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Angular structure factor of the hexatic-B liquid crystals: bridging theory and experiment†

Ivan A. Zaluzhnyy,^{ab} Ruslan Kurta,^c Michael Sprung,^a Ivan A. Vartanyants^{ad} and Boris I. Ostrovskii^{ef}

We report results of X-ray scattering studies of the angular structure factor of liquid crystal hexatic-B films. According to the sixfold rotational symmetry of the hexatic-B phase, its characteristic scattering splits into six reflections. The shape of the radial and angular cross-sections of these reflections and their temperature evolution are analyzed. We find that over a wide temperature range of the hexatic-B phase existence the angular profiles of the in-plane X-ray scattering are well fitted by the Voigt function, which is a convolution of the Gaussian and Lorentzian functions. This result is supported by the known theoretical considerations of the hexatic structure factor below the smectic-hexatic phase transition temperature. Similar predictions for the angular shape of the hexatic peak in the vicinity of the smectic-hexatic phase transition temperature follow from the multicritical scaling theory of the hexatic-B phase in three dimensions. We find that the specific shape of the hexatic structure factor can be explained by the interplay of two distinct contributions to the free energy of the system, a liquid-like density term and a coupling term between the bond-orientational order and short-range density fluctuations.

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1 Introduction

The hexatic-B (Hex-B) liquid crystal (LC) is a unique stacked three-dimensional (3D) analogue of the notable two-dimensional (2D) hexatic phase that was predicted as an intermediate state in the process of 2D crystal melting.^{1–3} The Hex-B phase is closely related to the well-known smectic-A (Sm-A) phase, therefore LC substances often exhibit a temperature induced Sm-A–Hex-B phase transition. Both the Hex-B and Sm-A phases consist of a stack of parallel equidistant molecular layers, in which elongated molecules are oriented on average along the layer normal, and exhibit the short-range positional correlations within the layers (Fig. 1(a and b)). The Hex-B phase differs from the Sm-A phase by a presence of a long-range bond-orientational (BO) order.^{4–6} The Hex-B phase is a truly 3D

system: the orientation of intermolecular bonds is the same in all layers in the stack, while the shear modulus in the plane of layers is zero, so there are no positional correlations between them (Fig. 1(b)).^{2–5} The BO order in the Hex-B phase is characterized by a sixfold rotational symmetry and can be described by the complex-valued local ordering field $\Psi(\mathbf{r}) = |\Psi(\mathbf{r})|e^{i6\psi(\mathbf{r})}$, where $|\Psi(\mathbf{r})|$ is a magnitude of the two-component BO order parameter and $\psi(\mathbf{r})$ is its phase that corresponds to an angle between the intermolecular bonds and some reference axis. In addition to classical thermotropic LCs, the 3D hexatic phase can be also found in some other systems, such as lipids,^{7,8} DNA solutions,^{9,10} smectic mesophases,¹¹ self-assembled micellar polymers¹² and columnar assemblies of colloidal discs.¹³

The structure of a homogeneous system, such as Sm-A or Hex-B phase, can be described in terms of the pair density correlation function $G(\mathbf{r}) = \langle \delta\rho(\mathbf{r} + \mathbf{R})\delta\rho(\mathbf{R}) \rangle_{\mathbf{R}}$, which gives the probability to find a particle at the separation \mathbf{r} from an arbitrary chosen reference particle positioned at \mathbf{R} .^{14,15} Here $\delta\rho(\mathbf{r})$ is a local deviation of the electron density from the average value, and the averaging $\langle \dots \rangle_{\mathbf{R}}$ is performed over the positions of a reference particle \mathbf{R} . The structure factor $S(\mathbf{q})$ measured in an X-ray scattering experiment is a Fourier transform of the pair density correlation function

$$S(\mathbf{q}) = \int G(\mathbf{r})e^{i\mathbf{q}\cdot\mathbf{r}}d\mathbf{r} = \left\langle |\delta\rho(\mathbf{q})|^2 \right\rangle, \quad (1)$$

where \mathbf{q} is the momentum transfer vector and $\delta\rho(\mathbf{q})$ is a Fourier transform of $\delta\rho(\mathbf{r})$. Hence, the information on the character of

^a Deutsches Elektronen-Synchrotron DESY, Notkestraße 85, 22607, Hamburg, Germany. E-mail: ivan.vartanyants@desy.de

^b Institut für Angewandte Physik, Universität Tübingen, Auf der Morgenstelle 10, 72076, Tübingen, Germany

^c European XFEL, Holzkoppel 4, D-22869 Schenefeld, Germany

^d National Research Nuclear University MEPhI (Moscow Engineering Physics Institute), Kashirskoe Shosse 31, 115409 Moscow, Russia

^e Federal Scientific Research Center “Crystallography and photonics”, Russian Academy of Sciences, Leninskii prospect 59, 119333 Moscow, Russia. E-mail: ostrenator@gmail.com

^f Institute of Solid State Physics, Russian Academy of Sciences, Academician Ossipyan str. 2, 142432 Chernogolovka, Russia

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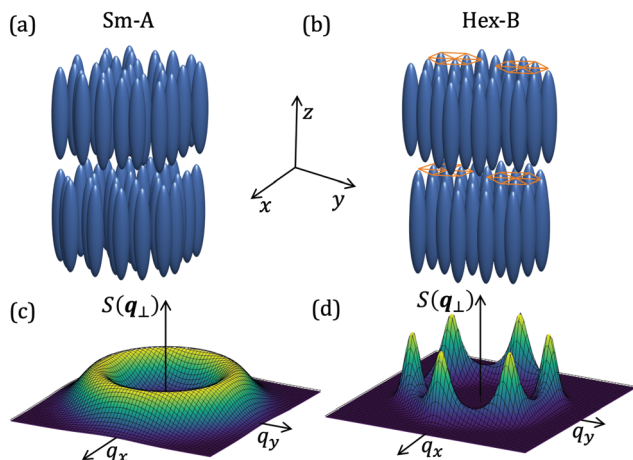


Fig. 1 (a and b) Layered molecular structures of the Sm-A (a) and Hex-B (b) phases. Two molecular layers are shown with the liquid-like short-range positional order within a layer. Orange hexagons indicate the bond-orientational order in the Hex-B phase. (c and d) Corresponding structure factor $S(\mathbf{q}_\perp)$ for the Sm-A (c) and Hex-B (d) phases.

the orientational and positional order, as well as the symmetry of the system can be obtained either directly from $S(\mathbf{q})$ or by reconstructing the pair density correlation function $G(\mathbf{r})$ from the scattering data.^{16,17}

The general feature of all layered (smectic, hexatic) LC phases in reciprocal space is a set of sharp $(00n)$ quasi-Bragg peaks in the direction along the layer normal (q_z) at wave-vector transfer $q_{zn} = 2\pi n/d$ (where d is a layer spacing). The in-plane positional order is short-range and can be determined from the position and width of the broad diffuse peaks in the \mathbf{q}_\perp plane ($\mathbf{q}_\perp = (q_\perp, \varphi)$ in polar coordinates) (Fig. 1(c and d)). In the following we consider the structure factor $S(\mathbf{q}_\perp)$ in the plane of the layers ($q_z \ll 2\pi/d$), which allows one to identify the Hex-B phase and to analyse its structural features.

In our previous studies of the hexatic films, we have obtained numerous data on the structure factor $S(\mathbf{q}_\perp)$ over the range of existence of the Hex-B phase, including the vicinity of the phase transition from a high temperature Sm-A to the Hex-B phase.^{18–20} In this work, we go beyond the standard analysis of the in-plane structure factor of the Hex-B phase and present the X-ray data from thick free-standing films of four different LC materials. We confirm that the radial profiles of the scattering peaks are well approximated by the Lorentzian function as expected for systems with a liquid-like in-plane order.^{16,17} The angular profiles of the scattering peaks in the Hex-B phase are more complex and can be described by a combination of the central Gaussian peak with the extended Lorentzian wings. Our experimental data show that the angular profiles of the hexatic reflections are well fitted by the Voigt profile, which is a convolution of the Gaussian and Lorentzian functions, at all temperatures within the region of the Hex-B phase existence. The relative contribution to the angular scattering profile from both functions depends on the temperature deviation $\Delta T = T - T_C$ from the Sm-A to Hex-B transition point T_C and, accordingly, on the strength of the BO order.

The obtained result is of a fundamental origin and reflects the intrinsic symmetry of the Hex-B structure factor. It agrees very well with the theoretical study of Aeppli and Bruinsma²¹ where they have shown that the angular cross-section of the hexatic structure factor in a 3D hexatic LC can be described by a convolution of the Gaussian and Lorentzian functions (see also²²). Their model is based on the interaction between the BO order parameter $\Psi(\mathbf{r})$ and the short-range density fluctuations $\delta\rho(\mathbf{r})$, introduced by a cross-term of the Landau functional depending on both fields $\Psi(\mathbf{r})$ and $\delta\rho(\mathbf{r})$. Similar predictions for the angular shape of the structure factor follow from the multicritical scaling theory (MCST) developed by Aharony *et al.* for the 3D Hex-B phase.²³ Despite the fact that these theoretical predictions were known for several decades, they have never been thoroughly verified experimentally. To the best of our knowledge the first experimental work in this direction²⁴ was limited by only two LC compounds and the comparison with the theories of Aeppli and Bruinsma²¹ and MCST²³ was lacking. In this work, we analyse in detail the angular shape of the diffraction peaks in the Hex-B phase and compare the experimental data with the existing theoretical models.

2 Bond-orientational order and the hexatic structure factor

The in-plane correlations between the molecules in the Sm-A phase give rise to a broad diffuse scattering ring in the \mathbf{q}_\perp plane, Fig. 2(a). Due to the short-range positional order between the molecules in a smectic layer, the radial cross section of the ring has a Lorentzian profile with a half width at half maximum (HWHM) γ . This allows one to define a characteristic length $\xi = 1/\gamma$ over which positional correlations between the molecules in a layer decay exponentially. The radius of the ring is $q_{\perp 0} \cong 4\pi/a\sqrt{3}$ where a is the average lateral molecular separation. The emergence of the BO order in the Hex-B phase breaks the angular isotropy of the in-plane structure factor and leads to a six-fold modulation of the in-plane

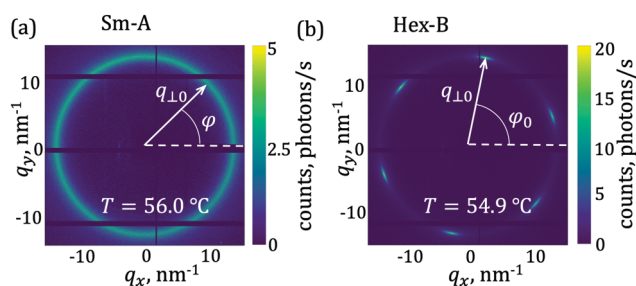


Fig. 2 Experimentally measured X-ray scattered intensity $I(\mathbf{q}_\perp) = I(q, \varphi)$ for the LC compound 54COOBC at temperature $T = 56^\circ\text{C}$ in the Sm-A phase (a) and $T = 54.9^\circ\text{C}$ in the Hex-B phase (b). $q_{\perp 0}$ is the radius of the diffuse scattering ring, φ is an angular variable, φ_0 defines the angular position of the diffraction peaks in the Hex-B phase.



scattering, Fig. 2(b). A Fourier expansion of the azimuthal scattering profile into cosine series gives⁵

$$I(q_{\perp 0}, \varphi) = I_0 \left[1 + 2 \sum_{m=1}^{+\infty} C_{6m} \cos(6m(\varphi - \varphi_0)) \right], \quad (2)$$

where I_0 is the angular-averaged scattered intensity, φ is the angle measured along the arc of the diffraction pattern, and φ_0 corresponds to the orientation of the diffraction pattern relative to the reference axis in the \mathbf{q}_{\perp} -plane. The coefficients C_{6m} are usually considered as the BO order parameters measuring the degree of the BO ordering in the Hex-B phase.^{5,18} In the Sm-A phase, the diffuse ring is uniform leading to all $C_{6m} = 0$ (Fig. 2(a)).

Upon cooling, just below the Sm-A–Hex-B phase transition temperature, only a subtle modulation of the intensity along the scattering ring appears. On further cooling, the scattering ring splits into six individual reflections when the BO order becomes strong enough (Fig. 2(b)). This corresponds to development of the BO order, which manifests itself both in an increase of the number of the nonzero BO order harmonics C_{6m} , and their magnitude.^{5,18–20} These X-ray reflections are asymmetric, each spot is narrow in the radial direction (but still diffuse) and elongated in the azimuthal direction.

Simultaneously with these changes in the azimuthal distribution of the scattered X-ray intensity, the radial width of the scattering peaks also decreases. This indicates a coupling between the positional correlations $G(\mathbf{r})$ defined by the correlation length ξ , and the BO order.^{19,20} The azimuthal scattering profile $I(\mathbf{q}_{\perp 0}, \varphi)$ in eqn (2), in the limit of infinite m , when all coefficients C_{6m} increase to unity, transforms to six sharp peaks in the form of delta-functions, $I \propto \sum_{n=1}^6 \delta(\varphi - \varphi_0 - \pi n/3)$. This

means that at a certain stage the increase of the positional in-plane order might lead to the formation of the 2D lattices that are locked together into the 3D crystalline structures. This should correspond to the formation of a hexagonal crystal which is characterized by a symmetrical sixfold pattern of true Bragg peaks.

As we mentioned above, the X-ray scattering technique measures the Fourier-transformed density–density correlation function $G(\mathbf{r})$ (see eqn (1)). The appearance of the BO order in the system by itself does not change the local density fluctuations responsible for X-ray scattering. The measurements of the evolution of the BO order are possible only due to coupling between the liquid density fluctuations $\delta\rho(\mathbf{r})$ and the hexatic order parameter $\Psi(\mathbf{r})$.^{21,22} This situation is similar to nematic LCs, where the long-range orientational order of the long axis of the molecules couples with the liquid-like density.¹⁶ This leads to the following generic expression for the hexatic structure factor^{21,22,25,26}

$$S(\mathbf{q}_{\perp}) = \left\langle |\delta\rho(\mathbf{q}_{\perp})|^2 \right\rangle \propto \left\langle \frac{k_B T}{A(q_{\perp}) + F(q_{\perp}, |\Psi|, \cos[6(\varphi - \varphi_0)])} \right\rangle. \quad (3)$$

Here $k_B T$ is a thermal energy with k_B being a Boltzmann constant, and $A(q_{\perp}) = \xi^{-2} + (q_{\perp} - q_{\perp 0})^2$ is the contribution of

the short-range positional order into the free energy, where $q_{\perp 0}$ is the preferred in-plane wave vector. The term $A(q_{\perp})$ has a minimum at $q_{\perp} = q_{\perp 0}$, which corresponds to the maximum of the structure factor $S(\mathbf{q}_{\perp})$. The second term in the denominator of eqn (3), $F(q_{\perp}, |\Psi|, \cos[6(\varphi - \varphi_0)])$, depends on the BO order and takes into account the coupling between the BO order and the positional order in the hexatic phase. The average on the right-hand side of eqn (3) is over different configurations of the BO order. In accordance with the symmetry of the Hex-B phase, eqn (3) implies six maxima in the structure factor $S(\mathbf{q}_{\perp})$ at the reciprocal ring $|\mathbf{q}_{\perp}| = q_{\perp 0}$. The expression for $A(q_{\perp})$ suggests the Lorentzian profile of the scattering peak in the radial direction with a HWHM γ inversely proportional to the positional correlation length ($\gamma = 1/\xi$), which is typical for the LC phases with a short range in-plane positional order.^{16,17} The exact functional dependence of the term $F(q_{\perp}, |\Psi|, \cos[6(\varphi - \varphi_0)])$ is unknown, however some predictions regarding the shape of the angular (azimuthal) component of the Hex-B structure factor still can be made.

Aeppli and Bruinsma in their work^{21,22} derived the structure factor of the Hex-B phase using the Landau free energy expansion that combines a liquid-like density term and a term describing the coupling between the BO order and the short-range density fluctuations (see ESI,† Section 1). Using reasonable assumptions, the authors analysed eqn (3) and have shown that the Hex-B structure factor crucially depends on the mean square amplitude of fluctuations of the phase of the BO order parameter, $\langle \delta\psi(\mathbf{r})^2 \rangle$. It was shown that well below the Sm-A–Hex-B phase transition point, where the magnitude of $|\Psi(\mathbf{r})|$ can be considered constant and the fluctuations of the phase $\langle \delta\psi^2 \rangle$ are small, the angular cross-section of the hexatic structure factor can be written as

$$S(q_{\perp 0}, \varphi) \propto \int_{-\infty}^{+\infty} d\theta \frac{e^{-\frac{(\theta - \varphi)^2}{2\langle \delta\psi^2 \rangle}}}{\kappa^2 + \theta^2}, \quad (4)$$

where the value of the parameter κ is defined by the short-range positional correlations within a molecular layer and the coupling between the positional order and the BO order.

Eqn (4) states that the azimuthal cross section of the hexatic peak can be described by a Voigt function

$$V(\varphi; \sigma, \kappa) = G(\varphi; \sigma) * L(\varphi; \kappa), \quad (5)$$

that is a convolution (denoted by asterisk) of a Gaussian function

$$G(\varphi; \sigma) = \frac{1}{\sqrt{2\pi\sigma^2}} e^{-\frac{\varphi^2}{2\sigma^2}} \quad (6)$$

with $\sigma^2 = \langle \delta\psi^2 \rangle$, and a Lorentzian function

$$L(\varphi; \kappa) = \frac{\kappa}{\pi(\varphi^2 + \kappa^2)}. \quad (7)$$

In contrast to many other applications, including X-ray crystallography, where the Voigt function appears as a result of finite resolution of the experimental instruments, here it originates from the specific coupling between the BO order and the short-range density fluctuations in the hexatic phase.



The Gaussian contribution arises from the fluctuations of the BO order and is determined by the parameter $\langle \delta\psi^2 \rangle$, while the Lorentzian counterpart appears due to the coupling between the positional and BO orders and is determined by the value of the parameter κ (see ESI,† Section 1).

Aeppli and Bruinsma²¹ also analyzed the radial cross section of a hexatic structure factor and showed that far away from the Sm-A–Hex-B phase transition region it can be well described by the Lorentzian profile with a HWHM $\gamma = 1/\xi$. However, due to strong fluctuations of the phase of the BO order parameter in the vicinity of the Sm-A–Hex-B phase transition point, the radial intensity profile changes its shape and is better approximated by the square root of a Lorentzian.^{21,27}

Another way of theoretical description of the BO order in hexatic LCs is the multicritical scaling theory (MCST) developed by Aharony *et al.*,²³ and based on the proximity of the Sm-A–Hex-B phase transition to the tricritical point. The authors have shown that the BO parameters C_{6m} defined in eqn (2) satisfy the following scaling law

$$C_{6m} = (C_6)^{\sigma_m}, \quad (8)$$

where

$$\sigma_m = m + \lambda m(m - 1). \quad (9)$$

Theoretical considerations^{23,28,29} suggest that the coefficient λ is equal to $\lambda = 1$ for 2D hexatic films, and $\lambda = 0.3$ for the 3D systems. These predictions are in good agreement with the experimental results.^{5,30–33} In our recent publications we have applied the angular X-ray cross-correlation analysis (AXCCA)^{34–36} to study the BO order in the single-domain Hex-B films. By using this approach, the individual harmonics C_{6m} of the BO order in the Hex-B phase can be directly determined without applying a fitting procedure to the angular intensity distributions. Our results support the validity of the MCST for LCs of various compositions.^{19,20}

It can be shown (see ESI,† Section 2) that eqn (8) and (9) of the MCST already suggest a specific angular shape of the hexatic diffraction peak. It can be demonstrated that these scaling relations for the Fourier components can be only satisfied for a certain type of functions. Indeed, for 2D hexatics, $\lambda = 1$ and $\sigma_m = m^2$ in eqn (9), which allows to rewrite the scaling law (8) as $C_{6m} = e^{-18m^2\sigma^2}$, where the parameter $18\sigma^2 = -\ln C_6$. This expression for the Fourier coefficients immediately implies the Gaussian shape of the scattering peaks in the direction along the ark.^{29,30,37}

In the 3D hexatic phase, $\lambda = 0.3$ that corresponds to a more complex profile of the structure factor in the azimuthal direction. In this case, the scaling law of MCST can be rewritten as a product

$$C_{6m} = (C_6)^{\sigma_m} = e^{-18m^2\sigma^2} e^{-6|m|\kappa}, \quad (10)$$

where the parameters σ and κ depend on the value of the coefficient C_6 in the angular Fourier series and the parameter λ . The first term $\exp(-18m^2\sigma^2)$ in eqn (10) corresponds to the Fourier transform of a Gaussian function, and the second term

$\exp(-6|m|\kappa)$ corresponds to the Fourier transform of a Lorentzian function. In this case, the convolution theorem states that the angular profile of a single hexatic peak is a Voigt function.

3 Experimental results and discussion

The X-ray scattering experiments were carried out at the coherence beamline P10 of the PETRA III synchrotron source at DESY. The details of the X-ray setup, LC materials, sample preparation, and experimental conditions can be found in ref. 18–20 and 38. Upon cooling, all LC materials undergo the Sm-A–Hex-B phase transition. We have studied thick (with the thickness of 1–15 μm) free standing films of four different LC compounds: 75OBC,^{18,20,24} 3(10)OBC,^{19,20,24} PIRO6²⁰ and 54COOBC.³⁸

The free-standing films are especially suitable for an X-ray scattering study: they have a controlled thickness, the films are not influenced by substrate interactions and their free surfaces promote an almost perfect alignment of the smectic layers.^{39,40} The examples of the experimentally measured scattered X-ray intensities for 54COOBC compound are shown in Fig. 2. The azimuthally uniform diffuse ring in Fig. 2(a) corresponds to the liquid-like structure of the smectic layers in Sm-A phase. Upon cooling to the Hex-B phase, the scattering ring splits into six distinct arcs manifesting the appearance of the BO order within the large-scale single hexatic domains (Fig. 2(b)). The observation of the 2D diffraction pattern with six symmetrical spots indicates that in all layers in the stack, the bond angles between the neighbouring molecules are the same (Fig. 1(b)). This means that the average phase of the BO order parameter $\langle \psi(\mathbf{r}) \rangle$ has the same value for all layers in the film. In our X-ray experiments, we have measured the intensity distribution within the single hexatic reflections in the radial and azimuthal directions. The resolution in our experiments was determined by the pixel size of the area detector for the given sample-detector distance. It was high enough ($\Delta q_{\perp}^{\text{res}} < 0.04 \text{ nm}^{-1}$, $\Delta\varphi^{\text{res}} < 6 \times 10^{-4} \text{ deg}$) to resolve the profiles of the narrowest X-ray scattering lines both in the radial and azimuthal directions.^{19,20} Below, we, first, shortly discuss the radial profile of the Hex-B structure factor, and then focus on its angular shape.

3.1 In-plane positional order and the radial intensity distribution

The in-plane positional order in the Sm-A and Hex-B phases of various LC compounds was studied using the shape analysis of the radial cross-sections of the diffraction patterns. The intensity distribution $I(q_{\perp}, \varphi_0) = I(\mathbf{q}_{\perp})$ in the radial direction through the maximum of one of the sixfold diffraction peaks in the Hex-B phase was averaged over different positions in the film and then fitted with the Lorentzian and square root Lorentzian (SRL) functions. In the case of 3(10)OBC and PIRO6 samples the SRL function gives better approximation for $I(\mathbf{q}_{\perp})$ in the temperature interval $-0.8 \text{ }^{\circ}\text{C} \leq \Delta T \leq 0 \text{ }^{\circ}\text{C}$ and $-2.7 \text{ }^{\circ}\text{C} \leq \Delta T \leq 0.3 \text{ }^{\circ}\text{C}$, respectively, which is in the close vicinity of the Sm-A–Hex-B phase transition temperature. Outside of this range the Lorentzian function fits the experimental data



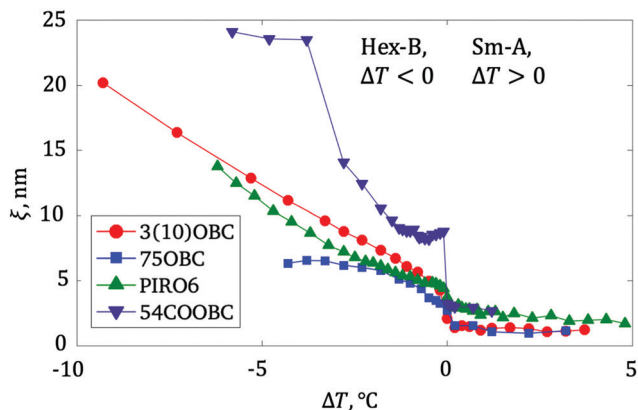


Fig. 3 Temperature dependence of the positional correlation length ξ . The temperature $\Delta T = T - T_C$ is measured relative to the phase transition point.

better.²⁰ For the 75OBC sample, the quality of the data does not allow us to unambiguously determine which function, the Lorentzian or SRL, provides a better fit. In the case of 54COOBC films the situation is even more ambiguous, due to the presence of the two-phase coexistence region near the first-order Sm-A–Hex-B phase transition in this compound.³⁸ For consistency, below we present the results of the fitting with the Lorentzian function for all temperatures.

In Fig. 3 the in-plane positional correlation length $\xi = 1/\gamma$ is shown for four hexatic samples as a function of the relative temperature ΔT . Here γ is the HWHM of the radial cross section of the scattered intensity determined by the fitting with the Lorentzian function. In the Sm-A phase, ξ is about 1–2 nm and does not change with the temperature. On cooling, upon approaching the Hex-B phase the correlation length ξ starts to increase and shows the maximum of the derivative $d\xi/dT$ at the phase transition point T_C (see for additional experimental data ref. 20 and 27). The temperature dependence of the positional correlation length differs from one hexatic compound to another due to differences in material properties and the values of a coupling constant determining the strength of the interaction between the lateral density fluctuations and the BO order. Thus, the correlation length ξ either saturates deeper in the Hex-B phase upon cooling as in 75OBC and 54COOBC, or linearly increases as in 3(10)OBC and PIRO6 hexatic films.^{19,20} Close to the crystallization temperature the in-plane positional correlation length increases by about an order of magnitude. In the crystal phase, the positional order becomes long-range, and the diffraction peaks become so narrow that the analysis of the shape is limited by the resolution of the detector (see, for example, ref. 18).

3.2 Orientational order and the azimuthal intensity distribution

Now we analyze the shape and the width of individual hexatic scattering peaks in the azimuthal direction, $I(q_{\perp 0}, \varphi) = I(\varphi)$. In Fig. 4, examples of the temperature evolution of the angular profiles for all four Hex-B compounds are shown.^{18–20,38} It is

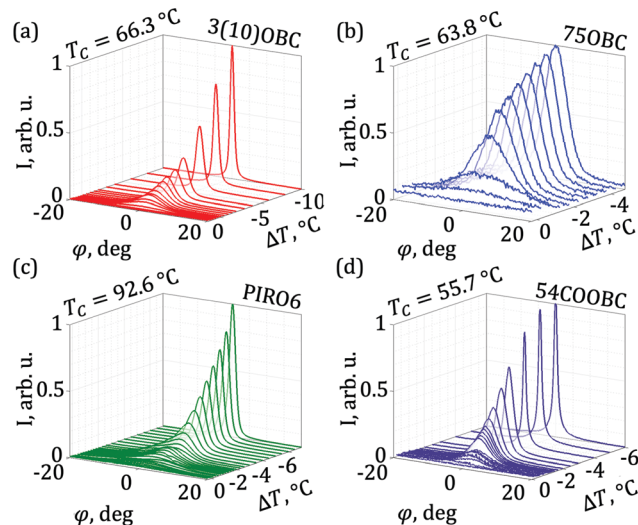


Fig. 4 Azimuthal dependence of the normalized scattered intensity $I(\varphi)$ of a single hexatic peak for four different LC compounds in the vicinity of the Sm-A – Hex-B phase transition and in the Hex-B phase: (a) 3(10)OBC, (b) 75OBC, (c) PIRO6, and (d) 54COOBC. The phase transition temperature T_C is indicated for each compound.

readily seen that close to the phase transition point the peaks are relatively broad, which corresponds to the presence of only few first harmonics in the Fourier expansion (2). Upon the temperature decrease, the peaks become more intense and narrower. However, even far away from the Sm-A–Hex-B transition the well-distinguishable Lorentzian-like wings still persist. The origin of these wings lies in the coupling between the positional and BO order (see eqn (4) and the ESI,† Section 1).

To quantify the temperature evolution of the azimuthal peak profile in various hexatic LCs, we show in Fig. 5 their full width at half maximum (FWHM) as a function of temperature deviation ΔT from the Sm-A–Hex-B phase transition point. The FWHM of the angular profiles decreases with the temperature by about an order of magnitude, and nearly saturates deep in the Hex-B phase. The temperature-induced changes in the positional and BO order, *i.e.* abrupt changes close to the phase

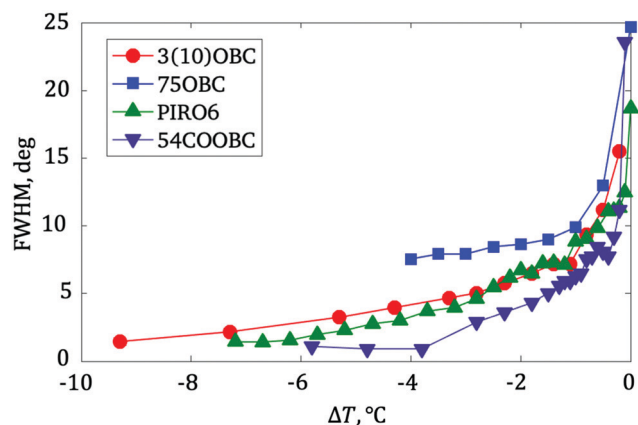


Fig. 5 Temperature dependence of the azimuthal width (FWHM) of a single hexatic peak for different LC compounds.



transition temperature $\Delta T = 0$, and a subsequent gradual evolution deeper in the Hex-B phase, occur almost synchronously as it can be seen in Fig. 3 and 5. It further proves the strong coupling between the short-range positional order and long-range BO order.

In order to get more insight on the temperature dependences of the azimuthal peak shapes in the Hex-B phase $I(\varphi)$, we have approximated these experimentally obtained profiles with three distinct functions, namely the Lorentzian, Gaussian and Voigt functions, defined in eqn (5)–(7). To fit the experimental data, the normalized functions (5)–(7) were scaled by a constant factor C and then a constant background B was added. One can see in Fig. 6, that pure Lorentzian and Gaussian functions do not provide good fitting results of the experimental data. On the other hand, in the whole range of temperatures corresponding to the Hex-B phase the Voigt function provides a good fit of the experimental angular profiles, including the wings of the hexatic peak, which is especially clear at low temperatures, when the hexatic peaks become narrow. The reason for this is more substantial than the mere fact that the Voigt function has one more free fitting parameter compared to the Lorentzian or Gaussian functions, as it was discussed in Section 2. The central part of the azimuthal intensity profile clearly has a Gaussian-like shape, as a result of the small fluctuations of the BO order parameter in 3D hexatic films. The contribution from the Lorentzian function has its origin in the short-range positional order and it is absolutely necessary to describe the long wings of the experimental profiles $I(\varphi)$.

Next, we analyze how the shape of the Voigt function changes with temperature. In Fig. 7 the parameters σ and κ of the Voigt function are plotted for four different LC

compounds. The general trend is clear: the values of σ and κ monotonically decrease on cooling within the Hex-B phase, while the decay law is different for the two parameters. The value of the parameter κ is lower than σ ; it shows a steep decrease at small deviations ΔT from the Sm-A–Hex-B phase transition point, and then stays almost constant at lower temperatures.

To characterize the relative contributions of the Gaussian and Lorentzian functions to the angular profile of the diffraction peak fitted by the Voigt function we have used a dimensionless parameter $\varepsilon = \Delta L / (\Delta L + \Delta G)$, where $\Delta L = 2\kappa$ and $\Delta G = 2\sqrt{2 \ln 2} \sigma$ are the FWHM of the Lorentzian and Gaussian functions, respectively.⁴¹ Parameter ε varies from zero to unity, where $\varepsilon = 0$ corresponds to the pure Gaussian, and $\varepsilon = 1$ corresponds to the pure Lorentzian functions.

The temperature dependence of the parameter ε for all four LC hexatic compounds is shown in Fig. 8. The behavior of the parameter ε is different for various LC compounds. This reflects the difference in material properties and the values of the coupling constants as a function of T determining the strength of the interaction between the short-range density fluctuations and the BO order (see Section 2 and Fig. 3). It is instructive to compare the data from Fig. 3, showing the temperature dependence of the positional correlation length ξ , with the temperature behavior of the parameter ε shown in Fig. 8 for various hexatic materials.

For the compounds 75OBC and PIRO6, the value of ε is slightly below 0.2 close to the Hex-B–Sm-A phase transition point. Such a behavior indicates that at small temperature deviations ΔT from the phase transition temperature the Gaussian contribution to the hexatic angular structure factor dominates. This is due to the fact that the mean square amplitude of fluctuations of the phase $\langle \delta\psi^2 \rangle$ of the BO order parameter is large in the vicinity of T_c , and the wide Gaussian dominates over the Lorentzian function in the convolution integral (eqn (4)). However, starting from the certain temperature deviation ΔT , the compounds 75OBC and PIRO6 behave differently: the value of ε steeply increases for PIRO6, while for 75OBC it remains at the same level. These observations correlate with the data from Fig. 3, where the value of the positional correlation length ξ in PIRO6 continues to grow linearly upon ΔT increase, while for 75OBC it nearly saturates. We attribute the increase of the parameter ε in PIRO6 with the increase of the coupling constant determining the strength of the interaction between the short-range density fluctuations and the BO order that enhances the Lorentzian contribution relative to the Gaussian in the convolution integral.

Different type of behavior is observed for the 3(10)OBC compound (Fig. 8). In this case, the contribution of the Gaussian and Lorentzian profiles in the total structure factor in the vicinity of the phase transition point is nearly equal (parameter ε is only slightly below 0.5) and only deeper in the Hex-B phase the Lorentzian profile slowly starts to dominate. Again, this observation correlates with the linear increase of the positional correlation length ξ with ΔT in 3(10)OBC, Fig. 3. The maturing of the in-plane positional order in the Hex-B phase of 3(10)OBC

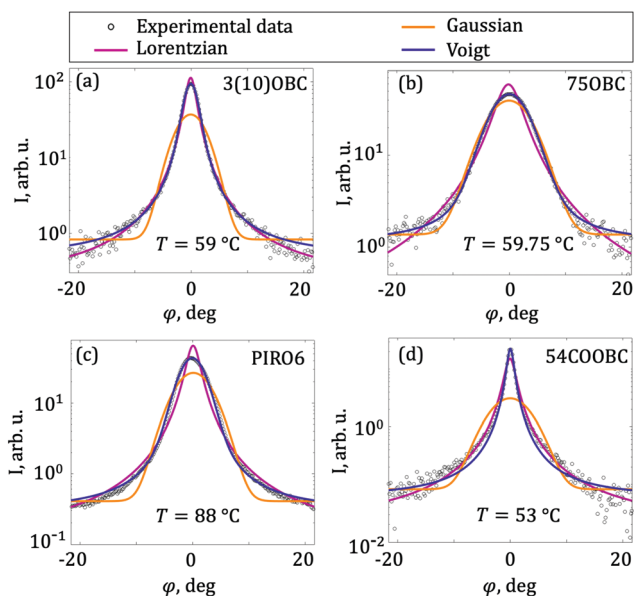


Fig. 6 Examples of the fitting of the angular shape of a single hexatic peak in four different LC compounds with the Lorentzian $I(\varphi) = B + C \cdot L(\varphi; \kappa)$, Gaussian $I(\varphi) = B + C \cdot G(\varphi; \sigma)$, and Voigt $I(\varphi) = B + C \cdot V(\varphi; \sigma, \kappa)$ functions, defined in eqn (5)–(7). The constant background B , the scaling factor C , and the angular variables σ and κ were the free parameters of the fitting.



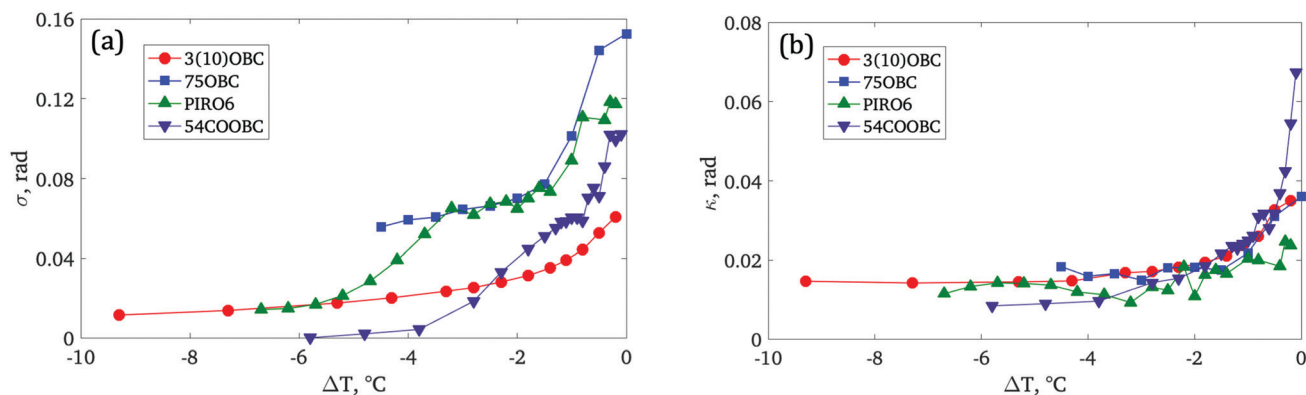


Fig. 7 (a and b) Temperature dependence of the parameters σ (a) and κ (b) for different LC compounds obtained from the fitting of the azimuthal profiles with the Voigt function $V(\varphi; \sigma, \kappa)$.

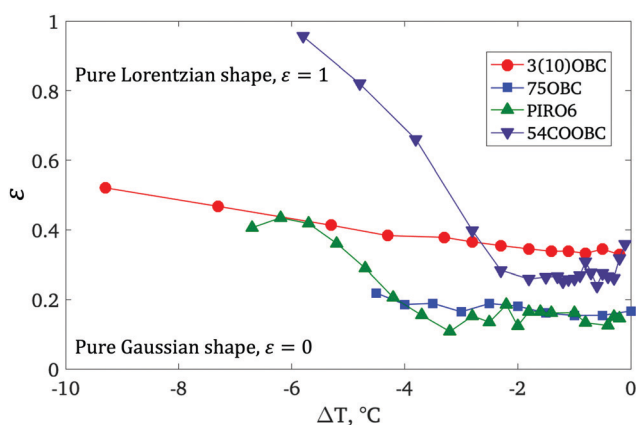


Fig. 8 Temperature dependence of the parameter ϵ obtained from the fitting of the angular profile of a single hexatic peak with the Voigt function for four hexatic LC compounds (eqn (5)–(7)). The value $\epsilon = 0$ corresponds to the pure Gaussian function and $\epsilon = 1$ – to the pure Lorentzian function.

leads to the slight prevailing of the Lorentzian function over the Gaussian.

The temperature behavior of the 54COOBC compound is quite exceptional: in the vicinity of the Sm-A–Hex-B phase transition region, the Gaussian contribution to the angular structure factor clearly dominates. However, upon further cooling into the Hex-B phase the Lorentzian contribution steeply increases and becomes the determining factor in the angular Voigt profile. We attribute such a behavior to the first-order character of the Sm-A–Hex-B phase transition in 54COOBC. The mean square amplitude of fluctuations of the phase of the BO order parameter $\langle \delta\psi^2 \rangle$ rapidly decreases with temperature deviation from T_C , making the Lorentzian contribution to the convolution integral dominating. This conclusion correlates with the steep increase of the positional correlation length ξ in 54COOBC (Fig. 3), which starts nearly at the same temperature deviation, $\Delta T \approx -2$ $^{\circ}\text{C}$, as shown in Fig. 8.

An interesting observation can be made by comparing the values of parameters $\gamma = 1/\xi$ determined from the radial cross section of the hexatic peak (Fig. 3), and κ determined from the

Lorentzian part of the azimuthal cross section of the same peak (see eqn (7) and Fig. 7(b)). Since both parameters are related to the short-range in-plane positional order, one can expect that their ratio is temperature independent. In order to analyze the relationship between the parameters γ and κ , we used the experimental data from Fig. 3 and 7(b) and plotted both parameters (Fig. 9(a)) and the dimensionless ratio $\gamma/(\kappa q_{\perp 0})$ as a function of the temperature deviation ΔT (Fig. 9(b)). It is readily seen from Fig. 9(a) that for three different LC compounds (3(10)OBC, 75OBC and PIRO6) exhibiting a second-order Sm-A–Hex-B phase transition there is a strong correlation between the parameters γ and κ . This is valid for a whole temperature range of the Hex-B phase existence in these compounds. We note that the large values of the parameters γ and κ correspond to the temperature range close to the Sm-A–Hex-B phase transition point, where the hexatic angular and radial profiles are broad. On the other hand, upon approaching the smaller values of these parameters we are going deeper in the hexatic-B phase, which corresponds to the increase of the short-range in-plane positional order. For 54COOBC compound showing the first-order Sm-A–Hex-B phase transition correlations between the parameters γ and κ are observed only far away from the phase transition point T_C (lower left corner of Fig. 9(a)). It is indicative that on approaching the phase transition region any correlations between the parameters γ and κ in this compound disappear.

The tendencies outlined above are also clearly seen in Fig. 9(b) that displays the temperature dependence of the dimensionless ratio $\gamma/(\kappa q_{\perp 0})$. For 3(10)OBC, 75OBC and PIRO6 compounds, the ratio $\gamma/(\kappa q_{\perp 0})$ attains the value of 0.5–0.6 just below the phase transition point and stays nearly constant. Such a behavior is in accordance with the data in Fig. 9(a), indicating that the parameters γ and κ are strongly correlated, *i.e.* their ratio is constant. The temperature behavior of the ratio $\gamma/(\kappa q_{\perp 0})$ for 54COOBC compound is again quite exceptional. In the phase transition region ($\Delta T \geq -2$ $^{\circ}\text{C}$), it increases from the low values to about 0.4–0.5 at which the saturation is almost achieved. Thus, in compliance with the data in Fig. 9(a), there is no correlation between the parameters γ and κ in the vicinity of



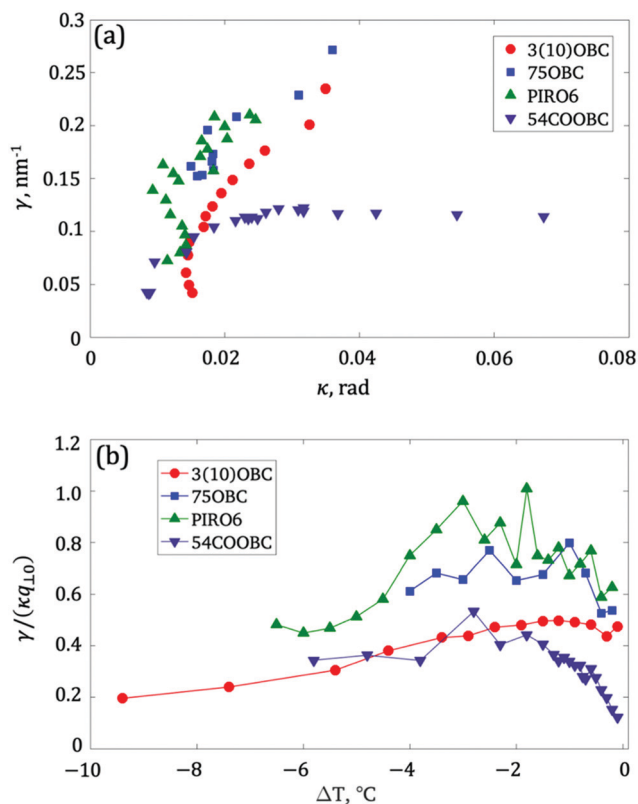


Fig. 9 (a) Correlations between the parameters $\gamma = 1/\xi$, where ξ is the positional correlation length shown in Fig. 3, and the Lorentzian parameter κ (see eqn (7) and Fig. 7(b)). Each point of the plot shows values of the parameters γ and κ at a given temperature ΔT below the Sm-A–Hex-B phase transition. (b) Temperature dependence of the dimensionless ratio $\gamma/(\kappa q_{\perp 0})$. The values of the in-plane wave vector $q_{\perp 0}$ for all four LC compounds are about 14 nm⁻¹ and were taken from ref. 20 and 38.

the Sm-A–Hex-B phase transition for this compound. An important outcome of the data presented in Fig. 9(b) is that all four compounds under study tend to approach nearly the same ratio $\gamma/(\kappa q_{\perp 0}) = 0.4\text{--}0.5$ deep in the Hex-B phase where the short-range positional order is well established. Finally, we would like to point out close similarity between the character of mutual correlations of the parameters γ and κ , as well as the temperature dependence of the ratio $\gamma/(\kappa q_{\perp 0})$ for three different LC compounds (3(10)OBC, 75OBC and PIRO6) showing a second-order Sm-A–Hex-B phase transition. This might be an indication of some universal law governing the coupling between the positional and the BO order in the vicinity of the Sm-A–Hex-B transition of the second order. Such a universal behavior is clearly absent for the case of 54COOBC compound exhibiting the first-order Sm-A–Hex-B phase transition.

4 Conclusions

The most important result of our study is the analysis of the shape of the angular structure factor in the Hex-B phase for different LC compounds. We analysed the results of the synchrotron X-ray scattering measurements of the free-standing LC films showing

the characteristic in-plane pattern with the sixfold reflections from the single hexatic domains. It was found that the angular profiles of the in-plane hexatic scattering are well fitted by the Voigt function that is a convolution of the Gaussian and Lorentzian profiles. The relative contribution to the angular scattering profile from both functions depends on the temperature deviation from the Sm-A–Hex-B transition point.

We showed that two theoretical approaches shed light on the origin of the peculiar Hex-B angular profiles. The first is the multicritical scaling theory, MCST, which works remarkably well in pinning together successive harmonics of the BO order and their temperature evolution. We show that the scaling law for the three-dimensional hexatics indeed implies the Voigt function as an angular shape of the scattering peak. At the same time, Aeppli and Bruinsma have developed the Landau theory of the hexatic order in LCs. It was shown that the unique shape of the hexatic angular structure factor is due to the coexistence of two terms in the free energy of the hexatic phase: a liquid-like density term and a mixed term corresponding to the coupling between the BO order and the short-range density fluctuations. Within this approach it can be shown that the angular cross-section of an individual hexatic peak is also a convolution of the Gaussian and Lorentzian functions. On the basis of this theory, we have explicitly shown that the relative contribution of the Gaussian and Lorentzian profiles in the Hex-B angular structure factor is determined by the interplay between the mean square amplitude of fluctuations of the phase $\langle \delta\psi^2 \rangle$ of the BO order parameter, as well as parameters characterizing the short-range in-plane positional order. The intermediate values of the dimensionless parameter ε , characterizing the ratio between the Gaussian and Lorentzian contributions to the Hex-B angular profile, indicate that the coupling between the BO and the positional order leads to the angular profile of the hexatic peak described by the Voigt function.

Our experimental X-ray data for various hexatic LC compounds are in good agreement with the above theoretical predictions. However, there are new findings in this work, which still wait for an appropriate theoretical interpretation. For three different hexatic compounds possessing a second-order Sm-A–Hex-B phase transition we have observed close similarities in the character of correlations between the parameters γ and κ , as well as in the temperature dependence of the ratio $\gamma/(\kappa q_{\perp 0})$. Both these parameters, γ and κ , characterize the short-range in-plane positional order, however their origin is different. The first one, γ , is determined from the radial scan of the hexatic single peak, while the second, κ , is extracted from the azimuthal cross-section of the same peak. This finding might point out to universal law governing the coupling between the positional and BO order in the vicinity of Sm-A–Hex-B transition of the second order.

Author contributions

I. A. V. and B. I. O. conceived of the presented idea. I. A. Z., R. K., M. S., I. A. V., and B. I. O. performed the X-ray experiments.



I. A. Z. and R. K. processed the experimental data, I. A. Z. conducted the analytical calculations. I. A. Z., I. A. V. and B. I. O. wrote the paper. All authors read and agreed on the final text of the paper.

Conflicts of interest

There are no conflicts to declare.

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