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An efficient blue-excitable broadband $\text{Y}_3\text{ScAl}_4\text{O}_{12}:\text{Ce}^{3+}$ garnet phosphor for WLEDs†

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Most commercial phosphor-converted white light-emitting diodes (pc-WLEDs) are manufactured with blue LED chips and yellow-emitting $\text{Y}_3\text{Al}_5\text{O}_{12}:\text{Ce}^{3+}$ (YAG: Ce^{3+}) garnet phosphor, but the lack of blue-green light in the spectrum results in a low color rendering index (CRI). In this paper, we synthesized $\text{Y}_3\text{ScAl}_4\text{O}_{12}:\text{Ce}^{3+}$ (YSAG: Ce^{3+}) by replacing Al^{3+} in YAG: Ce^{3+} with Sc^{3+} . The introduction of Sc^{3+} with a larger ionic radius through a cation substitution strategy causes lattice expansion, elongation of the Y–O bond, and ultimately a decrease in Ce^{3+} 5d level crystal field splitting. As a consequence, the emission spectrum undergoes a blue-shift of 10 nm. Furthermore, the YSAG: Ce^{3+} phosphor exhibits good thermal stability, and its emission intensity at 423 K is about 58% of that at 303 K. Moreover, the analysis of Eu^{3+} emission spectra demonstrates that the introduction of Sc^{3+} resulted in a slight reduction of the dodecahedral lattice symmetry. YSAG: Ce^{3+} effectively compensates for the lack of the blue-green region, and WLEDs with high color rendering index (90.1), low color temperature (4566 K) and high luminous efficiency (133.59 lm W^{-1}) were prepared using the combination of YSAG:0.08 Ce^{3+} , $\text{CaAlSiN}_3:\text{Eu}^{2+}$ and 450 nm blue chips. These findings indicate that YSAG: Ce^{3+} garnet phosphor has potential to be used in high quality WLEDs.

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1 Introduction

Phosphor-converted white light-emitting diodes are widely used in solid-state lighting because of their high luminous efficiency, low energy consumption, long service life and environmental friendliness.^{1–3} Some commercialized pc-WLEDs are obtained by combining YAG: Ce^{3+} phosphors with blue light chips. This method has the advantages of simple operation and high luminous efficiency. However, the blue-green and red regions of the spectrum are insufficient in the WLED, leading to a low color rendering index (CRI < 80) and a high correlated color temperature (CCT > 5000 K).^{4–6} The addition of red phosphors such as $(\text{Ca},\text{Sr})\text{AlSiN}_3:\text{Eu}^{2+}$ and $\text{Sr}_2\text{Si}_5\text{N}_8:\text{Eu}^{2+}$ can improve CRI and reduce CCT.⁷ Nonetheless, there is still a lack of high-efficiency and thermally stable blue-

green phosphors, which makes it a challenge to achieve an exceptionally high color rendering index (CRI > 90).

Rare earth Ce^{3+} ions are known to be excellent activators for WLED phosphors due to their typical parity-allowed $4f \leftrightarrow 5d$ transitions.⁸ When Ce^{3+} is doped into a suitable matrix, the Ce^{3+} activated phosphor exhibits strong absorption in the near ultraviolet to blue region due to the $4f \rightarrow 5d$ transition. In addition, these phosphors can produce wide emission bands as the parity-allowed $5d \rightarrow 4f$ transition of Ce^{3+} ions.⁹ Since the $5d \rightarrow 4f$ transition of Ce^{3+} ion is susceptible to the influence of crystal field environment, the luminescent properties of Ce^{3+} doped in different matrix materials exhibit noticeable differences. Consequently, selecting the appropriate matrix material is crucial to attain the desired emission characteristics of Ce^{3+} ions.^{10–12}

Garnet-type phosphors belong to the cubic crystal system, the space group is $Ia\bar{3}d$, and its chemical general formula can be written as $\text{A}_3\text{B}_2\text{C}_3\text{O}_{12}$, where A, B and C are cations in dodecahedron, octahedron and tetrahedron, respectively, which are interconnected by O atoms to form a stable garnet structure.^{13,14} The cations in the garnet structure can be flexibly replaced to achieve structural modulation and derivation of new materials, exhibiting remarkable modulation of the spectrum.¹⁵ $\text{Y}_3\text{Al}_5\text{O}_{12}$ is one of the best-known garnet phosphor substrates, Y^{3+} corresponds to A^{3+} and occupies the dodecahedral site, two Al^{3+} occupy the octahedral site corres-

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ponding to B^{3+} , and the remaining three Al^{3+} occupy the tetrahedral site corresponding to C^{3+} .^{16,17} When Sc^{3+} is introduced into the matrix, the Sc^{3+} replaces an Al^{3+} located in the octahedron. Some studies have shown that the introduction of Sc^{3+} will lead to the increase of the disorder of the matrix structure, which is conducive to the broadening of the emission spectrum of rare earth ions doped in the matrix.^{18–22}

In this paper, $YSAG:Ce^{3+}$ and $YAG:Ce^{3+}$ phosphors were successfully synthesised by high-temperature solid-phase method and their luminescent properties were investigated. Compare with the emission spectrum of $YAG:Ce^{3+}$, the emission spectrum of $YSAG:Ce^{3+}$ is blue-shifted by 10 nm. Moreover, Eu^{3+} was used as a fluorescence probe to calculate the red-orange ratio in different substrates, and the similar ratio indicates that the introduction of Sc^{3+} ions has little effect on the dodecahedral symmetry. Finally, the WLED device was prepared and its parameters were tested.

2 Experimental section

2.1 Materials and synthesis

A series of $YSAG:xCe^{3+}$ ($x = 0.02, 0.05, 0.08, 0.11, 0.14$ and 0.17) and $YAG:0.08Ce^{3+}$ samples were prepared by high-temperature solid-state method. The raw materials were Y_2O_3 (4N), Sc_2O_3 (4N), Al_2O_3 (4N) and CeO_2 (4N), all of which were accurately weighed according to the stoichiometric ratio, and H_3BO_3 at 1% of the total mass of the raw material was added as a flux. The raw materials were carefully ground in an agate mortar for 20 min, mixed well and transferred to an alumina crucible, which was then placed in a tube furnace under a reductive atmosphere (90% N_2 + 10% H_2). The temperature of the tube furnace was slowly increased to 1530 °C and held at this temperature for 4 h. After cooling to room temperature, the crucible was removed and the agglomerated product was reground into powder to obtain bright yellow phosphor. $YSAG:0.03Eu^{3+}$ and $YAG:0.03Eu^{3+}$ samples were also prepared by high-temperature solid-state method. The raw materials were Y_2O_3 (4N), Sc_2O_3 (4N), Al_2O_3 (4N) and Eu_2O_3 (4N), all of which were accurately weighed according to the stoichiometric ratio, and H_3BO_3 at 1% of the total mass of the raw material was added as a flux. The raw materials were carefully ground in an agate mortar for 20 min, mixed well and transferred to an alumina crucible, which was then placed in a chamber furnace. The temperature was slowly increased to 1530 °C and held at this temperature for 4 h. After cooling to room temperature, the crucible was removed and the agglomerated product was reground into powder to obtain the target phosphor.

2.2 Characterization

The phase purity and crystal structure of the samples were determined by Bruker D8 Advance X-ray diffractometer (XRD) with Cu as the radiation source, 40 mA current and 40 kV voltage, and XRD data were collected in the range of 10–80°. The XRD data were Rietveld refined using the General Structural Analysis System software (GSAS-II). Schottky field

emission electron microscope (JSM-IT800) was used to obtain the micromorphology images and EDS mapping. The photoluminescence (PL) and photoluminescence excitation (PLE) spectra were recorded with an F-7100 spectrophotometer. The variable temperature spectra (303–483 K) were recorded on an Edinburgh FLS-1000 fluorescence spectrometer with an additional heating controller. Fluorescence lifetimes of Ce^{3+} were tested with a fluorescence spectrometer (Edinburgh FLS-1000) and an additional 450 nm laser was fitted as an excitation light source. The performance of the packaged LEDs was measured under a Starspec SSP6612 measuring system. Tests without special instructions were performed at room temperature.

3 Results and discussion

3.1 Phase identification and crystal structure analysis

Fig. 1a shows the XRD pattern of $YSAG:xCe^{3+}$. Compared with the standard card $Y_3Al_5O_{12}$ (PDF#88-2048), the diffraction peak experiences a shift towards smaller angles, which is caused by the introduction of Sc^{3+} with a larger ionic radius to replace Al^{3+} in the matrix (0.745 Å for Sc^{3+} and 0.535 Å for Al^{3+} when the coordination number is 6). The degree of shift of the main peak to a small angle becomes more pronounced with increasing Ce^{3+} doping concentration, indicating the incorporation of a higher amount of Ce^{3+} into the lattice of the matrix. The diffraction peaks are all able to match with the standard card, indicating that the synthesized samples are all pure phases. Fig. S1† shows the XRD patterns of the $YSAG:0.03Eu^{3+}$, $YAG:0.03Eu^{3+}$ and $YAG:0.08Ce^{3+}$ samples, all of which are in good agreement with the standard card and no impurity peaks appear. Fig. 1b shows the schematic crystal structure of $YSAG$, where Y^{3+} ions occupy the dodecahedral sites, the octahedral sites are shared by Sc^{3+} and Al^{3+} ions, and the tetrahedral sites are occupied entirely by Al^{3+} ions. The Ce^{3+} ions occupy the dodecahedral sites of Y^{3+} ions in the matrix lattice.

In order to further confirm the crystal structure and phase purity of the samples, structural refinement of the $YSAG:0.08Ce^{3+}$ and $YAG:0.08Ce^{3+}$ samples were carried out by GSAS-II software, and the refinement results are shown in Fig. 1c and d. The values of R_{wp} , R_p and χ^2 are relatively small, indicating that the obtained results are reliable. Both $YSAG:0.08Ce^{3+}$ and $YAG:0.08Ce^{3+}$ have the same garnet structure. And the obtained parameters are listed in Table 1. The introduction of Sc^{3+} with large ionic radius into the YAG matrix lattice results in an increase in the lattice constants (a , b , c), the lattice volume (V), and the Y–O bond length. Although Sc^{3+} is introduced solely at the octahedral site, the dodecahedron occupied by Ce^{3+} is co-prismatic with the octahedron, and the introduction of Sc^{3+} leads to the expansion of the dodecahedron.

Fig. 2 shows the scanning electron microscope (SEM) image of the phosphor, revealing grain sizes of approximately 4 μm . In order to provide additional insights into the composition and elemental distribution of the sample, elemental analysis



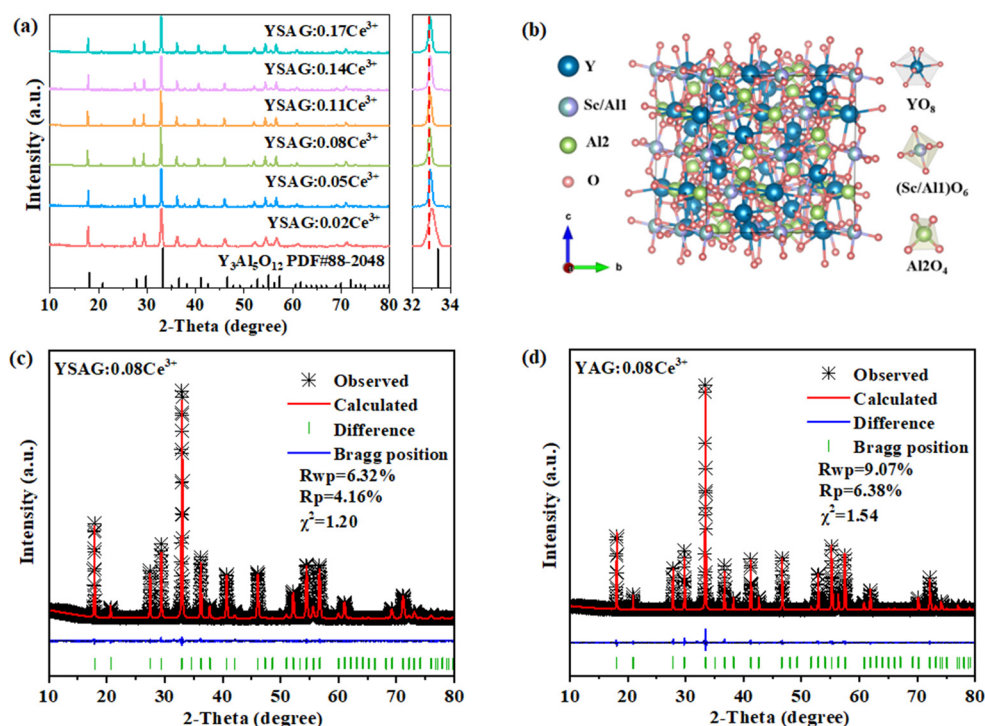


Fig. 1 (a) XRD patterns and magnified XRD patterns around 32°–34° of the YAG:*x*Ce³⁺. (b) Crystal structure diagram of the YAG. Rietveld refinement of (c) YAG:0.08Ce³⁺ and (d) YAG:0.08Ce³⁺.

Table 1 Crystallographic data for YAG:0.08Ce³⁺ and YAG:0.08Ce³⁺

Samples	$a = b = c$ (Å)	V (Å ³)	$\alpha = \beta = \gamma$ (°)	Y–O (Å)
YAG:0.08Ce ³⁺	12.00	1729.4	90	2.37
YAG:0.08Ce ³⁺	12.14	1791.8	90	2.41

was conducted on the sample. In the EDS element mapping model, different colors represent different elements, and all elements (Y, Sc, Al, O, and Ce) are uniformly distributed on

the particle, indicating that the phosphor YAG:0.08Ce³⁺ is well synthesized.

3.2 Photoluminescence properties and thermal Stability

Fig. 3a shows the photoluminescence excitation (PLE) and photoluminescence (PL) spectra of YAG:0.08Ce³⁺. The PLE spectrum contains three excitation bands at 254, 348 and 450 nm, respectively. The excitation band at 254 nm belongs to the absorption of the matrix, and the double bands at 348 and 450 nm are typical excitation bands of the Ce³⁺ ions, which

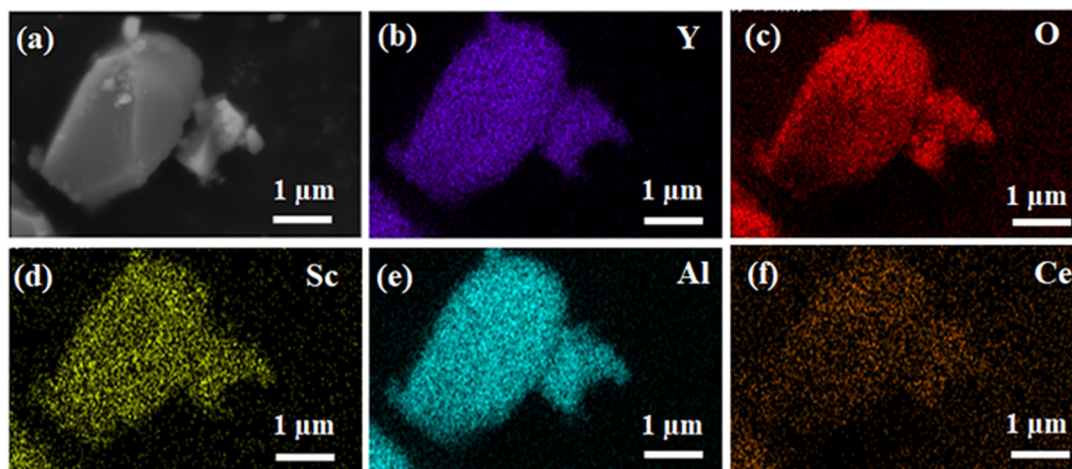


Fig. 2 (a) SEM image and (b-f) EDS elemental mapping of the YAG:0.08Ce³⁺.



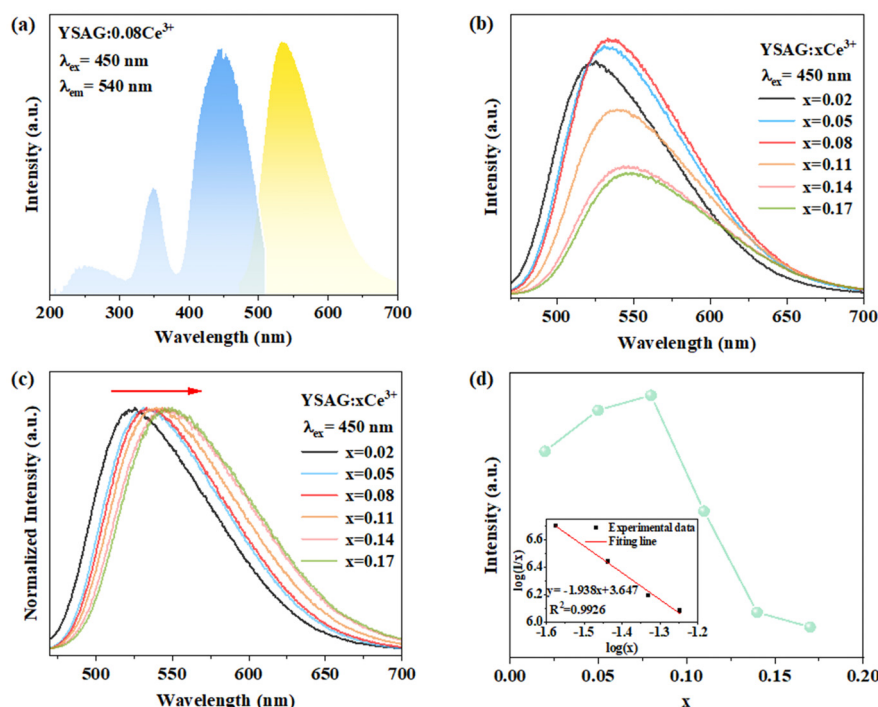


Fig. 3 (a) PLE and PL spectra of YSAAG:0.08Ce³⁺. (b) Variation of PL spectrum with Ce³⁺ doping concentration. (c) Normalized PL spectra. (d) The dependence of emission intensity on Ce³⁺ concentration.

correspond to the transition of the 4f electrons in the ground state to the 5d energy level of the Ce³⁺ ions. The PL spectrum of the Ce³⁺ ions is a broadband emission with a peak of 540 nm, due to the 5d to ²F_{5/2} and ²F_{7/2} transition of the Ce³⁺ ions. Fig. 3b shows the PL spectra of YSAAG:xCe³⁺. With an increase in doping Ce³⁺-concentration, the emission intensity exhibits a continuous increase, reaching its maximum value at $x = 0.08$. However, beyond this optimal doping concentration, the emission intensity gradually decreases due to concentration quenching. Since the distance between Ce³⁺ ions decreases as the doping concentration increases, this reduction in distance enhances the probability of nonradiative energy transfer between Ce³⁺ ions, consequently increasing the probability of electrons entering the quenching center. Fig. 3c shows the normalized PL spectra of YSAAG:0.08Ce³⁺. With an increase in Ce³⁺ doping concentration, the emission peak undergoes a red-shift from 524 to 549 nm. This shift is attributed to the increased probability of the energy transfer from the Ce³⁺ ions with the higher 5d energy levels to lower 5d energy levels of the Ce³⁺ ions. As a result, there is a decrease in emission intensity on the high-energy side of the spectrum.^{23,24}

The critical energy transfer distance (R_c), proposed by Blasse, is defined as the distance where the probability of energy transfer equals that of radiation emission by the activator, and can be estimated geometrically from the following equation:²⁵

$$R_c \approx 2 \left[\frac{3V}{4\pi x_c N} \right]^{\frac{1}{3}}$$

According to the refinement results, the cell volume V is 1791.81 Å³, the critical concentration x_c is 0.08, and the N is 8. The calculated R_c is 17.49 Å, which is significantly larger than 5.0 Å. This result reveals that the energy transfer mechanism among Ce³⁺ ions is governed by electric multipole interactions. The type of electric multipole interactions can be determined by the following equation:²⁶

$$\log\left(\frac{I}{x}\right) = -\frac{\theta}{3}\log(x) + A$$

where I is the PL intensity, x is the activator ion concentration, which is not less than the critical concentration, and θ is an indication of the type of electric multipolar interactions. $\theta = 6, 8$ and 10 correspond to dipole–dipole (d–d), dipole–quadrupole (d–q) and quadrupole–quadrupole (q–q) interactions, respectively. The linear relationship between $\log(I/x)$ and $\log(x)$ is given in the inset of Fig. 3d. On the basis of the fitting results, the slope of the fitted line yields a value of θ as 5.8, which is close to 6, indicating that dipole–dipole interactions dominate the energy transfer among Ce³⁺ ions.

In order to better evaluate the luminescent properties of the prepared phosphor, YAG:0.08Ce³⁺ was also prepared. The comparison of the emission spectra is shown in Fig. 4a. The emission peak position of YSAAG:0.08Ce³⁺ is blue-shifted by 10 nm, compared with YAG:0.08Ce³⁺. Additionally, there is a small difference in the full width half maximum with values of 91.7 nm and 91.4 nm for the YSAAG:0.08Ce³⁺ and YAG:0.08Ce³⁺ phosphors, respectively. The blue-shift of the emission spectrum can be attributed to the reduction of crystal field split-



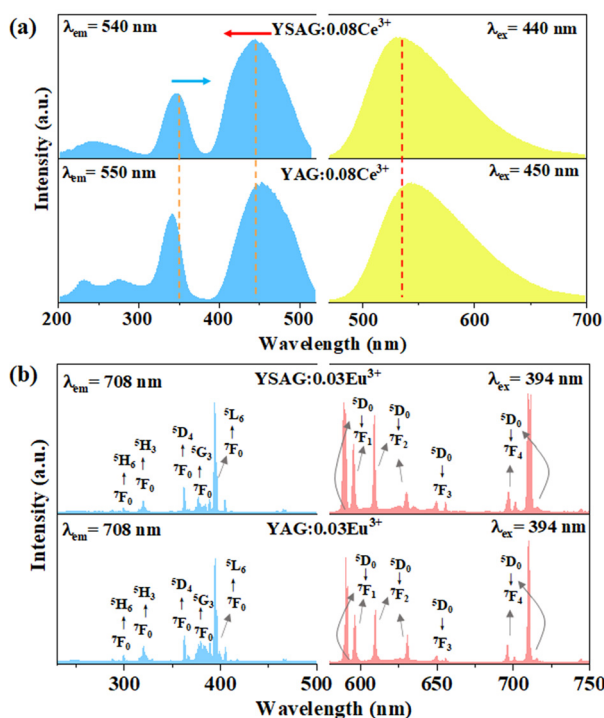


Fig. 4 (a) PLE and PL spectra of YAG:0.08Ce³⁺ and YAG:0.08Ce³⁺. (b) PLE and PL spectra of YAG:0.03Eu³⁺ and YAG:0.03Eu³⁺.

ting, which can be estimated using the following equation:²²

$$D_q = \frac{1}{6} Z e^2 \frac{r^4}{R^5}$$

where Z and e are the number of valence electrons and charge of the anion, respectively, and R and r are the average bond length and the radius of the d-wave function, respectively. In YAG and YAG, the ligands did not change, only the bond lengths increased (Table 1), which would lead to a decrease in D_q , causing a blue-shift in the emission spectra. The change in the excitation spectrum of YAG:Ce³⁺ further supports the aforementioned conclusion. As shown in Fig. 4a, the excitation peaks of the Ce³⁺ ions at 348 and 450 nm in YAG exhibit a red-shift and blue-shift, respectively, compared with YAG:Ce³⁺. This result demonstrates that the crystal field become weakens when Al³⁺ is replaced by Sc³⁺ ions.

The Eu³⁺ ion is commonly employed as spectroscopic probe for monitoring the occupancy of lattice sites.²⁸ In order to compare and study this phenomenon, YAG:0.03Eu³⁺ and YAG:0.03Eu³⁺ samples were synthesized. Fig. 4b shows their PLE and PL spectra. As presented, the excitation spectra of YAG:0.03Eu³⁺ and YAG:Eu³⁺ consist of a series of f–f transition peaks located at 299, 320, 361, 381 and 393 nm, which are assigned to the ${}^7F_0 \rightarrow {}^5H_6$, ${}^7F_0 \rightarrow {}^5H_3$, ${}^7F_0 \rightarrow {}^5D_4$, ${}^7F_0 \rightarrow {}^5G_3$ and ${}^7F_0 \rightarrow {}^5L_6$ transitions, respectively. Under 394 nm excitation, the emission spectrum consists of a series of peaks located at around 590, 596, 609, 630, 650, 696, and 708 nm, assigned to the ${}^5D_0 \rightarrow {}^7F_J$ ($J = 1-4$) transitions. It is worth noting that the ${}^5D_0 \rightarrow {}^7F_1$ transition of Eu³⁺ is a magnetic-

dipole transition, its intensity is largely independent of the surrounding chemical environment of the Eu³⁺ ion. On the other hand, the ${}^5D_0 \rightarrow {}^7F_2$ transition belongs to the hypersensitive electric-dipole transition, its intensity is greatly influenced by the lattice environment, and the higher the lattice symmetry the weaker the emission, and it does not emit light in the lattice with central symmetry. Therefore, the lattice symmetry can be detected by the intensity ratio of red and orange emission ($R/O = I({}^5D_0 \rightarrow {}^7F_2)/I({}^5D_0 \rightarrow {}^7F_1)$).^{29,30} Red emission at 613 nm can be observed in the PL spectra of both YAG:Eu³⁺ and YAG:Eu³⁺, indicating that Eu³⁺ occupies the lattice position without centrosymmetric. The red-orange ratios in YAG and YAG were calculated to be 68.9% and 70.6%, respectively. This indicates that the dodecahedral lattice symmetry is slightly reduced by the introduction of Sc³⁺ at the octahedral position. Therefore, the changes in the excitation and emission spectra of Ce³⁺ ions caused by the decrease in the symmetry of the dodecahedron are also small. In addition, Fig. 5 demonstrates the variable temperature spectrum of YAG:Eu³⁺, where the intensity of each emission peak of Eu³⁺ increases synchronously with decreasing temperature, suggesting that the emission of Eu³⁺ comes from the same dodecahedral lattice site and that there is no second dodecahedral lattice site. Therefore, the broadband emission of Ce³⁺ ions in this matrix is also from one dodecahedral lattice site.

Fig. 6a shows the fluorescence decay curves of YAG:xCe³⁺ samples measured by monitoring at 540 nm after 450 nm pulse excitation. The decay lifetime was able to be calculated according to the following equation:²⁷

$$\tau = \frac{\int_0^{\infty} t I(t) dt}{\int_0^{\infty} I(t) dt}$$

In the YAG:xCe³⁺ samples, the fluorescence lifetime decay curves of Ce³⁺ can be well fitted with the single exponential function, and the corresponding decay lifetimes were 56.32, 50.27, 49.34, 46.78, 43.51, and 41.74 ns when x was 0.02, 0.05, 0.08, 0.11, 0.14 and 0.17, respectively. Generally, the fluorescence lifetime of Ce³⁺ becomes shorter with an increase in doping concentration. This reduction in lifetime can be attrib-

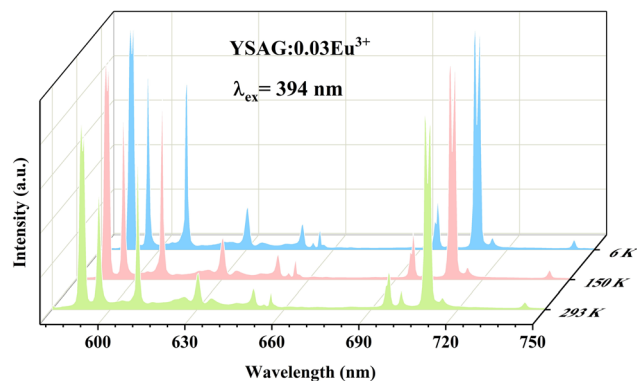


Fig. 5 Variable temperature spectra of YAG:0.03Eu³⁺.



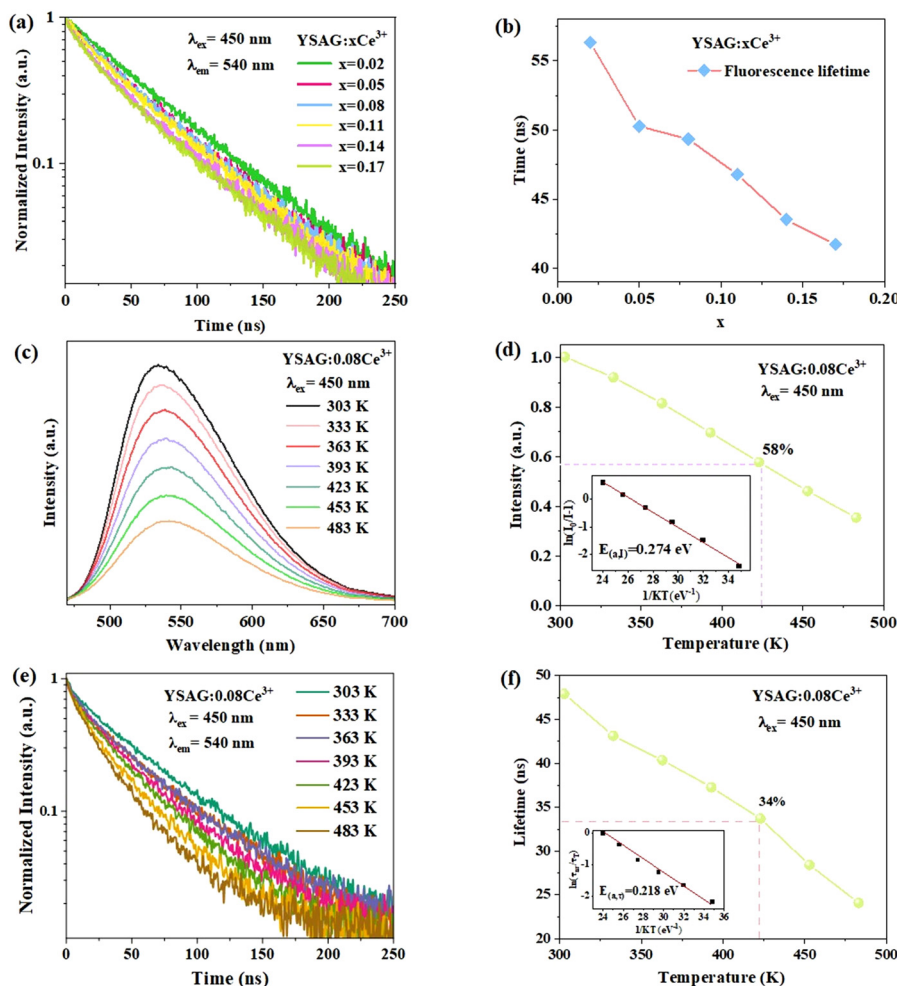


Fig. 6 (a) Fluorescence decay curves of the YAG:xCe³⁺. (b) Trends in the fluorescence lifetime of Ce³⁺. (c) PL spectra of YAG:0.08Ce³⁺ with increasing temperature. (d) The trend of temperature dependent emission intensity of YAG:0.08Ce³⁺. (e) Fluorescence lifetime of YAG:0.08Ce³⁺ with increasing temperature. (f) The trend of temperature dependent fluorescence lifetime of YAG:0.08Ce³⁺.

uted to an increase in the probability of non-radiation transition of Ce³⁺ to the quenching center.

Thermal stability is an important index for evaluating phosphors. In general, one can explain the thermal quenching behavior of Ce³⁺ by the conformational coordinate diagram in Fig. S2.† Under blue light excitation, the 4f electron of Ce³⁺ is excited to the 5d orbital, the electron in the excited state returns to the 4f ground state *via* process ① and simultaneously radiates yellow light. Additionally, excited state electrons can also cross the energy barrier (E_a) and return to the ground state *via* process ② without radiation. As the temperature increases, the probability of the excited electrons returning to the ground state *via* the process ② increases, leading to a decrease in the phosphor's emission intensity as well as a reduction in the radiation lifetime. The emission spectra of YAG:0.08Ce³⁺ sample were tested in the range of 303–483 K. As shown in Fig. 6c, its emission intensities decreased continuously with increasing temperature, and the emission intensity of YAG:0.08Ce³⁺ at 423 K retained 58% of its value

at 303 K. Thermal activation energy (E_a) can be utilized to assess the thermal stability of the luminescent material. A larger E_a indicates that electrons face greater difficulty in reaching the energy level intersection, resulting in a smaller probability of returning to the ground state *via* non-radiative transitions. The value of $E_{(a,T)}$ can be estimated using the Arrhenius equation, which is expressed as follows:³¹

$$I_T = \frac{I_0}{1 + Ae^{(-E_{(a,T)}/k_B T)}}$$

where I_0 is the initial PL intensity of the phosphor at room temperature, I_T is the PL intensity at a given temperature T , K is the Boltzmann constant with a value of 8.617×10^{-5} eV K⁻¹, and A is a constant. The emission spectra of YAG:0.08Ce³⁺ sample in the wavelength range of 470–700 nm at different temperatures were integrated as the emission intensity, and $E_{(a,T)}$ was determined by fitting the data to the equation, as shown in Fig. 6d, the activation energy of YAG:0.08Ce³⁺ is 0.27 eV.



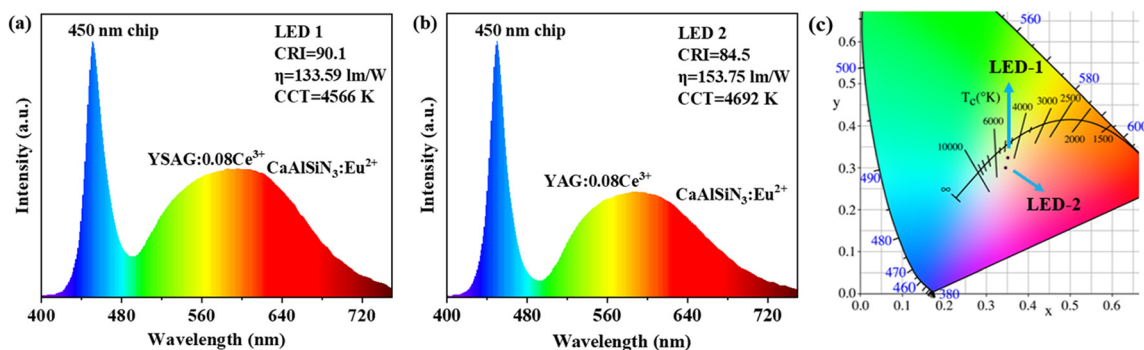


Fig. 7 (a) EL spectra of WLED devices with YSAG:0.08Ce³⁺. (b) EL spectra of WLED devices with YAG:0.08Ce³⁺. (c) CIE chromaticity diagram of YSAG:0.08Ce³⁺ and YAG:0.08Ce³⁺.

Alternatively, the activation energy can be calculated from the variable temperature lifetime, the fluorescence lifetime of YSAG:0.08Ce³⁺ was tested as a function of temperature, and it is seen in Fig. 6e that as the temperature increases from 303 to 483 K the fluorescence lifetime decreases continuously, and their values are 47.87, 43.08, 40.31, 37.25, 33.69, 28.42, and 24.10 ns, in that order. Plotting the decay of the fluorescence lifetime with temperature based on the Arrhenius equation below:³²

$$\tau(T) = \frac{\tau_r}{1 + (\tau_r/\tau_{nr}) \exp(-E_{(a,\tau)}/KT)}$$

where τ_{nr} is the non-radiative times, τ_r is the radiative times. K is the well-known Boltzmann constant. The value of the thermal activation energy $E_{(a,\tau)}$ is 0.22 eV.

3.3 Applications in pc-WLEDs

To demonstrate the application of the prepared phosphors, WLEDs were prepared using YSAG:0.08Ce³⁺ and YAG:0.08Ce³⁺ combined with commercial red CaAlSiN₃:Eu²⁺ phosphor and 450 nm blue chips, respectively. Their electroluminescence (EL) spectra are presented in Fig. 7a and b, where the CRI, CCT and the luminous efficiency (η) were also given. It is evident that the WLED prepared with YSAG:0.08Ce³⁺ has a higher CRI (90.1) when the value of CCT is similar with the WLED prepared with YAG:0.08Ce³⁺. This improvement in CRI is attributed to the effective blue-shift of the emission peak of YSAG:0.08Ce³⁺, which helps to fill the previously missing blue-green region of the spectrum. Fig. 6c shows their CIE chromaticity diagrams, the color coordinates of the WLED devices are (0.352, 0.324) and (0.346, 0.301), respectively.

4 Conclusions

In conclusion, YSAG:*x*Ce³⁺ phosphors were synthesized by a high-temperature solid-state reaction. The introduction of Sc³⁺ reduces the field splitting of the 5d energy level of Ce³⁺, resulting in a 10 nm blue-shift of the YSAG:Ce³⁺ emission spectrum compared with that of YAG:Ce³⁺. The blue-shift of the spectrum efficiently fills the blue-green region, WLED was fabri-

cated by combining YSAG:0.08Ce³⁺, CaAlSiN₃:Eu²⁺ with 450 nm blue chips, exhibiting high color rendering index (90.1), low color temperature (4566 K) and high luminous efficiency (133.59 lm W⁻¹). These results indicate that YSAG:0.08Ce³⁺ has potential application in higher color rendering index WLEDs.

Author contributions

Data curation, Hanwei Zhao; Formal analysis, Dashuai Sun, and Zeyu Lyu; Investigation, Lixuan Wang, Zheng Lu, and Jianhui Wang; Methodology, Sida Shen, and Luhui Zhou; Writing – original draft, Hanwei Zhao; Writing – review & editing, Dashuai Sun, Jinhua He, and Hongpeng You. All authors have read and agree to the published version of the manuscript.

Conflicts of interest

There are no conflicts to declare.

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