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# Understanding the limits to short-range order suppression in many-component disordered rock salt lithium-ion cathode materials†

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Suppressing unfavourable short-range ordering in disordered rock salt lithium-ion cathode materials is seen as a key research goal on their route to commercialisation. In this study we use cluster-expansion-driven Monte Carlo simulations of a model 3d-transition metal disordered rock salt oxyfluoride system to investigate the effect of many component cation substitution on the suppression of short-range ordering in disordered rock salt cathode materials. We confirm that many-cation substitution is effective in suppressing short-range ordering, but has diminishing returns on increasing the number of component transition metals, or alternatively, increasing the size of the long-range lithium diffusion network as the number of transition metals increases. We particularly emphasize the critical role of lithium excess and fluorine content in the success of the “high-entropy” cation substitution strategy: short-range ordering is strongly influenced by cation–anion bonding preferences, underscoring the need to consider the full composition of the target system when designing high entropy lithium-ion cathode materials.

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## 1. Introduction

The exploration and optimisation of novel chemistries for lithium-ion cathode materials remains a vast area of research for the materials science community<sup>1–5</sup> owing to the crucial role of electrochemical energy storage in transitioning to a low carbon economy. Disordered rock salt lithium-ion cathodes are an emerging materials class in this area.<sup>6–9</sup> These materials share the close-packed oxygen framework present in most commercial lithium-ion cathode materials such as the layered rock salt cobalt- and nickel-rich structures<sup>10,11</sup> conferring an excellent mass-to-charge ratio and volumetric efficiency. While performant, layered cobalt and nickel rich structures have poor long-term techno-economic prospects.<sup>12</sup> This is in contrast to – for example – manganese and iron which are cheaper and more abundant, but their layered structural analogues are not stable over many intercalation cycles.<sup>13–15</sup>

Embracing cation disorder enables the use of novel cathode chemistries that do not need to maintain a layered ordering on cycling.<sup>7,16–23</sup> The performance of any intercalation cathode is limited by the ability of lithium to reversibly intercalate into the host framework which is dependent on atomistic diffusion processes. Fig. 1 illustrates a cubic close-packed oxygen lattice.

In both the layered and disordered systems lithium diffusion occurs from an octahedral hole in the O<sup>2–</sup> framework *via* an octahedral face, into an interstitial tetrahedron. The diffusing ion then occupies a new octahedral site *via* another shared octahedral–tetrahedral face (an “oct–tet–oct” hop).<sup>6</sup> The energy barrier for lithium hopping is then strongly influenced by the stability of lithium in the tetrahedral hole.<sup>24</sup> The stability of lithium in this tetrahedral transition state depends on factors such as the identity of the face-sharing cations: too many face-sharing transition metals destabilises the lithium in the tetrahedral transition state increasing the energy barrier for lithium diffusion.<sup>6,25</sup> In disordered rock salt cathodes a percolating network of lithium sites connected by tetrahedral holes with no face-sharing transition metals (referred to as “0-TM” sites) is generally found to be necessary for effective long-range diffusion.<sup>6,8,25</sup> The pathway for single-ion lithium hopping and a range of potential transition state local environments in disordered rock salts are shown as a 2D schematic in Fig. 2.

In a fully random disordered rocksalt with the general formula Li<sub>1+x</sub>M<sub>1-x</sub>O<sub>2</sub> – where M is used as a generic symbol for any transition metal – a lithium excess of  $x \geq 0.09$  is required to ensure a long range percolating 0-TM network.<sup>6,26</sup> It is found the connectivity of the lithium diffusion network is poorer than expected in a totally random system in almost all cases, and substantial lithium excesses are required to give good long-range connectivity of the 0-TM network.<sup>25</sup> Minimising the lithium excess is desirable, as the additional lithium substitutes for the redox-active transition metal, reducing the capacity or increasing the anion-redox activity which can cause voltage hysteresis and structural degradation in disordered rock salt cathodes.<sup>27</sup>

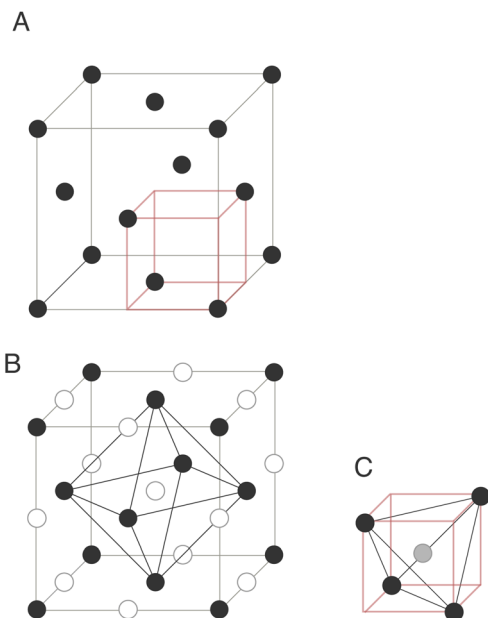
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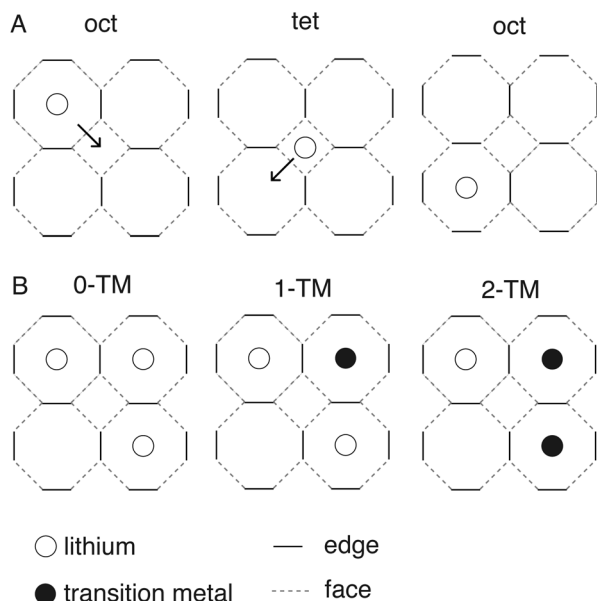
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**Fig. 1** (A) Face-centered cubic (FCC) lattice, such as the one exhibited by  $O^{2-}$  ions in many lithium-transition metal oxides. (B) Octahedral holes in the FCC lattice (white) that may be filled by cations in a lithium-transition metal oxide. (C) Tetrahedral holes (grey) which form the transition states for lithium diffusion in a disordered rock salt cathode material.

The cause of the reduced lithium percolation in disordered rock salt materials relative to the random limit is understood to be a consequence of the short-range ordering preferences of the cations, which are comprised of contributions from bonding



**Fig. 2** (A) 2D representation of the oct–tet–oct hop of a lithium ion diffusing in a rock salt structured cathode material. (B) 2D representation of the different face sharing ions for the transition-state tetrahedral site in the oct–tet–oct hop. The energy barrier for the hop will increase as a function of the number of transition metals face-sharing with the tetrahedral site.

preferences and electrostatics.<sup>8,25,28,29</sup> Thus, finding ways to suppress unfavourable short-range ordering is crucial for improving the capacity and overall performance of disordered rock salt cathode materials. Fluorine substitution allows for the manipulation of short-range ordering preferences on the anion framework by introducing bond-strength considerations between the cations and the anions which will compete with the cation–cation interactions to determine the overall short-range order.<sup>30</sup> Transition metals bind more strongly to oxygen, creating locally lithium–fluorine-rich regions and transition-metal–oxygen-rich regions, which alters the short-range ordering. Incorporating enough fluorine can therefore create lithium-rich regions with high numbers of 0-TM sites, and good lithium transport. However, high levels of fluorine incorporation requires energy-intensive mechanochemical synthesis<sup>31–33</sup> and can lead to phase separation.<sup>30</sup>

An alternative strategy for suppressing unfavourable short-range ordering in disordered rock salt cathode materials is to substitute many different species on the cation sublattice, creating a “high-entropy” material. This approach, first proposed by Lun *et al.*,<sup>34</sup> was inspired by the homogeneous cation distribution observed in high-entropy oxides.<sup>35,36</sup> While not the first work to successfully synthesize high-entropy disordered rock salt lithium-ion cathode materials,<sup>37</sup> Lun *et al.* were the first to link the performance of these materials to short-range order suppression. By increasing the number of species on the cation sublattice, they observed an increase in first-cycle capacity. This was linked a suppression of short-range ordering validated by electron diffraction experiments which showed a reduction in diffuse short-range order scattering as the number of transition metals in the sample was increased. Alternatively, this can be seen as an increase in cation configurational entropy. The short-range order suppression was then linked to an increase in the size of the long-range lithium percolation network by cluster-expansion-driven Monte Carlo simulations.

To extend the current understanding of short range order suppression in many-component disordered rock salt transition metal oxyfluoride cathode materials, we examine the lithium-transition metal ordering in rock-salt-structured 3d-transition-metal lithium oxyfluorides using cluster-expansion-driven Monte Carlo simulations. Such simulations have proved to be an excellent tool in unravelling short range ordering preferences in lithium-ion cathode materials.<sup>25,27,34,38</sup> We examine how the lithium percolation network evolves as a function of the number of transition metals on the cation sublattice and varying lithium excess and fluorine content. In doing so, we identify compositional guidelines for designing effective many-component disordered rock salt lithium-ion cathode materials, showing that any short range order suppression from an increase in cation configurational entropy is highly dependent on the lithium excess and fluorine content.

## 2. Methodology

A cluster expansion model, trained from density functional theory (DFT) calculations was used to drive Monte Carlo simulations to sample the short-range ordering behaviour of compositionally



complex disordered rock salt lithium-ion cathode materials. The cluster expansion basis was taken as a rock salt structure with  $a = 4.1 \text{ \AA}$  with  $\text{Li}^+$ ,  $\text{Mn}^{2+}$ ,  $\text{Mn}^{3+}$ ,  $\text{Mn}^{4+}$ ,  $\text{Ni}^{2+}$ ,  $\text{Ti}^{4+}$ ,  $\text{Fe}^{3+}$ , and  $\text{Cr}^{3+}$  allowed to occupy the cation site, and  $\text{O}^{2-}$  and  $\text{F}^-$  allowed to occupy the anion site. Ancillary analysis relied on the pymatgen,<sup>39</sup> seaborn,<sup>40</sup> pandas,<sup>41</sup> scipy,<sup>42</sup> numpy<sup>43</sup> and scikit-learn<sup>44</sup> packages. Visualisation relied on the use of perceptually uniform scientific color maps. Training data and fitted cluster expansions can be found in the associated dataset.<sup>45–47</sup>

## 2.1 Training data generation

All DFT calculations were performed using the plane-wave DFT Vienna *Ab initio* Simulation Package (VASP).<sup>48–50</sup> Valence electrons were described by a plane-wave basis set with a cutoff of 680 eV. Interactions between core and valence electrons were described using the projector-augmented wave (PAW) method.<sup>51,52</sup> Electronic exchange and correlation were approximated using the Perdew–Burke–Ernzerhof functional, revised for solids and surfaces (PBEsol)<sup>53</sup> with rotationally averaged Hubbard  $U$  corrections of 3.9 eV, 3.7 eV, 5.3 eV, 6.2 eV and 4.2 eV applied to the Mn, Cr, Fe, Ni and Ti 3d orbitals respectively to approximately correct for the self-interaction error (SIE).<sup>54–56</sup> Reciprocal space was sampled with a discretisation of  $0.35 \text{ \AA}^{-1}$ . All calculations were initialised with transition metal ions in a ferromagnetic configuration.

## 2.2 Cluster expansion fitting

The final cluster expansion was trained on a total of 7215 unique DFT calculations after filtering out structures which relax to off-lattice configurations, and possess transition metals in oxidation states not included in the model. Oxidation states were assigned using calculated magnetic moments. The variable oxidation state of Mn was assigned using a Bayesian black-box optimization process.<sup>57,58</sup> We employed the icet and trainstation software packages to train the cluster expansion.<sup>44,59,60</sup> The fitting used a pair-cluster cutoff of  $7.1 \text{ \AA}$  and a triplet-cluster cutoff of  $4.8 \text{ \AA}$  (quadruplet clusters were found to make only a very minor contribution to the model performance for increased complexity, so were excluded, as has been noted in other high-component cluster expansion studies of lithium transition metal oxides<sup>57</sup>). The final model has a cross validation score of 9.95 meV. More information on the fitting procedure and training set generation, including a discussion of limitations of the DFT+ $U$  methodology used in this paper can be found in the ESI.†

## 2.3 Monte Carlo simulations

Monte Carlo simulations were performed using the mchammer package.<sup>59</sup> Simulations of the lithium-excess simulations used a 990 atom expansion of the primitive cell described by the transformation matrix

$$\begin{bmatrix} 8 & 0 & -1 \\ -4 & 7 & 1 \\ -3 & -3 & 9 \end{bmatrix}.$$

This structure was selected using an algorithm to determine a cell with optimal cubic character, as implemented in the atomic simulation environment.<sup>61</sup> The simulations of the  $\text{LiMO}_2$  system used a 1000 atom cell,

$$\begin{bmatrix} 7 & 1 & -1 \\ -4 & 8 & 0 \\ -3 & -4 & 9 \end{bmatrix}.$$

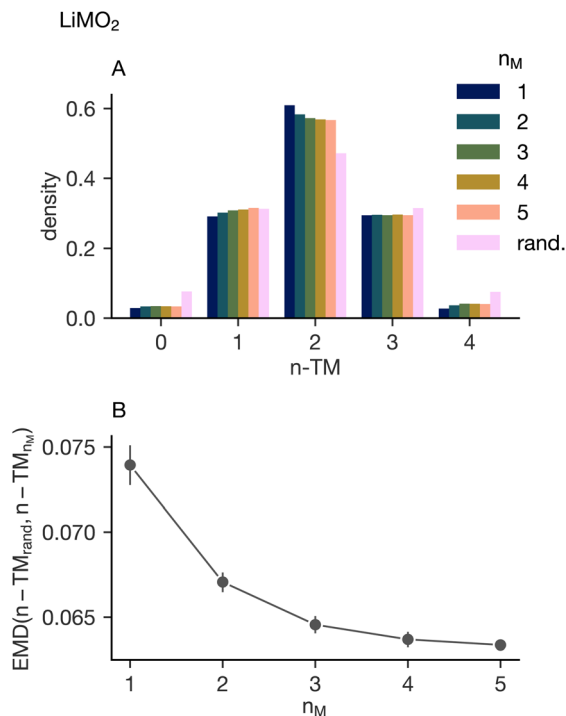
Each Monte Carlo simulation was run at 2000 K. 2000 K falls within the range of empirical simulation temperatures commonly used for Monte Carlo simulations of disordered rock salt structures<sup>27,34,62</sup> and we expect it to capture experimentally relevant trends in the short range order formation. Each simulation ran for at least 5 000 000 trial steps. Values reported from the Monte Carlo simulations are each averaged over five simulations of each composition presented, with 1000 structures extracted from each Monte Carlo simulation, *i.e.* each value is averaged over 5000 data points. The crystal-torture package was used to extract percolation statistics from the Monte Carlo simulations.<sup>63</sup>

## 3. Results and discussion

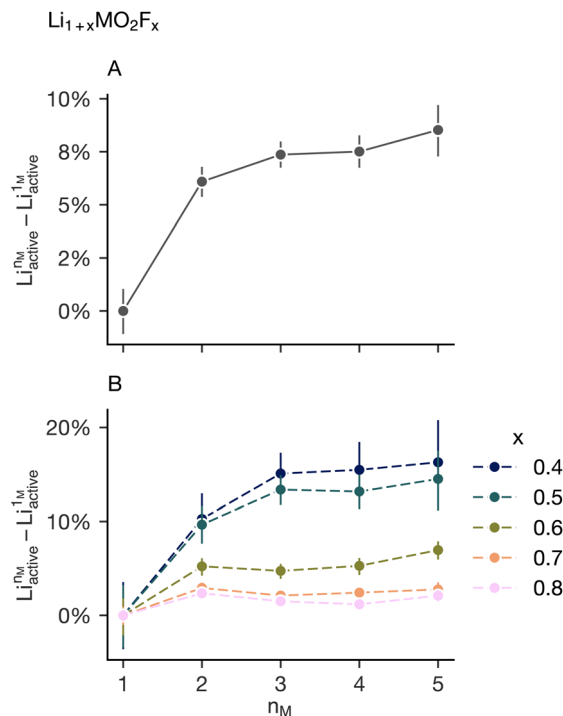
First, we confirm that an increase in configurational entropy of the cation sublattice suppresses the short range ordering. Fig. 3A shows a histogram of the occurrences of the different  $n$ -TM sites in  $\text{LiMO}_2$  where M represents all equimolar charge-balanced compositions of  $\text{Mn}(\text{III})$ ,  $\text{Cr}(\text{III})$ ,  $\text{Fe}(\text{III})$ ,  $\text{Ni}(\text{II})$  and  $\text{Ti}(\text{IV})$ . These are then separated into 1, 2, 3, 4, and 5 transition metal simulations. Visual inspection of the histogram indicates that higher-component systems result in an  $n$ -TM probability mass function more similar to a random one. However, even when mixing five transition metals, the difference between the simulated and random distribution appears significant. Fig. 3B shows the Earth Mover's Distance (EMD) between the random and simulated distributions. The EMD is a measure of dissimilarity between two probability distributions with a smaller EMD value indicating a closer match.<sup>64</sup>

The EMD taken between the random and simulated systems supports the observations that can be made from the histogram. The distribution of  $n$ -TM sites becomes increasingly similar to a random distribution as a function of the number of transition metals in the simulation and supports the hypothesis that increasing cation configurational entropy leads to a more homogeneous distribution of cations in lithium transition metal oxide disordered rock salts. However, there is a decrease in the rate at which the EMD value decreases as the number of transition metals increases, indicating that the distribution becomes more similar to a random distribution at a slower rate. “Diminishing returns” are observed in the short-range order suppression with respect to the number of transition metals. This indicates that increased cation configurational entropy can improve mixing to an extent, but is unlikely to fully overcome the electrostatic and elastic effects that drive cation short-range ordering in lithium transition metal oxides.<sup>8,25</sup> This implies that





**Fig. 3** (A) Histogram showing the occurrence of the different possible  $n$ -TM sites in Monte Carlo simulated  $\text{LiMO}_2$ , where M is all equimolar charge-balanced combinations of Mn(III), Cr(III), Fe(III), Ni(II) and Ti(IV). (B) The Earth Mover's Distance (EMD) between the Monte Carlo sampled distributions of  $n$ -TM sites, and the fully random distribution. The error bars show the 95% confidence interval in the EMD.



**Fig. 4** (A) The mean change in active lithium in a  $\text{Li}_{1+x}\text{MO}_2\text{F}_x$  as a function of the number of transition metals in the Monte Carlo simulation, relative to the average percolating lithium in single transition metal systems decomposed by the lithium excess and fluorine content. The error bars show the 95% confidence interval in the active lithium fraction. (B) Showing the same data, but decomposed per value of  $x$ .

moderate entropy systems will see most of the short-range order suppression benefit of the many-transition-metal systems.

To explore the relationship between the number of transition metals in the simulation and short-range order suppression in a more experimentally relevant compositional space we run additional Monte Carlo simulations on fluorinated lithium-excess disordered rock salt materials. We simulate all equimolar charged balanced compositions maintaining an average transition metal oxidation state of 3+ in the compositional range  $\text{Li}_{1+x}\text{MO}_2\text{F}_x$ , where  $x \in \{0.4, 0.5, 0.6, 0.7, 0.8\}$ . The transition metals included remain Mn(III), Cr(III), Fe(III), Ni(II) and Ti(IV). The lithium that is expected to be diffusive, and therefore contribute the capacity, is that which is connected *via* a percolating network of 0-TM sites,<sup>6,8,25</sup> we label this  $\text{Li}_{\text{active}}$ . Fig. 4A shows the calculated change in  $\text{Li}_{\text{active}}$  from the Monte Carlo simulation as the number of transition metals increases, averaged over all values of  $x$ . Increasing the number of transition metals on the cation sublattice generally results in an increase in  $\text{Li}_{\text{active}}$ . However, the increase in the size of the lithium diffusion network levels off as the number of transition metals increases mirroring the short-range order suppression behavior highlighted in Fig. 3. This suggests that the observation of “diminishing returns” of short-range order suppression noted for the  $\text{LiMO}_2$  system is extensible to the lithium-excess oxyfluorides.

Examining only the mean behaviour in the change in activate lithium as a function of  $n_M$  disguises the strongly different

responses for each value of  $x$ . Fig. 4B shows the average value decomposed into the individual values of  $x$ : lower values of  $x$  have a much stronger increase in the percolating lithium fraction as a function of the number of transition metals in the simulation cell. This suggests any short-range order suppression observed in “high-entropy” disordered rock salt cathode will be highly dependent on the lithium and fluorine content and warrants further investigation. This dependence is highlighted in Fig. 5 which shows analysis of simulations on the system  $\text{Li}_{1+x}\text{MO}_2\text{F}_x$  over a wider value of  $x$  ( $x \in \{0.1, 0.2, 0.3, 0.4, 0.5, 0.6, 0.7, 0.8, 0.9, 1\}$ ). This is done for  $M = \text{Mn}$  ( $n_M = 1$ ) and  $M = \text{Mn}_{0.2}\text{Cr}_{0.2}\text{Fe}_{0.2}\text{Ni}_{0.2}\text{Ti}_{0.2}$  ( $n_M = 5$ ). The increase in  $\text{Li}_{\text{active}}$  as a function of  $x$  for the two systems is shown in Fig. 5A, with the difference between these two series highlighted in Fig. 5B. Low values of  $x$  ( $x < 0.3$ ) and high values of  $x$  ( $x > 0.6$ ) see very little change in the amount of active lithium on increasing the number of transition metals, compatible with the observations made from Fig. 4.

To explain these observations we note the lithium excess and fluorine substitution level are not necessarily coupled and extend the previous Monte Carlo simulations on the system  $\text{Li}_{1+x}\text{MO}_2\text{F}_y$ , *i.e.* the lithium excess,  $x$ , and fluorine content,  $y$ , are now varied independently between 0.2 and 1 in the  $n_M = 1$  and  $n_M = 5$  systems examined previously. Fig. 6A and B confirm that the active lithium fraction in the 5-TM system for the



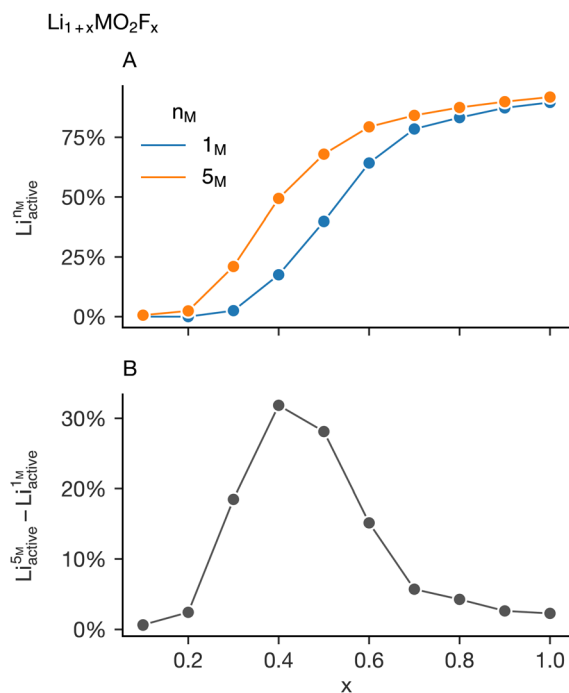


Fig. 5 (A) The change in the mean active lithium percentage as a function of  $x$  in  $\text{Li}_{1+x}\text{MnO}_2\text{F}_x$  and  $\text{Li}_{1+x}\text{Mn}_{0.2}\text{Cr}_{0.2}\text{Fe}_{0.2}\text{Ni}_{0.2}\text{Ti}_{0.2}\text{O}_2\text{F}_x$  as a function of  $x$ . (B) The difference between the two trends shown in plot (A).

equivalent  $x$  and  $y$  is generally higher. The difference between Fig. 6A and B, shown in Fig. 6C, highlights that the increase in active lithium fraction from the 1-TM system to the 5-TM system is non-uniform as  $x$  and  $y$  are varied. The increase is minimal at high  $x$  and  $y$ , and shows a coupled dependence on lithium and fluorine contents at intermediate compositions. Increasing the fluorine content (increasing  $y$ ) decreases the value of  $x$  for which the change between the 1-TM and 5-TM systems is at a maximum.

This strong dependence on short-range order suppression by increased cation substitution on the lithium-excess and fluorine content can be understood by considering the cation–anion bonding preferences. The orbital overlap between transition metal d-states and oxygen 2p states will be much stronger than any interaction between the transition metal ions and the fluorine. This general chemical insight is reflected in the cation–coordination preferences observed in multiple previous Monte Carlo studies of transition metal oxyfluoride disordered rock salt cathode materials, with lithium having an above-average number of fluorine nearest neighbors, and transition metals coordinating to a higher number of oxide ions.<sup>27,30,62</sup>

Fig. 7 shows the average coordination number of lithium and the transition metals with fluorine, for two cases: a Li–F poor case,  $\text{Li}_{1.5}\text{MO}_2\text{F}_{0.5}$ , and a Li–F rich case,  $\text{Li}_{1.8}\text{MO}_2\text{F}_{0.8}$ . The coordination numbers are displayed as a function of the number of transition metals in the simulation,  $n_M$ . As anticipated, lithium coordinates much more strongly to fluorine than the transition metals do. Among the transition metals,  $\text{Ti}(\text{IV})$  and  $\text{Cr}(\text{III})$  most strongly favour  $\text{O}^{2-}$  coordination. This

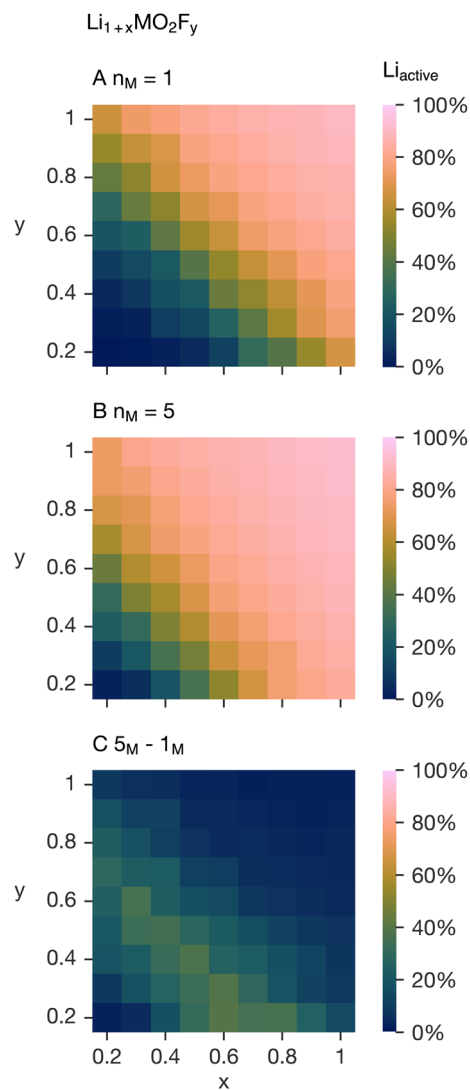


Fig. 6 (A) Active lithium fraction in  $\text{Li}_{1+x}\text{MnO}_2\text{F}_y$  as a function of  $x$  and  $y$  (B) active lithium fraction in  $\text{Li}_{1+x}\text{Mn}_{0.2}\text{Cr}_{0.2}\text{Fe}_{0.2}\text{Ni}_{0.2}\text{Ti}_{0.2}\text{O}_2\text{F}_y$  as a function of  $x$  and  $y$ , (C) difference map between plots (A) and (B). Off-diagonal compositions on these plots are not strictly charge balanced, however we expect these fixed-composition simulations to capture local bonding preferences which determine the short range order.

preference, due to the high charge of  $\text{Ti}(\text{IV})$  and the  $d^3$  electron configuration of  $\text{Cr}(\text{III})$ ,<sup>34</sup> results in O-rich coordination environments.  $\text{Ni}(\text{II})$ , with its comparatively low charge, exhibits the highest average fluorine coordination among the transition metals. In the Li–F poor case,  $\text{Ni}(\text{II})$  takes on higher-than-random F-coordination, influenced by the other transition metals stronger  $\text{O}^{2-}$  preferences, driving  $\text{Ni}(\text{II})$  to coordinate with more fluorine, despite its inherent inclination towards oxygen. In the Li–F rich case, the heightened Li–F content serves to bring down the Ni–F-coordination below that expected in a random distribution. The nickel–fluorine average coordination decreases with increased  $n_M$  in these simulations; for charge conservation, there is always an equal  $\text{Ti}(\text{IV})$  content. Introducing  $\text{Mn}(\text{III})$  and  $\text{Fe}(\text{III})$ , with a weaker preference for  $\text{O}^{2-}$  coordination than  $\text{Ti}(\text{IV})$  and  $\text{Cr}(\text{III})$ , enables  $\text{Ni}(\text{II})$  to occupy





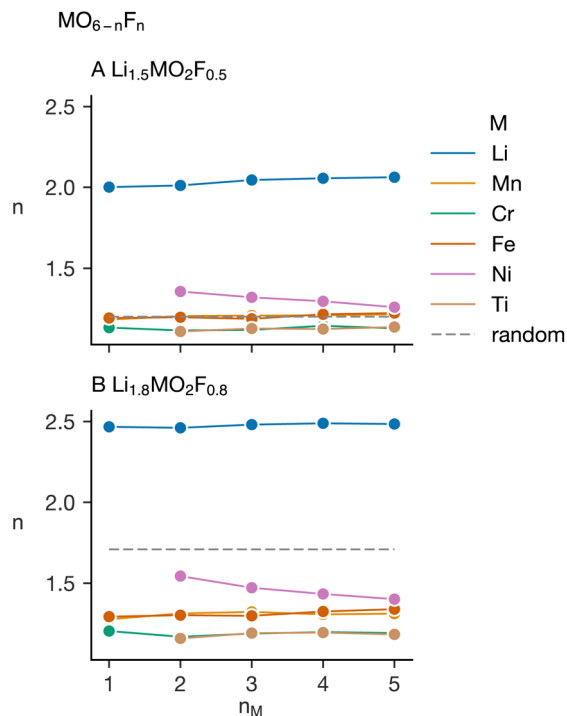


Fig. 7 (A) A plot of metal–fluorine coordination numbers in  $\text{Li}_{1.5}\text{MO}_2\text{F}_{0.5}$  as a function of the number of transition metals included in the simulation. (B) A plot of metal–fluorine coordination numbers in  $\text{Li}_{1.8}\text{MO}_2\text{F}_{0.8}$  as a function of the number of transition metals included in the simulation. In both cases the grey dashed line shows the average metal–fluorine coordination number in the fully random system.

cation sites with comparatively high oxygen coordination. Importantly, these plots underscore that multi-component cation substitution may somewhat mitigate short-range ordering preferences in these materials, but the strong preferences for transition-metal–oxygen bond formation persist.

Bringing together the observations made in this study, we are able to suggest some guidelines for designing compositionally complex disordered rock salt cathode materials. Our simulations support that substituting many transition metals on the cation sublattice is a route to suppressing short-range order in lithium transition metal oxides, but diminishing returns are observed for short-range order suppression as the number of transition metals in the simulation is increased. This is likely due to extensive transition metal substitution not being able to fully counteract electrostatic and elastic interactions between the transition metals and the lithium which drive short-range order formation in disordered rock salt cathodes. This suggests that most of the benefit of short-range order suppression observed in many component systems should be observed in more “moderate entropy” disordered rock salts.

In addition, the relationship between short-range order suppression and the number of constituent cations is strongly influenced by the lithium-excess and fluorine content. While cation mixing is able to strongly increase the active lithium fraction in systems with comparatively low lithium–fluorine content, when the lithium excess is high and the fluorine

content increases, the short range ordering formation is dominated by transition-metal–oxygen and lithium–fluorine bonding preferences as opposed to any improved cation mixing from many-transition-metal substitution.

It is worth noting that there may be beneficial effects in high-entropy disordered rock salt cathode materials with respect to lithium mobility. A “more disordered” system will likely see a wider range of site energies within the lithium transport network improving the conductivity as has been observed in a range of solid electrolyte materials;<sup>65–68</sup> these concepts are clearly extensible to intercalation cathodes.<sup>69,70</sup> However, we postulate that any increase in mobility will also be minimal at high fluorine and lithium contents as the formation of Li–F rich regions will minimise the impact of any local distortions that influence the site energies in the lithium diffusion pathway from introducing more transition metals.

The optimization of disordered rock salt cathode materials is a delicate balance between optimising the capacity of the material and the size of the percolating lithium network, while also considering the cycling stability and the synthesis methods. High lithium and fluorine content can lead to a large capacity from a large lithium-percolation network<sup>30,71</sup> (and redox participation from oxygen<sup>7</sup>) but poor cycling stability<sup>27</sup> and high energy intensive synthesis methods.<sup>7,32</sup> In addition, in rock salt structured transition metal oxides, electronic conductivity is mediated by small polaron hopping between adjacent redox-active transition metal sites.<sup>72–75</sup> While understudied in comparison to lithium ion percolation in disordered rock salt cathodes, the electronic carriers must also be able to percolate. As such, reducing the redox-active transition metal content *via* increasing the lithium content (and  $d^0$  spectator ion content) may significantly reduce the electronic conductivity of the cathode,<sup>76</sup> reducing, for example, the rate performance of the battery. These issues could be alleviated by utilising many-component cation mixing as an additional variable in the compositional design of disordered rock salt cathode materials, keeping lithium excesses and fluorine contents low while maintaining a large percolating lithium network. However, our results show that bonding interactions between cations and anions in many-component disordered rock salt cathode materials must be taken into account for optimal short-range order suppression.

## Conflicts of interest

There are no conflicts to declare.

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