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Sulfur–arene interactions: the S··· π and S–H··· π interactions in the dimers of benzofuran···sulfur dioxide and benzofuran···hydrogen sulfide†

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Non-covalent interactions between sulfur centers and aromatic rings play important roles in biological chemistry. We examined here the sulfur–arene interactions between the fused aromatic heterocycle benzofuran and two prototype sulfur divalent triatomics (sulfur dioxide and hydrogen sulfide). The weakly-bound adducts were generated in a supersonic jet expansion and characterized with broadband (chirped-pulsed) time-domain microwave spectroscopy. The rotational spectrum confirmed the detection of a single isomer for both heterodimers, consistent with the computational predictions for the global minima. The benzofuran···sulfur dioxide dimer exhibits a stacked structure with sulfur closer to benzofuran, while in benzofuran···hydrogen sulfide the two S–H bonds are oriented towards the bicycle. These binding topologies are similar to the corresponding benzene adducts, but offer increased interaction energies. The stabilizing interactions are described as S··· π or S–H··· π , respectively, using a combination of density-functional theory calculations (dispersion corrected B3LYP and B2PLYP), natural bond orbital theory, energy decomposition and electronic density analysis methods. The two heterodimers present a larger dispersion component, but nearly balanced by electrostatic contributions.

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Introduction

Sulfur-containing compounds are key components of living organisms, playing important biochemical and metabolic roles. As an example, the sulfur amino acids methionine and cysteine are incorporated into proteins, with the first one initiating the synthesis of virtually all eukaryotic proteins and the latter critically influencing protein folding by formation of disulfide bonds.¹ For this reason, the analysis of sulfur-centered non-covalent interactions² is of fundamental interest to understand the biological functions and reactivity of these compounds. Among the multiple non-covalent interactions observed in protein tertiary structures, the presence of sulfur–arene contacts was recognized early as a significant contributor,^{3,4} as later confirmed from crystal data mining from the Brookhaven Protein Data Bank (PDB) and

the Cambridge Crystallographic Database (CCD).⁵ However, it may be difficult to extract accurate information from individual intermolecular interactions using large macromolecular structures with conventional (*ca.* 2.5 Å) crystallographic resolution. In these cases, both computational modelling and gas-phase experimental studies can provide complementary chemically specific structural information and a physical description of the molecular forces at play. Previous computational studies concerning the interaction between aromatic rings and sulfur compounds have focused on small prototype heterodimers like methanethiol···benzene⁶ or hydrogen sulfide···benzene,^{7,8} characterized by a S–H··· π weak hydrogen bond. The use of sulfur non-covalent interactions in drug design and larger examples of arene complexation have also been reviewed elsewhere.^{9,10} However, the number of high-resolution gas-phase spectroscopic studies of aggregation between sulfur compounds and π -systems is small. Examples include the dimers of hydrogen sulfide with benzene,^{11,12} phenylacetylene,¹³ phenol,^{14,15} *p*-cresol¹⁶ or indole,¹⁷ or the sulfur dioxide dimers with benzene,^{18–20} toluene²¹ or furan,²² loosely classified as van der Waals adducts. Here, we will explore the interactions between divalent sulfur and aromatic rings using broadband (chirped-pulse) microwave spectroscopy and jet expansions, extending previous rotationally resolved investigations on sulfur dioxide aggregation²³ and thiol microsolvation^{24,25} and dimerization.^{26,27} We chose heterodimers of sulfur dioxide (SO₂) and hydrogen

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sulfide (H_2S) to simultaneously explore both the $\text{S}\cdots\pi$ and the hydrogen bonded $\text{S}\cdots\text{H}\cdots\pi$ interactions. In the present work the fused heterocycle benzofuran was selected as aggregation partner because of the presence of three plausible binding sites at the two six- and five-membered rings and the oxygen heteroatom. The observation of sulfur clusters is expected to offer comparison with the equivalent oxygen partners, gauging the influence of the lower electronegativity, larger polarizability and reduced hydrogen bond capacity of sulfur. This information will contribute to the description of weak non-covalent interactions involving chalcogens, which has expanded considerably in the last decade.^{28,29}

Methods

Experiment

The dimers of benzofuran with sulfur dioxide ($\text{BF}\cdots\text{SO}_2$) and hydrogen sulfide ($\text{BF}\cdots\text{H}_2\text{S}$) were prepared by supersonic expansion of the two precursors, highly diluted with neon as carrier gas. For this purpose, benzofuran was heated to *ca.* 310 K inside a solenoid-driven pulsed-nozzle (Parker Series 9) and pressurized with a gas line containing each of the sulfur compounds at typical concentrations of 0.5%. The expansion from stagnation pressures of *ca.* 2–3 bar to the vacuum chamber ultimate pressure of *ca.* 10^{-6} – 10^{-7} mbar created a vertically moving molecular jet. The nozzle orifice was 1 mm. All chemical samples were obtained commercially.

The microwave spectrum was recorded in the 2–8 GHz region with a direct-digital broadband (chirped-pulse) Fourier transform microwave spectrometer.³⁰ This spectrometer uses a fast (20 GS/s) digital arbitrary waveform generator to create a series of short (4 μs) microwave chirped-pulses, which are amplified to 20 W. The exciting pulse propagates perpendicularly to the jet molecular expansion, inducing a broadband fast-passage³¹ electric-dipole excitation. The resulting transient emission or free-induced decay is amplified, collected in the time-domain with a digital oscilloscope for *ca.* 40 μs per pulse, and Fourier-transformed to yield the frequency spectrum. The resonance frequencies have typical linewidths (FWHM) below 100 kHz after apodization with a Kaiser–Bessel window. The spectra were averaged for about 1 M cycles. The final frequency measurements have estimated accuracies of 10 kHz or better.

Computations

Initial guess structures were generated with a simple molecular mechanics force field (MMFF).³² These structures were then optimized with quantum mechanical density functional theory (DFT) methods. In this work, the adduct structures were first obtained using the B3LYP hybrid method and D3 two-body empirical dispersion corrections³³ (Becke–Johnson damping function³⁴). The most stable structures were later reoptimized with the B2PLYP double-hybrid method.³⁵ The Ahlrich's triple- ζ polarized basis set (def2-TZVP),³⁶ which performs reasonably well in similar clusters,²⁷ was used initially in combination with an ultrafine integration grid and normal convergence criteria. For comparison, the calculations were repeated using Dunning's

correlation-consistent cc-pVTZ basis set³⁷ and the jun-cc-pVTZ variation,³⁸ removing the highest angular momentum diffuse functions. The vibrational frequencies were evaluated within the harmonic approximation. For the calculation of complexation energies the basis-sets superposition errors were accounted for within the Boys–Bernardi approximation.³⁹ The natural-bond-orbital analysis⁴⁰ used NBO version 3. All calculations were implemented in Gaussian 16.⁴¹ For the non-covalent interactions, the Johnson–Contreras analysis⁴² of the reduced electronic density function used NCIPLOT.⁴³ An energy decomposition analysis was conducted with symmetry-adapted perturbation theory⁴⁴ (SAPT) at the SAPT 2 + 3(CCD)/aug-cc-pVDZ level, implemented in PSI4.⁴⁵

Results and discussion

Conformational landscape

The predictions for $\text{BF}\cdots\text{SO}_2$ confirmed the preference for stacked structures, similar to the heterodimers with benzene¹⁸ or furan.²² However, the larger number of non-equivalent binding sites on benzofuran results in four different isomers below 5 kJ mol^{-1} (B3LYP-D3(BJ)), shown in Fig. 1. For the B2PLYP method the last isomer converged to the global minimum. In these structures the sulfur atom of SO_2 sits on top of the benzene unit, with the oxygen atoms pointing outwards the ring (rotatable 3D Fig. S1, ESI[†]). The SO_2 moiety is not coplanar with the heterocycle, instead displaying a tilted orientation with sulfur closer to the ring and the two oxygens pointing outwards the ring. The predicted rotational and energetic parameters are presented in Table 1 and Tables S1, S2 (ESI[†]). In the case of the $\text{BF}\cdots\text{H}_2\text{S}$ dimer the situation is different, as the predictions in Table 2 and Tables S3, S4 (ESI[†]) suggest a tendency to form $\text{S}\cdots\text{H}\cdots\pi$ hydrogen bonds. In the most stable isomer of Fig. 2 the sulfur atom is located in between the two

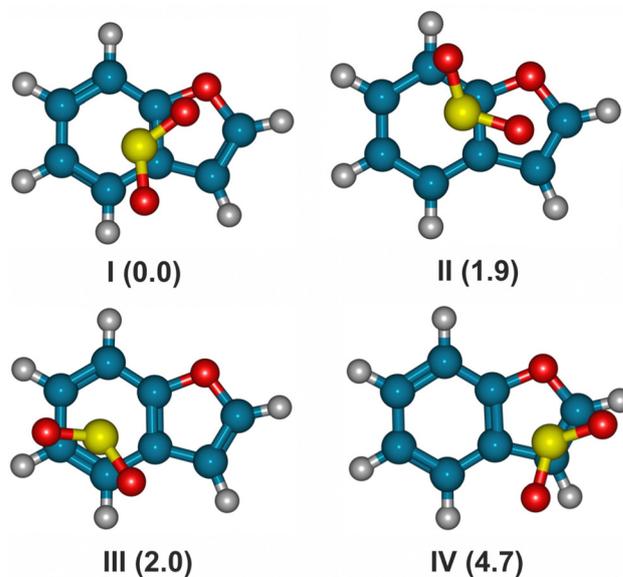


Fig. 1 The predicted most stable isomers of the dimer benzofuran \cdots sulfur dioxide (B3LYP-D3(BJ)/def2-TZVP electronic energies in parentheses, see Table 1 and Tables S1, S2, ESI[†]).



Table 1 Rotational parameters of the benzofuran...sulfur dioxide dimer

	Theory ^d			Experiment ^e	
	Isomer I	Isomer II	Isomer III	³² S species	³⁴ S species
A/MHz ^a	1070.1	1066.0	1244.1	1056.036(3)	1054.90(1)
B/MHz	830.5	830.1	666.0	811.9523(8)	802.077(1)
C/MHz	658.9	629.4	559.1	642.2063(9)	635.9851(9)
Δ _J /kHz	0.53	0.62	0.53	0.26(2)	[0.26]
Δ _{JK} /kHz	-0.16	2.0	-1.9	2.7(1)	[2.7]
Δ _K /kHz	-0.19	-2.3	2.9	-2.7(3)	[-2.7]
δ _J /kHz	0.003	-0.04	0.18	[0.0] ^f	[0.0]
δ _K /kHz	0.54	3.7	0.90	1.1(1)	[1.1]
μ _a /D	1.5	1.8	2.2	+	+
μ _b /D	0.8	0.7	0.2	+	+
μ _c /D	0.1	1.5	0.1	-	-
N ^b				53	17
σ/kHz				8.0	8.2
ΔE/kJ mol ^{-1c}	0.0	1.9	0.2		
ΔE ₀ /kJ mol ⁻¹	0.0	1.7	2.1		
ΔG/kJ mol ⁻¹	1.5	0.3	0.0		
ΔE _c /kJ mol ⁻¹	-21.30	-19.25	-19.41		

^a Rotational constants (A, B, C), Watson's A-reduction centrifugal distortion constants (Δ_J, Δ_{JK}, Δ_K, δ_J, δ_K) and electric dipole moments (μ_α, α = a, b, c).

^b Number of measured transitions (N) and standard deviation of the fit (σ).

^c Relative energies uncorrected (ΔE) and corrected with the zero-point energy (ΔE₀), Gibbs energy (ΔG, 298K, 1 atm) and complexation energy (ΔE_c, including BSSE corrections). ^d Calculations using B2PLYP-D3(BJ) and the def2-TZVP basis set, see also Tables S1 and S2 (ESI[†]). ^e Standard errors in parentheses in units of the last digit. ^f Values fixed in the fit.

rings, with the two S-H bonds pointing down, respectively, to the six- and the five-membered ring (rotatable 3D Fig. S2, ESI[†]). A second isomer *ca.* 2 kJ mol⁻¹ less stable has only one hydrogen atom oriented to the ring.

Rotational spectra

The experimental microwave spectra of the two benzofuran heterodimers are shown in Fig. S3 and S4 (ESI[†]), with the most

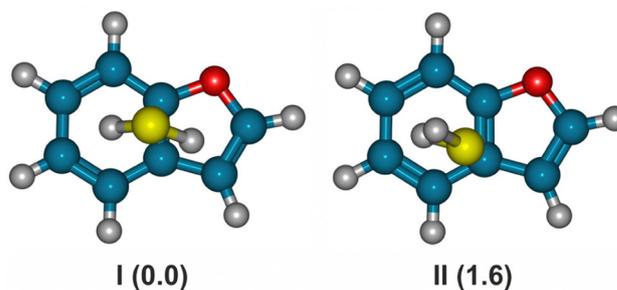


Fig. 2 The predicted most stable isomers of the dimer benzofuran...hydrogen sulfide (B3LYP-D3(BJ)/def2-TZVP electronic energies in parentheses, see Table 2 and Tables S3 and S4, ESI[†]).

prominent transitions being easily attributable to the benzofuran monomer.^{46–48} The transitions of the two adducts, generally of 1/20–1/30 intensity with respect to the monomer, had appreciable signal-to-noise ratios (SNR < 150 : 1) and revealed a single isomer for both aggregates. The spectrum of BF...SO₂, was made of R- ($J + 1 \leftarrow J$) and Q-branch ($J \leftarrow J$) μ_a- and μ_b-type transitions ($J = 1–5$, $K_{-1} = 0–4$). As shown in Fig. 3, the spectral intensity permitted the additional assignment of the ³⁴S isotopologue in natural abundance (*ca.* 4%). In both cases, the Watson A-reduced Hamiltonian reproduced the experiment to experimental accuracy (Table 1). Conversely, the spectrum of BF...H₂S showed evidence of a large-amplitude motion within the cluster, as the rotational transitions were split into two asymmetric hyperfine components, illustrated in Fig. 4. The exchange of two half-spin fermions would result in a 1:3 statistical weight for the 0⁺ and 0⁻ torsional states, respectively, which is close to the observations. For this reason, the tunneling splitting is ascribed to the exchange of the two hydrogen atoms by internal rotation of H₂S around its symmetry axis. This effect has been observed in microsolvation clusters where

Table 2 Rotational parameters of the benzofuran...hydrogen sulfide dimer

	Theory		Experiment			
	Isomer I	Isomer II	³² S species		³⁴ S species	
			ν = 0 ⁺	ν = 0 ⁻	ν = 0 ⁺	ν = 0 ⁻
A/MHz ^a	1167.9	1197.7	1174.3380(9)	1173.5247(9)	1172.814(4)	1172.033(4)
B/MHz	1087.7	1112.7	1037.3505(7)	1037.7417(6)	1004.705(1)	1005.059(1)
C/MHz	788.6	812.4	763.8964(5)	764.0018(4)	745.390(1)	745.495(2)
Δ _J /kHz	0.64	2.5	1.17(2)		1.10(4)	
Δ _{JK} /kHz	4.1	-6.9	2.47(8)		3.6(2)	
Δ _K /kHz	-4.5	4.8	-2.6(1)		-3.8(4)	
δ _J /kHz	-0.20	-0.28	0.449(9)		[0.449]	
δ _K /kHz	18.2	-13.0	1.81(4)		[1.81]	
μ _a /D	1.4	0.1	+		+	
μ _b /D	0.2	0.5	-		-	
μ _c /D	0.9	0.2	+		+	
N			30	30	15	19
σ/kHz			3.9		1.3	
ΔE/kJ mol ⁻¹	0.0	1.3				
ΔE ₀ /kJ mol ⁻¹	0.0	1.5				
ΔG/kJ mol ⁻¹	0.0	1.0				
ΔE _c /kJ mol ⁻¹	-14.02	-13.10				

^a Parameter definition as in Table 1. The theoretical predictions used the B2PLYP-D3(BJ) method and the def2-TZVP basis set, see also Tables S3 and S4 (ESI).



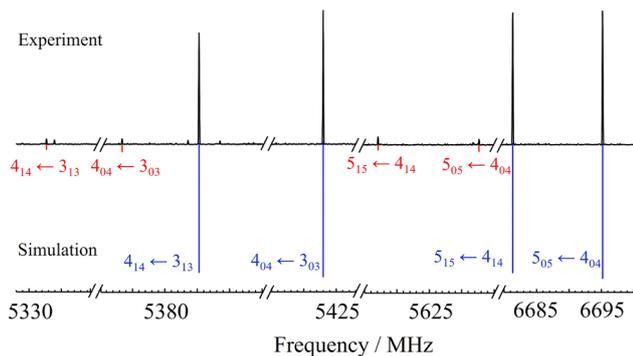


Fig. 3 A section of the microwave spectrum of the dimer benzofuran...sulfur dioxide with transitions from the isotopologues with ^{32}S (blue) and ^{34}S (red). See full spectrum in Fig. S3 (ESI †).

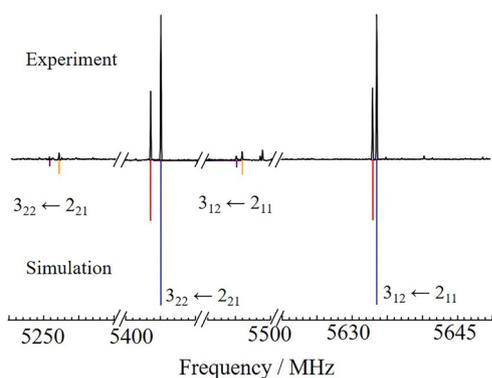


Fig. 4 A section of the microwave spectrum of the dimer benzofuran...hydrogen sulfide, showing the tunnelling splittings for the isotopologues with ^{32}S (red/blue) and ^{34}S (black/orange). See full spectrum in Fig. S4 (ESI †).

water behaves as weak proton donor,⁴⁹ as in benzene...water.⁵⁰ In our case the two torsional components were fitted independently to semirigid models, floating the rotational constants of each state, but sharing the centrifugal distortion parameters. The observations included R-branch μ_a - and μ_c -transitions ($J = 0-4$, $K_{-1} = 0-3$) for the parent and the ^{34}S isotopologue in natural abundance. The results of the fit are collected in Table 2. The observed transitions are listed in Tables S5–S8 (ESI †). No other dimer species were detected in the spectra.

The observed selection rules, rotational constants (Tables 1 and 2) and sulfur substitution coordinates⁵¹ (Tables S9 and S10, ESI †) were consistent with the observation of the global minimum species for the two dimers (predicted coordinates in Tables S11–S14, ESI †). Relative differences between the observed and calculated rotational constants are shown in Tables S1–S4 (ESI †). Discrimination of the different isomers is quite unambiguous. However, for the SO_2 heterodimer the rotational constants of isomers I and II are relatively close. In this case, the absence of μ_c spectrum confirmed the observation of isomer I. Fig. S5 (ESI †) illustrates the increase of μ_c in isomer II of $\text{BF}\cdots\text{SO}_2$. A comparison of the centrifugal distortion constants is not significant in this case because not all

Table 3 Effective molecular structures of the dimers of benzofuran with sulfur dioxide and hydrogen sulfide and B2PLYP-D3(BJ) predictions

	BF...SO ₂		BF...H ₂ S	
	r_0^b	r_e^c	r_0	r_e
$r_{\text{CM}} = r(\text{X}_S\text{-X}_{\text{BZ}})/\text{\AA}^a$	3.307(2) ^d	3.237	3.660(3)	3.567
$\angle (\text{X}_S\text{-X}_{\text{BZ}}\text{-X})/\text{deg}$	90.7(8)	95.1	87.7(2)	88.9
$\tau (\text{X}_S\text{-X}_{\text{BZ}}\text{-X-C}_5)/\text{deg}$	-78.7(2)	-80.8	-85.7(4)	-88.2
$r(\text{H/O-X}_S)/\text{\AA}$	[1.286] ^e	1.304	[1.303]	1.304
$\angle (\text{H/O-X}_S\text{-X}_{\text{BZ}})/\text{deg}$	[108.7]	108.6	[45.9]	45.8
$\tau (\text{H/O-X}_S\text{-X}_{\text{BZ}}\text{-X})/\text{deg}$	[80.3]	79.6	[-0.6]	-0.6
$r(\text{S-H/O})/\text{\AA}$	[1.431]	1.458	[1.336]	1.341
$\angle (\text{S-H/O-X}_S)/\text{deg}$	[14.0]	14.3	[1.7]	1.7
$\tau (\text{S-H/O-X}_S\text{-X}_{\text{BZ}})/\text{deg}$	[-66.9]	-66.4	[-164.4]	-164.5

^a X_S , X_{BZ} and X denote, respectively, the center of mass of the sulfur triatomic, the center of mass of the benzofuran monomer and the middle point between atoms C5 and C6 of benzofuran. ^b Monomer structures in ref. 47 and 53. ^c Equilibrium structures according to B2PLYP-D3(BJ)/cc-pVTZ. ^d Standard errors in parentheses in units of the last digit. ^e Derived parameters in square brackets.

parameters could be determined experimentally and because of the limitations of the harmonic approximation. The experiment also permitted the determination of a partial effective structure⁵² for both dimers, presented in Table 3. Assuming that the two moieties of the dimer^{47,53} remain unperturbed, six structural parameters are required to locate the triatomics over the benzofuran ring.²² However, since only two isotopic species are available, we fitted only the distance between the two centers of mass ($r_{\text{CM}} = r(\text{X}_S\text{-X}_{\text{BZ}})$) and the elevation ($\angle (\text{X}_S\text{-X}_{\text{BZ}}\text{-X})$) and dihedral ($\tau (\text{X}_S\text{-X}_{\text{BZ}}\text{-X-C}_5)$) of each triatomic with respect to benzofuran. Other parameters relating to the orientation of the triatomics were kept fixed to the B2PLYP structure. The resulting structures reproduce the observed rotational constants within 0.5 MHz for $\text{BF}\cdots\text{SO}_2$ and 0.4 MHz for $\text{BF}\cdots\text{H}_2\text{S}$ (Tables S15 and S16, ESI †). The effective structures are compared with theory in Table 3 and Tables S11, S13 (ESI †). The fitted center of mass distance in $\text{BF}\cdots\text{SO}_2$ (3.307(2) Å) is slightly shorter than those observed in the SO_2 dimers with benzene (3.485(1) Å¹⁸) or furan (3.432(1) Å²²), but they share a general stacked pattern characteristic of the $\text{S}\cdots\pi$ interaction, with the sulfur end of the triatomic closer to the ring. For the $\text{BF}\cdots\text{H}_2\text{S}$ dimer the center of mass distance (3.660(3) Å) is also shorter than in benzene... H_2S (3.771 Å¹¹), where the $\text{S-H}\cdots\pi$ interaction produces a symmetric top. The shorter distances are indicative of stronger interaction energies in the dimers with benzofuran, as suggested by the computational binding energies below.

The presence of a single isomer for the two dimers contrasts with the prediction of low energy isomers below 2 kJ mol⁻¹.



This fact may be attributed to the presence of conformational relaxation pathways in the jet expansion,⁵⁴ associated to small interconversion barriers. Considering the weak intermolecular forces (discussed below), and that the conversion between the stable isomers involves only small structural adjustments of the triatomics, the potential barriers are expected *a priori* to be small. We confirmed this fact with the calculation of interconversion barriers between the two most stable isomers using DFT (B3LYP-D3(BJ)/cc-pVTZ). The predicted direct/(reverse) barriers of 2.6/(0.5) and 2.4/(0.9) kJ mol⁻¹ for BF $\cdot\cdot$ SO $_2$ and BF $\cdot\cdot$ H $_2$ S, respectively, are much smaller than the empirical conversion thresholds of 5–7 kJ mol⁻¹ for collisional relaxation,⁵⁴ offering quantitative support for this hypothesis.

The non-covalent interactions in the observed dimers were analysed using several computational methods. The natural-bond-orbital (NBO) calculations in Tables S17 and S18 (ESI \dagger) use atomic orbital overlap and density matrices to form a basis of orthonormal bond orbitals ordered by a maximum occupancy criterion. The charge transfer interactions are associated with occupancy changes from the filled to the unfilled orbitals. The most important donor–acceptor orbital associations in the BF $\cdot\cdot$ SO $_2$ dimer correspond to $\pi \rightarrow \sigma^*(S)$ interactions, involving the donor bonding orbital BD(C $_{3a}$ –C $_4$) of benzofuran and three anti-bonding orbitals of SO $_2$ (perturbative interaction energies $E^{(2)} = 1.0$ – 2.7 kJ mol⁻¹). This interaction is complemented by $\pi(S) \rightarrow \sigma^*(C_{3a})$ and $n(S) \rightarrow \sigma^*(C_{3a})$ back donation from SO $_2$ ($E^{(2)} = 1.0$ kJ mol⁻¹), so they can be properly categorized as S $\cdot\cdot$ π bonding. For BF $\cdot\cdot$ H $_2$ S the main interactions are $\pi \rightarrow \sigma^*(SH)$, involving three π donor orbitals in the ring and non-bonding orbitals of H $_2$ S ($E^{(2)} = 0.6$ – 1.7 kJ mol⁻¹). Weaker $\pi\cdot\cdot\pi^*$ (in BF $\cdot\cdot$ SO $_2$) and $n\cdot\cdot\pi^*$ interactions originated by sulfur additionally contribute to dimer formation. As expected, the total interaction energies are much larger (*ca.* $\times 1.5$) for the sulfur dioxide dimer (B2PLYP-D3(BJ): -20.5 to -22.2 kJ mol⁻¹; B3LYP-D3(BJ): -25.4 to -25.7 kJ mol⁻¹, see Table 1 and Tables S1, S2, ESI \dagger) than for hydrogen sulfide (B2PLYP-D3(BJ): -14.0 to -14.8 kJ mol⁻¹; B3LYP-D3(BJ): -16.0 to -16.3 kJ mol⁻¹, see Table 2 and Tables S3, S4, ESI \dagger). A similar trend is observed for a calculation of complexation energies for the benzene dimers with SO $_2$ (B2PLYP-D3(BJ)/def2-TZVP: -17.9 kJ mol⁻¹) and H $_2$ S (-11.6 kJ mol⁻¹). Experimental complexation energies are available for the dimer of benzene $\cdot\cdot$ SO $_2$,²⁰ permitting an estimation of uncertainties for the computational values. The dissociation energy of $-18.4(12)$ kJ mol⁻¹ obtained from photoionization techniques suggests that the B2PLYP values are slightly underestimated (*ca.* 3%) respect to the experiment. To the best of our knowledge, no experimental dissociation energies for H $_2$ S dimers with benzene or furan have been determined.

A spatial distribution of the non-bonding interactions in the dimer was obtained from a Johnson–Contreras⁴² electronic density analysis, based on a density (ρ) over gradient power ($\nabla\rho$) ratio. The reduced density gradient $s \left(= \frac{1}{2(3\pi^2)^{1/3}} \times \frac{|\nabla\rho|}{\rho^{4/3}} \right)$ isosurfaces in Fig. 5 permit distinguishing attractive or repulsive regions by the sign of the second Hessian eigenvalue of the electronic density (sign($\lambda_2\rho$)), mapping the surfaces associated to

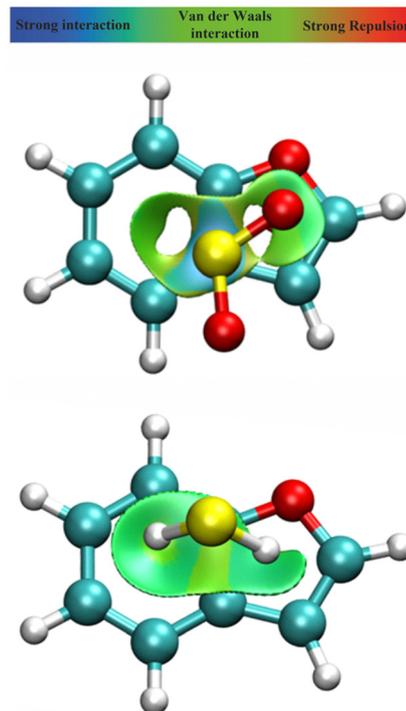


Fig. 5 NCI plot^{39,40} for the dimers of benzofuran with sulfur dioxide (upper panel) and hydrogen sulfide (lower panel), producing a spatial description of attractive interactions in the dimer (isosurface $s = 0.5$ a.u.).

attractive (coloured in blue) or weakly attractive (coloured in green) non-covalent interactions. Fig. 5 shows large delocalized regions between the two rings for both BF $\cdot\cdot$ SO $_2$ and BF $\cdot\cdot$ H $_2$ S, compatible with the interactions noticed in the NBO analysis. For BF $\cdot\cdot$ SO $_2$ the main interaction region is located above the C $_{3a}$ –C $_4$ and C $_{7a}$ –C $_{3a}$ bonds, associated to the S $\cdot\cdot$ π interactions, while both the C $_{3a}$ –C $_{7a}$ bonds and the ring centers present more intense interactions in BF $\cdot\cdot$ H $_2$ S, corresponding to the SH $\cdot\cdot$ π interactions.

The results of the SAPT energy decomposition analysis are presented in Table 4, giving insight into the components of the interaction energy. The two dimers similarly present a larger dispersion component (41% and 55%, respectively, of the total attractive energy) but nearly balanced by electrostatic contributions (34% and 35%, respectively). This energetic distribution is probably associated to the prevailing role in these complexes of the π ring electronic density and the large and polarizable

Table 4 The SAPT energy decomposition analysis of the non-covalent interactions in the benzofuran $\cdot\cdot$ sulfur dioxide and benzofuran $\cdot\cdot$ hydrogen sulfide dimers (all values in kJ mol⁻¹) using SAPT2 + 3(CCD)/aug-cc-pVDZ method

	BF $\cdot\cdot$ SO $_2$	BF $\cdot\cdot$ H $_2$ S
Electrostatics	-24.6 [33.6%] ^a	-12.7 [35.4%]
Induction	-18.8 [25.6%]	-3.4 [9.4%]
Dispersion	-29.9 [40.8%]	-19.8 [55.2%]
Exchange	47.1	21.1
Total	-26.2	-14.8

^a Relative percentages with respect to the total attractive energy in square brackets.



molecular orbital distribution in the sulfur triatomics. In this sense both the $S \cdots \pi$ and $S-H \cdots \pi$ interactions occupy an intermediate position between the dominant dispersive character observed in π -stacking^{26,27,55} and dimers involving a primary hydrogen bond, particularly alcohols,^{56–58} but also some thiols like the prototypic hydrogen sulfide dimer.⁵⁹ On the other hand, the dimer of benzyl mercaptan,⁶⁰ which combines $S-H \cdots S$ and $S-H \cdots \pi$ interactions, presents an energy distribution close to the benzofuran dimers.

Conclusions

Sulfur dioxide and hydrogen sulfide form weak dimers with benzofuran, as observed in a jet-cooled expansion by broadband rotational spectroscopy. In both aggregates a single most stable species is observed experimentally, despite DFT quantum mechanical predictions anticipating alternative low energy isomers below 2 kJ mol⁻¹. Because the conversion between stable isomers was shown to proceed through small potential barriers (<5–7 kJ mol⁻¹), the collisional relaxation mechanisms can explain the presence of a single isomer.⁵⁴ The carriers of the spectrum are reproduced to experimental uncertainty with semirigid rotor Hamiltonian models. Unlike similar SO₂ dimers of higher symmetry, like those with benzene or furan,^{18,22} no tunnelling motions associated to internal large-amplitude motions were observed in benzofuran \cdots sulfur dioxide. Conversely, a tunnelling motion was detected for the H₂S heterodimer, associated to the internal rotation in the triatomic. The experimental moments of inertia provided ground-state effective structures, consistent with the predicted global minima for the two dimers. In particular, SO₂ adopts a stacked position with the sulfur atom closer to the ring, while H₂S is depicted as hydrogen bonded to benzofuran. DFT D3-dispersion-corrected calculations at B3LYP and B2PLYP levels satisfactorily reproduced the observed molecular structures. The relative deviations for the rotational constants in the SO₂ dimer are kept under 2.6% for both B3LYP and B2PLYP (Tables S1–S4, ESI†), despite the larger computational cost of the later method. Small structural variations are also observed for the three used basis sets. However, the stabilization energies are predicted differently, with B3LYP giving complexation energies *ca.* 14–18% larger than B2PLYP for BF \cdots SO₂. The structural predictions are worse for the H₂S dimer, but kept under 4.8% for all six calculations methods. In the H₂S dimer the B3LYP complexation energies are *ca.* 8–13% larger than for B2PLYP. Computational predictions using natural bond orbitals explain the dimer interactions as primarily originated from $S \cdots \pi$ and $S-H \cdots \pi$ interactions. This description, despite an increase in the binding energy, shows large similarities with the prototypic benzene dimers.^{11,12,18} New experiments are in progress to explore the intermolecular interactions between larger sulfur compounds and aromatic systems.

Author contributions

Conceptualization, G. F. and A. L.; methodology, G. F. and A. L.; software, Y. J., W. L., R. T. S., M. J. and C. P.; validation, Y. J., W. L.,

R. T. S., M. J. and C. P.; formal analysis, Y. J., W. L., R. T. S., M. J., C. P., A. L., G. F.; investigation, Y. J., W. L., R. T. S., M. J.; resources, A. L. and G. F.; data curation, A. L. and G. F.; writing—original draft preparation, A. L. and G. F.; writing—review and editing, A. L., and G. F.; visualization, Y. J., W. L.; supervision, A. L. and G. F.; project administration, A. L. and G. F.; funding acquisition, A. L. and G. F.

Conflicts of interest

There are no conflicts to declare.

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