sub>2</sub>, the photocurrent density of n-Si/TiO<sub>2</sub> photoelectrodes is 20–25% higher, which benefits from the Z-scheme junction of n-Si/TiO<sub>2</sub> that promotes photo-generated carrier separation. As shown in Fig. 22b, after the contact between n-Si and n-TiO<sub>2</sub>, the n-Si energy band bends upward, and the n-TiO<sub>2</sub> energy band bends downward at the interface. Driven by the built-in electric field, the photogenerated holes in the n-Si valence band and the photogenerated electrons in the n-TiO<sub>2</sub> conduction band move toward the interface and recombine at the interface. The holes in the valence band of n-TiO<sub>2</sub> move to the surface to oxidize OH<sup>-</sup> to O<sub>2</sub> (OER). The electrons in the n-Si conduction band flow through the external circuit to the counter electrode for HER. As for the p-Si/n-TiO<sub>2</sub> structure, due to the lower Fermi energy level position of p-Si, the energy band bending direction of p-Si/n-TiO<sub>2</sub> is exactly opposite to that of n-Si/n-TiO<sub>2</sub> (Fig. 22c); thus the process as mentioned above of photogenerated carrier separation and transfer cannot occur, and the photocurrent is lowered.

Considering the very similar energy band structures of ZnO and TiO<sub>2</sub>, the formation of Z-scheme heterojunctions with ZnO and n-Si has also been investigated.<sup>144–146</sup> It is worth mentioning that Peng *et al.* introduced Au nanoparticles acting as catalysts and plasma photosensitizers in the ZnO layer to achieve Au-sensitized n-Si/ZnO core-shell nanowire array photoanode for the unassisted water splitting.<sup>145</sup> The doping of N in ZnO and the deposition of NiOOH co-catalyst on the ZnO surface also facilitate the performance of n-Si/ZnO 1D nanostructured array photoanodes for water splitting.<sup>146</sup>

Wang and Peng *et al.* fabricated n-Si/Fe<sub>2</sub>O<sub>3</sub> nanowire array photoelectrode, also based on the Z-scheme band alignment and charge transportation, to achieve an enhanced performance of the PEC OER.<sup>147,148</sup> Qi *et al.* prepared n-Si/ $\alpha$ -Fe<sub>2</sub>O<sub>3</sub>

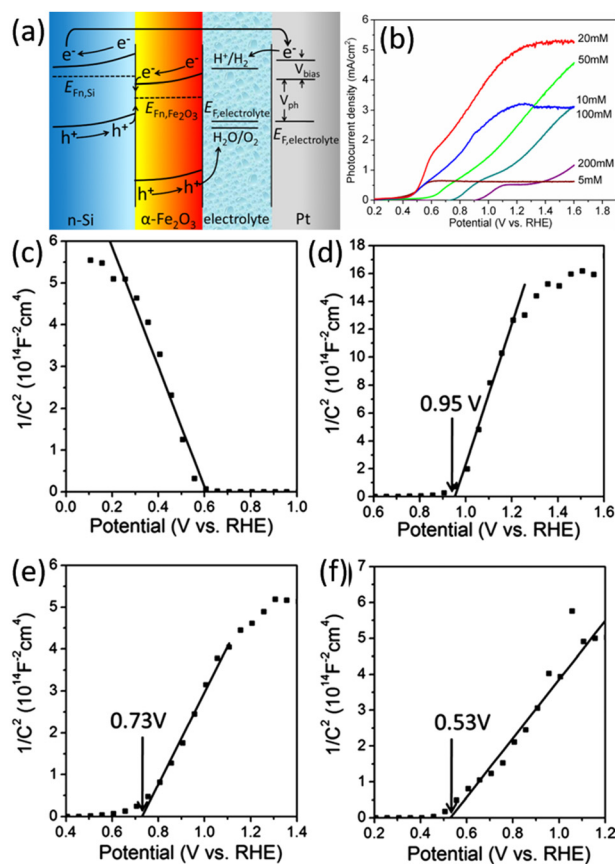


Fig. 23 (a) Schematic band alignment of the n-Si/ $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> heterostructure. (b) LSV curves of the n-Si/ $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> core/shell nanowire arrays prepared at different iron nitrate concentrations. Mott-Schottky curves recorded from photoelectrodes with different  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> thickness: (c) p-Si/ $\alpha$ -Fe<sub>2</sub>O<sub>3</sub>-20 mM, (d) n-Si/ $\alpha$ -Fe<sub>2</sub>O<sub>3</sub>-200 mM, (e) n-Si/ $\alpha$ -Fe<sub>2</sub>O<sub>3</sub>-20 mM, (f) n-Si/ $\alpha$ -Fe<sub>2</sub>O<sub>3</sub>-5 mM. Reproduced with permission from ref. 149, Copyright © The Royal Society of Chemistry 2014.

core/shell nanowire arrays used as photoanode for PEC water splitting.<sup>149</sup> The onset potential for OER was as low as 0.5 V vs. RHE, and the photocurrent was up to 5.28 mA cm<sup>-2</sup> at a potential of 1.23 V vs. RHE. Fig. 23a clearly illustrates the energy band structure of n-Si/ $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> and the separation and transfer process of photogenerated carriers under illumination. The thickness of the  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> layer on the Si surface is regulated by adjusting the concentration of iron nitrate in the preparation process. As shown in Fig. 23b, the different thickness of  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> significantly influences on the onset potential and current density of the photoelectrode. The onset potential becomes more positive with the increase of  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> thickness. Mott-Schottky test results show that the flat-band potential of the n-Si/ $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> photoelectrodes increases with the increase of  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> thickness (Fig. 23c–f). This indicates that the band edge position of n-Si at the n-Si/ $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> interface is higher in the samples with a smaller thickness of  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub>, so the quasi-Fermi energy level of the photogenerated electrons in n-Si is also higher after illumination. A smaller applied voltage is needed to assist the photogenerated electrons in n-Si to meet the energy required to reduce water to hydrogen. Hence, the onset





potential of the photocurrent is more positive in the samples with a smaller thickness of  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub>. Researchers have proposed several methods to improve further the performance of the n-Si/ $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> core/shell nanowire array photoelectrodes. NiO and Co<sub>3</sub>O<sub>4</sub> co-catalysts were loaded on the  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> surface to improve the OER reaction kinetics.<sup>150</sup> By doping  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> with Sn and forming oxygen vacancies, the charge transfer resistance is reduced, and the donor concentration is increased.<sup>151</sup>

The heterojunction formed by Si and other semiconductors effectively increases the photogenerated voltage of 1D Si photoelectrodes. However, the large number of defects at the Si/semiconductor interface will act as recombination centers and lead to recombination of the photogenerated carriers. A passivation layer has been introduced between Si and semiconductor to solve this problem. For example, an Al<sub>2</sub>O<sub>3</sub> interlayer was inserted between n-Si and  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub>.<sup>152,153</sup> Al<sub>2</sub>O<sub>3</sub> can passivate the defects on the surface of Si microwires and the bottom surface of  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub>, promote the separation of photogenerated carriers and reduce the photoanode surface recombination rate. In addition, Al was doped into  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> to improve the charge transfer efficiency on the  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> surface. TiO<sub>2</sub> interlayer was also introduced to the n-Si/ $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> heterojunction.<sup>154,155</sup> Solution-processed mesoporous TiO<sub>2</sub> layer covered on the surface of the Si nanowires, and then  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> nanorods were grown on TiO<sub>2</sub> through a hydrothermal method.<sup>155</sup> Fig. 24a and b show the SEM image and schematic diagram of the n-Si nanowires/TiO<sub>2</sub>/ $\alpha$ -Fe<sub>2</sub>O<sub>3</sub>. It was found that the introduction of mesoporous TiO<sub>2</sub> significantly enhanced the OER photocurrents of the photoelectrodes, which was attributed to the reduced charge recombination rate at the surface of  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub>, as illustrated in the band alignment of Fig. 24c and d.

In addition, a 1D Si nanostructured heterojunction with a combined Z-scheme and Type II mechanism was also constructed. Liu *et al.* reported a multijunction of ZnIn<sub>2</sub>S<sub>4</sub> nanosheet/TiO<sub>2</sub> film/Si nanowire for PEC water splitting.<sup>156</sup> Carrier transport between TiO<sub>2</sub> and ZnIn<sub>2</sub>S<sub>4</sub> follows the Type

II mechanism, while carrier transport between TiO<sub>2</sub> and Si follows the Z-scheme mechanism. The flat-band potential of the multi-junction becomes much more negative than that of the Si nanowire, indicating a higher carrier concentration and a more efficient charger transport process in the multi-junction. Furthermore, there are reports on the formation of Z-scheme heterojunctions of n-Si with other semiconductor materials, such as Si nanowires/WO<sub>3</sub>,<sup>157</sup> Si nanowires/Co<sub>3</sub>O<sub>4</sub>,<sup>158</sup> *etc.*

### Si/polymer heterojunction

In addition to inorganic semiconductors, 1D Si nanostructures can also form heterojunctions with semiconducting polymers. The Si/polymer heterojunctions offer advantages such as large-area coverage, low-temperature process capability, easy preparation, low cost, and so on.<sup>159</sup> Zhang *et al.* reported the modification of Si nanowires with semiconducting polymers, including doped poly(3,4-ethylenedioxythiophene) (PEDOT)<sup>160,161</sup> and poly((2,9,16,23-tetraaminophthalocyanine)Cu) (CuTAPc).<sup>162</sup> As shown in Fig. 25a, the Si/PEDOT junction helps separate the photogenerated carriers.<sup>160</sup> The electrons transferred to the counter electrode for HER, and the holes transported to the surface of PEDOT to oxidize water with the assistance of the CuPc co-catalyst. As a result, the PEDOT/Si nanowire photoelectrodes exhibit significantly improved PEC performance (Fig. 25b). It was also reported that the methylation of Si could improve the PEC performance by dealing with the dangling bonds on Si, suppressing the carrier recombination, increasing the donor concentration, decreasing the charge transfer

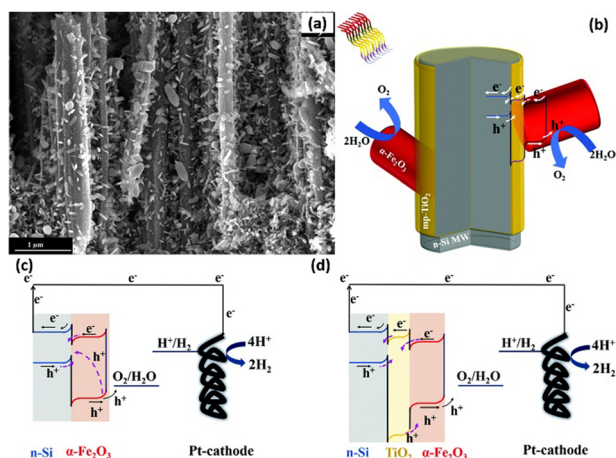


Fig. 24 SEM image (a) and schematic diagram (b) of the n-Si nanowires/TiO<sub>2</sub>/ $\alpha$ -Fe<sub>2</sub>O<sub>3</sub>. Band alignment of n-Si/ $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> (c) and n-Si/TiO<sub>2</sub>/ $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> (d). Reproduced with permission from ref. 155, Copyright © The Royal Society of Chemistry 2020.

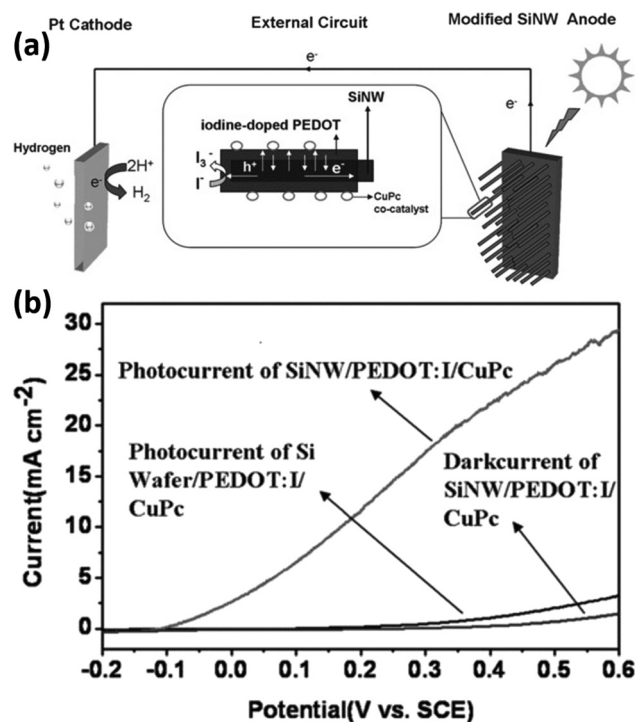


Fig. 25 (a) Schematic diagram of the hydrogen-production PEC cell. (b) PEC performance comparison of the different photoelectrodes. Reproduced with permission from ref. 160, Copyright © 2012 Wiley-VCH Verlag GmbH & Co. KGaA, Weinheim.



resistance, raising the flat band, and building a friendly and tight interface between the  $\text{CH}_3\text{-SiNWs}$  array and PEDOT.<sup>163</sup>

## 1D Si nanostructure homojunction

The narrow band gap of Si affords it excellent solar absorption properties but also determines the small photovoltage it can generate. As mentioned above, the use of other semiconductors to form heterojunctions with Si not only promotes the separation and transportation of photogenerated carriers, but also may compensate for the deficiencies caused by the energy level structure of Si. In addition to this, Si can be easily doped in different types and concentrations by well-established processes. With them, Si homojunctions (*e.g.*, pn junctions, p-i-n junctions, *etc.*) can be constructed to substantially increase the photovoltage of Si-based devices, which have been successfully applied to solar cells. Nakato *et al.* reported that preparing a layer of heavily doped p<sup>+</sup>-Si (or n<sup>+</sup>-Si) on the surface of n-Si (or p-Si) can effectively increase the photovoltage of Si photoelectrodes and achieve PEC hydrogen production at a lower overpotential.<sup>164–166</sup> This strategy has been applied to the 1D Si nanostructured photoelectrodes to increase the photovoltage and reduce the overpotential for PEC water splitting. Nathan S. Lewis *et al.* proposed the introduction of a buried metallurgical n<sup>+</sup>p junction into the Si wires.<sup>68,167</sup> Fig. 26a illustrates the preparation process of the n<sup>+</sup>p-Si wire arrays loaded with Pt co-catalysts.<sup>68</sup> When p-Si comes into contact with the electrolyte, the chemical potential (Fermi level) of p-Si is different from that of the H<sup>+</sup>/H<sub>2</sub> redox pair, and charge transfer between p-Si and the electrolyte occurs until equilibrium, and the energy band is bent downward by the formation of a depletion layer on the surface of p-Si, as shown in Fig. 26b. When intense light is illuminated to the p-Si/electrolyte system, the electrons and holes generated by photoexcitation are not transferred to the solution in time to be consumed or transferred to the Si interior. The newly generated net electrons will replenish the depletion layer, reducing the degree of band bending. The Fermi level splits into two “quasi-Fermi levels,”  $E_{f,n}$  and  $E_{f,p}$ , to describe the electrons and holes, respectively. The potential difference between  $E_{f,n}$  and  $E_{f,p}$  is the maximum photovoltage that the photoelectrode can generate. For the p-Si/n<sup>+</sup>-Si photoelectrodes, the Fermi energy level of n<sup>+</sup>-Si is consistent with the redox level of the H<sup>+</sup>/H<sub>2</sub> couple because the surface n<sup>+</sup>-Si layer is heavily doped and degenerates at room temperature.<sup>165</sup> The energy band structure of p-Si/n<sup>+</sup>-Si in contact with the electrolyte when it reaches equilibrium is shown in Fig. 26c. Under illumination, the potential difference between  $E_{f,n}$  and  $E_{f,p}$  is significantly larger than that of p-Si/electrolyte, indicating that p-Si/n<sup>+</sup>-Si/electrolyte can generate a higher photovoltage. It was confirmed that the photovoltage increased from 0.16 V to 0.54 V after introducing the n<sup>+</sup>-Si thin layer on the p-Si surface.<sup>68</sup> Similarly, p-i-n junction Si nanowire arrays were prepared for photoelectrocatalytic water splitting for hydrogen production.<sup>168</sup>

The photoelectrochemical properties of the Si nanowire arrays are averaged over numerous individual Si nanowires. Each Si nanowire may differ in geometry, doping, defects, and

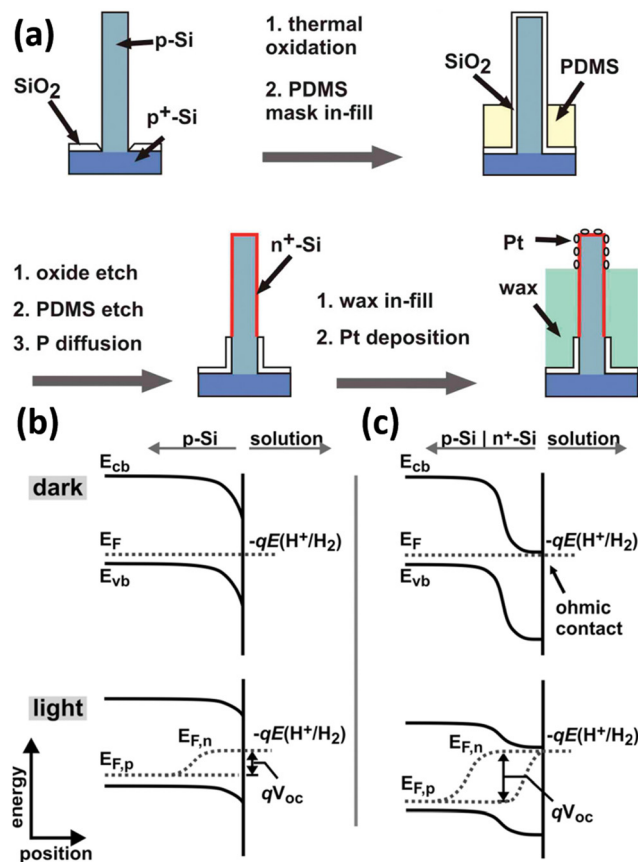


Fig. 26 (a) Schematic diagram of the preparation process of Pt/n<sup>+</sup>p-Si wire arrays. Band bending of (b) p-Si and (c) n<sup>+</sup>p-Si photocathodes in contact with electrolyte. Reproduced with permission from ref. 68, Copyright © 2011, American Chemical Society.

catalyst loading. This makes it challenging to correlate the average performance measured on the Si nanowire arrays well with these characteristics of the individual nanowires. In this regard, Yang *et al.* investigated the photoelectrochemistry of a single p-Si nanowire with an n<sup>+</sup>-Si thin layer on the surface.<sup>169</sup> The Si nanowires were prepared by the VLS mechanism with a length of 8–20 μm and a diameter of 500–800 nm. The p-Si and n<sup>+</sup>p-Si nanowires were obtained by post-treatment of the Si nanowires and then loaded with Pt nanoparticles as co-catalysts. The PEC response curve shows that the HER onset potential of n<sup>+</sup>p-Si nanowires is more positive than that of p-Si. Statistical results for several single nanowires tested show that the HER onset potential of n<sup>+</sup>p-Si is positively shifted by about 250 mV.

By constructing buried pn junctions, increasing the photovoltage of 1D Si nanowire photoelectrodes is possible. However, the photovoltage generated by 1 pn junction is less than 0.6 V. The theoretical voltage for water splitting is 1.23 V, and the actual voltage required is as high as 1.7 V after considering the overpotential. To this end, Teitworth *et al.* designed and fabricated Si nanowires encoded with p-i-n superlattice that can generate photovoltage of up to 10 V under 1 sun illumination.<sup>170</sup> The Si nanowires were prepared by the VLS process. CoO<sub>x</sub> as OER catalyst and Pt as HER catalyst were



modified at both ends of the Si nanowires to undergo water splitting, respectively. The photovoltage increased with the increasing number of p-i-n junctions, which was proved by the measurement of Si nanowires with various junctions. It was found that the photovoltage of a single junction is about 0.27 V. The photovoltage of the multijunction Si nanowires increases linearly with the number of p-i-n junctions. When the number of junctions was 40, the open-circuit voltage of the multijunction Si nanowire reached above 10 V under 1 sun illumination. Unassisted solar water splitting for hydrogen production was achieved based on such multi-junctional Si nanowires.

## Conclusions and perspectives

1D Si nanostructured photoelectrodes have shown great potential in PEC water splitting for hydrogen production. It also has some problems, such as high overpotential due to poor catalytic activity and inappropriate energy band structure. Several strategies have been developed in recent years to reduce the overpotential of 1D Si nanostructured photoelectrodes for PEC water splitting. The modification of 1D Si nanostructures with co-catalysts, including noble metal catalysts (Pt, Ru, Ag, *etc.*), non-precious metal catalysts (Ni, Ni-Mo, NiCo, NiFe, *etc.*), MoS<sub>x</sub> (Mo<sub>3</sub>S<sub>4</sub>, MoS<sub>2</sub>, *etc.*) and other metal compounds (metal sulfides, metal phosphides, metal borides, metal oxides, *etc.*), carbon materials (N-doped graphene, reduced graphene oxide, carbon nanowires, *etc.*), could effectively improve the kinetics of the PEC water splitting reactions on the surface of 1D Si nanostructures and reduces the overpotential. Moreover, Si can also be combined with other semiconductor materials to construct 1D Si nanostructured heterojunction photoelectrodes, which can facilitate carrier utilization and enhance photovoltage through energy band engineering, and ultimately reduce overpotential for water splitting reaction. In addition, the construction of 1D Si nanostructured homojunctions has also been used to reduce the overpotential by increasing the photovoltage of the photoelectrode using the pn junction. In particular, multiple p-i-n junctions are formed in a single Si nanowire, and the photovoltage can be regulated by adjusting the number of p-i-n junctions, which can generate photovoltage up to 10 V. Using such multi-junction Si nanowires, hydrogen production from water under illumination without external voltage has been successfully achieved.

Although significant progress has been achieved in reducing the overpotential of 1D nanostructured photoelectrodes for PEC hydrogen production, there are still many issues that deserve in-depth study. Currently, the modification of metals on the surface of 1D nanostructures is intended to be used as a co-catalyst to enhance the catalytic activity of the photoelectrodes. Schottky junctions formed by combining Si with metals can also be utilized to decrease the overpotential of the photoelectrodes. The choice of the appropriate metal and the proper modulation of the metal/Si interface can significantly increase the photovoltage.<sup>171-173</sup> In addition, some highly promising catalysts emerging in recent years, such as high-entropy

alloy catalysts,<sup>174</sup> triple ionic-electronic conducting oxides,<sup>175</sup> single-atom catalysts,<sup>176</sup> *etc.*, have not yet been applied in 1D Si nanostructured photoelectrodes. These new catalysts may further enhance the performance of 1D Si nanostructured photoelectrodes.

The slow kinetics of the OER half-reaction leads to an increase in the overpotential for the overall water splitting reaction. To address this problem, replacing the OER reaction with other reactions with lower oxidation potential and faster reaction kinetics is an effective strategy to reduce the overpotential of PEC hydrogen production.<sup>4,177-179</sup> Especially, coupling HER with some organic reactions that produce high-value chemicals can reduce the overpotential while lowering the cost of hydrogen production.<sup>180,181</sup> Although this strategy has not yet been applied in 1D Si nanostructured photoelectrodes, we believe it would be an effective way to reduce the overpotential of 1D Si nanostructured photoelectrodes for solar hydrogen production and reduce the cost of hydrogen production, which would help to improve the competitiveness of solar hydrogen production technology compared with the conventional technologies.

Computational methods based on density functional theory have successfully predicted the catalytic activity of catalysts, explained catalytic phenomena, and found new high-performance catalysts.<sup>182-184</sup> In addition, artificial intelligence (AI) techniques, which have emerged in recent years, are gradually being applied to materials and chemistry research fields.<sup>185,186</sup> The importance of these new computational-based tools in scientific research is becoming increasingly evident. Applying these new technologies in 1D Si nanostructure photovoltaic research is rare. It is conceivable that the development and application of high-performance 1D Si nanostructured photoelectrodes could be significantly facilitated if theoretical computing and AI technologies are utilized and fully exploited.

## Author contributions

This article was completed with the joint contribution of all authors. Guangwei She was responsible for the conceptualization and writing of the article. Xiaopeng Qi participated in the conception of the article and wrote and revised parts of it. Wensheng Shi was in overall charge of structuring and revising the article. Lixuan Mu was involved in discussing and revising parts of the article. Jingxuan Ma, Xue Hao, Changzhou Ru, and Haoyue Zhang conducted the literature review, analysis, and article revision.

## Conflicts of interest

There are no conflicts to declare.

## Acknowledgements

This work was supported by National Key R & D Program of China (2022YFB3803600) and the National Natural Science Foundation of China (21975269, 52172107).





## References

- Z. W. Seh, J. Kibsgaard, C. F. Dickens, I. Chorkendorff, J. K. Nørskov and T. F. Jaramillo, *Science*, 2017, **355**, eaad4998.
- T. Hisatomi and K. Domen, *Nat. Catal.*, 2019, **2**, 387–399.
- M. F. Lagadec and A. Grimaud, *Nat. Mater.*, 2020, **19**, 1140–1150.
- X. M. C. Ta, R. Daiyan, T. K. A. Nguyen, R. Amal, T. Tran-Phu and A. Tricoli, *Adv. Energy Mater.*, 2022, **12**, 2201358.
- M. Chatenet, B. G. Pollet, D. R. Dekel, F. Dionigi, J. Deseure, P. Millet, R. D. Braatz, M. Z. Bazant, M. Eikerling, I. Staffell, P. Balcombe, Y. Shao-Horn and H. Schäfer, *Chem. Soc. Rev.*, 2022, **51**, 4583–4762.
- G. Segev, J. Kibsgaard, C. Hahn, Z. J. Xu, W.-H. Cheng, T. G. Deutsch, C. Xiang, J. Z. Zhang, L. Hammarström, D. G. Nocera, A. Z. Weber, P. Agbo, T. Hisatomi, F. E. Osterloh, K. Domen, F. F. Abdi, S. Haussener, D. J. Miller, S. Ardo, P. C. McIntyre, T. Hannappel, S. Hu, H. Atwater, J. M. Gregoire, M. Z. Ertem, I. D. Sharp, K.-S. Choi, J. S. Lee, O. Ishitani, J. W. Ager, R. R. Prabhakar, A. T. Bell, S. W. Boettcher, K. Vincent, K. Takanabe, V. Artero, R. Napier, B. R. Cuenya, M. T. M. Koper, R. Van De Krol and F. Houle, *J. Phys. D: Appl. Phys.*, 2022, **55**, 323003.
- A. Aissa, N. A. A. Qasem, A. Mourad, H. Laidoudi, O. Younis, K. Guedri and A. Alazzam, *Appl. Therm. Eng.*, 2023, **220**, 119663.
- A. Shah, P. Torres, R. Tscharnner, N. Wyrsh and H. Keppner, *Science*, 1999, **285**, 692–698.
- H. Nishiyama, T. Yamada, M. Nakabayashi, Y. Maehara, M. Yamaguchi, Y. Kuromiya, Y. Nagatsuma, H. Tokudome, S. Akiyama, T. Watanabe, R. Narushima, S. Okunaka, N. Shibata, T. Takata, T. Hisatomi and K. Domen, *Nature*, 2021, **598**, 304–307.
- Q. Wang, C. Pornrungrroj, S. Linley and E. Reisner, *Nat. Energy*, 2022, **7**, 13–24.
- N. S. Lewis, *Science*, 2016, **351**, aad1920.
- P. Zhou, I. A. Navid, Y. Ma, Y. Xiao, P. Wang, Z. Ye, B. Zhou, K. Sun and Z. Mi, *Nature*, 2023, **613**, 66–70.
- M. R. Shaner, H. A. Atwater, N. S. Lewis and E. W. McFarland, *Energy Environ. Sci.*, 2016, **9**, 2354–2371.
- A. Heller, *Science*, 1984, **223**, 1141–1148.
- A. J. Shih and S. M. Haile, *Science*, 2022, **376**, 348–349.
- H. Song, S. Luo, H. Huang, B. Deng and J. Ye, *ACS Energy Lett.*, 2022, **7**, 1043–1065.
- K. Rajeshwar, R. McConnell and S. Licht, *Solar hydrogen generation: toward a renewable energy future*, Springer, 2008.
- K. Sivula and R. van de Krol, *Nat. Rev. Mater.*, 2016, **1**, 15010.
- A. Fujishima and K. Honda, *Nature*, 1972, **238**, 37–38.
- X. Qi, G. She, Y. Liu, L. Mu and W. Shi, *Chem. Commun.*, 2012, **48**, 242–244.
- L. Cheng, Q. Xiang, Y. Liao and H. Zhang, *Energy Environ. Sci.*, 2018, **11**, 1362–1391.
- X. Qi, G. She, M. Wang, L. Mu and W. Shi, *Chem. Commun.*, 2013, **49**, 5742–5744.
- C. A. Mesa, L. Francàs, K. R. Yang, P. Garrido-Barros, E. Pastor, Y. Ma, A. Kafizas, T. E. Rosser, M. T. Mayer, E. Reisner, M. Grätzel, V. S. Batista and J. R. Durrant, *Nat. Chem.*, 2020, **12**, 82–89.
- T. W. Kim and K.-S. Choi, *Science*, 2014, **343**, 990–994.
- V. Andrei, G. M. Ucoski, C. Pornrungrroj, C. Uswachoke, Q. Wang, D. S. Achilleos, H. Kasap, K. P. Sokol, R. A. Jagt, H. Lu, T. Lawson, A. Wagner, S. D. Pike, D. S. Wright, R. L. Z. Hoyer, J. L. MacManus-Driscoll, H. J. Joyce, R. H. Friend and E. Reisner, *Nature*, 2022, **608**, 518–522.
- S. Hu, M. R. Shaner, J. A. Beardslee, M. Lichterman, B. S. Brunshwig and N. S. Lewis, *Science*, 2014, **344**, 1005–1009.
- D. Kang, J. L. Young, H. Lim, W. E. Klein, H. Chen, Y. Xi, B. Gai, T. G. Deutsch and J. Yoon, *Nat. Energy*, 2017, **2**, 17043.
- M. J. Kenney, M. Gong, Y. Li, J. Z. Wu, J. Feng, M. Lanza and H. Dai, *Science*, 2013, **342**, 836–840.
- Y. Li, Y. Xiao, C. Wu, D. Zhang, J. Huang, Z. Zhang, J. He and C. Li, *ACS Mater. Lett.*, 2022, **4**, 779–804.
- K. Sun, S. Shen, Y. Liang, P. E. Burrows, S. S. Mao and D. Wang, *Chem. Rev.*, 2014, **114**, 8662–8719.
- Z. M. Ronglei Fan and M. Shen, *Opt. Express*, 2019, A51–A80.
- Z. Luo, T. Wang and J. Gong, *Chem. Soc. Rev.*, 2019, **48**, 2158–2181.
- S. Chandrasekaran, T. Nann and N. H. Voelcker, *Nano Energy*, 2015, **17**, 308–322.
- D. Liu, J. Ma, R. Long, C. Gao and Y. Xiong, *Nano Today*, 2017, **17**, 96–116.
- J.-W. Lee, K.-H. Cho, J.-S. Yoon, Y.-M. Kim and Y.-M. Sung, *J. Mater. Chem. A*, 2021, **9**, 21576–21606.
- J. Deng, Y. Su, D. Liu, P. Yang, B. Liu and C. Liu, *Chem. Rev.*, 2019, **119**, 9221–9259.
- J. M. Foley, M. J. Price, J. I. Feldblyum and S. Maldonado, *Energy Environ. Sci.*, 2012, **5**, 5203–5220.
- E. L. Warren, H. A. Atwater and N. S. Lewis, *J. Phys. Chem. C*, 2014, **118**, 747–759.
- B. Zhang, J. Jie, X. Zhang, X. Ou and X. Zhang, *ACS Appl. Mater. Interfaces*, 2017, **9**, 34527–34543.
- K.-Q. Peng, X. Wang, L. Li, Y. Hu and S.-T. Lee, *Nano Today*, 2013, **8**, 75–97.
- A. Jena, C.-J. Chen, H. Chang, S.-F. Hu and R.-S. Liu, *J. Mater. Chem. A*, 2021, **9**, 3767–3785.
- S. A. Lee, S. Choi, C. Kim, J. W. Yang, S. Y. Kim and H. W. Jang, *ACS Mater. Lett.*, 2020, **2**, 107–126.
- C. Cheng, W. Zhang, X. Chen, S. Peng and Y. Li, *Energy Sci. Eng.*, 2022, **10**, 1526–1543.
- M. G. Walter, E. L. Warren, J. R. McKone, S. W. Boettcher, Q. Mi, E. A. Santori and N. S. Lewis, *Chem. Rev.*, 2010, **110**, 6446–6473.
- E. M. Stuve, in *Encyclopedia of Applied Electrochemistry*, ed. G. Kreysa, K.-I. Ota and R. F. Savinell, Springer, New York, New York, NY, 2014, pp. 1445–1453, DOI: [10.1007/978-1-4419-6996-5\\_330](https://doi.org/10.1007/978-1-4419-6996-5_330).
- S. L. S. David Tilley and Roel van de Krol, *Advances in Photoelectrochemical Water Splitting: Theory, Experiment and Systems Analysis*, The Royal Society of Chemistry, 2018.



- 47 R. Memming, *Semiconductor Electrochemistry*, Wiley-VCH Verlag GmbH, 2nd edn, 2015.
- 48 *Semiconductor Electrodes and Photoelectrochemistry*, ed. Allen J. Bard, M. Stratmann, S. Licht, 2002.
- 49 X. G. Zhang, *Electrochemistry of Silicon and Its Oxide*, Springer, 2001.
- 50 C. Huo, J. Wang, H. Fu, X. Li, Y. Yang, H. Wang, A. Mateen, G. Farid and K.-Q. Peng, *Adv. Funct. Mater.*, 2020, **30**, 2005744.
- 51 L. Liu, K.-Q. Peng, Y. Hu, X.-L. Wu and S.-T. Lee, *Adv. Mater.*, 2014, **26**, 1410–1413.
- 52 K. Peng, A. Lu, R. Zhang and S.-T. Lee, *Adv. Funct. Mater.*, 2008, **18**, 3026–3035.
- 53 K. Peng, M. Zhang, A. Lu, N.-B. Wong, R. Zhang and S.-T. Lee, *Appl. Phys. Lett.*, 2007, **90**, 163123.
- 54 H. Chen, H. Wang, X.-H. Zhang, C.-S. Lee and S.-T. Lee, *Nano Lett.*, 2010, **10**, 864–868.
- 55 A. Mateen, G. Farid and K.-Q. Peng, *ECS J. Solid State Sci. Technol.*, 2022, **11**, 054006.
- 56 K.-Q. Peng, X. Wang, L. Li, X.-L. Wu and S.-T. Lee, *J. Am. Chem. Soc.*, 2010, **132**, 6872–6873.
- 57 W. A. Bryant, *J. Mater. Sci.*, 1977, **12**, 1285–1306.
- 58 S. N. Mohammad, in *Synthesis of Nanomaterials: Mechanisms, Kinetics and Materials Properties*, ed. S. N. Mohammad, Springer International Publishing, Cham, 2020, pp. 69–99, DOI: [10.1007/978-3-030-57585-4\\_5](https://doi.org/10.1007/978-3-030-57585-4_5).
- 59 V. Schmidt and U. Gösele, *Science*, 2007, **316**, 698–699.
- 60 W. Lu and C. M. Lieber, *J. Phys. D: Appl. Phys.*, 2006, **39**, R387.
- 61 Y. Wu and P. Yang, *J. Am. Chem. Soc.*, 2001, **123**, 3165–3166.
- 62 B. M. Kayes, M. A. Filler, M. C. Putnam, M. D. Kelzenberg, N. S. Lewis and H. A. Atwater, *Appl. Phys. Lett.*, 2007, **91**, 103110.
- 63 T. Shimizu, T. Xie, J. Nishikawa, S. Shingubara, S. Senz and U. Gösele, *Adv. Mater.*, 2007, **19**, 917–920.
- 64 R. Q. Zhang, Y. Lifshitz and S. T. Lee, *Adv. Mater.*, 2003, **15**, 635–640.
- 65 A. T. Heitsch, D. D. Fanfair, H.-Y. Tuan and B. A. Korgel, *J. Am. Chem. Soc.*, 2008, **130**, 5436–5437.
- 66 B. Fuhrmann, H. S. Leipner, H.-R. Höche, L. Schubert, P. Werner and U. Gösele, *Nano Lett.*, 2005, **5**, 2524–2527.
- 67 C.-M. Hsu, S. T. Connor, M. X. Tang and Y. Cui, *Appl. Phys. Lett.*, 2008, **93**, 133109.
- 68 S. W. Boettcher, E. L. Warren, M. C. Putnam, E. A. Santori, D. Turner-Evans, M. D. Kelzenberg, M. G. Walter, J. R. McKone, B. S. Brunschwig, H. A. Atwater and N. S. Lewis, *J. Am. Chem. Soc.*, 2011, **133**, 1216–1219.
- 69 J. R. McKone, E. L. Warren, M. J. Bierman, S. W. Boettcher, B. S. Brunschwig, N. S. Lewis and H. B. Gray, *Energy Environ. Sci.*, 2011, **4**, 3573–3583.
- 70 M. D. Kelzenberg, S. W. Boettcher, J. A. Petykiewicz, D. B. Turner-Evans, M. C. Putnam, E. L. Warren, J. M. Spurgeon, R. M. Briggs, N. S. Lewis and H. A. Atwater, *Nat. Mater.*, 2010, **9**, 239–244.
- 71 K. Peng, X. Wang and S.-T. Lee, *Appl. Phys. Lett.*, 2008, **92**, 163103.
- 72 X. Wang, K.-Q. Peng, X.-J. Pan, X. Chen, Y. Yang, L. Li, X.-M. Meng, W.-J. Zhang and S.-T. Lee, *Angew. Chem., Int. Ed.*, 2011, **50**, 9861–9865.
- 73 J. Zhu, Z. Yu, G. F. Burkhard, C.-M. Hsu, S. T. Connor, Y. Xu, Q. Wang, M. McGehee, S. Fan and Y. Cui, *Nano Lett.*, 2009, **9**, 279–282.
- 74 K. Peng, Y. Yan, S. Gao and J. Zhu, *Adv. Funct. Mater.*, 2003, **13**, 127–132.
- 75 K. Peng, Y. Wu, H. Fang, X. Zhong, Y. Xu and J. Zhu, *Angew. Chem., Int. Ed.*, 2005, **44**, 2737–2742.
- 76 R. P. Srivastava and D.-Y. Khang, *Adv. Mater.*, 2021, **33**, 2005932.
- 77 J.-Y. Jung, Z. Guo, S.-W. Jee, H.-D. Um, K.-T. Park and J.-H. Lee, *Opt. Express*, 2010, **18**, A286–A292.
- 78 S. Huang, H. Zhang, Z. Wu, D. Kong, D. Lin, Y. Fan, X. Yang, Z. Zhong, S. Huang, Z. Jiang and C. Cheng, *ACS Appl. Mater. Interfaces*, 2014, **6**, 12111–12118.
- 79 X. Li, Y. Xiao, K. Zhou, J. Wang, S. L. Schweizer, A. Sprafke, J.-H. Lee and R. B. Wehrspohn, *Phys. Chem. Chem. Phys.*, 2015, **17**, 800–804.
- 80 Y. J. Hwang, C. H. Wu, C. Hahn, H. E. Jeong and P. Yang, *Nano Lett.*, 2012, **12**, 1678–1682.
- 81 I. Oh, J. Kye and S. Hwang, *Nano Lett.*, 2012, **12**, 298–302.
- 82 *Electrochemical Dictionary – 2nd, Revised and Extended Edition*, ed. Allen J. Bard, G. Inzelt, F. Scholz, 2012.
- 83 J. K. Nørskov, T. Bligaard, A. Logadottir, J. R. Kitchin, J. G. Chen, S. Pandalov and U. Stimming, *J. Electrochem. Soc.*, 2005, **152**, J23.
- 84 J. Ji, H. Zhang, Y. Qiu, Y. Wang, Y. Luo and L. Hu, *J. Mater. Sci.: Mater. Electron.*, 2013, **24**, 4433–4438.
- 85 L. Qiao, M. Zhou, Y. Li, A. Zhang, J. Deng, M. Liao, P. Xiao, Y. Zhang and S. Zhang, *J. Electrochem. Soc.*, 2014, **161**, H458–H463.
- 86 P. Dai, J. Xie, M. T. Mayer, X. Yang, J. Zhan and D. Wang, *Angew. Chem., Int. Ed.*, 2013, **52**, 11119–11123.
- 87 N. P. Dasgupta, C. Liu, S. Andrews, F. B. Prinz and P. Yang, *J. Am. Chem. Soc.*, 2013, **135**, 12932–12935.
- 88 L. Zhu, Q. Cai, F. Liao, M. Sheng, B. Wu and M. Shao, *Electrochem. Commun.*, 2015, **52**, 29–33.
- 89 F. Liao, B. Jiang, W. Shen, Y. Chen, Y. Li, Y. Shen, K. Yin and M. Shao, *ChemCatChem*, 2019, **11**, 2126–2130.
- 90 U. Dadwal, N. Ranjan and R. Singh, *AIP Conf. Proc.*, 2016, **1731**, 050014.
- 91 X. Gao, S. Wu, J. Yan, X. Zhai and X. Li, *ACS Appl. Mater. Interfaces*, 2016, **8**, 30072–30078.
- 92 E. L. Warren, J. R. McKone, H. A. Atwater, H. B. Gray and N. S. Lewis, *Energy Environ. Sci.*, 2012, **5**, 9653–9661.
- 93 M. R. Shaner, J. R. McKone, H. B. Gray and N. S. Lewis, *Energy Environ. Sci.*, 2015, **8**, 2977–2984.
- 94 C.-W. Tung, C.-H. Hou, H.-T. Lin, Y. Zheng, Y.-P. Huang, Y.-F. Liao, J.-J. Shiue and H. M. Chen, *Solar RRL*, 2020, **4**, 2000028.
- 95 W. Vajselaar, P. Westerik, J. Veerbeek, R. M. Tiggelaar, E. Berenschot, N. R. Tas, H. Gardeniers and J. Huskens, *Nat. Energy*, 2018, **3**, 185–192.



- 96 C.-W. Tung, Y. Chuang, H.-C. Chen, T.-S. Chan, J.-Y. Li and H. M. Chen, *Part. Part. Syst. Charact.*, 2018, **35**, 1700321.
- 97 T. F. Jaramillo, K. P. Jørgensen, J. Bonde, J. H. Nielsen, S. Horch and I. Chorkendorff, *Science*, 2007, **317**, 100–102.
- 98 B. Hinnemann, P. G. Moses, J. Bonde, K. P. Jørgensen, J. H. Nielsen, S. Horch, I. Chorkendorff and J. K. Nørskov, *J. Am. Chem. Soc.*, 2005, **127**, 5308–5309.
- 99 P. D. Tran, S. S. Pramana, V. S. Kale, M. Nguyen, S. Y. Chiam, S. K. Batabyal, L. H. Wong, J. Barber and J. Loo, *Chem. – Eur. J.*, 2012, **18**, 13994–13999.
- 100 L. Zhang, C. Liu, A. B. Wong, J. Resasco and P. Yang, *Nano Res.*, 2015, **8**, 281–287.
- 101 L. Qiao, M. Liao, K. Fang, X. He and Y. Zhang, *Silicon*, 2019, **11**, 1963–1970.
- 102 Q. Ding, F. Meng, C. R. English, M. Cabán-Acevedo, M. J. Shearer, D. Liang, A. S. Daniel, R. J. Hamers and S. Jin, *J. Am. Chem. Soc.*, 2014, **136**, 8504–8507.
- 103 S. M. Younan, Z. Li, M. P. Fairchild, N. B. Williams, Y. Huang and J. Gu, *Adv. Mater. Interfaces*, 2022, **9**, 2200178.
- 104 L. Qiao, M. Liao, J. Wu, Q. Jiang, Y. Zhang and Y. Li, *Int. J. Hydrogen Energy*, 2018, **43**, 22235–22242.
- 105 Z. Huang, C. Wang, L. Pan, F. Tian, X. Zhang and C. Zhang, *Nano Energy*, 2013, **2**, 1337–1346.
- 106 G.-L. Zang, G.-P. Sheng, C. Shi, Y.-K. Wang, W.-W. Li and H.-Q. Yu, *Energy Environ. Sci.*, 2014, **7**, 3033–3039.
- 107 S. H. Hsieh, S. T. Ho and W. J. Chen, *J. Nanomater.*, 2016, **2016**, 6974646.
- 108 X.-Q. Bao, D. Y. Petrovykh, P. Alpuim, D. G. Stroppa, N. Guldris, H. Fonseca, M. Costa, J. Gaspar, C. Jin and L. Liu, *Nano Energy*, 2015, **16**, 130–142.
- 109 Y. Hou, B. L. Abrams, P. C. K. Vesborg, M. E. Björketun, K. Herbst, L. Bech, A. M. Setti, C. D. Damsgaard, T. Pedersen, O. Hansen, J. Rossmeisl, S. Dahl, J. K. Nørskov and I. Chorkendorff, *Nat. Mater.*, 2011, **10**, 434.
- 110 Z. Huang, C. Wang, Z. Chen, H. Meng, C. Lv, Z. Chen, R. Han and C. Zhang, *ACS Appl. Mater. Interfaces*, 2014, **6**, 10408–10414.
- 111 C.-J. Chen, P.-T. Chen, M. Basu, K.-C. Yang, Y.-R. Lu, C.-L. Dong, C.-G. Ma, C.-C. Shen, S.-F. Hu and R.-S. Liu, *J. Mater. Chem. A*, 2015, **3**, 23466–23476.
- 112 S. Gopalakrishnan, G. Paulraj, M. K. Eswaran, A. Ray, N. Singh and K. Jeganathan, *Chemosphere*, 2022, **302**, 134708.
- 113 E. J. Popczun, C. G. Read, C. W. Roske, N. S. Lewis and R. E. Schaak, *Angew. Chem., Int. Ed.*, 2014, **53**, 5427–5430.
- 114 J. Tian, Q. Liu, A. M. Asiri and X. Sun, *J. Am. Chem. Soc.*, 2014, **136**, 7587–7590.
- 115 X.-Q. Bao, M. F. Cerqueira, P. Alpuim and L. Liu, *Chem. Commun.*, 2015, **51**, 10742–10745.
- 116 S. M. Thalluri, J. Borme, K. Yu, J. Xu, I. Amorim, J. Gaspar, L. Qiao, P. Ferreira, P. Alpuim and L. Liu, *Nano Res.*, 2018, **11**, 4823–4835.
- 117 C. W. Roske, E. J. Popczun, B. Seger, C. G. Read, T. Pedersen, O. Hansen, P. C. K. Vesborg, B. S. Brunshwig, R. E. Schaak, I. Chorkendorff, H. B. Gray and N. S. Lewis, *J. Phys. Chem. Lett.*, 2015, **6**, 1679–1683.
- 118 C. Lv, Z. Chen, Z. Chen, B. Zhang, Y. Qin, Z. Huang and C. Zhang, *J. Mater. Chem. A*, 2015, **3**, 17669–17675.
- 119 Z. Huang, Z. Chen, Z. Chen, C. Lv, H. Meng and C. Zhang, *ACS Nano*, 2014, **8**, 8121–8129.
- 120 H. Zhang, Q. Ding, D. He, H. Liu, W. Liu, Z. Li, B. Yang, X. Zhang, L. Lei and S. Jin, *Energy Environ. Sci.*, 2016, **9**, 3113–3119.
- 121 S. Lee, S. Cha, Y. Myung, K. Park, I. H. Kwak, I. S. Kwon, J. Seo, S. A. Lim, E. H. Cha and J. Park, *ACS Appl. Mater. Interfaces*, 2018, **10**, 33198–33204.
- 122 G. She, S. Zhang, S. Li, F. Qu, L. Mu and W. Shi, *J. Photochem. Photobiol., A*, 2018, **355**, 55–59.
- 123 G. Yan, Y. Dong, T. Wu, S. Xing and X. Wang, *ACS Appl. Mater. Interfaces*, 2021, **13**, 52912–52920.
- 124 S. Zhang, G. She, S. Li, M. Lixuan and S. Wensheng, *Nanotechnology*, 2018, **29**, 035601.
- 125 S. Zhang, G. She, L. Mu and W. Shi, *Mater. Lett.*, 2017, **197**, 131–134.
- 126 Y. Yang, M. Wang, P. Zhang, W. Wang, H. Han and L. Sun, *ACS Appl. Mater. Interfaces*, 2016, **8**, 30143–30151.
- 127 S. Y. Kim, R. Sivasankaran, M. Arunachalam, G. H. Lee, D. H. Youn and S. H. Kang, *J. Electrochem. Soc.*, 2021, **168**, 066519.
- 128 S. Chandrasekaran, T. Nann and N. H. Voelcker, *Nanomaterials*, 2016, **6**, 144.
- 129 X. Zhang and J. He, *J. Nanosci. Nanotechnol.*, 2018, **18**, 5674–5678.
- 130 U. Sim, T.-Y. Yang, J. Moon, J. An, J. Hwang, J.-H. Seo, J. Lee, K. Y. Kim, J. Lee, S. Han, B. H. Hong and K. T. Nam, *Energy Environ. Sci.*, 2013, **6**, 3658–3664.
- 131 U. Sim, J. Moon, J. An, J. H. Kang, S. E. Jerng, J. Moon, S.-P. Cho, B. H. Hong and K. T. Nam, *Energy Environ. Sci.*, 2015, **8**, 1329–1338.
- 132 H. Meng, K. Fan, J. Low and J. Yu, *Dalton Trans.*, 2016, **45**, 13717–13725.
- 133 J. Moon, U. Sim, D. J. Kim, H.-Y. Ahn, J. An, H. Ha, K. S. Choi, C. Jeon, J. Lee, K. T. Nam and B. H. Hong, *Nanoscale*, 2018, **10**, 13936–13941.
- 134 R. R. Devarapalli, C. K. Kamaja and M. V. Shelke, *J. Mater. Chem. A*, 2014, **2**, 13352–13358.
- 135 Z. Chen, G. Ma, Z. Chen, Y. Zhang, Z. Zhang, J. Gao, Q. Meng, M. Yuan, X. Wang, J.-M. Liu and G. Zhou, *Appl. Surf. Sci.*, 2017, **396**, 609–615.
- 136 S. Gopalakrishnan, G. M. Bhalerao and K. Jeganathan, *J. Power Sources*, 2019, **413**, 293–301.
- 137 S. Gopalakrishnan, G. M. Bhalerao and K. Jeganathan, *ACS Sustainable Chem. Eng.*, 2019, **7**, 13911–13919.
- 138 S. Bai, J. Jiang, Q. Zhang and Y. Xiong, *Chem. Soc. Rev.*, 2015, **44**, 2893–2939.
- 139 W. Baicheng, X. Fenghua and X. Jianguang, *Nanotechnology*, 2014, **25**, 455402.
- 140 R. Yang, Y. Ji, Q. Li, Z. Zhao, R. Zhang, L. Liang, F. Liu, Y. Chen, S. Han, X. Yu and H. Liu, *Appl. Catal., B*, 2019, **256**, 117798.
- 141 S. Riyajuddin, J. Sultana, S. A. Siddiqui, S. Kumar, D. Badhwar, S. S. Yadav, S. Goyal, A. Venkatesan,





- S. Chakraverty and K. Ghosh, *Sustainable Energy Fuels*, 2022, **6**, 197–208.
- 142 X. Li, C. Garlisi, Q. Guan, S. Anwer, K. Al-Ali, G. Palmisano and L. Zheng, *Mater. Today*, 2021, **47**, 75–107.
- 143 Y. J. Hwang, A. Boukai and P. Yang, *Nano Lett.*, 2009, **9**, 410–415.
- 144 M. Shi, X. Pan, W. Qiu, D. Zheng, M. Xu and H. Chen, *Int. J. Hydrogen Energy*, 2011, **36**, 15153–15159.
- 145 F.-Q. Zhang, Y. Hu, R.-N. Sun, H. Fu and K.-Q. Peng, *Front. Chem.*, 2019, **7**, 206.
- 146 I. V. Bagal, M. Arunachalam, A. Abdullah, A. Waseem, M. A. Kulkarni, S. Hyung Kang and S.-W. Ryu, *J. Energy Chem.*, 2022, **71**, 45–55.
- 147 M. T. Mayer, C. Du and D. Wang, *J. Am. Chem. Soc.*, 2012, **134**, 12406–12409.
- 148 X. Wang, K.-Q. Peng, Y. Hu, F.-Q. Zhang, B. Hu, L. Li, M. Wang, X.-M. Meng and S.-T. Lee, *Nano Lett.*, 2013, **14**, 18–23.
- 149 X. Qi, G. She, X. Huang, T. Zhang, H. Wang, L. Mu and W. Shi, *Nanoscale*, 2014, **6**, 3182–3189.
- 150 R. R. Devarapalli, C. K. Kamaja and M. V. Shelke, *Chemistry-Select*, 2017, **2**, 2544–2551.
- 151 Z. Zhou, S. Wu, L. Qin, L. Li, L. Li and X. Li, *J. Mater. Chem. A*, 2018, **6**, 15593–15602.
- 152 Y. Qu, F. Li, P. Zhang, L. Zhao, J. Liu, X. Song and L. Gao, *Appl. Surf. Sci.*, 2019, **471**, 528–536.
- 153 Z. Zhou, L. Li, Y. Niu, H. Song, X.-S. Xing, Z. Guo and S. Wu, *Dalton Trans.*, 2021, **50**, 2936–2944.
- 154 W. Zhang, H. Chen, L. Zhang, P. Zhang, E. Dong, J. Ma and G. Wang, *J. Alloys Compd.*, 2019, **773**, 597–604.
- 155 A. Kawde, A. Annamalai, A. Sellstedt, J. Uhlig, T. Wågberg, P. Glatzel and J. Messinger, *Phys. Chem. Chem. Phys.*, 2020, **22**, 28459–28467.
- 156 Q. Liu, F. Wu, F. Cao, L. Chen, X. Xie, W. Wang, W. Tian and L. Li, *Nano Res.*, 2015, **8**, 3524–3534.
- 157 Z. Chen, M. Ning, G. Ma, Q. Meng, Y. Zhang, J. Gao, M. Jin, Z. Chen, M. Yuan, X. Wang, J.-M. Liu and G. Zhou, *Nanotechnology*, 2017, **28**, 275401.
- 158 K.-Q. Lu, X. Lin, Z.-R. Tang and Y.-J. Xu, *Catal. Today*, 2019, **335**, 294–299.
- 159 S.-C. Shiu, J.-J. Chao, S.-C. Hung, C.-L. Yeh and C.-F. Lin, *Chem. Mater.*, 2010, **22**, 3108–3113.
- 160 T. Yang, H. Wang, X.-M. Ou, C.-S. Lee and X.-H. Zhang, *Adv. Mater.*, 2012, **24**, 6199–6203.
- 161 C. Duan, H. Wang, X. Ou, F. Li and X. Zhang, *ACS Appl. Mater. Interfaces*, 2014, **6**, 9742–9750.
- 162 Z. Liu, H. Wang, X.-M. Ou, C.-S. Lee and X.-H. Zhang, *Chem. Commun.*, 2012, **48**, 2815–2817.
- 163 X. Zhang and J. He, *J. Nanosci. Nanotechnol.*, 2018, **18**, 2844–2849.
- 164 Y. Nakato, A. Tsumura and H. Tsubomura, *Chem. Lett.*, 1982, 1071–1074.
- 165 Y. Nakato, Y. Egi, M. Hiramoto and H. Tsubomura, *J. Phys. Chem.*, 1984, **88**, 4218–4222.
- 166 Y. Nakato, M. Yoshimura, M. Hiramoto, A. Tsumura, T. Murahashi and H. Tsubomura, *Bull. Chem. Soc. Jpn.*, 1984, **57**, 355–360.
- 167 M. R. Shaner, M. T. McDowell, A. Pien, H. A. Atwater and N. S. Lewis, *J. Electrochem. Soc.*, 2016, **163**, H261–H264.
- 168 S. Riyajuddin, S. Kumar, D. Badhwar, S. A. Siddiqui, J. Sultana and K. Ghosh, *Sustainable Energy Fuels*, 2021, **5**, 3160–3171.
- 169 Y. Su, C. Liu, S. Brittan, J. Tang, A. Fu, N. Kornienko, Q. Kong and P. Yang, *Nat. Nano*, 2016, **11**, 609–612.
- 170 T. S. Teitsworth, D. J. Hill, S. R. Litvin, E. T. Ritchie, J.-S. Park, J. P. Custer, A. D. Taggart, S. R. Bottum, S. E. Morley, S. Kim, J. R. McBride, J. M. Atkin and J. F. Cahoon, *Nature*, 2023, **614**, 270–274.
- 171 H. Zhang, G. She, J. Xu, S. Li, Y. Liu, J. Luo and W. Shi, *J. Mater. Chem. A*, 2022, **10**, 4952–4959.
- 172 S. Li, H. Zhang, G. She, J. Xu, S. Zhang, Y. Deng, L. Mu, Q. Zhou, Y. Liu, J. Luo and W. Shi, *ACS Appl. Energy Mater.*, 2021, **4**, 11574–11579.
- 173 S. Li, G. She, J. Xu, S. Zhang, H. Zhang, L. Mu, C. Ge, K. Jin, J. Luo and W. Shi, *ACS Appl. Mater. Interfaces*, 2020, **12**, 39092–39097.
- 174 Y. Luo, Z. Zhang, M. Chhowalla and B. Liu, *Adv. Mater.*, 2022, **34**, 2108133.
- 175 M. Papac, V. Stevanović, A. Zakutayev and R. O'Hayre, *Nat. Mater.*, 2021, **20**, 301–313.
- 176 L. Zhang, Y. Ren, W. Liu, A. Wang and T. Zhang, *Natl. Sci. Rev.*, 2018, **5**, 653–672.
- 177 B. Mei, G. Mul and B. Seger, *Adv. Sustainable Sys.*, 2017, **1**, 1600035.
- 178 A. Y. Ru Ng, B. Boruah, K. F. Chin, J. M. Modak and H. S. Soo, *ChemNanoMat*, 2020, **6**, 185–203.
- 179 Y. Xu and B. Zhang, *ChemElectroChem*, 2019, **6**, 3214–3226.
- 180 G. Li, G. Han, L. Wang, X. Cui, N. K. Moehring, P. R. Kidambi, D.-E. Jiang and Y. Sun, *Nat. Commun.*, 2023, **14**, 525.
- 181 Z. Li, Y. Yan, S.-M. Xu, H. Zhou, M. Xu, L. Ma, M. Shao, X. Kong, B. Wang, L. Zheng and H. Duan, *Nat. Commun.*, 2022, **13**, 147.
- 182 A. Jain, Y. Shin and K. A. Persson, *Nat. Rev. Mater.*, 2016, **1**, 15004.
- 183 J. K. Nørskov, T. Bligaard, J. Rossmeisl and C. H. Christensen, *Nat. Chem.*, 2009, **1**, 37–46.
- 184 J. Resasco, F. Abild-Pedersen, C. Hahn, Z. Bao, M. T. M. Koper and T. F. Jaramillo, *Nat. Catal.*, 2022, **5**, 374–381.
- 185 A. Mazheika, Y.-G. Wang, R. Valero, F. Viñes, F. Illas, L. M. Ghiringhelli, S. V. Levchenko and M. Scheffler, *Nat. Commun.*, 2022, **13**, 419.
- 186 P.-L. Kang, Y.-F. Shi, C. Shang and Z.-P. Liu, *Chem. Sci.*, 2022, **13**, 8148–8160.

