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Tuning the optical and magnetic properties of lanthanide single-ion magnets using nitro-functionalized trispyrazolylborate ligands†

Christopher Hossack, Folasade Abdul, Christopher Cahill And Claire Besson

We report the synthesis, crystal structures, photophysical and magnetic properties of 11 novel lanthanide complexes with the asymmetrically functionalized trispyrazolylborate ligand 4-nitrotrispyrazolylborate, $4-NO_2Tp^-$: $[Ln(4-NO_2Tp)_3]$ (Ln = La-Dy, except Pm). In-depth photophysical characterization of the ligands *via* luminescence, reflectance and absorption spectroscopic techniques, decay lifetimes, quantum yields supported by time-dependent density functional theory (TD-DFT) and natural bond order (NBO) analysis reveal that $n-NO_2Tp^-$ ligands are dominated by intra-ligand charge transfer (ILCT) transitions and that second-sphere interactions are critical to the stabilization of the T_1 state of $n-NO_2Tp^-$ ligands and hence their ability to sensitize Ln^{3+} emission. The luminescence properties of the complexes indicate that $4-NO_2Tp^-$ is a poor sensitizer of Ln^{3+} emission, unlike $3-NO_2Tp^-$. Moreover, $[Nd(4-NO_2Tp)_3]$ (crystallized as a hexane solvate) displays single-molecule magnet (SMM) properties, with longer relaxation times and larger barrier than the non-functionalized $[NdTp_3]$, attributed to the addition of the NO_2 -group and subsequent rigidification of the molecular structure.

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Introduction

Lanthanide (Ln) hybrid materials with tunable optical and magnetic properties have been attractive targets in fields such as opto-electronics and molecular magnetism. Hills there has been great success with tailoring the optical properties of the lanthanides towards specific applications such as metal ion and explosives sensors, hermometers, information storage devices etc., the ability to tune the molecular magnetic properties of the lanthanide ions remains quite elusive. The intrinsic, magnetic properties of 4f ions, namely, their large single-ion anisotropy and high degree of spin-orbit coupling allow for Ln-materials to have applications such as high density, molecule-based information storage, spintronic devices and quantum computing. Recent reports on lanthanide molecular magnets that seek to harness the large single-ion anisotropy of the lanthanides, do so by engineering a wide

variety of chemical environments including radical-bridged complexes to promote superexchange interactions, $^{10-15}$ low coordination number (CN < 8) complexes, $^{16-19}$ and many dicyclopentadienyl (metallocene) complexes. $^{20-22}$

There has, however, been comparatively very little work done on optimizing both the optical and magnetic properties of the lanthanides simultaneously, or using the lessons learned from harnessing the optical properties to inform design of lanthanide molecular magnets.²³ Bi-functional Lnmaterials that display useful photophysical and magnetic properties are advantageous because they could display unusual properties, *e.g.* optical polarization of nuclear spin,²⁴ and yield insights into the fundamental properties of f-electrons by studying excited state dynamics in a magnetic field. Such work would require a ligand platform that can: (i) exert a large degree of control over the primary coordination sphere of the lanthanides, (ii) sensitize lanthanide luminescence and (iii) accommodate a wide variety of ring substituents to facilitate tunability of properties.²³

Pyrazolylborate ligands ($H_{4-x}BPz_x$) meet all these criteria. They are widely used to make lanthanide chelates owing to their ability to control the spherical coordination environment of f-block elements via steric bulk/chelate effect, ^{25,26} and can accommodate a very diverse array of ring substituents. ^{27–29} Moreover, lanthanide pyrazolylborates, $[Ln(H_{4-x}BPz_x)_y]$, have been characterized as light-emitting devices (LEDs), ^{30,31} single-

Department of Chemistry, The George Washington University, 800 22nd Street, NW, Washington, D.C. 20052, USA. E-mail: chossack@gwu.edu, cahill@gwu.edu, claire besson@gwu.edu

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molecule magnets (SMMs)^{32,33} and even magneto-luminescent materials that display both sensitized luminescence and SMM properties.³⁴ Despite the breadth of work done on harnessing the desirable, opto-electronic and magnetic properties of the f-elements using pyrazolylborate ligands, our understanding of how ring substituents affect the properties of $[Ln(H_{4-x}BPz_x)_y]$ materials is lacking, especially when compared to Ln-benzoate or Ln-β-diketonate systems.³⁵

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Latva et al., 36 Tsaryuk et al., 37,38 de Bettencourt-Dias et al.,39,40 Raymond et al.41 among others, have detailed descriptions on how to tune the optical properties of Ln hybrid materials using aromatic ring substituents to control the antenna effect. The antenna effect describes the energy transfer from the S₁ or T₁ states of coordinated "antenna" ligands (organic chromophores) to lanthanide ions' ^{2S+1}L_I states by the Förster or Dexter energy transfer mechanisms. The influence of both electron-donating and electron-withdrawing groups, as well as the impact of the number of substituents and ring position on the photophysical properties of lanthanide complexes is well understood and has been implemented with great success.³⁵ Conversely, less explored is the influence of said substituent effects on slow magnetic relaxation dynamics (SMM properties), considering both structural changes (intermolecular interactions) and perturbations of the crystal field experience by the Ln³⁺ ion. We are aware of only two reports, by Sushila et al. and Gálico et al, on the influence of halo-substitution on the optical and magnetic properties of lanthanide amino-bisphenol and β-diketonate complexes respectively, that addresses the impact of substituent (halogen) effects on optical and magnetic properties of Ln-materials. 42,43 Given the variety of lanthanide molecular magnets with pyrazolylborates and our lab's recent work with selective functionalization of trispyrazolylborate (Tp⁻) ligands, 29,44 we postulate that asymmetrically functionalized Tp ligands (Fig. 1) are an ideal platform for establishing how aromatic ring substituents influence the properties of lanthanide molecular magnets.

We previously reported five families of lanthanide 3-nitrotripyrazolylborates, $Ln(3-NO_2Tp)_x$ complexes that displayed tunable nuclearity and charge transfer mediated optical properties, and we elucidated the impact of strongly electron-with-

Fig. 1 Diagrams of asymmetric trispyrazolylborate ligands, left - 3-nitrotrispyrazolylborate (3-NO₂Tp $^-$) and right - 4-nitrotrispyrazolylborate (4-NO₂Tp $^-$).

drawing substituents on the coordination chemistry and photophysical properties of $Ln(H_{4-x}Pz_x)_y$ complexes.⁴⁵ We noted that the drastic difference in structural types and photophysical properties, when compared to other homoleptic Ln $(H_{4-x}BPz_x)_y$ complexes was primarily due to the presence of new binding site via the oxygen of the nitro-group (Ln-O_{NO₂}) on the 3-NO₂Tp⁻. With this in mind, we substituted 3-NO₂Tp⁻ for the 4-nitrotrispyrazolylborate ligand, 4-NO₂Tp⁻, to investigate the influence of the nitro-group on lanthanide pyrazolylborates without Ln-O_{NO}, coordination. Herein, we report that the coordination compounds of 4-NO₂Tp⁻ with Ln³⁺ ions are indeed 9-coordinate homoleptic complexes, [Ln(4-NO₂Tp)₃] (Ln = La-Dy, except Pm) analogous to the previously reported [LnTp₃]. Moreover, the inclusion of the non-coordinating nitro-group imparts control over the crystal packing and crystallographic symmetry of [Ln(4-NO₂Tp)₃] which may be used as a tool to influence the slow magnetic relaxation of [Nd(4-NO₂Tp)₃]. The photophysics of [Ln(4-NO₂Tp)₃] feature Ln3+-selective singlet sensitization of Eu3+ and Tb3+ and various intra-ligand charge transfer (ILCT) transitions between the S_n and T_n states of the boron-bridged pyrazole and 4-nitropyrazole rings. We also provide a comparative analysis of the performance of Tp⁻ vs. n-NO₂Tp⁻ (n = 3 or 4) as Ln³⁺ sensitizers and SMM-generating ligands (Tp vs. 4-NO₂Tp).

Results

Synthesis of Ln(4-NO₂Tp)₃·solvent

The reported complexes, $[Ln(4-NO_2Tp)_3]$ (Ln = La-Dy, except Pm, 1-9) were synthesized by combining methanolic solutions of LnCl₃ salts with [TBA][4-NO₂Tp]. The solid crude product may be readily isolated by removing the methanol in vacuo and washing away tert-butylammonium chloride in isopropanol. The complexes are stable in acetonitrile (as indicated by their well-defined NMR spectra, see Fig. S7-S14†). They are also soluble in common solvents such as dicholoromethane and acetone, but only sparingly soluble in methanol and ethanol and insoluble in benzene and isopropanol. This differs drastically from the non-functionalized [LnTp3] complexes which are insoluble in all solvents when Ln = La-Gd and only sparingly soluble in benzene and dicholoromethane for Ln = Dy and Tb. 46 Crude [Ln(4-NO₂Tp)₃] can be recrystallized as one of two polymorphs, a hexane solvate, [Ln(4-NO₂Tp)₃]·C₆H₁₄ (a) and a benzene solvate, $[Ln(4-NO_2Tp)_3]\cdot 7C_6H_6$ (b).

Crystal structures

Structural description of [Ln(4-NO₂Tp)₃]-C₆H₁₄. Compounds **1a–9a** are all isomorphous and hence only the crystal structure of **4a** will be discussed in detail. **4a** contains one crystallographically unique Nd^{3+} ion coordinated to 9 nitrogen atoms from three unique 4-NO₂Tp⁻ ligands, with the Nd^{3+} center exhibiting a tricapped trigonal prismatic coordination geometry of approximately D_{3h} site symmetry. Ln–N bonds lengths range from 2.565(2) to 2.579(3) Å for the nitrogen atoms of the trigonal prism and 2.739(3) to 2.753(2) Å for the capping nitro-

gen atoms. The mononuclear [Ln(4-NO₂Tp)₃] complex displays an approximate C_{3h} point symmetry, analogous to that of the 9-coordinate [LnTp₃] complexes, 46 with a C_3 axis passing through the Nd³⁺ center perpendicular to the plane defined by the three nitropyrazoles. However, unlike in [LnTp₃], there are three unique NO₂···H-C hydrogen bonds between adjacent [Ln(4-NO₂Tp)₃] molecules, whereas [LnTp₃] does not feature any significant intermolecular interactions. The NO2···H-C hydrogen bonds as shown Fig. 2(b) range from 3.1916(1) Å to 3.4382(1) Å, where two of the three H bonds occur between pyrazole and 4-nitropyrazole rings and the last between two 4-nitropyrazole rings. Compound 4a crystallizes in the C2/c space group and it contains two highly disordered hexane sites in the lattice. The hexane molecule does not significantly separate the [Ln(4-NO₂Tp)₃] molecules, resulting in fairly dense crystal packing in 4a with the shortest intermolecular Nd³⁺-to-Nd³⁺ distance being 9.471(6) Å, similar to [LnTp₃] at 9.686(1) Å. Interestingly this packing features sheets of [Ln(4-NO₂Tp)₃] molecules, along the [001] direction, such that a given sheet is sandwiched between a parallel one and another that is slighted canted by 19° (Fig. S1†).

Structural description of $[Ln(4-NO_2Tp)_3]$ -7C₆H₆ (5b, Ln = Sm and 8b, Ln = Tb). Compounds 5b and 8b constitute a second solvate polymorph of $[Ln(4-NO_2Tp)_3]$, crystalized with 7 solvent benzene molecules per molecule of $[Ln(4-NO_2Tp)_3]$. Since 5b and 8b are isostructural and only the structure of 5b will be discussed. The $[Ln(4-NO_2Tp)_3]$ complex of 5b (Ln = Sm) is the same as 4a (Ln = Nd) (Fig. S2a†), with Sm–N bond lengths that range from 2.5427(11)–2.5440(13) Å for the trigonal prism nitrogrens and 2.7408(11) Å for the capping nitrogens, however it crystallizes in the $R\bar{3}$ space group. Additionally, the crystallographic 3-fold rotoinversion axis overlaps with the

molecular C_3 that through Sm^{3+} , resulting in only a third of the $[\mathrm{Ln}(4\text{-NO}_2\mathrm{Tp})_3]$ complex being present in the asymmetric unit of **5b**, while the entire molecule is contained with the asymmetric unit of **4a**. There are 7 lattice benzene molecules per molecule of $[\mathrm{Ln}(4\text{-NO}_2\mathrm{Tp})_3]$, resulting in significantly less dense crystal packing of the $[\mathrm{Ln}(4\text{-NO}_2\mathrm{Tp})_3]$ units (Fig. S2b†), with fairly long intermolecular Sm^{3+} -to- Sm^{3+} distances, the shortest being 12.667(1) Å. All of the $[\mathrm{Sm}(4\text{-NO}_2\mathrm{Tp})_3]$ complexes lie perfectly parallel to each other and there are no significant non-covalent interactions between the monomers (Fig. 3).

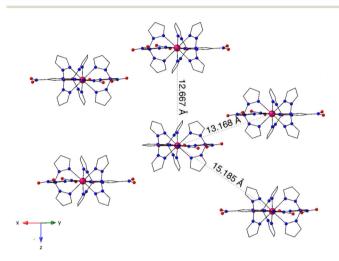


Fig. 3 Packing diagram of 5b (solvent benzene molecules omitted for clarity) where the $[Sm(4-NO_2Tp)_3]$ molecules align collinearly in parallel planes.

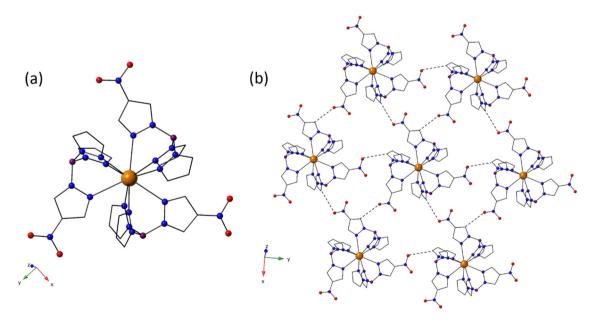


Fig. 2 (a) Asymmetric unit of 4a (with the hexane omitted for clarity) featuring the $[Ln(4-NO_2Tp)_3]$ complex. (b) $NO_2\cdots H-C$ hydrogen bonding between $[Ln(4-NO_2Tp)_3]$ complexes forming a 2-D layer in the xy plane.

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Photophysical properties of 4-NO₂Tp⁻ and [Ln(4- $NO_2Tp)_3$ - C_6H_{14}

Optical properties of 4-NO₂Tp⁻. We utilized diffuse reflectance and luminescence spectroscopy techniques, in addition to time-dependent density functional theory (TD-DFT) to fully characterize the photophysical properties of 4-NO₂Tp⁻. Diffuse reflectance measurements for K[4-NO₂Tp] (Fig. S16†) reveal a broad absorption band, spanning the UV to blue region,

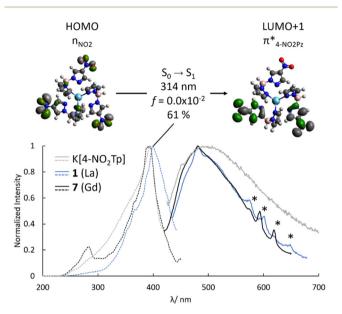


Fig. 4 Top: Calculated $S_0 \rightarrow S_1$ transition of Calc-1 as predicted by TD-DFT. Bottom: Room temperature luminescence excitation (left, broken lines) and emission (right, solid lines) of K[4-NO₂Tp] (λ_{exc} = 393 nm, λ_{em} = 490 nm), **1a** (λ_{exc} = 400 nm, λ_{em} = 481 nm) and **7a** (at 77 K, $\lambda_{\rm exc}$ = 394 nm, $\lambda_{\rm em}$ = 481 nm). The excitation peak at 400 nm is assigned as $S_0 \rightarrow S_1$ of 4-NO₂Tp⁻. * – background signal from the Xe arc lamp observable owing to very weak emission intensity of 1a and 7a.

between 225 and 430 nm. The solid-state luminescence spectra of K[4-NO₂Tp], 1a (fluorescence) and 7a (phosphorescence) spectra respectively, for 4-NO₂Tp⁻, are summarized in Fig. 4. La and Gd complexes are well known to yield emission spectra that display decay of S_n and T_n states almost exclusively 45,47,48 and are included to aid in characterization of those transitions/states with respect to 4-NO₂Tp⁻. The similar profiles for all three (emission) spectra, even 7a at 77 K, indicate that the decay of the T_n states of 4-NO₂Tp⁻ is non-radiative (NR) and/or that T₁ state is unstable and excitation energy is readily lost. The most probable cause would be the NO2-group intra-ligand charge transfer (ICLT) quenching pathway, described by Tsaryuk et al., where excitation energy is lost due to the non-radiative π^* - n_{NO_2} transition. 37,38

To assign the $S_0 \rightarrow S_1$ and $T_1 \rightarrow S_0$ transitions of 4-NO₂Tp⁻ associated with the antenna effect, we carried out TD-DFT calculations on a geometry optimized model of 1, Calc-1 (Fig. 4 and Fig. S26†). The population analysis of Calc-1 indicates that the HOMOs and LUMOs of the ligand are localized at the π systems of aromatic pyrazole (Pz) and 4-nitropyrazole (4-NO₂Pz) rings. There is good agreement between the calculated absorption spectrum of Calc-1 and the solution absorption spectrum of crude 1 in acetonitrile (Fig. 5). However, the low energy absorptions in the excitation spectra of K[4-NO₂Tp] and 1a at ~ 400 nm, that lead to radiative decay (ca. 485 nm) are not predicted by TD-DFT nor are they present in the solution absorption spectrum. We attribute this to the nature of collecting luminescence excitation spectra where only transitions that lead to emission can be observed and not necessarily the transitions with large ε values, which is what is typically observed in an absorption spectrum. Since reflectance spectroscopy is sensitive to transitions with very low molar absorptivity/f-oscillator strengths (see Ln3+ absorptions in Fig. S19-S25†), we note that the 400 nm transition of 4-NO₂Tp⁻ (which has a very low ε); is present in the reflectance spectrum of 1a. Given that the higher energy transitions with

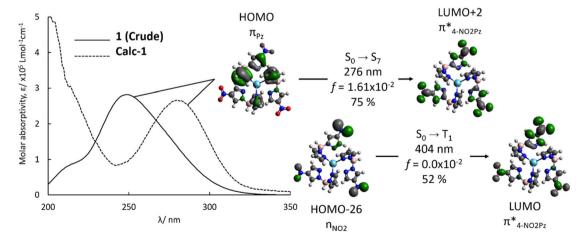


Fig. 5 Left – Electronic absorption spectrum (solid line) of crude 1 in acetonitrile (1.03 × 10⁻⁵ M) and calculated UV-VIS of Calc-1 in the gas phase. Right – TD-DFT output of the assigned $S_0 \to S_7$ as an ILCT $\pi_{Pz} \to \pi^*_{4-NO_9Pz}$ transition and the calculated $S_0 \to T_1$ transition associated with the

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larger oscillator strengths do not lead to emission from 4-NO₂Tp⁻, assignment of the peaks in the excitation spectra is quite difficult. However, based on the TD-DFT calculations (Fig. 5) and our previous work with the 3-NO₂Tp⁻ ligand and we assign the 4-NO₂Tp⁻ transitions as:

excitation peak at 400 nm: $S_0 \rightarrow S_1$, $n_{NO_2} \rightarrow \pi^*_{4-NO_2Pz}$ ICLT, emission at 485 nm: $S_1 \rightarrow S_0$, $\pi^*_{4\text{-NO,Pz}} \rightarrow n_{NO}$, ICLT, theoretical emission of $T_1 \rightarrow S_0$: $\pi^*_{4\text{-NO},Pz} \rightarrow n_{NO}$, ILCT.

Luminescence properties of [Ln(4-NO₂Tp)₃]·C₆H₁₄. All of the reported complexes (except 2a and 7a) display characteristic Ln³⁺ emission in the UV-visible region (Fig. 6). The profiles and relative intensities for 3a (Pr), 4a (Nd), 5a (Sm), 8a (Tb) and 9a (Dy) emission spectra are typical for the line-like, f-f transitions of the trivalent lanthanides. 34,45 The europium analogue, 6a, features an emission spectrum with the expected $^5\mathrm{D_0} \to {}^7\mathrm{F}_I \, (J=1\text{--}4)$ transitions. Induced electric dipole ($^5\mathrm{D_0} \to$ $^{7}F_{2}$, $^{5}D_{0} \rightarrow ^{7}F_{4}$) transitions are allowed in **6a**, given the lack of a center of symmetry, while fully forbidden transitions, ⁵D₀ \rightarrow $^{7}\text{F}_{0}$ and $^{5}\text{D}_{0} \rightarrow ^{7}\text{F}_{3}$ are either not observable ($^{5}\text{D}_{0} \rightarrow ^{7}\text{F}_{0}$) or very weak (${}^5D_0 \rightarrow {}^7F_3$). The intensity of magnetic dipole (${}^5D_0 \rightarrow {}^7F_1$) transition is independent on any Eu³⁺-ligand interactions, ⁴⁹ and just happens to coincide with ${}^5D_0 \rightarrow {}^7F_2$. The difference in intensity between J = 2 and J = 4 is somewhat unexpected, as the former transition is usually considered hypersensitive and thus expected to be more intense. 49 This can be attributed to two factors: (i) lack of electric quadrupole contributions to I = 2 and (ii) a static coupling mechanism between Eu³⁺ and 4-NO₂Tp⁻, *i.e.* a weakly covalent Eu³⁺-4-NO₂Tp⁻ bond.^{50,51} A preliminary Judd-Ofelt analysis of the luminescence data performed on 6a (Table S2†) provides some evidence for the latter with a low value for Ω_2 .⁵⁰ Finally, the crystal field splitting observed for the ${}^{5}\mathrm{D}_{0} \rightarrow {}^{7}\mathrm{F}_{I} (J = 1 - 4)$ transitions is in agreement with the high local symmetry (D_{3h}) of the europium center.⁵²

The observed emission spectra in Fig. 6 is due primarily to direct f-f transitions and not sensitization via the 4-NO₂Tp⁻ ligand. The excitation spectra of 1a-9a (except 2a and 7a) and the weak (or in some cases absent) absorption associated with 4-NO₂Tp⁻, highlights that 4-NO₂Tp⁻ is an inefficient sensitizer for Ln³⁺ emission. Interestingly, the presence of several direct f-f absorption peaks in Fig. 6 reveal that 4-NO₂Tp⁻ does not directly quench Ln³⁺ luminescence via Ln³⁺-to-4-NO₂Tp⁻ backenergy transfer and that direct f-f absorption and/or Ln3+-to-Ln³⁺ energy transfer (energy migration) are fairly efficient sensitization mechanisms. 2 likely does not display room temperature luminescence due to a fast $d_{Ce} \rightarrow \pi^*_{\text{4-NO}_2Pz}$ metalto-ligand charge transfer (MLCT) and subsequent nonradiative $\pi^*_{\text{4-NO}_2\text{Pz}} \rightarrow n_{\text{NO}_2}$ transition, analogous to monomeric, [Ce(3-NO₂Tp)₂(NO₃)].⁴⁵

6a and 8a are the only compounds that feature a significant absorption band ca. 330-340 nm that could be attributed to ligand absorption and sensitization via 4-NO₂Tp⁻. This band coincides somewhat with the absorption at 284 nm in the absorption spectrum of crude 1 and a $\pi_{Pz} \to \pi^*_{\text{4-NO},Pz}$ transition predicted at 276 nm for Calc-1 and therefore, this peak is assigned as the $S_0 \to S_7$, $\pi_{Pz} \to \pi^*_{4\text{-NO},Pz}$ ICLT transition. The disparity between the transitions observed in the absorption

spectrum versus the excitation and emission spectra may be attributed to two distinct pathways associated with luminescence vs. energy transfer to Eu³⁺/Tb³⁺; usually these pathways are identical (to be discussed later, see Fig. 10). The ability of 4-NO₂Tp⁻ to weakly sensitize Tb³⁺ and Eu³⁺ emission suggests that the T₁ state of 4-NO₂Tp⁻ (which we could not determine experimentally), may have enough energy to sensitize Eu^{3+} (${}^{5}D_{0}$, 17 500 cm⁻¹) and Tb (${}^{5}D_{4}$, 21 500 cm⁻¹) emission, yet may be too short lived to participate in energy transfer for the lanthanide ions studied, and that the rate of non-radiative decay via a fast $\pi^*_{4\text{-NO}_2Pz} \rightarrow n_{NO_2}$ transition is large $(k_{\text{nonrad}} \gg k_{\text{rad}})$. Another explanation for the sensitized emission in 6a and 8a and not the other lanthanides, may be a singlet energy transfer pathway^{53,54} via the S₀ \rightarrow S₇, $\pi_{Pz} \rightarrow \pi^*_{\text{4-NO,Pz}}$ ICLT transition that typically only Eu³⁺ and Tb³⁺ can participate in, whereas triplet sensitization is usually necessary for more weakly luminescent lanthanides $(Pr^{3+}, Sm^{3+}, Dy^{3+}).$

Magnetic properties of [Ln(4-NO₂Tp)₃]·C₆H₁₄

Static magnetic measurements. The magnetic susceptibility of 4a, 8a and 9a were studied under a static field of 1000 Oe (0.1 T) as shown in Fig. 7. The room temperature $\chi_{\rm m}T$ values were 1.65 cm³ K mol⁻¹, 11.85 cm³ K mol⁻¹ and 13.93 cm³ K mol⁻¹ for 4a, 8a and 9a respectively. These values agree well with those anticipated for free uncoupled Ln³⁺ ions at 1.64 cm³ K mol⁻¹ (Ln = Nd, ${}^{4}I_{9/2}$, S = 3/2, L = 6, g = 8/11), 11.82 cm³ K mol⁻¹ (Ln = Tb, ${}^{7}F_{6}$, S = 3, L = 3, g = 3/2) and 14.17 cm³ K mol⁻¹ (Ln = Dy, ${}^{6}H_{15/2}$, S = 5/2, L = 5, g = 4/3). For all three compounds, $\chi = T$ decreases with T, owing to the depopulation of higher energy m_I states generated by crystal field splitting of the ground state. 34 Assuming a perfect D_{3h} coordination sphere for the lanthanide center, the crystal field Hamiltonian can be expressed as:

$$\hat{H}_{CF} = B_2^0 \theta_2 \hat{O}_2^0 + B_4^0 \theta_4 \hat{O}_4^0 + B_6^0 \theta_6 \hat{O}_6^0 + B_6^6 \theta_6 \hat{O}_6^6$$

where B_{k}^{q} are the crystal field parameters (including orbital reduction parameters), θ_k , are the operator equivalent factors and \hat{O}_{i}^{q} the Stevens operator equivalents. A survey of the crystal field parameters between -5000 and +5000 cm⁻¹ indicated that the best fits of the experimental data around $B_6^0 = B_6^6 = 0$ for all three cations. While the negligible value of the sole equatorial term B₆ is in agreement with the substantially longer Ln—N_{NO,Pz} distances in the plane of the tricapped trigonal prismatic coordination sphere, the B_6^0 value is more surprising for 4f centers. Attempts to determine values for the B_2^0 and B_4^0 parameters were unsuccessful owing to overparameterization of the systems (Fig. S27-S29†). Even when further narrowing the possibilities by assuming a linear dependence of the crystal field parameters on the f-electron count of the metal center, two sets of parameters remained equally suited to account for the DC magnetization data (Table 1, Fig. 7 and S30†).

Dynamic magnetic measurements. Dynamic (ac) magnetic susceptibility experiments were collected on 4a, 8a and 9a in

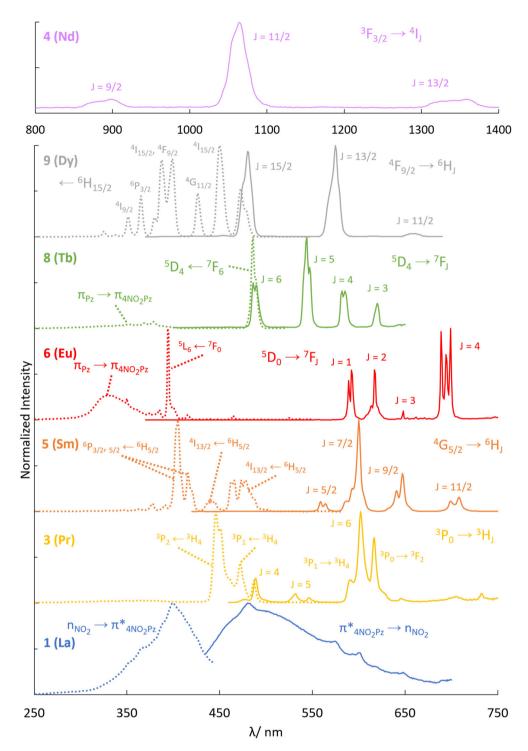


Fig. 6 298 K temperature solid-state luminescence excitation (left, broken lines) and emission (right, solid lines) of 1a ($\lambda_{\rm exc}=400$ nm, $\lambda_{\rm em}=481$ nm), 3a ($\lambda_{\rm exc}=446$ nm, $\lambda_{\rm em}=602$ nm), 4a ($\lambda_{\rm exc}=350$ nm), 5a ($\lambda_{\rm exc}=405$ nm, $\lambda_{\rm em}=600$ nm), 6a ($\lambda_{\rm exc}=394$ nm, $\lambda_{\rm em}=700$ nm), 8a ($\lambda_{\rm exc}=486$ nm, $\lambda_{\rm em}=544$ nm) and 9a ($\lambda_{\rm exc}=450$ nm, $\lambda_{\rm em}=575$ nm) with assignments of the various $4-NO_2Tp^-$ or direct f-to-f Ln³⁺ transitions. The emission spectrum of 4a collected using 350 nm may be a $4-NO_2Tp^-$ transition or Nd^{3+} absorption at ~350 nm ($^2I_{11/2}$). 55

the 0–1000 Hz frequency range between 1.8 and 9.8 K. None of the complexes displayed an out-of-phase signal in the absence of a magnetic field bias, but the Nd³⁺ complex, 4a did so, under a small applied field (H = 400 Oe). $\chi'_{\text{m}} vs. f$ and $\chi''_{\text{m}} vs. f$

isotherms are presented in Fig. 8(a) and (b) respectively and the frequency dependence of $\chi''_{\rm m}$ confirm that **4a** is indeed a single-molecule magnet. The time dependence of the relaxation times, extracted using the CCFit2 program, ⁵⁸ can be

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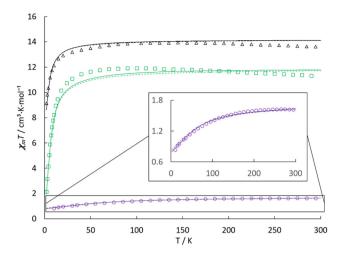


Fig. 7 Temperature dependence of $\chi_m T$ product at 1000 Oe for 4a (Nd,)), **8a** (Tb, \square) and **9a** (Dy, Δ). Solid lines correspond to simulations using the crystal field parameters of set 1 in Table 1, dashed lines to those using set 2.

Table 1 Crystal field parameters suitable for simulations of the static magnetization

Crystal field parameters (cm ⁻¹)	4a	8a	9a
$B_2^0/B_4^0 ext{ (set 1)} B_2^0/B_4^0 ext{ (set 2)}$	960/-300	-220/-70	-460/-20
	-560/340	-610/10	-620/-50

modelled using eqn (1), which contains a direct term (AT) and a two phonons Raman term.⁵⁹

$$\tau^{-1} = AT + \tau_0^{-1} \frac{e^{\frac{-E_{\text{vib}}}{k_B T}}}{\left(e^{\frac{-E_{\text{vib}}}{k_B T}} - 1\right)^2}$$
(1)

Since 4a is a field-induced SMM, no quantum tunneling parameter was included when modelling the relaxation data. Full details on all of the terms and parameters used for the fit are outlined in Fig. 11 and S30.† Initial inclusion of a Raman term (CTⁿ) yielded a power dependence n term of \sim 3.5; since this is a significant departure from the idealized Raman relaxation mechanism, we excluded any direct Raman contribution. Moreover, the use of a model with an Orbach $\left({{ au _0}^{ - 1} \cdot e^{rac{{ - {U_{
m eff}}}}{{{k_{
m B}}T}}}} \right.$

and a direct term (AT) yielded a sufficient model, however the Ueff was significantly smaller than what is expected for a true Orbach process at 14 cm⁻¹. We therefore turned to a model with direct and two-phonons Raman terms, which yielded a nearly identical U_{eff} , or rather E_{vib} at 13 cm⁻¹ (Fig. S30†), corresponding to phonon-coupled under-barrier relaxation process rather than a true Δm_I transition/Orbach process.

Discussion

Ligand design for Tp-based antenna ligands

Comparison of the optical properties of the [Ln(4-NO₂Tp)₃] complexes and those of previously published complexes $[LnTp_3]$ and $[Ln(3-NO_2Tp)_2(NO_3)]$ shows the markedly inferior sensitizing efficiency of the 4-NO₂Tp⁻ ligand. Given the similarity of 3-NO₂Tp⁻ and 4-NO₂Tp⁻, including the prevalence in both ligands of $\pi_{Pz} \to {\pi^*}_{n\text{-NO},Pz}$ charge transfer transitions, predicted by TD-DFT and confirmed experimentally, this difference is rather surprising. To elucidate the origin of the disparity between the n-NO₂Tp⁻ (n = 3 or 4) ligands, we determined the quantum yield (Φ_{exp}) , emission lifetime (τ_{Eu}) and derived their intrinsic quantum yield (Φ_{Eu}^{Eu}), the rates of radiative (k_{rad}) and nonradiative decay (k_{nonrad}) as well as their efficiency of sensitization (η_{sens}) (Table 2). The radiative lifetime for Eu³⁺ is similar for all three complexes, with a slight reduction in [Eu(3-NO₂Tp)₂(NO₃)]. Expectedly, 10-coordinate [Eu(3-NO₂Tp)₂(NO₃)] has a lower intrinsic quantum yield than the two 9-coordinate complexes, which show nearly identical $\Phi_{\rm Eu}^{\rm Eu}$. The primary difference between the EuTp-complexes is their quantum yields. Consistent with the reduction in sensitization efficiency often associated with nitrated antenna ligands, 38 the quantum yield and radiative decay rate of 6a are about half of those of [EuTp₃] (~2.4% vs. 4.6% for $\Phi_{\rm exp}$ and ~13 s⁻¹ vs. 25 s⁻¹ for $k_{\rm rad}$), however the quantum yield and radiative decay rate for [Eu(3-NO₂Tp)₂(NO₃)] are 4% and 21 s⁻¹ respectively, approximately equal to those of [EuTp₃]. The reduction in the radiative decay of 6a as compared to [EuTp₃]

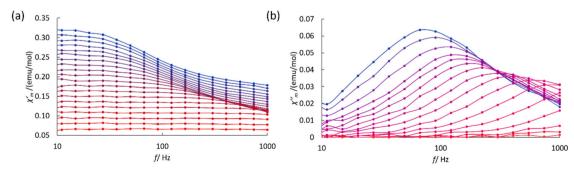


Fig. 8 (a) In-phase molar susceptibility (χ'_m) vs. frequency and (b) out-of-phase molar susceptibility (χ''_m) vs. frequency of 4a between 2 and 9 K at H = 400 Oe.

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Table 2 Photophysical data comparing [EuTp₃], [Eu(3-NO₂Tp)₂(NO₃)] and 6a

Compounds	$ au_{\mathrm{Eu}}/\mathrm{ms}$	$ au_{ ext{ligand (T1)}}/ ext{ms}$	$\Phi_{\mathrm{Eu}}^{\mathrm{Eu}}{}^a/\%$	$\Phi_{\mathrm{exp}}\left(\Phi_{\mathrm{Eu}}^{\mathrm{Tp}^{*}}\right)\!/\%$	$k_{\rm rad}^{\ b}/{\rm s}^{-1}$	$k_{ m nornrad}^{\ \ c}/{ m s}^{-1}$	$\eta_{ m sens}^{\ a}/\%$
$ \begin{split} & [EuTp_3] \\ & [Eu(3\text{-NO}_2Tp)_2(NO_3)] \\ & [Eu(4\text{-NO}_2Tp)_3] \cdot C_6H_{14} \ (\textbf{6a}) \end{split} $	1.84 (ref. 60) 1.90 \pm 0.27 1.82 \pm 0.03 ^e 1.94 \pm 0.20 ^f	1.90 (ref. 60), 9.00 (ref. 34) ^d 0.113 ± 0.07 ^d N/A	43 (ref. 60) 22 45	4.6 (ref. 60) 4.00 1.8 ^e 2.9 ^f	25 21 ± 3 9.9 ± 0.2^{e} 16 ± 1.6^{f}	518 506 ± 72 539 ± 8^{e} 395 ± 53^{f}	10.7 18.2 4.0^{e} 6.4^{f}

 $a_{\text{Nens}} = (\Phi_{\text{exp}}/\Phi_{\text{Eu}}^{\text{Eu}})$ according to de Sá et al. $a_{\text{rad}}^{\text{61}} = \Phi_{\text{exp}}/\tau_{\text{Eu}}$. $a_{\text{Nonrad}}^{\text{62}} = (1/\tau_{\text{Eu}}) - k_{\text{rad}}$. $a_{\text{Lifetime}}^{\text{62}} = (1/\tau_{\text{Eu}}) - k_{\text{rad}}$. $a_{\text{Cir}}^{\text{62}} = (1/\tau_{\text{Eu}}) - k_{\text{Cir}}$. Measured with a 340 nm excitation. f Measured with a 395 nm excitation.

can be attributed to the fast decay/short lifetime of the T₁ state of 4-NO₂Tp⁻ and inefficient energy transfer to Eu³⁺ owing to the presence of the nitro-group. However, the nitro group appears to have no impact in [Eu(3-NO₂Tp)₂(NO₃)].

Examination of the differences between [Eu(3- $NO_2Tp_2(NO_3)$ and 6a seem to indicate that the triplet state of 3-NO₂Tp⁻ is more stable/longer-lived than that of 4-NO₂Tp⁻, given that the emission from 4-NO₂Tp⁻ could not be observed at even 77 K, while we recorded room temperature T1 emission from 3-NO₂Tp⁻ in [Gd(3-NO₂Tp)₂(NO₃)].⁴⁵ TD-DFT calculations predict that both 3-NO₂Tp⁻ and 4-NO₂Tp⁻ should have similar triplet energies (422 and 404 nm respectively), indicating that the enhanced stability of 3-NO₂Tp⁻ must be due to external factors, outside of the energy of the molecular orbitals. One possible explanation is revealed by a comparison of the noncovalent interactions in the crystal structures $[Eu(3-NO_2Tp)_2(NO_3)]$ and **6a**. These interactions are $\pi-\pi$ stacking and NO₂···H-C hydrogen bonding in [Eu(3-NO₂Tp)₂(NO₃)] and 6a respectively, both interactions involving the n-NO₂Pz $(n = 3 \text{ or } 4) \text{ rings. Natural bond order (NBO) derived stabiliz$ ation energies associated with each NCI (π - π stacking or NO₂···H-C hydrogen bonding) (Fig. 9) and indicates that the π - π stacking interactions are significantly stronger than NO₂···H-C hydrogen bonding. We postulate that the stronger π - π stacking NCIs of [Eu(3-NO₂Tp)₂(NO₃)] increase the rigidity of 3-NO₂Tp⁻. The coordination of the 3-NO₂ group to the Ln³⁺ centers is likely to also contribute to an increased rigidity, accounting overall for the higher sensitization efficiency observed in [Eu(3-NO₂Tp)₂(NO₃)] vs. 6a. This is consistent with recent reports on NCIs stabilizing the triplet state of aromatic ligands. 62-65 The enhanced stability afforded to the T₁ state of 3-NO₂Tp⁻ in [Eu(3-NO₂Tp)₂(NO₃)] likely increased its lifetime enough that it could undergo substantial energy transfer with Eu³⁺ (Fig. 10), while the T₁ state of 4-NO₂Tp⁻ undergoes a rapid non-radiative decay. By changing the position of the ring substituent in n-NO₂Tp⁻ (n = 3 or 4) ligands, we can control the type and strength of NCIs in Ln(n-NO₂Tp)_x complexes and consequently tune the efficiency of n-NO2Tp-to-Ln3+ energy transfer(s) using second-sphere interactions.

Ligand design for Tp-based single-molecule magnets

The capped trigonal prismatic coordination sphere provided by the 4-NO₂Tp⁻ ligands corresponds to a relatively strongly axial crystal. Therefore, single molecule magnet properties are more likely to be obtained with lanthanide ions with an oblate electronic density along with a relatively high J value.8 Those

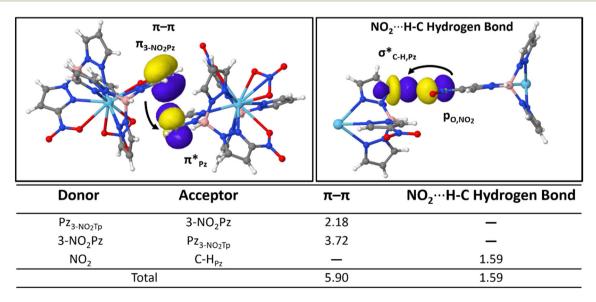


Fig. 9 Top: Representative natural bond orbitals (NBOs) of the $\pi-\pi$ stacking interactions in [Eu(3-NO₂Tp)₂(NO₃)] and the NO₂····H-C hydrogen bonding in 6a. Bottom: Table summarizing the donor and acceptor pairs and associated contribution to the stabilization energy in kJ mol $^{-1}$.

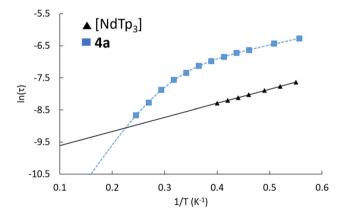
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Energy/ **Photon Absorption** x103 cm-1 Inter-systems Crossing Energy Transfer to Ln3+ M→ Non-radiative Decay S₇ICLT $(\pi_{Pz}\pi^*_{4NO2Pz})$ 30 S_1^{ICLT} T_1 ILCT $(\pi_{Pz}\pi^*_{3NO2Pz})$ (n_{NO2}π*_{4NO2Pz}) $(n_{NO2}\pi^*_{4NO2Pz})$ T_1^{ILCT} Tb3+ (5D. $(n_{NO2}\pi^*_{3NO2Pz})$ Pr3+ 20 Tb³⁺ 5D₄ Sm³⁺ 4G_{5/2} Eu³⁺ 445 nm $^{5}D_{0}$ 483 nm 404* nm 470 nm 10 400 nm unstable T₁ π-π stabilized T₁ **Energy transfer** Luminescence Luminescence = Energy transfer **Pathway Pathway** 3-NO₂Tp⁻

Fig. 10 Jablonski diagram comparing the luminescence and energy transfer pathways in 3-NO₂Tp⁻ and 4-NO₂Tp⁻.

include neodymium (J = 9/2), terbium (J = 6) and dysprosium (J = 15/2). Magnetic studies on 4a (Nd), 8a (Tb) and 9a (Dy) offer a preliminary confirmation of the axiality of the crystal field, as no B_{k}^{q} term with $q \neq 0$ appears in the modelisation of the dc magnetization data. However, either set of crystal field parameters (Table 1) obtained from those fits indicate ground states with $m_I = 0$ for 8a and $m_I = \frac{1}{2}$ for 9a, in agreement with the absence of single molecule magnetism observed for those complexes. Of the two sets obtained for 4a, the first corresponds to a ground state $m_I = 3/2$ compatible with the observed slow magnetic relaxation and should therefore be preferred to the second set which would indicate a $m_I = \frac{1}{2}$ ground state.

In 4a, as in the unsubstituted analogue [NdTp₃],³² slow relaxation can only be observed under a small external magnetic field. This indicates substantial quantum tunneling of the magnetization at zero field, and in turn the presence of some small equatorial contribution to the crystal field, albeit one that could not be quantified from the dc data. Interestingly, relaxation is about 5 times slower in 4a than in [NdTp₃] at the same temperature (Fig. 11) despite the virtually identical geometry of the two complexes (and higher crystallographic symmetry in [NdTp3]). Moreover, 4a's relaxation barrier, Evib is about 6 times larger than NdTp3's "Orbach" barrier (12.7 vs. 2.8 cm⁻¹). [NdTp₃] was previously reported to have a small relaxation barrier ($U_{\text{eff}} = 2.84 \text{ cm}^{-1}$, $\tau_0 = 4.2 \times 10^{-5} \text{ s}$), ³² likely not a "true" Orbach relaxation barrier but rather a value corresponding to under barrier relaxation due to vibronic coupling. Given that dilution of NdTp₃ with diamagnetic La³⁺ ions, did not significantly affect the relaxation behavior, 32 the slight increase in intermolecular Ln3+-to-Ln3+ distances in 4a is likely not responsible for the improvement either, nor is the



4-NO₂Tp⁻

Fig. 11 Temperature dependence of the magnetic relaxation times, under a small bias field, of [NdTp3] (black triangles, 100 Oe) with the reported Orbach fit (black line),32 and 4a (blue squares, 400 Oe) with the fit (blue broken line) generated using egn (1). From the fitted AC data (solid lines), [NdTp₃]: $U_{\text{eff}} = 2.84(2) \text{ cm}^{-1}$, $\tau_0 = 4.2(2) \times 10^{-4} \text{ s}$; 4a: A = 318(8) K, $U_{\rm eff} = 12.7(1.6) \ {\rm cm}^{-1}$ and $\tau_0 = 4.1(1.6) \times 10^{-5} \ {\rm s}$.

change from the parallel anisotropy axes in [NdTp₃]⁶⁶ to the slightly canted layers of 4a, which one would usually assume to be detrimental to slow relaxation. A change in vibrational levels is therefore the most likely mechanism for the observed difference, linked to the lattice rigidifying effect of the NO2···H-C hydrogen bonding in 4a. This marked increase in relaxation times and under-barrier in 4a over [NdTp3], signify that the former is a better SMM than the latter owing to the nitro-functionalization. The apparent similarity in the primary coordination sphere of 4a and [NdTp₃] facilitates the first example of the influence of a strongly electron-withdrawing group on relaxation dynamics via a direct comparison of the two Nd³⁺ chelates.

Conclusion

Eleven new lanthanide complexes with an asymmetric nitrofunctionalized trispyrazolylborate ligand, 4-NO₂Tp⁻ were synthesized and their optical and magnetic properties are presented. The [Ln(4-NO₂Tp)₃] complexes display crystalline polymorphism based on the solvent used for recrystallization. The choice of lattice solvent can be used as a method of controlling crystal packing and supramolecular assembly (or lack thereof) via NO₂···H-C hydrogen bonding. An in-depth analysis of the optical properties of 4-NO₂Tp⁻ revealed two distinct pathways for luminescence and energy transfer comprised of several ILCT transitions such as $n_{NO_2} \rightarrow \pi^*_{4\text{-NO}_2Pz}$ and $\pi_{Pz} \rightarrow \pi^*_{4\text{-NO}_2Pz}$ absorptions and $\pi^*_{\text{4-NO,Pz}} \rightarrow n_{\text{NO,}}$ and $\pi^*_{\text{4-NO,Pz}} \rightarrow \pi_{\text{Pz}}$ emissions. Moreover, 4-NO₂Tp⁻ is a poor sensitizer for Ln³⁺ emission (vs. Tp and 3-NO₂Tp), as a direct result of adding a non-coordinating nitro-group. With regard to the magnetic properties of the reported complexes, the opposite is true, whereby adding a nitro-group increased the relaxation times by a factor of 5 as well as the under-barrier of [Nd(4-NO₂Tp)₃] over [NdTp₃] by one order of magnitude. Our findings on the ability to control the crystal packing/non-covalent interactions, diverse charge transfer optical properties and enhancement of relaxation barriers using 4-NO₂Tp⁻ over Tp⁻, in context with our previous work on [Ln(3-NO₂Tp)_x] complexes, highlights the utility of asymmetric Tp⁻ ligands for tuning the properties of the lanthanides to generate and study magneto-luminescent materials and design better molecular magnets.

Experimental

Materials

Lanthanide chloride salts, $LnCl_3 \cdot 7H_2O$ ($Ln = La^{3+}$, Ce^{3+} , Fisher Scientific, 99.9%), $PrCl_3 \cdot 7H_2O$ (Strem Chemicals, 99.9%), $LnCl_3 \cdot 6H_2O$ ($Ln = Nd^{3+}$, $Sm^{3+}-Dy^{3+}$, Aldrich, 99.9%) are commercially available and were used as received.

Synthesis of tetrabutylammonium 4-nitrotrispyrazolylborate, $[Bu_4N][4-NO_2Tp]$

 $[Bu_4N][4\text{-NO}_2Tp]$ was prepared as described in the literature. After each synthesis of $[Bu_4N][4\text{-NO}_2Tp]$, a standard solution, typically 0.3 M, was prepared by dissolving the reaction product in methanol.

Synthesis of potassium 4-nitrotrispyrazolylborate, $K[4-NO_2Tp]$

 $K[4\text{-NO}_2\text{Tp}]$ was prepared by dissolving 4-nitropyrazole (0.4485 g, 4 mmol) and potassium trispyrazolylborate (1.000 g, 4 mmol) in dry ethyl acetate and heating under reflux for 2 hours. Pentane (50 mL) was added to the cloudy mixture which was then placed in a refrigerator at 5 $^{\circ}\text{C}$ overnight. Two

distinct phases were always isolated from this reaction, a white powder of $K[4\text{-}NO_2\text{Tp}]$ in very low yield (>5%) and a yellow powdery mixture of $K[4\text{-}NO_2\text{Tp}]$ and 4-nitropyrazole. Owing to the low yields, the solvent-free synthesis with the tetrabutyl-ammonium salt was always used to generate the $4\text{-}NO_2\text{Tp}^-$.

Synthesis of Ln(4-NO₂Tp)₃·solvent (1-9)

Complexes of **1–9** were synthesized using the same procedure differing only in the recrystallization conditions. A standard solution of $LnCl_3$ was generated by dissolving an appropriate amount of the salt, *e.g.* $LaCl_3$ (0.4910 g, 2.002 mmol) in methanol (20 mL, 0.1 M). Solutions of $LaCl_3$ (1 mL, 0.1 M, 1 mmol) and $[Bu_4N][4\text{-NO}_2\text{Tp}]$ (1 mL, 0.3 M, 3 mmol) were combined in a scintillation vial. The solvent was removed *in vacuo* and the oil-like residue was suspended in 3 mL 1:1 cyclohexane: isopropanol; vacuum filtration and washing with of isopropanol (3 × 3 mL) yielded dry, crude $La(4\text{-NO}_2\text{Tp})_3$ as a white powder.

Recrystallized solvates were prepared by dissolving the crude product in of methanol (3 mL, La–Eu) or 1:1 methanol: ethanol (3 mL, Gd–Dy). Hexanes or benzene for (6 mL) was layered on top of the solution. Single crystals of hexane (1a–9a) or benzene (5b, 8b) solvates were observed after 1 day and the final product was isolated and collected after 3 days *via* vacuum filtration and washing with hexanes (3 × 3 mL). Crystals of the hexane polymorph had a tendency to re-dissolve in the mother liquor after 5 or 6 days. X-ray diffraction quality single crystals were collected when the recrystallization was carried out in a sealed scintillation vial, while a microcrystalline powder was obtained when the synthesis was done in a parafilmed 25 mL Erlenmeyer flask.

[La(4-NO₂Tp)₃]-C₆H₁₄ (1a). White single crystals were isolated from methanol/hexanes. Yield (recrystallized): 51%. Elemental Anal. Calc. for LaC₃₃H₄₁O₆N₂₁B₃: C 39.67, H 4.11, N 29.45. Found: C 39.40, H 3.59, N 29.05.

[Ce(4-NO₂Tp)₃]·C₆H₁₄ (2a). Pale yellow single crystals were isolated from methanol/hexanes. Yield (recrystallized): 52%. Elemental Anal. Calc. for $CeC_{33}H_{41}O_6N_{21}B_3$: C 39.62, H 4.10, N 29.41. Found: C 39.21, H 3.88, N 29.55.

[Pr(4-NO₂Tp)₃]- C_6H_{14} (3a). Green single crystals were isolated from methanol/hexanes. Yield (recrystallized): 45%. Elemental Anal. Calc. for $PrC_{33}H_{41}O_6N_{21}B_3$: C 39.60, H 4.10, N 29.39. Found: C 39.45, H 3.94, N 29.59.

 $[Nd(4-NO_2Tp)_3]\cdot C_6H_{14}$ (4a). Lilac single crystals were isolated from methanol/hexanes. Yield (recrystallized): 51%. Elemental Anal. Calc. for $NdC_{33}H_{41}O_6N_{21}B_3$: C 39.47, H 4.09, N 29.29. Found: C 39.29, H 3.94, N 29.13.

[Sm(4-NO₂Tp)₃]·C₆H₁₄ (5a). Pale yellow single crystals were isolated from methanol/hexanes. Yield (recrystallized): 36%. Elemental Anal. Calc. for SmC₃₃H₄₁O₆N₂₁B₃: C 39.22, H 4.06, N 29.11. Found: C 39.11, H 3.91, N 29.30. τ = 0.116 ms ($\lambda_{\rm exc}$ = 340 nm (6 P_{3/2}), $\lambda_{\rm em}$ = 604 nm).

 $[Eu(4-NO_2Tp)_3]\cdot C_6H_{14}$ (6a). White single crystals were isolated from methanol/hexanes. Yield (recrystallized): 30%. Elemental Anal. Calc. for $EuC_{33}H_{41}O_6N_{21}B_3$: C 39.16, H 4.05, N

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29.07. Found: C 38.87, H 3.87, N 29.07. τ = 1.80 ms ($\lambda_{\rm exc}$ = 394 nm ($^{5}L_{\rm 6}$), $\lambda_{\rm em}$ = 700 nm).

 $[Gd(4-NO_2Tp)_3]\cdot C_6H_{14}$ (7a). White single crystals were isolated from methanol/hexanes. Yield (recrystallized): 33%. Elemental Anal. Calc. for $GdC_{33}H_{41}O_6N_{21}B_3$: C 38.96, H 4.03, N 28.92. Found: C 38.91, H 3.97, N 29.04.

[**Tb**(4-NO₂**Tp**)₃]-C₆H₁₄ (8a). White single crystals were isolated from methanol/hexanes. Yield (recrystallized): 50%. Elemental Anal. Calc. for TbC₃₃H₄₁O₆N₂₁B₃: C 38.90, H 4.03, N 28.87. Found: C 38.71, H 4.00, N 28.92. τ = 1.87 ms ($\lambda_{\rm exc}$ = 486 nm (5 D₄), $\lambda_{\rm em}$ = 542 nm).

 $[Dy(4-NO_2Tp)_3]\cdot C_6H_{14}$ (9a). White single crystals were isolated from methanol/hexanes. Yield (recrystallized): 36%. Elemental Anal. Calc. for $DyC_{33}H_{41}O_6N_{21}B_3$: C 38.76, H 4.01, N 28.77. Found: C 38.78, H 3.92, N 28.66.

 $[Sm(4-NO_2Tp)_3]$ - $7C_6H_6$ (5b). Pale yellow single crystals were isolated from methanol/benzene. Yield (recrystallized): 44%. Elemental Anal. Calc. for $SmC_{69}H_{69}O_6N_{21}B_3$: C 56.33, H 4.69, N 20.00. Found: C 55.41, H 4.66, N 20.61.

[Tb(4-NO₂Tp)₃]·7C₆H₆ (8b). White single crystals were isolated from methanol/benzene. Yield (recrystallized): 40%. Elemental Anal. Calc. for TbC₆₉H₆₉O₆N₂₁B₃: C 56.01, H 4.67, N 19.88. Found: C 55.40, H 4.70, N 20.43.

X-ray structure determination

Crystals of 1a-9a, 5b and 8b, were harvested from mother liquors and mounted on 50 µm MiTeGen mounts. All measurements were made using monochromated microfocus Mo Kα (λ = 0.71073) radiation on a Bruker D8 Quest, equipped with a Photon II detector. All reflection data were collected at 100(2) K with 0.5° φ and ω scans. The data were reduced using SAINT, 67 and empirical absorption corrections were applied using SADABS, 68 for 1a-9a, 5b and 8b. Structure solutions solved using intrinsic phasing were performed using the ShelXT package⁶⁹ in APEX III. All data were subsequently refined using SHELXL-2014 in the program SHELXle.⁷⁰ All atoms were refined anisotropically. The reported CIFs for 1a-9a, feature structural models that were refined with scattering contributions from the disordered hexane sites removed from the diffraction data using the bypass procedure in PLATON.⁷¹ The electron count from the "squeezed" model converged in good agreement with a single hexane molecule. Aromatic hydrogen atoms were placed in idealized positions and allowed to ride on the coordinates of the parent atom with isotropic thermal parameters (U_{iso}) fixed at $1.2U_{eq}$ for all carbon atoms and at $1.5U_{\rm eq}$ for all boron atoms. Details of the X-ray diffraction experiments and crystal data are summarized in Table 3.

Powder X-ray diffraction

Powder X-ray diffraction (PXRD) data on the bulk recrystallization products from each sample were collected on a Rigaku Miniflex (Cu K α 2 θ = 5–60) and analyzed using the Match! software program. The PXRD patterns of the bulk products for **1a**–**9a**, **5b** and **8b** were used to check purity and reproducibility and are provided in ESI Fig. S3 and S4.†

Infrared spectroscopy

Infrared spectra were collected from 650 to 4000 cm⁻¹ using a PerkinElmer Frontier FT-IR spectrophotometer with a diamond attenuated total reflectance (ATR) sample holder. The IR spectra of the bulk, recrystallized products for **1a-9a**, **5b** and **8b** were used to check purity and reproducibility and are provided in ESI Fig. S5 and S6.†

¹H nuclear magnetic resonance spectroscopy

NMR spectra of crude **1–9** (no solvates) were recorded on a Varian 400 MHz spectrometer. Spectra of crude **1–9** (except 7) are provided in the ESI Fig. S4–S14.†

Thermogravimetric analysis

Thermogravimetric analysis (TGA) was performed on a PerkinElmer Pyris 1 TGA. Samples were heated from 30-500 °C at 20 °C min⁻¹ in air. TGA curves of **5a** and **5b** are given in Fig. S15† indicating thermal stability of [Ln(4-NO₂Tp)₃] up to ~240 K, with solvent loss at 185 °C for **a** (hexanes) solvate and 77 °C for **b** (benzene) solvate.

Elemental analysis

Elemental analyses were performed by ALS Environment for 1a and 9a, and Atlantic Microlab for 2a-8a, 5b and 8b.

Photophysical measurements

Visible and NIR solid-state luminescence measurements were obtained at room temperature for 1a-9a (and 77 K for 6). Luminescence spectra were collected with a Horiba Jobin Yvon Fluorolog-3 spectrophotometer using a 450 W xenon arc lamp combined with a double excitation monochromator and double emission monochromator. For spectra in the visible region, a photomultiplier tube at 950 V was used as the emission detector, whereas for spectra in the near-IR region, a liquid nitrogen cooled, Symphony II NIR InGaAs diode array detector was used as the emission detector. Data were collected and analyzed using the FluorEssence software package. The solid samples were mounted on a J1933 Solid Sample Holder using non-emitting high vacuum grease for room temperature scans. Low temperature luminescence measurements were collected on solid samples under vacuum using a Janis VPF-100 cryostat equipped with UV-grade fused silica windows coupled with a Lakshore model 325 temperature controller. Lifetime measurements were collected with a Horiba Jobin Yvon Flurolog-3 spectrophotometer adapted for time-correlated single photon counting (TCSPC) measurements using a xenon flash lamp as the light source. Lifetime profiles for 5a, 6a and 8a were obtained using the TCSPC module and the data were fit using DAS6 software. Quantum yields measurements were collected in duplicate using a Horiba PTI QM-400 fluorometer using a PTFE powder holder under ambient conditions and an 8.9 cm integrating sphere with a Spectralon fluropolymer coating. The samples were crushed with a mortar and pestle, diluted with KBr in a 1:50 ratio of sample: KBr prior to data collection. Diffuse reflectance spectra were col-

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Table 3 Crystallographic data for 1a-9a, 5b and 8b

	1a	2a	3a	4a	5a	5 b
CCDC no.	2226885	2226886	2226887	2226888	2226889	2226890
Formula	$LaN_{21}O_6B_3C_{27}H_{27}$	$CeN_{21}O_6B_3C_{27}H_{27}$	$PrN_{21}O_{6}B_{3}C_{27}H_{27}$	$NdN_{21}O_{6}B_{3}C_{27}H_{27}$	$SmN_{21}O_6B_3C_{27}H_{27}$	$SmN_{21}O_{6}B_{3}C_{69}H$
Formula weight	913.04	914.25	915.04	918.37	924.49	1471.24
Crystal system	Monoclinic	Monoclinic	Monoclinic	Monoclinic	Monoclinic	Trigonal
Space group	C2/c	C2/c	C2/c	C2/c	C2/c	$R\bar{3}$
a, Å	24.473(5)	24.400(8)	24.369(6)	24.2443(7)	24.324(15)	21.6878(4)
b, Å	13.830(3)	13.801(3)	13.783(3)	13.7711(3)	13.756(10)	21.6878(4)
c, Å	29.778(4)	28.778(10)	28.792(̈́7)	28.8113(11)	28.833(15)	25.7723(8)
α, °	90	90	90	90	90	90
β , \circ	116.898(5)	116.759(12)	112.453(8)	112.537(1)	112.49(2)	90
γ, °	90	90	90	90	90	90
<i>V</i> , Å ³	8988(3)	8954(5)	8937(4)	8921.3(5)	8913(10)	10 498.2(5)
$Z^{'}$	8	8	8	8	8	6
<i>T</i> , K	100	100	100	100	100	100
$\rho_{\rm calc}$, g cm ⁻³	1.349	1.356	1.360	1.367	1.378	1.396
μ , mm ⁻¹	1.011	1.077	1.150	1.224	1.378	0.907
λ, Μο Κα	0.71073	0.71073	0.71073	0.71073	0.71073	0.71073
$R_{ m int}$	0.0483	0.0590	0.0455	0.0639	0.0673	0.0403
Residuals: ^a R; R _w	0.0219; 0.0477	0.0236; 0.0546	0.0252; 0.0638	0.0295; 0.0651	0.0268; 0.0613	0.0175; 0.0426
Goodness of fit	1.022	1.042	1.049	1.007	0.996	1.071
	6a	7a	8a	8	lb	9a

	6a	7a	8a	8b	9a
CCDC no.	2226891	2226892	2226893	2227039	2226894
Formula	$EuN_{21}O_6B_3C_{27}H_{27}$	$GdN_{21}O_6B_3C_{27}H_{27}$	$TbN_{21}O_6B_3C_{27}H_{27}$	$TbN_{21}O_6B_3C_{69}H_{69}$	$DyN_{21}O_6B_3C_{27}H_{27}$
Formula weight	926.10	931.38	933.06	1479.80	936.63
Crystal System	Monoclinic	Monoclinic	Monoclinic	Trigonal	Monoclinic
Space group	C2/c	C2/c	C2/c	$R\bar{3}$	C2/c
a, Å	24.313(16)	24.296(6)	24.248(2)	21.6423(9)	24.291(5)
b, Å	13.724(10)	13.742(2)	13.7248(17)	21.6423(9)	13.695(3)
c, Å	29.722(15)	29.768(4)	28.843(4)	25.7873(17)	29.707(8)
α , \circ	90	90	90	90	90
β , \circ	115.893(13)	116.414(8)	112.506(4)	90	115.032(5)
γ, °	90	90	90	90	90
V , $\mathring{\mathbf{A}}^3$	8922(10)	8901(3)	8867.9(18)	10 460.3(11)	8954(4)
Z	8	8	8	6	8
<i>T</i> , K	100	100	100	100	100
$\rho_{\rm cale}, {\rm g~cm}^{-3}$	1.379	1.390	1.398	1.409	1.390
μ , mm ⁻¹	1.466	1.551	1.656	1.082	1.729
λ, Μο Κα	0.71073	0.71073	0.71073	0.71073	0.71073
$R_{ m int}$	0.0643	0.0608	0.0394	0.0800	0.0477
Residuals: R; R _w	0.0430; 0.0872	0.0251; 0.0629	0.0323; 0.0713	0.0355; 0.0689	0.0426; 0.0852
Goodness of fit	0.975	1.020	1.003	1.118	1.163

 $^{^{}a}R = R_{1} = \sum |F_{0}| - |F_{c}||/\sum |F_{0}|$ for observed data only. $R_{w} = wR_{2} = \{\sum [w(F_{0}^{2} - F_{c}^{2})^{2}]/\sum [w(F_{0}^{2})^{2}]\}^{1/2}$ for all data.

lected on solid samples at 298 K. The light source was a Mikropack DH-2000-BAL deuterium and halogen light source coupled with an Ocean Optics Flame detector. Scattered light was collected with a fiber-optic cable. Spectra were referenced with BaSO₄. Data were processed using OceanView spectroscopy software. The electronic absorption spectrum of 1 was collected on a SPECORD 600 UV-VIS diode array spectrophotometer.

Magnetic SQUID measurements

Crystalline samples for static SQUID magnetometry were prepared inside a glovebox where a measured amount of sample was added to a half-sealed quartz tube using a glass pipet followed by a measured amount of eicosane. The top of the tube was then fitted to an Ultra Torr Swagelok adapter. This was taken out of the glovebox and the eicosane was melted using hot water to fix the sample. The assembly was then attached to a Schlenk

line, and the top of the tube was sealed using an H₂/O₂ torch while the sample was under vacuum. The sealed tube was taped to a straw using Kapton tape and loaded onto the instrument. For the dynamic measurements, the samples were enclosed in a copper foil pouch and mounted on a quartz holder with Kapton tape. The molar diamagnetic susceptibilities of the compounds were estimated from their molar mass $(\chi_{dia} (cm^3 mol^{-1}) = -[MW]$ $(g \text{ mol}^{-1}) \times 10^{-6}$]/2) and subtracted from the experimental value. Data was collected at 0.1 T for static measurements and 0.04 T for dynamic measurements for using a Quantum Design SQUID MPMS3 magnetometer.

Computational methods

Computational studies were conducted using the high-performance computing cluster at the George Washington University. The input structure for Calc-1 was derived from the crystal structure of 1a. The frontier molecular orbitals and the

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UV/VIS spectrum of Calc-1 were computed using density functional theory (DFT) in the Gaussian 16 software (Gaussian Inc.).⁷² A ground state optimization was performed on Calc-1, using the B3LYP^{73,74} level of theory with the modified scalarrelativistic effective core potential (ECP) basis set def2-TZVP as implemented in the software with the def2-TZVP pseudopotential applied to La^{3+,75-78} Geometry optimizations were performed on the input structure without symmetry constraints in the gas phase. Subsequent frequency calculations were performed on Calc-1 to confirm that the optimized structure was its global minimum. No imaginary frequencies were present in the calculated IR and Raman spectra. Time-dependent DFT calculations were then conducted to determine the expected absorption spectra of Calc-1 considering singlet $(S_0 \rightarrow S_n)$ and triplet $(S_0 \to T_n)$ transitions. 100 singlet and triplet states each were identified for Calc-1. NBO calculations were performed using NBO7 on two models. Model A consisted of two molecules of [La(3-NO₂Tp)₂(NO₃)] (generated from crystallographidetermined atomic coordinates of [Gd(3- $NO_2Tp)_2(NO_3)$ $\frac{1}{4}H_2O)$ participating in a π - π stacking interaction. Model B comprised two [La(4-NO₂Tp)]²⁺ units participating in a NO2···H-C hydrogen bond (generated from crystallographically determined atomic coordinates in 1a). NBO second-order perturbation theory was applied to quantify the magnitude of the donor-acceptor interaction between adjacent complexes as well as to identify the atomic and molecular orbitals involved with Jmol as the visualization software.⁷⁹

Author contributions

The manuscript was written through contributions of all authors.

Conflicts of interest

The authors declare no competing financial interests.

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References

- 1 K. Binnemans, Lanthanide-Based Luminescent Hybrid Materials, Chem. Rev., 2009, 109, 4283-4374.
- 2 A. Dey, P. Kalita and V. Chandrasekhar, Lanthanide(III)-Based Single-Ion Magnets, ACS Omega, 2018, 3, 9462-9475.
- 3 C. A. P. Goodwin, Blocking like it's hot: a synthetic chemists' path to high-temperature lanthanide single molecule magnets, Dalton Trans., 2020, 49, 14320-14337.
- 4 J. A. Smith, M. A. Singh-Wilmot, K. P. Carter, C. L. Cahill and J. A. Ridenour, Lanthanide-2,3,5,6-Tetrabromoterephthalic Acid Metal-Organic Frameworks: Evolution of Halogen...Halogen Interactions across the Lanthanide Series and Their Potential as Selective Bifunctional Sensors for the Detection of Fe3+, Cu2+, and Nitroaromatics, Cryst. Growth Des., 2019, 19, 305-319.
- 5 J. A. Smith, M. A. Singh-Wilmot, K. P. Carter, C. L. Cahill and J. A. Ridenour, Supramolecular assembly of lanthanide-2,3,5,6-tetrafluoroterephthalic acid coordination polymers via fluorine...fluorine interactions: a platform for luminescent detection of Fe3+ and nitroaromatic compounds, New J. Chem., 2020, 44, 12317-12330.
- 6 J. H. S. K. Monteiro, F. A. Sigoli and A. D. Bettencourt-Dias, A water-soluble TbIII complex as a temperature-sensitive luminescent probe, Can. J. Chem., 2018, 96, 859-864.
- 7 J. F. C. B. Ramalho, S. F. H. Correia, L. Fu, L. M. S. Dias, P. Adão, P. Mateus, R. A. S. Ferreira and P. S. André, Super modules-based active QR codes for smart trackability and IoT: a responsive-banknotes case study, npj Flexible Electron., 2020, 4, 11.
- 8 J. D. Rinehart and J. R. Long, Exploiting single-ion anisotropy in the design of f-element single-molecule magnets, Chem. Sci., 2011, 2, 2078-2085.
- 9 S. G. McAdams, A.-M. Ariciu, A. K. Kostopoulos, J. P. S. Walsh and F. Tuna, Molecular single-ion magnets based on lanthanides and actinides: Design considerations and new advances in the context of quantum technologies, Coord. Chem. Rev., 2017, 346, 216-239.

10 D. I. Alexandropoulos, K. R. Vignesh, B. S. Dolinar and K. R. Dunbar, End-to-end azides as bridging ligands in lanthanide coordination chemistry: Magnetic and magnetocaloric properties of tetranuclear Ln4 (Ln = Gd, Dy) complexes exhibiting a rare rhombus topology, *Polyhedron*, 2018, 151, 255–263.

Dalton Transactions

- 11 S. Demir, I.-R. Jeon, J. R. Long and T. D. Harris, Radical ligand-containing single-molecule magnets, *Coord. Chem. Rev.*, 2015, **289**, 149–176.
- 12 S. Demir, M. I. Gonzalez, L. E. Darago, W. J. Evans and J. R. Long, Giant coercivity and high magnetic blocking temperatures for N23- radical-bridged dilanthanide complexes upon ligand dissociation, *Nat. Commun.*, 2017, 8, 2144.
- 13 C. A. Gould, E. Mu, V. Vieru, L. E. Darago, K. Chakarawet, M. I. Gonzalez, S. Demir and J. R. Long, Substituent Effects on Exchange Coupling and Magnetic Relaxation in 2,2'-Bipyrimidine Radical-Bridged Dilanthanide Complexes, *J. Am. Chem. Soc.*, 2020, 142, 21197–21209.
- 14 R. Ishikawa, S. Michiwaki, T. Noda, K. Katoh, M. Yamashita, K. Matsubara and S. Kawata, Field-Induced Slow Magnetic Relaxation of Mono- and Dinuclear Dysprosium(III) Complexes Coordinated by a Chloranilate with Different Resonance Forms, *Inorganics*, 2018, 6, 7.
- 15 R. Ishikawa, S. Michiwaki, T. Noda, K. Katoh, M. Yamashita and S. Kawata, Series of Chloranilate-Bridged Dinuclear Lanthanide Complexes: Kramers Systems Showing Field-Induced Slow Magnetic Relaxation, *Magnetochemistry*, 2019, 5(2), 30.
- 16 T. P. Latendresse, V. Vieru, A. Upadhyay, N. S. Bhuvanesh, L. F. Chibotaru and M. Nippe, Trends in trigonal prismatic Ln-[1]ferrocenophane complexes and discovery of a Ho3+ single-molecule magnet, *Chem. Sci.*, 2020, 11, 3936–3951.
- 17 K. R. Meihaus, J. D. Rinehart and J. R. Long, Dilution-Induced Slow Magnetic Relaxation and Anomalous Hysteresis in Trigonal Prismatic Dysprosium(III) and Uranium(III) Complexes, *Inorg. Chem.*, 2011, **50**, 8484–8489.
- 18 K. R. Meihaus, S. G. Minasian, W. W. Lukens, S. A. Kozimor, D. K. Shuh, T. Tyliszczak and J. R. Long, Influence of Pyrazolate vs N-Heterocyclic Carbene Ligands on the Slow Magnetic Relaxation of Homoleptic Trischelate Lanthanide(III) and Uranium(III) Complexes, *J. Am. Chem. Soc.*, 2014, **136**, 6056–6068.
- 19 Y.-S. Ding, N. F. Chilton, R. E. P. Winpenny and Y.-Z. Zheng, On Approaching the Limit of Molecular Magnetic Anisotropy: A Near-Perfect Pentagonal Bipyramidal Dysprosium(III) Single-Molecule Magnet, Angew. Chem., Int. Ed., 2016, 55, 16071–16074.
- 20 K. R. McClain, C. A. Gould, K. Chakarawet, S. J. Teat, T. J. Groshens, J. R. Long and B. G. Harvey, High-temperature magnetic blocking and magneto-structural correlations in a series of dysprosium(III) metallocenium singlemolecule magnets, *Chem. Sci.*, 2018, 9, 8492–8503.
- 21 C. A. Gould, K. R. McClain, J. M. Yu, T. J. Groshens, F. Furche, B. G. Harvey and J. R. Long, Synthesis and Magnetism of Neutral, Linear Metallocene Complexes of

- Terbium(II) and Dysprosium(II), *J. Am. Chem. Soc.*, 2019, **141**, 12967–12973.
- 22 C. A. Goodwin, F. Ortu, D. Reta, N. F. Chilton and D. P. Mills, Molecular magnetic hysteresis at 60 kelvin in dysprosocenium, *Nature*, 2017, 548, 439–442.
- 23 J. Long, Y. Guari, R. A. S. Ferreira, L. D. Carlos and J. Larionova, Recent advances in luminescent lanthanide based Single-Molecule Magnets, *Coord. Chem. Rev.*, 2018, 363, 57–70.
- 24 K. S. Kumar, D. Serrano, A. M. Nonat, B. Heinrich, L. Karmazin, L. J. Charbonnière, P. Goldner and M. Ruben, Optical spin-state polarization in a binuclear europium complex towards molecule-based coherent light-spin interfaces, *Nat. Commun.*, 2021, 12, 2152.
- 25 E. M. Matson, J. J. Kiernicki, N. H. Anderson, P. E. Fanwick and S. C. Bart, Isolation of a uranium(III) benzophenone ketyl radical that displays redox-active ligand behaviour, *Dalton Trans.*, 2014, 43, 17885–17888.
- 26 S. Isabel and N. Marques, Recent advances in the chemistry of f-element poly(pyrazolyl)borate complexes, *New J. Chem.*, 1995, 19, 55–571.
- 27 A. J. Amoroso, A. M. C. Thompson, J. C. Jeffery, P. L. Jones, J. A. McCleverty and M. D. Ward, Synthesis of the new tripodal ligand tris-[3-(2'-pyridyl)pyrazol-1-yl]hydroborate, and the crystal structure of its europium(III) complex, *J. Chem. Soc., Chem. Commun.*, 1994, 2751–2752, DOI: 10.1039/C39940002751.
- 28 J. Fleming, E. Psillakis, S. Couchman, J. Jeffery, J. McCleverty and M. Ward, Complexes of the potentially hexadentate ligand bis{3-[6-(2,2'-bipyridyl)]pyrazol-1-yl} hydroborate with representative s-, p-, d- and f-block metal ions: factors promoting formation of mononuclear or double-helical dinuclear complexes, *Dalton Trans.*, 1998, 537–544, DOI: 10.1039/a707936b.
- 29 H. Flötotto, T. Secker, P. Kögerler and C. Besson, Amine-Functionalized Spin Crossover Building Blocks, Eur. J. Inorg. Chem., 2019, 2019, 4621–4624.
- 30 P. Fang, L. Wang, G. Zhan, W. Yan, P. Huo, A. Ying, Y. Zhang, Z. Zhao, G. Yu, Y. Huang, S. Gong, L. Duan, Z. Liu, Z. Bian and C. Huang, Lanthanide Cerium(III) Tris (pyrazolyl)borate Complexes: Efficient Blue Emitters for Doublet Organic Light-Emitting Diodes, ACS Appl. Mater. Interfaces, 2021, 13, 45686–45695.
- 31 W. Yan, L. Wang, H. Qi, G. Zhan, P. Fang, Z. Liu and Z. Bian, Highly Efficient Heteroleptic Cerium(III) Complexes with a Substituted Pyrazole Ancillary Ligand and Their Application in Blue Organic Light-Emitting Diodes, *Inorg. Chem.*, 2021, **60**, 18103–18111.
- 32 J. D. Rinehart and J. R. Long, Slow magnetic relaxation in homoleptic trispyrazolylborate complexes of neodymium (III) and uranium(III), *Dalton Trans.*, 2012, **41**, 13572–13574.
- 33 J. D. Rinehart, K. R. Meihaus and J. R. Long, Observation of a Secondary Slow Relaxation Process for the Field-Induced Single-Molecule Magnet U(H2BPz2)3, *J. Am. Chem.* Soc., 2010, 132, 7572–7573.

Paper

34 E. A. Mikhalyova, M. Zeller, J. P. Jasinski, R. J. Butcher, L. M. Carrella, A. E. Sedykh, K. S. Gavrilenko, S. S. Smola, M. Frasso, S. C. Cazorla, K. Perera, A. Shi, H. G. Ranjbar, C. Smith, A. Deac, Y. Liu, S. M. McGee, V. P. Dotsenko, M. U. Kumke, K. Müller-Buschbaum, E. Rentschler, A. W. Addison and V. V. Pavlishchuk, Combination of single-molecule magnet behaviour and luminescence properties in a new series of lanthanide complexes with tris (pyrazolyl)borate and oligo(β-diketonate) ligands, *Dalton Trans.*, 2020, 49, 7774–7789.

- 35 M. L. P. Reddy and S. Sivakumar, Lanthanide benzoates: a versatile building block for the construction of efficient light emitting materials, *Dalton Trans.*, 2013, 42, 2663–2678.
- 36 M. Latva, H. Takalo, V.-M. Mukkala, C. Matachescu, J. C. Rodríguez-Ubis and J. Kankare, Correlation between the lowest triplet state energy level of the ligand and lanthanide(III) luminescence quantum yield, *J. Lumin.*, 1997, 75, 149–169.
- 37 V. Tsaryuk, K. Zhuravlev, V. Zolin, P. Gawryszewska, J. Legendziewicz, V. Kudryashova and I. Pekareva, Regulation of excitation and luminescence efficiencies of europium and terbium benzoates and 8-oxyquinolinates by modification of ligands, *J. Photochem. Photobiol.*, A, 2006, 177, 314–323.
- 38 V. Tsaryuk, V. Kudryashova, P. Gawryszewska, R. Szostak, A. Vologzhanina, K. Zhuravlev, Z. Klemenkova, J. Legendziewicz and V. Zolin, Structures, luminescence and vibrational spectroscopy of europium and terbium nitro- and dinitro-substituted benzoates. Nitro groups as quenchers of Ln3+ luminescence, *J. Photochem. Photobiol.*, A, 2012, 239, 37–46.
- 39 S. Viswanathan and A. De Bettencourt-Dias, Eu(III) and Tb (III) Luminescence Sensitized by Thiophenyl-Derivatized Nitrobenzoato Antennas, *Inorg. Chem.*, 2006, **45**, 10138–10146.
- 40 A. D. Bettencourt-Dias and S. Viswanathan, Nitrofunctionalization and luminescence quantum yield of Eu (III) and Tb(III) benzoic acid complexes, *Dalton Trans.*, 2006, 4093–4103, DOI: **10.1039/B606970C**.
- 41 A. P. S. Samuel, J. Xu and K. N. Raymond, Predicting Efficient Antenna Ligands for Tb(III) Emission, *Inorg. Chem.*, 2009, **48**, 687–698.
- 42 Sushila, R. Siddiqui, S. Patra, K. Shivam, A. Sil, B. Guchhait, H. Tian, R. Kataria, S. Goswami, P. Venugopalan and R. Patra, Halogen Bond Mediated Self-Assembly of Mononuclear Lanthanide Complexes: Perception of Supramolecular Interactions, Slow Magnetic Relaxation, and Photoluminescence Properties, *Inorg. Chem.*, 2022, 61, 11484–11496.
- 43 D. A. Gálico, R. Marin, G. Brunet, D. Errulat, E. Hemmer, F. A. Sigoli, J. O. Moilanen and M. Murugesu, Triplet-State Position and Crystal-Field Tuning in Opto-Magnetic Lanthanide Complexes: Two Sides of the Same Coin, *Chem. Eur. J.*, 2019, 25(64), 14625–14637.
- 44 C. Ma and C. Besson, Precise control of the degree and regioselectivity of functionalization in nitro- and amino-

- functionalized di(trispyrazolylborato)iron(II) spin crossover complexes, *Dalton Trans.*, 2021, **50**, 18077–18088.
- 45 C. H. Hossack, R. J. Butcher, C. L. Cahill and C. Besson, Structural Diversity of Lanthanide 3-Nitrotrispyrazolylborates: Tunable Nuclearity and Intra-Ligand Charge Transfer Sensitization of Visible and NIR Ln3+ Emission, *Inorg. Chem.*, 2021, 60, 15724–15743.
- 46 C. Apostolidis, J. Rebizant, B. Kanellakopulos, R. von Ammon, E. Dornberger, J. Müller, B. Powietzka and B. Nuber, Homoscorpionates (hydridotris(1-pyrazolyl) borato complexes) of the trivalent 4fions. The crystal and molecular structure of [(HB(N2C3H3)3]3LnIII, (Ln = Pr, Nd), *Polyhedron*, 1997, **16**, 1057–1068.
- 47 G. Crosby, R. Whan and R. Alire, Intramolecular energy transfer in rare earth chelates. Role of the triplet state, *J. Chem. Phys.*, 1961, 34, 743–748.
- 48 W. Sager, N. Filipescu and F. Serafin, Substituent effects on intramolecular energy transfer. I. Absorption and phosphorescence spectra of rare earth β-diketone chelates, *J. Phys. Chem.*, 1965, 69, 1092–1100.
- 49 K. Binnemans, Interpretation of europium(III) spectra, *Coord. Chem. Rev.*, 2015, 295, 1–45.
- 50 J. G. Santos, J. D. L. Dutra, S. A. Junior, R. O. Freire and N. B. da Costa Junior, Theoretical Spectroscopic Study of Europium Tris(bipyridine) Cryptates, *J. Phys. Chem. A*, 2012, 116, 4318–4322.
- 51 A. De Bettencourt-Dias, in *Luminescence of Lanthanide Ions* in *Coordination Compounds and Nanomaterials*, 2014, pp. 1–48, DOI: 10.1002/9781118682760.ch01.
- 52 P. A. Tanner, Some misconceptions concerning the electronic spectra of tri-positive europium and cerium, *Chem. Soc. Rev.*, 2013, 42, 5090–5101.
- 53 C.-L. Liu, R.-L. Zhang, C.-S. Lin, L.-P. Zhou, L.-X. Cai, J.-T. Kong, S.-Q. Yang, K.-L. Han and Q.-F. Sun, Intraligand Charge Transfer Sensitization on Self-Assembled Europium Tetrahedral Cage Leads to Dual-Selective Luminescent Sensing toward Anion and Cation, J. Am. Chem. Soc., 2017, 139, 12474–12479.
- 54 K. D. Matthews, S. A. Bailey-Folkes, I. A. Kahwa, G. L. McPherson, C. A. O'Mahoney, S. V. Ley, D. J. Williams, C. J. Groombridge and C. A. O'Connor, Luminescence Dynamics and ¹³C NMR Characteristics of Dlnuclear Complexes Exhibiting Coupled Lanthanide(III) Cation Pairs, J. Phys. Chem., 1992, 96, 7021–7027.
- 55 K. Binnemans and C. Görller-Walrand, On the color of the trivalent lanthanide ions, *Chem. Phys. Lett.*, 1995, 235, 163–174.
- 56 E. Bauer and M. Rotter, in *Properties and Applications of Complex Intermetallics*, 2009, pp. 183–248, DOI: 10.1142/9789814261647_0005.
- 57 N. F. Chilton, R. P. Anderson, L. D. Turner, A. Soncini and K. S. Murray, PHI: A powerful new program for the analysis of anisotropic monomeric and exchange-coupled polynuclear d- and f-block complexes, *J. Comput. Chem.*, 2013, 34, 1164–1175.
- 58 D. Reta and N. F. Chilton, Uncertainty estimates for magnetic relaxation times and magnetic relaxation parameters, Phys. Chem. Chem. Phys., 2019, 21, 23567–23575.

Dalton Transactions Paper

- 59 A. Lunghi and S. Sanvito, The Limit of Spin Lifetime in Solid-State Electronic Spins, *J. Phys. Chem. Lett.*, 2020, 11(15), 6273–6278.
- 60 M. Bortoluzzi, G. Paolucci, S. Polizzi, L. Bellotto, F. Enrichi, S. Ciorba and B. S. Richards, Photoluminescence studies on europium-based scorpionate-complex, *Inorg. Chem. Commun.*, 2011, 14, 1762–1766.
- 61 G. F. de Sá, O. L. Malta, C. de Mello Donegá, A. M. Simas, R. L. Longo, P. A. Santa-Cruz and E. F. da Silva, Spectroscopic properties and design of highly luminescent lanthanide coordination complexes, *Coord. Chem. Rev.*, 2000, 196, 165–195.
- 62 S.-N. Liu, K.-N. Tong, Y. Zhao, J.-F. Cheng, M.-K. Fung and J. Fan, Efficient red phosphorescent Ir(III) complexes based on rigid ligands with high external quantum efficiency and low efficiency roll-off, *J. Mater. Chem. C*, 2020, **8**, 6168–6175.
- 63 J.-C. Jin, N.-N. Shen, Y.-P. Lin, L.-K. Gong, H.-Y. Tong, K.-Z. Du and X.-Y. Huang, Modulation of the Structure and Photoluminescence of Bismuth(III) Chloride Hybrids by Altering the Ionic-Liquid Cations, *Inorg. Chem.*, 2020, **59**, 13465–13472.
- 64 J.-C. Jin, Y.-P. Lin, Y.-H. Wu, L.-K. Gong, N.-N. Shen, Y. Song, W. Ma, Z.-Z. Zhang, K.-Z. Du and X.-Y. Huang, Long lifetime phosphorescence and X-ray scintillation of chlorobismuthate hybrids incorporating ionic liquid cations, *J. Mater. Chem. C*, 2021, **9**, 1814–1821.
- 65 A. C. Marwitz, A. D. Nicholas, L. M. Breuer, J. A. Bertke and K. E. Knope, Harnessing Bismuth Coordination Chemistry to Achieve Bright, Long-Lived Organic Phosphorescence, *Inorg. Chem.*, 2021, 60, 16840–16851.
- 66 C. Apostolidis, A. Kovács, A. Morgenstern, J. Rebizant and O. Walter, Tris-{Hydridotris(1-pyrazolyl)borato}lanthanide Complexes: Synthesis, Spectroscopy, Crystal Structure and Bonding Properties, *Inorganics*, 2021, 9(6), 44.
- 67 SAINT, version 8.34a, Bruker AXS Inc.: Madison, WI, 2013.
- 68 G. Sheldrick, Bruker AXS Inc.: Madison, WI, 2005.
- 69 G. Sheldrick, A short history of SHELX, *Acta Crystallogr.*, *Sect. A: Found. Crystallogr.*, 2008, **64**, 112–122.
- 70 C. B. Hubschle, G. M. Sheldrick and B. Dittrich, ShelXle: a Qt graphical user interface for SHELXL, J. Appl. Crystallogr., 2011, 44, 1281–1284.

- 71 A. Spek, Platon/Squeeze, *Acta Crystallogr.*, *Sect. D: Biol. Crystallogr.*, 2009, **65**, 148–155.
- 72 M. J. Frisch, G. W. Trucks, H. B. Schlegel, G. E. Scuseria, M. A. Robb, J. R. Cheeseman, G. Scalmani, V. Barone, G. A. Petersson, H. Nakatsuji, X. Li, M. Caricato, A. V. Marenich, J. Bloino, B. G. Janesko, R. Gomperts, B. Mennucci, H. P. Hratchian, J. V. Ortiz, A. F. Izmaylov, J. L. Sonnenberg, D. Williams-Young, F. Ding, F. Lipparini, F. Egidi, J. Goings, B. Peng, A. Petrone, T. Henderson, D. Ranasinghe, V. G. Zakrzewski, J. Gao, N. Rega, G. Zheng, W. Liang, M. Hada, M. Ehara, K. Toyota, R. Fukuda, J. Hasegawa, M. Ishida, T. Nakajima, Y. Honda, O. Kitao, H. Nakai, T. Vreven, K. Throssell, J. A. Montgomery Jr., J. E. Peralta, F. Ogliaro, M. J. Bearpark, J. J. Hevd, E. N. Brothers, K. N. Kudin, V. N. Staroverov, T. A. Keith, R. Kobayashi, J. Normand, K. Raghavachari, A. P. Rendell, J. C. Burant, S. S. Iyengar, J. Tomasi, M. Cossi, M. Millam, M. Klene, C. Adamo, R. Cammi, J. W. Ochterski, R. L. Martin, K. Morokuma, O. Farkas, J. B. Foresman and D. J. Fox, Gaussian 16 Rev. A.03, 2016.
- 73 C. Lee, W. Yang and R. G. Parr, Development of the Colle-Salvetti correlation-energy formula into a functional of the electron density, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1988, 37, 785–789.
- 74 A. D. Becke, Density-functional thermochemistry. III. The role of exact exchange, *J. Chem. Phys.*, 1993, **98**, 5648–5652.
- 75 M. Dolg, H. Stoll, A. Savin and H. Preuss, Energy-adjusted pseudopotentials for the rare earth elements, *Theor. Chim. Acta*, 1989, 75, 173–194.
- 76 R. Gulde, P. Pollak and F. Weigend, Error-Balanced Segmented Contracted Basis Sets of Double-ζ to Quadruple-ζ Valence Quality for the Lanthanides, *J. Chem. Theory Comput.*, 2012, **8**, 4062–4068.
- 77 M. Dolg, H. Stoll and H. Preuss, A combination of quasirelativistic pseudopotential and ligand field calculations for lanthanoid compounds, *Theor. Chim. Acta*, 1993, 85, 441–450.
- 78 F. Weigend and R. Ahlrichs, Balanced basis sets of split valence, triple zeta valence and quadruple zeta valence quality for H to Rn: Design and assessment of accuracy, *Phys. Chem. Chem. Phys.*, 2005, 7, 3297–3305.
- 79 Jmol: an open-source Java viewer for chemical structures in 3D, https://www.jmol.org/.Journal.