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*In situ* protonated-phosphorus interstitial doping induces long-lived shallow charge trapping in porous  $C_{3-x}N_4$  photocatalysts for highly efficient  $H_2$  generation

Professor Z. X. Guo and Professor D. L. Phillips with co-workers report a fundamental discovery of an *in situ* proton-mediated mechanism to promote the photodynamics and hence charge separation in an interstitial-phosphorus doped g- $C_{3-x}N_4$  catalyst. The *in situ* protonation of the phosphorous site creates an effective atomic heterojunction with active shallow trap states to enhance the life-time of the active hot electrons for effective solar water-splitting to generate “net-zero” hydrogen molecules.

### As featured in:



See David Lee Phillips, Zhengxiao Guo *et al.*, *Energy Environ. Sci.*, 2023, **16**, 460.

## PAPER

View Article Online  
View Journal | View IssueCite this: *Energy Environ. Sci.*, 2023, 16, 460

# *In situ* protonated-phosphorus interstitial doping induces long-lived shallow charge trapping in porous $C_{3-x}N_4$ photocatalysts for highly efficient $H_2$ generation†

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Efficient photocatalytic solar-to- $H_2$  conversion is pivotal to zero-carbon energy supply. Graphitic carbon nitride ( $g-C_3N_4$ ) is a promising visible-light photocatalyst but suffers from intrinsic electron-hole recombination and deep-charge trapping, limiting its efficiency. Here, we show a synergistic strategy of porosity, vacancy and shallow(trapping)-state engineering to enrich catalytic sites and promote the lifetime of active electrons by thermochemical treatment and phosphorus-interstitial-doping. The latter enhances the electron delocalization in the  $\pi$ -conjugate polymeric structure. The optimized photocatalyst shows a  $\sim 800\%$  increase in  $H_2$  generation ( $6323 \mu\text{mol h}^{-1} \text{g}^{-1}$ ) and an about 5-fold increase in quantum efficiency ( $QE_{420 \text{ nm}} = 5.08\%$ ). The superior performance is attributed to the long-lived shallow charge trapping, as a result of proton-feeding to the coordinated phosphorus site during the photocatalytic reaction, which enhances the photogenerated carrier lifetime and positively optimizes the band structure of the catalyst. Femtosecond transient absorption spectroscopy reveals a doubling lifetime of shallow-trapped charges ( $\sim 405.5 \text{ ps}$ ), favoring high mobility for electron-involved photocatalytic  $H_2$  generation. This work provides a new mechanism for improving charge carrier dynamics and photocatalytic performance.

Received 19th August 2022,  
Accepted 28th November 2022

DOI: 10.1039/d2ee02680e

rsc.li/ees

## Broader context

Photocatalytic water-splitting holds great potential for the cost-effective large-scale production of green hydrogen ( $H_2$ ) from abundant solar energy and water molecules. The key challenges are to develop a highly efficient and stable photocatalyst that can promote the effective photodynamics of charge carriers, *i.e.* electrons/holes ( $e^-/h^+$ ), to drive their sufficient separation and bulk-to-surface transfer, while maintaining its performance durability. Herein, we report an effective atomic heterojunction formed by porosity-stabilised interstitial phosphorus(P)-doping, along with its *in situ* protonation, to induce electron shallow-trap states in a carbon-lean graphitic carbon nitride ( $g-C_{3-x}N_4$ ) photocatalyst. Experiments and theoretical calculations show that the key roles of the *in situ* protonated interstitial P are in (1) restraining undesirable deep charge trapping, (2) prolonging the lifetime of the active charges and (3) forming a spin-enhanced structure. The resulting catalyst achieved 800% enhancement of the  $H_2$  generation rate at  $6323 \mu\text{mol h}^{-1} \text{g}^{-1}$  with an internal quantum efficiency of 31.8% under visible-light excitation. Due to the durable “ $-N-P=N-$ ” junction for efficient charge transfer and the abundant proton source in water, the as-synthesized catalyst shows superior photostability for  $H_2$  production, compared to its pristine counterpart. This report provides new insights into the proton-/electron-involved photocatalytic  $H_2$  production mechanism by transient time-resolved spectroscopic studies, and provides a novel *in situ* proton-assisted mechanistic model for effective solar-to- $H_2$  conversion in solution.

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† Electronic supplementary information (ESI) available: Details of experimental methods and calculation parameters; SEM and STEM images; EDX elements' linear distribution and XPS spectra; FTIR profiles; comparison of photocatalytic  $H_2$  production activity; Tafel plots; the fs-TAS and ns-TAS setup; exponential function fitted parameters of fs-TAS and ns-TAS; and calculated DOS and formation energy. See DOI: <https://doi.org/10.1039/d2ee02680e>





## Introduction

Photocatalytic water-splitting is highly desirable to generate green  $H_2$  with net-zero carbon emission.<sup>1–5</sup> Considerable efforts have been devoted to the process since its inception in 1970s,<sup>6</sup> where both thermodynamic requirements (*e.g.*, band-gap, redox potential *vs.*  $H^+/H_2$ ) and dynamic efficiency (*e.g.*, charge separation/accumulation) are highly important.<sup>7–10</sup> Many investigations have focused on the development of semi-conductive materials with suitable band structures that can accelerate charge carrier migration and enhance optical absorption, as outlined recently.<sup>11–14</sup> The engineering of single-component photocatalysts, *e.g.* by vacancy-assisted electronic structure tuning, phase transformation and single-atom doping, has recently been rekindled to boost photocatalytic performance with well-tuned band edges and enriched catalytic sites. However, even with well-aligned band-structures, single-component photocatalysts often show insufficient driving force for water-splitting due to their photodynamic losses, *e.g.*, by intrinsic emission and deep charge trapping. On the other hand, two-component photocatalysts (*e.g.*,  $TiO_2/MoS_2$ ,  $WO_3/BiVO_4$  and  $MoS_2/WS_2$ ) are under popular consideration to enhance efficiency by fast interfacial charge carrier transfer *via* a built-in heterojunction.<sup>15–17</sup> Unfortunately, as with most of the Type-II and Z-scheme heterojunction photocatalytic systems, the enhanced dynamic efficiency comes at the cost of a reduced redox potential when two coupled catalysts exhibit staggered band edges. It is very challenging to develop a catalyst for the dynamically effective conversion of solar energy into  $H_2$  with high redox capability and stability.

With due consideration of redox-potential tuning and enrichment of catalytic centers, it is found that two-dimensional conjugated polymeric materials are attractive candidates for high-efficiency catalysis.<sup>18</sup> For instance, graphitic carbon nitride ( $g-C_3N_4$ ) shows tremendous promise for cost-effective photocatalytic water-splitting due to its superior visible-light activity and tunable atomic structure.<sup>19–22</sup> However, its overall performance for water-splitting is limited by the rapid decay of the excited states *via* charge carrier recombination, deep trapping and poor redox potential, as the band-edge is close to the  $H^+/H_2$  potential.<sup>23</sup> The intrinsic deep-charge trapping features of the pristine  $g-C_3N_4$  constrain active electrons transport, directly reducing the dynamic efficiency.<sup>24</sup> A viable strategy to suppress deep trapping is to create long-lived active electrons *via* promotion of electron shallow trapping at a relatively high energy level. In our previous studies, heterojunctions and near-field assistance were shown to improve the charge separation lifetime and also promote shallow charge trapping of  $g-C_3N_4$ , respectively, leading to efficient visible-light-driven photocatalytic  $H_2$  production.<sup>25–27</sup> The enhancement of hot-carrier production was verified by Xie *et al.* over an order-to-disorder transformation of a  $g-C_3N_4$  photocatalyst, *via* accelerated exciton dissociation.<sup>28</sup> In addition, defect engineering can be an effective way of adjusting the Fermi level of semiconductors to improve interfacial charge carrier transfer and electron-hole separation. Ye *et al.* have reported a widened Fermi level mismatch strategy by tuning the N-vacancy density of a  $g-C_3N_4$ -based material for efficient

photocatalysis.<sup>29</sup> Moreover, atomic vacancies of  $g-C_3N_4$  play an important role in exposing catalytic and binding sites to enhance heteroatomic coordination in new chemical interactions. For example, Zeng and co-workers impart nitric oxide photo-conversion activity to C-vacancy mediated Pd-doped  $g-C_3N_4$ , as a result of ready formation and stabilization of a Pd–N bond.<sup>30</sup> Hence, defect engineering is an effective approach to bind heteroatoms and tune band-edge positions. However, it is unclear if such an atomic defect may promote high-frequency and/or lifetime of shallow trapping and change the charge carrier transport pathway.

Recently, considerable effort has also been devoted to non-metallic heteroatom (*e.g.*, S, O and B) doping of  $g-C_3N_4$  to boost photocatalytic water-splitting by the engineering of effective catalytic sites and band edges.<sup>31–33</sup> For example, Shen *et al.* reported a self-Z-scheme photocatalyst of boron-doped and pristine  $g-C_3N_4$  with a well-tuned band structure to drive photocatalytic overall water-splitting in the presence of Pt and  $Co(OH)_2$  co-catalysts.<sup>34</sup> Notably, this boron doping strategy not only enhances the visible-light absorption efficiency of  $g-C_3N_4$  but also accelerates the interfacial charge carrier mobility *via* band edge alignment. Other research suggested that the ideal dopants should not perturb the ordered microstructure of the semiconducting host in order to extend the delocalization of the  $\pi$ -electrons.<sup>35</sup> Different from boron, phosphorus (P) is theoretically predicted to occupy preferably the “interstitial sites” of  $g-C_3N_4$  between two adjacent tri-s-triazine units, creating a P–N bridged bond, which may act as a new carrier migration channel to promote charge mobility between the two adjacent heptazine units.<sup>36</sup> Moreover, it is also predicted that this interstitially P-doped structure can result in several potentially beneficial effects: (1) a downshift of the CB edge by  $\sim 0.25$  eV without altering the VB edge, leading to a narrowed band-gap of 2.3 eV, extending the absorption edge by about 80 nm, though this also reduces the driving force for proton reduction; (2) an isolated P 3p state localized  $\sim 0.15$  eV below the LUMO of the host  $g-C_3N_4$ , which may induces shallow trapping; and (3) delocalized/noncoplanar HOMO and LUMO distributions around the P and the adjacent N atoms, respectively, which facilitate charge separation and inhibit charge recombination. However, such potentially beneficial effects have not been fully realized in practice, to the best of our knowledge. In addition, another important fundamental issue is how protons are involved in photocatalysis (not just as a reactant), as  $H_2$  is usually produced by electron-involved proton reduction in an aqueous solution.

With full consideration of the above, we resorted to a heuristic strategy of tuning the band structure and charge carrier migration by both porosity engineering and phosphorus-interstitial doping of  $g-C_3N_4$  *via* a facile thermochemical treatment method. The resulting P-interstitially-doped carbon-defective holey  $g-C_3N_4$  (Pi-Ho@ $C_{3-x}N_4$ ) is an effective photocatalyst with enriched catalytic sites, substantially improved redox capacity and photo-dynamics for  $H_2$  generation. The increased surface area and pore distribution were demonstrated by the holey and ultrathin structure of  $g-C_3N_4$  with abundant carbon vacancies and catalytic

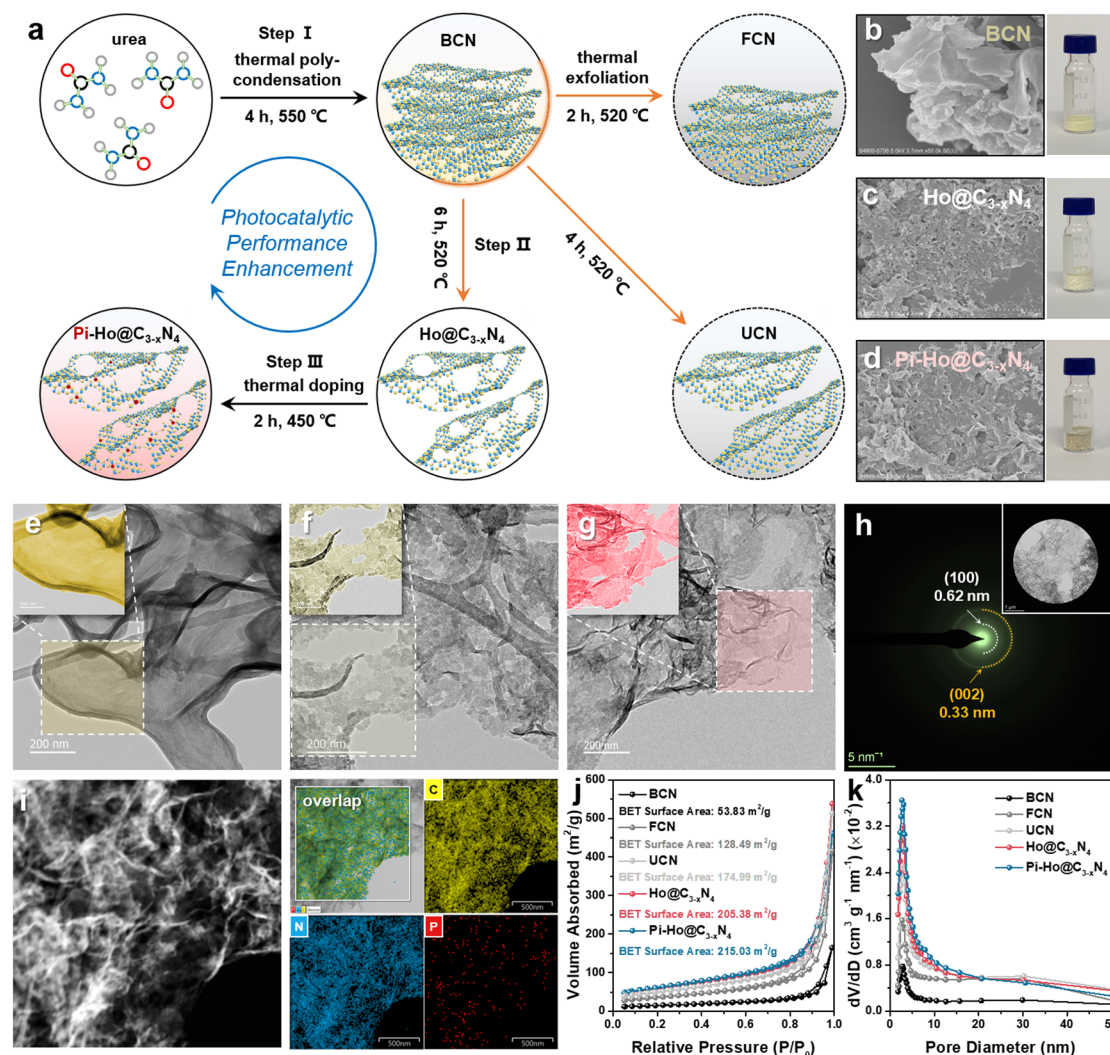


centres. Efficient charge separation was achieved in  $\text{Pi-Ho@C}_{3-x}\text{N}_4$ , compared with the bulk  $\text{g-C}_3\text{N}_4$ . Atomic phosphorus embedded in the interstitial sites of the adjacent heptazine units synergistically promotes the shallow trapping of the hot electron carriers and increases the conduction band potential to enhance both the kinetic and thermodynamic driving forces for  $\text{H}_2$  generation. Interestingly, we found that the *in situ* protonation plays a significant role in optimizing the band structure of the  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  catalyst and prolonging the photogenerated carrier lifetime after forming a proton-phosphorus coordinated ligand. Density Functional Theory (DFT) based simulations reveal the disappearance of a deep trap state and the emergence of a shallow trap state in a protonated phosphorus interstitial bridge. Femtosecond transient absorption spectroscopy (fs-TAS) *in situ* confirms the highly efficient charge carrier separation and accumulation of photocatalytically long-lived shallow trap states in water as well. As a

result, the photoactivity of  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  was significantly enhanced by over 800% ( $6323 \mu\text{mol h}^{-1} \text{g}^{-1}$ ), compared with that of bulk  $\text{g-C}_3\text{N}_4$  ( $790 \mu\text{mol h}^{-1} \text{g}^{-1}$ ). The design approach provides an effective pathway for the engineering of highly efficient and stable photocatalytic systems for visible-light-driven  $\text{H}_2$  generation.

## Results and discussion

Fig. 1a shows the facile thermochemical process and the resulting structures for bulk  $\text{g-C}_3\text{N}_4$  (BCN), holey  $\text{g-C}_3\text{N}_4$  ( $\text{Ho@C}_{3-x}\text{N}_4$ ) and  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  photocatalyst synthesis. Briefly, in step I, the BCN sample was synthesized by conventional thermal polycondensation of urea in a muffle furnace (Fig. S1a, see the ESI†). Then, the sample was thermochemically exfoliated at  $520^\circ\text{C}$  in an air atmosphere to prepare few-layer



**Fig. 1** (a) Schematic illustration of the synthesis process of  $\text{Pi-Ho@C}_{3-x}\text{N}_4$ ; SEM images of (b) BCN, (c)  $\text{Ho@C}_{3-x}\text{N}_4$  and (d)  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  and the corresponding digital photograph of each catalyst (20 mg); STEM images of (e) BCN, (f)  $\text{Ho@C}_{3-x}\text{N}_4$  and (g)  $\text{Pi-Ho@C}_{3-x}\text{N}_4$ ; (h) SAED pattern of  $\text{Pi-Ho@C}_{3-x}\text{N}_4$ ; (i) STEM EDX-mapping area for C, N and P elements distributed in the  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  sample; (j) Nitrogen adsorption-desorption isotherms; (k) Corresponding pore diameter distribution of BCN, FCN, UCN,  $\text{Ho@C}_{3-x}\text{N}_4$  and  $\text{Pi-Ho@C}_{3-x}\text{N}_4$ .





(FCN), ultrathin (UCN) and  $\text{Ho@C}_{3-x}\text{N}_4$  samples, respectively (step II). Finally, red phosphorus vapour doping was conducted under an inert (Ar) atmosphere (step III Fig. S1b, see the ESI†) to produce  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  catalysts. Both scanning electron microscopy (SEM) and scanning transmission electron microscopy (STEM) images revealed the morphology of the as-synthesized samples, as shown in Fig. 1b–i. Briefly, the SEM and STEM images of BCN in Fig. 1b and e show relatively thick 2D micro-sheets. The few-layer and the ultrathin structures of  $\text{g-C}_3\text{N}_4$  were obtained after thermochemical exfoliation for 2 and 4 hours, respectively (Fig. S2, see the ESI†). After the thermochemical treatment for 6 h, it is clearly seen that a holey structure was achieved with clean and transparent surfaces (Fig. 1c and f). Similarly, thus a holey feature persisted in the  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  sample (Fig. 1d and g). The ringed SAED pattern indicates the existence of an amorphous structure of the  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  sample, with only short-range orders; the main diffraction “rings” corresponding to the (002) and (100) crystal planes with d-spacing values of  $\sim 0.33$  and  $0.62$  nm, respectively (Fig. 1h).<sup>37,38</sup> X-ray maps and the high-resolution image show the uniform distribution of C, N and P elements and the micro-sheet structure of  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  (Fig. 1i). The linear scan shows strong C and N signals (Fig. S3, see the ESI†), and notably, a weak P signal, which reflects well the general chemical structure and the relatively low concentration of phosphorus. The photographic insets (Fig. 1b–d) show the expanded volumes of the non-exfoliated BCN, exfoliated  $\text{Ho@C}_{3-x}\text{N}_4$  and  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  samples at the same mass (20 mg), suggesting the successful exfoliation of  $\text{g-C}_3\text{N}_4$  from the bulk to the holey ultrathin structure. In addition, the colour of the resulting powders changes from light yellow for BCN, white for  $\text{Ho@C}_{3-x}\text{N}_4$  to light grey for  $\text{Pi-Ho@C}_{3-x}\text{N}_4$ , demonstrating that the phosphorus doping did affect the optical properties and the electronic structure of the photocatalyst.

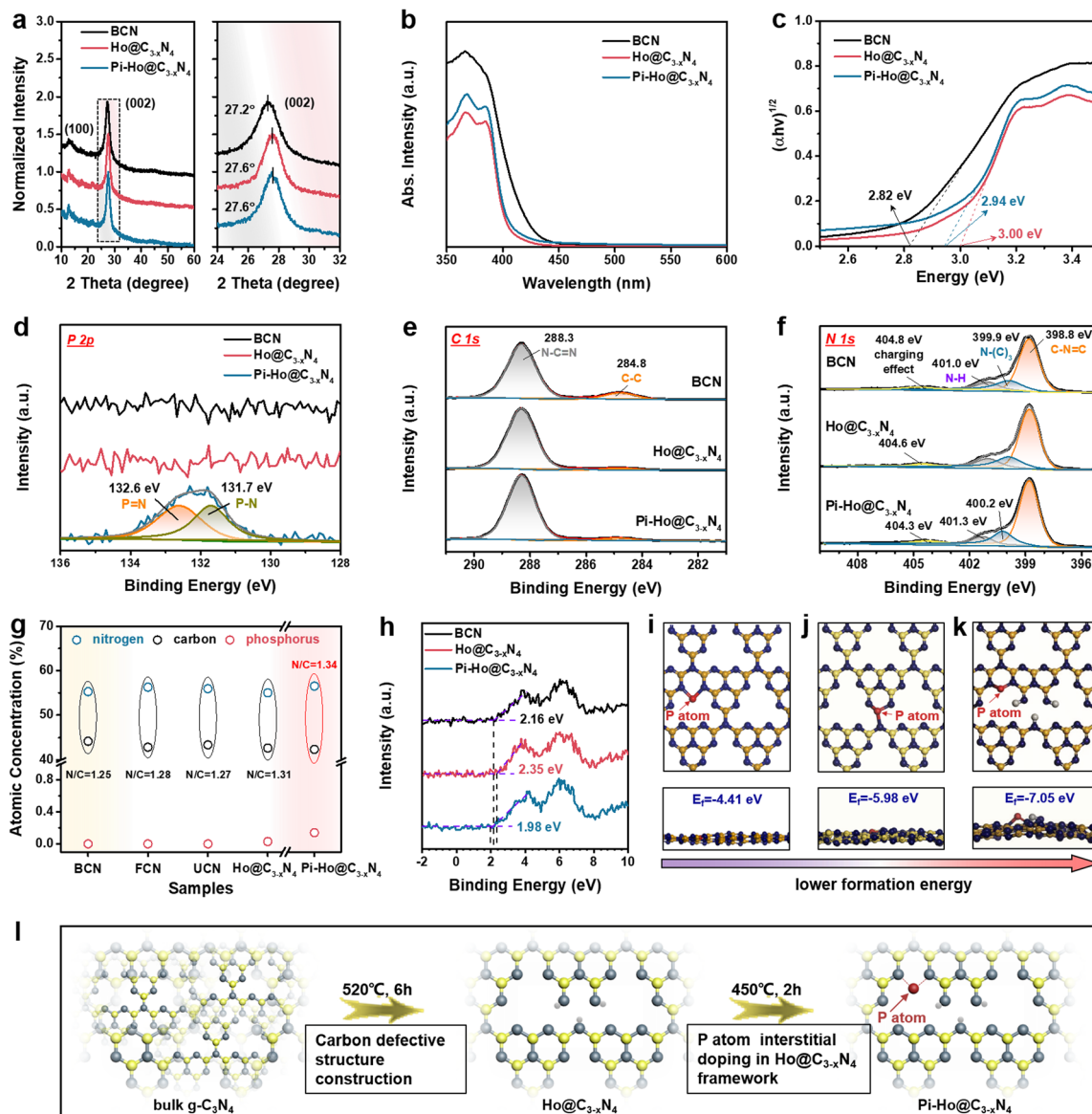
Nitrogen adsorption-desorption isotherms demonstrated an increased surface area in the order from  $\text{Pi-Ho@C}_{3-x}\text{N}_4 \approx \text{Ho@C}_{3-x}\text{N}_4 > \text{UCN} > \text{FCN} > \text{BCN}$  (Fig. 1j). The specific surface area of  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  ( $215.03 \text{ m}^2 \text{ g}^{-1}$ ) and  $\text{Ho@C}_{3-x}\text{N}_4$  ( $205.38 \text{ m}^2 \text{ g}^{-1}$ ) is the highest, compared with the ultrathin  $\text{g-C}_3\text{N}_4$  ( $174.99 \text{ m}^2 \text{ g}^{-1}$ ) and the few-layer  $\text{g-C}_3\text{N}_4$  ( $128.49 \text{ m}^2 \text{ g}^{-1}$ ), and is almost four times greater than that of the BCN sample ( $53.83 \text{ m}^2 \text{ g}^{-1}$ ). It is well known that the higher the surface area, the more the catalytic centres and the faster the mass transfer. Due to the abundant pores, both  $\text{Ho@C}_{3-x}\text{N}_4$  and  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  reveal a Type-IV isotherm of a highly mesoporous structure. Compared with the BCN, the pore diameters of the exfoliated samples show various levels of increase at  $\sim 2.6$  nm and towards  $\sim 5$  nm. In particular, both  $\text{Ho@C}_{3-x}\text{N}_4$  and  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  display far great levels of the mesopores by both the broadness and sharpness of the curves (Fig. 1k). X-ray diffraction (XRD) patterns of BCN,  $\text{Ho@C}_{3-x}\text{N}_4$  and  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  show two distinct peaks around  $27.2^\circ$  and  $12.8^\circ$ , which can be indexed to the  $\pi$ - $\pi$  interlayer stacking (002) and tri-s-triazine unit (100) planes, respectively (Fig. 2a).<sup>39,40</sup> This suggests that the conjugated aromatic groups maintain periodic stacking. Importantly, a slight right-shift ( $\sim 0.5^\circ$ ) of the (002) diffraction was observed

for  $\text{Ho@C}_{3-x}\text{N}_4$  and  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  samples, indicating a reduced interlayer distance.<sup>41</sup> UV-Vis diffuse reflectance spectroscopy (UV-Vis DRS) revealed the direct influence of the structural disorder on photon absorption, as shown in Fig. 2b. Compared to BCN,  $\text{Ho@C}_{3-x}\text{N}_4$  shows an evident blue shift of the absorption edge toward a smaller wavelength, which was ascribed to the following two reasons: (i) a quantum confinement effect and (ii) a conjugation unit reduction caused by the in-plane holes.<sup>42</sup> In contrast,  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  shows stronger visible-light absorption than that of  $\text{Ho@C}_{3-x}\text{N}_4$  due to P doping. The bandgap energy ( $E_g$ ) was calculated to be 2.82 eV for BCN, 3.00 eV for  $\text{Ho@C}_{3-x}\text{N}_4$  and 2.94 eV for  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  using Tauc plots (Fig. 2c).

X-ray photoelectron spectroscopy (XPS) was employed to investigate the surface chemical structures of BCN,  $\text{Ho@C}_{3-x}\text{N}_4$  and  $\text{Pi-Ho@C}_{3-x}\text{N}_4$ , as shown in Fig. 2d–f. XPS wide spectra showed that the N and C are the main components of the BCN,  $\text{Ho@C}_{3-x}\text{N}_4$  and  $\text{Pi-Ho@C}_{3-x}\text{N}_4$ , whereas no clear phosphorus signal can be found in the  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  framework due to its relatively low concentration (Fig. S4, see the ESI†). However, the high resolution and curve fitting of P 2p spectra show two P-relevant binding energies around 131.7 and 132.6 eV in  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  (Fig. 2d), corresponding to the P–N and P=N coordination, respectively. Furthermore, the C 1s and N 1s XPS spectra (Fig. 2e and f) show similar characteristic peaks of  $\text{g-C}_3\text{N}_4$ . As shown in Fig. 2e, the two peaks centered at 284.8 and 288.3 eV of the C 1s can be assigned to graphitic carbon (C–C coordination) and  $\text{sp}^2$ -hybridized carbon (N–C=N), respectively.<sup>43</sup> The N 1s spectrum can be discretely divided into four peaks, respectively, attributed to the  $\text{sp}^2$ -bonded nitrogen in C–N=C (398.8 eV), the tertiary N atoms (399.9 eV), the terminal N–H group (401.0 eV) and the positive charge localization in the heterocycles (404.8 eV) (Fig. 2f).<sup>44</sup> Notably, the slight shift of the N–(C)<sub>3</sub>, N–H and the positive charge localization are likely due to the phosphorus-doping induced charge redistribution in the  $\text{g-C}_3\text{N}_4$  matrix. Conventionally, a P dopant prefers to substitute the bay-carbon sites in a triazine unit, forming three P–N bonds.<sup>45–47</sup> However, for P interstitial doping, the P atom would only bond to two N atoms. Because of the higher electronegativity of the N atom than that of the P atom, fewer P–N bonds imply higher electron density remaining at the P atom,<sup>48</sup> or a lower P–N binding energy, as indeed observed for the  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  (131.7 and 132.6 eV for the P–N and P=N bonds, respectively, Fig. 2d), compared to its conventional phosphorus substitutionally doped counterpart (131.9 and 133.0 eV for the P–N and P=N bonds, Fig. S5, see the ESI†). The observation indicates that P-doping is mainly at the interstitial site in  $\text{Pi-Ho@C}_{3-x}\text{N}_4$ .

Fourier transformation infrared (FTIR) spectroscopy was used to investigate further the chemical structure and phosphorus-interstitial-doping of the as-prepared catalysts (Fig. S6, see the ESI†). The characteristic vibration bands at  $1200$ – $1700 \text{ cm}^{-1}$  are from the stretching mode of the C–N vibrations (Fig. S6a, see the ESI†).<sup>49</sup> Furthermore, the sharp peak centred at  $800 \text{ cm}^{-1}$  corresponds to the intrinsic heptazine unit of the carbon nitride structure.<sup>50</sup> In addition, the





**Fig. 2** (a) XRD patterns; (b) UV-Vis DRS; (c) Tauc plots derived from (b); (d) high resolution XPS spectra of P 2p; (e) high resolution XPS spectra of C 1s. (f) High resolution XPS spectra of N 1s. (g) Atomic concentration of C, N and P in BCN, FCN, UCN and Ho@C<sub>3-x</sub>N<sub>4</sub> and Pi-Ho@C<sub>3-x</sub>N<sub>4</sub>, respectively. (h) VB XPS spectra of BCN, Ho@C<sub>3-x</sub>N<sub>4</sub> and Pi-Ho@C<sub>3-x</sub>N<sub>4</sub>; (i–k) atomic schemes (top and side views) of Pi-Ho@C<sub>3-x</sub>N<sub>4</sub> with phosphorus atoms at the most energetically favourable sites. (l) Proposed atomic structure evolution.

broad absorption bands at 3000–3400 cm<sup>−1</sup> are caused by the absorbed H<sub>2</sub>O molecules and the uncondensed amino group.<sup>51</sup> The vibration peak is around 667 cm<sup>−1</sup>, which is close to the range observed for g-C<sub>3</sub>N<sub>4</sub> in the literature,<sup>47</sup> and hence attributable to the stretching mode of the P–N bond. The slight shift (to a higher frequency) indeed reflects the different (stronger) bonding environment for the P-doping at the interstitial site (Fig. S6b, see the ESI†). The neighbouring signal, ~646 cm<sup>−1</sup>, comes from the heptazine structure, which exists in all the samples. Importantly, both P–N and P=N coordination was found from the high resolution XPS spectra for P 2p in Pi-Ho@C<sub>3-x</sub>N<sub>4</sub> (Fig. 2d), which is a valid indication that the phosphorus atoms may be located at the interstitial site between two adjacent tri-s-triazine units, forming a new “–N–P=N–” bridging channel. To further

clarify the phosphorus substitution site in g-C<sub>3</sub>N<sub>4</sub>, both the interstitial and the C-vacancy sites were evaluated for P doping, by the Vienna Ab initio Simulation Package (VASP). The corresponding atomic schemes (top and side views) of the P-doped Ho@C<sub>3-x</sub>N<sub>4</sub> with the most energetic configurations are shown in Fig. 2i–k. It is noted that the formation energy  $E_f$  of the interstitial P-doped g-C<sub>3</sub>N<sub>4</sub> ( $E_f = -7.05$  eV) is the most negative among that of the conventional P-doped g-C<sub>3</sub>N<sub>4</sub> (via replacing carbon atom,  $E_f = -5.98$  eV) and the interstitial P-doped without a carbon vacancy ( $E_f = -4.41$  eV), respectively.<sup>52</sup> Hence, the phosphorus atoms are preferentially located at the interstitial site with a new “–N–P=N–” connection.

Further elemental analysis of the XPS results revealed deep insight into the existence of carbon vacancies, as evident from





the atomic ratio of N/C (Fig. 2g). The N/C ratio increases from 1.25 for BCN, 1.28 for FCN, 1.27 for UCN to 1.31 for  $\text{Ho@C}_{3-x}\text{N}_4$ , which can be ascribed to the removal of intrinsic carbon impurities and creation of the carbon vacancies during thermochemical exfoliation. Therefore, Fig. 2e shows negligible C–C coordination in  $\text{Ho@C}_{3-x}\text{N}_4$  and  $\text{Pi-Ho@C}_{3-x}\text{N}_4$ , respectively, compared to the BCN. Moreover, both  $\text{Ho@C}_{3-x}\text{N}_4$  and  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  show a higher N/C ratio, due to the enriched carbon vacancy structure. As illustrated in Fig. 2l, the  $\text{Ho@C}_{3-x}\text{N}_4$  sample was obtained under 520 °C, followed by the phosphorus-interstitial-doping in the  $\text{g-C}_3\text{N}_4$  structure. This interstitial doped phosphorus can provide a new “–N–P–N–” chemical channel on the  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  sheets to boost charge separation. To investigate the band positions of the as-synthesized catalysts, valence band (VB) XPS spectra were used to estimate VB positions (Fig. 2h). Compared with bulk  $\text{g-C}_3\text{N}_4$  (BCN,  $\sim 2.16$  eV),  $\text{Ho@C}_{3-x}\text{N}_4$  displayed the most positive potential edge at  $\sim 2.35$  eV. For  $\text{Pi-Ho@C}_{3-x}\text{N}_4$ , the phosphorus interstitial doping induced a VB shift to 1.98 eV.

Fig. 3a shows the visible-light-driven ( $\lambda > 400$  nm)  $\text{H}_2$  production characteristics of the as-prepared catalysts with

TEOA (10 vol%) as the hole scavenger and Pt (co-catalyst, 3 wt%) as the electron acceptor. Pt nanoclusters were anchored on the  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  surface by the photoreduction method as shown in Fig. S7 (see the ESI<sup>†</sup>), which shows the uniform distribution of Pt nanoclusters and the nanosheet structure of  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  by STEM EDX-mapping. As expected,  $\text{Ho@C}_{3-x}\text{N}_4$  presented the highest photocatalytic hydrogen generation activity ( $2746 \mu\text{mol h}^{-1} \text{g}^{-1}$ ) than those of BCN ( $790 \mu\text{mol h}^{-1} \text{g}^{-1}$ ), FCN ( $1014 \mu\text{mol h}^{-1} \text{g}^{-1}$ ) and UCN ( $2187 \mu\text{mol h}^{-1} \text{g}^{-1}$ ), in correspondence with the improvements from the increased surface area, catalytic site density and high redox capability. By optimization of the P-doping level, the photocatalytic  $\text{H}_2$  production rate reached a maximum of  $6323 \mu\text{mol h}^{-1} \text{g}^{-1}$  over the  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  photocatalyst (Fig. 3a), which is up to 8 times greater than the BCN sample. As shown in Fig. S8 (see the ESI<sup>†</sup>), EDX measurement reflected the actual P-doping content in the  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  samples. However, an excess amount of P dopants may form potential recombination centers for electron–hole pairs.<sup>53</sup> Hence, the  $\text{H}_2$  production rate decreases to  $4249 \mu\text{mol h}^{-1} \text{g}^{-1}$  when the

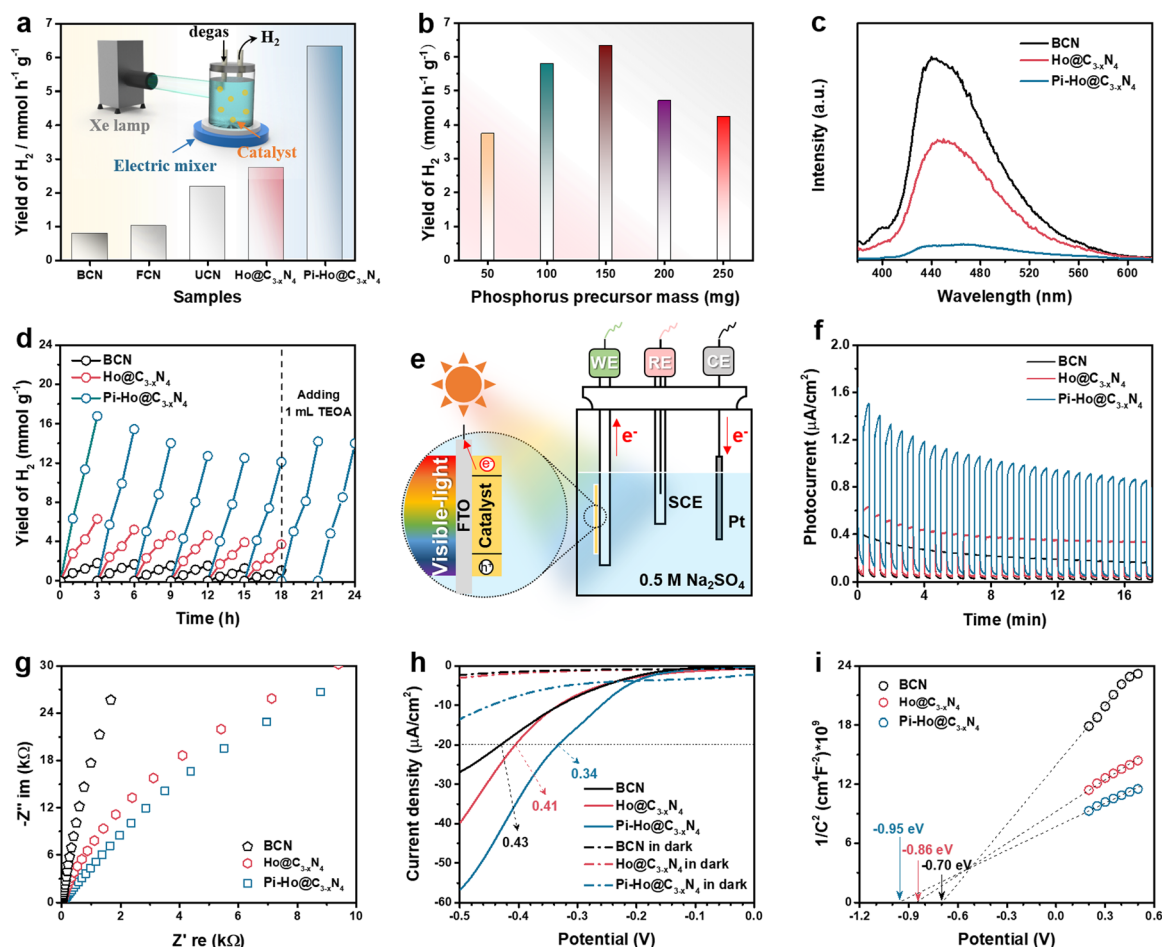


Fig. 3 (a) Comparison of the  $\text{H}_2$  evolution rates of different catalysts and inset: Schematic of the photocatalytic measurement system; (b)  $\text{H}_2$  evolution rate as a function of red P precursor mass (50–250 mg); (c) steady-state photoluminescence (PL) emission under 355 nm laser flash; (d) cyclic runs of photocatalytic  $\text{H}_2$  generation over the BCN,  $\text{Ho@C}_{3-x}\text{N}_4$  and  $\text{Pi-Ho@C}_{3-x}\text{N}_4$ , respectively; (e) schematic of a photoelectrochemical cell (working electrode: photocatalyst/FTO; counter electrode: Pt foil; reference electrode: SCE), and corresponding photoelectrochemical measurements: (f) photocurrent, (g) Nyquist plots, (h) LSV curves and (i) Mott–Schottky plots.



phosphorus precursor was increased to 250 mg (Fig. 3b). The control experiment of the thermal doping temperature was evaluated in Fig. S9 (see the ESI†), showing that the optimal temperature is 450 °C. The photocatalytic H<sub>2</sub> production characteristics under simulated sunlight (AM 1.5) and visible-light irradiation ( $\lambda > 400$  nm) are compared in Fig. S10 (see the ESI†). The histograms show that the H<sub>2</sub> production rate reduces in the order  $\text{Pi-Ho@C}_{3-x}\text{N}_4 > \text{Ho@C}_{3-x}\text{N}_4 > \text{BCN}$ .

On the other hand, the photocatalytic activity is also strongly influenced by charge carrier kinetics and separation. Fig. 3c shows the steady-state photoluminescence (PL) emission spectra of the BCN,  $\text{Ho@C}_{3-x}\text{N}_4$  and  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  in pure water for the investigation of electron/hole recombination. The radiative peaks around 450 and 400 nm are related to the lowest singlet exciton and excimer emissions, respectively.<sup>54</sup> Notably, the dramatic PL quenching of the  $\text{Pi-Ho@C}_{3-x}\text{N}_4$ , relative to BCN and  $\text{Ho@C}_{3-x}\text{N}_4$ , implies effective survival of the excited states with a high population of active species, which can only be attributed to efficient charge transport boosted by the “-N-P=N-” channel. The quantum efficiency (QE) at 420 nm of BCN,  $\text{Ho@C}_{3-x}\text{N}_4$  and  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  was estimated to be 0.87, 2.64 and 5.08%, respectively. The internal quantum efficiency (IQE) of  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  photocatalyst (3 wt% Pt as a co-catalyst) reaches 31.8% at 400 nm, compared to 5.7% for the BCN under the same conditions. In addition, the photostability of the as-prepared catalysts was also investigated, as noted in Fig. 3d. The photo-activity was largely recovered, from ~72% to 83%, after addition of only 1.0 mL of fresh electron donor (TEOA) into the system at the end of the sixth cycle. This indicates that much of the reduction in H<sub>2</sub> generation is due to the exhaustion of the electron donor (TEOA), rather than the degradation of the catalyst. The almost unchanged FTIR spectra (Fig. S11, see the ESI†) after the 18 h reaction further confirm the high stability of the  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  structure for visible-light-driven H<sub>2</sub> production. Morphological studies from SEM and high-resolution STEM (HRSTEM) demonstrate the holey structure and the transparent surfaces after the photostability test (Fig. S12a–c, see the ESI†). In addition, Pt nanoclusters, as co-catalyst, were deposited on the  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  under visible-light irradiation as shown in Fig. S12b (see the ESI†). The HRSTEM image in Fig. S12c (see the ESI†) demonstrates that these 5–10 nm highly crystallized Pt nanoclusters prefer to grow along the (111) crystal planes.<sup>55</sup> Indeed the postmortem structural analysis of the catalyst by SEM and STEM (Fig. S12, ESI†) shows little change in the structural integrity, which further supports the long-term operational potential of the catalyst. A comparison of the H<sub>2</sub> production rate with the reported conventional phosphorus substitutional doped and other heteroatom doped g-C<sub>3</sub>N<sub>4</sub> catalysts (mostly using Pt as a co-catalyst) in the literature is shown in Table S1 and S2 (see the ESI†), respectively. It is clear that  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  shows much more effective H<sub>2</sub> production than that of the other (g-C<sub>3</sub>N<sub>4</sub>)-based photocatalysts under similar experimental conditions.

To understand the intrinsic charge transport of the as-synthesized photocatalysts, photoelectrochemical properties were characterized by the photocurrent, electrochemical

impedance spectroscopy (EIS), linear sweep voltammetry (LSV) and Tafel plot, using a 0.5 M Na<sub>2</sub>SO<sub>4</sub> aqueous solution without a hole scavenger or a co-catalyst (Fig. 3f–i). Fig. 3e shows a schematic diagram of the conventional three-electrode system in a photoelectrochemical cell. According to past reports, the highest photocurrent manifests the most efficient electron–hole separation.<sup>56</sup> The long-term transient current (*i*–*t*) curves were measured to show the photo-response and stability in Fig. 3f. Due to the interaction equilibrium required between the electrolyte and the electrode in the early stage, all samples show an apparent decline of photocurrent density (PCD), followed by good stability after about 12 min. As expected, the strongest PCD value was observed at the  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  electrode, suggesting the most efficient electron–hole separation under visible-light irradiation. Tafel, EIS and polarization curves were employed to investigate the charge transfer efficiency and kinetics. Typically, for the Nyquist plots, the largest diameter reveals the highest charge transportation resistance.<sup>57</sup> Fig. 3g shows clearly that the diameter of the Nyquist plot is in the following order:  $\text{Pi-Ho@C}_{3-x}\text{N}_4 < \text{Ho@C}_{3-x}\text{N}_4 < \text{BCN}$ , suggesting the lowest resistance of the  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  for the highest carrier transportation efficiency. Moreover, the polarization profiles of the catalysts were obtained from LSV under dark and visible light irradiations, respectively (Fig. 3h). The different photocurrent densities illustrate the enhanced visible-light-driven photocatalytic kinetics for water splitting in comparison with that under dark conditions. The overpotential (at  $-20 \mu\text{A cm}^{-2}$ ) of the H<sub>2</sub> production of the  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  (0.34 V) is much lower than that of both the BCN (0.43 V) and the  $\text{Ho@C}_{3-x}\text{N}_4$  (0.41 V) under visible-light exposure, demonstrating the important roles of the phosphorus-interstitial-doping and the defective holey carbon structure for radically improving photocatalytic activity and kinetics.

To gain further insight into the relationship between the carrier density ( $N_D$ ) and the flat band potential ( $E_{\text{FB}}$ ), Mott–Schottky plots were obtained by fitting the *x*-axis intercept of the linear section of  $1/C^2$  vs. potential, as shown in Fig. 3i. The  $E_{\text{FB}}$  values of the BCN,  $\text{Ho@C}_{3-x}\text{N}_4$  and  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  were determined to be  $-0.70$ ,  $-0.86$ , and  $-0.95$  V vs. SCE (saturated calomel electrode), corresponding to the  $-0.76$ ,  $-0.92$  and  $-1.01$  eV of the conduction band (CB) positions, respectively (Fig. 4a). Therefore, owing to the most negative CB potential, the  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  photocatalyst possesses the most efficient driving force for photocatalytic proton reduction among the catalysts under consideration. Moreover, the exchange current density ( $J_0$ ) was investigated by Tafel-polarization (Fig. S13, see the ESI†).  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  possesses the highest  $J_0$  at the cathodic branch, favouring the solar-driven reduction reaction and the carrier kinetics. Taking the bandgap and the UV-vis DRS into consideration, the remarkable difference in the H<sub>2</sub> production rate may be attributed to the improvements of the band-edge, charge carrier separation (*via* phosphorus-interstitial-doping) and active site density (*via* defective porosity engineering). In addition, the valence band (VB) potentials of the BCN (2.06 eV),  $\text{Ho@C}_{3-x}\text{N}_4$  (2.08 eV) and  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  (1.93 eV) were further





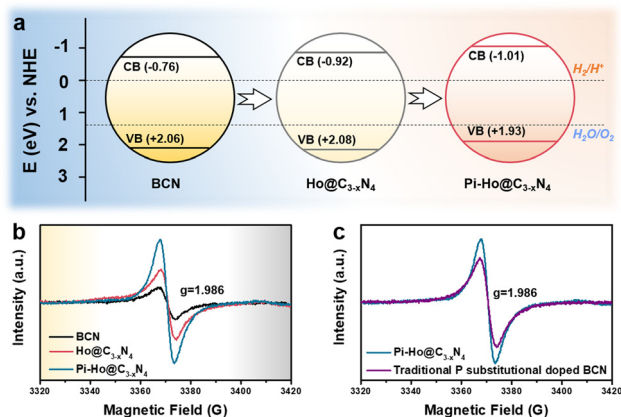


Fig. 4 (a) Energy level structure diagram of the BCN, Ho@C<sub>3-x</sub>N<sub>4</sub> and Pi-Ho@C<sub>3-x</sub>N<sub>4</sub>, respectively; (b and c) EPR spectra of BCN, Ho@C<sub>3-x</sub>N<sub>4</sub>, Pi-Ho@C<sub>3-x</sub>N<sub>4</sub>; and P substitutionally doped BCN samples at low-temperature.

estimated from CB and  $E_g$  (Fig. 4a), which are coincident with the VB XPS spectra results. Moreover, the capacitance and carrier density were calculated using the Mott-Schottky equation:<sup>58,59</sup>

$$\frac{1}{C^2} = \left( \frac{2}{N_D e_0 \epsilon_0 \epsilon} \right) \left[ (E - E_{FB}) - \frac{kT}{e_0} \right]$$

$$N_D = \frac{2}{e_0 \epsilon_0 \epsilon} \left( \frac{dE}{d\left(\frac{1}{C^2}\right)} \right)$$

where  $C$ ,  $N_D$ ,  $E$  and  $E_{FB}$  refer to the charge capacitance, carrier density, applied potential and flat band potential, respectively;  $e_0$ ,  $\epsilon$  and  $\epsilon_0$  correspond to the charge constant ( $1.60 \times 10^{-19}$  C), dielectric constant (7–8 for g-C<sub>3</sub>N<sub>4</sub>) and permittivity of a vacuum ( $8.85 \times 10^{-12}$  F m<sup>-1</sup>);  $\kappa$  is the Boltzmann constant ( $1.38 \times 10^{-23}$  J K<sup>-1</sup>); and  $T$  is the temperature.<sup>60</sup> The  $N_D$  value of Pi-Ho@C<sub>3-x</sub>N<sub>4</sub> ( $3.1 \times 10^{19}$  cm<sup>-3</sup>) was the highest among those considered here, BCN ( $9.2 \times 10^{18}$  cm<sup>-3</sup>) and Ho@C<sub>3-x</sub>N<sub>4</sub> ( $2.3 \times 10^{19}$  cm<sup>-3</sup>), which indicates the efficient separation of the charge carriers to produce active electrons for H<sub>2</sub> generation.

The paramagnetic properties of the BCN, Ho@C<sub>3-x</sub>N<sub>4</sub> and Pi-Ho@C<sub>3-x</sub>N<sub>4</sub> were investigated by electron paramagnetic resonance (EPR) at low temperature (about 40 K, Fig. 4b). A noticeably enhanced intensity of the EPR spectra of the Pi-Ho@C<sub>3-x</sub>N<sub>4</sub> was observed, compared to that of the BCN and the Ho@C<sub>3-x</sub>N<sub>4</sub>, indicating that an increased concentration of the unpaired electrons and/or the delocalized electrons in the  $\pi$ -conjugated aromatic rings of the Pi-Ho@C<sub>3-x</sub>N<sub>4</sub>.<sup>61,62</sup> The effect can be attributed to the phosphorus interstitial doping, which leads to more unpaired electrons in the adjacent nitrogen atom.<sup>63</sup> In addition, the much higher intensity of the Ho@C<sub>3-x</sub>N<sub>4</sub> than that of the BCN (Fig. 4b), implies more abundant defective sites in the Ho@C<sub>3-x</sub>N<sub>4</sub>.<sup>64</sup> As depicted in Fig. 4c, the P substitutionally doped BCN sample shows a lower intensity than the Pi-Ho@C<sub>3-x</sub>N<sub>4</sub>. Hence, we propose that the P interstitial doping prefers to form an electron-enriched

structure. Therefore, the synergetic effect of the porosity and the interstitial doping can effectively optimize electron delocalization, to enhance photocatalytic performance.

Herein, we carried out density functional theory (DFT) calculations to clarify the shallow and deep trap states change of g-C<sub>3</sub>N<sub>4</sub> derived photocatalysts. The details of the calculation procedures are given in the ESI†. The density of states (DOS) of a g-C<sub>3</sub>N<sub>4</sub> monolayer (denoted as pristine g-C<sub>3</sub>N<sub>4</sub>) was shown in Fig. 5a<sub>1</sub>, in comparison with those of the Ho@C<sub>3-x</sub>N<sub>4</sub> (Fig. 5b<sub>1</sub>), and Pi-Ho@C<sub>3-x</sub>N<sub>4</sub> (Fig. S14a, see the ESI†). Note that the estimated band gaps ( $\sim 1.1$  eV for the pristine g-C<sub>3</sub>N<sub>4</sub>) were much smaller than the experimental value ( $\sim 2.8$  eV), due to the well-known limitation of the band gap by the generalized gradient approximation method.<sup>65</sup> Such underestimation can be ascribed to the delocalization error of the approximation functional.<sup>66</sup> On the other hand, the relative mechanistic effects of the vacancy and the phosphorus-interstitial-doping remain valid here. Notably, C-vacancy induced an apparent upshift of CB with a higher driving force (Fig. 5b<sub>1</sub>). The CB and VB of both pristine g-C<sub>3</sub>N<sub>4</sub> and Ho@C<sub>3-x</sub>N<sub>4</sub> are mainly composed of C 2p and N 2p orbitals, respectively (Fig. 5a<sub>2</sub> and b<sub>2</sub>). The photogenerated electrons will be transferred from N 2p to C 2p, followed by the re-transfer from the C atom to the N atom for the photocatalytic reaction.<sup>67</sup> This is also the reason for the high recombination rate in the pristine g-C<sub>3</sub>N<sub>4</sub>. Similarly, the CB of the Pi-Ho@C<sub>3-x</sub>N<sub>4</sub> is mainly composed of C 2p orbital (Fig. S14a, see the ESI†). However, rather differently, the Pi-Ho@C<sub>3-x</sub>N<sub>4</sub> shows two trap states (shallow and deep states, respectively) close to the CB after the phosphorus-interstitial-doping. A comparison of the DOSs of Pi-Ho@C<sub>3-x</sub>N<sub>4</sub> and H-Pi-Ho@C<sub>3-x</sub>N<sub>4</sub> (proton-terminated Pi-Ho@C<sub>3-x</sub>N<sub>4</sub> at the P site), as protons widely exist in aqueous solution, demonstrates two significant differences: (1) disappearance of the deep state after constructing a proton-Pi-Ho@C<sub>3-x</sub>N<sub>4</sub> interaction bond with a lower formation energy (Fig. S15, see the ESI†), reflecting that the natural proton assists shallow trap state survival as the deep state is photocatalytically inactive for H<sub>2</sub> production in water (Fig. 5c<sub>1</sub>); (2) this shallow state is located at the Fermi level, and thus should benefit the electronic conductivity (Fig. 5c<sub>1</sub>).<sup>68</sup> By *in situ* construction of the proton-terminated Pi-Ho@C<sub>3-x</sub>N<sub>4</sub>, the model shows the disappearance of the deep trap state, which is below the Fermi level (Fig. 5c<sub>1</sub>).

To obtain further insight into the  $S_0 \rightarrow S_1$  transition of the g-C<sub>3</sub>N<sub>4</sub> derived photocatalysts, further theoretical calculations were conducted on molecular models that include the pristine tri-s-triazine unit, one with a carbon vacancy and another with both the carbon vacancy and the interstitial-doping of phosphorus (Fig. 5d<sub>1</sub>–d<sub>3</sub>). The green and blue meshes represent the charge density difference (CDD), respectively representing the “accumulation” electrons and the holes. The CDD of the pristine g-C<sub>3</sub>N<sub>4</sub> indicates that the excited electrons and holes are distributed at the C and N atoms during the  $S_0 \rightarrow S_1$  transition (Fig. 5d<sub>1</sub>). The different distributions of the excited electrons are reflected by the electron-hole pair distance ( $\Delta D$ ) values.<sup>69</sup> The shorter  $\Delta D$  of pristine g-C<sub>3</sub>N<sub>4</sub> (0.079 Å) than that



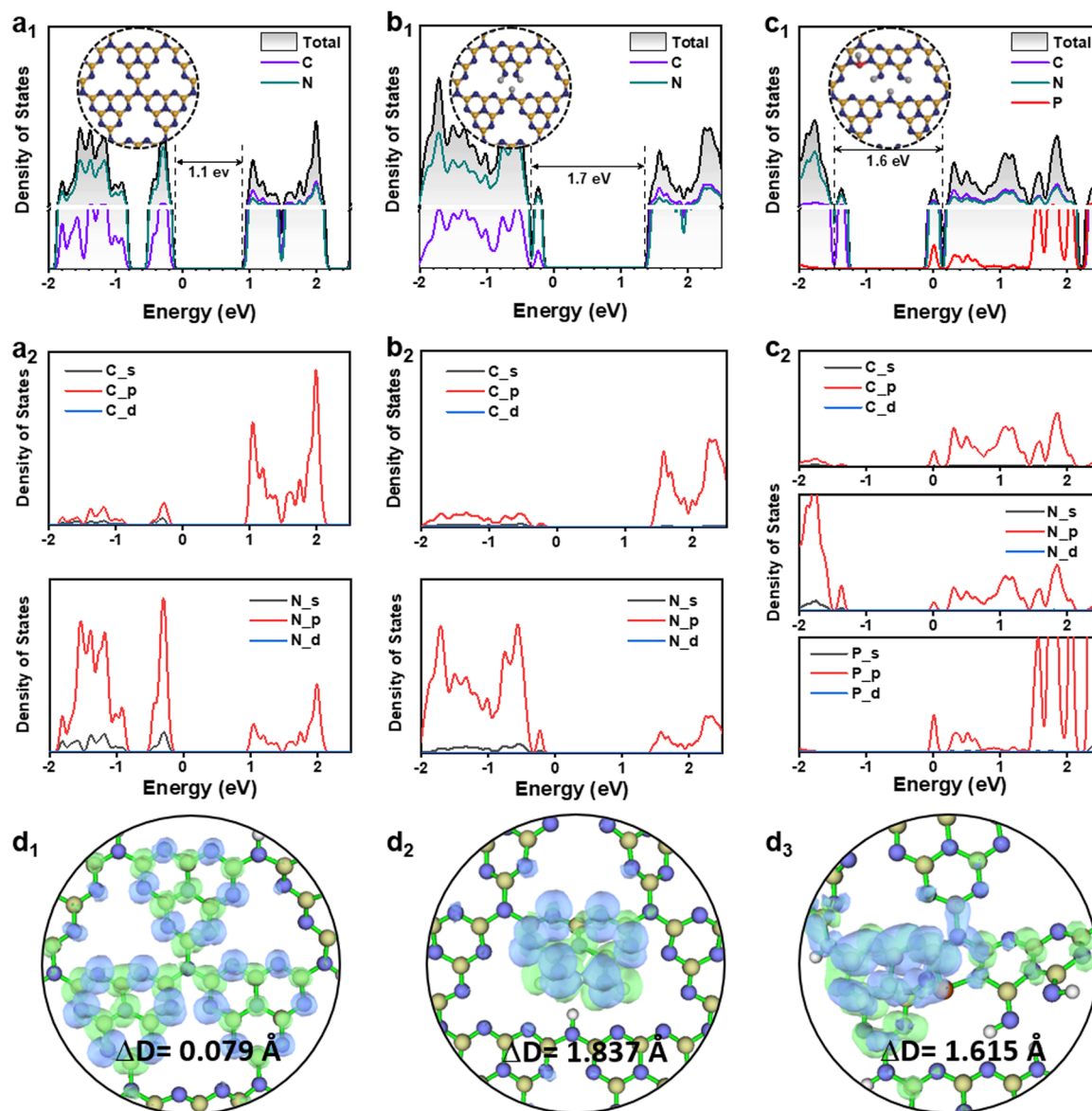


Fig. 5 Calculated density of states (DOS) of (a<sub>1</sub>–a<sub>2</sub>) pristine g-C<sub>3</sub>N<sub>4</sub>, (b<sub>1</sub>–b<sub>2</sub>) Ho@C<sub>3–x</sub>N<sub>4</sub>, (c<sub>1</sub>–c<sub>2</sub>) *in situ* proton-terminated Pi-Ho@C<sub>3–x</sub>N<sub>4</sub> at the P site; charge density difference of (d<sub>1</sub>) pristine g-C<sub>3</sub>N<sub>4</sub>, (d<sub>2</sub>) Ho@C<sub>3–x</sub>N<sub>4</sub>, (d<sub>3</sub>) *in situ* proton-terminated Pi-Ho@C<sub>3–x</sub>N<sub>4</sub> during the S<sub>0</sub>–S<sub>1</sub> transition, calculated using the B3LYP functional with the 6-31G basis set; the green and blue meshes represent the accumulation of electrons and the holes during the S<sub>0</sub>–S<sub>1</sub> transition, respectively.

of Ho@C<sub>3–x</sub>N<sub>4</sub> (1.837 Å), suggesting the less charge transfer of pristine g-C<sub>3</sub>N<sub>4</sub>. As expected, Ho@C<sub>3–x</sub>N<sub>4</sub> shows the separated electrons and holes near the heptazine unit, as shown in Fig. 5d<sub>2</sub>. Strong charge transport is noted in the phosphorus interstitially doped carbon-defective g-C<sub>3</sub>N<sub>4</sub>. As a result, the excited charges are distributed around the interstitially doped phosphorus with a longer  $\Delta D$  (2.736 Å, Fig. S14b, see the ESI†). However, proton-terminated Pi-Ho@C<sub>3–x</sub>N<sub>4</sub> shows a shorter  $\Delta D$  value (1.615 Å, Fig. 5d<sub>3</sub>), corresponding to the P–H bond localising in the excited charge carriers for the H<sub>2</sub> production reaction.

To further investigate the effect of *in situ* protonation, PL tests were carried out in an inert medium (acetonitrile), and the PL spectra are shown in Fig. S16 (see the ESI†). Compared with

the BCN and the Ho-C<sub>3–x</sub>N<sub>4</sub> sample, Pi-Ho@C<sub>3–x</sub>N<sub>4</sub> catalyst shows the weakest emission signal in acetonitrile, corresponding to the most efficient electron–hole separation. However, this emission process was significantly suppressed in water (Fig. 3c) due to the *in situ* protonation and results in the enhanced charge trapping process. To clarify further the influence of proton concentration in the photocatalytic system, the photocatalytic H<sub>2</sub> production reaction and PL spectra were obtained for different pH solutions, respectively (Fig. S17, see the ESI†). Compared to photoactivity in water, the maximum photocatalytic H<sub>2</sub> production rate reaches 9705 μmol h<sup>–1</sup> g<sup>–1</sup> in a buffer solution (pH = 5). However, excessive protons in the solution led to a reduction in catalytic performance. This pH-dependent photocatalysis implies that the Pi-Ho@C<sub>3–x</sub>N<sub>4</sub>





shows saturated protonation at a low pH level. PL spectra in Fig. S17b (see the ESI†) demonstrates the insignificant effect of pH values for electron–hole recombination. Femtosecond transient absorption spectroscopy (fs-TAS, excited by 400 nm laser flash pulses of light, from 200 fs to 2.5 ns) as straightforward evidence was carried out for a further understanding of the intrinsic photocatalytic dynamics and charge transfer routes in water without the use of acceptor–donor agents (*e.g.*, TEOA and Pt). Fig. S18a (see the ESI†) provides detailed information on the fs-TAS setup. As observed in Fig. 6a, the time-dependent absorption of the BCN sample shows a broad negative absorption feature from 450 to 650 nm. Interestingly, both the  $\text{Ho@C}_{3-x}\text{N}_4$  and  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  catalysts displayed two polarized characteristic absorption areas at (i) 450–550 nm (negative absorption) and (ii) 550–700 nm (positive absorption), respectively, as shown in Fig. 6b and c. As previously reported by Schlenker and Luo,<sup>70,71</sup> this negative feature is due to the ground state bleaching or stimulated emission (SE) and the positive signal is assigned to excited-state absorption (ESA). The UV-Vis DRS and steady-state PL clearly show a weak absorption and a strong emission peak after 450 nm, respectively. Therefore, the negative absorption feature can be assigned to SE, which promotes the recombination of electron–hole pairs and hence reduces the photocatalytic efficiency. The ESA feature in

the visible-light wavelength region was primarily caused by photogenerated electrons or holes or electron–hole pairs in the  $\text{g-C}_3\text{N}_4$ .<sup>72</sup> Fig. 6d–f show selected fs-TAS profiles of BCN,  $\text{Ho@C}_{3-x}\text{N}_4$  and  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  from 1.0 to 2000 ps. Compared with the BCN, the  $\text{Ho@C}_{3-x}\text{N}_4$  sample displays a marked enhancement in the ESA under the fs-TAS, corresponding to the C-vacancy induced charge separation. Note that  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  shows the weakest SE and the strongest ESA among those BCN and  $\text{Ho@C}_{3-x}\text{N}_4$  samples, suggesting the significant role of phosphorus-interstitial-doping towards the highly efficient and long-lived charge separation.

Representative kinetic decay traces for the ESA (630 nm) and SE (460 nm) are shown in Fig. 6g and h and the relevant fitted parameters are listed in Table S3 (see the ESI†). Concerned about the existence of a crossover of ESA and SE in the fs-TAS, the ESA evolution was first analysed by fitting the kinetic decay at 630 nm to clarify the kinetics of the excited states (Fig. 6g). Generally, the recombination of the charge carriers occurs within a few picoseconds (ps) due to the extremely strong Coulombic force between the electrons and holes. This phenomenon is the culprit for the loss of dynamical efficiency in photocatalysis.<sup>7</sup> In addition, the photogenerated electrons can easily fall into deep trap states, followed by photoactivity quenching, as previously reported by Durrant.<sup>73</sup> Deep trap

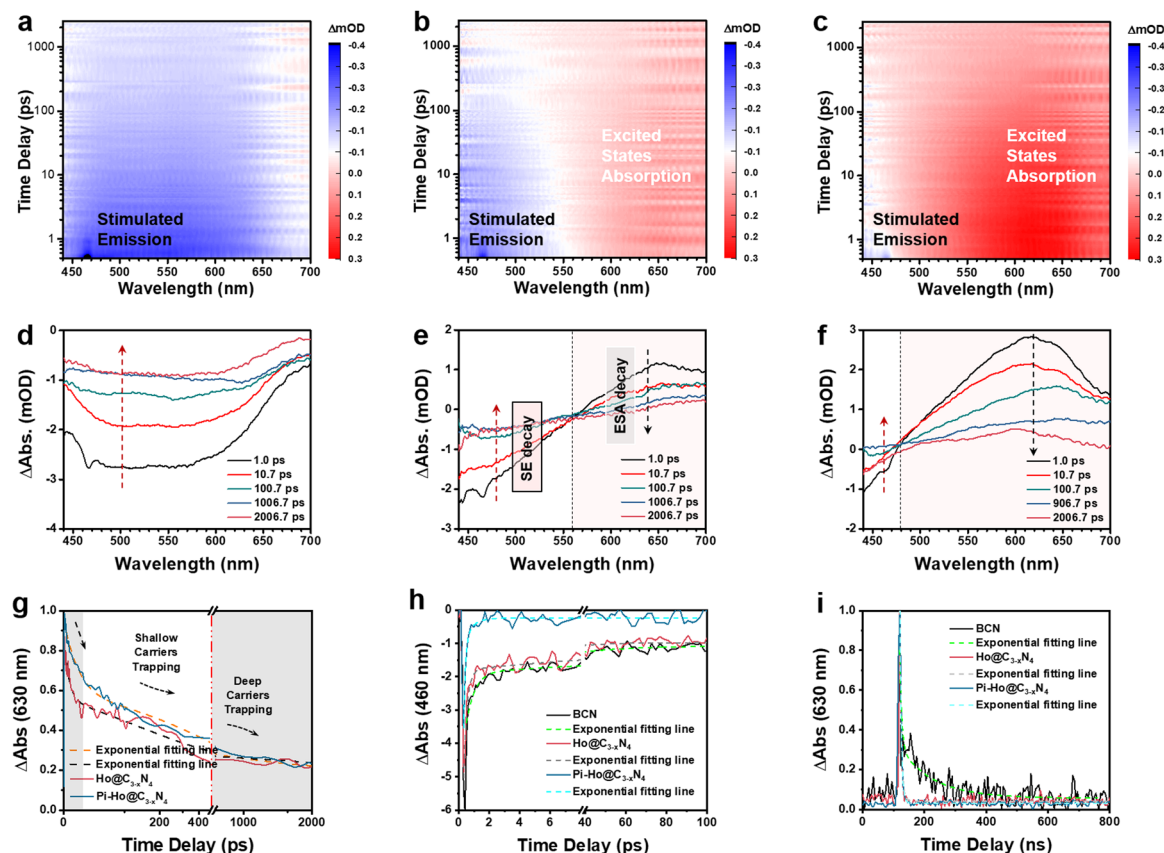


Fig. 6 Time-dependent 3D contour plots of fs-TAS in water (pump laser: 400 nm): (a) BCN, (b)  $\text{Ho@C}_{3-x}\text{N}_4$  and (c)  $\text{Pi-Ho@C}_{3-x}\text{N}_4$ ; absorption decay (Abs.) from fs-TAS: (d) BCN, (e)  $\text{Ho@C}_{3-x}\text{N}_4$  and (f)  $\text{Pi-Ho@C}_{3-x}\text{N}_4$ . Kinetics decay process from fs-TAS observed at (g) 630 nm and (h) 460 nm, with correspondingly fitted exponential-function curves; (i) ns-TAS kinetics decay process observed at 630 nm with the fitted exponential function curves.



states are usually located at energies of  $\sim 1.0$  eV or more below the CB of  $g\text{-C}_3\text{N}_4$ , which is lower than the  $\text{H}^+/\text{H}_2$  reduction potential. More importantly, Majima *et al.* have demonstrated that the long-lived shallow charge trapping is photocatalytically active for effective excited-state survival and proton reduction.<sup>66</sup> DFT simulation results show that both  $\text{Ho@C}_{3-x}\text{N}_4$  and  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  has a hole-trapping state closes to VB. For the  $\text{Ho@C}_{3-x}\text{N}_4$ , the three time constants at 6.0 ps (46.5%), 189.9 ps (39.8%) and 2 ns (13.7%) can be assigned to the photo-generated electron-hole recombination ( $\tau_1$ ), hole trapping ( $\tau_2$ ) and deep electron trapping processes ( $\tau_3$ ), respectively (Table S3, see the ESI<sup>†</sup>). Compared to  $\text{Ho@C}_{3-x}\text{N}_4$ ,  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  shows slightly prolonged recombination (17.9 ps) but much long-lived charge trapping (405.5 ps) with increased portions of 48.9% and 45.6%, respectively. The slower decay and enhanced portion of  $\tau_2$  increases the propensity and population of active electron transport for the photocatalytic reduction reaction. Interstitial-doping sites provide the efficient shallow electron trapping in the  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  for the active electrons' survival for high-probability photocatalytic  $\text{H}_2$  production. To clarify further the important role of *in situ* protonation for photocatalysis, further fs-TAS control experiments were conducted in inert solvents (acetonitrile and ethanol) and the results are shown in Fig. S21 (a<sub>1</sub>–b<sub>1</sub>: acetonitrile and a<sub>2</sub>–b<sub>2</sub>: ethanol) (ESI<sup>†</sup>). In comparison, both of the fs-TAS in the inert cases show relatively strong stimulated emissions from 440 to 530 nm, mainly due to the strong undesirable electron-hole recombination. Excited state absorption is noted in a similar range (550–700 nm) to the water case, Fig. S21 a<sub>1</sub> and b<sub>1</sub> (ESI<sup>†</sup>), but is much weaker than in the water case (Fig. 6f). Furthermore, the enhanced lifetime of  $\tau_2$  (405.5, 156.6, 112.3 ps for the water, acetonitrile and ethanol cases, respectively) demonstrates clearly the significant role of *in situ* protonation in water to enhance the shallow trapping contribution to photodynamics. In addition, an ultralong-lived deep charge trapping ( $\tau_3 > 2000$  ps) was observed in the  $\text{Pi-Ho@C}_{3-x}\text{N}_4$ , which is beyond the time delay window of our fs-TAS system. To further investigate this deep charge trapping, nanosecond transient absorption spectroscopy (ns-TAS, excited by 355 nm laser flash pulse; probing time from 5 to 500 ns) was carried out in water (Fig. S18b, see the ESI<sup>†</sup>) and the results are shown in Fig. S20 (see the ESI<sup>†</sup>). As observed in Table S4 (see the ESI<sup>†</sup>) the shortened lifetime (to  $\sim 5$  ns) of  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  suggest the efficient inhibition of electron trapping in the deep trap state. The fitting results of SE evolution of  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  reveal that the two-time constants with  $<150$  fs (instrument response resolution) and 0.5 ps, correspond to hot-exciton recombination ( $\tau_1$ ) and photogenerated electron emission ( $\tau_2$ ) from CB to VB.

Taking together the DFT analysis and the TAS observations, we propose a proton-fed shallow charge trapping mechanism here for the photocatalytic  $\text{H}_2$  production of the  $\text{Pi-Ho@C}_{3-x}\text{N}_4$ , as illustrated in Fig. 7. Due to the enhanced driving force and charge separation in the  $\text{Pi-Ho@C}_{3-x}\text{N}_4$ , by the synergistic effect of the high redox capacity (*via* C-vacancy) and the “dynamic” charge trapping (*via* natural proton assists shallow trap state survival), the long-lived active electrons approaching

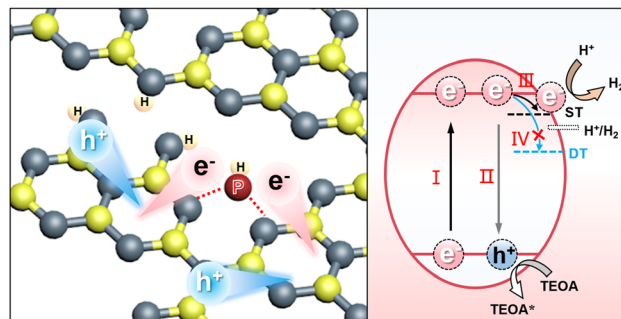


Fig. 7 Photocatalytic  $\text{H}_2$  generation mechanisms over the  $\text{Pi-Ho@C}_{3-x}\text{N}_4$ : (Step I) visible-light excitation and electron-hole separation; (Step II) electron-hole recombination; (Step III) electron-trapping by shallow trap state (ST); (Step IV) electron-trapping by deep trap state (DT).

the shallow states possess sufficient potential and accumulated charge density to react efficiently with adsorbed protons to produce  $\text{H}_2$ . The simulation results, along with the extensive characterizations, clearly confirm the effectiveness of the design approach and explain the experimentally observed enhancement of photocatalytic  $\text{H}_2$  generation efficiency over the  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  photocatalyst.

## Conclusions

In summary, a highly efficient  $g\text{-C}_3\text{N}_4$ -based photocatalyst has been designed and synthesized to incorporate strong synergistic effects of mesoporosity and phosphorous interstitial doping. Moreover, the *in situ* proton-assisted phosphorus interstitial bridge enhances shallow charge trapping in holey  $g\text{-C}_3\text{N}_4$ . The resulting  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  photocatalyst greatly enhances the photodynamics of visible-light-driven  $\text{H}_2$  production. Pore-enriched  $g\text{-C}_3\text{N}_4$  directly improves the redox capacity for the photocatalytic water reduction with more exposed catalytic sites and carbon vacancies. Importantly, the phosphorus interstitial doping provides a new charge migration channel ( $-\text{N}-\text{P}=\text{N}-$ ), still with a robust driving force to guarantee highly efficient spatial separation of the charge carriers and long-lived active intermediates. Theoretical simulation and *in situ* transient absorption spectroscopy clearly show the long lifetime of the shallow-trapped charges and efficient kinetics in H-terminated catalysts. Proton-induced kinetics effectively promote the active species to participate in the photocatalytic reactions before being recombined. As a result, the as-prepared  $\text{Pi-Ho@C}_{3-x}\text{N}_4$  photocatalyst shows an 8-fold enhancement in photocatalytic  $\text{H}_2$  production ( $6323$  vs.  $790 \mu\text{mol h}^{-1} \text{g}^{-1}$ ), and a 5.8-fold increase ( $5.08\%$  vs.  $0.87\%$  at  $420$  nm) in quantum efficiency, compared with bulk  $g\text{-C}_3\text{N}_4$ . The proposed synergistic effect of water and interstitial P-doping offers a promising strategy for solar-to-hydrogen conversion with balanced redox capability and effective charge carrier separation.

## Author contributions

W. Wang: investigation, formal analysis, data curation, and writing of the original draft. L. Du: investigation, data curation,





and writing – review & editing. R. Xia: investigation. R. Liang: investigation and data curation. T. Zhou: investigation. Z. Yan: investigation. C. Shang: investigation and revision. H. Lee: investigation. D. L. Phillips: supervision, conceptualization, and writing – review & editing. Z. X. Guo: supervision, funding acquisition, conceptualization, project administration, and writing – review & editing. The manuscript was written by W. Wang, Z. Guo and D. L. Phillips with contributions from the other co-authors. All authors participated in the discussion of the research and have given approval to the final version of the manuscript.

## Conflicts of interest

There are no conflicts to declare.

## Acknowledgements

We gratefully acknowledge the following sponsorship and support from the Hong Kong UGC-TRS (T23-713/22-R) award, the Environment and Conservation Fund (ECF 2021-152), the RGC-EU Collaborative Programme initiative (E-HKU704/19), the RGC-GRF (17302419, 17316922), the Natural Science Foundation of Zhejiang Province (LY19E020008), the Key-Area Research and Development Program of Guangdong Province (2020B0101370003), the Joint Laboratory Funding Scheme (JLFS/P-704/18), the “Hong Kong Quantum AI Lab Ltd” funded by the AIR@InnoHK, launched by the Innovation and Technology Commission (ITC), and the URC Platform Technology Fund and the start-up support from the University of Hong Kong. Publication made possible in part by support from the HKU Libraries *via* the Open Access Author Fund.

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