



Cite this: *Chem. Soc. Rev.*, 2023, 52, 6417

## Comparison and assessment of methods for cellulose crystallinity determination

Khandoker Samaher Salem, \*<sup>ab</sup> Nitesh Kumar Kasera, †<sup>ac</sup>  
 Md. Ashiqur Rahman, <sup>ad</sup> Hasan Jameel, <sup>b</sup> Youssef Habibi, <sup>e</sup> Stephen J. Eichhorn, <sup>f</sup>  
 Alfred D. French, <sup>g</sup> Lokendra Pal \*<sup>b</sup> and Lucian A. Lucia \*<sup>bhii</sup>

The degree of crystallinity in cellulose significantly affects the physical, mechanical, and chemical properties of cellulosic materials, their processing, and their final application. Measuring the crystalline structures of cellulose is a challenging task due to inadequate consistency among the variety of analytical techniques available and the lack of absolute crystalline and amorphous standards. Our article reviews the primary methods for estimating the crystallinity of cellulose, namely, X-ray diffraction (XRD), nuclear magnetic resonance (NMR), Raman and Fourier-transform infrared (FTIR) spectroscopy, sum-frequency generation vibrational spectroscopy (SFG), as well as differential scanning calorimetry (DSC), and evolving biochemical methods using cellulose binding molecules (CBMs). The techniques are compared to better interrogate not only the requirements of each method, but also their differences, synergies, and limitations. The article highlights fundamental principles to guide the general community to initiate studies of the crystallinity of cellulosic materials.

Received 17th February 2023

DOI: 10.1039/d2cs00569g

[rsc.li/chem-soc-rev](http://rsc.li/chem-soc-rev)

<sup>a</sup> Department of Applied Chemistry and Chemical Engineering, University of Dhaka, Dhaka-1000, Bangladesh. E-mail: samaher.salem@du.ac.bd

<sup>b</sup> Department of Forest Biomaterials, North Carolina State University, Raleigh, NC, USA. E-mail: lpal@ncsu.edu, lalucia@ncsu.edu

<sup>c</sup> Department of Biological and Agricultural Engineering, North Carolina State University, Raleigh, NC, USA

<sup>d</sup> National Institute of Textile Engineering and Research, University of Dhaka, Dhaka-1000, Bangladesh

<sup>e</sup> Sustainable Materials Research Center (SUSMAT-RC), University Mohamed VI Polytechnic (UM6P), Lot 660, Hay Moulay Rachid, Benguerir, 43150, Morocco

<sup>f</sup> Bristol Composites Institute, School of Civil, Aerospace, and Mechanical Engineering, University of Bristol, Bristol, BS8 1TR, UK

<sup>g</sup> United States Department of Agriculture, Agricultural Research Service, Southern Regional Research Center USDA ARS SRRRC, New Orleans, LA 70124, USA

<sup>h</sup> Department of Chemistry, North Carolina State University, Raleigh, NC 27695-8204, USA

<sup>i</sup> State Key Laboratory of Biobased Materials & Green Papermaking, Qilu University of Technology/Shandong Academy of Sciences, Jinan, 250353, P. R. China

† Author KSS and NKK have equally contributed to the conceptualizations, writing, drawing figures and editing.



**Khandoker Samaher Salem**

cellulosic materials for barrier application and hygiene products, such as paper towel, tissue paper, and synthesis of biomaterials by upcycling leather waste and their valorization through sustainable applications.

*Khandoker Samaher Salem is Assistant Professor of Department of Applied Chemistry and Chemical Engineering at the University of Dhaka, Bangladesh. He also maintains his appointments as a visiting scholar in the Department of Forest Biomaterials, North Carolina State University. He completed his doctoral degree from the Forest Biomaterials, North Carolina State University. His research interests focus on surface modification of ligno-*



**Nitesh Kumar Kasera**

*Nitesh Kasera completed his PhD in Biological and Agricultural Engineering from North Carolina State University in 2022. His research interests include modifying the properties of biopolymers and biochars for environmental applications with an emphasis on structure–activity relationships. Nitesh is a graduate of Applied Chemistry and Chemical Engineering from the University of Dhaka, Bangladesh. He achieved the 2nd runner-up position in the Boyd Scott Graduate Research Competition at the 2021 Annual International Meeting of the American Society of Agricultural and Biological Engineers for a part of his doctoral research. Nitesh will be joining as a Post-Doctoral Fellow in Auburn University's Biosystems Engineering Department starting in August 2023.*



## 1. Introduction

Societal awareness and governmental regulations are steering the world towards a sustainable bio-economy through optimized exploitation of natural resources.<sup>1–5</sup> Thus, biomaterials based on cellulose, the most abundant global biopolymer and one of the most important polysaccharides, are gaining significance and importance.<sup>6,7</sup> Cellulose occurs primarily as the major component of plant cell walls in a hierarchy of structures (Fig. 1A). It is present as a high molecular weight linear-chain polymer containing three hydroxyl groups per glucose residue unit.<sup>8</sup> An extended shape arises from the preferred torsion angles about the C<sub>1</sub>–O and O–C<sub>4</sub> bonds of the  $\beta$ -(1 → 4)-glycosidic linkage.<sup>9</sup> These linear (*i.e.*, unbranched) macromolecules

comprise 1000–10 000  $\beta$ -glucose residues (the correct descriptor of the repeating unit of cellulose is a glucose residue, which is often confused with anhydroglucose and sometimes with a cellobiose repeating unit).<sup>10–13</sup> The shape of the cellulose chains allows them to closely approach each other and their hydroxyl groups form extensive intra- and intermolecular hydrogen bond networks. Further, the most common shape, that of a flat ribbon, also permits van der Waals forces and C–H···O hydrogen bonds to form.<sup>14</sup> Together, these forces result in various ordered crystalline arrangements (Fig. 1B).<sup>15</sup> The hydrogen bonds augment the stiff and straight chain nature of the cellulose molecules. Another critical factor is the presence of the C6 primary alcohol group adjacent to the linkage.<sup>16</sup>



**Youssef Habibi**

*Professor Youssef Habibi holds a dual PhD in Polymer Chemistry from Joseph Fourier University and University of Cadi Ayyad. He joined the University Mohammed VI Polytechnic (UM6P) as Chair in Sustainable Materials, is a member of the ACS, a Fellow of the RSC, and a Fellow of the IAAM. His research focuses on new bio-derived polymers and (nano)fillers, biomaterials development, and high-performance lignocellulosic nanocomposites, biomass conversion, recycling, and novel analytical tools. He authored over 120 articles, is a Global Highly Cited Researcher, authored over 20 book chapters, edited one book, and has an H-index of 54 with over 22 000 citations.*



**Stephen J. Eichhorn**

*Stephen Eichhorn is Professor of Materials Science and Engineering at the University of Bristol, within the Bristol Composites Institute. He has research interests in the physical properties of cellulosic materials, with particular emphasis on Raman spectroscopy. He was awarded the Swinburne Medal in 2020 from the Institute of Materials, Minerals and Mining (IOM3), the Hayashi Jisuke Prize in 2017, and the Rosenhain Medal and Award from IOM3 in 2012 for his research. Professor Eichhorn is also a Fellow of the Royal Society of Chemistry.*



**Alfred D. French**

*Alfred Dexter French is the son of Dexter French, who pioneered X-ray diffraction studies of starch and carbohydrate enzymology. Alfred is unofficially an Emeritus Research Scientist and serves as Editor-in-Chief of the journal *Cellulose*. Al is known for his work on molecular modeling of the building blocks of polysaccharides and the connection between the models and observed crystal structures. His recent work has been on powder XRD of cellulose using the Rietveld approach. He has served as the Chair of both the ACS Divisions of Cellulose and Renewable Materials and of Carbohydrate Chemistry & Chemical Glycobiology.*



**Lokendra Pal**

*Lokendra Pal is the E. J. Woody Rice Professor in the Department of Forest Biomaterials at NC State University. He has over 22 years of global experience in the research and technological advancement of bio-based materials. He has been granted 24 patents. Pal received a PhD in Paper and Imaging Science and Engineering from Western Michigan University. He is the NC State University Faculty Scholar and TAPPI Fellow and received the University Outstanding Teacher Award. His current research is focused on sustainable materials research and advanced manufacturing innovation, utilizing green and sustainable chemistry, engineering, and artificial intelligence/machine learning.*



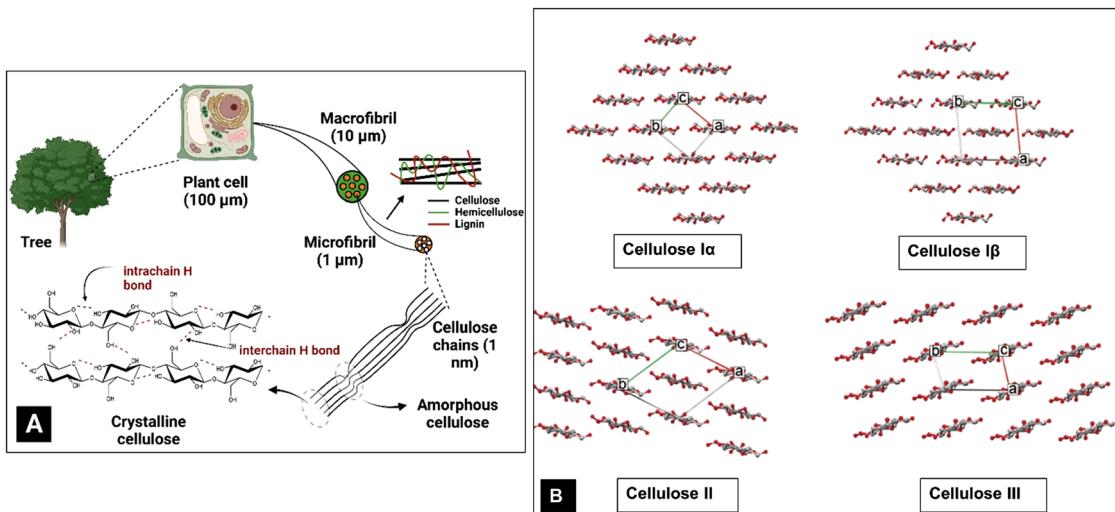


Fig. 1 (A) Hierarchical structure of cellulose and (B) different crystalline forms of cellulose.

The forces also organize the chains in parallel arrangements of crystallites, which are the elements of the secondary structure of native cellulose fibrils and cellulose fibres.<sup>11</sup> Cellulose fibres are typically composed of nanofibrils with lengths of several micrometers and diameters ranging from 2 to 20 nm, where the cellulose crystallites are linked to each other by amorphous cellulose domains.<sup>11</sup> Eight different crystalline allomorphs, referred to as cellulose I $\alpha$ , I $\beta$ , II, III $_1$ , III $_{11}$ , IV $_1$ , IV $_{11}$  and X, have been identified using different techniques, with I $\alpha$  and I $\beta$  being the naturally occurring or native crystalline forms. Cellulose II has also been discovered in rare examples of algal<sup>17,18</sup> and bacterial cellulose<sup>19,20</sup> but more typically results from treatments with strong NaOH in the laboratory<sup>21,22</sup> (mercerization) or dissolution and regeneration, as in the industrial production of rayon. The other forms are mostly of

scientific interest, although there is commercial interest in cellulose III $_1$  which results from treatment of cellulose I with amines.

Carl von Nägeli first discovered the ordered arrangement of cellulose fibrils in 1858 using polarized light microscopy,<sup>23</sup> which was later confirmed by X-ray diffraction (XRD) patterns of ramie<sup>24</sup> in 1926. This was more than a decade after the first diffraction pattern of cellulose was recorded in 1913, just a year after the technique itself was discovered.<sup>25–27</sup> The 1926 paper proposed the polymeric nature of cellulose, but had 1,1 and 4,4 linkages instead of the 1,4 linkages known today. A correction to the linkages soon followed along with changes in the unit cell axes. This resulted in a structure with the correct basic nature, but left un-resolved issues of whether the two molecular chains of the unit cell had parallel or antiparallel packing. Another question left to the 2000s era was the orientation of the rotatable C6-OH group. The cellulose molecules aggregate to adopt crystalline fibrils with well-defined shapes, sizes and crystal allomorphs.<sup>26</sup> The nature of these crystal structures, their interaction with the surrounding molecules and their further assembly into higher order structures to form long slender fibres profoundly affect the overall properties of cellulosic materials. Therefore, an understanding of cellulose crystallinity is one of the most important requisites because crystallinity affects the physical, mechanical, chemical, and morphological properties of cellulose and its surface interaction with the surroundings and hence its use in a wide range of applications including, but not limited to, biomedicine, textiles, and electronics.<sup>28–34</sup>

The crystallinity index (CI) is commonly used to measure the relative amount of the crystalline fraction of cellulosic materials and to quantify their modification following a variety of physico-chemical and biological treatments.<sup>8,31,35–40</sup> This parameter is based on a two-phase model of the material – namely a crystalline phase and an amorphous phase. This may or may not be the most correct understanding. An alternative would be to



*Lucian Lucia is Professor and Director of The Laboratory of Soft Materials & Green Chemistry at NC State University. He maintains appointments in Biomedical Engineering, Chemistry, and Forest Biomaterials, while serving as faculty in the Fiber & Polymer Sciences program. He received his PhD in physical organic chemistry from the University of Florida with Professor Kirk Schanze, served as a postdoctoral fellow at the NSF*

*Center for Photoinduced Charge Transfer at the University of Rochester with Professor David Whitten, and is a Fellow of the Royal Society of Chemistry. His research includes biopolymer functionalization, green chemistry reactions, and sustainable materials characterization.*

**Lucian A. Lucia**

consider crystallinity based on the size, not perfection, of the crystals in the sample. Smaller crystallites would be more like amorphous material in the sense that they have a larger proportion of surface molecules than large crystallites and therefore faster reactivity. Rather than having just two distinct phases, amorphous and crystalline, the actual situation could be more of a continuous variation in the molecular order. In naturally deposited cellulose structures, this is a somewhat less attractive view, as the biosynthetic mechanisms are thought to generally create the fibrils very repetitiously.

The CI (often called CrI) is simply the percentage of the crystalline phase that is present. In native cellulose, the non-crystalline or amorphous materials are thought to be located on the surfaces of the cellulose elementary fibres as well as between the crystalline regions as indicated in Fig. 1A. The primary evidence for the interruption of the crystalline regions along the elementary fiber axes is the “leveling-off degree of polymerization”.<sup>41</sup> Namely, hydrolysis in 2 N HCl at 80 °C rapidly reduces the molecular weight until it stops dropping so rapidly. Also, the crystallite length is much shorter than the initial molecular weight.<sup>42</sup> Some NMR studies have indicated that the surface molecules on the cellulose crystallites have different structural details from those in the interior, and the surface molecules are expected to be less ordered based on computational molecular dynamics studies.

Several techniques have been used to measure the CI, including X-ray diffraction (XRD),<sup>21,26,36,43–51</sup> nuclear magnetic resonance (NMR) spectroscopy,<sup>44,52–57</sup> Fourier transform (FTIR) infrared spectroscopy,<sup>44,45,49,55,58,59</sup> Raman spectroscopy,<sup>55,60</sup> sum frequency generation (SFG) vibration spectroscopy,<sup>61,62</sup> differential scanning calorimetry (DSC)<sup>63</sup> and enzymatic cellulose binding molecules (CBMs).<sup>64</sup> However, a common feature of all these techniques is that the crystal disorder and complex heterogeneous nature of the components in lignocellulosic biomass lead to issues in interpretation for most established methods estimating cellulose crystallinity.<sup>26,60</sup>

The amorphous and crystalline structures in cellulose have been found to govern its chemical and enzymatic reaction kinetics.<sup>65–67</sup> Therefore, the measurement of CI has become crucial to ascertain cellulose structure, which consequently would help to design experimental conditions to achieve the desired properties of the final products for different applications. In light of this, our present review assesses the available techniques to measure cellulose crystallinity and establishes basic governing principles of XRD, NMR, IR, Raman, SFG spectroscopy, DSC thermogram, and CBM methods and their applications for cellulose structure characterization. The first part performs an individual analysis of each technique, including working principles, capabilities, and limitations. The latter section provides a comprehensive comparison among techniques and elaborates on the idea of the lack of an absolute method for measuring cellulose crystallinity. It also discusses the use of diffractograms and spectroscopic techniques complementing each other to measure cellulose crystallinity. In all cases, the review introduces fundamental principles and the application of different techniques for the measurement of cellulose

crystallinity. These enabling techniques can help us better evaluate the structure of cellulosic materials, but they also have their limitations and pitfalls in applications if not undertaken with judicious care and attention.

## 2. X-Ray diffraction (XRD)

### 2.1. XRD fundamentals

The phenomenon of diffraction by crystals is based on the concept that a crystal effectively has infinite repetition of a small geometric unit in three dimensions. A small parallelepiped, the unit cell, contains all the atoms and their three-dimensional arrangement to permit the crystal to be generated along three linear axes by simple translation. It is not required to have intact atoms or molecules within the unit cell. For many crystalline substances, the atoms and molecules only become whole when the unit cell is repeated. In Fig. 1B, for example, the unit cells, indicated by the white, red, and green lines, could be cut out from the paper and moved along the unit cell boundary lines to create the rest of the crystal. A simple inorganic or metallic unit cell may only have atoms at the corners of the unit cell, but organic molecules generally will not be able to have such a simple arrangement. These unit cells have, over the years, been described under several conventions as to whether the angles among the axes at the cell origin are obtuse or acute and which axis, *x*, *y*, or *z*, should correspond to the molecular axis in the case of a polymer like cellulose. Our review will consider the *z*-axis (the *c* unit cell dimension) as the molecular axis, with the shorter of the other two axes being *x*, and all three axes being related by the “right-hand rule”. Crystal symmetry also plays a role; the triclinic *P*1 cell used herein is as published by Nishiyama *et al.* for cellulose I $\alpha$ .<sup>68</sup> The remaining space groups for neat cellulose are all monoclinic *P*2<sub>1</sub> (*P*112<sub>1</sub> indicates the unique *c*-axis) with the obtuse monoclinic angle  $\gamma$  being between the *a*- and *b*-axes.

In the earliest decades of diffraction crystal structure studies, it was not properly understood that a polymer could fit in the small unit cells indicated by the diffraction data. For example, the longest dimension of a cellulose unit cell is approximately a nanometer, while a polymer of the glucose molecule would be many times that long. The quandary was resolved by understanding that the observed unit cell dimensions along the molecular axis were the predominantly repeated unit and there were so few ends, that the additional H atom and OH group at the ends of the molecule would not be detected by the experiment.

Once a proper context of the unit cell is established, the concept of crystal planes can be discussed. These planes are useful for the identification of the various peaks on a diffraction pattern. For example, the (200) plane intersects the *a*-axis at one half its length; the “2” in “200” represents the reciprocal of one half. The reciprocals of the other two “Miller indices” are infinity, indicating that the (200) plane never intersects the *b*- or *c*-axis. Another important plane in cellulose crystallinity is often the (110) plane, which intersects the *a*-axis one unit away



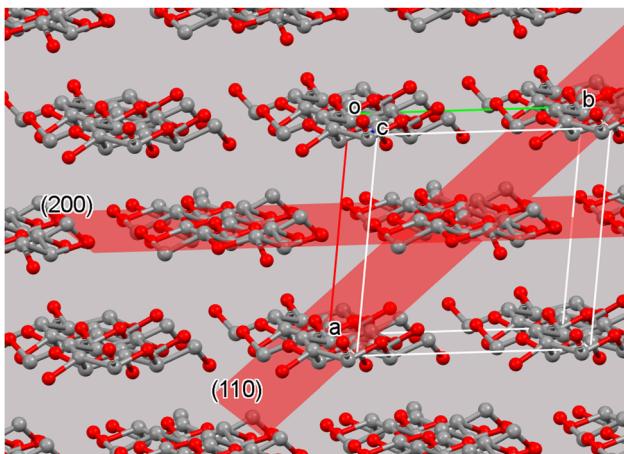


Fig. 2 The cellulose I $\beta$  unit cell and surrounding molecules with the crystal (200) and (110) planes. The planes extend the entire dimensions of the crystal and are found for each unit cell. The chemically repeated unit of cellulose is the glucose residue. In neat cellulose crystals the two-fold (or pseudo-two-fold) symmetry results in a crystallographic repeat distance that includes two glucose residues. Other cellulose structures have different repeat distances.<sup>12</sup>

from the origin of the unit cell. It also intersects the  $b$ -axis at one unit away in the negative direction from the origin (Fig. 2). The (110) and (110) planes are equivalent, but the convention is to list the negative Miller index value on the  $b$  term. Another important plane in cellulose, the (110) is shown in Fig. 2.

X-Ray diffraction (XRD), being a powerful non-invasive analytical technique for characterizing crystalline materials, can provide information on structures (3-D spatial representations of the molecular structure), phases, crystal orientations, crystallinity, and other structural parameters such as average grain

size, strain, and crystal defects. The very first applications of XRD techniques used polychromatic X-rays which produced 'smeared' diffraction patterns.<sup>27</sup> This issue was later resolved by using monochromatic sources that resulted in distinct diffraction peaks from what are termed 'planes of reflection' in Fig. 2. Although the peaks result from diffraction planes, the Miller indices for planes are enclosed in parentheses, but the peaks of a diffraction pattern are not. That convention is often ignored, however. Distinct XRD peaks are produced by coherent interference of a monochromatic beam of X-rays which are scattered by atoms in periodic lattice planes in a sample.<sup>69</sup> The atomic positions within the unit cell contribute to the peak intensities – often called the lattice and molecular transforms. It is also important to note that the diffraction pattern observed is then a combination of two transforms – the lattice and the molecular transforms. Therefore the intensities contain information about the relative atomic positions and thus about the shapes of the molecules at those lattice positions. Only crystal lattices can have planes that in turn have Miller indices; it is never appropriate to refer to the Miller indices of amorphous materials.

The diffraction of X-rays by crystal planes shows how lattice spacings are formed in the plane using Bragg's law ( $n\lambda = 2d_{hkl} \sin \theta$ ) as illustrated in Fig. 3A. Here,  $n$  is an integer called the order of the reflection,  $\lambda$  is the wavelength of the X-rays used,  $d$  is the characteristic spacing between the crystal planes of a given specimen and  $\theta$  is the incident angle of light (X-ray beam). By measuring the angles,  $\theta$ , under which the constructively interfering X-rays leave the crystal, the interplanar spacings,  $d_{hkl}$ , of these planes can be determined. In the case of cellulose powder diffraction, however, the peaks from the various planes are mostly overlapped and such determinations are more difficult. Fig. 3B shows the schematic of a typical XRD

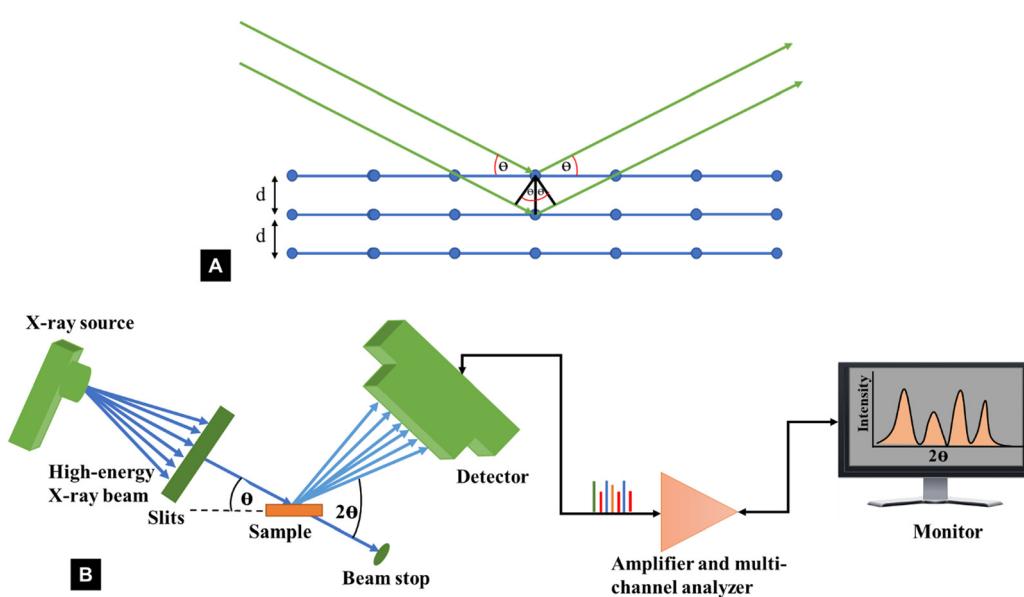


Fig. 3 Schematic representation of (A) Bragg's law and (B) XRD instrumentation set up typically used for powder diffractometer determinations of the crystallinity index of cellulose.

Bragg-Brentano (reflection) set up used for a powder diffractometer, where X-ray diffraction results from planes that satisfy the Bragg condition.

Again, in the first forty years of X-ray diffraction studies, the X-ray peaks were detected on photographic film. If the sample was a bundle of aligned fibres, the recorded pattern was a two-dimensional array (a “fibre pattern”) that furnishes more information about the structure. See, for example, the diffraction patterns for tunicate cellulose or highly crystallized cellulose II published by Nishiyama, Langan, and Chanzy and by Langan, Nishiyama, and Chanzy.<sup>56,70</sup> For many samples, however, especially nanocellulose and wood pulp, it is more convenient to use powder diffraction. Unfortunately, that results in a substantial loss of information, particularly as it is important to capture as many intensities as possible to get an accurate structural analysis. The XRD part of our review is focused on powder diffraction with a powder diffractometer, which is distinct from flat plate or 2D diffraction studies. (2D data of various types can be converted to 1D data typical of powder diffractometers with software.)

Since the discovery of XRD, its non-invasive nature and relative accessibility have made it a ubiquitous technique for studying the crystallinity of cellulosic materials.<sup>2,15,64,71–77</sup> The diffraction pattern for cellulose displays several intensities typically superimposed on a broad background. The “broad background” is composed of environmental scatter (air scatter, sample holder, shot noise, <sup>78</sup> etc.). There is also inelastic or Compton scattering that goes into the background as well as other wavelengths that get through the filter that removes K beta radiation. However, XRD requires a certain minimum crystallite size to consider a fraction in a solid as being crystalline, in which the width of the peaks in the diffractograms is known to be inversely proportional to the lateral width of the crystallites.<sup>79,80</sup> Thus, in cellulosic materials, the crystalline peaks get broader for smaller sized crystallites.<sup>69</sup> The material’s surface also affects the crystallinity in XRD analysis. Though the surface atoms have greater mobility, and a lower degree of order compared to the bulk of the crystallites, they are captured as crystalline material during the XRD measurements. This is also the case for minor defects within crystalline regions that are too small to be detected as regions of molecular disorder.<sup>71</sup>

So, it is worth noting at this point that not all disorder can be accurately determined for cellulosic samples, if the scale of this disordered material is sufficiently small.

As previously mentioned, the crystallinity index (CI) has been used to describe the relative amount of crystalline material in cellulose. It is calculated as the contribution of the crystalline portion of the material relative to the combined areas of the crystalline and amorphous regions.<sup>55</sup> Several XRD techniques have been proposed to determine the CI. In this sense, comparing literature data for cellulose crystallinity is challenging due to the lack of standard methods. However, the most widely applied XRD methods for studying cellulose crystallinity are Segal peak height, peak fitting (also referred to as peak deconvolution), and amorphous subtraction.<sup>15,71</sup> However, there is a substantial amount of Rietveld crystallinity work being currently undertaken which is very attractive because it can consider all of the inputs to a diffraction pattern in a physically relevant manner. Although intellectually satisfying, it is however noted that the Rietveld method has been widely applied to inorganic materials (e.g. ceramics and metals) that make better crystals and provide more diffraction data. Even in their application to these materials that form much larger crystals, issues regarding “accuracy” of fit, and thereby calculations of peak intensity do occur. The theoretical implications behind each method, as well as limitations, are discussed in the following sections. Comparing the different techniques will elucidate how CI values change with different methods and enable an understanding of the simplest and the most widely used XRD method to measure the CI of cellulosic materials.

**2.1.1. The Segal peak height method.** The Segal peak height method was developed by Segal and coworkers to measure the change in the crystalline content of cotton cellulose I following chemical or mechanical treatment.<sup>50</sup> This method was developed for rapid comparison of the crystallinity of cellulose samples determined using the relatively new-at-the-time X-ray powder diffractometer.<sup>36,81</sup> The height of the 200 peak located between the scattered angles of  $2\theta = 22^\circ$  to  $23^\circ$  represents the sum of the crystalline and amorphous components, and the amorphous-only component is represented in this empirical method by the intensity at the minimum of

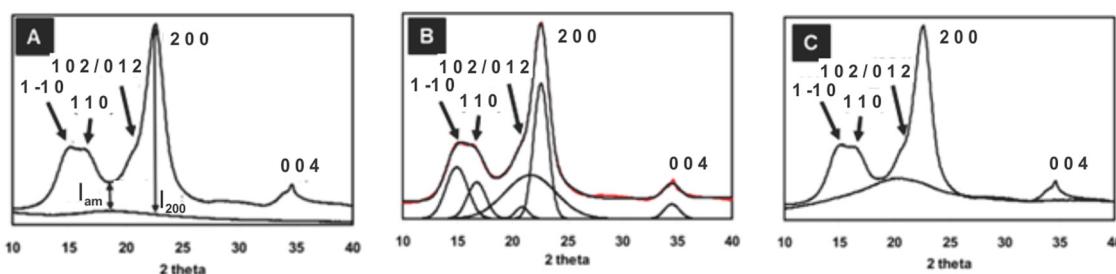


Fig. 4 The XRD pattern for Avicel PH-101 showing three methods of crystallinity determination: (A) background-corrected Segal peak height method, (B) peak fitting (peak deconvolution) method using Voigt functions for fitting five crystalline peaks and the relatively broad amorphous contribution,<sup>14,66</sup> and (C) the amorphous subtraction method where a ball milled amorphous sample is used as a reference<sup>15,71</sup> (adapted from and used with permission from ref. 15. Copyright Springer-Nature, 2010). Note: the authors of the original drawing in ref. 15 opted to use the convention of *b* as the unique axis, rather than the typical *c*-axis convention for polymer diffraction that is used here. Thus, the original peak labels were 101, 10̄1, 021 and 120, 002, and 040.



about  $18^\circ$  between the 110 and 200 peaks, shown in Fig. 4A. The reasoning was that there are no significant crystal lattice peaks expected at  $2\theta = 18^\circ$ , so all intensity must arise from amorphous scattering. Segal's exhaustively ball-milled cotton samples had a maximum at this same point, so all the intensity there was attributed to the amorphous content. The amorphous scattering, indicated in Fig. 4C, does not decrease much as it goes under the 200 peak, so the net crystalline material was indicated by the height of the 200 peak ( $I_{200}$ ) – the intensity at  $18^\circ$  ( $I_{AM}$ ). Thus, the equation of Segal *et al.*, expressed in the nomenclature of this review, is

$$\text{Crystallinity index (CrI)} = 100 \times (I_{200} - I_{AM})/I_{200} \quad (1)$$

Fig. 4A indicates a background line, and the heights of  $I_{200}$  and  $I_{AM}$  are measured from that line. In fact, Segal recommended against using a background correction and given his equipment and technology, it was perhaps a reasonable choice. Because of the variation in modern instrumentation and sample holders, it is preferable to make such a correction. Environmental background is not amorphous scattering from the sample, so it seems reasonable to us to record a blank run and subtract it from the experimental sample data. This applies to other XRD methods as well. Admittedly, the subtraction of a blank is controversial in the powder diffraction community. A sure sign that something is wrong would be finding that the intensity values of the corrected pattern sometimes fall below zero. Other methods of background correction are problematic as well. Computed routines in several diffraction software packages draw lines that connect the minima and then algorithmically pull the minima down to the line of zero intensity. These methods are completely unsuited to cellulose because the observed minima are actually the intensities of broad overlapping peaks that result from small crystals. Because of such differences in the chosen approaches, care should be taken when replicating approaches without reference to the original work. This is noted with reference to the work of Çetin *et al.*<sup>10</sup>

A major advantage of the Segal method is its simplicity. It could be easily used on the paper chart output from the early diffractometers; no computers were needed (or otherwise available). However, there are numerous criticisms of the Segal method in the literature. Despite its simplicity, some researchers misapply the Segal method by taking the  $1\bar{1}0/110$  peak (Fig. 4A) height to represent the amorphous fraction, whereas it is the trough height that must be used. Of course, that is not Segal's fault. A major intrinsic fault is that it assumes that all the intensity at the  $18^\circ$  point (modern work finds the minimum closer to  $18.6^\circ$ ) is from amorphous scattering and none of it arises from the overlap of the broad peaks that result from the small crystallite sizes of cellulose. Because of peak overlap, calculated patterns for perfect but small crystals (*e.g.*, a size similar to cotton crystallites) would only be 90% crystalline, and an experimental sample will never give a value of 0% crystallinity because  $I_{200}$  and  $I_{AM}$  would have to be identical. Another little-mentioned defect in the assumptions is that samples can suffer from non-random orientation of the

crystallites, based on the sample preparation and placement in the X-ray beam. Situations with preferred orientation of the crystallites place a particular plane with respect to the incident X-ray beam and therefore exaggerate those intensities, changing the ratio of  $I_{200}$  to  $I_{AM}$  despite identical crystallinity. Preferred orientation is especially prevalent in samples of bacterial cellulose as well as nanocellulose that is dried in various ways.

High-end diffractometers often have spinning sample holders to assist with randomization of the crystallite orientation, but they will not generally solve all such problems. For example, samples of fibers cut to very short lengths still tend to lie down in the plane of the sample holder. Spinning the sample will make sure that the pattern does not suffer from a secondary orientation of the fibres. The fibres will be directed uniformly in  $360^\circ$ . However, a third type of orientation can still arise. Namely, the crystallites might be rotated about their long axes in a particular orientation. Spinning sample holders do not deal with the first or the third type of preferred orientation. It is good practice to do this anyway for all samples since the beam size may sample regions of a supposedly random sample, but local orientation effects may occur. This approach is not only valid for the Segal method, but also for all powder XRD approaches. Of course, complete characterization of the sample would include information on the extent of preferred orientation. In any case, such effects can result in invalid crystallinity measurements if not taken into consideration, and the Segal method cannot easily deal with them.

Segal's method is empirical and was not developed to compare samples of different nature, but rather to quantify differences among a single set of samples.<sup>15,71</sup> Therefore, Segal's method is better suited to study relative changes in crystallinity resulting from exposure of a sample to different treatments and not for estimating absolute values of crystallinity and amorphous fractions in cellulosic materials.<sup>2</sup>

**2.1.2. The peak fitting, peak area, or peak deconvolution method.** Peak deconvolution accounts for the observed pattern by fitting the visible peaks of the XRD pattern plus an additional broad peak (or peaks) to account for the amorphous material, as shown in Fig. 4B.<sup>82</sup> Sometimes Gaussian,<sup>83</sup> Voigt,<sup>84</sup> or Lorentzian<sup>85</sup> peaks will be inserted at the expected positions based on the known unit cell dimensions. During optimization of the peak parameters with curve-fitting software, their sum will be adjusted to match the experimental pattern.

As many as five crystalline peaks indexed  $1\bar{1}0$ ,  $110$ ,  $012/102$ ,  $200$ , and  $004$  are often considered while using Gaussian and Voigt fit functions. The peaks indexed  $1\bar{1}0$ ,  $110$  and  $200$  are clearly visible in most cellulose samples, whereas the  $012/102$  overlapped reflection might appear as a shoulder; however, these peaks might overlap with each other and could not be resolved so distinctly for some samples, *e.g.*, bleached eucalyptus.<sup>84</sup> Another diffraction peak is often visible at approximately  $34.5^\circ$  which gives a somewhat broad peak indexed as  $004$  although it is a composite of several neighbouring peaks.<sup>47</sup> The Voigt (often the pseudo-Voigt) function has been used to fit the five peaks generated in the diffractogram of cellulose



samples. Fig. 4B is a typical example of such a fit. However, four crystalline peaks indexed  $1\bar{1}0$ ,  $110$ ,  $200$ , and  $004$  might be used when a Lorentzian fit function is applied since the  $012/102$  is difficult to resolve as it appears as a shoulder in the diffractogram. The crystallinity index is calculated from the ratio of the area of all crystalline peaks to the total area using the equation

$$CI = \frac{A_{cr}}{A_{sample}} = \frac{\left( \int_{2\theta_1}^{2\theta_2} I_{cr} d2\theta \right)}{\left( \int_{2\theta_1}^{2\theta_2} I_{sample} d2\theta \right)} \quad (2)$$

where  $CI$  is the crystallinity index,  $A_{cr}$  is the area of all the crystalline peaks,  $A_{sample}$  is the total area of the sample,  $I_{cr}$  is the integral sum of the intensities of the crystalline peaks, and  $I_{sample}$  is the sum of the intensities of all the diffraction space. The amorphous contribution is typically modeled by the insertion of a very broad peak somewhere between  $18^\circ$  and  $21^\circ$ , using one of several choices in either general peak-fitting software or software more dedicated to XRD work. In a notable effort, Nishiyama *et al.*<sup>86</sup> studied ball-milled samples of numerous cellulose and related molecules with diffraction and found extreme similarities among them. They created a Fourier series expression for the amorphous contribution that was refined with MATLAB software. That expression was then used with PeakFit software for studies of cellulose crystallinity.

Because more peaks are considered than in the Segal method, a possible advantage of the peak fitting method is that it could more safely be used to compare the crystallinity index of different samples. Concern has been expressed that the peak area method does not account for all the factors that contribute to peak broadening, including crystalline disorder, crystallite size and non-uniform strain within the crystal.<sup>84</sup> Other work had indicated that crystallite size is the dominating factor.<sup>87</sup> Another problem is that the pattern cannot be resolved into many narrow diffraction contributions as the cellulose intensities are very broad with overlapping peaks.<sup>84,88,89</sup>

While the Segal method has an intrinsic fault, *i.e.*, assuming that all the intensity at the  $18^\circ$  intensity minimum in Fig. 4A is from amorphous scattering, problems with peak deconvolution are more practical. The first is that there are many smaller peaks that are not considered. This is shown very clearly by patterns calculated assuming crystallites large enough to give very sharp peaks (see Fig. 5).<sup>47</sup> When the size is 100 nm or more the peaks that contribute to the crystalline intensity become visible but when the modeled crystallites are realistically small, these peaks blend into what is often mistaken for background or amorphous scattering. Thus, these smaller peaks in the higher  $2\theta$  range give a reasonable basis for that broad but low intensity. No observed or calculated intensity is shown in that area of Fig. 4B despite its presence in 4A. This indicates that the intensity in that region has been subtracted as background, as done in ref. 82. In that work, the authors mention that “a detailed and rigorous treatment would include a full description of the amorphous background.” As mentioned above, it is important to not equate “amorphous scattering” with “environmental and intrinsic background.” An intrinsic limitation of

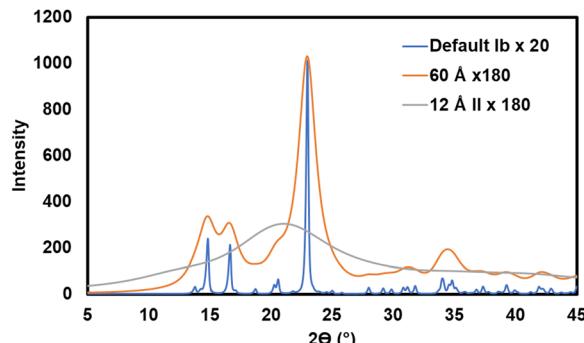


Fig. 5 Calculated diffraction patterns. The blue pattern for cellulose Ib represents the default parameters in the MAUD<sup>90</sup> Rietveld software including a nearly infinite crystal size of 1000 Å (the incident X-ray beam has been scaled by a factor of 20 to reach 1000 counts). To prevent excess noise in experimental data, at least 2000 counts should be accumulated at the maximum intensity for any report of the XRD pattern, with more being helpful for Rietveld work. (10 000 to 100 000 is not excessive with a modern XRD detector.) The orange pattern, which represents a moderately large cellulose crystal of 60 Å diameter, was scaled by a factor of 180. The grey line is a calculated pattern for cellulose II, except that it has a very small crystallite size of 12 Å, representing amorphous cellulose. It was scaled for the same incident beam intensity as the orange pattern. The area under the amorphous pattern is 88% that of the 60 Å pattern.

peak deconvolution is that there is no crystallographic science used to link the intensities of the various peaks.

**2.1.3. The amorphous subtraction method.** Introduced by Ruland, this method measures the crystallinity of the cellulose sample by deducting the diffraction of an amorphous region of a standard material from the diffraction pattern from the sample under investigation.<sup>84</sup> The amorphous standard is chosen so that it matches the crystalline contribution of the sample. Fig. 4C is a typical example of the amorphous subtraction method used to determine cellulose crystallinity where a ball milled amorphous standard sample is used as a reference. The scale of the amorphous component is adjusted so that the intensities of the amorphous and experimental patterns are equal at the  $18^\circ$  point that Segal assigned to  $I_{AM}$ .

However, if the intensity of the selected amorphous standard does not match the experimental amorphous contribution, then the fitting is allowed by some workers to exceed the experimental intensities marginally at some scattering angles to improve the fit. The area under the amorphous contribution and the total area of the samples are measured, and the crystallinity is calculated using the equation

$$CI = 1 - \frac{A_{am}}{A_{sample}} = 1 - \frac{\left( \int_{2\theta_1}^{2\theta_2} I_{am} d2\theta \right)}{\left( \int_{2\theta_1}^{2\theta_2} I_{sample} d2\theta \right)} \quad (3)$$

where  $CI$  is the crystallinity index,  $A_{am}$  is the area of all the amorphous peaks,  $A_{sample}$  is the total area of the sample,  $I_{am}$  is the integral sum of the intensities of the amorphous peaks, and  $I_{sample}$  is the sum of the intensities of all the diffraction space in the sample.

A wide scattering angle from  $2\theta_1 \sim 13^\circ$  to  $2\theta_2 \sim 50^\circ$  is used to calculate the area. The selection of an ideally amorphous



standard is very challenging which limits the use of this method.<sup>15</sup> Moreover, amorphous subtraction has the same limitation as Segal's method, which is due to assigning the intensities to the amorphous fraction without considering the presence of overlapping peaks. Nonetheless, ball-milled cellulose, regenerated cellulose, and xylan or lignin powder are widely used among several standards to fit in the measured amorphous region.<sup>40,51</sup> As seen in Fig. 4C, the experimental data above  $25^\circ 2\theta$  has been considered to be amorphous, except for the point near the 004 peak. A test that should be applied to results from both the deconvolution and the amorphous subtraction results is to plot just the pattern that results from what is considered as the crystalline material. Here, the subtraction would just be the amorphous from the experimental, and for deconvolution, it would be the sum of all the crystalline peaks at each  $2\theta$  step. If the data had been corrected for environmental background, then the removal of amorphous data from the pattern should result in an ideal pattern that can be compared with ideal patterns for crystallites of a given size. Alternatively, calculated ideal patterns can be analyzed by the selected method to check for 100% crystallinity.

**2.1.4. The Rietveld method.** The Rietveld method<sup>91</sup> is another way to determine cellulose crystallinity. Although first intended for determining the atomic positions for materials that could not grow crystals large enough for single crystal methods, it has tools to compensate for multiple phases (crystalline and amorphous) and sample texture (aspects such as crystallite size and preferred orientation). Application of these tools gives a report on the difference between the sample and an ideal structure. The analysis starts with an ideal XRD pattern based on a reliable previous determination of  $x$ ,  $y$ , and  $z$  atomic coordinates for a sufficiently similar structure.<sup>91,92</sup> The ideal way to view the Rietveld analysis is that the analyst is presented with a calculated, ideal diffraction pattern. This pattern is initially defined as having an infinite size and perfectly ordered crystals, randomly oriented particles of spherical shape, and unit cell dimensions and atomic coordinates similar to those of the sample. The experimental pattern is placed on the same  $2\theta$  coordinate axis, and the calculated pattern is adjusted by least squares fitting of variable parameters. For example, the unit cell dimensions must be varied because it is well known that the peak positions, especially for cellulose, will not be exactly correct otherwise.<sup>93</sup> Of course, the height of the pattern will need to be scaled to match the experimental data, and even if a blank run is subtracted for removal of environmental background, an additional quadratic background correction may be needed. The crystallite size is related to the width of the peaks to a first approximation by the Scherrer equation and is another obvious variable.<sup>80</sup>

There are several lessons in Fig. 5, which shows, in blue, the ideal pattern that might greet the user at the start of the Rietveld calculation. The first is that there are many small peaks in that blue pattern that blur out when the crystallite size is diminished to match an experimental pattern. Another is that a pattern with lower peak intensities is not necessarily less crystalline in terms of crystal perfection, but in this case, it took

nine times as many incident photons for the 200 peak on the  $60\text{ \AA}$  pattern to reach the same height as the  $1000\text{ \AA}$  pattern.

Different Rietveld software programs, widely available in free and commercial products, have different philosophies for handling an amorphous phase. One strategy uses a very small crystal of the same material as a model. For cellulose, it seems at present that a small crystal of cellulose I is not the best approach. Instead, both cellulose II and cellulose IV structures are used. Each can be justified. The cellulose II model is supported by the resemblance of patterns calculated from  $12\text{ \AA}$  crystallites (Fig. 5) to the ball-milled or quickly regenerated samples that are regarded as amorphous models. This is not arbitrary. Amorphous samples share the dominant orientation of O6 with cellulose II and soaking of amorphous material in water can result in cellulose II.<sup>94</sup> Additionally, nanocellulose samples can have cellulose II that is not visible on the diffraction pattern. However, if the size of the model crystal is allowed to change during refinement, it can rise to  $30\text{ \AA}$ , indicating the presence of cellulose II. The cellulose IV model is supported by the finding that the primary cell wall is often composed of that allomorph. Another approach is to use different functions that result in a pattern resembling ball-milled or regenerated cellulose, or to include an additional arbitrary broad peak as needed.<sup>92</sup> This is currently the frontier area of research.

The correct analysis of texture (crystallite size and orientation) also affects the CI (percentages of the crystalline and amorphous phases) of cellulose samples.<sup>95</sup> The Rietveld method optimizes the CI measurement by using a curve fitting method, where it considers all the variables including the lower intensity peaks as well. Thus, Rietveld is a complete peak area fitting method that does not have the limitations of the simple methods discussed above. It is worth noting that this method is well-established for more crystalline inorganic and metallic materials. More pertinently, the issues with obtaining an accurate fit to diffraction data for these more crystalline materials are also well-known, even when these materials' patterns appear to have a relative absence of pronounced amorphous backgrounds. Peak fitting itself is fraught with difficulty when it comes to diffraction patterns, especially when signal to noise and backgrounds are considered. This is a subject which has been noted in the literature but is often ignored in analysis. Another issue is on the consensus of what parameter to use to quantify the degree of fit, and the fact that an overinterpretation of small values of commonly used parameters such as the  $R$ -factor can lead to erroneous conclusions. Having reliable structural determinations on a database is key to this approach, and without due care in using these database structures, mistakes can easily be made, *e.g.*, the use of the wrong allomorph for a native cellulose, and the fact that lattice distortions (seen as shifts in peak positions) can occur for woody plant materials *versus* more crystalline forms of cellulose. An over-reliance on a good fit between the refined model and the experimental data can lead to the wrong conclusions.

The precision of the Rietveld analysis has however been improved by the Cellulose Rietveld Analysis for Fine Structure (CRAFS) software which uses a crystallographic model along



with a computational algorithm having only the variables which are deemed to be most practical for cellulose studies with the limited amount of information. The other advantage is that CRAFS as well as some other Rietveld software, can also handle 2-dimensional data where the 2-D fibre patterns, as mentioned above, can yield much more information than the 1-D patterns discussed herein.<sup>95,96</sup> The advantage of 2-D patterns is shown in ref. 90 that compares both 1-D and 2-D calculated patterns. The spots on a 2-D pattern fall on different layer lines so that they are better resolved, providing a greater number of independent data points. That work also compared the 1-D patterns for both ideal crystallites of different sizes and their counterparts that had been subjected to energy minimization and molecular dynamics studies and therefore they lost their exact periodic order. As a result of this approach, it has been observed that many of the diffraction peaks formerly credited to background scattering or amorphous cellulose are actually coming from adjacent crystalline peaks.<sup>47,92</sup> However, it is still early to say that Rietveld can produce the most accurate and reliable crystallinity index value for a broad range of cellulose allomorphs due to its capacity to consider all the crystalline peaks. However, in principle, it has much to offer if better standards are adopted and if the pitfalls encountered for more regular materials are considered. It is fair to say that the approach is in its infancy for cellulose compared to more crystalline materials, and much work is needed to ensure that robust approaches are adopted without an over-reliance on erroneous underlying assumptions.

The biggest problem with Rietveld determination is the limited amount of experimental data. Based on the structure in Fig. 6b, the variation of atomic coordinates for cellulose is futile, and simplified models of the crystallite size anisotropy must be used. Unlike single crystal diffraction studies, the ratio of data to variable parameters is not clearly defined for Rietveld work on cellulose diffraction patterns.

## 2.2. Comparison of different XRD methods

XRD is a commonly used method to characterize the crystallinity of cellulosic materials. However, each of the above-discussed methods has some limitations. One of the limitations associated with the methods is the dependence of the crystallinity index on the crystallite size.<sup>71</sup> For fully crystalline models, all the methods have been found to show a positive linear correlation with the crystallite size.<sup>71</sup> Segal CI calculations on model patterns showed a very strong, but non-linear correlation.<sup>48</sup> However, all the methods showed variation in the crystallinity index as shown in Table 1. The peak fitting method exhibits the highest variation (from the 'Maximum Difference' parameter), while the Rietveld method the least variation.<sup>71</sup>

From Table 1, it is also observed that the crystallinity index measured using all the methods is less than the ideal value. The crystallinity index measured using the Segal method is higher than that of the other methods, and this trend has been reported for several different cellulosic samples, as shown in Table 2.<sup>15,51,71,97</sup>

The Segal method yields the highest values for crystallinity indices and has been reported to show a mean error of over 20% for crystallinity values for known samples.<sup>43</sup> In contrast, the crystallinity indices obtained by Rietveld and peak fitting methods for different cellulosic species have been shown to be consistent.<sup>99</sup> The high values of the crystallinity index of the cellulosic samples obtained by the Segal method are thought to be due to several reasons. The main reason is that the peaks associated with the amorphous regions in the XRD are very broad, and a simple height comparison ignoring the variation in peak width cannot provide a correct estimate of the crystallinity index.<sup>84</sup> Besides, the trough height also varies which is not considered during the CI measurement. Therefore, when the crystallinity index is measured using the Segal method without considering the above mentioned discrepancies, it

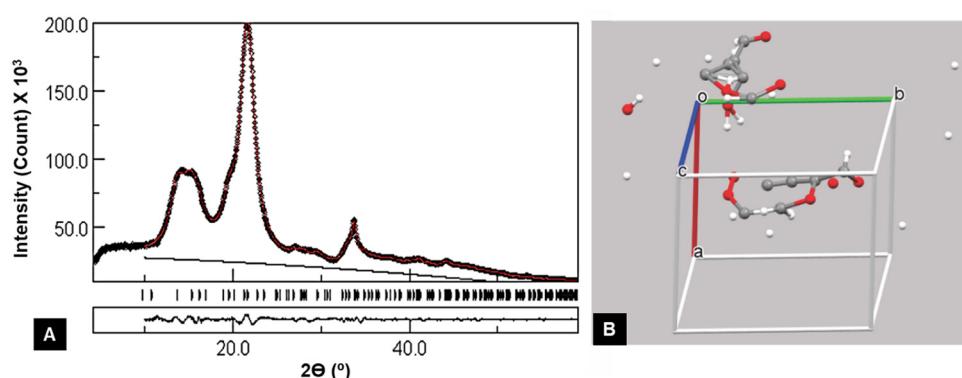
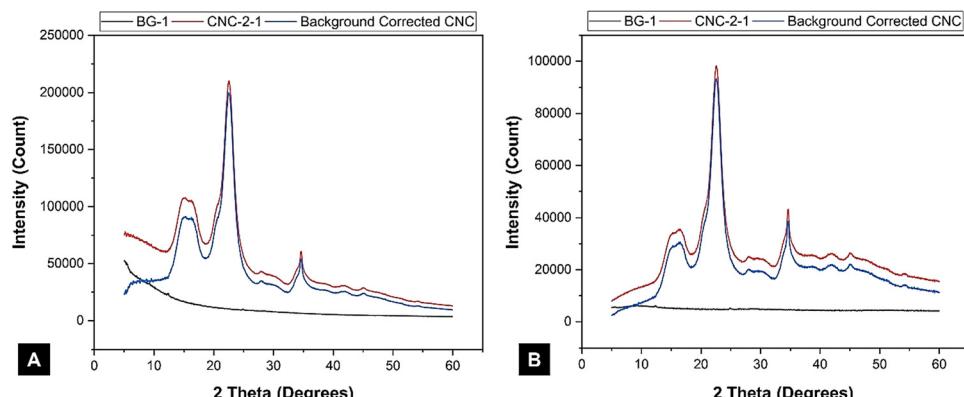


Fig. 6 Example of a badly underdetermined structure, *i.e.*, too many variables for the available data, that gave an excellent agreement between the observed and calculated intensities. (A) Output from the MAUD Rietveld software. The black circles represent the experimental data, and the red line shows the fit line for the adjusted cellulose I $\beta$  pattern. The black line is the added background needed despite the previous subtraction of environmental background obtained with a blank run (see Fig. 7a). The black and red lines show that data before  $20 = 10^\circ$  were not included in the analysis. The short vertical lines to the right of "Tuncate" indicate all of the peaks that contributed to the red calculated pattern. Output from the MAUD Rietveld software. (B) However, the "refined" structure with variable parameters for the unit cell dimensions and atomic coordinates resulted in a physically meaningless underlying structure. Unit cell and contents from the MAUD refinement that included atomic positions, unit cell dimensions, background, preferred orientation, and anisotropic crystallite size, displayed with the Mercury program. The 147 variables were far too many for the limited data and the results were non-physical.





**Fig. 7** Effect of the (A) fixed slit and (B) variable slit X-ray diffractometer on the background curve and the resulting corrected intensity of a cellulose nanocrystal (CNC) sample. Rietveld analysis software typically assumes a fixed-slit diffractometer configuration. These fixed-slit data were converted from variable-slit data with the Malvern-Panalytical HighScore software.

**Table 1** Variation of the cellulose crystallinity value based on 100% crystalline models with different measuring techniques<sup>8,71</sup>

XRD method	CI mean (%)	St. dev. (%)	Max. difference (%)
Segal peak height	~93	3.7	10.3
Peak fitting	77–84	3.5	13.3
Amorphous subtraction	80–82	4.0	10.8
Rietveld method	0–100	N/A	N/A
Ideal values	100	0	0

Maximum difference is the difference between the highest and lowest crystallinity values.

**Table 2** Crystallinity index of cellulose obtained from different sources using different XRD methods<sup>51,71,98</sup>

Sample	Crystallinity index (%)			
	Segal peak height method	Peak fitting method	Amorphous subtraction method	Rietveld method
Juniper	34–38	19–22	22–23	36
Moso bamboo	45–46	22–24	20–21	28
Hemp	75	40	49	60
Nata de coco	77	44–47	49–52	61
Spruce-pine	75–77	49–50	42–48	58
Avicel PH-102	76–82	49–53	60–65	67
Cotton linter	85–87	55–56	62–67	72.1

leads to overestimation of the intensity of the crystalline region and hence the total crystallinity index.<sup>100–104</sup>

The calculated degree of cellulose crystallinity can also be affected by the orientation of samples, which can be assessed by measuring the same sample using various diffractometer geometries, as well as the rotation of the samples already mentioned. For example, it was found that when a medium-density balsa sample was measured with a different geometry, the symmetrical transmission geometry produced a crystallinity index of 48%, whereas when the symmetrical reflection geometry of the same sample was used, the measured crystallinity index was almost 100%.<sup>2</sup> This was due to optimal scattering of the sample's 200 reflection, which resulted in overestimating the contribution from the scattering pattern.<sup>105</sup> Thus, a symmetrical

transmission geometry yielded a higher Segal crystallinity index for balsa. Another issue is the presence of "environmental background" which is independent of the sample and tends to overlap with the diffraction pattern of the amorphous scattering contributed by a low intensity region. Therefore, the background and amorphous scattering must be separated to get a more reliable and meaningful crystallinity index value.<sup>39</sup>

Despite different techniques resulting in different crystallinity indices of the cellulosic samples, a good linear correlation ( $r^2 \geq 0.9$ ) has been reported between samples when their relative difference in the crystallinity index was large.<sup>71</sup> The Segal method showed a linear correlation with other methods when a single sample set was measured; nonetheless, the linearity disappeared when various samples were measured.<sup>71</sup> This is thought to occur because the Segal method is an empirical method used to study relative changes within a single sample set. The peak fitting method yields lower crystallinity values compared to the Segal method since the peak fitting method can give unrealistic amorphous contributions if the fitting limits are too strict. Amorphous subtraction would also yield a crystallinity index founded on an erroneous assumption if an appropriate standard amorphous model cannot be used since the standard model cannot go beyond the sample intensity even if their intensities do not match in some regions of the selected scattering angle.<sup>71</sup> To overcome this limitation, occasionally amorphous fitting is used, allowing mismatch between the intensities of the standard model and the sample. But due to different factors like the amorphous model, scattering angle, background corrections, and subtraction, the amorphous fitting model cannot be used to determine absolute values, *i.e.*, whether a sample is 100% crystalline or not.

The Rietveld method is emerging as a useful tool to measure the crystallinity index of cellulose samples by XRD. The CRAFS software analyzes the 2D diffraction pattern and generates a few data such as unit cell parameters, crystallite size, crystalline index/intensity, and crystallite orientation. This method could overcome the two major limitations, such as overlapping of diffraction peaks and preferred crystallite orientation, which were issues with the other three techniques.<sup>95</sup> The increased



accuracy was possible because the Rietveld software usually can adjust the model pattern for preferred orientation of the cellulose, and some software also can compensate for crystallites that are anisotropic (they have different dimensions along different axes). This allows the Rietveld method to result in a plausibly more reliable determination of crystallinity than the Segal height, peak deconvolution, and amorphous subtraction techniques. However, issues exist in its application, and the choice of amorphous standard can change the results significantly. Scattering from various amorphous samples varies, and so a single, invariant standard might result in an incorrect estimation of the crystallinity index since the amorphous nature in the sample to be tested is unknown. Proper selection also of the standard crystallographic allomorph of cellulose is obviously a prerequisite to obtain accurate fits to diffraction data. While standard for most well-studied inorganic materials, Rietveld analysis is still in its infancy in the field of cellulose crystallinity determination and should be pursued with caution. One way to be cautious is to introduce the variables in several individual refinements rather than to try to refine all variables in just one refinement. Individual variables will tend to take unrealistic values and should be monitored when judging the success of the refinement.

Another factor that affects all the methods discussed above is the intensity of the radiation on the sample at different  $2\theta$  angles. It has been observed that the calculated lines are smooth, but the experimental lines contain noise which arises from the photon counting statistics. The noise in the data can be reduced when the sample is irradiated with a high intensity beam, *e.g.*, a rotating anode tube or synchrotron radiation, or by counting for a longer time at each step with a conventional diffractometer. The standard deviation (noise) for the photon count is proportional to the square root of the number of photons counted (the intensity). The high counts that gave very smooth experimental data in Fig. 7 were obtained with a conventional laboratory diffractometer with a modern detector in a half-hour. Moreover, using fixed-slit diffractometer, a standard method to collect the data, can produce a background with quite a radical curvature, as shown in Fig. 7A. In contrast, variable slits cause the radiation to hit the sample with different intensity which produces an almost horizontal background (Fig. 7B). Researchers should also make sure that the chosen Rietveld software is programmed for the slit type and reduce the impact of the environmental background by subtracting a blank run.

XRD techniques have been widely used to determine cellulose crystallinity, although it is very challenging to use these approaches to compare the crystallinity of different samples if identical measurement and analysis protocols are not followed. All the diffraction peaks of amorphous and crystalline regions must be considered for complete quantification and optimum assessment of cellulose crystallinity. The presence of preferred orientation can be confirmed by using different sample orientations although in principle, it is not necessary as the Rietveld software can include corrections for preferred orientation. Though any XRD technique can provide relative crystallinity

indices of a single set of samples, it might be beneficial to use the Rietveld method to determine and compare values for different sets of samples because it has the capability, if applied correctly, to produce reliable results if enough data are provided.

### 3. Nuclear magnetic resonance (NMR)

Nuclear magnetic resonance (NMR) ranks among the widely used tools to measure the crystallinity of cellulose and nanocellulose.<sup>106–112</sup> The technique makes use of many different instrumentation types that are available for the characterization of the cellulose crystallinity index.<sup>110</sup> The crystallinity index measured by NMR can result in more meaningful data than those obtained by DSC or FTIR since NMR is able to detect multiple conformations, whereas DSC and FTIR techniques cannot determine secondary structures or domain movements. In this sense, crystallinity measurement by NMR relates more meaningfully to a discrete structure of cellulose.

#### 3.1. NMR fundamentals

Nuclear magnetic resonance is based on the magnetic properties of the atomic nucleus. The resonance phenomenon occurs at a characteristic frequency (Larmor frequency) in the electromagnetic spectrum when the nuclei of certain atoms are placed in a strong magnetic field.<sup>28</sup> The nuclei of certain atoms are very sensitive to slight variations in the magnetic field shielding applied by the surrounding atoms.<sup>28</sup> Characteristic, yet slight variations occur in the electrons when the external magnetic field near the nucleus is modified.<sup>113</sup> These slight variations that occur are called a “chemical shift” and provide critical information about the molecules. The information obtained indicates the characteristics of the type of nucleus and the atom’s position within a molecule.<sup>28,113</sup> When a constant magnetic field of strength  $B_0$  is applied, the polarization of the magnetic nuclear spin splits into two energy states ( $2I + 1$ ) (Fig. 8A).<sup>28</sup>  $I$  is the magnetic spin number and  $I = 1/2$  is from the atomic nucleus, where the sum of the proton and neutron number is odd such as  $^1\text{H}$  and  $^{13}\text{C}$ . One of the states has the same magnetic moment as the  $B_0$  field, and the other state has a momentum opposite to the  $B_0$  field.<sup>114</sup> The magnetic moment, parallel to the external field, has a lower energy ( $m = +1/2$ ). In contrast, the magnetic moment with the anti-parallel alignment has higher energy ( $m = -1/2$ ).<sup>114</sup> In the presence of a magnetic field  $B_0$ , the nucleus moves in precession around the applied field at an angular velocity defined by the Larmor frequency  $\omega_0 = \gamma B_0$  (Fig. 8B).<sup>115,116</sup> The constant  $\gamma$  is called the magnetogyric ratio and has a unique value for different nuclei.<sup>116</sup> When the electromagnetic waves with frequency ( $\nu$ ) irradiate the system, the precessing nuclei can undergo a transition between the two energy states by the absorption of energy.<sup>28</sup> The nuclei in the lower energy state ( $m = +1/2$ ) are promoted to the higher energy state ( $m = -1/2$ ). Because of the precession resonance ( $\nu = \omega_0$ ) and the frequencies of the applied radiation, the transition process is defined as nuclear magnetic resonance (NMR).<sup>116</sup> Based on



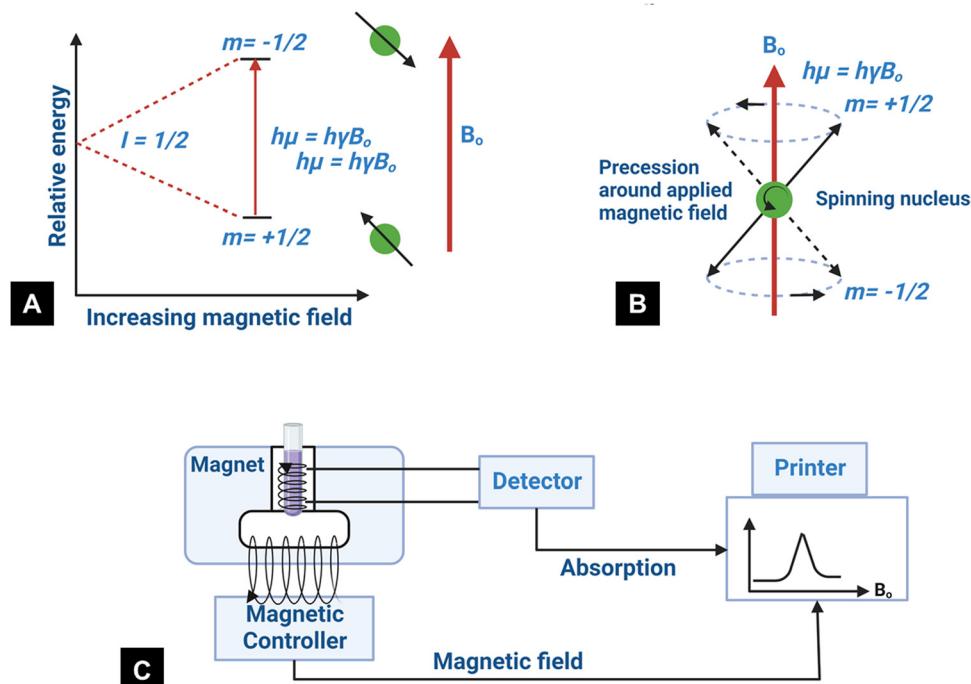


Fig. 8 (A) State of nuclear spin under the influence of the external magnetic field, (B) graphical representation of the precession for spin-up and spin-down states at the Larmor frequency (redrawn from ref. 117), and (C) schematic of the NMR technique.<sup>118</sup>

the fundamentals of NMR, when a magnetogyric ratio with a specific frequency is used to pulse the sample to a high energy state, the atomic nucleus in the sample can give a unique signal for the analytical spectra. The schematic of typical NMR instrumentation is shown in Fig. 8C.

### 3.2. NMR to measure crystallinity of cellulose

Cellulose exhibits limited solubility and complex chemical bonds due to its crystallinity. Thus, NMR can be used for studying cellulose crystallinity in addition to providing information about bonding, functionality, and the chemical environment.  $^{13}\text{C}$  CP/MAS NMR is the summation of three techniques, namely, cross-polarization, magic angle spinning, and high-power coupling.<sup>119</sup> Cross polarization increases the  $^{13}\text{C}$  intensity (as it is low because of the low natural cluster and magnetogyric ratio) by transferring magnetization from the proton to the carbon. The adaptation of the proton–carbon polarization (CP) method has been found to solve the weak signal problem by increasing the  $^{13}\text{C}$  sensitivity and annihilating the dipolar interactions. Moreover, magic angle spinning (MAS), which spins the sample about the alleged magic angle, has solved the issue of anisotropy of solid samples in the magnetic field.<sup>28</sup> Using  $^{13}\text{C}$  CP/MAS solid-state NMR, cellulose derived from variable processes such as kraft pulp, low-DP acid-hydrolyzed cellulose, and tunicate-extracted specimens exhibiting highly crystalline structures can be analyzed. These materials have been widely studied by NMR and their crystallinity is measured by calculating the variation in the chemical shift.<sup>52,120–122</sup>

Several different NMR techniques are used to study cellulose crystallinity. However, one of the most widely used NMR methods is the C4 peak separation of the cellulose molecules

using  $^{13}\text{C}$  CP/MAS NMR spectroscopy.<sup>54,120,121,123,124</sup> In this case, doublets of C4 peaks represent two different signals in the  $^{13}\text{C}$  spectrum, and each peak is assigned to each cellulose structure. One of the peaks, located at  $\sim 89$  ppm, indicates the C4 carbon in ordered cellulose structures in the interior of the crystallites, and another, located at  $\sim 84$  ppm, reveals the nature of amorphous cellulose.<sup>120,121</sup> Amorphous cellulose has been assigned by the region of amorphous C4 of cellulose.<sup>52</sup> The reason is that the surface material of cellulose crystallites appears as a sharper doublet in the amorphous region, and xylan deposited on cellulose crystallites appears as a peak at 82.1 ppm.<sup>125,126</sup> In addition, crystalline mannan has been reported to resonate in that area.<sup>127</sup> The crystallinity index (CI) of cellulose can be obtained by dividing the area of the crystalline peak ( $\sim 89$  ppm) by the total area of the C4 peaks,<sup>75</sup> where the  $\sim 89$  ppm and  $\sim 84$  ppm signals represent the crystalline and amorphous (disordered) regions respectively. The crystallinity index (CrI) as a percentage can be calculated by the sum of the integral area below these signals according to the equation

$$\text{CrI} = \frac{a}{a+b} \times 100\% \quad (4)$$

where  $a$  is the integral area of the crystalline region and  $b$  is the integral area of the amorphous region. Mustapha *et al.* investigated cellulosic samples obtained from cellulose pulp, regenerated cellulose pulp, and bacterial cellulose by  $^{13}\text{C}$ -CP/MAS NMR.<sup>128</sup> Spectral deconvolution of the C4 region of untreated sugarcane bagasse showed that the crystalline region possessed a narrow fitting of  $\alpha$  at 90 ppm,  $\beta$  at 86 ppm, a mixture of  $\alpha + \beta$  at



89 ppm, and a paracrystalline region at 88 ppm.<sup>129</sup> An amorphous region was found from an intense characteristic doublet with two peaks at 84 and 85 ppm.<sup>130</sup>

The <sup>13</sup>C-CP/MAS NMR spectra of different cellulose samples have been used to understand the cellulose crystallinity by the spectral analysis of chemical shifts of different polymorphs as shown in Fig. 9A.<sup>131</sup>

In the spectrum of the sample, two outer and inner signals at 105 ppm represents the cellulose I<sub>β</sub> and I<sub>α</sub> are apparent, respectively.<sup>125,133</sup> The presence of the I<sub>α</sub> allomorph is confirmed by up-field wings in the C4 signals of the fibres (see Table 3).

A sharper resonance of C1 and C4 signals represents the crystalline region that overlaps with the up-field wings of the less ordered regions, having two different environments – one is the crystallite surface and the other is the amorphous regions.<sup>133</sup> However, NMR spectral intensities play a valuable role in determining crystallinity, as the signals from different nuclei reduce the magnetic moment induction due to the application of a radio frequency pulse during the measurements. As the proton–carbon polarization (CP) dynamics and proton relaxation are different for different phases, the total intensities of the spectrum do not necessarily represent the actual number of nuclei under the spectrum. To overcome this situation, and get reliable results from CP/MAS NMR, the time constant is maintained in the order  $T_{\text{CH}} \ll T_{\text{CP}} \ll T_{1\text{pH}}$ , where  $T_{\text{CH}}$ ,  $T_{\text{CP}}$ , and  $T_{1\text{pH}}$  are respectively proton carbon (CP) time constants, CP time intervals in the process, and proton

Table 3 Signal intensities of cellulose allomorphs in the <sup>13</sup>C-CP/MAS NMR spectra of bamboo fibres<sup>88</sup>

Allomorphs	Chemical shift (ppm)		
	C1	C4	C6
Amorphous	—	86.3–80.1	63.5–59.5
Cellulose I <sub>β</sub>	105.6, 104.0	89.4, 88.2	65.4, 64.9
Cellulose I <sub>α</sub>	105.2	90.0, 89.4	—
Paracrystalline	104.8	88.8	—

spin-lattice relaxation times in the rotation frame. Thus, to avoid errors due to intensity distortion, the time constant is kept in the order  $25T_{\text{CH}} \ll 5T_{\text{CP}} \ll T_{1\text{pH}}$ .<sup>132,134</sup> The effect of the nature of the intensity of NMR spectra of microcrystalline cellulose (with various CP times up to 20 ms) on its crystallinity is shown in Table 4.

The calculation has been carried out following the logarithmic scale described for the high-resolution NMR of solids, using the equation<sup>135</sup>

$$\beta(t) = \frac{1}{1-\gamma} \left( 1 - e^{-\frac{(1-\gamma)}{T_{\text{CH}} t}} \right) e^{-\frac{1}{T_{1\text{pH}} t}} \quad (5)$$

where  $\gamma$  is the magnetogyric ratio,  $T_{\text{CH}}$  is the proton carbon time constant, and  $T_{1\text{pH}}$  is the proton spin-lattice relaxation time in the rotation frame.

Another limiting factor is the proton–carbon (CP) polarization time. It has been found that a longer CP time is responsible for the increase in signal intensity difference.<sup>132</sup> The <sup>13</sup>C

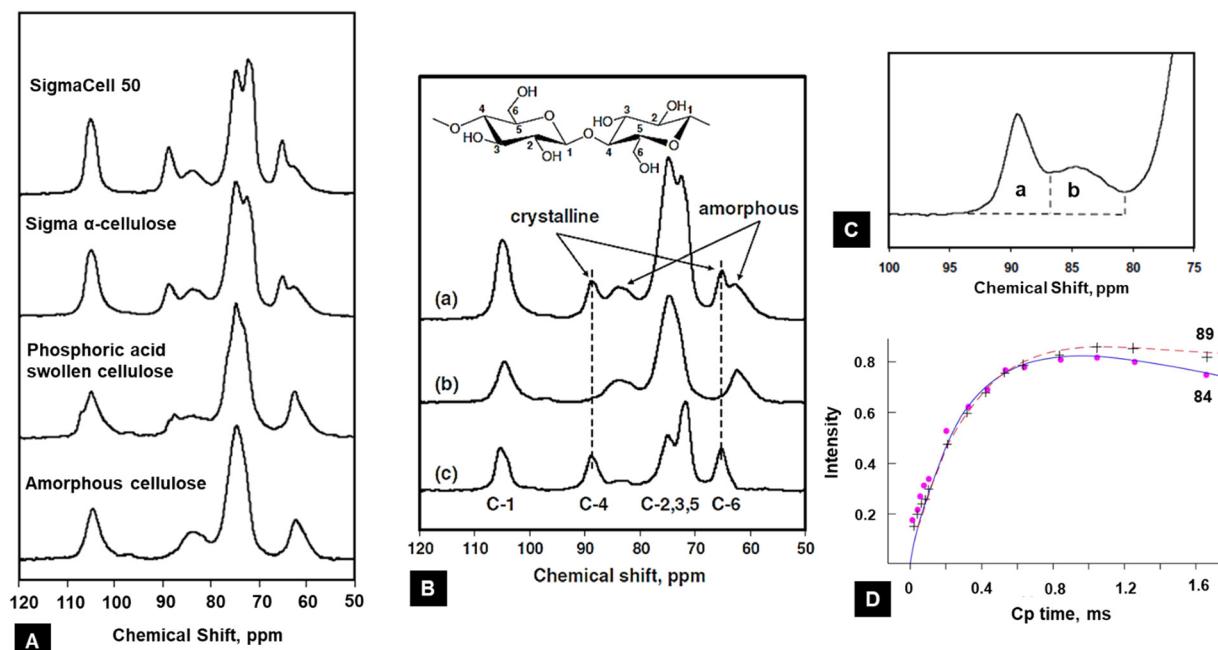


Fig. 9 (A) NMR method to calculate the crystallinity of cellulose from different sources.<sup>131</sup> (B) Example of the procedure of the subtraction method using <sup>13</sup>C NMR spectra: (a) sigma α-cellulose, (b) 100% amorphous cellulose, and (c) crystalline fraction<sup>131</sup> (reprinted with permission from ref. 131. Copyright Springer-Nature 2009). (C) Magnified version to show the peaks assigned to C4 in cellulose<sup>15</sup> (adapted from and used with permission from ref. 15. Copyright Springer-Nature, 2010). (D) <sup>13</sup>C CP NMR spectral signal intensities of 84 and 89 ppm lines upon CP time for microcrystalline cellulose. Curves are normalized to equal spin temperature and the number of protons at  $t = 0$ <sup>132</sup> (reprinted with permission from ref. 132. Copyright Springer-Nature, 1987).



Table 4 Effect of the NMR spectral intensity on cellulose crystallinity

Spectral lines	$T_{1pH}$	$T_{CH}$
Crystalline	11.8	0.31
Amorphous	6.4	0.29

NMR intensities have been determined using eqn (5), and the  $^{13}\text{C}$  NMR signal intensities vary with different CP times as shown in Fig. 9D. The data signify no intensity distortion below 0.5 ms, but there is an increasing signal intensity difference between  $\sim 89$  ppm and  $\sim 84$  ppm above 0.5 ms CP time (Fig. 9D). Therefore, the CP time must be below 0.5 ms for reliable intensity characterization. The determination of spectral lines is often carried out using a curve fitting process, much the same as for XRD analysis. Curve fitting has been carried out using Gaussian and Lorentzian functions by the variation of the shapes of spectral lines. The crystallinity has then been determined by considering the line located at  $\sim 89$  ppm to the simulated C4 total intensity line.<sup>131</sup>

### 3.3. Other methods to measure cellulose crystallinity using NMR

The crystallinity index can also be determined from the average lateral microfibril size ( $L$ ) or crystallite size as it depends on the peak intensities of C4 signals of NMR spectra by using the following equations:

$$\text{CrI} = \left( \frac{n-2}{n} \right)^2 \quad (6)$$

$$L = 0.57n \quad (7)$$

where CrI is the crystallinity index, the average lateral microfibril size is  $L$ , 0.57 refers to the average inter-chain distance (in nm),  $n$  represents the number of sides of microfibrils per chain, and  $(n-2)$  is the total number of glucose units.<sup>136</sup>

Another methodology based on  $^{13}\text{C}$  NMR spectroscopy is the subtraction of a standard for amorphous cellulose.<sup>137</sup> For this method, a fully amorphous cellulose sample is prepared to estimate the signal that might be expected of the amorphous regions of the cellulose. The amount of crystalline cellulose in this sample can be obtained by subtracting the area of the amorphous cellulose from the original signal. After the subtraction, a scale factor is applied to the amorphous cellulose spectrum to eliminate a negative signal in the residual spectrum, as shown in Fig. 9(B and C). The CI is finally acquired by calculating the area ratio between the crystalline contribution and the total area.<sup>131</sup> Fig. 9(B and C) depicts the subtraction method used to calculate cellulose crystallinity. The spectrum of the  $\alpha$ -cellulose (Fig. 9B(a)) and 100% amorphous cellulose (Fig. 9B(b)) was measured, and crystallinity fractions were calculated by subtracting the amorphous value from the  $\alpha$ -cellulose (Fig. 9C(c)). After subtraction, a small intensity peak located at  $\sim 84$  ppm showed that the ordered crystallites present in the amorphous region had been normalized.<sup>138</sup> The peak located at  $\sim 89$  ppm represents the crystallites of the internal

crystalline region. The area under these two peaks is calculated, and the crystallinity index is calculated using eqn (4).

It was determined that amorphous C4-carbons are more than likely located at distances less than about 1 nm from the crystalline C4-carbons and that those amorphous C4-carbons are not solely localized to the microfibril surface.<sup>139</sup> Contradictory NMR spin-diffusion results were published by Fernandes *et al.* that seem to support a core–shell model for the cellulose microfibril (isolated from spruce wood), where the core is more ordered than the shell.<sup>140,141</sup>

### 3.4. Comparison of different NMR methods

One of the major concerns with the peak separation technique is that it only deals with the C4 carbon signals and the separation in energy between the two spin states. This signal strongly depends on the content of lignin and hemicellulose and the external magnetic field strength, the latter of which is relatively weak.<sup>131</sup> The presence of residual lignin and hemicellulose causes overlapping peaks in  $^{13}\text{C}$  NMR spectra, which makes it harder to separate the signals into ordered and disordered peaks when explicitly using the C4 peak separation technique. Specifically, the NMR technique relies on the size of the crystallites because only the cellulose within the interior of the crystallites is regarded as crystalline.<sup>44,52</sup> The larger the crystallites, the more intense the NMR peak appears.

Moreover, ambient Boltzmann thermal energy ( $E \sim kT$ ) is sufficient to promote large portions of nuclei into a less favorable higher energy state.<sup>115</sup> For example, the Boltzmann distribution predicts that for one million hydrogen nuclei in a magnetic field of  $B_0 = 1.4$  T, only a few could contribute to the NMR signal. This explains the limitation of the sensitivity of NMR. Therefore, NMR needs a significant sample size (usually more than 100 mg for most MWs).<sup>28,122</sup> NMR measurements also yield relatively weak signals that are averaged over a long time with enough signal-to-noise ratio to produce spectra. The subtraction method has some advantages over the peak separation method in measuring the crystallinity index of cellulose as this method can be used for determining this parameter for a wide range of materials.

## 4. Vibrational spectroscopy methods

Along with other widely used analytical techniques such as XRD and NMR for the characterization of cellulose structure, vibrational spectroscopy methods (infrared and Raman) are also useful for investigating cellulose-like polymer structures and properties.<sup>142–146</sup> A vibrational spectrum is obtained due to the absorption or emission of electromagnetic radiation at specific frequencies that correlate to the energy states of the molecules, which are the quantum mechanical states of electron waves confined in specific chemical bond coordinates.<sup>28</sup> However, there are some fundamental differences in the mechanisms that give rise to IR and Raman intensities. The former relies on an existing dipole between the atoms, and the latter is due to a dipole that is induced by radiation.



#### 4.1. Vibrational spectroscopy fundamentals

The basic principle of molecular vibrations responsible for the characteristic bands observed in infrared and Raman spectroscopy can be derived by comparing a di-atomic chemical bond with a simple harmonic oscillator model (Fig. 10A).<sup>117</sup> In this model, a chemical bond can be considered as a spring with its ends attached to two atomic species with masses  $m_1$  and  $m_2$  separated by a distance  $d$ , having a spring constant  $k$ . Hooke's law can be used to express the chemical potential near the equilibrium bond distance.<sup>106</sup> After solving Schrödinger's wave equation, the wave functions of electrons involved in this chemical bond yield discrete vibration energy levels, which are expressed as

$$E_v = \frac{h}{2\pi} \left( \frac{k}{\mu} \right)^{\frac{1}{2}} \left( v + \frac{1}{2} \right) \quad (8)$$

where  $h$  is Planck's constant,  $\mu$  is the reduced mass ( $m_1 m_2 / (m_1 + m_2)$ ),  $v$  is the vibration quantum number and  $k$  is the spring constant. This equation links the force constant between two atoms (or molecular fragments), the mass of the interacting atoms (or molecular fragments), and the vibration frequency.<sup>28,147</sup>

IR spectroscopy is the measure of the intensity of a beam of electromagnetic radiation, usually a laser, observed through the absorption of radiation whose specific oscillating radiation frequency matches the natural frequency of a particular normal vibrational mode, *i.e.*, vibrating dipole of a molecule. The spectrum is obtained by plotting the intensity (absorbance or

transmittance) *versus* the wavenumber or frequency, which is proportional to the energy ( $E$ ) difference between the ground and the excited vibrational states, mainly from the  $v = 0$  state to the  $v = 1$  state, which can be expressed as

$$\omega_{\text{IR}} = E_{v=1} - E_{v=0} \quad (9)$$

where  $\omega_{\text{IR}}$  is the energy of the IR photon which is equal to the energy difference between two vibrational energy levels,  $v = 1$  and  $v = 0$ . For energy to be transferred from the IR photon to the molecule *via* absorption, the molecular vibration must cause a net change in the molecule's dipole moment or a functional group.<sup>28,148</sup> Initially, dispersive instruments were used to obtain infrared spectra. However, in more recent times, Fourier-transform infrared spectrometers have replaced the dispersive instruments as they have dramatically improved the acquisition of infrared spectra. This approach is based on the idea that the interference of radiation between two beams yields an interferogram, a signal that is produced as a function of the change of path length between the two beams. The radiation generated from the source is passed through an interferometer to the sample before reaching a detector. Upon amplification of the signal, the data are converted to a digital form by an analog-to-digital converter and transferred to the computer for Fourier transformation, a method in which two domains of distance and frequency are interconvertible.<sup>147</sup> The basic components of an FTIR spectrometer are shown schematically in Fig. 10B.

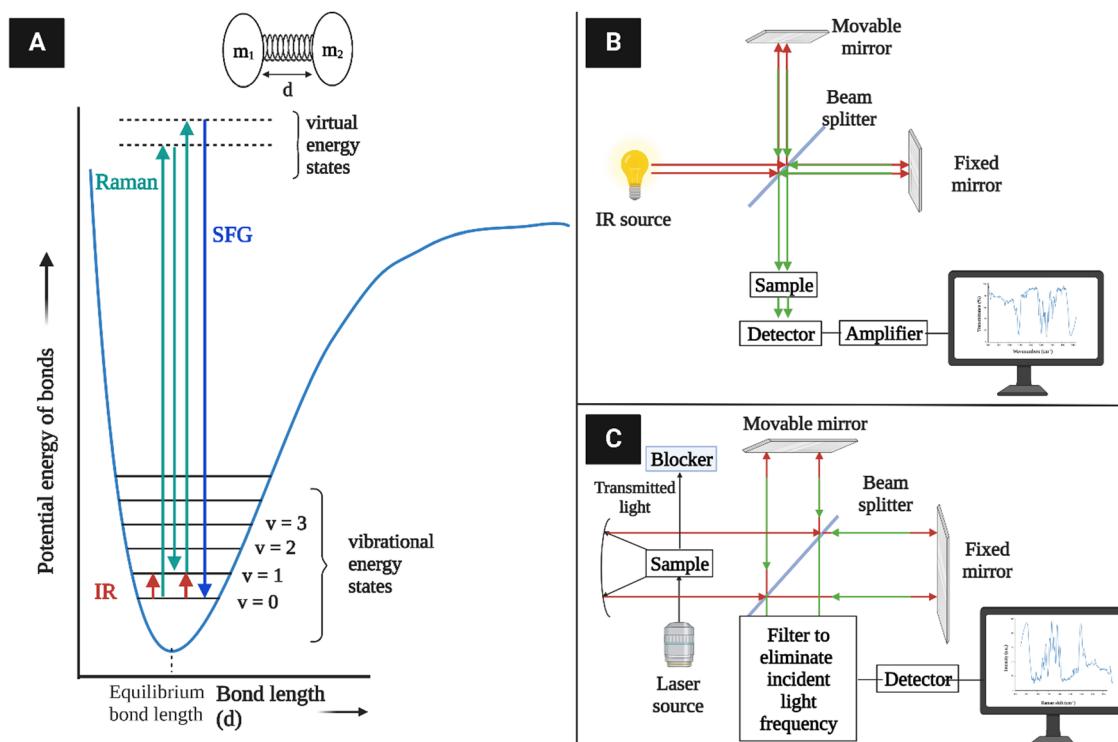


Fig. 10 (A) Vibration energy states of diatomic molecules and electronic transitions involved in IR absorption and Raman scattering. The diatomic molecules can be modeled as two masses ( $m_1$  and  $m_2$ ) connected with a spring constant  $k$ <sup>117</sup> (redrawn from ref. 117). (B) A simplified schematic of typical IR instrumentation. (C) A Raman spectrometer setup.



As the excited electrons tend to revert to the ground electronic state after being excited to a virtual electronic state, most of them come back to the vibrational ground state ( $\nu = 0$ ) in elastic scattering, but a small fraction returns to the  $\nu = 1$  state in inelastic scattering. The latter case is known as Raman scattering, where the energy difference between the originally absorbed photon (excitation) and the emitted photon (scattering) corresponds to the vibration energy level difference according to the equation

$$\omega_{\text{Raman}} = \omega_{\text{excite}} - \omega_{\text{scatter}} \quad (10)$$

The fundamental requirement for Raman activity is a net change in bond polarizability during the electronic transition upon the scattering of electromagnetic radiation.<sup>28,149</sup> The basic components of an FT-Raman spectrometer are shown schematically in Fig. 10C. According to the Raman selection rule, the polarizability change will be the result of the displacement of the constituent atoms from the equilibrium positions due to molecular vibrations.<sup>150</sup> A Raman spectrometer consists of a light source, monochromator, sample holder, and detector. Raman spectroscopy methods like dispersive Raman spectroscopy and Fourier transform Raman spectroscopy differ in their laser sources and detection methods. Several types of lasers with discrete monochromatic light sources are typically available, *e.g.*, krypton ion (530.9 and 647.1 nm), He:Ne (632.8 nm), Nd:YAG (1064 nm and 532 nm), argon ion (488.0 and 514.5 nm), diode laser (630 and 780 nm), *etc.* It is also possible to now have tunable laser systems that cover the full range from ultraviolet (UV) to near infrared (NIR). FT-Raman typically uses a 1064 nm near-IR (NIR) excitation laser, which typically results in reduced fluorescence compared to visible wavelength lasers.<sup>151</sup> This is important with respect to cellulose since fluorescence is a common problem with plant materials that may contain lignin. Also, given that cellulose has a relatively large band-gap<sup>152</sup> (much larger than the energy of the laser) it does not produce resonant Raman effects, which are more typically seen for materials such as carbon nanotubes.<sup>153</sup> With this in mind, long acquisition times are often needed to obtain spectra with low signal to noise ratios. Approaches to overcoming fluorescence in other polymers have been the use of photobleaching, *i.e.* long exposure under the laser, and also chemical bleaching using hydrogen peroxide for cellulose fibres.<sup>154</sup> The latter of these may impact the properties of the material and should be implemented with caution. Beam damage under the laser is also an issue for many fibres,<sup>155</sup> so a significantly lower power should be used for acquisition of spectra.

#### 4.2. Measurement of crystallinity of cellulose substrates using IR spectroscopy

According to the traditional two-phase cellulose model, the polymeric chains contain both crystalline (ordered) and amorphous (less ordered) regions. A clear relationship between the interaction of hydroxyl groups and crystallinity in cellulose has been established by using FTIR.<sup>156</sup> An empirical “crystallinity index” has been developed using the absorption ratio at  $1429 \text{ cm}^{-1}$  and  $893 \text{ cm}^{-1}$ .<sup>157</sup> As already discussed, the FTIR

absorption band located at  $\sim 1430 \text{ cm}^{-1}$  is assigned to asymmetric  $\text{CH}_2$  bending vibration, and is known as the “crystallinity band.” On the other hand, the FTIR absorption band located at  $\sim 898 \text{ cm}^{-1}$  is assigned to C–O–C stretching of  $\beta$ -(1 → 4)-glycosidic linkages and has been designated as the “amorphous” absorption band. It has been found that crystallinity decreases as the intensity of the absorption at  $\sim 1429 \text{ cm}^{-1}$  decreases, whereas the intensity of the absorption at  $893 \text{ cm}^{-1}$  increases.<sup>158–160</sup> Nelson and O’Connor introduced another ratio index, “total crystalline index,” which is the ratio of the intensities at  $\sim 1372 \text{ cm}^{-1}$  and  $\sim 2900 \text{ cm}^{-1}$ . Bands centered near  $\sim 2900 \text{ cm}^{-1}$  are C–H stretching vibration bands, which conform to the characteristics of amorphous cellulose, and the absorbance band at  $\sim 1372 \text{ cm}^{-1}$  is attributed to C–H bending.<sup>161</sup> It has been reported that the band located at  $\sim 1336 \text{ cm}^{-1}$  showed the most significant differences between crystalline cellulose II and amorphous cellulose. Later it was demonstrated that these bands were not characteristic of the amorphous regions.<sup>161,162</sup> Both of these ratios are now used to measure the crystallinity of cellulose where these indices are assigned as

$$\text{Cr.}R_1 = \frac{A_{1372}}{A_{2900}} \quad (11)$$

$$\text{Cr.}R_2 = \frac{A_{1430}}{A_{893}} \quad (12)$$

where  $\text{Cr.}R_1$  and  $\text{Cr.}R_2$  are respectively the absorbance ratios of cellulose  $A_{1372}/A_{2900}$  and cellulose  $A_{1430}/A_{893}$  which are also used to measure relative cellulose crystallinity. Additionally, the ratio  $A_{1370}/A_{670}$  has been invoked for investigating the transformation of cellulose I to cellulose II.<sup>163</sup> The origins of the bands used in these determinations are shown in Table 5.<sup>44</sup> To calculate the crystallinity in the pulp after being produced, the following equation has been proposed:

$$\text{Yield corrected crystalline cellulose (\%)} = A_{1370}/A_{670} \times \text{amount of cellulose} \times \text{pulp yield} \times 0.1 \quad (13)$$

where the multiplication factor (0.1) is included to give numerical values for crystalline cellulose. Thus, it has been seen that the proportion of crystalline cellulose in pulp fibres increases as the yield decreases. This is attributed to the removal of lignin and hemicelluloses from the pulp. Such determinations give only relative crystallinity due to the lack of a 100% crystalline standard and the poor resolution of the region of the spectrum for OH groups.<sup>44</sup>

Hydrogen bond intensity (HBI) can also be used to interpret qualitative changes in cellulose crystallinity; for example, crystallinity decreases with the increase of HBI. An increase in HBI represents an increase in hydrogen bonding between certain hydroxyl functions in the cellulose, which is typical of the conversion of cellulose I to cellulose II; even though this represents a decrease in the overall crystallinity index.<sup>45</sup>

FTIR spectra in attenuated total reflectance (ATR) mode offer a complementary method for providing information on chemical compounds and hydrogen bonding characteristics to



Table 5 Peak assignments in infrared spectra of cellulose

Peak (cm <sup>-1</sup> )	Type of vibration
1370	CH bending
1335	OH in-plane bending
1320	CH <sub>2</sub> wagging
670	OH out-of-plane bending
893	Non-symmetric out-phase ring
1429	CH <sub>2</sub> symmetric bending
2900	CH stretching

confirm XRD findings.<sup>164</sup> FTIR spectra of freeze-dried bacterial cellulose samples in attenuated total reflectance (ATR) mode have been used to obtain crystallinity indices by using the two conventional methods: (1) ratios of the absorption intensities at wavenumber positions of  $\sim 1428$  and  $\sim 897$  cm<sup>-1</sup> ( $Cr_1 = A_{1428}/A_{897}$ ), and (2) ratios of the absorption intensities at wavenumber positions of  $\sim 1372$  and  $\sim 2898$  cm<sup>-1</sup> ( $Cr_2 = A_{1372}/A_{2898}$ )<sup>165</sup> (Fig. 11).

#### 4.3. Measurement of crystallinity of cellulose substrates using Raman spectroscopy

Raman or its more advanced form, FT-Raman spectroscopy, has emerged as a rapid and convenient analytical method to determine cellulose crystallinity. The principle of the Raman method is to use spectral features whose intensity, bandwidth, and position are affected by cellulose crystallinity.<sup>58</sup>

The FT-Raman-based quantitation method for cellulose I crystallinity is developed using the weak bands located at  $\sim 1462$  and  $\sim 1481$  cm<sup>-1</sup> (CH<sub>2</sub> bending modes) in conjunction

with spectral deconvolution. However, the intensities of the selected bands are relatively low, and the deconvolution process is not free of the band fitting problem; hence a better approach has been sought.<sup>55</sup> In this better approach, the crystallinity of cellulose I samples has been determined based on the intensity ratio of the bands located at  $\sim 380$  and  $\sim 1096$  cm<sup>-1</sup>.<sup>58</sup> The peak height was calculated by choosing a minimum intensity wavenumber position near the peak (e.g.,  $\sim 358$  and  $\sim 944$  cm<sup>-1</sup> for the  $\sim 380$  and  $\sim 1,096$  cm<sup>-1</sup> bands, respectively) and by drawing a horizontal line (from that wavenumber position) under the peak. This is called the baseline method and using this method, in the remaining spectra for bands located at  $\sim 380$  and  $\sim 1096$  cm<sup>-1</sup>, the same values of band minima (at  $\sim 358$  and  $\sim 944$  cm<sup>-1</sup>) were used for drawing the baselines. After taking the peak heights, the intensity data were used to calculate Raman band intensity ratios, and regression models were developed. To generate a calibration curve, cellulosic samples were ball-milled for different time-spans to observe a change in their crystallinity. Significant effects on the peak shape and intensity at those peak regions were observed because of the mechanical treatment.

#### 4.4. Sum-frequency generation vibrational spectroscopy (SFG)

Another vibrational method, SFG vibrational spectroscopy, is a non-linear optical phenomenon demonstrated to be an effective technique for probing interfaces.<sup>166–169</sup> It is highly sensitive to non-centrosymmetric crystalline materials, e.g., piezocrystals and non-centrosymmetrically arranged functional groups in a medium.<sup>170,171</sup> Due to its selective vibrational modes and

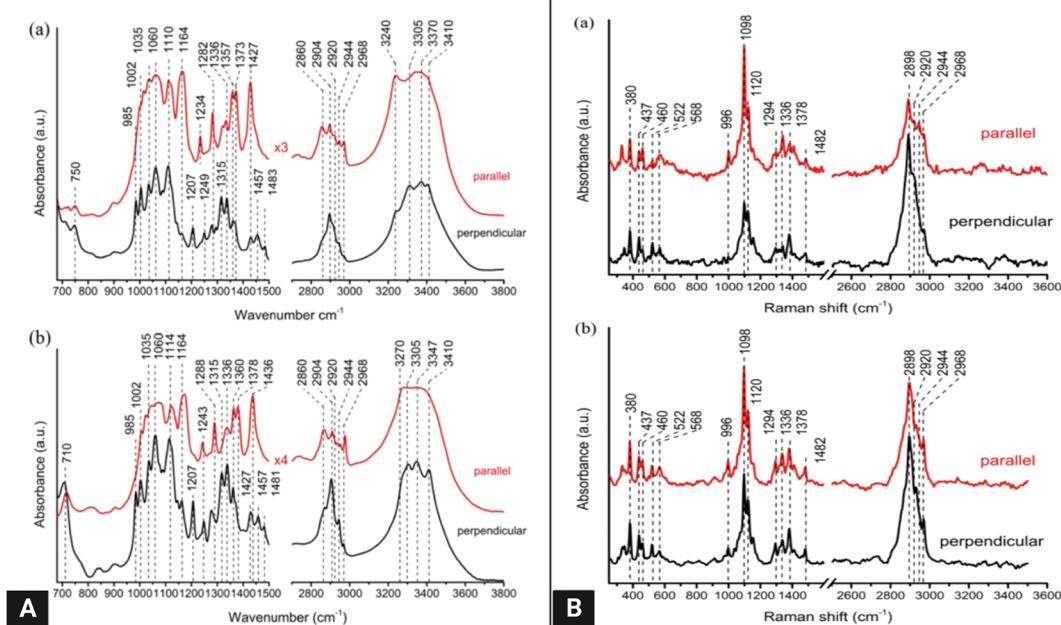


Fig. 11 (A) Polarized FT-IR spectra of laterally packed and aligned films of (a) cellulose I<sub>α</sub> and (b) I<sub>β</sub> nanocrystals. Spectra were taken with the IR polarization perpendicular to the nanocrystals' alignment direction.<sup>117</sup> (B) Polarized FT-Raman spectra of laterally packed and aligned films of (a) cellulose I<sub>α</sub> and (b) I<sub>β</sub> nanocrystals. Spectra were taken with the excitation laser polarization perpendicular to the nanocrystals' alignment direction. The peaks discussed in the text are highlighted with dotted lines<sup>117</sup> (reprinted with permission from ref. 117. Copyright 2013 Korean Journal of Chemical Engineering).



sensitivity for chiral molecules in biomaterials, it is also used to probe biological molecules such as DNA and other helical proteins.<sup>172</sup> Recently, it has been shown to selectively detect crystalline cellulose from lignocellulosic biomass.<sup>61</sup>

The principle of SFG vibrational spectroscopy depends on the non-linear optical response of sum frequency generation. In this process, temporal and spatial overlap between two pulsed laser beams occurs at an interface. One of these beams has a fixed visible frequency ( $\omega_{\text{VIS}}$ ) and the other has an adjustable infrared frequency ( $\omega_{\text{IR}}$ ). Thus, the light is emitted as the summed frequency ( $\omega_{\text{SFG}}$ ) of the two frequencies *via* the equation<sup>173</sup>

$$\omega_{\text{SFG}} = \omega_{\text{VIS}} + \omega_{\text{IR}} \quad (14)$$

Resonance occurs when the vibrational frequency of the interface molecules coincides with the frequency of the tunable infrared radiation, and thus the intensity of the light is enhanced. This sum-frequency light is plotted against infrared frequency to obtain a vibrational spectrum.<sup>173</sup> A schematic of the SFG process is shown in Fig. 12A. When light propagates through a medium, the valence electrons of the molecules of the medium experience a force due to its electric field,  $E(\omega)$ . A dipole moment is created as the sum of all molecular electric dipoles, which is called polarization of the medium,  $P(\omega)$ , according to the equation<sup>28</sup>

$$P(\omega) = \epsilon_0(X^{(1)}E(\omega) + X^{(2)}E(\omega)^2 + X^{(3)}E(\omega)^3 + \dots) \quad (15)$$

where  $\epsilon_0$  is the dielectric permittivity of a vacuum and  $\chi^{(n)}$  is the electric susceptibility of the  $n$ th order. Mathematically,  $\chi$  is a

tensor with a  $3^{n+1}$  component.<sup>174</sup> The higher order  $\chi$  terms are extremely small compared to  $\chi^{(1)}$  when a material is irradiated with typical IR or Raman irradiation sources. However, these terms become significant when the material is irradiated with high-intensity laser beams.<sup>28</sup> As SFG depends on second-order susceptibility, we will restrict the discussion to the second-order polarization term. The electric field,  $E(\omega)$ , can be expressed as the sum of two cosine waves using the equation

$$E(\omega) = E_1 \cos \omega_1 t + E_2 \cos \omega_2 t \quad (16)$$

where  $E(\omega)$  is the electric field,  $E_1$  and  $E_2$  are two components and  $t$  is time. Thus, the second-order polarization term can be expressed as

$$\begin{aligned} P^{(2)} &= \epsilon_0 \chi^{(2)} (E_1 \cos \omega_1 t + E_2 \cos \omega_2 t)^2 \\ &= \epsilon_0 \chi^{(2)} (E_1^2 \cos^2 \omega_1 t + 2E_1 E_2 \cos(\omega_1 t) \\ &\quad \times \cos(\omega_2 t) + E_2^2 \cos^2 \omega_2 t) \\ &= \epsilon_0 \chi^{(2)} \left( \frac{E_1^2 (1 + \cos 2\omega_1 t)}{2} + 2E_1 E_2 \right. \\ &\quad \times \left. \frac{\cos(\omega_1 + \omega_2)t + \cos(\omega_1 - \omega_2)t}{2} + \frac{E_2^2 (1 + \cos 2\omega_2 t)}{2} \right) \\ &= \epsilon_0 \chi^{(2)} \left( \frac{1}{2} (E_1^2 + E_2^2) + \frac{1}{2} (E_1^2 \cos 2\omega_1 t + E_2^2 \cos 2\omega_2 t) \right. \\ &\quad \left. + E_1 E_2 \cos(\omega_1 + \omega_2)t + E_1 E_2 \cos(\omega_1 - \omega_2)t \right) \end{aligned} \quad (17)$$

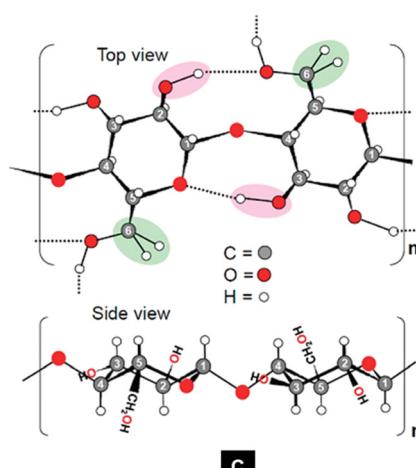
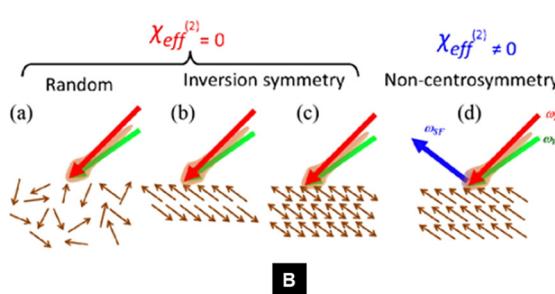
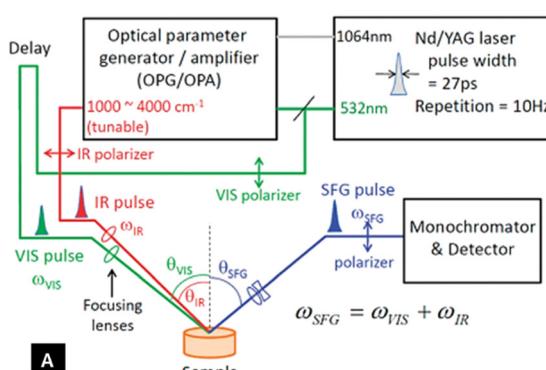


Fig. 12 (A) Schematic diagram of an SFG system<sup>61</sup> (reprinted with permission from ref. 61. Copyright 2011 American Chemical Society). (B) Schematic of the noncentrosymmetry requirement of  $X_{\text{eff}}^{(2)}$  in the SFG process.<sup>117</sup> The directionality of vibration modes is depicted by brown arrows (reprinted with permission from ref. 117. Copyright 2013 Korean Journal of Chemical Engineering). (C) Schematic view of two glucopyranose units in a cellulose chain. The SFG-active  $\text{CH}_2$  and  $\text{OH}$  stretch peaks are highlighted in green and pink, respectively (reprinted with permission from ref. 61. Copyright 2011 American Chemical Society).



Both the incident laser beams give rise to a DC field (independent frequency term,  $E_1^2 + E_2^2$ ), simple harmonic generation (SHG) for both the input beams (twice the frequency of both input frequencies,  $2\omega_1$  and  $2\omega_2$ ), difference frequency generation (DFG,  $\omega_1 - \omega_2$ ) and SFG ( $\omega_1 + \omega_2$ ).<sup>173</sup> SFG vibrational spectroscopy focuses on the wave emitted from the sample having the sum frequency of the incident beams, in this case, the frequency of IR and visible laser beams. The intensity of the SFG signal,  $I(\omega_{\text{SFG}})$ , can be expressed as

$$I(\omega_{\text{SFG}}) \propto |X_{\text{eff}}^{(2)}|^2 I(\omega_{\text{IR}}) I(\omega_{\text{VIS}}) \quad (18)$$

where  $I(\omega_{\text{IR}})$  and  $I(\omega_{\text{VIS}})$  are the intensities of IR and visible beams, respectively, and  $\chi_{\text{eff}}$  is the effective second-order nonlinear susceptibility that can be expressed as

$$\chi_{\text{eff}}^{(2)} = \frac{N \sum_{\alpha, \beta, \gamma} \langle M_{\alpha\beta} A_{\gamma} \rangle}{\varepsilon_0 (\omega_{\text{IR}} - \omega_q - i\Gamma)} \quad (19)$$

where  $N$  is the number density of the SFG signal-producing molecules,  $\langle M_{\alpha\beta} A_{\gamma} \rangle$  is the angle average of the product of IR and Raman tensors,  $\omega_q$  is the frequency of normal vibration modes, and  $\Gamma$  is the damping constant. As the expression includes the product of both IR and Raman tensors, the vibration modes must be Raman and IR active to show the SFG signal. Another important selection rule for the SFG process is that the irradiated medium must not have inversion symmetry.<sup>173,175,176</sup>

As shown in Fig. 12B, when the vibration modes (shown by brown arrows) are arranged either randomly (a), or with inversion symmetry (b), or the molecule has inversion symmetry itself (c), the SFG signal cannot be generated. The only way that  $X_{\text{eff}}^{(2)}$  becomes non-zero is when the molecules with no inversion symmetry are arranged non-centrosymmetrically in space.<sup>28,173</sup> Non-centrosymmetry exists over several orders of length scales in crystalline cellulose.<sup>177</sup> At the molecular level, its carbon centers (C1, C2, C3, C4, and C5) in the glucose ring are chiral. The crystal unit cells manifest non-centrosymmetry (space groups  $P1$  and  $P2_1$ ).<sup>56,68</sup> Also, cellulose  $I_{\alpha}$  and  $I_{\beta}$  have parallel glucan chains, unlike cellulose II, in which the chains have an antiparallel arrangement.<sup>178</sup> As these glucan chains in cellulose I are unidirectional, the  $\text{CH}_2$  group of exocyclic  $\text{C}_6\text{H}_2\text{OH}$  side chains and OH groups that have interchain H-bonding form net dipoles across the entire crystal or multiple closely packed crystals.<sup>61,179</sup> As a result, the characteristic SFG peaks located at  $2945 \text{ cm}^{-1}$  and  $3320\text{--}3250 \text{ cm}^{-1}$  are observed for  $\text{CH}_2$  and OH vibration modes, respectively.<sup>62,178</sup> It has been observed that, the SFG peak intensity at  $2945 \text{ cm}^{-1}$  varies monotonically, but non-linearly with crystalline cellulose samples,<sup>97</sup> based on which an Avicel-based calibration curve using XRD, IR, and Raman techniques to determine the amount of Avicel-equivalent crystalline cellulose in wood samples has been developed.

Other C-H and O-H than the aforementioned peaks are not observed because of the non-centrosymmetric selection rule. The same reason is also attributed to the SFG inactivity of hemicellulose and lignin because of their random arrangement in biomass.<sup>61</sup> Moreover, these noncentrosymmetric crystals could be dispersed in a random, centrosymmetric, or

noncentrosymmetric fashion within an amorphous matrix, which can also affect the SFG spectra over the characteristic length of SFG.<sup>180</sup> The SFG technique is becoming increasingly popular in probing the crystal structure of cellulose. Huang *et al.* (2020) employed SFG vibrational spectroscopy to correlate cellulose enzyme complexes' shape and size to cellulose's crystal size using freshwater algae *Microasterias*.<sup>181</sup> The process was also able to distinguish between the surface and bulk hydroxyl groups in cellulose nanocrystals owing to its non-linear scattering.<sup>182</sup> The SFG technique is also sensitive to changes in the inter-fibrillar distance between cellulose microfibrils and the distance between crystals.<sup>183,184</sup> The process can distinguish between chain polarity and determine the orientation in the direction of growth of the cellulose microfibrils (unidirectional or bidirectional).<sup>28,185<sup>\text{--}</sup>187</sup> This technique is also able to distinguish different cellulose polymorphs.<sup>178</sup> Moreover, the technique has enabled researchers to probe into the mesoscale assembly of cellulose microfibrils.<sup>179,180</sup>

The SFG technique can be used to quantify cellulose crystallinity, given that the crystalline packing pattern of cellulose remains reasonably constant.<sup>97</sup> Current quantification of cellulose crystallinity by the SFG method involves a calibration curve between a known cellulose concentration and the SFG peak intensity.<sup>188,189</sup> However, the quantification becomes challenging when the packing of crystalline cellulose changes drastically with different samples.<sup>177,180</sup>

## 5. Differential scanning calorimetry (DSC)

Calorimetry measures thermal energy absorbed or released when a change in state variables is encountered due to phase transitions, chemical reactions, or physical changes. The simple equation for measuring heat change as a function of temperature can be expressed as

$$dQ = m \times c \times \Delta T \quad (20)$$

where  $dQ$  is the heat change,  $m$  is the sample mass,  $c$  is the specific heat capacity, and  $\Delta T$  is the temperature change. Two methods, namely heat-flux and power-compensated DSC, are typically used.

### 5.1. Heat-flux DSC

In a heat-flux DSC system, the reference and the sample are heated in a single furnace, where two separate sensors are used to measure their temperatures. A purge gas allows a uniform heat distribution and expels unwanted substances like oxygen. Thus, both the sample and reference are given equal heat transfer, and their temperature difference is measured. The temperature difference is translated using the heat flow equation to heat flux.<sup>190</sup> However, the heat-flux DSC systems suffer from the disadvantage of the significant time constant. Hence, they are typically not used to study semi-crystalline polymers.<sup>191</sup> A simplified schematic is shown in Fig. 13A.



## 5.2. Power-compensated DSC

In power-compensated DSC systems, two separate chambers are used for the sample and the reference. Two individual heaters heat the sample and the reference through a differential power controller to maintain a zero-temperature difference between them. The temperature is increased linearly with time. Thus, in power-compensated DSC systems, the energy flow rate varies to keep the same temperature in both chambers. The purge gas is used to allow for uniform heat distribution and the removal of unwanted substances. The difference in power can be determined using the equation

$$\Delta P = \frac{dQ_{\text{Sample}}}{dt} - \frac{dQ_{\text{Reference}}}{dt} = I \times \Delta V = \frac{dH}{dt} \quad (21)$$

where  $\Delta P$  is the required change in power,  $dQ/dt$  is the heat absorbed or released by the sample or the reference per unit time,  $I$  is the total current,  $\Delta V$  is the voltage difference, and  $dH/dt$  is the change of enthalpy per unit time. A simplified schematic of power-compensated DSC is shown in Fig. 13B.

Both methods produce an output of the heat-flow curve with respect to temperature. From the curve, an inference about material properties can be drawn. Fig. 13C shows the general DSC output of a semi-crystalline polymer along with possible interpretations. The enthalpy change is proportional to the area under the peaks. The enthalpy change ( $\Delta H$ ) can be expressed as

$$\Delta H = K \times A \quad (22)$$

where  $K$  is the calorimetric constant and  $A$  is the area under the curve. Thus, the DSC technique can be instrumental in gaining information about the thermodynamic functions by measuring

the specific heat of substances.<sup>191,192</sup> The technique and principle for measuring specific heat have been previously published.<sup>193</sup> This technique can also be used to assess the thermal behavior of polymers and structural information, including crystallinity.<sup>191,194</sup>

## 5.3. Measurement of crystallinity of cellulose using DSC

Conventionally, relative crystallinity in polymers is determined by taking the ratio of the melting enthalpy of the sample obtained from a DSC run to the melting enthalpy of a crystalline sample.<sup>194,195</sup> It can be expressed as

$$X_c = \frac{\Delta H_m}{\Delta H_m^\circ} \times 100\% \quad (23)$$

where  $X_c$  is the relative amount of the crystalline fraction in the polymer,  $\Delta H_m$  is the melting enthalpy of the polymer sample, and  $\Delta H_m^\circ$  is the melting enthalpy of a 100% crystalline sample of the same polymer.

Cellulose has a mixture of well-defined crystalline and amorphous structures and the amorphous regions allow aqueous reagents to penetrate inside the structure.<sup>196</sup> The crystalline domains are very closely packed and resist water absorption. Hence, it is thought that the free hydroxyl groups present in the amorphous regions are responsible for almost complete water absorption. Based on these considerations, Dale proposed a method to determine the cellulose crystallinity that considers the dehydration endothermic peak area.<sup>196</sup> It is reported that cellulose fibres typically show an endothermic peak between 273 and 400 K, which is attributed to dehydration.<sup>197</sup> Thus, the authors hypothesized that the endothermic peak of cellulose dehydration is related to the amorphous fraction of

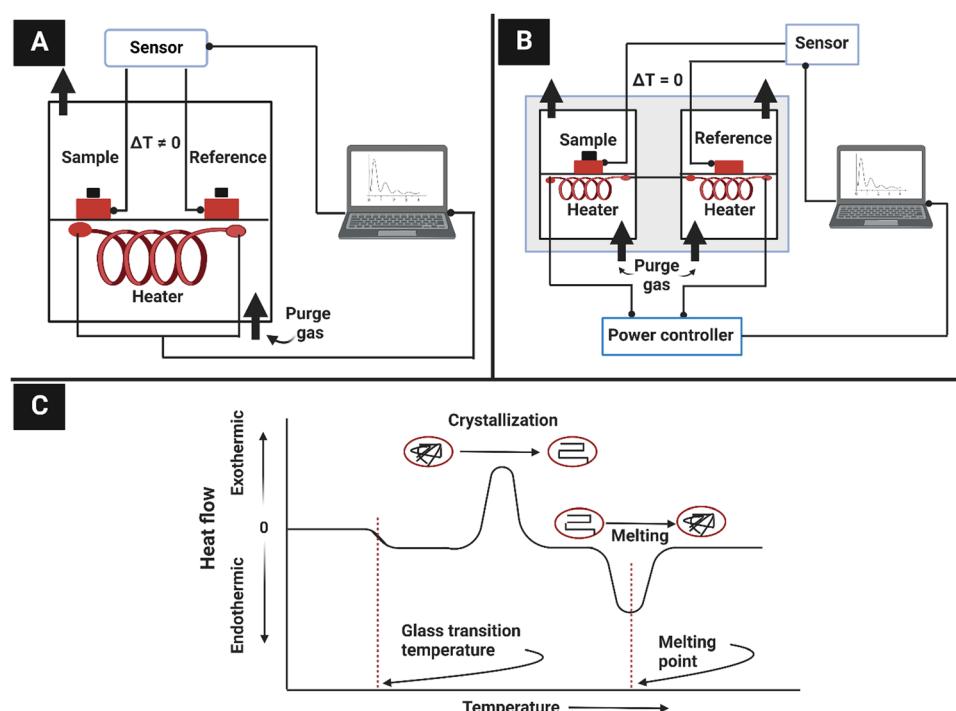


Fig. 13 (A) Schematic of heat-flux DSC, (B) schematic diagram of power-compensated DSC, and (C) general DSC output for semi-crystalline polymers.



cellulose.<sup>196</sup> The cellulose crystallinity (CC%) could be determined from the equation

$$CC\% = \frac{\Delta H_0 - \Delta H_s}{\Delta H_0} \times 100 \quad (24)$$

where  $\Delta H_0$  is the heat required to remove water from a standard cellulose, which is entirely amorphous and exposed to a specific relative humidity until a constant weight is obtained, and  $\Delta H_s$  is the heat needed to dehydrate the cellulose sample which has also been exposed to the same conditions as the standard. However, it is to be noted that the use of wet amorphous cellulose could be problematic as it becomes recrystallized in the presence of moisture. Moreover, to develop a general protocol to assess cellulose crystallinity with this method, it is necessary to know the relative humidity at which the samples sorbed water vapor and the preservation technique before dehydration.

There are some other limitations associated with DSC since this method assumes that the amorphous regions of cellulose are entirely responsible for water absorption. However, a small portion of water adheres to the crystalline part of cellulose, which cannot be avoided for samples with high crystallinity. Thus, the method is most suitable for low crystallinity cellulose samples.<sup>196</sup> It is also difficult to estimate the baseline for the peak area measurement.<sup>191,194,198</sup> Although the DSC technique has gained attention among polymer scientists regarding its use in determining crystallinity, it has been used in only a limited way in cellulose research. The main reason for this is that in conventional semicrystalline polymers one can measure the heat of crystallization through melting. Since cellulose does not melt, it is not possible to measure this property. Moreover, the transitions of cellulose due to moisture are also influenced by other factors and other polymers present in the samples. It has been used to probe only the thermal behavior, which can be understood from the literature survey on determining cellulose crystallinity by the DSC method.<sup>199,200</sup> Further research is necessary to address the limitations associated with the DSC technique of assessing cellulose crystallinity.

## 6. Cellulose crystallinity determination by carbohydrate-binding modules (CBM)

Biochemical routes have also been developed to determine the cellulose crystallinity in addition to getting information on the morphology of cellulosic structures. Initially, these methods were developed to quantitatively determine the accessibility of cellulose fibres subjected to enzymatic hydrolysis. However, the degree of cellulose crystallinity may not influence the yield during enzymatic hydrolysis but can affect the kinetics of hydrolysis.<sup>201</sup> For example, it has been shown that paracrystalline cellulosic regions hydrolyze faster than crystalline regions.<sup>202</sup> However, as cellulose crystallinity is one of the critical factors influencing cellulose accessibility, qualitative information on the crystalline cellulosic part can be obtained through these methods.

These methods take advantage of carbohydrate-binding modules (CBMs), which are polypeptide sequences found in

carbohydrate-active enzymes with specific binding abilities to different carbohydrate structures. Depending on the molecular structures and amino acid sequences, CBMs can be classified into 68 families and based on their three-dimensional structures and activity, they are classified as type A, type B, and type C.<sup>201,203</sup> Type A CBMs are surface binding type proteins with a greater affinity towards crystalline cellulose. Unlike type A, type B proteins show affinity towards the amorphous cellulosic regions by binding to glycan chains. Yet, type B CBMs cannot recognize glycan chains consisting of less than three sugar units while type C CBMs exhibit binding specificity for 1–3 sugar units.<sup>201,203</sup>

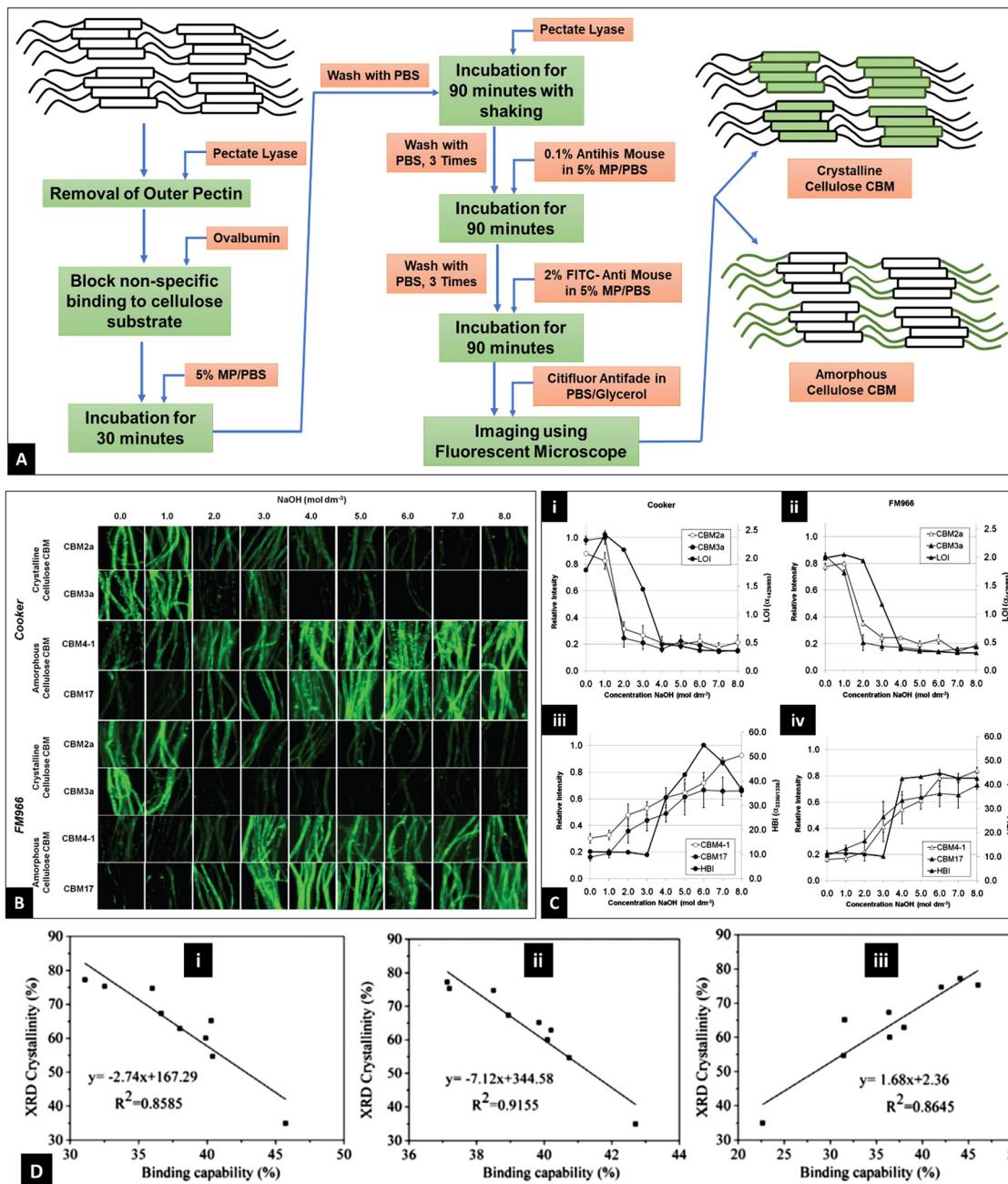
Moreover, fusion proteins have also been developed by combining CBMs with green fluorescent proteins (GFPs) to visualize cellulose morphology easily. For example, Hong *et al.* (2007) prepared one such fusion protein to quantitatively determine cellulose accessibility to cellulase with a hydrolysis temperature of 50 °C.<sup>204</sup> Similarly, Li *et al.* (2018) designed a recombinant protein by combining CBMs with a di-green fluorescent protein to quantitatively determine cellulose accessibility to cellulase.<sup>205</sup> Kljun *et al.* (2011) used CBMs to assess changes in cellulose crystallinity during the mercerization process of cotton fibres (Fig. 14A).<sup>206</sup> The authors used two type A and two type B CBMs coupled with fluorescence imaging to visualize the changes in the crystallinity at various stages of the mercerization process, which agreed with the results obtained from traditional XRD and FTIR techniques as represented in Fig. 14B.<sup>206</sup> The binding intensity of the CBMs was calculated by measuring the total dark areas for the amorphous region and total green/bright areas for the crystalline region. The total fibre area was also measured, and the bright or dark area ratio to the total fibre area was taken to calculate the percentage of bound and unbound CBMs to the fibre. They also estimated the lateral order index (LOI) and hydrogen bond intensity (HBI) from an FTIR trace. LOI and HBI were determined from the intensity of the absorbance bands located at 1429 to 893 cm<sup>-1</sup> and the ratio of the intensities of the 3336 to 1336 cm<sup>-1</sup> bands, respectively.<sup>207,208</sup> They found that the LOI followed a similar trend as the relative intensity of the binding capacity of the CBMs onto the crystalline cellulose for both the Coker and FM966 fibres. On the other hand, HBI increased with an increase in the amorphous region of the fibre samples. It also displayed similar trends to the relative intensity of the CBM which shows affinity towards the amorphous cellulosic region, as shown in Fig. 14C.

Guo *et al.* (2019) constructed a CBM-GFP-directed protein probe to quantitatively determine cellulose crystallinity. The binding concentrations of the different probes were determined using a fluorescence spectrophotometer, and the binding capability of the probes to substrates was calculated from the relative intensity using the equation<sup>209,210</sup>

$$\text{Binding capability (\%)} = \frac{(\text{FI}_{\text{control}} - \text{FI}_{\text{final}})}{\text{FI}_{\text{control}}} \times 100 \quad (25)$$

where  $\text{FI}_{\text{control}}$  and  $\text{FI}_{\text{final}}$  represent the supernatant's fluorescence intensity after incubation without and with cellulosic substrates, respectively. The authors showed a positive linear





**Fig. 14** (A) Schematic of cellulose crystallinity determination using CBMs.<sup>206,208</sup> (B) Fluorescence imaging of the binding of four CBMs to different cotton fibres after treatments with a range of NaOH concentrations (0–8 M). (C) Quantification of CBM binding to cotton fibres and FTIR analyses in response to NaOH treatments. (i) and (ii) Comparison of LOI with the binding intensity of crystalline directed CBMs, (iii) and (iv) comparison of HBI with the binding intensity of amorphous directed CBMs. (Parts B and C are reprinted with permission from ref. 206. Copyright 2011 American Chemical Society.) (D) Correlation between CBM binding capacity and crystallinity index of different cellulose samples where (i) and (ii) CBMs show affinity towards amorphous cellulose and (iii) CBM has an affinity for crystalline cellulose (reprinted with permission from ref. 209, 2019. All rights reserved by Springer International Publishing).

correlation of the binding ability of their fusion proteins, CG2a (type A CBM, affinity towards crystalline cellulose) and CBM17 and CBM28 (type B CBMs, affinity towards amorphous cellulose), with the XRD crystallinity data as shown in Fig. 14D. Based on their correlation equation, the authors calculated the cellulose crystallinity of representative samples and reported the same trends as observed from XRD measurements.<sup>209</sup>

## 7. Other methods

Several physicochemical properties of cellulose such as specific gravity,<sup>211,212</sup> enthalpy of wetting,<sup>213</sup> water vapor sorption,<sup>214</sup> iodine sorption,<sup>215</sup> chemical reactivity,<sup>216</sup> and others could also be used to estimate the crystalline fraction in the cellulose polymer. Some of these are discussed in the following.

## Specific gravity

Specific gravity is defined as the density of a substance compared to the density of water. Direct measurement of specific gravity of amorphous cellulose can be challenging due to the dependence of cellulose structure on factors such as the presence of impurities and residual moisture, different oxidation degrees, and drying methods. Thus, cellulose samples with known degrees of crystallinity can be used to estimate the crystallinity of the unknown samples. It is known that the relationship between the specific gravity of semi-crystalline polymers and their degree of crystallinity ( $X$ ) can be expressed as follows:

$$X = \frac{(\rho_c/\rho)(\rho - \rho_a)}{(\rho_c - \rho_a)} \quad (26)$$

where  $\rho_c$  and  $\rho_a$  are the specific gravities of the crystalline and amorphous regions of the semicrystalline polymer, respectively. Eqn (26) can be converted and rearranged to obtain

$$\frac{1}{\rho} = \frac{1}{\rho_a} - X \left[ \left( \frac{1}{\rho_a} \right) - \left( \frac{1}{\rho_c} \right) \right] \quad (27)$$

which resembles a straight-line equation of the form

$$V = V_a - KX \quad (28)$$

where  $\frac{1}{\rho} = V$  is the sample's specific volume and  $\frac{1}{\rho_a} = V_a$  is the specific volume of the amorphous region, and  $K = \left( \frac{1}{\rho_a} \right) - \left( \frac{1}{\rho_c} \right)$ .

Thus, by measuring the specific gravity of some cellulose samples of known crystallinity, it is possible to obtain a straight line and calculate the specific gravity of the crystalline and amorphous regions. Loelovich *et al.* (2010) used  $\rho_c = 1.62 \text{ g cm}^{-3}$  and  $\rho_a = 1.44 \text{ g cm}^{-3}$  to measure the cellulose crystallinity with this method. The results obtained agreed with the XRD results.<sup>211</sup>

## Wetting enthalpy

It has been mentioned before that, the cellulose–water interaction is attributed to the amorphous regions of cellulose. Thus, the associated exothermic heat generation with this phenomenon or wetting enthalpy ( $\Delta H$ ) is directly proportional to the amorphous content ( $Y$ ) of cellulose.

$$\Delta H = \Delta H_{\text{Am}} Y \quad (29)$$

$$Y = \frac{\Delta H}{\Delta H_{\text{Am}}} \quad (30)$$

$$X = 1 - \left( \frac{\Delta H}{\Delta H_{\text{Am}}} \right) \quad (31)$$

where  $X$  is the content of the crystalline domain in the cellulose sample and  $\Delta H_{\text{Am}}$  is the wetting enthalpy of completely amorphous cellulose. Experimental observations have shown that the wetting enthalpy ( $\Delta H$ ) increases linearly with the maximum vapor sorption value ( $A_o$ ) of the cellulose sample and a constant specific enthalpy has been reported  $\left( h = \frac{\Delta H}{A_o} \simeq 335 \text{ J g}^{-1} \text{ water} \right)$ .

Thus, considering maximal absorption for amorphous cellulose,  $\Delta H_{\text{Am}}$  was determined as follows:

$$\Delta H_{\text{Am}} = h A_o \quad (32)$$

Substituting  $\Delta H_{\text{Am}}$  value into eqn (31) yields the cellulose crystalline fraction. The authors reported that, the crystallinity values obtained by this method were higher than those obtained by the XRD technique.<sup>214</sup>

## 8. Comparison among different techniques for measuring cellulose crystallinity

Cellulose crystallinity has long been playing an important role in defining a material's physical and chemical properties. Thus, the study of the crystalline structure of cellulose has received attention for its diversified application. Different techniques have also been used to measure the cellulose crystallinity index, all of which yield different values. These are now discussed in this section and compared, as shown in Fig. 15.

Cellulose pulps from different sources having different purity were taken and analyzed using XRD, NMR, IR, Raman, SFG, and DSC to measure their crystallinity index.<sup>177</sup> Cotton linter (one of the purest sources of cellulose), which contain more than 95% of cellulose by dry mass, Avicel with a low degree of polymerization,  $\alpha$ -cellulose (mostly cellulose I<sub>B</sub>), bleached hardwood (HW), and softwood (SW) pulps were taken for comparison. Fig. 15 shows the graphical representation of data produced while measuring the CI of cellulose using a range of different techniques, namely XRD, NMR, IR, Raman, and SFG. The CI values measured from XRD were calculated using the peak height (PH) and amorphous subtraction (AS) method, as shown in Table 6. For <sup>13</sup>C NMR, the CI was measured by the relative intensity of the C4 peak chemical shifts of amorphous (80–85 ppm) and crystalline (85–89 ppm) cellulose. It can be observed from Table 6 that the methods used to measure the cellulose CI values show the same qualitative trends for samples, *i.e.*, CI is the highest for cotton linter and lowest for  $\alpha$ -cellulose. However, the absolute CI values differ for different methods, limiting direct quantitative comparison of the CI values estimated using different approaches. So, the cellulose crystallinity index measured for different samples using any technique can only be compared qualitatively.

The CI values of HW and SW measured using the XRD AS method are very close to those measured using the PH method. However, there is a large difference in the CI values for other cellulose sources, as shown in Table 6. This was due to the high intensity of the amorphous background of the HW and SW fibres at a  $2\theta$  angle greater than  $25^\circ$ , whereas the other cellulosic sources exhibited low intensity in this region. As a result of this, the AS method was not able to subtract the amorphous background in the high  $2\theta$  region, and instead it effectively added the difference in the crystalline portion, which led to much higher CI values than the true mass fraction of



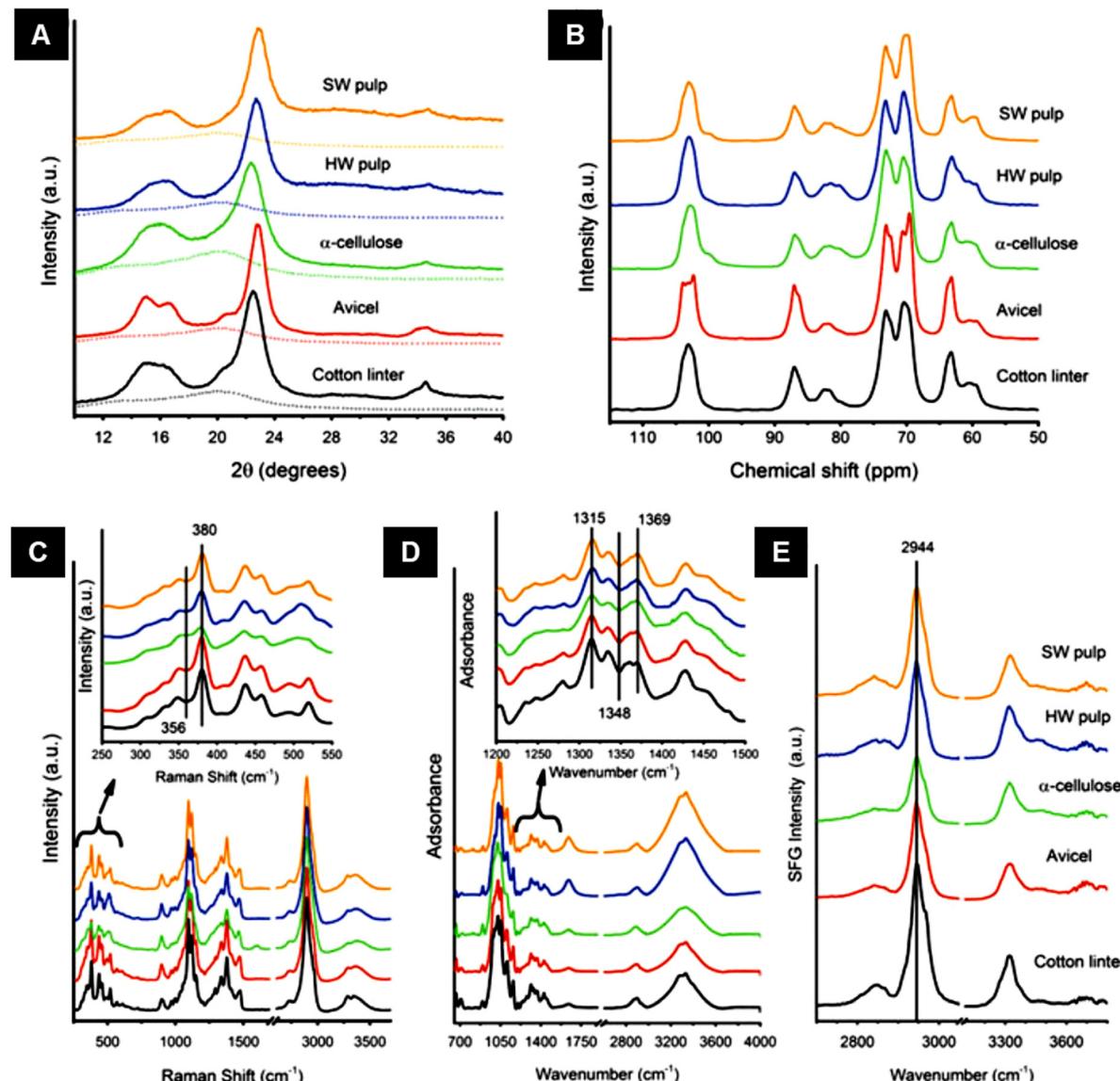


Fig. 15 Crystallinity measurement of cellulose samples using (A) XRD, (B)  $^{13}\text{C}$  NMR, (C) Raman, (D) IR, and (E) SFG. Samples are cotton linter (black), Avicel (red),  $\alpha$ -cellulose (green), hardwood (HW) pulp (blue), and softwood (SW) pulp (orange). In (A), the amorphous standard is plotted as a dotted line. The insets in figures (C) and (D) show the spectral regions used for CI calculations. All spectra are offset for clarity (Reprinted with permission from ref. 217, 2015. All rights reserved by Springer International Publishing).

crystalline cellulose over the total mass for HW and SW than the other samples.

Different CI values are obtained from different techniques for several reasons. NMR is likely to yield a much lower CI value compared to that calculated from XRD, even if the CI is measured by the AS method, as shown in Table 6. It has been found from different literature sources that the XRD method gives more precise data for the crystalline phase and shows a lesser sensitivity to the amorphous phase, consequently producing data having an inevitable inclination towards crystalline phases.<sup>15,217</sup> The cellulose structure also contains paracrystalline cellulose, a more complicated structure than the crystalline and amorphous cellulose, which also interferes with the CI value measurement.<sup>125</sup> In contrast, NMR shows equal

sensitivity to cellulose structure's crystalline and amorphous phases. The difference in detection sensitivity between these two methods affects the CI values and is responsible for the discrepancy, as shown in Table 6.

The CI values measured using DSC were very close to those measured using XRD. Still, DSC has some limitations as it only takes account of the moisture present in the amorphous region of the cellulose and completely neglects the presence of moisture at the surface of the crystalline regions. Another limiting factor of DSC is that it measures only the sum of all heat flow events occurring at a particular time and temperature, which might result in hiding the exothermic crystallization of others. Thus, the method is most suitable for cellulose having low crystallinity.<sup>196</sup> Therefore, it is clear from the discussion that

**Table 6** Crystallinity index of various cellulose-containing samples measured using different techniques

Sample	Calculated crystallinity <sup>a</sup> (%)						Ref.
	XRD(PH)	XRD(AS)	NMR	Raman	IR	SGF	
Cotton linter	85	72	65	78	80	77	83.8 196,217
Avicel	81	70	58	77	75	71	80.37 196,217
$\alpha$ -Cellulose	65	48	44	65	67	65	63.73 196,217
Hardwood fibre	84	81	45	73	72	71	79.75 159,217
Softwood fibre	84	84	52	76	73	73	82.23 159,217

<sup>a</sup> Values obtained from Raman, IR, and SFG were correlated relative to the XRD values obtained using the peak height method (PH) and the amorphous subtraction method (AS).

none of the crystallinity measuring techniques can be used to obtain an absolute value. However, all the methods display similar trends in measuring the crystallinity values of cellulose samples. Therefore, the techniques overviewed can be used to provide qualitative data and a relative measurement of cellulose crystallinity.

## 9. Conclusions

The importance of defining the crystallinity of cellulose has been discussed in many reports and over many decades since the very early beginning of our understanding of this complex and ubiquitous polymer. Understanding the structure of cellulose facilitates physical modification to render high-performance products and increase product yields. This review paper has covered the different techniques used to determine cellulose crystallinity, namely X-ray diffraction (XRD), nuclear magnetic resonance (NMR), Fourier-transform infrared spectroscopy (FTIR), Raman, differential scanning calorimetry (DSC), sum-frequency generation vibrational spectroscopy (SFG), cellulose binding molecule (CBM) methods, specific gravity, and wetting enthalpy. Each crystallinity index can take different values depending on the technique selected for the analysis. Typically, CI values obtained from NMR are lower than values obtained from XRD measurements. However, relative crystallinity values and crystallinity trends are common for all techniques. For example, cotton linter, known as almost a fully crystalline structure, yields the highest CI according to the different methods presented. Moreover, other factors such as size of the crystallite, intensity of the incident beam (in the case of XRD), source and pretreatment of the sample before the tests also affect the crystallinity value. In this sense, despite many efforts to determine cellulose crystallinity, there is no agreement in the literature on the absolute technique for determining the crystallinity value and studying the cellulose structure. This fact reflects that cellulose crystallinity cannot be estimated using a single technique; therefore, combining different techniques to calculate the CI is likely the best approach. The use of software or computational analysis might be a vital tool in addition to the experimental data since the software can be designed to calculate and predict the intricate details and refine the data which is not possible experimentally. Consequently, they can complement the

experimental findings by assisting the interpretation of the experimental data.

## Conflicts of interest

The authors have no conflicts of interest to declare that are relevant to the content of this article.

## Acknowledgements

This work is in part supported by McIntire-Stennis grant number NI21MSCFRXXXG054 from the USDA National Institute of Food and Agriculture. The work of Dr. Lokendra Pal is in part supported by E. J. "Woody" Rice Professorship Endowment. The authors also acknowledge the funding support from Ministry of Science and Technology (Project ID: SRG-222388), Bangladesh and University Grants Commission's research grant (2022-2023) for the faculty of University of Dhaka while carrying out the work of Dr. Salem partially in the Department of Applied Chemistry and Chemical Engineering, University of Dhaka. We also wish to acknowledge Dr. Peter W. Hart of WestRock for his generous financial support for the current review. The authors gratefully thank Mrittika Debnath for her support during the writing of this manuscript.

## References

- 1 M. Brodin, M. Vallejos, M. T. Opedal, M. C. Area and G. Chinga-Carrasco, *J. Cleaner Prod.*, 2017, **162**, 646–664.
- 2 K. S. Salem, H. R. Starkey, L. Pal, L. Lucia and H. Jameel, *ACS Sustainable Chem. Eng.*, 2020, **8**, 1471–1478.
- 3 K. S. Salem, V. Naithani, H. Jameel, L. Lucia and L. Pal, *Global chall.*, 2021, **5**, 2000065.
- 4 K. S. Salem, V. Naithani, H. Jameel, L. Lucia and L. Pal, *Carbohydr. Polym.*, 2022, **295**, 119856.
- 5 C. Chen, Y. Kuang, S. Zhu, I. Burgert, T. Keplinger, A. Gong, T. Li, L. Berglund, S. J. Eichhorn and L. Hu, *Nat. Rev. Mater.*, 2020, **5**, 642–666.
- 6 S. J. Eichhorn, A. Dufresne, M. Aranguren, N. E. Marcovich, J. R. Capadona, S. J. Rowan, C. Weder, W. Thielemans, M. Roman, S. Renneckar, W. Gindl, S. Veigel, J. Keckes, H. Yano, K. Abe, M. Nogi, A. N. Nakagaito, A. Mangalam, J. Simonsen, A. S. Benight, A. Bismarck, L. A. Berglund and T. Peijs, *J. Mater. Sci.*, 2010, **45**, 1–33.
- 7 A. Etale, A. J. Onyianta, S. R. Turner and S. J. Eichhorn, *Chem. Rev.*, 2023, **123**, 2016–2048.
- 8 X. Ju, M. Bowden, E. E. Brown and X. Zhang, *Carbohydr. Polym.*, 2015, **123**, 476–481.
- 9 S. Rongpipi, D. Ye, E. D. Gomez and E. W. Gomez, *Front. Plant Sci.*, 2019, **9**, 1894.
- 10 N. S. Çetin, P. Tingaut, N. Özmen, N. Henry, D. Harper, M. Dadmun and G. Sèbe, *Macromol. Biosci.*, 2009, **9**, 997–1003.
- 11 W. Stelte and A. R. Sanadi, *Ind. Eng. Chem. Res.*, 2009, **48**, 11211–11219.



12 A. D. French, *Cellulose*, 2017, **24**, 4605–4609.

13 E. Toumpanaki, D. U. Shah and S. J. Eichhorn, *Adv. Mater.*, 2021, **33**, 2001613.

14 A. D. French, M. Concha, M. K. Dowd and E. D. Stevens, *Cellulose*, 2014, **21**, 1051–1063.

15 S. Park, J. O. Baker, M. E. Himmel, P. A. Parilla and D. K. Johnson, *Biotechnol. Biofuels*, 2010, **3**, 1–10.

16 A. D. French, D. W. Montgomery, N. T. Prevost, J. V. Edwards and R. J. Woods, *Carbohydr. Polym.*, 2021, **264**, 118004.

17 Y. B. Park, K. Kafle, C. M. Lee, D. J. Cosgrove and S. H. Kim, *Cellulose*, 2015, **22**, 3531–3540.

18 W. A. Sisson, *Science*, 1938, **87**, 350.

19 A. Hirai, M. Tsuji and F. Horii, *Cellulose*, 2002, **9**, 105–113.

20 A. Hirai, M. Tsuji and F. Horii, *Cellulose*, 1997, **4**, 239–245.

21 J. Hayashi, A. Sufoka, J. Ohkita and S. Watanabe, *J. Polym. Sci., Polym. Lett. Ed.*, 1975, **13**, 23–27.

22 E. S. Gardiner and A. Sarko, *Can. J. Chem.*, 1985, **63**, 173–180.

23 J. S. Wilkie, *Nature*, 1961, **190**, 1145–1150.

24 O. L. Sponsler, *J. Gen. Physiol.*, 1926, **9**, 677–695.

25 K. H. Meyer and L. Misch, *Helv. Chim. Acta*, 1937, **20**, 232–244.

26 A. D. French and P. Langan, *Cellulose*, 2014, **21**, 1087–1089.

27 S. Nishikawa and S. Ono, *Proc. Tokyo Math.-Phys. Soc.*, 1913, **7**, 131.

28 S. H. Kim, C. M. Lee and K. Kafle, *Korean J. Chem. Eng.*, 2013, **30**, 2127–2141.

29 P. J. Weimer, J. M. Hackney and A. D. French, *Biotechnol. Bioeng.*, 1995, **48**, 169–178.

30 Y. Habibi, L. A. Lucia and O. J. Rojas, *Chem. Rev.*, 2010, **110**, 3479–3500.

31 S. Andersson, R. Serimaa, T. Paakkari, P. Saranpää and E. Pesonen, *J. Wood Sci.*, 2003, **49**, 531–537.

32 M. Debnath, K. S. Salem, V. Naithani, E. Musten, M. A. Hubbe and L. Pal, *Cellulose*, 2021, **12**, 7981–7994.

33 X. Sun, S. Agate, K. S. Salem, L. Lucia and L. Pal, *ACS Appl. Bio Mater.*, 2020, **4**, 140–162.

34 P. Tyagi, K. S. Salem, M. A. Hubbe and L. Pal, *Trends Food Sci. Technol.*, 2021, **115**, 461–485.

35 S. Al-Zuhair, *Bioresour. Technol.*, 2008, **99**, 4078–4085.

36 Y. Cao and H. Tan, *Enzyme Microb. Technol.*, 2005, **36**, 314–317.

37 N. Lavoine, I. Desloges, A. Dufresne and J. Bras, *Carbohydr. Polym.*, 2012, **90**, 735–764.

38 M. Hall, P. Bansal, J. H. Lee, M. J. Realff and A. S. Bommarius, *FEBS J.*, 2010, **277**, 1571–1582.

39 A. French, *BioResources*, 2022, **17**, 5557–5561.

40 Z. Ling, T. Wang, M. Makarem, M. Santiago Cintrón, H. N. Cheng, X. Kang, M. Bacher, A. Potthast, T. Rosenau, H. King, C. D. Delhom, S. Nam, J. Vincent Edwards, S. H. Kim, F. Xu and A. D. French, *Cellulose*, 2019, **26**, 305–328.

41 M. Borrega, P. Ahvenainen and E. Kontturi, *Cellulose*, 2018, **25**, 6811–6818.

42 Y. Nishiyama, U.-J. Kim, D.-Y. Kim, K. S. Katsumata, R. P. May and P. Langan, *Biomacromolecules*, 2003, **4**, 1013–1017.

43 P. Bansal, M. Hall, M. J. Realff, J. H. Lee and A. S. Bommarius, *Bioresour. Technol.*, 2010, **101**, 4461–4471.

44 R. Evans, A. F. A. Wallis, R. H. Newman, U. C. Roick and I. D. Suckling, *Holzforschung*, 1995, **49**, 498–504.

45 A. Kljun, T. A. S. Benians, F. Goubet, F. Meulewaeter, J. P. Knox and R. S. Blackburn, *Biomacromolecules*, 2011, **12**, 4121–4126.

46 C. Driemeier and G. A. Calligaris, *J. Appl. Crystallogr.*, 2011, **44**, 184–192.

47 A. D. French, *Cellulose*, 2014, **21**, 885–896.

48 A. D. French and M. Santiago Cintrón, *Cellulose*, 2013, **20**, 583–588.

49 R. C. Rowe, A. G. McKillop and D. Bray, *Int. J. Pharm.*, 1994, **101**, 169–172.

50 L. Segal, J. J. Creely, A. E. Martin and C. M. Conrad, *Text. Res. J.*, 1959, **29**, 786–794.

51 A. Thygesen, J. Oddershede, H. Lilholt, A. B. Thomsen and K. Ståhl, *Cellulose*, 2005, **12**, 563–576.

52 T. Liitiä, S. L. Maunu and B. Hortling, *Holzforschung*, 2000, **54**, 618–624.

53 H. M. Rietveld, *J. Appl. Crystallogr.*, 1969, **2**, 65–71.

54 F. Horii, A. Hirai and R. Kitamaru, *Macromolecules*, 1987, **20**, 2117–2120.

55 S. Park, J. O. Baker, M. E. Himmel, P. A. Parilla and D. K. Johnson, *Biotechnol. Biofuels*, 2010, **3**, 1–10.

56 Y. Nishiyama, P. Langan and H. Chanzy, *J. Am. Chem. Soc.*, 2002, **124**, 9074–9082.

57 A. Kirui, Z. Ling, X. Kang, M. C. Dickwella Widanage, F. Mentink-Vigier, A. D. French and T. Wang, *Cellulose*, 2019, **26**, 329–339.

58 U. P. Agarwal, R. S. Reiner and S. A. Ralph, *Cellulose*, 2010, **17**, 721–733.

59 S. H. D. Hullemans, J. M. Van Hazendonk and J. E. G. Van Dam, *Carbohydr. Res.*, 1994, **261**, 163–172.

60 U. P. Agarwal, R. R. Reiner and S. A. Ralph, *J. Agric. Food Chem.*, 2013, **61**, 103–113.

61 A. L. Barnette, L. C. Bradley, B. D. Veres, E. P. Schreiner, Y. B. Park, J. Park, S. Park and S. H. Kim, *Biomacromolecules*, 2011, **12**, 2434–2439.

62 C. M. Lee, N. M. A. Mohamed, H. D. Watts, J. D. Kubicki and S. H. Kim, *J. Phys. Chem. B*, 2013, **117**, 6681–6692.

63 M. S. Bertran and B. E. Dale, *J. Appl. Polym. Sci.*, 1986, **32**, 4241–4253.

64 A. Kljun, T. A. S. Benians, F. Goubet, F. Meulewaeter, J. P. Knox and R. S. Blackburn, *Biomacromolecules*, 2011, **12**, 4121–4126.

65 M. M. Rehman, M. Zeeshan, K. Shaker and Y. Nawab, *Cellulose*, 2019, **26**, 9057–9069.

66 H.-J. Kim, S. Roy and J.-W. Rhim, *J. Environ. Chem. Eng.*, 2021, **9**, 106043.

67 L. T. Fan, Y.-H. Lee and D. H. Beardmore, *Biotechnol. Bioeng.*, 1980, **22**, 177–199.

68 Y. Nishiyama, J. Sugiyama, H. Chanzy and P. Langan, *J. Am. Chem. Soc.*, 2003, **125**, 14300–14306.

69 A. D. French, *Cellulose*, 2020, **27**, 5445–5448.

70 P. Langan, Y. Nishiyama and H. Chanzy, *Biomacromolecules*, 2001, **2**, 410–416.



71 P. Ahvenainen, I. Kontro and K. Svedström, *Cellulose*, 2016, **23**, 1073–1086.

72 K. S. Salem, H. Jameel, L. Lucia and L. Pal, *Mater. Today Sust.*, 2023, **22**, 100342.

73 W. Yao, Y. Weng and J. M. Catchmark, *Cellulose*, 2020, **27**, 5563–5579.

74 W. R. Kunusa, I. Isa, L. A. Laliyo and H. Iyabu, *J. Phys.: Conf. Ser.*, 2018, **1028**, 012199.

75 S. Y. Oh, I. Y. Dong, Y. Shin, C. K. Hwan, Y. K. Hak, S. C. Yong, H. P. Won and H. Y. Ji, *Carbohydr. Res.*, 2005, **340**, 2376–2391.

76 R. P. Oliveira and C. Driemeier, *J. Appl. Crystallogr.*, 2013, **46**, 1196–1210.

77 K. Das, D. Ray, N. R. Bandyopadhyay and S. Sengupta, *J. Polym. Environ.*, 2010, **18**, 355–363.

78 M. Holtz, C. Hauf, J. Weisshaupt, A.-A. H. Salvador, M. Woerner and T. Elsaesser, *Struct. Dyn.*, 2017, **4**, 054304.

79 V. W. Tripp, *Cellulose and cellulose derivatives*, 1971, vol. 5, pp. 305–323.

80 P. Scherrer, *Nachr. Ges. Wiss. Goettingen, Math.-Phys. Kl.*, 1918, **1918**, 98–100.

81 Y. Wang, Y. Zhao and Y. Deng, *Carbohydr. Polym.*, 2008, **72**, 178–184.

82 G. V. Gusev, *Polym. Sci. U.S.S.R.*, 1978, **20**, 1295–1297.

83 L. E. Hult, T. Iversen and J. Sugiyama, *Cellulose*, 2003, **10**, 103–110.

84 C. J. Garvey, I. H. Parker and G. P. Simon, *Macromol. Chem. Phys.*, 2005, **206**, 1568–1575.

85 J. He, S. Cui and S-Y. Wang, *J. Appl. Polym. Sci.*, 2008, **107**, 1029–1038.

86 Y. Nishiyama, J. Sugiyama, H. Chanzy and P. Langan, *J. Am. Chem. Soc.*, 2003, **125**, 14300–14306.

87 H.-P. Fink, D. Hoffmann and B. Philipp, *Cellulose*, 1995, **2**, 51–70.

88 J. He, S. Cui and S. Wang, *J. Appl. Polym. Sci.*, 2008, **107**, 1029–1038.

89 E. Gümüskaya, M. Usta and H. Kirci, *Polym. Degrad. Stab.*, 2003, **81**, 559–564.

90 L. Lutterotti, M. Bortolotti, G. Ischia, I. Lonardelli and H.-R. Wenk, *Zeitschrift für Kristallographie Supplements*, 2007, **26**, 125–130.

91 H. M. Rietveld, *J. Appl. Crystallogr.*, 1969, **2**, 65–71.

92 A. D. French, *Cellulose*, 2020, **27**, 5445–5448.

93 M. Wada, T. Okano and J. Sugiyama, *J. Wood Sci.*, 2001, **47**, 124–128.

94 A. Isogai and R. H. Atalla, *J. Polym. Sci., Part A: Polym. Chem.*, 1991, **29**, 113–119.

95 C. Driemeier, *Cellulose*, 2014, **21**, 1065–1073.

96 Y. Nishiyama, G. P. Johnson and A. D. French, *Cellulose*, 2012, **19**, 319–336.

97 A. L. Barnette, C. Lee, L. C. Bradley, E. P. Schreiner, Y. B. Park, H. Shin, D. J. Cosgrove, S. Park and S. H. Kim, *Carbohydr. Polym.*, 2012, **89**, 802–809.

98 S. Wu, X. Jiang, H. Jiang, S. Wu, S. Ding and Y. Jin, *Cellulose*, 2021, **28**, 1947–1959.

99 N. Terinte, R. Ibbett and K. C. Schuster, *Lenzinger Ber.*, 2011, **89**, 118–131.

100 U. P. Agarwal, S. A. Ralph, C. Baez, R. S. Reiner and S. P. Verrill, *Cellulose*, 2017, **24**, 1971–1984.

101 S. Nam, A. D. French, B. D. Condon and M. Concha, *Carbohydr. Polym.*, 2016, **135**, 1–9.

102 Z. Ling, S. Chen, X. Zhang, K. Takabe and F. Xu, *Sci. Rep.*, 2017, **7**, 1–11.

103 Z. Ling, S. Chen, X. Zhang and F. Xu, *Bioresour. Technol.*, 2017, **224**, 611–617.

104 X. Guo, F. Yang, H. Liu, Y. Hou, Y. Wang, J. Sun, X. Chen, Y. Liu and X. Li, *Macromol. Res.*, 2019, **27**, 377–385.

105 T. Paakkari, M. Blomberg, R. Serimaa and M. Järvinen, *J. Appl. Crystallogr.*, 1988, **21**, 393–397.

106 W. Ren, F. Guo, J. Zhu, M. Cao, H. Wang and Y. Yu, *Cellulose*, 2021, **28**, 5993–6005.

107 S. Nomura, Y. Kugo and T. Erata, *Cellulose*, 2020, **27**, 3553–3563.

108 S. Haslinger, S. Hietala, M. Hummel, S. L. Maunu and H. Sixta, *Carbohydr. Polym.*, 2019, **207**, 11–16.

109 M. Tyufekchiev, A. Kolodziejczak, P. Duan, M. Foston, K. Schmidt-Rohr and M. T. Timko, *Green Chem.*, 2019, **21**, 5541–5555.

110 T. Sparrman, L. Svenningsson, K. Sahlin-Sjövold, L. Nordstierna, G. Westman and D. Bernin, *Cellulose*, 2019, **26**, 8993–9003.

111 H. Zhao, J. H. Kwak, Z. Conrad Zhang, H. M. Brown, B. W. Arey and J. E. Holladay, *Carbohydr. Polym.*, 2007, **68**, 235–241.

112 M. Foston, *Curr. Opin. Biotechnol.*, 2014, **27**, 176–184.

113 N. E. Jacobsen, *NMR spectroscopy explained: simplified theory, applications and examples for organic chemistry and structural biology*, John Wiley & Sons, 2007.

114 J. W. Akitt and B. E. Mann, *NMR and Chemistry: An introduction to modern NMR spectroscopy*, CRC Press, 2017.

115 D. C. Apperley, R. K. Harris and P. Hodgkinson, *Solid-state NMR: Basic principles and practice*, Momentum Press, 2012.

116 L. Schröder and C. Faber, *In vivo NMR Imaging*, Humana Press, 2011.

117 S. H. Kim, C. M. Lee and K. Kafle, *Korean J. Chem. Eng.*, 2013, **30**, 2127–2141.

118 K. Zia, T. Siddiqui, S. Ali, I. Farooq, M. S. Zafar and Z. Khurshid, *Eur. J. Dent.*, 2019, **13**, 124–128.

119 Introduction to Solid-State NMR Spectroscopy|Wiley, <https://www.wiley.com/en-us/Introduction+to+Solid+State+NMR+Spectroscopy-p-9781405109147>, accessed April 4, 2023.

120 W. L. Earl and D. L. VanderHart, *Macromolecules*, 1981, **14**, 570–574.

121 D. L. VanderHart and R. H. Atalla, *Macromolecules*, 1984, **17**, 1465–1472.

122 M. Foston, *Curr. Opin. Biotechnol.*, 2014, **27**, 176–184.

123 R. H. Newman and J. A. Hemmingson, *Holzforschung*, 1990, **44**, 351–356.

124 H. Yang, T. Wang, D. Oehme, L. Petridis, M. Hong and J. D. Kubicki, *Cellulose*, 2018, **25**, 23–36.

125 P. T. Larsson, K. Wickholm and T. Iversen, *Carbohydr. Res.*, 1997, **302**, 19–25.



126 K. Wickholm, P. T. Larsson and T. Iversen, *Carbohydr. Res.*, 1998, **312**, 123–129.

127 C. Y. Liang and R. H. Marchessault, *J. Polym. Sci.*, 1959, **39**, 269–278.

128 M. El Hariri El Nokab, M. H. Habib, Y. A. Alassmy, M. M. Abduljawad, K. M. Alshamrani and K. O. Sebakhy, *Polymers*, 2022, **14**, 1049.

129 M. Ghosh, B. P. Prajapati, R. K. Suryawanshi, K. Kishor Dey and N. Kango, *Chem. Phys. Lett.*, 2019, **727**, 105–115.

130 A. Palme, A. Idström, L. Nordstierna and H. Breliid, *Cellulose*, 2014, **21**, 4681–4691.

131 S. Park, D. K. Johnson, C. I. Ishizawa, P. A. Parilla and M. F. Davis, *Cellulose*, 2009, **16**, 641–647.

132 R. Teeäär, R. Serimaa and T. Paakkari, *Polym. Bull.*, 1987, **17**, 231–237.

133 R. H. Atalla and D. L. VanderHart, *Solid State Nucl. Magn. Reson.*, 1999, **15**, 1–19.

134 F. Horii, A. Hirai and R. Kitamaru, *J. Carbohydr. Chem.*, 1984, **3**, 641–662.

135 S. Sternhell and M. Mehrling, *Org. Magn. Reson.*, 1983, **21**, 770.

136 C. Rondeau-Mourot, B. Bouchet, B. Pontoire, P. Robert, J. Mazoyer and A. Buléon, *Carbohydr. Polym.*, 2003, **53**, 241–252.

137 L. R. Schroeder, V. M. Gentile and R. H. Atalla, *J. Wood Chem. Technol.*, 1986, **6**, 1–14.

138 R. H. Newman and J. A. Hemmingson, *Cellulose*, 1995, **2**, 95–110.

139 K. Masuda, M. Adachi, A. Hirai, H. Yamamoto, H. Kaji and F. Horii, *Solid State Nucl. Magn. Reson.*, 2003, **23**, 198–212.

140 A. N. Fernandes, L. H. Thomas, C. M. Altaner, P. Callow, V. T. Forsyth, D. C. Apperley, C. J. Kennedy and M. C. Jarvis, *Proc. Natl. Acad. Sci. U. S. A.*, 2011, **108**, E1195–E1203.

141 M. Foston, R. Katahira, E. Gjersing, M. F. Davis and A. J. Ragauskas, *J. Agric. Food Chem.*, 2012, **60**, 1419–1427.

142 I. Jankowska, P. Ławniczak, K. Pogorzelec-Glaser, A. Łapiński, R. Pankiewicz and J. Tritt-Goc, *Mater. Chem. Phys.*, 2020, **239**, 122056.

143 G. Cavallaro, A. Agliolo Gallitto, L. Lisuzzo and G. Lazzara, *Cellulose*, 2019, **26**, 8853–8865.

144 A. L. P. Queiroz, B. M. Kerins, J. Yadav, F. Farag, W. Faisal, M. E. Crowley, S. E. Lawrence, H. A. Moynihan, A.-M. Healy, S. Vucen and A. M. Crean, *Cellulose*, 2021, **28**, 8971–8985.

145 U. P. Agarwal, S. A. Ralph, C. Baez and R. S. Reiner, *Cellulose*, 2021, **28**, 9069–9079.

146 U. P. Agarwal, S. A. Ralph, C. Baez and R. S. Reiner, *Biomacromolecules*, 2021, **22**, 1357–1373.

147 B. Stuart, *Analytical techniques in the sciences. Infrared Spectroscopy: Fundamentals and Applications*, 2004.

148 C. N. Banwell and M. Elaine, *Fundamentals of molecular spectroscopy*, McGraw-Hill, New York, 1994.

149 P. Larkin, *Infrared and Raman spectroscopy: principles and spectral interpretation*, Elsevier, 2017.

150 C. L. Haynes, A. D. McFarland and R. P. Van Duyne, *Anal. Chem.*, 2005, 338–346.

151 P. Rostron, S. Gaber and D. Gaber, *Laser*, 2016, **21**, 24.

152 D. Srivastava, M. S. Kuklin, J. Ahopelto and A. J. Karttunen, *Carbohydr. Polym.*, 2020, **243**, 116440.

153 A. Jorio, R. Saito, J. H. Hafner, C. M. Lieber, M. Hunter, T. McClure, G. Dresselhaus and M. S. Dresselhaus, *Phys. Rev. Lett.*, 2001, **86**, 1118–1121.

154 S. J. Eichhorn and R. J. Young, *Compos. Sci. Technol.*, 2004, **64**, 767–772.

155 R. J. Young, D. Lu and R. J. Day, *Polym. Int.*, 1991, **24**, 71–76.

156 J. Široký, R. S. Blackburn, T. Bechtold, J. Taylor and P. White, *Cellulose*, 2010, **17**, 103–115.

157 A. A. Saapan, S. H. Kandil and A. M. Habib, *Text. Res. J.*, 1984, **54**, 863–867.

158 R. T. O'Connor, E. F. DuPré and D. Mitcham, *Text. Res. J.*, 1958, **28**, 382–392.

159 D. Ciolacu, F. Ciolacu and V. I. Popa, *Cellul. Chem. Technol.*, 2011, **45**, 13.

160 M. Fan, D. Dai and B. Huang, *Fourier transform-materials analysis*, 2012.

161 M. L. Nelson and R. T. O'Connor, *J. Appl. Polym. Sci.*, 1964, **8**, 1325–1341.

162 T. Kondo and C. Sawatari, *Polymer*, 1996, **37**, 393–399.

163 U. Richter, T. Krause and W. Schempp, *Die Angewandte Makromolekulare Chemie, Appl. Macromol. Chem. Phys.*, 1991, **185**, 155–167.

164 J. Li, X. Wei, Q. Wang, J. Chen, G. Chang, L. Kong, J. Su and Y. Liu, *Carbohydr. Polym.*, 2012, **90**, 1609–1613.

165 I. Reiniati, A. N. Hrymak and A. Margaritis, *Biochem. Eng. J.*, 2017, **127**, 21–31.

166 M. J. Shultz, *Advances In Multi-Photon Processes And Spectroscopy*, 2008, pp. 133–199.

167 M. Makarem, C. M. Lee, K. Kafle, S. Huang, I. Chae, H. Yang, J. D. Kubicki and S. H. Kim, *Cellulose*, 2019, **26**, 35–79.

168 J. Choi, J. Lee, M. Makarem, S. Huang and S. H. Kim, *J. Phys. Chem. B*, 2022, **126**, 6629–6641.

169 I. Chae, S. M. Q. Bokhari, X. Chen, R. Zu, K. Liu, A. Borhan, V. Gopalan, J. M. Catchmark and S. H. Kim, *Carbohydr. Polym.*, 2021, **255**, 117328.

170 Y. R. Shen, *Nature*, 1989, **337**, 519.

171 S. A. Denev, T. T. A. Lummen, E. Barnes, A. Kumar and V. Gopalan, *J. Am. Ceram. Soc.*, 2011, **94**, 2699–2727.

172 H. C. Hieu, N. A. Tuan, H. Li, Y. Miyauchi and G. Mizutani, *Appl. Spectrosc.*, 2011, **65**, 1254–1259.

173 A. G. Lambert, P. B. Davies and D. J. Neivandt, *Appl. Spectrosc. Rev.*, 2005, **40**, 103–145.

174 V. Cedric, C. Yves and P. Andre, *Biosensors – Emerging Materials and Applications*, 2011.

175 G. L. Richmond, *Chem. Rev.*, 2002, **102**, 2693–2724.

176 C. T. Williams and D. A. Beattie, *Surf. Sci.*, 2002, **500**, 545–576.

177 O. J. Rojas, *Cellulose Chemistry and Properties: Fibers, Nanocelluloses and Advanced Materials*, 2016.

178 C. M. Lee, A. Mittal, A. L. Barnette, K. Kafle, Y. B. Park, H. Shin, D. K. Johnson, S. Park and S. H. Kim, *Cellulose*, 2013, **20**, 991–1000.



179 Y. B. Park, C. M. Lee, B. W. Koo, S. Park, D. J. Cosgrove and S. H. Kim, *Plant Physiol.*, 2013, **163**, 907–913.

180 C. M. Lee, K. Kafle, Y. B. Park and S. H. Kim, *Phys. Chem. Chem. Phys.*, 2014, **16**, 10844–10853.

181 S. Huang, S. N. Kiemle, M. Makarem and S. H. Kim, *Cellulose*, 2020, **27**, 57–69.

182 M. Makarem, C. M. Lee, D. Sawada, H. M. O'Neill and S. H. Kim, *J. Phys. Chem. Lett.*, 2018, **9**, 70–75.

183 T. De Assis, S. Huang, C. E. Driemeier, B. S. Donohoe, C. Kim, S. H. Kim, R. Gonzalez, H. Jameel and S. Park, *Biotechnol. Biofuels*, 2018, **11**, 1–11.

184 M. Makarem, D. Sawada, H. M. O'Neill, C. M. Lee, K. Kafle, Y. B. Park, A. Mittal and S. H. Kim, *J. Phys. Chem. C*, 2017, **121**, 10249–10257.

185 X. Chen, C. M. Lee, H. F. Wang, L. Jensen and S. H. Kim, *J. Phys. Chem. C*, 2017, **121**, 18876–18886.

186 K. Kafle, R. Shi, C. M. Lee, A. Mittal, Y. B. Park, Y. H. Sun, S. Park, V. Chiang and S. H. Kim, *Cellulose*, 2014, **21**, 2219–2231.

187 L. Kong, C. Lee, S. H. Kim and G. R. Ziegler, *J. Phys. Chem. B*, 2014, **118**, 1775–1783.

188 P. P. Handakumbura, D. A. Matos, K. S. Osmont, M. J. Harrington, K. Heo, K. Kafle, S. H. Kim, T. I. Baskin and S. P. Hazen, *BMC Plant Biol.*, 2013, **13**, 1–6.

189 K. Kafle, C. M. Lee, H. Shin, J. Zoppe, D. K. Johnson, S. H. Kim and S. Park, *BioEnergy Res.*, 2015, **8**, 1750–1758.

190 J. Xie, Y. Li, W. Wang, S. Pan, N. Cui and J. Liu, *Adv. Mech. Eng.*, 2020, **44**, 125001.

191 C. Schick, *Anal. Bioanal. Chem.*, 2009, **395**, 1589–1611.

192 G. W. H. Hljhne, *Thermochim. Acta*, 1991, **187**, 283–292.

193 H. J. Flammersheim, W. F. Hemminger and G. W. H. Hohne, *Differential Scanning Calorimetry*, 2003.

194 Y. Kong and J. N. Hay, *Polymer*, 2002, **43**, 3873–3878.

195 S. Chapi, S. Raghu and H. Devendrappa, *Ionics*, 2016, **22**, 803–814.

196 M. S. Bertran and B. E. Dale, *J. Appl. Polym. Sci.*, 1986, **32**, 4241–4253.

197 T. Hatakeyama and H. Hatakeyama, *Thermal Properties of Green Polymers*, Hot Topics in Thermal Analysis and Calorimetry, 2005.

198 D. J. Blundell, D. R. Beckett and P. H. Willcocks, *Polymer*, 1981, **22**, 704–707.

199 V. B. F. Mathot and M. F. J. Pijpers, *Thermochim. Acta*, 1989, **151**, 241–259.

200 L. Szcześniak, A. Rachocki and J. Tritt-Goc, *Cellulose*, 2008, **15**, 445–451.

201 K. Aïssa, Doctoral dissertation, University of British Columbia, 2020.

202 G. T. Beckham, J. F. Matthews, B. Peters, Y. J. Bomble, M. E. Himmel and M. F. Crowley, *J. Phys. Chem. B*, 2011, **115**, 4118–4127.

203 S. Gao, C. You, S. Renneckar, J. Bao and Y.-H. P. Zhang, *Biotechnol. Biofuels*, 2014, **7**, 24.

204 J. Hong, X. Ye and Y.-H. P. Zhang, *Langmuir*, 2007, **23**, 12535–12540.

205 T. Li, N. Liu, X. Ou, X. Zhao, F. Qi, J. Huang and D. Liu, *Biotechnol. Biofuels*, 2018, **11**, 105.

206 A. Kljun, T. A. S. Benians, F. Goubet, F. Meulewaeter, J. P. Knox and R. S. Blackburn, *Biomacromolecules*, 2011, **12**, 4121–4126.

207 J. Široký, R. S. Blackburn, T. Bechtold, J. Taylor and P. White, *Cellulose*, 2010, **17**, 103–115.

208 J. Široký, T. A. S. Benians, S. J. Russell, T. Bechtold, J. Paul Knox and R. S. Blackburn, *Carbohydr. Polym.*, 2012, **89**, 213–221.

209 X. Guo, F. Yang, H. Liu, Y. Hou, Y. Wang, J. Sun, X. Chen, Y. Liu and X. Li, *Macromol. Res.*, 2019, **27**, 377–385.

210 T. Kawakubo, S. Karita, Y. Araki, S. Watanabe, M. Oyadomari, R. Takada, F. Tanaka, K. Abe, T. Watanabe, Y. Honda and T. Watanabe, *Biotechnol. Bioeng.*, 2010, **105**, 499–508.

211 M. Ioelovich, A. Leykin and O. Figovsky, *BioResources*, 2010, **5**, 3.

212 J. W. Hearle and W. E. Morton, *Physical properties of textile fibres*, Elsevier, 2008.

213 M. Ioelovich, *Am. J. BioSci.*, 2014, **2**, 6.

214 M. Ioelovich, *ChemXpress*, 2016, **9**, 245–251.

215 L. E. Hessler and R. E. Power, *Text. Res. J.*, 1954, **24**, 822–827.

216 M. Ioelovich, *BioResources*, 2009, **4**, 3.

217 C. Lee, K. Dazen, K. Kafle, A. Moore, D. K. Johnson, S. Park and S. H. Kim, in *Cellulose Chemistry and Properties: Fibers, Nanocelluloses and Advanced Materials*, ed. O. J. Rojas, Springer International Publishing, Cham, 2016, pp. 115–131.

