



Cite this: *Phys. Chem. Chem. Phys.*,  
2023, 25, 25619

# Formation of $\text{H}_3\text{O}^+$ and OH by $\text{CO}_2$ and $\text{N}_2\text{O}$ trace gases in the atmospheric environment†

Daniele Catone,<sup>a</sup>  <sup>a</sup> Mattea Carmen Castrovilli,<sup>b</sup>  <sup>b</sup> Francesca Nicolanti,<sup>c</sup>  
Mauro Satta  <sup>\*de</sup> and Antonella Cartoni  <sup>\*e</sup>

The impact of cosmic rays' energetic subatomic particles on climate and global warming is still controversial and under debate. Cosmic rays produce ions that can trigger fast reactions affecting chemical networks in the troposphere and stratosphere especially when a large amount of relevant trace gases such as carbon dioxide, methane, sulfur dioxide and water are injected by volcanic eruptions. This work focuses on synchrotron experiments and an *ab initio* theoretical study of the ion chemistry of carbon dioxide and nitrous oxide radical cations reacting with water. These molecules catalyze a fast exothermic formation of hydronium ions  $\text{H}_3\text{O}^+$  and the hydroxyl radical OH, the main oxidant in the atmosphere. Moreover, theoretical calculations demonstrate that at the end of the catalytic cycle,  $\text{CO}_2$  and  $\text{N}_2\text{O}$  are produced vibrationally excited and subsequently they quench in the microsecond time scale by collision with the surrounding atmospheric molecules at the pressure and temperature of the upper-troposphere/stratosphere. The chemistry involved in these reactions has a strong impact on the oxidant capacity of the atmosphere, on the sulfate aerosol production, on the cloud formation and eventually on the chemical networks controlling climate and global warming models.

Received 26th May 2023,  
Accepted 11th August 2023

DOI: 10.1039/d3cp02427j

rsc.li/pccp

## 1. Introduction

The role of ionic processes occurring in the atmosphere is still a challenging and controversial topic because of their potential effect on the formation and depletion of many chemical species relevant to climate models.<sup>1</sup> Despite the fundamental role of neutrals in many chemical processes, ion–molecule reactions<sup>2</sup> are also important because these reactions destroy or produce species through faster reactions. In the troposphere and stratosphere ions are formed by Cosmic rays' energetic subatomic particles that enter Earth's atmosphere from space.<sup>3</sup> The ionization rate due to cosmic rays is at its maximum value of about  $20\text{--}40\text{ cm}^{-3}\text{ s}^{-1}$  in the upper troposphere/lower stratosphere layer (as shown for instance in Fig. 1 of ref. 3). Ionization of

gaseous air molecules, mainly  $\text{N}_2$  and  $\text{O}_2$ , produces primary ions which afterwards react with other atmospheric molecules. Moreover, secondary particles such as protons, electrons and muons, lose their energy ionizing species as they propagate downwards. Svensmark<sup>4</sup> has studied for many years the connection between cosmic rays, clouds, and climate. These studies are based on the observation that when there is a change in solar activity and hence in the flux of cosmic rays there is a change in climate because cosmic rays influence cloud formations which affects Earth's temperature. Indeed, cosmic rays can produce ions that stabilize and favor aerosol formation which leads to clouds. CLOUD Project Experiments performed in Ginevra at CERN from 2006 were also prepared to understand the influence of galactic cosmic rays (GCRs) on aerosols and clouds, and their implications for climate.<sup>5</sup> They want to find out which trace gases and ions are involved in aerosol particle formation. However, this topic, due to its complex nature, remains the object of controversial studies based on laboratory measurements and on satellite data analysis.<sup>6</sup> In the atmosphere, composed of a mixture of gases, the main components are  $\text{N}_2$  and  $\text{O}_2$  but also trace gases such as carbon dioxide, methane, water, nitrous oxide, ozone and sulfur dioxide are relevant in the atmospheric chemical networks. These gases are present in very little amount but cannot be neglected. They are emitted in the atmosphere by natural or anthropogenic sources and could have a key role in

<sup>a</sup> Istituto di Struttura della Materia - CNR (ISM-CNR), Area della Ricerca di Roma 2, Via del Fosso del Cavaliere 100, 00133, Rome, Italy

<sup>b</sup> Istituto di Struttura della Materia - CNR (ISM-CNR), Area della Ricerca di Roma 1, Monterotondo Scalo 00015, Italy

<sup>c</sup> Department of Physics, Sapienza University of Rome, P. le Aldo Moro 5, Rome, 00185, Italy

<sup>d</sup> Institute for the Study of Nanostructured Materials-CNR (ISMN-CNR), Department of Chemistry, Sapienza University of Rome, P. le Aldo Moro 5, Rome, 00185, Italy. E-mail: mauro.satta@cnr.it; Tel: + 39 06 49913381

<sup>e</sup> Department of Chemistry, Sapienza University of Rome, P. le Aldo Moro 5, Rome, 00185, Italy. E-mail: antonella.cartoni@uniroma1.it; Tel: + 39 06 49913678

† Electronic supplementary information (ESI) available. See DOI: <https://doi.org/10.1039/d3cp02427j>



chemical reactions that change the budget of climate relevant species with consequent alteration of the atmospheric composition with potentially social and health severe effects. Chlorofluorocarbon molecules are a classic example of such trace gases. They interact with sun light and produce chlorine radicals that catalytically destroy the ozone molecules.<sup>7</sup> One of the main cation present in the upper atmosphere is hydronium ion, the seed of water organic cluster whose climate effects are debated.<sup>8</sup> The formation of  $\text{H}_3\text{O}^+$  should be considered especially in regions of the atmosphere, as upper troposphere<sup>9</sup> where water and trace gases are present and secondary subatomic particles can trigger ionization. Indeed, the atmosphere is a complex system subject to sudden and heterogeneous spatial and temporal changes in its chemistry and dynamics. For instance, recently, a large amount of water vapor has been injected into the stratosphere by the eruption of the submarine volcano Hunga Tonga–Hunga Ha’apai that has increased the amount of global stratospheric water vapor by more than 5–10%.<sup>10,11</sup> This rare event can affect the chemistry of the stratosphere for several years and it should be considered, given that water vapor is the most abundant trace gas able to influence the climate by absorption and emission of infrared radiant energy.<sup>12,13</sup> Volcanic eruptions also eject other trace gases, such as carbon dioxide, carbon monoxide and sulfur dioxide that could have a role in alternative chemical networks active in this modified atmosphere.

In this scenario, the present manuscript focuses on an experimental and theoretical study of the ion chemistry of carbon dioxide and nitrous oxide radical cation with water. These molecules, as observed in the reaction of sulfur dioxide radical cation with water,<sup>14–16</sup> can catalyze a fast production of hydronium ions and hydroxyl radical OH, the main oxidant present in the atmosphere.

## 2. Methods

### 2.1. Synchrotron experiments

The radiation available at the Circular Polarization beamline (CiPo) at the synchrotron Elettra (Trieste) has been used to produce carbon dioxide and nitrous oxide radical cations with different internal energies. The beamline already described in previous works,<sup>17–19</sup> is equipped with a Normal Incidence Monochromator (NIM) and an electromagnetic elliptical undulator/wiggler providing monochromatized radiation photons in the 8–40 eV energy range. In these experiments, the aluminium grating of the NIM, operating in the energy range 8–18 eV has been used, with an energy resolution of about 20 meV. The photon energy was calibrated against the autoionization features observed in the Ar total photoionization cross-section between the 3p spin orbit components.<sup>20</sup> Carbon dioxide or nitrous oxide was introduced in the ion source through a leak valve and the molecules were ionized by synchrotron radiation at a pressure of about  $10^{-6}$ – $10^{-5}$  mbar. The  $\text{CO}_2^{\bullet+}$  or  $\text{N}_2\text{O}^{\bullet+}$  were then guided into the octupole reaction cell with several optical lenses at the nominal collision energy (CE) of 0 eV with an energy spread of about 100–150 meV. Water was introduced

into the reaction cell (octupole) at room temperature and at different nominal pressures. The photoionization efficiencies curve (PIEC) of both  $\text{N}_2\text{O}$  and  $\text{CO}_2$  were acquired by scanning the photon energy from their ionization energies to about 17 eV with a step of 0.02 eV and an acquisition time of 15 seconds per point. The PIEC of  $\text{NO}^+$  from  $\text{N}_2\text{O}^+$  is also reported and acquired in the 14.0–18.0 eV photon energy range with a 0.02 eV step and an acquisition time of 10 seconds per point. Mass spectra were acquired at the photon energies of 13.5 ( $\text{N}_2\text{O}$ ) and 14.0 ( $\text{CO}_2$ ) eV before and after the introduction of water in the reaction cell and in the mass over charge ( $m/z$ ) range 10–50. The measurements have been performed at different water pressures in the  $10^{-6}$ – $10^{-5}$  mbar range and with acquisition time ranging from 1 to 5 seconds per point. Both reactions produce protonated molecules  $\text{CO}_2\text{H}^+$  ( $m/z = 45$ ),  $\text{N}_2\text{OH}^+$  ( $m/z = 45$ ) and  $\text{H}_2\text{O}^{\bullet+}$  ( $m/z = 18$ ) with a tiny signal at  $m/z = 19$  due to hydronium ion  $\text{H}_3\text{O}^+$  that become the most intense peak at high pressures. The intensities of the reagent and product ions were then acquired by scanning the photon energy in the range of 13.8–15.0 eV and 12.9–13.7 eV for  $\text{CO}_2$  and  $\text{N}_2\text{O}$  respectively, with a step of 0.10 eV and an acquisition time of 30 seconds per point. The pressure measurements are affected by an error of about 30%. For the sake of clarity, the  $\bullet$  in radicals and radical cation is omitted in the following sections. Linearity of the product/reagent ratio with the pressure has been verified at different photon energies up to about  $8.0 \times 10^{-5}$  mbar.

**Materials.** All the samples were used at room temperature. Carbon dioxide  $\text{CO}_2$  was purchased from SIAD with purity > 99.99% and  $\text{N}_2\text{O}$  from Aldrich with purity of 99%.  $\text{H}_2\text{O}$ , suitable for HPLC, was purchased from Sigma-Aldrich.

### 2.2. Theoretical calculations

**2.2.1. Potential energy surfaces of the reactions.** Both the charge (CT) and hydrogen (HT) transfer reactions of  $\text{CO}_2^+$  and  $\text{N}_2\text{O}^+$  with water have been studied by analysing their energetics and their branching-ratio with *ab initio* calculations and micro-canonical transition state theory. The formation of the two molecular adducts  $[\text{H}_2\text{O}-\text{CO}_2]^+$  and  $[\text{N}_2\text{O}-\text{H}_2\text{O}]^+$  is assumed to be in the Langevin regime since the literature experimental overall rate coefficient<sup>21</sup> (HT plus CT processes) are quite in agreement with their respective Langevin rate coefficient as discussed in the next paragraph. Therefore, the bottleneck of the reaction is the formation of the molecular adducts as controlled by capture theory.<sup>22</sup> Hence, the Minimum Energy Path (MEP) relevant for calculating the branching ratio of the HT vs. CT channels is the one that connects the molecular adduct with the products of the charge and hydrogen transfer reactions.

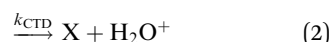
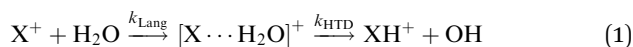
In particular the *ab initio* calculations have employed the Density Functional Theory Double-hybrid approach, which is able to describe the radical nature of the species reacting during the charge and hydrogen transfer processes. The functional used in all calculations is the B2PLYP of Grimme<sup>23</sup> and the basis set employed is the aug-cc-pVTZ which is based on triple zeta, including polarization, tight



core and diffuse functions.<sup>24</sup> All frequency calculations have been performed using the harmonic approximation and the electronic structure calculations have been done with the Gaussian code.<sup>25</sup>

The present reactions occur *via* open shell radical paths which are properly described with the level of calculation here used, as reported by the partial charge and spin analysis in the next paragraph. The scan of the reactive potential energy surfaces has been accomplished by scanning both the HO–H and H<sub>2</sub>O–OCO (H<sub>2</sub>O–ONN) coordinates with a variable step whose minimum value has been taken as 0.02 Å. All the other geometrical coordinates, except the scanning coordinates, have been optimized during the scans. The spin and charge population are described with the Mulliken analysis of the electron density.<sup>26</sup>

**2.2.2. Microcanonical branching ratio.** The HT and CT reactions can be written as two parallel reactions (X = CO<sub>2</sub>, N<sub>2</sub>O):



where  $k_{\text{Lang}}$  is the Langevin rate coefficient for the formation of the capture complex, and  $k_{\text{HTD}}$  and  $k_{\text{CTD}}$  are the microcanonical unimolecular dissociation rate coefficients for the HT and CT channels respectively. The overall rate coefficients for the HT or CT (denoted by HT/CT) reactions are given by:  $k_{\text{HT/CT}} = k_{\text{Lang}} \cdot k_{\text{HTD/CTD}}$ .

These two last unimolecular rate coefficients can be written in the adiabatic rotation approximation of the microcanonical transition state theory<sup>27</sup> as:

$$k_{\text{HT/CT}} = \frac{k_{\text{Lang}}}{\rho(E_{\text{complex}})} \cdot N_{\text{OH/X}}^v(E_{\text{OH/X}}^v) \cdot N_{\text{HX}^+/\text{H}_2\text{O}^+}^v(E_{\text{HX}^+/\text{H}_2\text{O}^+}^v) \cdot Q_{\text{OH/X}}^{\text{rot}} \cdot Q_{\text{HX}^+/\text{H}_2\text{O}^+}^{\text{rot}} \cdot N_{\text{HT/CT}}^{\text{tras}}(E_{\text{HT/CT}}^{\text{kin}}) \quad (3)$$

where  $\rho(E_{\text{complex}})$  is the density of states of the reactive complex at its  $E_{\text{complex}}$  energy,  $N^v$  are the number of vibrational states for the products calculated at the vibrational energy  $E^v$ , and  $Q^{\text{rot}}$  are their rotational partition functions at  $T = 298$  K,  $N^{\text{tras}}$  are the number of translational states for the relative motion of the products for the two reaction channels HT and CT, with their respective reduced mass  $\mu_{\text{HT/CT}}$ :

$$N_{\text{HT/CT}}^{\text{tras}}(E_{\text{HT/CT}}^{\text{kin}}) = \frac{2}{3}\pi \cdot \left( \frac{\sqrt{(2\mu_{\text{HT/CT}})}}{h} \right)^3 \times \left( \sqrt{E_{\text{HT/CT}}^{\text{kin}}} \right)^3 + 1 \quad (4)$$

The kinetic energy is given by three terms:

$$E_{\text{HT/CT}}^{\text{kin}} = \frac{4k_{\text{B}}T}{\pi} + E_{\text{exp}}^{\text{tras}} + E_{\text{HT/CT}}^{\text{tras}} \quad (5)$$

where, the first term is the thermal energy of the reagents, the second term is the experimental collision energy (75 meV) and the third term  $E_{\text{HT/CT}}^{\text{tras}}$  is the relative kinetic energy produced during the reaction, and it is given by energy conservation with respect to the total reaction energy available:

$$E_{\text{INT}}(h\nu) + E_{\text{HT/CT}} = E_{\text{OH/H}_2\text{O}^+}^v + E_{\text{HX}^+/\text{X}}^v + E_{\text{HT/CT}}^{\text{tras}} \quad (6)$$

where  $E_{\text{INT}}(h\nu)$  is the internal energy which remains in the X<sup>+</sup> ion at the photon energy  $h\nu$  (see the ESI† Section S1 and Fig. S1 for N<sub>2</sub>O and Section S3 of the ESI† of ref. 28 for CO<sub>2</sub>), and  $E_{\text{HT/CT}}$  is the reaction energies and  $E^v$  is the vibrational energy of the species.

The branching ratio (BR) of the HT channels have been analysed to get information on the vibrational–translational energy partition of the products by comparing the experimental and theoretical BR. The theoretical BR of HT channel can be written as:

$$\text{BR}_{\text{HT}} = \frac{k_{\text{HT}}}{k_{\text{HT}} + k_{\text{CT}}} \quad (7)$$

which can be explicitly written without the density of states of the reactive complex:

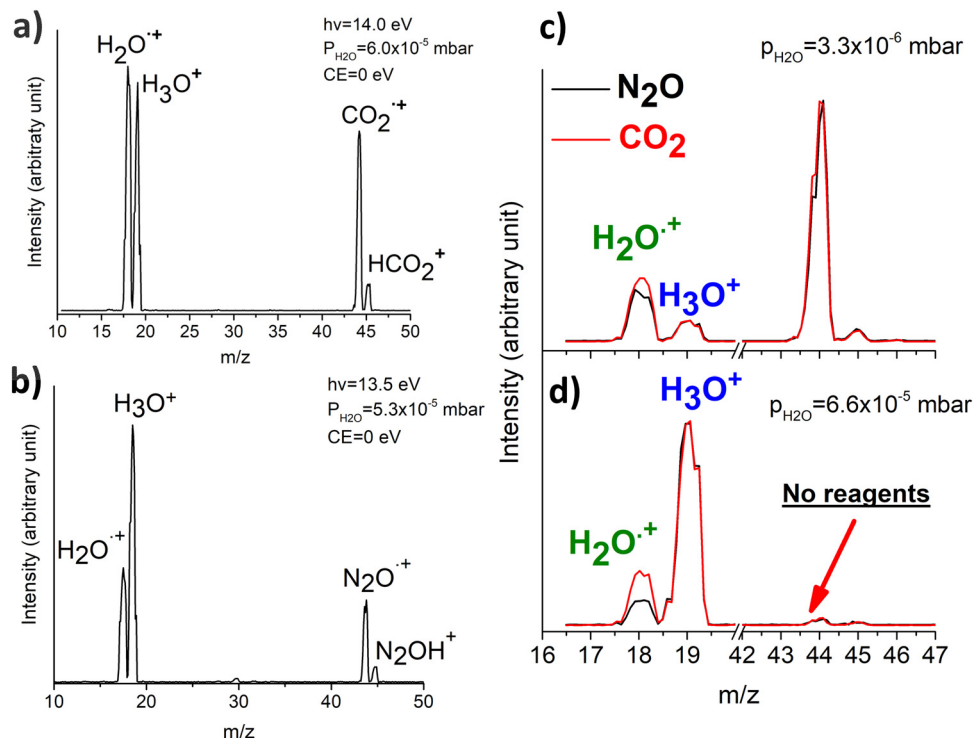
$$\text{BR} = \frac{(N_{\text{OH}}^v(E_{\text{OH}}^v) \cdot N_{\text{HX}^+}^v(E_{\text{HX}^+}^v) \cdot Q_{\text{OH}}^{\text{rot}} \cdot Q_{\text{HX}^+}^{\text{rot}} \cdot N_{\text{HT}}^{\text{tras}}(E_{\text{HT}}^{\text{kin}})) / (N_{\text{OH}}^v(E_{\text{OH}}^v) \cdot N_{\text{HX}^+}^v(E_{\text{HX}^+}^v) \cdot Q_{\text{OH}}^{\text{rot}} \cdot Q_{\text{HX}^+}^{\text{rot}} \cdot N_{\text{HT}}^{\text{tras}}(E_{\text{HT}}^{\text{kin}}) + N_{\text{H}_2\text{O}^+}^v(E_{\text{H}_2\text{O}^+}^v) \cdot N_{\text{X}}^v(E_{\text{X}}^v) \cdot Q_{\text{H}_2\text{O}^+}^{\text{rot}} \cdot Q_{\text{X}}^{\text{rot}} \cdot N_{\text{CT}}^{\text{tras}}(E_{\text{CT}}^{\text{kin}}))}{(N_{\text{OH}}^v(E_{\text{OH}}^v) \cdot N_{\text{HX}^+}^v(E_{\text{HX}^+}^v) \cdot Q_{\text{OH}}^{\text{rot}} \cdot Q_{\text{HX}^+}^{\text{rot}} \cdot N_{\text{HT}}^{\text{tras}}(E_{\text{HT}}^{\text{kin}})) / (N_{\text{OH}}^v(E_{\text{OH}}^v) \cdot N_{\text{HX}^+}^v(E_{\text{HX}^+}^v) \cdot Q_{\text{OH}}^{\text{rot}} \cdot Q_{\text{HX}^+}^{\text{rot}} \cdot N_{\text{HT}}^{\text{tras}}(E_{\text{HT}}^{\text{kin}}) + N_{\text{H}_2\text{O}^+}^v(E_{\text{H}_2\text{O}^+}^v) \cdot N_{\text{X}}^v(E_{\text{X}}^v) \cdot Q_{\text{H}_2\text{O}^+}^{\text{rot}} \cdot Q_{\text{X}}^{\text{rot}} \cdot N_{\text{CT}}^{\text{tras}}(E_{\text{CT}}^{\text{kin}}))} \quad (8)$$

We have calculated the BR for each combination of vibrational–translational energies of the products (energy triplet:  $E_{\text{OH/H}_2\text{O}^+}^v, E_{\text{HX}^+/\text{X}}^v, E_{\text{HT/CT}}^{\text{tras}}$ ) at each photon energy  $h\nu$ , and with a 1.0 cm<sup>-1</sup> energy scan. The optimal energy triplet was selected such that the theoretical BR was the closest to the experimental BR (with a maximum deviation of 2%), whose values were evaluated with the rate coefficients obtained by the procedure described in the ESI† (Section S2 see later).

### 3. Results and discussion

In Fig. 1(a) and (b) the mass spectra of the reactions involving CO<sub>2</sub><sup>+</sup> + H<sub>2</sub>O and N<sub>2</sub>O<sup>+</sup> + H<sub>2</sub>O are shown, respectively. In Fig. 1(a) the peaks at  $m/z = 44, 45, 18$  and  $19$  are due to the ions CO<sub>2</sub><sup>+</sup>, HCO<sub>2</sub><sup>+</sup>, H<sub>2</sub>O<sup>+</sup> and H<sub>3</sub>O<sup>+</sup>, respectively, while in Fig. 1(b) the same  $m/z$  data are referred to N<sub>2</sub>O<sup>+</sup>, N<sub>2</sub>OH<sup>+</sup>, H<sub>2</sub>O and H<sub>3</sub>O<sup>+</sup>. The small peak at  $m/z = 30$  (NO<sup>+</sup>) is due to a minor second order radiation contribution. In Fig. 1(c) and (d) the enlargement of the interested zone is shown for the two systems in the experiments performed under the same experimental conditions and during the same beam time, in two sequential measurements by changing only the neutral gas reagent in the ion source (CO<sub>2</sub> in red and N<sub>2</sub>O in black) and keeping the water pressures fixed at the same nominal values of  $3.3 \times 10^{-6}$  and  $6.6 \times 10^{-5}$  mbar. At high pressure is evident that both the CO<sub>2</sub><sup>+</sup> and N<sub>2</sub>O<sup>+</sup> ions react completely (Fig. 1(d)).





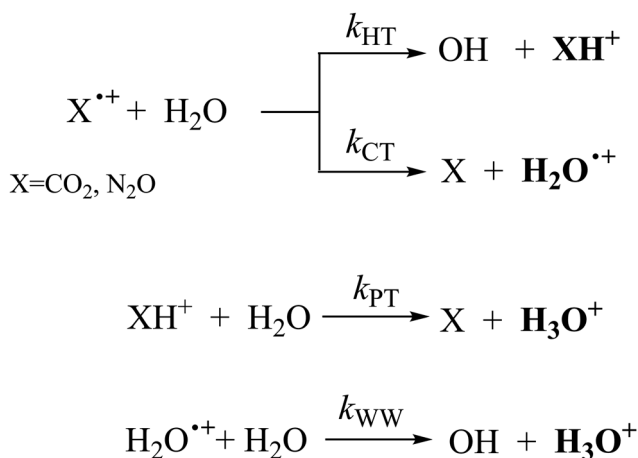
**Fig. 1** (a) Mass spectrum of the reaction between the two reagents  $\text{CO}_2^+ + \text{H}_2\text{O}$  acquired at the  $\text{CE} = 0$ ,  $P_{\text{H}_2\text{O}} = 6.0 \times 10^{-5}$  mbar and photon energy of 14.0 eV. (b) Mass spectrum of the reaction between the two reagents  $\text{N}_2\text{O}^+ + \text{H}_2\text{O}$  acquired at the  $\text{CE} = 0$ ,  $P_{\text{H}_2\text{O}} = 5.3 \times 10^{-5}$  mbar and photon energy of 13.5 eV. (c) and (d) represent reactions of  $\text{CO}_2^+$  and  $\text{N}_2\text{O}^+$  with the water pressure inside the reaction cell of  $3.3 \times 10^{-6}$  mbar and  $6.6 \times 10^{-5}$  mbar, respectively.

By considering the observed ions in the mass spectra (Fig. 1) and the data from literature<sup>21</sup> it is reasonable to assert that the reactions network is depicted as in Scheme 1.

Where PT is the proton transfer process and  $k_{\text{WW}}$  is the rate coefficient of the hydrogen transfer reaction between two water molecules. The literature experimental rate coefficients  $k$  at 300 K are reported in Table 1.<sup>21</sup> These data demonstrate that CT, PT and WW reactions are very fast and at the high pressure all products of the first two reactions,  $\text{XH}^+$  and  $\text{H}_2\text{O}^+$ , react

completely with water leading to the sink ion  $\text{H}_3\text{O}^+$ . The thermochemical  $\Delta H_f^\circ$  data from the literature<sup>29</sup> allow for the calculations of the thermochemistry of all reactions in Scheme 1. In the case of  $\text{CO}_2$  the HT, CT and PT are exothermic by 60.0, 113.0 and 150.1  $\text{kJ mol}^{-1}$  while in the case of  $\text{N}_2\text{O}$  they are exothermic by 10.1, 26.9 and 113.9  $\text{kJ mol}^{-1}$ . Finally, reaction between  $\text{H}_2\text{O}^+$  and  $\text{H}_2\text{O}$  is exothermic by 97.3  $\text{kJ mol}^{-1}$ . It is noteworthy to say that in the HT reaction with  $\text{X} = \text{N}_2\text{O}$  only the oxygen atom  $[\text{HONN}]^+$  has been considered as the HT on N  $[\text{HNNO}]^+$  is endothermic by 15.4  $\text{kJ mol}^{-1}$ .

In order to check that there are not reaction barriers in the dissociation of the reactive complex, the Minimum Energy Path of the HT and CT reactions, both with carbon dioxide and nitrous oxide, have been calculated starting from the structure of the minimum adduct  $[\text{X} \cdots \text{H}_2\text{O}]^+$  (Fig. 2(a) and (b)) at the energies of 208.7 and 136.6  $\text{kJ mol}^{-1}$  for  $\text{CO}_2$  and  $\text{N}_2\text{O}$ , respectively. Once formed, the adduct ions can lead to CT or HT, with CT faster than HT. The reaction coordinate of CT is the distance between the oxygen atoms of water and nitrous oxide  $[\text{H}_2\text{O} \cdots \text{ONN}]^+$  and one of the O of  $\text{CO}_2$   $[\text{H}_2\text{O} \cdots \text{OCO}]^+$ . In the HT reaction, the O-H bond of  $\text{H}_2\text{O}$  has been considered as the reaction coordinate. The spin and charge evolution of the two systems along the MEP from the adduct  $[\text{X} \cdots \text{H}_2\text{O}]^+$  ions are reported in Fig. 2(c) and (d). In the CT reaction, the charge and spin are mostly transferred to the water molecule in both adduct ions  $[\text{H}_2\text{O} \cdots \text{X}]^+$ . In the case of HT, the spin is entirely on OH while the charge, up to 2 Å,



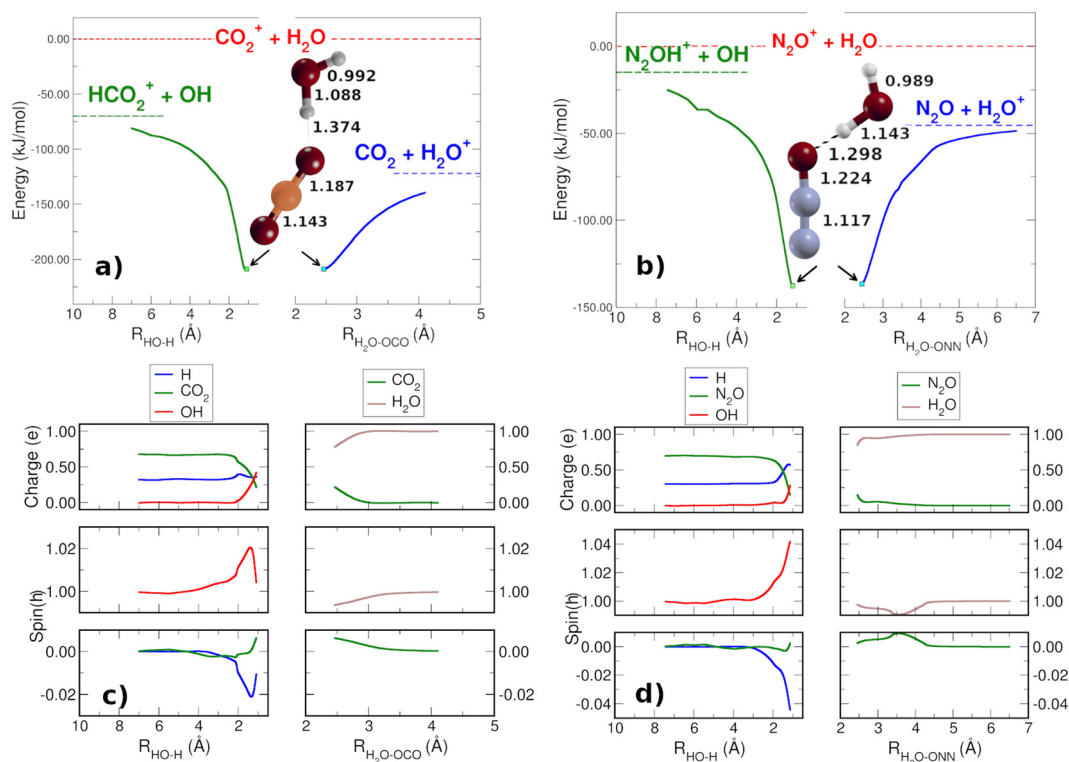
**Scheme 1** Reaction networks when  $\text{CO}_2^+$  and  $\text{N}_2\text{O}^+$  collide with  $\text{H}_2\text{O}$  at  $\text{CE} = 0$ .



**Table 1** Literature experimental<sup>29</sup> and theoretical (in bold)  $\Delta H_r^\circ$  (kJ mol<sup>-1</sup>) for the reactions in Scheme 1. Literature experimental rate coefficients  $k$  in cm<sup>3</sup> s<sup>-1</sup> molecule<sup>-1</sup> at 300 K for the four reactions in Scheme 1.  $k_{\text{HT}}$  and  $k_{\text{CT}}$  have an error of 15%.  $k_{\text{PT}}$  has an error of 20 and 25% in the reaction with CO<sub>2</sub> and N<sub>2</sub>O, respectively and  $k_{\text{WW}}$  of 10%

Molecule	$\Delta H_r^\circ$ (HT)	$\Delta H_r^\circ$ (CT)	$\Delta H_r^\circ$ (PT)	$\Delta H_r^\circ$ (WW)	$k_{\text{HT}}$	$k_{\text{CT}}$	$k_{\text{PT}}$	$k_{\text{WW}}$
CO <sub>2</sub>	-60.0 <b>-60.8<sup>a</sup></b>	-113.0 <b>-112.0<sup>a</sup></b>	-150.1	-97.3	$6.00 \times 10^{-10}$	$1.80 \times 10^{-9}$	$2.65 \times 10^{-9}$	$2.05 \times 10^{-9}$
N <sub>2</sub> O	-10.1 <b>-14.3<sup>a</sup></b>	-26.9 <b>-29.6<sup>a</sup></b>	-113.9	-97.3	$2.10 \times 10^{-10}$	$1.89 \times 10^{-9}$	$2.83 \times 10^{-9}$	$2.05 \times 10^{-9}$

<sup>a</sup> CCSD(T)/aug-cc-pVTZ energies with B2PLYP/aug-cc-pVTZ geometries and zero point correction.



**Fig. 2** MEPS, evaluated at the B2PLYP/aug-cc-pVTZ level of theory and with ZPE corrections, for the reactive dissociation of the molecular adducts [CO<sub>2</sub>-H<sub>2</sub>O]<sup>+</sup> (panel a) and [N<sub>2</sub>O-H<sub>2</sub>O]<sup>+</sup> (panel b). In red are the entrance channels of the reactions to which energies are referred. In green are indicated the energies of the HT products, while in blue are shown the CT energies. On the bottom panel c (CO<sub>2</sub> reaction) and d (N<sub>2</sub>O reaction) are reported the Mulliken population analysis relative to spin and charge during the reactive dissociation of the molecular adducts. See main text for further details.

is shared among X, the hydrogen involved in HT (blue H in Fig. 2(c) and (d)) and OH.

These pictures demonstrate that in HT the redistribution of the charge is a factor that, together with the lower reaction energy, probably makes HT slower than CT, where, in the adduct ions, both charge and spin are already in the final product H<sub>2</sub>O.

We also investigated the two reactions by changing the photon energies of the synchrotron radiation used to ionize the neutrals CO<sub>2</sub> and N<sub>2</sub>O. Carbon dioxide and nitrous oxide have different ionization energies being ( $X^2\Pi_g$ )  $13.777 \pm 0.001$  eV and ( $X^2\Pi$ )  $12.889 \pm 0.004$  eV, respectively.<sup>30</sup> Their PIEC, reported in the ESI<sup>†</sup> in Fig. S3 and S4, have been acquired together with the PIEC of NO<sup>+</sup> (Fig. S5, ESI<sup>†</sup>), the main fragment of nitrous oxide radical cation. In the case of CO<sub>2</sub>

the ionization efficiency is quite low at about 15.0 eV and no fragments have been observed up to this energy,<sup>21,31</sup> and it is known that the first excited state  $A^2\Pi_u$  is at 17.331 eV.<sup>32</sup> As regard N<sub>2</sub>O, the picture is more complex since several autoionizing resonances between the ground  $X^2\Pi$  and the first excited state  $A^2\Sigma^+$  (16.38 eV)<sup>33</sup> of N<sub>2</sub>O<sup>+</sup> start appearing at 13.9 eV. At about 14.77 eV the appearance of NO<sup>+</sup> is observed (Fig. S5, ESI<sup>†</sup>) due to the dissociation of excited N<sub>2</sub>O\* which autoionizes into N<sub>2</sub>O<sup>+</sup>( $X^2\Pi$ ) and predissociates.<sup>34</sup>

For the above discussed reasons the reaction involving CO<sub>2</sub><sup>+</sup> + H<sub>2</sub>O and N<sub>2</sub>O<sup>+</sup> + H<sub>2</sub>O have been studied in the photon energy range from 13.8 to 15.0 eV and 12.9 to 13.7 eV, respectively, to investigate the effect of the photon energy on the reactivity without the presence of fragments, resonances or electronic excited states.



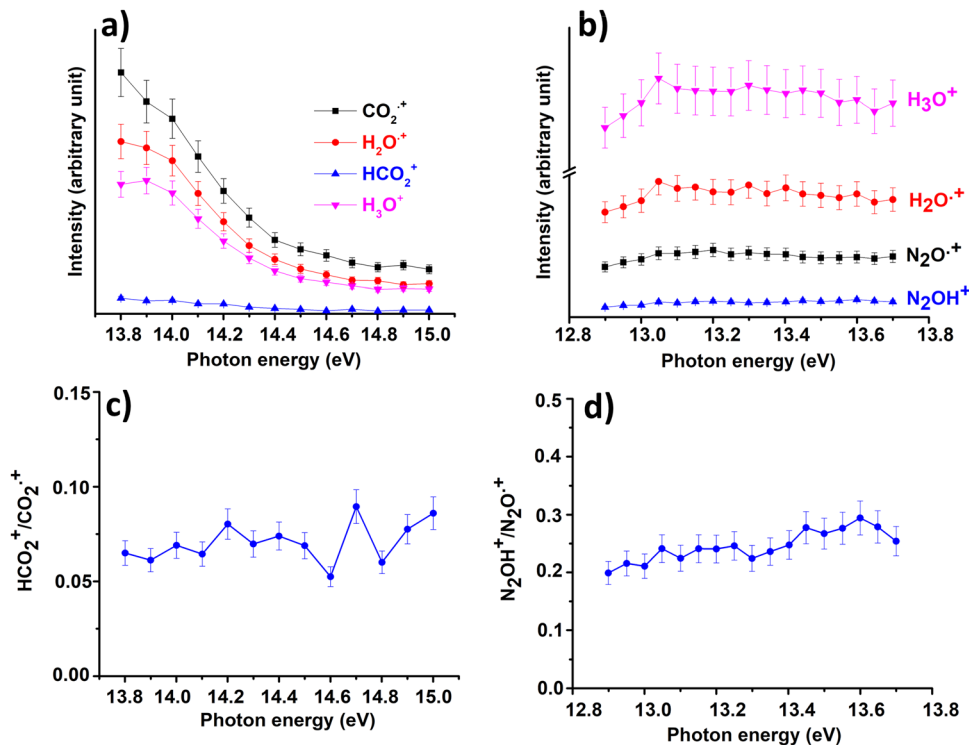


Fig. 3 Intensities of the ions involved in the reactions of Scheme 1 with  $X = \text{CO}_2$  (a) and  $X = \text{N}_2\text{O}$  (b) and the ratio  $\text{HCO}_2^+/\text{CO}_2^{*+}$  (c) and  $\text{N}_2\text{OH}^+/\text{N}_2\text{O}^{*+}$  (d) as a function of photon energy at the water pressure of  $4.3 \times 10^{-5}$  mbar and  $8.2 \times 10^{-5}$  mbar, respectively and  $\text{CE} = 0$ . Error bars are 10% of the values.

In Fig. 3(a) and (b) the intensities of all ions involved in the reactions of  $\text{CO}_2^+ + \text{H}_2\text{O}$  and  $\text{N}_2\text{O}^+ + \text{H}_2\text{O}$  together with the ratio  $\text{HCO}_2^+/\text{CO}_2^{*+}$  (Fig. 3(c)) and  $\text{N}_2\text{OH}^+/\text{N}_2\text{O}^{*+}$  (Fig. 3(d)) are reported as a function of photon energies and at  $\text{CE} = 0$ .

These data show that the slowest HT channel could be weakly influenced by the internal energy of the  $\text{CO}_2^+$  and  $\text{N}_2\text{O}^+$  ions. These measurements do not directly provide the experimental rate coefficients  $k$  of processes HT and CT as a function of  $h\nu$  as well as the internal energy of the ions acquired during the interaction with photons. Nevertheless, these important quantities can be derived by the experimental mass intensities recorded as a function of photoionization energy (Fig. 3(a) and (b)) together with the experimental photoionization spectra (Fig. S3 and S4 shown in ESI†). These data allow the development of an appropriate kinetic model, described in detail in the ESI† (Section S2).

The experimental rate coefficients  $k_{\text{HT}}$  and  $k_{\text{CT}}$  as a function of photon energy derived for  $\text{CO}_2^+ + \text{H}_2\text{O}$  and  $\text{N}_2\text{O}^+ + \text{H}_2\text{O}$  reactions are shown in Fig. 4 and 5, respectively.

In both HT reactions, the derived experimental rate coefficients seem to be not influenced by  $h\nu$  in the explored photon energy range. This is also true for the CT in  $\text{N}_2\text{O}^+ + \text{H}_2\text{O}$ , but the picture is more complex in the case of the reaction of water with carbon dioxide where a sinusoidal trend for  $k_{\text{CT}}$  vs.  $h\nu$  is derived. In this case the rate coefficient increases from the threshold up to 14.0 eV, then goes down up to 14.6 eV and finally grows again up to 15.0 eV. These evidences suggest that the well-known complex vibronic structure of the ionic ground  $\text{CO}_2^+(\text{X}^2\Pi_g)$  state in the energy range 13.8–14.7 eV could affect

the CT reactivity.<sup>35</sup> Probably the HT is also influenced by the manifold vibronic band structures of  $\text{CO}_2^+$ , but errors do not allow highlighting this effect. Nevertheless, although our measurements are not state-selective for a specific vibrational level of the reagent ions, in the case of  $\text{CO}_2^+$  is quite evident that vibrations in the ground state are involved in the photoionization process and affect its reactivity. Early studies,<sup>36–38</sup> have already shown that the vibrational excitation of  $\text{CO}_2^+$  has different effect, such as increasing the rate coefficient of CT with  $\text{O}_2$  and  $\text{NO}$ , while it decreases the rate coefficient in the HT with  $\text{H}_2$ .

The vibrational content of the products ( $E_{\text{vib}}$ ) at the different photon energies, together with their relative kinetic energy ( $E_{\text{tras}}$ ) have been also evaluated following the procedure described in Section 2.2.2, and discussed in the following next Sections 3.1 and 3.2.

### 3.1. HT and CT Products in the reactions of $\text{CO}_2^{*+}$ with $\text{H}_2\text{O}$

In Fig. 6(a) and (b) the energies content of the products in the HT (panel a) and CT (panel b) reactions of  $\text{CO}_2^+$  with  $\text{H}_2\text{O}$  is reported. In the HT reaction, the OH is always in its vibrational ground level, and all the vibrational energy is in the  $\text{HCO}_2^+$  ion. There is an almost constant trend in the energy partition between the vibrational and relative kinetic energies: the total energy ( $\sim 6000\text{--}7000\text{ cm}^{-1}$ , evaluated with eqn (6)) at each  $h\nu$  partially goes to the relative motion of the products ( $E_{\text{HT}}^{\text{tras}}$  between  $3100$  and  $4072\text{ cm}^{-1}$ ) whereas  $\text{HCO}_2^+$  is vibrationally excited between  $2382$  and  $3412\text{ cm}^{-1}$  (see Tables S3 and S4 in the ESI†).



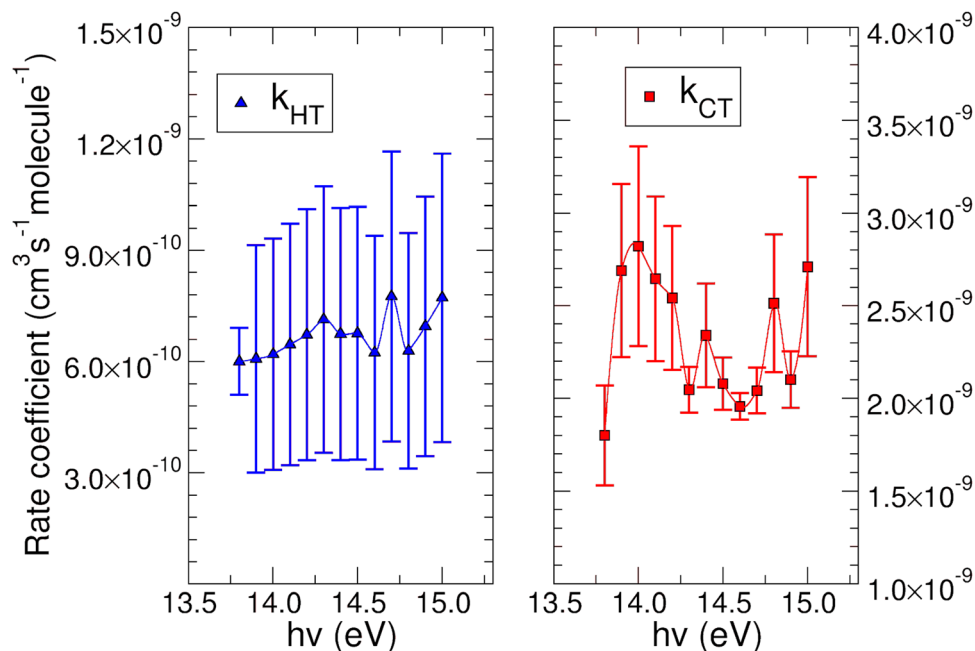


Fig. 4 Rate coefficient for the reaction of  $\text{CO}_2^+$  with water. On the left panel is reported the  $k$  for the HT channel as a function of the photon energy. On the right panel the rate coefficient for the CT process is shown. See main text for further details.

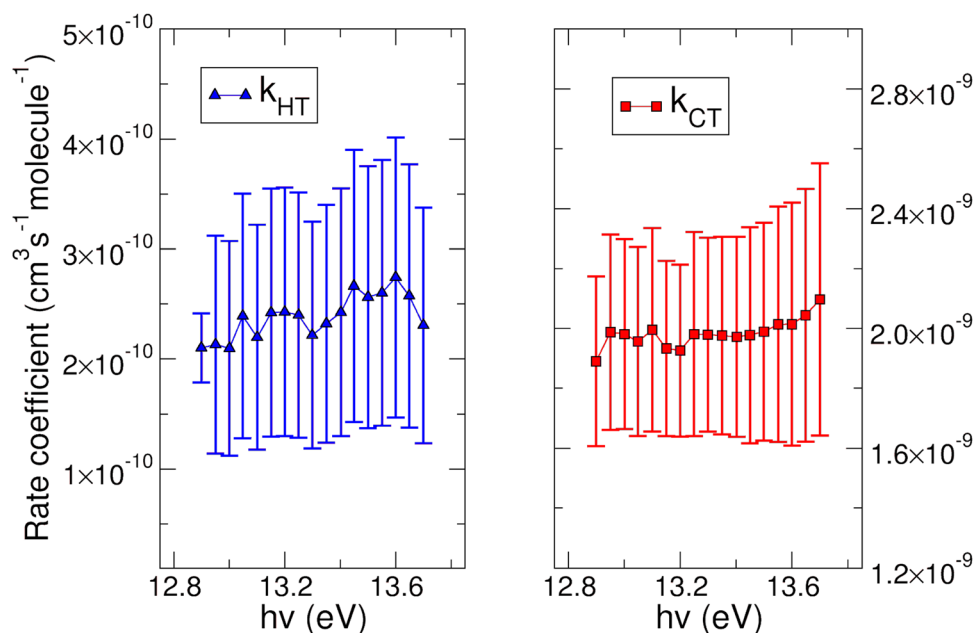


Fig. 5 Rate coefficient for the reaction of  $\text{N}_2\text{O}^+$  with water. On the left panel is reported the  $k$  for the HT channel as a function of the photon energy. On the right panel the rate coefficient for the CT process is shown. See main text for further details.

The CT channel exhibits quite a different behavior in the energy partition of the products: here the relative kinetic energies are much lower than in the HT channel, with values in between 884 and 2224  $\text{cm}^{-1}$ . The great part of the energy goes to the vibrational levels of both products  $\text{CO}_2$  and  $\text{H}_2\text{O}^+$ . Carbon dioxide is excited with vibrational energies between 4391 and 10132  $\text{cm}^{-1}$ , with oscillations as a function of  $h\nu$ , and

negatively correlated with the energy content of  $\text{H}_2\text{O}^+$ , which has vibrational levels excited between 0 and 4868  $\text{cm}^{-1}$  (see Tables S3 and S5 in the ESI<sup>†</sup>). This behavior resembles the oscillations observed in the rate coefficient as a function of  $h\nu$ . When the rate coefficient is higher, at  $h\nu$  around 14.0 and greater than 14.6 eV, the vibrational levels of the  $\text{H}_2\text{O}^+$  ion are excited. Overall the most substantial outcome of



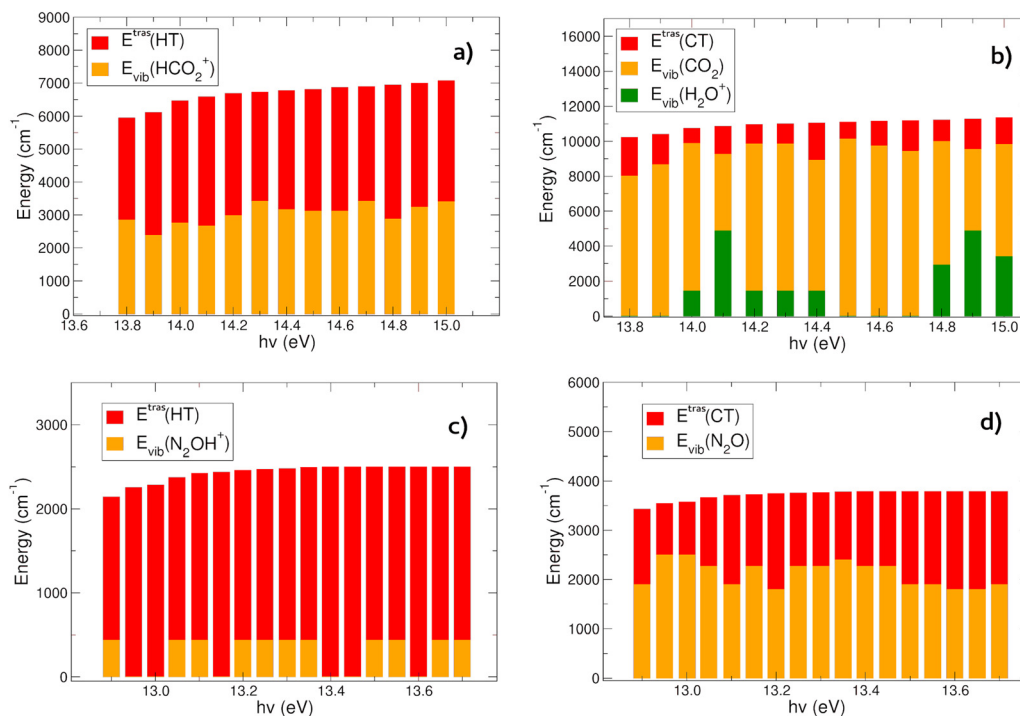


Fig. 6 Energy content of the products formed in the reaction involving  $\text{CO}_2^+ + \text{H}_2\text{O}$  in HT (a) and CT channels (b). Energy content of the products formed in the reaction involving  $\text{N}_2\text{O}^+ + \text{H}_2\text{O}$  in HT (c) and CT channels (d).

the  $\text{CO}_2^+ + \text{H}_2\text{O}$  reaction is the vibrational heating of the neutral carbon dioxide, and to a lesser extent the vibrational heating of the water cation product.

### 3.2. HT and CT in the reactions of $\text{N}_2\text{O}^+$ with $\text{H}_2\text{O}$

The CT and HT reactions of water with  $\text{N}_2\text{O}^+$  as a function of the photoionization energies reveal a different scenario in the energy partition of the products at variance with the  $\text{CO}_2$  reactions. In Fig. 6(c) and (d) the energy distributions of the products in the two channels HT and CT, respectively, are reported. The data reveal that in this case the kinetic energy of the relative motion of the products plays a major role in the HT reaction (Fig. 6(c)). Here the  $\text{N}_2\text{OH}^+$  is vibrationally excited only in its NNO in plane bending mode (see Tables S3 and S6, ESI<sup>†</sup>), while all the other frequency modes are spectator during the reactions. The relative kinetic energies here are within a range of 1706–2501  $\text{cm}^{-1}$ . The CT channel shows a very different trend in water cation vibrational excitation compared to the  $\text{CO}_2$  reaction. Here there is no involvement of the frequency modes of  $\text{H}_2\text{O}^+$  during the charge transfer process. The neutral  $\text{N}_2\text{O}$  molecule is vibrationally excited in a range between 1797 and 2495  $\text{cm}^{-1}$ , whereas the relative kinetic energies vary between 1048 and 1990  $\text{cm}^{-1}$  (see Table S7, ESI<sup>†</sup>).

In conclusion, the products energy distribution analysis shows that in the CT reactions of  $\text{CO}_2^+$  and  $\text{N}_2\text{O}^+$  with water, vibrationally hot neutrals  $\text{CO}_2$  and  $\text{N}_2\text{O}$  products are formed. Considering the relevance of these molecules in the atmospheric environments, hereafter their relaxation mechanisms will be discussed and the radiative and collision quenching will be considered.

The calculated vibrational energies of  $\text{CO}_2$  and  $\text{N}_2\text{O}$  (Tables S5 and S7, ESI<sup>†</sup>) at the photon energies investigated have been used to evaluate their vibrational radiative power and lifetime by using the high-temperature molecular spectroscopic database HITEMP (Section S3 in the ESI<sup>†</sup>).<sup>39</sup> The vibrational radiative emission has been found to have radiation lifetimes higher than one second (see Tables S8 and S9 of ESI<sup>†</sup>). On the other hand, the vibrational-translational relaxation due to collisions with surrounding atmospheric molecules has lifetimes of the order  $\mu\text{s}^{40-42}$  at the upper tropospheric/stratospheric pressures of about 0.1 atm.<sup>43</sup> Hence, the dominant quenching mechanism in this atmospheric environment is not radiative emission, but the collision of hot  $\text{CO}_2$  and  $\text{N}_2\text{O}$  with the surrounding molecules. This relaxation mechanism is essentially similar to the energy transfer which  $\text{CO}_2$  undergoes after earth's IR excitation, and it points to the role of cosmic rays as not only effective in global climate cooling by triggering the formation of clouds, but also as a warming agent.<sup>6</sup> On the other hand, the vibrational radiative emission mechanism becomes competitive with the collisional relaxation process only at lower pressures such as those occurring in the outer space such as exosphere and the interstellar medium ( $P \leq 10^{-7}$  atm).<sup>44</sup>

The results of this work are important for the role, sometime controversial and not yet well understood, of all these molecules on climate models, as for instance when they are injected into the stratosphere by volcanoes. Recently, the submarine volcano Hunga Tonga–Hunga Ha'apai eruption has increased the amount of global stratospheric water vapor by ca. 10%.<sup>11</sup> During the volcanic eruption a large amount of gases are

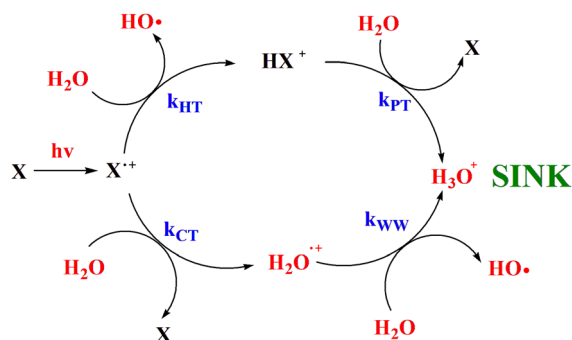


introduced into the atmosphere such as  $\text{CO}_2$ ,  $\text{CH}_4$ ,  $\text{H}_2\text{O}$ ,  $\text{SO}_2$  with a potential climate impact. Indeed, these molecules, which are greenhouse gases, can alter, among the others, the ozone layer, the stratospheric sulfate aerosol dynamic, cloud formation and in general the atmospheric chemistry.<sup>45–47</sup> More specifically the  $\text{H}_3\text{O}^+$  and  $\text{OH}$  species studied in the chemical network of the present work, have a relevant involvement in the atmospheric chemistry. Protonated mixed water organic clusters are among the most abundant cations in the stratosphere,<sup>48,49</sup> and alter the Earth's radiative balance due to IR absorption.<sup>8</sup> Hydroxyl radical  $\text{OH}$  is the key oxidant in the atmosphere<sup>50</sup> and affect the oxidation of methane, the budget of ozone and the sulfate aerosol formation by reacting with  $\text{SO}_2$  thus with a potential climate effect.<sup>51,52</sup>  $\text{OH}$ , primary formed by the reaction of  $\text{O}(^1\text{D})$  with  $\text{H}_2\text{O}$  triggered by sunlight dissociation of  $\text{O}_3$  in  $\text{O}(^1\text{D})$  and  $\text{O}_2$ , can also be formed by secondary processes. These processes are operative also during nighttime started either by neutral chemicals, mainly in the troposphere,<sup>50</sup> or by ions, such as  $\text{SO}_2^+$ ,  $\text{CO}_2^+$  and  $\text{N}_2\text{O}^+$  produced by cosmic rays. These ions can increase the rate coefficients of the reactions of several orders of magnitude<sup>53</sup> with respect to the relative neutral reaction, resulting in a fast formation of reactive  $\text{OH}$  radicals and “hot”  $\text{CO}_2$  and  $\text{N}_2\text{O}$  that can relax losing their vibrational energy by collisions.

## 4. Conclusions

In this work the mechanistic insights of the reactions involving  $\text{CO}_2^+ + \text{H}_2\text{O}$  and  $\text{N}_2\text{O}^+ + \text{H}_2\text{O}$  are reported and the rate coefficients as a function of photon energy are given. The results show that a chemical network is operative as shown in Scheme 2.

Overall these processes do not destroy the molecules  $\text{CO}_2$  and  $\text{N}_2\text{O}$ , rather, their ionization by cosmic rays' catalyze the transformation of  $\text{H}_2\text{O}$  in  $\text{H}_3\text{O}^+$  and  $\text{OH}$  as also observed with the molecule  $\text{SO}_2$  investigated in our previous works.<sup>15,16</sup> Moreover, theoretical calculations demonstrates that  $\text{CO}_2$  and  $\text{N}_2\text{O}$  molecules are reformed in excited vibrational levels that in microseconds time scale quench *via* collision with the



$\text{X} = \text{CO}_2, \text{N}_2\text{O}, \text{SO}_2$

**Scheme 2** Catalytic scheme proposed for the reaction of  $\text{CO}_2$ ,  $\text{N}_2\text{O}$  and  $\text{SO}_2$  with  $\text{H}_2\text{O}$  in the gas phase under ionization conditions.

surrounding molecules at the typical pressures and temperatures of the upper troposphere/stratosphere and warming the surrounding atmosphere. In conclusion this work highlights the relevance of ion chemistry of trace gases and of laboratories studies that can improve the knowledge of the role of ions to climate and global warming, a challenging scientific topic of our days.<sup>54,55</sup>

## Author contributions

Daniele Catone and Mattea Carmen Castrovilli: performed experiments and data analysis, review & editing. Francesca Nicolanti: review & editing. Mauro Satta: conceptualization, performed theoretical calculations, writing – original draft, review & editing. Antonella Cartoni: supervision, conceptualization, performed experiments and data curation, writing – original draft, review & editing.

## Conflicts of interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

## Acknowledgements

We thank Fabio Zuccaro for assistance during the experiments on the beamline CiPo of Elettra Synchrotron. We acknowledge Elettra Synchrotron Trieste for providing access to its synchrotron radiation facilities and for financial support. This article is based upon work from COST action CA18212 – Molecular Dynamics in the GAS phase (MD-GAS).

## References

- 1 D. Siingh and R. P. Singh, *J. Phys.*, 2010, **74**, 153–168.
- 2 D. Catone, M. Satta, M. C. Castrovilli, L. Avaldi and A. Cartoni, *Chem. Phys. Lett.*, 2021, **771**, 138467.
- 3 I. A. Mironova, K. L. Aplin, F. Arnold, G. A. Bazilevskaya, R. G. Harrison, A. A. Krivolutsky, K. A. Nicoll, E. V. Rozanov, E. Turunen and I. G. Usoskin, *Space Sci. Rev.*, 2015, **194**, 1–96.
- 4 H. Svensmark, M. B. Enghoff, N. J. Shaviv and J. Svensmark, *Nat. Commun.*, 2017, **8**, 1–9.
- 5 E. M. Dunne, H. Gordon, A. Kürten, J. Almeida, J. Duplissy, C. Williamson, I. K. Ortega, K. J. Pringle, A. Adamov and U. Baltenspergere, *et al.*, *Science*, 2016, **354**, 1119–1124.
- 6 V. Kumar, S. K. Dhaka, M. H. Hitchman and S. Yoden, *Sci. Rep.*, 2023, **13**, 3707.
- 7 S. Solomon, *Rev. Geophys.*, 1999, **37**, 275–316.
- 8 K. L. Aplin and R. A. McPheat, *J. Atmos. Sol.-Terr. Phys.*, 2005, **67**, 775–783.
- 9 T. Leblanc, T. D. Walsh, I. S. McDerimid, G. C. Toon, J.-F. Blavier, B. Haines, W. G. Read, B. Herman, E. Fetzer and S. Sander, *et al.*, *Atmos. Meas. Tech.*, 2011, **4**, 2579–2605.



- 10 H. Vömel, S. Evan and M. Tully, *Science*, 2022, **377**, 1444–1447.
- 11 S. Khaykin, A. Podglajen, F. Ploeger, J.-U. Groöf, F. Tence, S. Bekki, K. Khlopenkov, K. Bedka, L. Rieger and A. Baron, *et al.*, *Commun. Earth Environ.*, 2022, **3**, 316.
- 12 S. Jenkins, C. Smith, M. Allen and R. Grainger, *Nat. Clim. Change*, 2023, **13**, 127–129.
- 13 P. Sellitto, A. Podglajen, R. Belhadji, M. Boichu, E. Carboni, J. Cuesta, C. Duchamp, C. Kloss, R. Siddans and N. Bègue, *et al.*, *Commun. Earth Environ.*, 2022, **3**, 288.
- 14 G. de Petris, A. Cartoni, A. Troiani, G. Angelini and O. Ursini, *Phys. Chem. Chem. Phys.*, 2009, **11**, 9976.
- 15 A. Cartoni, D. Catone, P. Bolognesi, M. Satta, P. Markus and L. Avaldi, *Chem. – Eur. J.*, 2017, **23**, 6772–6780.
- 16 M. Satta, A. Cartoni, D. Catone, M. C. Castrovilli, P. Bolognesi, N. Zema and L. Avaldi, *ChemPhysChem*, 2020, **21**, 1146–1156.
- 17 M. Satta, P. Bolognesi, A. Cartoni, A. R. Casavola, D. Catone, P. Markus and L. Avaldi, *J. Chem. Phys.*, 2015, **143**, 244312.
- 18 A. Cartoni, A. R. Casavola, P. Bolognesi, S. Borocci and L. Avaldi, *J. Phys. Chem. A*, 2015, **119**, 3704.
- 19 A. Cartoni, P. Bolognesi, E. Fainelli and L. Avaldi, *J. Chem. Phys.*, 2014, **140**, 184307.
- 20 G. V. Marr and J. B. West, *At. Data Nucl. Data Tables*, 1976, **18**, 497–508.
- 21 V. G. Anicich, *An Index of the Literature for Bimolecular Gas Phase Cation-Molecule Reaction Kinetics*, JPL Publication, 2003, pp.3–19.
- 22 T. Su and W. J. Chesnavich, *J. Chem. Phys.*, 1982, **76**, 5183–5185.
- 23 S. Grimme, *J. Chem. Phys.*, 2006, **124**, 034108.
- 24 R. A. Kendall, T. H. Dunning and R. Harrison, *J. Chem. Phys.*, 1992, **96**, 6796–6806.
- 25 M. J. Frisch, G. W. Trucks, H. B. Schlegel, G. E. Scuseria, M. A. Robb, J. R. Cheeseman, G. Scalmani, V. Barone, G. A. Petersson, H. Nakatsuji, X. Li, M. Caricato, A. Marenich, J. Bloino, B. G. Janesko, R. Gomperts, B. Mennucci, H. P. Hratchian, J. V. Ortiz, A. F. Izmaylov, J. L. Sonnenberg, D. Williams-Young, F. Ding, F. Lipparini, F. Egidi, J. Goings, B. Peng, A. Petrone, T. Henderson, D. Ranasinghe, V. G. Zakrzewski, J. Gao, N. Rega, G. Zheng, W. Liang, M. Hada, M. Ehara, K. Toyota, R. Fukuda, J. Hasegawa, M. Ishida, T. Nakajima, Y. Honda, O. Kitao, H. Nakai, T. Vreven, K. Throssell, J. A. Montgomery, Jr., J. E. Peralta, F. Ogliaro, M. Bearpark, J. J. Heyd, E. Brothers, K. N. Kudin, V. N. Staroverov, T. Keith, R. Kobayashi, J. Normand, K. Raghavachari, A. Rendell, J. C. Burant, S. S. Iyengar, J. Tomasi, M. Cossi, J. M. Millam, M. Klene, C. Adamo, R. Cammi, J. W. Ochterski, R. L. Martin, K. Morokuma, O. Farkas, J. B. Foresman and D. J. Fox, *Gaussian 09, Revision A.02, ct technical report*, Gaussian, Inc., Wallingford, 2016.
- 26 R. S. Mulliken, *J. Chem. Phys.*, 1955, **23**, 1833–1840.
- 27 J. L. Bao and D. G. Truhlar, *Chem. Soc. Rev.*, 2017, **46**, 7548–7596.
- 28 M. Satta, D. Catone, M. C. Castrovilli, P. Bolognesi, L. Avaldi, N. Zema and A. Cartoni, *J. Phys. Chem. A*, 2022, **126**(22), 3463–3471.
- 29 B. Ruscic and D. H. Bross, *Active Thermochemical Tables (ATcT) Values Based on ver. 1.122p of the Thermochemical Network*, 2020, <https://atct.anl.gov/> (accessed 2023-02-01).
- 30 *NIST Chemistry WebBook*, ed. P. J. Linstrom and W. G. Mallard, NIST Standard Reference Database number 69; National Institute of Standards and Technology, Gaithersburg, MD, 2005, <http://webbook.nist.gov> (accessed January 2023).
- 31 J. H. D. Eland, *Int. J. Spectrom. Ion Phys.*, 1972, **9**, 397–406.
- 32 M. Baer and P. M. Guyon, *J. Chem. Phys.*, 1986, **85**, 4765.
- 33 E. Kinmonda, J. H. D. Eland and L. Karlsson, *Int. J. Mass Spectrom.*, 1999, **185–187**, 437–447.
- 34 J. L. Olivier, R. Locht and J. Momigny, *Chem. Phys.*, 1982, **68**, 201–211.
- 35 J. Liu, W. Chen, C.-W. Hsu, M. Hochlaf, M. Evans, S. Stimson and C. Y. Ng, *J. Chem. Phys.*, 2000, **112**(24), 22.
- 36 C. Nicolas, C. Alcaraz, R. Thissen, J. Žabka and O. Dutuit, *Planet. Space Sci.*, 2002, **50**, 877–887.
- 37 E. Alge, H. Villinger and W. Lindinger, *Plasma Chem. Plasma Process.*, 1981, **1**, 65–71.
- 38 G. Borodi, A. Luca and D. Gerlich, *Int. J. Mass Spectrom.*, 2009, **280**, 218–225.
- 39 L. S. Rothman, I. E. Gordon, R. J. Barber, H. Dothe, R. R. Gamache, A. Goldman, V. Perevalov, S. A. Tashkun and J. Tennyson, *J. Quant. Spectrosc. Radiat. Transfer*, 2010, **111**, 2139–2150.
- 40 Z. Baalbaki and H. Teitelbaum, *Chem. Phys.*, 1986, **104**, 83–106.
- 41 F. Cannemeyer and A. E. De Vries, *Physica*, 1974, **74**, 196–204.
- 42 R. M. Siddles, G. J. Wilson and C. J. S. M. Simpson, *Chem. Phys.*, 1994, **189**, 779–791.
- 43 M. Pidwirny, Atmospheric Pressure, *Fundamentals of Physical Geography*, 2nd edn, 2006.
- 44 K. M. Ferriere, *Rev. Mod. Phys.*, 2001, **73**, 1031.
- 45 J. G. Anderson, D. M. Wilmouth, J. B. Smith and D. S. Sayres, *Science*, 2012, **337**, 835–839.
- 46 M. M. Joshi and G. S. Jones, *Atmos. Chem. Phys.*, 2009, **9**, 6109–6118.
- 47 L. Millán, M. L. Santee, A. Lambert, N. J. Livesey, F. Werner, M. J. Schwartz, H. C. Pumphrey, G. L. Manney, Y. Wang, H. Su, L. Wu, W. G. Read and L. Froidevaux, *Geophys. Res. Lett.*, 2022, **49**, 1–10.
- 48 G. Henschen and F. Arnold, *Nature*, 1981, **291**, 211–213.
- 49 N. S. Shuman, D. E. Hunton and A. A. Viggiano, *Chem. Rev.*, 2015, **115**, 4542–4570.
- 50 J. Lelieveld, S. Gromov, A. Pozzer and D. Taraborrelli, *Atmos. Chem. Phys.*, 2016, **16**, 12477–12493.
- 51 V. Vaida, *J. Chem. Phys.*, 2011, **135**, 020901.
- 52 M. Li, E. Karu, C. Brenninkmeijer, H. Fischer, J. Lelieveld and J. Williams, *npj Clim. Atmos. Sci.*, 2018, **29**, 1–7.
- 53 J. H. Futrell and T. O. Tiernan, *Science*, 1968, **162**, 415–422.
- 54 J. B. Burkholder, J. P. D. Abbatt, I. Barnes, J. M. Roberts, M. L. Melamed, M. Ammann, A. K. Bertram, C. D. Cappa, A. G. Carlton and L. J. Carpenter, *et al.*, *Environ. Sci. Technol.*, 2017, **51**, 2519–2528.
- 55 M. Satta, M. C. Castrovilli, F. Nicolanti, A. R. Casavola, C. Mancini Terracciano and A. Cartoni, *Condens. Matter*, 2022, **46**, 1–30.

