



Cite this: *Phys. Chem. Chem. Phys.*,
2023, 25, 17869

Nanostructured supramolecular networks from self-assembled diamondoid molecules under ultracold conditions†

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Diamondoid molecules and their derivatives have attracted attention as fascinating building blocks for advanced functional materials. Depending on the balance between hydrogen bonds and London dispersion interactions, they can self-organize in different cluster structures with functional groups tailored for various applications. Here, we present a new approach to supramolecular aggregation where self-assembly of diamondoid acids and alcohols in the ultracold environment of superfluid helium nanodroplets (HNDs) was analyzed by a combination of time-of-flight mass spectrometry and computational tools. Experimentally observed magic numbers of the assembled cluster sizes were successfully identified and computed cluster structures gave valuable insights into a different conglomeration mode when compared to previously explored less-polar diamondoid derivatives. We have confirmed that functional groups acting as good hydrogen bond donors completely take over the self-organization process, resulting in fascinating pair-wise or cyclic supramolecular assemblies. Particularly noteworthy is that mono- and bis-substituted diamondoid derivatives of both series engage in completely different modes of action, which is reflected in differing non-covalent cluster geometries. Additionally, formed cyclic clusters with a polar cavity in the center and a non-polar diamondoid outer layer can be of high interest in porous material design and provide insights into the structural requirements needed to produce bulk materials with desired properties.

Received 23rd May 2023,
Accepted 13th June 2023

DOI: 10.1039/d3cp02367b

rsc.li/pccp

Introduction

Diamondoids,^{1,2} cage-like hydrocarbons that are nanometer sized and naturally occurring,³ are very rich in C–H bonds and can act as good London dispersion (LD)⁴ energy donors. They share with their diamond counterpart some excellent properties like thermal stability and rigidity. Through selective chemical functionalization their derivatives may serve for various applications in materials science and biomedical areas.⁵ Typically, diamondoid building blocks are tailored using the means of organic synthesis and then individual scaffolds are left to form, in a best case scenario, superstructures of desired properties. However, to truly achieve predictable and controllable self-assembly, one needs to understand the underlying

intermolecular binding modes. Our goal was to study fundamental aspects of the binding processes in self-organized nanosized networks of diamondoid derivatives under largely unperturbed conditions. Especially, the investigation of the interplay between hydrogen bonding and dispersion attraction effects requires the use of the least interacting non-polar environment possible. To this end we turned to helium as a solvent.

The highly non-polar medium of helium nanodroplets (HNDs) offers an opportunity to non-disruptively study weak intermolecular interactions.^{6–14} By expanding gaseous helium in a high vacuum at cryogenic temperatures,^{15,16} one obtains helium aggregates suitable for trapping introduced atoms and molecules. Cluster formation is a favorable process since the binding energy of aggregation can be dissipated through the evaporation of helium atoms. Some recent examples of studied molecular clusters in HNDs include (V₂O₅)_n,¹⁷ weakly bound alkali triplet-dimers¹⁸ and quartet-trimers,¹⁹ diamondoid molecules like adamantane^{20–22} and diamantane,²³ as well as their ether derivatives,²⁴ etc.

On the other hand, introduction of hydrogen bonding in a HND cluster system is expected to provide means for spatial

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† Electronic supplementary information (ESI) available: Figures and tables with experimental and computational details. See DOI: <https://doi.org/10.1039/d3cp02367b>

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control of cluster self-organization and thereby offer a way to design supramolecular assemblies with the exact and repeatable arrangement of its constituent molecules. Hydrogen bonds are present in practically all biological and chemical processes and are essential for the chemistry of life, from the structural and functional features of a simple water molecule to the behavior of complex molecules such as proteins or DNA. According to the latest IUPAC recommendation, the hydrogen bond is considered to be an attractive interaction between a hydrogen atom from a molecule or a molecular fragment X–H in which X is more electronegative than H, and an atom or a group of atoms in the same or a different molecule, in which there is evidence of bond formation.²⁵ Not surprisingly, hydrogen bonding is therefore an underlying phenomenon found in diverse areas spanning from supramolecular assemblies, organic synthesis, polymerization, medicinal chemistry, *etc.*, to name just a few. As hydrogen bonds are directional, they are often used as a strategic tool in the synthesis of novel supramolecular architectures with a goal to optimize molecular recognition processes that mimic nature's elegant designs.^{26,27} An important factor for our current considerations is that hydrogen bonding is especially favored in the non-polar environment, such as organic solvents.²⁸

We previously investigated the influence of dispersion interactions on the self-assembly of less-polar diamondoids in HNDs.^{23,24} Going one step further, in the scope of this study we have now studied cluster formation of diamondoid acids **1a–4a** and alcohols **1b–4b** (Fig. 1) under HND conditions. The experimental findings were supported by a computational analysis, with special emphasis on the identified cluster sizes with large abundances, *i.e.*, the magic numbers. Our computational approach was described previously^{23,24} and applied here for the structural analysis of the clusters. To the best of our knowledge, the aggregation behavior of diamondoids in HNDs has only been studied for adamantane, diamantane, and diamondoid ethers; other diamondoids or their derivatives have remained unexplored up to now. For diamondoid acids and alcohols, molecules that possess functional groups capable of efficient hydrogen bond formation, it was expected that network assembly and stability would indeed be predominately governed by hydrogen bonding interactions. Our current work therefore had the goal to assess a possible interplay of hydrogen bonding with dispersion interactions in diamondoid acid and alcohol assemblies which were formed by clean molecule-by-molecule aggregation at low temperature inside a superfluid helium environment.

Experimental details

Synthesis

1-Adamantanecarboxylic acid (**1a**) and 1-hydroxyadamantane (**1b**) were purchased from Sigma-Aldrich (99% purity) and were used as received. 1-Diamantanecarboxylic acid (**2a**),²⁹ 4-diamantanecarboxylic acid (**3a**),³⁰ and 4,9-diamantanedicarboxylic acid (**4a**)²⁹ as well as alcohols³¹ 1-hydroxydiamantane (**2b**), 4-hydroxydiamantane (**3b**) and 4,9-dihydroxydiamantane (**4b**) were synthesized according to previously published procedures.

Helium nanodroplets

The apparatus used for the generation of the HNDs is described in detail in ref. 12,17. In short, pressurized high purity He (99.9999%) is cooled to temperatures below 20 K using a closed-cycle refrigerator (Sumitomo RDK-408D2) and expanded through a 5 μm nozzle in a high vacuum. During this process, the gaseous He condenses into small superfluid droplets. Under the expansion conditions used in the experiments ($p_{\text{He}} = 60$ bars, $T_{\text{He}} = 11.5\text{--}12.5$ K), He droplets with a mean diameter of 40 to 60 nm, consisting of about 1×10^6 to 3×10^6 He atoms, are formed.^{15,16} Subsequently, the beam passes a skimmer and the helium droplets pick up the desired dopant species in a separately pumped chamber. Here, we dope the droplets with diamondoid derivatives using a heated gas-pickup cell (96 °C for **1a**, 200 °C for **2a**, 180 °C for **3a**, 45 °C for **1b**, 120 °C for **2b**, and 114 °C for **3b**), which is connected to a heated reservoir (70 °C) *via* a precision leak valve or using a heated crucible directly in the pickup chamber (320 °C for **4a** and 145 °C for **4b**). The doped He droplets then enter the differentially pumped analysis chamber, where a reflectron time-of-flight mass spectrometer KAESDORF RTF50 is utilized to record the mass spectra. Upon electron impact ionization, cluster ions are expelled from the He droplets and can be detected at the corresponding mass channel. The employed emission current (I_{em}) and ionization energy (E_{el}) are $I_{\text{em}} = 6.8$ μA and $E_{\text{el}} = 90$ eV, respectively.

Theoretical methods

Geometry optimizations of acids **1a**, **2a** and **4a** and alcohols **1b–4b** were performed using the Orca 5.0.3 program package^{32,33} using the B3LYP-gCP-D3(BJ)-ABC/def2-TZVPP level of theory,^{34,35} and the obtained minima were verified by frequency computations. A search for favorable cluster structures was done using the Conformer-Rotamer Ensemble

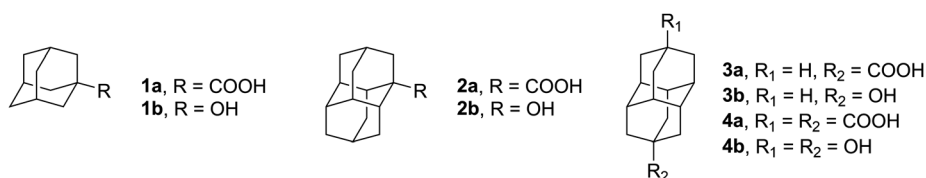


Fig. 1 Structures of diamondoid acids **1a–4a** and alcohols **1b–4b** used in HND studies.



Sampling Tool (CREST) based on the GFN methods^{36,37} by applying the iterative *meta*-dynamic sampling for non-covalently bound complexes, clusters or aggregates (NCI-iMTD mode). Single point computations on the identified best cluster structures were performed using the B3LYP-gCP-D3(BJ)-ABC/def2-TZVPP level of theory. The choice of the DFT method was based on our previous benchmarking of similarly bulky molecules in the HNDs.^{23,24} We found that using the D3(BJ) dispersion correction,^{34,35,38} the three-body dispersion contribution term implemented in Orca as well as the geometrical counterpoise (gCP) correction³⁹ successfully accounts for subtle intermolecular interactions and mitigates the basis-set superposition errors (BSSEs), respectively. Please see the ESI† for more computational details.

Results and discussion

In previous studies of self-organized HND aggregates consisting of diamondoid molecules such as adamantane,^{20–22} diamantane²³ and diamondoid ethers,²⁴ certain preferences for cluster abundances were found that belonged to particularly stable geometries. These observed privileged abundances were termed as magic numbers. Additionally, the intermolecular dispersion attraction between bulky hydrocarbons had a pronounced effect on supramolecular binding and cluster formation in the HNDs. We were interested in exploring how the introduction of hydrogen bond donating groups affects the aggregation of diamondoid molecules in a HND medium, unimpeded by the effect of common solvents. When the organic molecule is no longer a pure hydrocarbon but still of low polarity, like in the case of diamondoid ethers, dispersion attraction still acts as a driving force for cluster formation but hydrogen bonding between ether oxygens and trace water molecules leads to an eventual breakdown of the initial supramolecular aggregates. With this in mind, we went further up on the relative polarity scale and prepared diamondoid acids **1a–4a** and alcohols **1b–4b** (Fig. 1). Acid **1a** and alcohol **1b** consist of an adamantane scaffold, while acids **2a–4a** and alcohols **2b–4b** incorporate the diamantane cage carrier. In continuation, we will focus on each series of the prepared diamondoid derivatives.

Diamondoid acids

To date, investigations of organic acids in the helium environment were mostly limited to small molecules like formic and acetic acids. The HND studies of formic acid, the simplest of organic carboxylic acids, included characterization studies of monomers and clusters of both the parent molecule⁴⁰ and its deuterated form.⁴¹ Special attention was paid to its dimer which is the smallest organic prototype for a double-hydrogen-bonded system. Surprisingly, the preferred formation of an acyclic polar dimer in HNDs,^{42,43} a structure that is not a global energetic minimum, demonstrated yet again the uniqueness of this medium for the exploration of non-covalently bound clusters. Namely, instead of the double-hydrogen-

bonded acid dimer structure that has two strong and equivalent O–H···O=C interactions and is typically found in the gas phase, the observed structure of formic acid dimers in the HNDs adopted a molecular orientation where one strong O–H···O=C and one relatively weak C=O···H–C interaction co-existed. A similar observation was also found for acetic acid dimers,^{44–46} where the preferential formation of metastable acyclic pair-wise assemblies in the HNDs was in contrast to the cyclic species usually dominant in the gas phase. It was suggested that acyclic acid clusters with their asymmetric charge distributions interact more readily with He and thereby form stable clusters with helium atoms attached.⁴⁷ Such an influence of helium atoms on metastable dimer abundances was also recently demonstrated using molecular dynamics computations.⁴⁷ As the acid monomer is embedded into a helium medium, the collision with the second acid molecule results in the release of excess energy through helium atom evaporation and the formed dimer is prone to keep the initial relative orientations of the acid molecules, unlike in the gas phase where the collision energy is not dispersed in such a way and the distribution of internal energies can be very broad.

In the case of diamondoid acids, aside from the dimer structures, we also observed the formation of larger magic number clusters. Fig. 2(a) shows the course of the mass peak height with respect to the number of acid molecules per cluster for the investigated compounds. In panel Fig. 2(b) the integrated peak area is plotted with respect to the number of acid molecules per cluster. Both indicate that a cluster of five molecules for adamantane acid **1a** and diamantane medial acid **2a** is particularly stable as well as a cluster of four molecules for diamantane bis-apical diacid **4a**. This is also confirmed by the spectra shown in Fig. S1 (ESI†). There seems to be no preferable cluster size for diamantane mono-apical acid **3a**, for which the signal drop-off is also the strongest with an increasing cluster size, suggesting that clusters of **3a** are the weakest bound.

Moreover, for every cluster size, there is a series of additional peaks at larger mass numbers. These can be assigned to a specific number of water molecules added to the cluster. In Fig. S1 (ESI†), the peak areas with respect to the number of additional water molecules are indicated, with each line scaled by its highest value. It is striking that for small clusters of **1a** and **2a** no additional water molecule seems to be favorable and for larger clusters this shifts to one additional water molecule. For **1a** this transition is located between the cluster sizes eight and nine and for **2a** it is located between the cluster sizes seven and eight. This is probably due to the stabilization of the larger cluster by the water molecule. For **4a** something similar is observable for cluster sizes between six and seven; however, here the preferred number of water molecules changes from zero to two.

The computational analysis of the experimentally most abundant clusters formed by studied diamondoid acids in the HNDs, **1aCL5** ($n = 5$), **2aCL5** ($n = 5$) and **4aCLA** ($n = 4$), revealed informative details about structural organization (Fig. S3, ESI†). Building on our previous benchmarking for



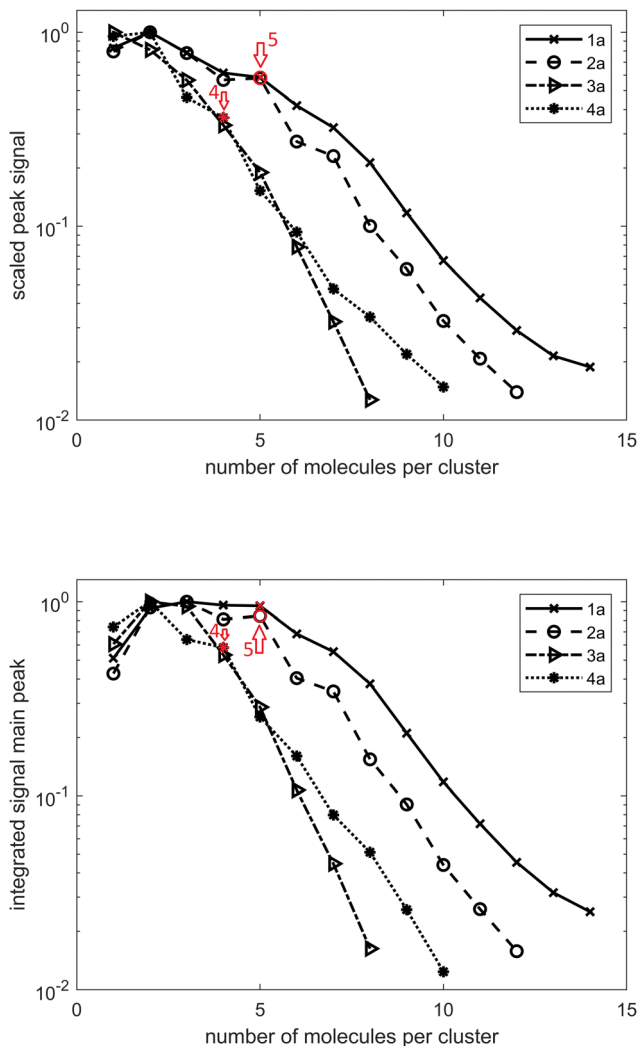


Fig. 2 TOF signal curves for adamantane acid **1a**, diamantane medial acid **2a**, diamantane mono-apical acid **3a**, and diamantane bis-apical diacid **4a**. (a) Peak height scaled by the strongest signal of each compound. (b) Integrated peak area scaled by the largest peak area of each compound.

similar systems,^{23,24} we performed a search for viable cluster structures using CREST based on the GFN methods,^{36,37} more specifically, by applying the iterative *meta*-dynamic sampling for non-covalently bound systems (NCI-iMTD mode). After identifying the best candidates, we performed single point computations on these geometries using the B3LYP-gCP-D3(BJ)-ABC/def2-TZVPP level of theory.

The obtained interaction energies (Table S3, ESI[†]) for the here studied acid clusters are in line with our previously obtained values for diamantane²³ and diamondoid ether²⁴ clusters, confirming that the aggregation of diamondoid acids in HNDs is indeed an energetically favorable process. However, in contrast to the aforementioned aggregates, intermolecular LD is no longer a decisive factor for cluster aggregation; this role here goes to hydrogen bonding. For example, by observing the spatial arrangement of acid molecules in the computed adamantane **1aCL5** and medial diamantane **2aCL5** clusters (Fig. S3, ESI[†]), one observes the acid pair formation of an

O–H···O=C pattern. As these magic number clusters are odd-numbered, the lone acid molecule alternates between the existing pairs and engages in a bridging interaction. Remarkably, cluster **4aCL4** possesses completely different structural features. As the parent molecule has two acid moieties substituted in the apical positions of the diamantane cage, the resulting cluster consisting of four **4a** molecules readily forms a cyclic supramolecular assembly. In this cavity-forming assembly held together by hydrogen bonding, two bridges engage in interactions of a double O–H···O=C type while the remaining two engage in a single O–H···O=C interaction between the two neighboring acid groups (Fig. S3, ESI[†]). Thus, such structural cluster arrangement directly showcases the power of hydrogen bonding in diamondoid acid agglomeration processes and confirms that neither LD interactions between numerous good dispersion energy donor scaffolds nor solvating stabilization achieved by weak helium interactions can compete with strong hydrogen bond forming functional groups like acid moieties. Unlike small acid dimers of formic acid or acetic acid, larger diamondoid acids seem to be less influenced by the collision energy distribution upon cluster formation and one in principle observes strong O–H···O=C bonding patterns between the molecules. This may be due to the physicochemical nature of diamondoid acids as they consist not only of hydrogen bond forming groups but also incorporate bulky and heavy cage subunits. Note that the principle of cavity tailoring of hydrogen-bonded molecular scaffolds using relatively weak intermolecular contacts was described previously and can be a useful tool for the design of functional supramolecular assemblies.⁴⁸

To confirm that our computational findings are valid beyond just gas phase conditions, we performed the optimization of the cluster consisting of five **1a** molecules (analogous to **1aCL5**) but also having in the beginning each **1a** molecule spatially isolated and surrounded by 100 He atoms. The proceeding of the optimization and the resulting cluster structure confirms that the agglomeration of diamondoid acids is indeed preferential to simple “solvation” of individual molecules with helium atoms (ESI,† **1aCL5_500He**). Moreover, the obtained final structure consisting of a diamondoid acid molecules core and an outer helium shell is comparable to the structure of cluster **1aCL5**, with similar hydrogen bonding patterns between acid moieties observed (Fig. 3). Note also that the emerged molecules core of the final computed **1aCL5_500He** structure (ESI[†]) contains no helium atoms at all dispersed between diamondoid subunits. This clearly demonstrates that not only does the hydrogen bonding region expectedly push out all He out of its reach of influence, but also that diamondoid cages themselves preferentially engage in the intermolecular dispersion interaction rather than in simple “solvation” with helium, in line with our previous findings.^{23,24}

Concerning the number of water molecules observed for acid clusters, smaller assemblies of **1a** and **2a** do not favor the presence of water and when increasing the cluster size one observes the incorporation of a single water molecule (Fig. S1, ESI[†]). From a structural standpoint that is to be expected since



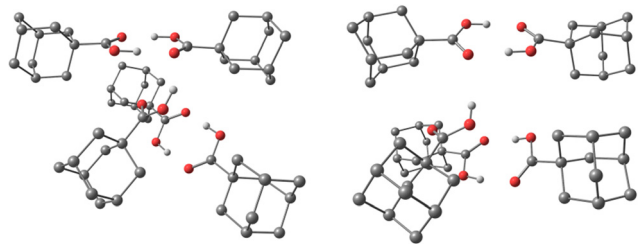


Fig. 3 Optimized structures of **1aCL5** (left) and **1aCL5_500He** (right) obtained from computations (*vide supra*). Color codes: C, grey; H, white; and O, red. Helium atoms and hydrogen atoms attached to carbon are removed for clarity.

the acid moieties engage in favorable intermolecular hydrogen bonding pair-wise patterns and no place for water incorporation is available. As the cluster size and the number of pairs grow, the network seems to benefit from stabilization accomplished by including additional water. In the case of **4a**, larger cluster sizes are capable of incorporating even more water molecules since this diacid is capable of forming cavities of alternating polarity (Fig. S3, ESI[†]), which enables easier access to hydrogen bond partners, *i.e.*, one available O or OH group per molecule.

Diamondoid alcohols

Turning now our attention to alcohols, one in principle expects weaker hydrogen bonding since the alcohol O–H bond is less polarized than its acid counterpart. To date, mostly small primary, secondary and tertiary alcohols were investigated in a HND environment^{49–57} as well as several combined systems, including NaCl–methanol,⁵⁸ fullerene and methanol or ethanol clusters,⁵⁹ anisole–methanol,⁶⁰ and trifluoromethoxybenzene–methanol complexes.⁶¹ However, as previously mentioned, no diamondoid alcohols in HNDs were characterized to date. We therefore performed a similar combined experimental and computational analysis for these alcohols in order to compare the self-assembly behaviors of the two studied classes of hydrogen bond forming groups bound to a diamondoid skeleton. One immediately notices that determination of the preferred cluster sizes for alcohols **1b–4b** is not as straightforward as it was for the analogous acids. The data depicted in Fig. 4 suggest a preference of seven molecules for **1b** and **3b**, four molecules for **2b** and eight molecules for **4b**. In addition, three molecules per cluster also seem to be candidates for **3b** and **4b**. Similar to the acids, for **1b** and **2b**, one can observe a preference for larger clusters to include one water molecule. For **1b**, this transition is analogous to **1a** between cluster sizes of eight and nine molecules; for **2b** it is between six and seven molecules. In contrast to **4a**, no influence of water absorption is observed with **4b**.

When analyzing the computed structures for the respective clusters of alcohols **1b–3b** (Fig. S4, ESI[†]), it becomes apparent that a strict pair-wise directionality seen in acid clusters is missing here. Namely, the alcohol molecules in the computed structures engage in cyclic hydrogen bond-held assemblies

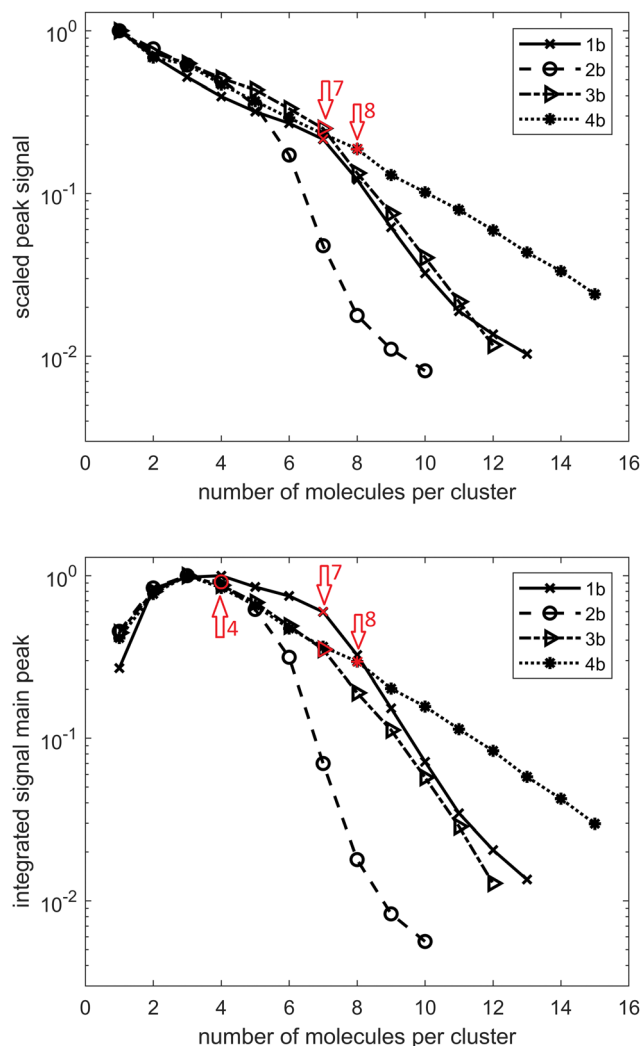


Fig. 4 TOF signal curves for adamantane alcohol **1b**, diamantane medial alcohol **2b**, diamantane mono-apical alcohol **3b**, and diamantane bis-apical diol **4b**. (a) Peak height scaled by the strongest signal of each compound. (b) Integrated peak area scaled by the largest peak area of each compound.

where the OH groups are oriented towards the center of the cluster and form a cavity while the diamondoid part is exposed to the helium environment. Such supramolecular structures can therefore be compared to reverse micelles where the polar parts of the molecules are pointed inwards and form a cavity while the non-polar parts are in contact with solvent molecules of low polarity, in our case helium. It therefore comes as no surprise that no strong preference for specific magic number clusters was observed; many cyclic arrangements of this type are feasible and energetically favorable. The formation of comparable cyclic clusters was also observed previously for methanol.⁶² However, although structurally less ordered and weaker in individual hydrogen bond interactions, assemblies of this type could be of even greater practical importance than their acid counterparts since they are capable of spontaneous polar cavity formation in a highly non-polar medium. As reverse micelles are often used as efficient nano-carriers capable of



targeted delivery and/or solubilization of small molecules, including drugs, production and characterization of more advanced supramolecular architectures in the HNDs is an appealing venue of investigation. For example, a two substrate system that would include doping of HNDs with both a diamondoid alcohol and a polar substrate of choice or a polar solvent molecule could provide a reverse micelle⁶³ with an incorporated nano-cargo, essentially making an assembly inside the assembly (HNDs are doped with diamondoid cavity self-assemblies that are in turn doped with selected polar molecules).

Clusters of alcohols **1b–3b** also demonstrate the capability of water inclusion and the process is proposed to be aided by the formed central cavity. However, note that the type of cavity present in such alcohol assemblies differs from the one found in the bis-acid cluster **4aCL4**. The bis-acid molecules therefore build a cyclic structure through an alternating pattern, *i.e.*, the rim is composed of repeating diamondoid cage subunits and acid groups. Despite alternating rim polarities, such an arrangement nevertheless fosters water incorporation in the cluster since, as mentioned, the two neighboring acid group pairs of the cluster engaged only in a single O–H...O=C interaction and are therefore free to interact with water molecules through their remaining oxygen groups. On the other hand, diamondoid alcohol clusters that resemble a reverse micelle possess a central cavity that is polar in nature and should be capable of favorably hosting water molecules in an inclusion binding event. Lastly, a marked difference is observed for clusters of bis-alcohol **4b**. Assemblies of this apically bis-substituted diamantane alcohol do not prefer to incorporate water molecules at all and the reason for that can again be elucidated from the respective cluster structure **4bCL8** (Fig. S4, ESI†). One observes that this diol does not form a central polar cavity like the mono-alcohols discussed so far but rather engages in random hydrogen bonding events with the surrounding diol molecules. The result is a poorly organized supramolecular structure that resembles more a network of randomly meshed molecules than an arranged self-assembled system. This goes to show that a feature (double substitution) beneficial for one type of non-covalent cluster formation (cavity of alternating polarity for **4a**) can be detrimental for another type (random organization of **4b** molecules).

Conclusions

We studied self-assembled clusters of both diamondoid acids and alcohols in superfluid helium nanodroplets by means of mass spectrometry and found that the observed magic number abundances differ when compared to both diamondoid hydrocarbons and ethers. While hydrocarbon sphere-like molecules pack much more condensely in HNDs, diamondoid ether packing and cluster abundances are strongly influenced by individual molecular geometries as well as the presence of oxygen atoms. In contrast to these predominately dispersion governed processes, the introduction of functional groups with

higher polarity, which are good hydrogen bond donors, to a diamondoid scaffold drastically changes the molecular self-organization in HNDs. Thus, the computational analysis revealed that mono-substituted diamondoid acids readily form strictly spatially directed networks built up through pair-wise intermolecular hydrogen bonding interactions, while analogous mono-alcohols generate a wide range of cyclic clusters resembling reverse micelles. Additionally, a change was observed in a cluster self-assembly mode when going from mono- to bis-derivatives for both series of compounds. Namely, diamantane bis-apical diacid **4a** forms a cavity having a pattern of alternating polarity, while structurally analogous bis-alcohol **4b** produces somewhat random structural patterns with no central polar cavity that was noted for all studied diamondoid mono-alcohols. The observed cavity formation and spontaneous and spatially ordered supramolecular network organization achieved with these hydrogen-bonding systems can have implications for use in porous material production, emphasizing factors to be taken into consideration when designing materials of tailor-made properties.

Author contributions

The manuscript was written through contributions of all authors. All authors have given approval to the final version of the manuscript.

Conflicts of interest

The authors declare no competing financial interest.

Acknowledgements

We acknowledge financial support from the Croatian Science Foundation (project UIP-2017-05-9653 “Diamondoid scaffolds containing heteroatoms – preparation and application in development of advanced materials”), the Austrian Science Fund (FWF) under a grant P30940-N36, the Croatian Ministry of Science and Education and the Austrian Agency for Education and Internationalisation (WTZ project HR 05/2022), and the COST Action CA21101. M. A. thanks the Austrian Academy of Science for financing a short-term research stay at TU Graz through the Joint Excellence in Science and Humanities (JESH) program. The computations were performed using the resources of the computer cluster Isabella based in SRCE – University of Zagreb, University Computing Centre.

References

- 1 H. Schwertfeger, A. A. Fokin and P. R. Schreiner, *Angew. Chem., Int. Ed.*, 2008, **47**, 1022–1036.
- 2 M. A. Gunawan, J.-C. Hierso, D. Poinot, A. A. Fokin, N. A. Fokina, B. A. Tkachenko and P. R. Schreiner, *New J. Chem.*, 2014, **38**, 28–41.



- 3 J. E. Dahl, S. G. Liu and R. M. K. Carlson, *Science*, 2003, **299**, 96–99.
- 4 J. P. Wagner and P. R. Schreiner, *Angew. Chem., Int. Ed.*, 2015, **54**, 12274–12296.
- 5 K.-W. Yeung, Y. Dong, L. Chen, C.-Y. Tang, W.-C. Law and G. C.-P. Tsui, *Nanotechnol. Rev.*, 2020, **9**, 650–669.
- 6 V. Mozhayskiy, M. N. Slipchenko, V. K. Adamchuk and A. F. Vilesov, *J. Chem. Phys.*, 2007, **127**, 094701.
- 7 W. K. Lewis, C. M. Lindsay and R. E. Miller, *J. Chem. Phys.*, 2008, **129**, 201101.
- 8 J. Poms, A. W. Hauser and W. E. Ernst, *Phys. Chem. Chem. Phys.*, 2012, **14**, 15158–15165.
- 9 P. Thaler, A. Volk, D. Knez, F. Lackner, G. Haberehner, J. Steurer, M. Schnedlitz and W. E. Ernst, *J. Chem. Phys.*, 2015, **143**, 134201.
- 10 A. Volk, P. Thaler, D. Knez, A. W. Hauser, J. Steurer, W. Grogger, F. Hofer and W. E. Ernst, *Phys. Chem. Chem. Phys.*, 2016, **18**, 1451–1459.
- 11 A. Mauracher, O. Echt, A. M. Ellis, S. Yang, D. K. Bohme, J. Postler, A. Kaiser, S. Denifl and P. Scheier, *Phys. Rep.*, 2018, **751**, 1–90.
- 12 W. E. Ernst and A. W. Hauser, *Phys. Chem. Chem. Phys.*, 2021, **23**, 7553–7574.
- 13 F. Lackner, in *Molecules in Superfluid Helium Nanodroplets: Spectroscopy, Structure, and Dynamics*, ed. A. Slenczka and J. P. Toennies, Springer International Publishing, Cham, 2022, pp. 513–560, DOI: [10.1007/978-3-030-94896-2_11](https://doi.org/10.1007/978-3-030-94896-2_11).
- 14 A. Schiller, F. Laimer and L. Tiefenthaler, in *Molecules in Superfluid Helium Nanodroplets: Spectroscopy, Structure, and Dynamics*, ed. A. Slenczka and J. P. Toennies, Springer International Publishing, Cham, 2022, pp. 67–153, DOI: [10.1007/978-3-030-94896-2_3](https://doi.org/10.1007/978-3-030-94896-2_3).
- 15 J. P. Toennies and A. F. Vilesov, *Angew. Chem., Int. Ed.*, 2004, **43**, 2622–2648.
- 16 C. Callegari and W. E. Ernst, in *Handbook of High-resolution Spectroscopy*, ed. F. Merkt and M. Quack, John Wiley & Sons, Chichester, 1st edn, 2011, vol. 3, pp. 1551–1594.
- 17 M. Lasserus, M. Schnedlitz, R. Messner, F. Lackner, W. E. Ernst and A. W. Hauser, *Chem. Sci.*, 2019, **10**, 3473–3480.
- 18 G. Auböck, J. Nagl, C. Callegari and W. E. Ernst, *J. Phys. Chem. A*, 2007, **111**, 7404–7410.
- 19 J. Nagl, G. Auböck, A. W. Hauser, O. Allard, C. Callegari and W. E. Ernst, *Phys. Rev. Lett.*, 2008, **100**, 063001.
- 20 M. Goulart, M. Kuhn, L. Kranabetter, A. Kaiser, J. Postler, M. Rastogi, A. Aleem, B. Rasul, D. K. Bohme and P. Scheier, *J. Phys. Chem. C*, 2017, **121**, 10767–10772.
- 21 L. Kranabetter, P. Martini, N. Gitzl, M. Kuhn, F. Saleem, B. Rasul, M. Mahmoodi Darian, E. Jabbour Al Maalouf, I. Sukuba, A. Kaiser, M. Goulart, D. K. Bohme and P. Scheier, *Phys. Chem. Chem. Phys.*, 2018, **20**, 21573–21579.
- 22 M. Kappe, A. Schiller, S. A. Krasnokutski, M. Ončák, P. Scheier and E. M. Cunningham, *Phys. Chem. Chem. Phys.*, 2022, **24**, 23142–23151.
- 23 J. Alić, R. Messner, F. Lackner, W. E. Ernst and M. Šekutor, *Phys. Chem. Chem. Phys.*, 2021, **23**, 21833–21839.
- 24 J. Alić, R. Messner, M. Alešković, F. Kustner, M. Rubčić, F. Lackner, W. E. Ernst and M. Šekutor, *Phys. Chem. Chem. Phys.*, 2023, **25**, 11951–11958.
- 25 E. Arunan, G. R. Desiraju, R. A. Klein, J. Sadlej, S. Scheiner, I. Alkorta, D. C. Clary, R. H. Crabtree, J. J. Dannenberg, P. Hobza, H. G. Kjaergaard, A. C. Legon, B. Mennucci and D. J. Nesbitt, *Pure Appl. Chem.*, 2011, **83**, 1637–1641.
- 26 J. W. Steed and J. L. Atwood, *Supramolecular chemistry*, John Wiley & Sons, 2022.
- 27 J. W. Steed, D. R. Turner and K. J. Wallace, *Core concepts in supramolecular chemistry and nanochemistry*, John Wiley & Sons, 2007.
- 28 J. L. Cook, C. A. Hunter, C. M. R. Low, A. Perez-Velasco and J. G. Vinter, *Angew. Chem., Int. Ed.*, 2007, **46**, 3706–3709.
- 29 M. C. Davis and S. Liu, *Synth. Commun.*, 2006, **36**, 3509–3514.
- 30 N. A. Fokina, B. A. Tkachenko, J. E. P. Dahl, R. M. K. Carlson, A. A. Fokin and P. R. Schreiner, *Synthesis*, 2012, 259–264.
- 31 N. A. Fokina, B. A. Tkachenko, A. Merz, M. Serafin, J. E. P. Dahl, R. M. K. Carlson, A. A. Fokin and P. R. Schreiner, *Eur. J. Org. Chem.*, 2007, 4738–4745.
- 32 F. Neese, *Wiley Interdiscip. Rev.: Comput. Mol. Sci.*, 2012, **2**, 73–78.
- 33 F. Neese, *Wiley Interdiscip. Rev.: Comput. Mol. Sci.*, 2018, **8**, e1327.
- 34 S. Grimme, J. Antony, S. Ehrlich and H. Krieg, *J. Chem. Phys.*, 2010, **132**, 154104.
- 35 S. Grimme, S. Ehrlich and L. Goerigk, *J. Comput. Chem.*, 2011, **32**, 1456–1465.
- 36 P. Pracht, F. Bohle and S. Grimme, *Phys. Chem. Chem. Phys.*, 2020, **22**, 7169–7192.
- 37 S. Grimme, *J. Chem. Theory Comput.*, 2019, **15**, 2847–2862.
- 38 S. Grimme, *Wiley Interdiscip. Rev.: Comput. Mol. Sci.*, 2011, **1**, 211–228.
- 39 H. Kruse and S. Grimme, *J. Chem. Phys.*, 2012, **136**, 154101.
- 40 M. Mahmoodi-Darian, L. Lundberg, S. Zöttl, P. Scheier and O. Echt, *J. Am. Soc. Mass Spectrom.*, 2019, **30**, 787–795.
- 41 P. Das, C. J. Knapp and W. Jäger, *J. Mol. Spectrosc.*, 2017, **341**, 17–22.
- 42 F. Madeja, M. Havenith, K. Nauta, R. E. Miller, J. Chocholoušová and P. Hobza, *J. Chem. Phys.*, 2004, **120**, 10554–10560.
- 43 K. A. E. Meyer, J. A. Davies and A. M. Ellis, *Phys. Chem. Chem. Phys.*, 2020, **22**, 9637–9646.
- 44 F. Ferreira da Silva, S. Jaksch, G. Martins, H. M. Dang, M. Dampe, S. Denifl, T. D. Märk, P. Limão-Vieira, J. Liu, S. Yang, A. M. Ellis and P. Scheier, *Phys. Chem. Chem. Phys.*, 2009, **11**, 11631–11637.
- 45 J. A. Davies, M. W. D. Hanson-Heine, N. A. Besley, A. Shirley, J. Trowers, S. Yang and A. M. Ellis, *Phys. Chem. Chem. Phys.*, 2019, **21**, 13950–13958.
- 46 J. A. Davies, N. A. Besley, S. Yang and A. M. Ellis, *J. Phys. Chem. Lett.*, 2019, **10**, 2108–2112.
- 47 A. M. Ellis, J. A. Davies, E. Yurtsever and F. Calvo, *J. Chem. Phys.*, 2022, **156**, 174304.



- 48 A. Navarro-Huerta, M. J. Jellen, J. Arcudia, S. J. Teat, R. A. Toscano, G. Merino and B. Rodríguez-Molina, *Chem. Sci.*, 2021, **12**, 2181–2188.
- 49 W. K. Lewis, B. E. Applegate, J. Sztáray, B. Sztáray, T. Baer, R. J. Bemish and R. E. Miller, *J. Am. Chem. Soc.*, 2004, **126**, 11283–11292.
- 50 S. Yang, S. M. Brereton, M. D. Wheeler and A. M. Ellis, *Phys. Chem. Chem. Phys.*, 2005, **7**, 4082–4088.
- 51 S. Yang, S. M. Brereton and A. M. Ellis, *Int. J. Mass Spectrom.*, 2006, **253**, 79–86.
- 52 A. Boatwright, J. Jeffs and A. J. Stace, *J. Phys. Chem. A*, 2007, **111**, 7481–7488.
- 53 M. Goulart, P. Bartl, A. Mauracher, F. Zappa, A. M. Ellis and P. Scheier, *Phys. Chem. Chem. Phys.*, 2013, **15**, 3577–3582.
- 54 F. J. Hernandez, J. T. Brice, C. M. Leavitt, G. A. Pino and G. E. Douberly, *J. Phys. Chem. A*, 2015, **119**, 8125–8132.
- 55 D. Mani, T. Fischer, R. Schwan, A. Dey, B. Redlich, A. F. G. Van der Meer, G. Schwaab and M. Havenith, *RSC Adv.*, 2017, **7**, 54318–54325.
- 56 D. Mani, N. Pal, M. Smialkowski, C. Beakovic, G. Schwaab and M. Havenith, *Phys. Chem. Chem. Phys.*, 2019, **21**, 20582–20587.
- 57 R. Messner, W. E. Ernst and F. Lackner, *J. Phys. Chem. Lett.*, 2021, **12**, 145–150.
- 58 A. M. Sadoon, G. Sarma, E. M. Cunningham, J. Tandy, M. W. D. Hanson-Heine, N. A. Besley, S. Yang and A. M. Ellis, *J. Phys. Chem. A*, 2016, **120**, 8085–8092.
- 59 M. Goulart, F. Zappa, A. M. Ellis, P. Bartl, S. Ralser and P. Scheier, *Phys. Chem. Chem. Phys.*, 2017, **19**, 24197–24201.
- 60 T. K. Roy, D. Mani, G. Schwaab and M. Havenith, *Phys. Chem. Chem. Phys.*, 2020, **22**, 22408–22416.
- 61 T. K. Roy, D. Mani, G. Schwaab and M. Havenith, *Phys. Chem. Chem. Phys.*, 2021, **23**, 25180–25187.
- 62 M. I. Sulaiman, S. Yang and A. M. Ellis, *J. Phys. Chem. A*, 2017, **121**, 771–776.
- 63 N. M. Correa, J. J. Silber, R. E. Riter and N. E. Levinger, *Chem. Rev.*, 2012, **112**, 4569–4602.

