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Easy and accurate computation of energy barriers for carbocation solvation: an expeditious tool to face carbocation chemistry†

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An expeditious procedure for the challenging computation of the free energy barriers (ΔG^\ddagger) for the solvation of carbocations is presented. This procedure is based on Marcus Theory (MT) and the popular B3LYP/6-31G(d)//PCM method, and it allows the easy, accurate and inexpensive prediction of these barriers for carbocations of very different stability. This method was validated by the fair mean absolute error (ca. 1.5 kcal mol⁻¹) achieved in the prediction of 19 known experimental barriers covering a range of ca. 50 kcal mol⁻¹. Interestingly, the new procedure also uses an original method for the calculation of the required inner reorganization energy (λ) and free energy of reaction (ΔG). This procedure should pave the way to face computationally the pivotal issue of carbocation chemistry and could be easily extended to any bimolecular organic reaction.

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1. Introduction

Carbocations are molecular organic ions containing positively charged carbon atoms, making them electron-deficient species and hence highly reactive electrophiles.^{1,2} They can be found as long-lived species in very acidic liquid-phase media,³ but they have also been found to exist in different environments such as planetary atmospheres and other astrophysical objects.⁴

Carbocations are pivotal in organic chemistry because of their role as reactive intermediates in many organic transformations, such as the common unimolecular nucleophilic substitution (S_N1) and elimination (E₁) reactions.^{1–5} Moreover, carbocations are involved in the biosynthetic routes of a plethora of natural products like polyketides and terpenes.⁶

Since solvation, especially water solvation, is an essential topic for understanding conventional carbocation chemistry in solution, there has been a good deal of experimental and theoretical work devoted to the determination of solvation–reaction barriers.

Theoretical work on this issue has implied a notorious computational effort. So far, many aspects of the chemistry

and properties of carbocations, such as thermochemistry, non-bonded interactions and reaction kinetics, have been studied using *ab initio* Density Functional Theory (DFT) and hybrid quantum mechanics-molecular mechanics (QM/MM) methods.⁶ However, the important subject of the reactivity with nucleophiles was performed not by computational methods but by linear free-energy relationships (LFERs),^{7,8} like the Mayr equation.⁹ The first successful computations of the pseudo-first order rate constants (k_w) for the water solvation of carbocations were performed by Guthrie *et al.* for benzyl and benzyl-like cations based on no barrier theory (NBT).¹⁰ Unfortunately, the NBT method requires the computationally expensive optimization of eight structures, four before and four after the supposed transition state (TS) structure, according to a cubic reaction diagram.¹⁰ In this context, the need for a simple and efficient computational method, allowing the accurate prediction of the experimental outcome of these processes, is obvious and has been sought for a long time ago. Herein, we demonstrate that the free energy barriers (ΔG^\ddagger) for the solvation of carbocations with very different stability can be accurately computed by using a new procedure based on Marcus Theory (MT),¹¹ which is much simpler than the NBT method. This simplicity should pave the way to computationally study the fundamental carbocation chemistry.

2. Theoretical background

We have previously proposed that the reaction of the *tert*-butyl cation with water (eqn (1); where R = *t*-Bu, and H_{aq}⁺ represents a water solvated proton) is a two-step process,^{12a} which is

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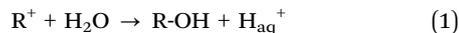
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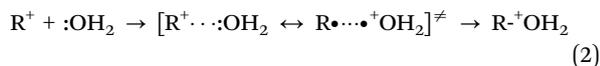
† Electronic supplementary information (ESI) available: Cartesian coordinates of the B3LYP optimized geometries for carbocations and carbocation–water complexes, as well as SMD energies (Table S1). See DOI: <https://doi.org/10.1039/d3cp03544a>



supported by the accurate computational prediction of the energy barrier for the proton transfer involved (second reaction step). This precise prediction was done by DFT calculations on the basis of the MT.^{12a}



The first step of this reaction is now being studied using DFT and MT, which could serve to easily predict energy barriers for carbocation solvation by water and other nucleophiles. For this purpose, we propose that such a first reaction step (see eqn (2)) consists of the pyramidalization of the water-solvated planar (sp^2) carbon carbocation, which occurs as an inner-sphere reorganization process according to the MT for adiabatic reactions.^{12a} Once the TS structure (shown in brackets in eqn (2)) is achieved, single electron transfer (SET) from the involved water molecule takes place without any structural modification according to the Franck-Condon principle.^{12a} After relaxation, the corresponding short-lived protonated alcohol, R-OH_2^+ , is generated as the final state (FS) of the reaction. It must be stressed here that the complexes shown within brackets in eqn (2) are resonance structures, involving C-O distances shorter than that corresponding to the final (relaxed) R-OH_2^+ species. This constitutes an unprecedented mechanistical view of this relevant process. Thus, eqn (2) implies a new mechanism for carbocation solvation by water (or other protic solvents), which is a fundamental process having implications in both Organic Chemistry and Biochemistry, as mentioned in the introduction section. Thus, this mechanism, based on a SET (see TS in eqn (2)), is interpreted by us in terms of a modified version of Marcus theory (MT), in order to be used in a different context to the one in which the original MT was devised.



It must be noted that the quadratic Marcus Equation (eqn (3)) gives the free energy of activation of a SET reaction (ΔG^\ddagger) as a nonlinear function of the free energy of reaction (ΔG) and other energy parameters (ω_{R} , ω_{P} and λ).^{12a,13} Usually, the difference between the work to bring the reactants together into the reactive complex (ω_{R}), and the negative work done when the products complex separates to afford the free products (ω_{P}), is negligible compared to ΔG .¹¹ Parameter λ is the barrier height when $\Delta G = 0$ (intrinsic barrier) and has no equivalent in variational transition state theory (VTST).¹⁴ The value of λ is given by the reorganization energy, λ , being $\lambda = 4A$.¹³ In our case, λ is the sum of the inner-sphere reorganization energy of the reactants (λ_{i}) and the outer-sphere electronic reorganization energy of the solvent (λ_0^{EL}).^{12a,13}

$$\Delta G^\ddagger = \lambda \left(1 + \frac{\Delta G - \omega_{\text{R}} + \omega_{\text{P}}}{4\lambda} \right)^2 \quad (3)$$

Fig. 1 schematically illustrates the adaptation of the generic Marcus energy diagram to the reaction mechanism proposed by us for the carbocation solvation reaction (eqn (2)).

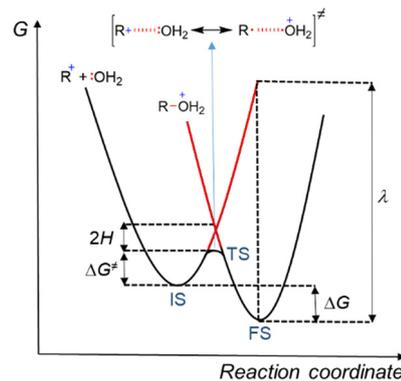


Fig. 1 Adiabatic crossing (TS) of the MT parabolas for the proposed exergonic SET-based reaction mechanism (see eqn (2)). The tunneling rate is proportional to $2H$.

A_{i} value is usually computed by the quadratic variational method, which requires knowledge of the geometries of the initial (IS) and final (FS) reaction states.^{15,16} However, determining the equilibrium geometry of the IS is challenging in the case of intermolecular reactions, as our case is, and usually performed either according to time dependent DFT (TD-DFT) or scanning methods.^{12a,d,e} These scanning methods usually guess the geometry of the TS and then determine the IS and FS geometries by following the intrinsic reaction coordinate (IRC) from the assumed TS in the proper direction.^{12e} To address better the determination of A_{i} , we now propose a new and very simple procedure on the basis of using eqn (4), where $G(\text{R}^+)$ and $G(\text{R}_{\text{F}}^+)$ are the free energies of the fully optimized carbocation and the corresponding frozen carbocation, respectively, the latter resulting from deleting the water molecule from the TS structure. This “frozen” carbocation represents the adaptation of the planar carbocation geometry to the sp^3 one, which is required to enable the proposed SET process (see eqn (2)) without any further geometrical changes. In all cases, the most stable conformation of R^+ is used in this approach.

$$A_{\text{i}} = G(\text{R}^+) - G(\text{R}_{\text{F}}^+) \quad (4)$$

On the other hand, the electronic reorganization energy of the solvent, $\lambda_0^{\text{EL}} = \lambda_0^{\text{EL}}/4$, can be easily calculated by eqn (5), according to the simplified two-sphere model,¹⁵ where parameter a is the radius of the corresponding reactant equivalent sphere, ϵ_0 is the vacuum permittivity, and ϵ_{op} and ϵ_{s} are the optical and static dielectric constants of water, respectively.

$$\lambda_0^{\text{EL}} = \frac{e^2}{4\pi\epsilon_0} \left(\frac{1}{\epsilon_{\text{op}}} - \frac{1}{\epsilon_{\text{s}}} \right) \frac{1}{2a} \quad (5)$$

Finally, for determining the ΔG value, we use eqn (6), which involves energies of frozen carbocations and corresponds to an exergonic reaction between non-equilibrium species.

$$\Delta G = G([\text{R}_{\text{F}}^+ \cdots \text{:OH}_2]^{\ddagger}) - [G(\text{R}_{\text{F}}^+) + G(\text{H}_2\text{O})] \quad (6)$$



It is interesting to note that, using instead standard conditions (eqn (7)), results in too high barriers (ΔG^\ddagger).

$$\Delta G = G(\text{R-OH}_2^+) - [G(\text{R}^+) + G(\text{H}_2\text{O})] \quad (7)$$

To determine the required optimal water-complex structure (eqn (6)), we have conducted a relaxed scanning of the involved C–O distance with a step size of 0.1 Å. In the case of the highly stable trimethoxymethyl cation, $(\text{MeO})_3\text{C}^+$, we found a C–O distance at the TS equal to 1.688 Å (see ESI†). In order to simplify our methodology, we used this distance for all the studied cases. As shown in the Results and Discussion section, this guess seems to be a reasonable choice in the frame of the MT.

3. Computational methods

The required computations were performed using the computationally inexpensive B3LYP/6-31G(d) method implemented in GAUSSIAN 09,¹⁷ since it has been successfully used by us in similar problems related to carbocation chemistry.^{12a–d} Nonetheless, considering that the main drawback of the B3LYP functional when computing large species seems to be due to self-interaction effects for long range electron–electron interactions,¹⁸ and this concern could be resolved using long-range corrected (LC) functionals,¹⁹ we have also used the well-known LC MPWPW91 method for comparison purposes.²⁰ Additionally, both methods were combined with the popular SMD and PCM solvation models, which are also implemented in GAUSSIAN 09,¹⁷ using the default water dielectric constant value (78.3553), which is usually accepted for this solvent at 25 °C.

4. Results and discussion

As a benchmark for proving our simplified MT procedure for the calculation of the energy barriers for carbocation solvation, we selected water solvation and an ample number of well-known carbocations (Fig. 2), ranging from highly stabilized carbocations (*e.g.*, $(\text{MeO})_3\text{C}^+$) to unstable ones (*e.g.*, benzyl cation, Bn^+). It must be noted here that, according to the classification of McClelland,²¹ highly stabilized carbocations are those whose pseudo-first order water solvation rates (k_w) can be experimentally determined by stopped-flow spectroscopy, whereas stabilized carbocation kinetics needs to be followed by laser flash photolysis. As limit cases of instability, we included the *tert*-butyl (*t*-Bu⁺) and Bn^+ carbocations, whose stabilities can be only determined experimentally by indirect methods (see references in Table 1).

4.1. B3LYP/6-31G(d) calculations

The B3LYP/6-31G(d)//PCM-calculated ΔG^\ddagger , ΔG and A values, and the corresponding experimental barriers ($\Delta G^\ddagger(\text{exp})$), are shown in Table 1. As mentioned in Section 2, we used the computed C–O distance at the TS of highly stable $(\text{MeO})_3\text{C}^+$ (1.688 Å) for facing the computations of all the selected

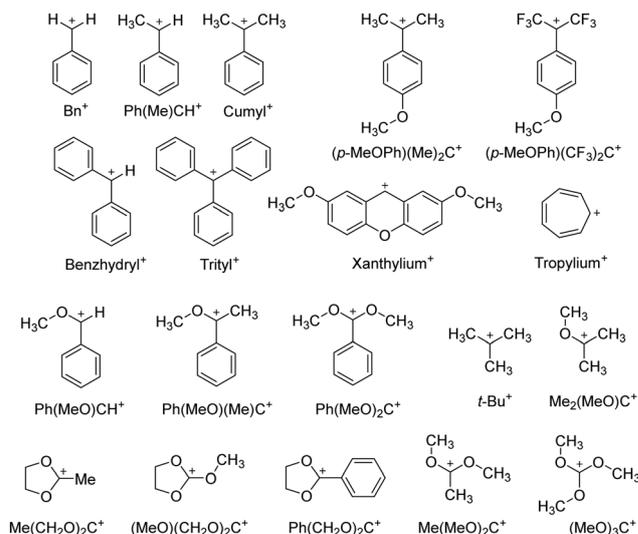


Fig. 2 Benchmark set of carbocations and used abbreviations.

Table 1 B3LYP/6-31G(d)//PCM calculated (ΔG^\ddagger) and experimental ($\Delta G^\ddagger(\text{exp})$) values for the water solvation of a selected set of carbocations (R^+ ; see Fig. 2), together with the corresponding calculated ΔG and A values required to use Marcus eqn (3), in water solution. Previously reported ΔG^\ddagger values calculated using the NBT procedure ($\Delta G^\ddagger(\text{NBT})$) are included for comparison purposes. All data in kcal mol⁻¹. The symmetry group of the most stable R^+ conformation used for the computations is given into the parenthesis

Entry	R^+	A_i	A_o^{EL}	ΔG	ΔG^\ddagger	$\Delta G^\ddagger(\text{NBT})^a$	$\Delta G^\ddagger(\text{exp})$
1	<i>t</i> -Bu ⁺ (C_{3v})	15.5	6.40	−49.8	4.03	—	3.90 ^b
2	Bn^+ (C_{2v})	11.0	5.71	−48.4	1.30	1.26	1.47 ^a
3	$\text{Ph}(\text{Me})\text{CH}^+$ (C_s)	14.3	5.37	−41.8	4.06	3.28	2.15 ^a
4	Benzhydryl ⁺ (C_2)	14.2	4.74	−33.3	5.94	—	5.02 ^c
5	Cumyl ⁺ (C_s)	13.8	4.85	−37.0	4.73	4.11	3.2 ^a
6	$(p\text{-MeOPh})(\text{CF}_3)_2\text{C}^+$ (C_s)	16.9	4.60	−37.2	8.60	4.80	7.07 ^c
7	$(p\text{-MeOPh})(\text{Me})_2\text{C}^+$ (C_s)	16.7	4.87	−30.4	9.06	9.91	7.85 ^c
8	Trityl ⁺ (C_{3v})	14.2	4.32	−20.3	9.74	—	10.3 ^c
9	Tropylium ⁺ (D_7)	23.5	5.35	−18.3	20.5	—	17.0 ^c
10	$\text{Ph}(\text{MeO})_2\text{C}^+$ (C_1)	15.7	5.00	−29.0	9.81	—	10.6 ^c
11	$\text{Ph}(\text{CH}_2\text{O})_2\text{C}^+$ (C_{2v})	14.0	4.90	−19.4	10.4	—	11.3 ^c
12	$(\text{MeO})_3\text{C}^+$ (C_{3h})	17.6	5.35	−20.2	14.0	—	11.8 ^c
13	$(\text{MeO})(\text{CH}_2\text{O})_2\text{C}^+$ (C_1)	18.8	5.90	−24.0	14.2	—	11.4 ^c
14	$\text{Me}(\text{MeO})_2\text{C}^+$ (C_1)	20.8	5.77	−27.7	12.4	—	10.5 ^c
15	$\text{Me}(\text{CH}_2\text{O})_2\text{C}^+$ (C_s)	16.2	5.60	−26.2	10.6	—	9.1 ^c
16	$\text{Ph}(\text{MeO})(\text{Me})\text{C}^+$ (C_1)	15.2	5.00	−27.2	8.86	—	6.92 ^c
17	$\text{Me}_2(\text{MeO})\text{C}^+$ (C_s)	16.2	6.08	−23.7	8.48	—	5.17 ^c
18	$\text{Ph}(\text{MeO})\text{CH}^+$ (C_s)	14.0	5.92	−34.0	6.28	—	4.85 ^c
19	Xanthylum ⁺ (C_{2v})	13.8	14.4	−15.4	10.1	—	9.58 ^d

^a Data from ref. 10a. ^b Data from ref. 22. ^c Data from ref. 21. ^d Mean value of those reported in ref. 21 and 23.

carbocations (Fig. 2). Curiously enough, this distance is the same as that computationally found by us for protonated *tert*-butyl alcohol (an unstable carbocation) at equilibrium. On the other hand, selecting shorter distances, for example 1.522 Å, corresponding to protonated methanol at the equilibrium, produces too high computed barriers for all the selected cases (*e.g.*, 16.8 kcal mol⁻¹ for $(\text{MeO})_3\text{C}^+$), whereas using larger



distances (for example 1.822 Å) affords too low barriers in all the cases (e.g., 9.7 kcal mol⁻¹ for (MeO)₃C⁺). However, just by using critical 1.688 Å allows the accurate computation of ΔG^\ddagger for the water solvation of the 19 selected carbocations (Fig. 2), covering an ample ΔG^\ddagger range (between 1.3 kcal mol⁻¹ and 20.5 kcal mol⁻¹), as shown in Table 1.

The mean absolute error (MAE) of our calculated ΔG^\ddagger values in relation to the experimental ones is 1.51 kcal mol⁻¹, within a range of ca. 20 kcal mol⁻¹. This ΔG^\ddagger range corresponds to k_w values between 6.16×10^{12} s⁻¹ and 14.3 s⁻¹, for Bn⁺ and xanthylum⁺, respectively, at 25 °C. Hence, the accordance between experimental and calculated values can be qualified as excellent, not far away from the so-called “chemical accuracy” (1.0 kcal mol⁻¹). It must be noted that the MAE of our computed ΔG^\ddagger values is the same as that obtained from previously reported data computed using the more expensive NBT procedure in a more restricted set of carbocations (see Table 1).

It is worth mentioning that the ΔG^\ddagger values computed according to our method (see Table 1) predict the kinetically stabilizing stereoelectronic effect exerted by alkyl (see Table 1 and cf. data in entries 2,3 and 5), phenyl (cf. entries 2 and 4) and alkoxy groups (cf. entries 5 and 10) when attached to the cationic centre. These effects are cast, in quantitative terms, by plotting $\log k_w$ vs. the sum of the Hammett–Brown σ^+ constants for the attached groups, as suggested by Kresge *et al.*²³ However, as highlighted by McClelland, a simple linear free-energy relationship (LFER) is not clearly evident.²¹ Thus, for example, selecting carbocations bearing H, alkyl and alkoxy groups attached to the carbocationic centre produces a linear correlation with a slope, ρ^+ , equal to 6.1; whereas selecting carbocations bearing a single phenyl produces a different LFER with a lower slope (5.5).²¹ Moreover, increasing the number of phenyl substituents in benzyl carbocation (Ph₂CH⁺ and Ph₃C⁺) generates points out of the latter line.²¹ In this context, it is clear that our simple computational procedure is free of these LFER inconsistencies. Even more, it is also striking that alkyl substitution at the carbocationic centre results in reactivity changes that follow neither the steric nor the electronic effects of the substituents.¹⁷

Another intriguing fact explained by our computations is the lack of kinetic destabilization associated with the involvement of electron-withdrawing groups directly attached to the carbocationic center. Thus, Richard *et al.* observed that (*p*-MeOPh)(CF₃)₂C⁺ is only slightly more reactive than (*p*-MeOPh)(Me)₂C⁺ (see ΔG^\ddagger (exp) in Table 1, entries 6 and 7).¹³ This fact was cleverly rationalized by Richard considering that the strong electron-withdrawing polar effect exerted by the trifluoromethyl groups in the case of (*p*-MeOPh)(CF₃)₂C⁺ (destabilizing charge-dipole interaction with the carbocationic center) is however almost relieved by an enhanced stabilizing charge-delocalization effect exerted by the methoxy group located at *para* position in the aromatic ring (resonance effect).¹³ However, our computation of the HOMOs of the water complexes of (*p*-MeOPh)(CF₃)₂C⁺ and (*p*-MeOPh)(Me)₂C⁺ at their solvation TSs shows that there is no enhanced resonance effect. Thus, these HOMOs (see Fig. 3 and ESI[†]) resemble the Ψ_3 orbital of benzene, showing a through-bond nodal plane which electronically isolates the methoxy group

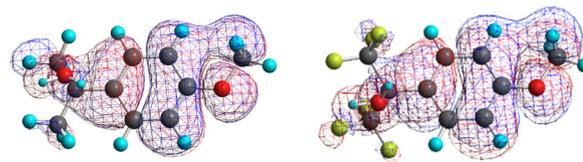


Fig. 3 Electron density contour maps ($\rho = 10^{-3}$) of HOMOs for the water complexes of (*p*-MeOPh)(Me)₂C⁺ (left) and (*p*-MeOPh)(CF₃)₂C⁺ (right) at the solvation reaction TS.

from the carbocation center. Moreover, our computations reveal another explanation for the unexpectedly high stability of the (*p*-MeOPh)(CF₃)₂C⁺ carbocation: the back-donation of the trifluoromethyl carbons to the cationic center, as supported by the significant π -bonding interaction found between the trifluoromethyl carbons and the cationic one in the computed HOMO (see Fig. 3). More importantly, our computations agree better with the experimental results than the NBT approach (cf. the corresponding ΔG^\ddagger , ΔG^\ddagger (NBT) and ΔG^\ddagger (exp) data in Table 1) and give a quantitative answer to the question: the destabilizing polar effect exerted by the trifluoromethyl groups is behind the high exergonicity of the (*p*-MeOPh)(CF₃)₂C⁺ water-solvation reaction when compared to that of (*p*-MeOPh)(Me)₂C⁺ ($\Delta G = -37.2$ kcal mol⁻¹ vs. -30.4 kcal mol⁻¹; see Table 1) and such an effect is compensated by the working back-donation effect, which results in just a slightly higher solvation reactivity for (*p*-MeOPh)(CF₃)₂C⁺ ($\Delta G^\ddagger = 8.60$ kcal mol⁻¹ vs. 9.06 kcal mol⁻¹; see Table 1).

The relative stability of *t*-Bu⁺ and Bn⁺, which is an often-discussed question, can be accounted for by our procedure. From calculations of an isodesmic reaction involving both carbocations in the gas-phase, it was previously concluded that *t*-Bu⁺ is ca. 6 kcal mol⁻¹ more stable than Bn⁺,²⁴ despite the highly stabilizing resonance effect exerted by the phenyl group in the case of the Bn⁺.

However, according to the relative ion-stability scale of Abboud *et al.*,⁸ the relative stabilities of both carbocations in the gas-phase (ΔG_r°) are almost the same, considering the error range ($\Delta G_r^\circ = -6.0 \pm 1.0$ and -5.8 ± 1.0 kcal mol⁻¹ for Bn⁺ and *t*-Bu⁺, respectively). This situation is very different in solution. Thus, from extrapolation of experimental solvolytic rate results of different substrates in different solvents, *tert*-butyl chloride should solvolyze ca. 176 times faster than benzyl chloride at 45 °C in 50% aqueous EtOH, also suggesting that *t*-Bu⁺ is more stable than Bn⁺ in solution too.²⁴ This fact could be explained by the strongly stabilizing inductive (polar) and hyperconjugative effects exerted by the methyl groups of *t*-Bu⁺, overweighing the said phenyl resonance effect. Our computations in water at 25 °C support these solvolytic conclusions. Thus, Bn⁺ should react with water (at 25 °C) ca. 100 times faster than *t*-Bu⁺, as calculated by us from the corresponding computed ΔG^\ddagger values (see Table 1, entries 1 and 2) using the Eyring equation.²⁵ In the gas phase, at the same temperature, this difference is predicted to be even larger (3.65×10^5 times faster; see Table 2). The main ground for this difference seems to be the high exergonicity of the reaction with water for Bn⁺ (C_{2v}) in the gas phase (see Table 1, entries 1 and 2, and Table 2, and cf. ΔG values).



The conformational analysis of *t*-Bu⁺ is especially complex. Thus, a previous study in the gas phase at the MP2(full)/6-31G** level, conducted by v. R. Schleyer *et al.*,²⁶ revealed that the energy barrier for CH₃-rotation is extremely low and that the C_s conformation, which is energetically very close to the C_{3h} one, is the global minimum, the latter being preferred by *ca.* 1.2 kcal mol⁻¹ to the C_{3v} conformation. According to our computations in the gas phase, C_{3h} is also 1.52 kcal mol⁻¹ more stable than the C_{3v} structure. However, in solution (PCM), we found that the barrier to CH₃-rotation is even lower than in the gas phase, but now the C_{3v} conformation is slightly (by only 0.4 kcal mol⁻¹) more stable than the C_{3h} one.

The effect of cyclization on the reaction rates of acyclic oxocarocations can also be detected by our method. Thus, McClelland explained the different water-solvation reactivity of acyclic Ph(OMe)₂C⁺ and cyclic Ph(CH₂O)₂C⁺ ($\Delta G^\ddagger(\text{exp}) = 10.6$ and 11.3 kcal mol⁻¹, respectively; see entries 10 and 11 in Table 1) on the basis of a differential steric interaction involving the phenyl group, and inhibiting its full stabilizing resonance effect in the acyclic case.²¹ Our computations agree with a higher reactivity for acyclic Ph(OMe)₂C⁺ (see Table 1). This differential effect does not exist in the related Me(MeO)₂C⁺/Me(CH₂O)₂C⁺ couple bearing methyl instead of phenyl, as evidenced by the corresponding $\Delta G^\ddagger(\text{exp})$ values (higher for the acyclic partner; *cf.* entries 14 and 15 in Table 1). Once again, our computational method reproduces this situation with computational accuracy (Table 1). Finally, the related (MeO)₃C⁺/(MeO)-(CH₂O)₂C⁺ couple can be considered an intermediate case between the previous ones, due to the presence of a methoxyl group instead of phenyl or methyl. This group, as the phenyl one, has the ability to exert a stabilizing resonance effect, but this effect should not be significantly different for both the cyclic and the acyclic cations, due to its small steric effect. As a matter of fact, the corresponding $\Delta G^\ddagger(\text{exp})$ values are very similar (11.8 and 11.4 kcal mol⁻¹, respectively; see entries 12 and 13 in Table 1), and our computations reproduce this similarity within the above-said accuracy as well.

Notably, our method can be extended to other nucleophiles instead of water. As an example, we have computed $\Delta G^\ddagger = 4.71$ kcal mol⁻¹ for the reaction of *t*-Bu⁺ with azide anion in water solution at 298 K. The calculated MT parameters are

Table 2 B3LYP/6-31G(d)//PCM ΔG^\ddagger values for the water solvation of *t*-Bu⁺ and Bn⁺ (see Fig. 2), together with the corresponding calculated ΔG and *A* values required to use Marcus eqn (3), in the gas phase (all data in kcal mol⁻¹), along with the resulting rate constants, *k* (in s⁻¹), at 25 °C. The symmetry group of the most stable R⁺ conformation used for the computations is given in parenthesis

R ⁺	<i>A</i> _i	<i>A</i> ₀ ^{EL,a}	ΔG	ΔG^\ddagger	<i>k</i> ^b
<i>t</i> -Bu ⁺ (C _{3v})	17.3	0.00	-5.40	14.7	1.03×10^2
Bn ⁺ (C _{2v})	13.1	0.00	-13.9	7.11	3.76×10^7

^a *A*₀^{EL} is 0 because the optical (ϵ_{op}) and static (ϵ_s) dielectric constants of vacuum are both equal to 1 (see eqn (6)). ^b As calculated by applying the Eyring equation from the corresponding ΔG^\ddagger values.

$A = 30.9$ kcal mol⁻¹, $\Delta G = -90.6$ kcal mol⁻¹ and $A_0^{\text{EL}} = 5.28$ kcal mol⁻¹. Based on the Eyring equation, such an energy barrier allowed the calculation of a pseudo-first order rate (*k*_w) for this reaction equal to 2.18×10^9 s⁻¹. This value is in good agreement with that found by Richie *et al.* (5×10^9 s⁻¹) for the pseudo-first order diffusion-controlled rate constant of any reactive carbocation with azide anions, which is based on the ‘‘azide clock method’’ for the experimental determination of rate constants for reactions of carbocations with water.^{21,27,28}

4.2. B3LYP/6-31G(d), B3LYP/6-31+G(d,p) and MPWPW91/6-31+G(d,p) calculations

The results obtained when applying B3LYP/6-31G(d)//SMD to our benchmark set of carbocations are summarized in Table S1 in the SI. In all cases, the obtained ΔG^\ddagger values are much higher than the experimental ones, and no correlation between both sets of values is found.

In order to study the possible relevance of the long-range (LR) interactions in the studied reaction, we used the well-known MPWPW91/6-31+G(d,p) functional, which takes into account this kind of interaction.²⁰ As an example, the hydration of Benzhydryl⁺ can be selected, since it exhibits intramolecular π - π and H-H interactions between the phenyl rings. For this case, MPWPW91/6-31+G(d,p)//PCM calculates $A_i = 17.8$ kcal mol⁻¹ and $\Delta G = -35.4$ kcal mol⁻¹, giving place to $\Delta G^\ddagger = 8.49$ kcal mol⁻¹ (Table 3). This computational barrier is significantly higher than the experimental one ($\Delta G^\ddagger(\text{exp}) = 5.02$ kcal mol⁻¹; see entry 4 in Table 1). To diminish this computational deviation, the C-O distance at the TS ($[\text{R}^+ \cdots \text{OH}_2]^\ddagger$) should be enlarged with respect to the standard 1.688 Å value (see above). Moreover, this enlargement is substrate-dependent at variance with that occurring when the B3LYP functional was used. All this represents a serious complication in the computations when using the LR MPWPW91/6-31+G(d,p)//PCM method and supports the simpler use of the B3LYP/6-31G(d)//PCM one.

Interestingly, as shown in Table 3, the selection of the solvation model (PCM *vs.* SMD) is much more critical than the inclusion of LR effects in the selected functional. The origin of this fact seems to be a failure in the SMD model, giving place to more positive ΔG values than the PCM one (see Table 1, and Table S1 in the ESI,[†] and *cf.* related ΔG values). In our opinion, it is not worth using LC-corrected methods for the calculation of the ΔG^\ddagger in systems as large as Trityl⁺ and Xanthylum⁺ (see Fig. 2) exhibiting no protobranching.²⁹

In order to shed light on the influence of the basis set, we have also computed the energy barrier for the water solvation of

Table 3 ΔG^\ddagger values (in kcal mol⁻¹) for the water solvation of C₂ benzhydryl⁺ (see Fig. 2), using different methods

Entry	Method	ΔG^\ddagger
1	B3LYP/6-31G(d)//PCM	5.94
2	B3LYP/6-31+G(d,p)//PCM	0.82
3	MPWPW91/6-31+G(d,p)//PCM	8.49
4	B3LYP/6-31G(d)//SMD	16.0
5	MPWPW91/6-31+G(d,p)//SMD	17.3



benzhydryl⁺ using the B3LYP/6-31+G(d,p)//PCM methodology, which includes diffuse functions, while keeping the mentioned selected C–O distance and the PCM. As shown in Table 3, the computed barrier (0.82 kcal mol⁻¹) is now extremely low, and far away from the experimental one (5.02 kcal mol⁻¹; see Table 1). Thus, in this case, a shorter C–O distance should be used to get a reliable result, as a consequence of including diffuse functions (*cf.* entries 1 and 2 in Table 3). From all these results, we conclude that there is a C–O distance appropriate for each model chemistry.

5. Conclusions

The prediction of the fundamental reactivity of carbocations with nucleophiles is usually performed by using LFERs. However, this can be faulty, owing to specific stereoelectronic interactions that cannot be easily parameterized. On the other hand, the alternative azide-clock predictive method requires accurate experimental measurements. Regarding computational methods, the only computational procedure available to date for the prediction of reaction rates of carbocation solvations was the computationally expensive four-dimensional one based on the NBT. However, our methodology has been proven to be a convenient alternative to the NBT one, because it is simpler and affords a similar MAE at a much lower computational cost. On the other side, the found low MAE when computing carbocation water-solvation barriers (1.5 kcal mol⁻¹ for 19 cases covering a very broad range of ΔG^\ddagger values), as well as the demonstrated accuracy when predicting kinetics for reactions involving other nucleophiles different of water (azide anion), validates our method as highly appropriate to establish accurate reactivity-structure relationships in carbocation chemistry, which is central in organic chemistry.

Moreover, a very simple procedure is established for computationally determining the internal reorganization energy and the free energy of the reaction of carbocations with water and other nucleophiles. In this context, it must be noted that the difficulty in determining these energies is the main drawback limiting the application of the Marcus equations to the study of a plethora of fundamental intermolecular reactions, mainly within the organic framework, which could be now rapidly and accurately addressed, from the computational point of view, on the basis of this procedure.

Also, a relevant point is the idea that the C–O distance for the transition states for carbocation water solvation shouldn't change significantly from a study case to another. As a result, this makes affordable the application of the methodology to a good deal of organic systems with a moderate computational effort. Finally, the good agreement of the computed energy barriers with the experimental ones supports the proposed SET-based mechanism for this fundamental reaction.

Author contributions

The manuscript was written through contributions of all authors. All authors have given approval to the final version of the manuscript.

Conflicts of interest

There are no conflicts of interest to declare.

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