

Cite this: *Chem. Sci.*, 2022, 13, 12560 All publication charges for this article have been paid for by the Royal Society of Chemistry

High spin Fe³⁺-related bonding strength and electron transfer for sensitive and stable SERS detection†

Xinlu Zheng, Xiao Wu, Letian Zhang, Jianjian Kang, Man Zhou, Yang Zhong, Jinlong Zhang  and Lingzhi Wang *

The intrinsic electronic states of transition metal-containing SERS substrates, especially the effect of spin state on the detection sensitivity, still remain unknown. Herein, we propose a simple co-precipitation approach to form trimetallic MIL-101(FeNiTi) with high-spin (HS) Fe³⁺ as a result of geometric distortion of the octahedral symmetry. Using methylene blue as a demonstration, the trimetallic MIL-101(FeNiTi) shows a high enhancement factor (EF) of 6.1×10^6 , a low detection limit of 10^{-9} M and excellent detection stability after long-term preservation. X-ray absorption fine structure and photoelectron spectra demonstrate that coupling between high-spin Fe³⁺ and aliovalent transition metals Ni²⁺ and Ti⁴⁺ with different filling degree of 3d e_g-orbitals results in electron delocalization. The DFT calculation suggests that MIL-101(FeNiTi) with high-spin Fe³⁺ favors molecular adsorption and the charge transfer from the molecule to MIL-101(FeNiTi) is promoted, benefitting from the enhanced electron delocalization, which both contribute to the distinguished SERS performance of MIL-101(FeNiTi). This finding provides in-depth mechanistic understanding of the effect of the spin state of transition metals on mediating SERS activity, which is expected to efficiently promote the development of SERS platforms based on non-noble metals.

Received 18th July 2022
Accepted 5th October 2022

DOI: 10.1039/d2sc03998b

rsc.li/chemical-science

Introduction

Surface enhanced Raman spectroscopy (SERS), as a non-destructive analytical technique with a fingerprint spectrum, has attracted tremendous interest in many research fields.^{1–3} Noble metal substrates (Au, Ag) based on the surface plasmon resonance (SPR) effect have been widely adopted due to the high enhancement efficiency.^{4–8} However, due to the unavoidable limitations of noble metals, including high cost, poor stability and low biocompatibility,^{9–11} increasing attention has been focused on semiconductors or other materials such as metal-organic frameworks (MOFs) and graphene.^{12–15} For these materials, the charge transfer (CT) process between the substrate and adsorbate is well accepted as the dominant SERS mechanism.^{16–19} To date, great progress has been achieved in developing sensitive SERS substrates based on non-noble metals. For example, photonic crystals with unique inverse opal structure and yolk-shell nanomaterials have been

exploited as superior SERS substrates *via* enhancing the light-matter interaction,^{20,21} however they are restricted by the tedious preparation process. Meanwhile, reducible metal oxides with abundant oxygen vacancies (MoO_{3-x}, TiO_{2-x}, WO_{3-x}) have proven efficient in promoting CT between the molecule and substrate due to the formation of extra energy levels and improved resonance absorption.^{22–25} However, surface defects on these materials are unstable under laser irradiation and easy to oxidize even under ambient conditions, which significantly deteriorates the practical SERS performance. Therefore, great effort still needs to be devoted to developing flexible strategies for robust SERS substrates.

Compared with defect engineering, the effect of other intrinsic electronic states of non-noble metal substrates including the orbital filling degree (spin state) and electron localization/delocalization on the SERS performance has been scarcely explored to date. It is known that the bonding strength and electron transfer between the crystal and adsorbates can be tuned through regulating the spin state of transition metal ions in the fields of electrocatalysis and photocatalysis. For example, Zou and coworkers proposed the crystal-size-restricting strategy for finely modulating the high-to-low spin transformation of Fe³⁺, which boosted the bonding strength with adsorbates and facilitated the electron transfer process.²⁶ Xie *et al.* realized delocalized spin states in Ni-based compounds through introducing structural distortion to the confined 2D atomic layers,

Shanghai Engineering Research Center for Multi-Media Environmental Catalysis and Resource Utilization, Key Lab for Advanced Materials and Joint International Research Laboratory of Precision Chemistry and Molecular Engineering, Feringa Nobel Prize Scientist Joint Research Center, Institute of Fine Chemicals, School of Chemistry and Molecular Engineering, East China University of Science & Technology, 130 Meilong Road, Shanghai, 200237, China. E-mail: wlz@ecust.edu.cn

† Electronic supplementary information (ESI) available. See DOI: <https://doi.org/10.1039/d2sc03998b>



leading to modification of the electronic structure and improved electron transport.²⁷ Compared with the commonly observed instability of vacancy defects, the inherent spin characteristics of transition metals are supposed to be more attractive since they simultaneously possess chemical stability under ambient conditions and field-induced reversibility.^{28–31} To date, the correlation between the SERS performance and electronic spin state of transition metal-based materials still remains unknown. There is a lack of efficient strategies to regulate the SERS activity through tuning the spin state.

Herein, we demonstrated that the low-spin to high-spin transition of Fe³⁺ in a series of MIL-101 composites has a significant effect on SERS performance, which was achieved *via* introducing the aliovalent metal dopants Ni²⁺ and Ti⁴⁺ with different electronic configuration. Using methylene blue (MB) as the model molecule, distinguished sensitivity with an enhancement factor (EF) of up to 6.1×10^6 and excellent stability is revealed for trimetallic MIL-101(FeNiTi) with high-spin (HS) Fe³⁺. By combining electronic structure analyses and density functional theory (DFT) calculations, the SERS performance and the spin state of Fe³⁺ were positively correlated. The SERS mechanism was explored regarding the bonding strength and CT between molecules and MIL-101 with different spin states of Fe³⁺. It is expected that the strategy based on spin-regulation would provide a brand-new idea for the development of sensitive non-noble metal SERS technology.

Results and discussion

Sample fabrication and characterization

Tuning of the spin state of MIL-101(Fe) was attempted through adding aliovalent transition metal ions with different electronic configuration and valence state. A series of MIL-101 materials with different ratios of Fe : Ni, Fe : Ti and Fe : Ni : Ti were fabricated through a co-precipitation method (Fig. S1, Table S1†). The X-ray diffraction (XRD) pattern indicates that trimetallic MIL-101(FeNiTi) and bimetallic MIL-101(FeNi) maintain the

characteristic peaks of MIL-101(Fe) (Fig. 1a),^{32,33} suggesting that the introduction of Ni²⁺ and Ti⁴⁺ ions does not alter the crystalline phase although MIL-101(FeNiTi) has a lower crystallinity than the other composites. The scanning electron microscope (SEM) image of trimetallic MIL-101(FeNiTi) presents a uniform shuttle shape (Fig. 1b), which is slightly different from MIL-101(Fe) and bimetallic MIL-101(FeNi) with an octahedral geometry (Fig. S2†).³⁴ Considering the similar particle size of *ca.* 1.2 μm of all three composites, the morphological change should be attributed to the decreased crystallinity of MIL-101(FeNiTi). Moreover, the transmission electron microscope (TEM) mapping of MIL-101(FeNiTi) reveals the uniform distribution of elements Fe, Ni, Ti, C and O (Fig. 1c).

Transmission ⁵⁷Fe Mössbauer spectroscopy was used to characterize the spin states of Fe in three kinds of MIL-101 considering its sensitivity towards the spin states of Fe species.^{35–37} Fig. 2a–c show the ⁵⁷Fe Mössbauer spectra of the three samples, and Table S2† lists the Mössbauer parameters as well as the corresponding areas of different spin states of Fe. For all MIL-101 samples, the spectra can be well fitted into two doublets. Based on the values of isomer shift (IS) and quadrupole splitting (QS), doublets D1 and D2 are assigned to the LS state and the HS state of Fe³⁺, respectively. It can be seen in Fig. 2a that the pristine MIL-101(Fe) mainly exhibits the LS state (the proportion is *ca.* 63.0%) with hyperfine magnetic field of 515.4 kOe, IS of 0.404 mm s⁻¹, and QS of 0.486 mm s⁻¹ (Table S1†). The proportion of the HS state in pristine MIL-101(Fe) is much smaller, and is only 37.0% with IS of 0.414 mm s⁻¹, and QS of 0.901 mm s⁻¹. The above results indicate that most Fe³⁺ in MIL-101(Fe) is in the LS state. But for bimetallic MIL-101(FeNi) (Fig. 2b), it can be clearly observed that the content of the HS state (D2) is increased from 37.0% to 67.9%, indicating that Fe³⁺ is transformed from the LS state to the HS state due to the introduction of Ni²⁺. As presented in Fig. 2c, trimetallic MIL-

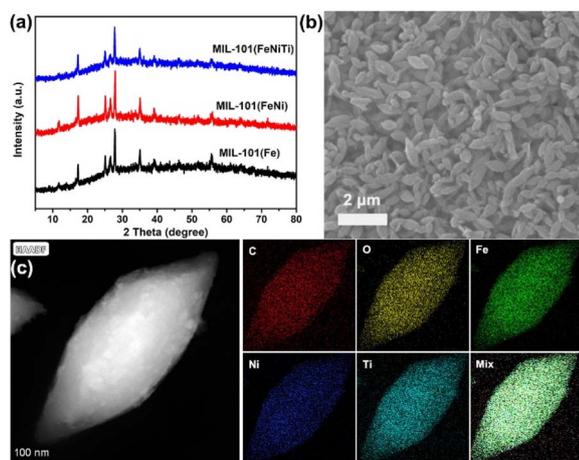


Fig. 1 Crystalline structure and morphology characterization of MIL-101. (a) XRD of three kinds of MIL-101; (b) SEM of trimetallic MIL-101(FeNiTi); (c) TEM-mapping of trimetallic MIL-101(FeNiTi).

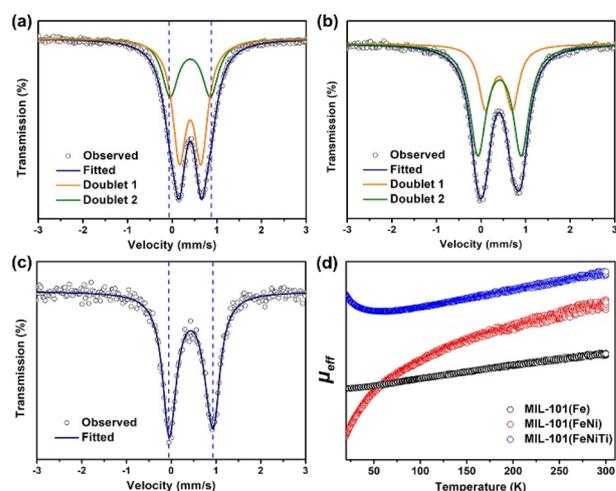


Fig. 2 Spin states of MIL-101(Fe), MIL-101(FeNi) and MIL-101(FeNiTi). ⁵⁷Fe-Mössbauer spectroscopy of (a) MIL-101(Fe); (b) MIL-101(FeNi) and (c) MIL-101(FeNiTi) at room temperature. (d) The comparison of μ_{eff} for the three kinds of MOFs.



101(FeNiTi) formed through the further introduction of Ti^{4+} exhibits the complete conversion of Fe^{3+} to the HS state. Moreover, the HS state of MIL-101(FeTi) (D2, 65.8%) occupies a higher proportion than the LS state (D1, 34.2%), but a lower proportion than the HS state in trimetallic MIL-101(FeNiTi), which suggests that both Ni and Ti elements play an indispensable role in the low-to-high spin transition of Fe^{3+} (Fig. S4†). Overall, 100% ferric iron can be modulated from the LS state to the HS state through the simultaneous introduction of Ni^{2+} and Ti^{4+} into MIL-101(Fe) (Fig. S5†).

To more explicitly demonstrate the spin state transformation of the prepared MIL-101, the 3d orbital electron configurations were described by performing the temperature-dependent magnetic susceptibility (M - T) measurement due to the high sensitivity of magnetism to unpaired electrons.^{38,39} Temperature-dependent paramagnetism can be seen for all MIL-101 materials (Fig. S6†). And the total effective magnetic moment (μ_{eff}), which intimately relates to the number of unpaired 3d electrons (n) of Fe^{3+} , can be further derived from linear fittings of the χ^{-1} - T plots according to the following equation:^{40,41}

$$2.828\sqrt{\chi_m T} = (\mu_{\text{eff}}) = \sqrt{n(n+2)}$$

It is worth noting that the value of μ_{eff} increases with the introduction of Ni and Ti elements (Fig. 2d), which confirms that the number of singly paired 3d electrons of Fe^{3+} in trimetallic MIL-101(FeNiTi) is indeed increased compared with that of the pure MIL-101(Fe) and bimetallic MIL-101(FeNi) counterparts, corresponding to an effectively improved local spin state of Fe^{3+} .

SERS performance of trimetallic MIL-101(FeNiTi)

Fig. 3a shows the SERS spectra of MB molecules adsorbed on MIL-101(Fe), bimetallic MIL-101(FeNi) and trimetallic MIL-101(FeNiTi) substrates under 532 nm irradiation. The strongest peak at 1625 cm^{-1} is attributed to the C-C stretching vibration mode of the thionine ring.⁴² It can be found that the characteristic peaks of MB can be obviously observed on all three substrates, and the SERS activity is positively correlated with the spin state of MIL-101. The EF was quantitatively estimated based on the peak at 1625 cm^{-1} using eqn (S1)–(S3).† The calculated EF of trimetallic MIL-101(FeNiTi) can reach as high as 6.1×10^6 , which is comparable to that of traditional noble metal-based SERS platforms. In addition, a 785 nm laser was further utilized to investigate the SERS performance of the prepared MIL-101 (Fig. S8†). It can be seen that the intensities under 785 nm irradiation are weaker than those under a 532 nm laser but the characteristic peaks assigned to MB are clearly observed. The limit of detection (LOD) was further investigated, as shown in Fig. 3b. It can be seen that the SERS signal of MB can still be detected even when the concentration of MB is decreased to 10^{-9} M , indicating that trimetallic MIL-101(FeNiTi) has excellent SERS detection ability. In addition, the universality of trimetallic MIL-101(FeNiTi) as a SERS substrate was verified for various dye molecules, including rose

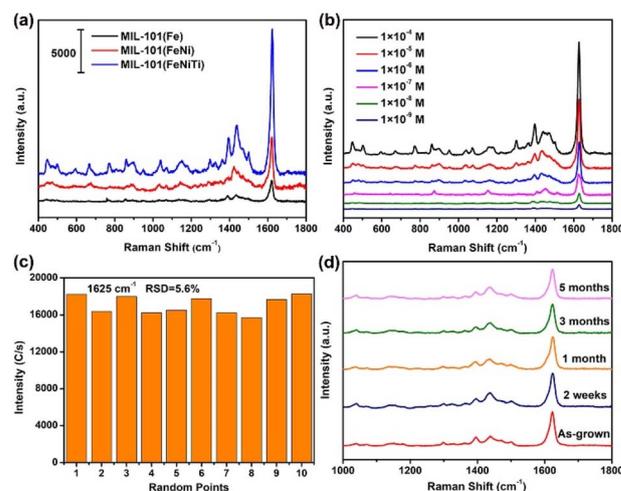


Fig. 3 SERS performance of MIL-101(Fe), MIL-101(FeNi) and MIL-101(FeNiTi). (a) The SERS spectra of MB (10^{-5} M) adsorbed on the three kinds of MOF under 532 nm irradiation; (b) the SERS spectra of MB with different concentrations on trimetallic MIL-101(FeNiTi); (c) the corresponding relative standard deviation (RSD) for the characteristic band (1625 cm^{-1}) based on 10 random positions; (d) the SERS spectra of MB on trimetallic MIL-101(FeNiTi) preserved for different periods. All Raman intensities of detected molecules were obtained by deducting the peaks of the MOF substrates.

bengal (RB, 10^{-5} M), acid blue (AB, 10^{-5} M) and Congo red (CR, 10^{-5} M) (Fig. S9†). It is widely considered that the detection uniformity is a significant criterion for evaluating the practical application performance of SERS-active platforms. Therefore, the SERS signals of MB molecules from 10 randomly selected positions were collected (Fig. S10†). The relative standard deviation (RSD) for the characteristic band (1625 cm^{-1}) was further calculated to quantitatively assess the uniformity of MIL-101(FeNiTi), which was determined to be *ca.* 5.6% (Fig. 3c), demonstrating the ideal detection uniformity of trimetallic MOFs. In addition, the SERS intensities obtained for MIL-101(FeNiTi) kept for 2 weeks, 1 month and 3 months were nearly the same as those of the as-grown trimetallic MIL-101(FeNiTi). Even when the duration was extended to 5 months, the Raman peaks of the MB molecule still did not show any obvious decrease, exhibiting the excellent long-term chemical stability of trimetallic MIL-101(FeNiTi) (Fig. 3d). Considering the easy aggregation of noble metal nanoparticles and the high oxidative susceptibility of defective oxides under ambient conditions, this outstanding stability indicates the robust application potential of MIL-101(FeNiTi) as a SERS substrate.

The origin of high-spin Fe in MIL-101(FeNiTi)

To understand the origin of the high-spin Fe in trimetallic MIL-101(FeNiTi) and its relationship with the SERS performance, X-ray absorption near edge structure (XANES) spectroscopy was utilized to characterize the electronic characteristics of the transition metal atoms. Fig. 4a presents the Fe K-edge XANES of Fe foil, Fe_2O_3 , MIL-101(Fe), bimetallic MIL-101(FeNi) and trimetallic MIL-101(FeNiTi). The pre-edge peak of MIL-101 is



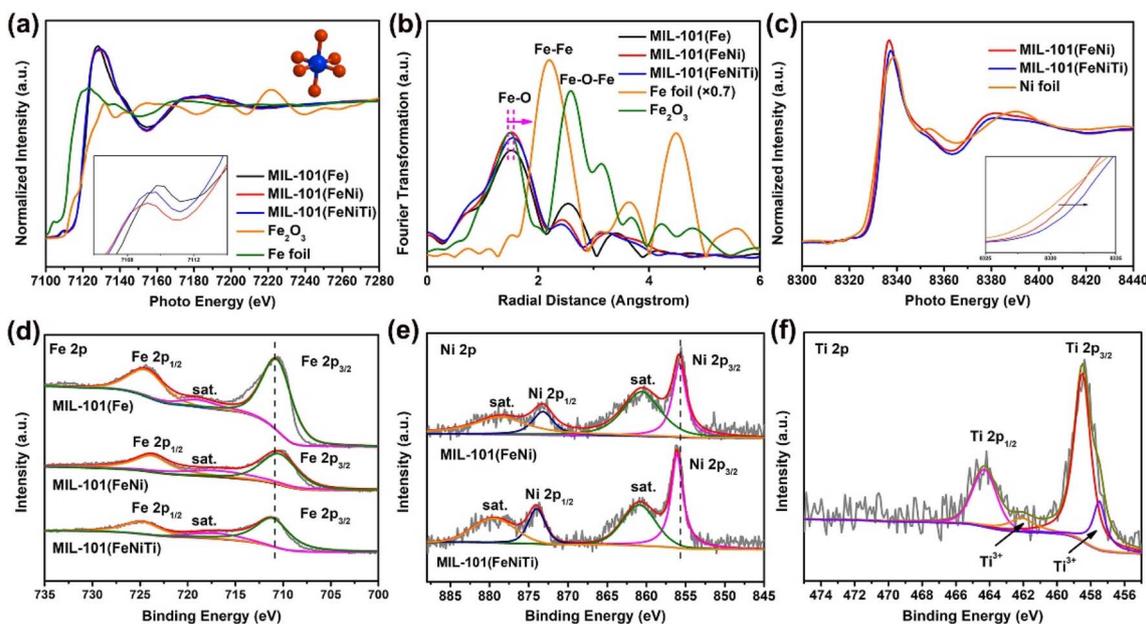


Fig. 4 Electronic structure characterization of MIL-101(Fe), MIL-101(FeNi) and MIL-101(FeNiTi). (a) Fe K-edge XANES spectra of Fe foil, Fe_2O_3 , MIL-101(Fe), bimetallic MIL-101(FeNi) and trimetallic MIL-101(FeNiTi); (b) Fourier transforms of R -space oscillations of Fe for the synthesized MOFs; (c) Ni K-edge XANES spectra of Ni foil, MIL-101(FeNi) and MIL-101(FeNiTi). XPS spectra of (d) Fe 2p, (e) Ni 2p and (f) Ti 2p for the prepared MOFs.

assigned to the 1s to 3d quadrupole transition of the octahedral Fe atom.⁴³ The absorption edges of MIL-101(Fe), bimetallic MIL-101(FeNi) and trimetallic MIL-101(FeNiTi) are closer to that of the Fe_2O_3 reference and far away from that of Fe foil, suggesting that the Fe species in these samples exist in the cationic state (Fig. 4a). Moreover, k^3 -weighted Fourier-transformed extended X-ray absorption fine structure (FT-EXAFS) spectroscopy was conducted to explore the differences in atomic information. Fe K-edge FT-EXAFS spectra show that the Fe–O bond in bimetallic MIL-101(FeNi) and trimetallic MIL-101(FeNiTi) is longer than that in MIL-101(Fe) (Fig. 4b), suggesting that the octahedral geometry of the Fe sites is distorted after the introduction of Ni and Ti elements,⁴⁴ which is also consistent with the XRD result. In addition, the Ni K-edge XANES spectra exhibit a slight shift towards higher energy (Fig. 4c), suggesting that Ni species with a higher valence state are formed due to the introduction of Ti atoms.⁴⁵ The Fourier transforms of R -space oscillations show that only the Ni–O bond appears in MIL-101(FeNi) and MIL-101(FeNiTi) (Fig. S11[†]), implying that Ni exists in modified MIL-101 as single atoms through replacing the Fe atoms.

X-ray photoelectron spectroscopy (XPS) was further applied to identify the electronic characteristics of different MIL-101 composites. In the Fe 2p XPS spectra (Fig. 4d), the signals of $\text{Fe}^{3+} 2p_{3/2}$ and $\text{Fe}^{3+} 2p_{1/2}$ are marked at 711.5 and 724.9 eV, respectively.⁴⁶ Compared to MIL-101(Fe), the characteristic band of Fe^{3+} in MIL-101(FeNi) is slightly shifted towards lower energy, while that of MIL-101(FeNiTi) is shifted towards higher energy, which means that Fe atoms tend to act as electron acceptors and electron donors in MIL-101(FeNi) and MIL-101(FeNiTi), respectively. Additionally, the presence of Ni^{2+} is verified clearly by the high-resolution Ni 2p XPS spectra of MIL-

101(FeNi) and MIL-101(FeNiTi) (Fig. 4e). The characteristic peaks at 855.6 and 872.9 eV and their respective shake-up satellites are attributed to $\text{Ni} 2p_{3/2}$ and $\text{Ni} 2p_{1/2}$, respectively.⁴⁷ It is worth noting that the binding energy of the $\text{Ni} 2p_{3/2}$ peak in MIL-101(FeNiTi) is positively shifted, suggesting that the electrons should be delocalized around the Ni sites after the introduction of the Ti dopant.⁴⁸ In addition, the successful introduction of Ti^{4+} is verified by the Ti 2p data; a signal for Ti^{3+} is also observed in the bands located at 462.6 eV and 457.5 eV (Fig. 4f),⁴⁹ demonstrating that the Ti dopant acts as an electron acceptor in MIL-101(FeNiTi). It is known that the spin state of transition metals is dependent on the electronic configuration and valence state.⁵⁰ Based on the above results, the introduction of Ni and Ti causes distortion of the octahedral symmetric Fe–O sites, which may reduce the crystal-field splitting energy, thus leading to the transition from the low-spin state to the high-spin state of Fe^{3+} .

Furthermore, density functional theory (DFT) calculations were also performed to investigate the spin transition of Fe^{3+} among the three kinds of MIL-101. The three models are displayed in Fig. S12.[†] As disclosed by the density of states (DOS) in Fig. 5a–c, the Fe ions in MIL-101(Fe) and MIL-101(FeNi) have d-band centers located at -0.775 eV and -1.159 eV, respectively, which are higher than for trimetallic MIL-101(FeNiTi) with a d-band center at -1.369 eV. A lower d band center is supposed to be less favorable for molecule adsorption.⁵¹ Moreover, the crystal-field splitting energy of the 3d orbitals was also calculated to understand the spin transition. As shown in Fig. 5d–f, the t_{2g} and e_g centers in MIL-101(Fe) are located at -1.039 and -0.378 eV with crystal-field splitting energy of 0.661 eV, while MIL-101(FeNi) and MIL-101(FeNiTi) have smaller crystal-field



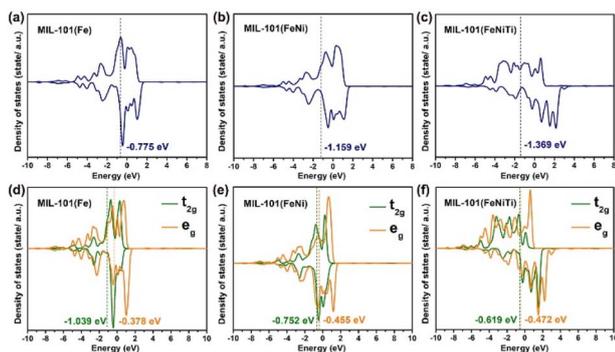


Fig. 5 DFT calculation of the electronic characteristics of different MIL-101 materials. Projected DOS of the Fe 3d orbitals of (a) MIL-101(Fe), (b) MIL-101(FeNi) and (c) MIL-101(FeNiTi). Projected DOS of the t_{2g} and e_g orbitals of the Fe atoms in (d) MIL-101(Fe), (e) MIL-101(FeNi) and (f) MIL-101(FeNiTi).

splittings of 0.297 and 0.147 eV, respectively. According to crystal-field theory,⁵² the Fe 3d electrons in trimetallic MIL-101(FeNiTi) preferentially occupy orbitals individually rather than in pairs, thus resulting in HS Fe atoms. The DFT calculation confirms that the crystal-field splitting energy of Fe atoms can be reduced through the introduction of aliovalent Ni and Ti ions into MIL-101(Fe), thus promoting the low-to-high spin state conversion of Fe^{3+} .

SERS mechanism of MIL-101(FeNiTi)

The possibility of resonance absorption for enhanced SERS activity was first definitively excluded (Fig. S13[†]) since there is no difference among the three kinds of MIL-101 material and no obvious resonance absorption peaks. The spin state of Fe^{3+} may influence the bonding strength between MIL-101 and the molecule since half-occupied (HS) and empty orbitals (LS) have different affinity to electrons.⁵³ In light of this speculation, the adsorption energies of different MOF-MB systems were calculated by DFT calculations. As presented in Fig. 6a–c, the MB molecule is chemically adsorbed on the surfaces of MOF models *via* forming an Fe–S bond with the length in the range of 0.21–0.27 nm. The energies of MB adsorbed over MIL-101(Fe), MIL-101(FeNi) and MIL-101(FeNiTi) are -0.96 , -1.46 and -2.16 eV, respectively (Table S3[†]). It turns out that the bonding strength between MB and MIL-101 follows the order of MIL-101(Fe) < MIL-101(FeNi) < MIL-101(FeNiTi), consistent with the ordering of the HS ratio. This result suggests that the metal-adsorbate interaction should be dependent on the spin state rather than the d-band center of Fe in MIL-101 since MIL-101(FeNiTi) has the lowest d band center but highest spin state.⁵⁴ The enhanced bonding strength should benefit the interfacial electronic interaction between MB and MIL-101,^{55,56} which is critical to the SERS performance of MIL-101.

The charge difference distribution of MIL-101 with adsorbed MB was further analyzed to understand the CT between them. The accumulation region (marked in yellow) for electron density is around the Fe atoms in the MIL-101 substrates, and the depletion region (marked in green) is near the S and C atoms in

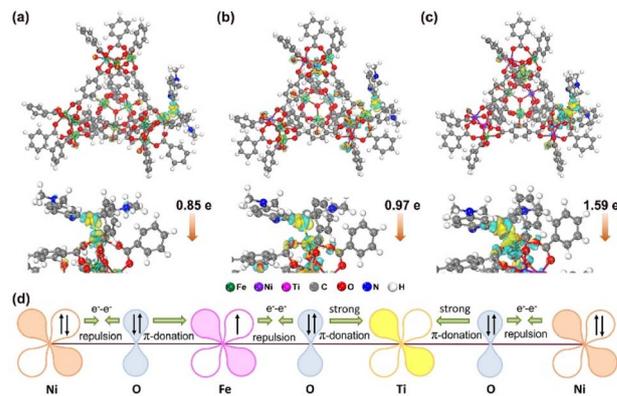


Fig. 6 The DFT calculations of MB@MIL-101(Fe), MB@MIL-101(FeNi) and MB@MIL-101(FeNiTi). The difference in charge distributions for (a) MB@MIL-101(Fe), (b) MB@MIL-101(FeNi) and (c) MB@MIL-101(FeNiTi). The accumulation region for electron density is marked in yellow and the depletion region is marked in green. The colors of Fe, Ni, Ti, C, O, N and H atoms are green, purple, pink, gray, red, blue and white, respectively; (d) schematic representation of the electronic coupling among Fe, Ni and Ti in MIL-101(FeNiTi).

the aromatic ring. According to the Hirshfeld population analysis,⁵⁷ 0.85e and 0.97e are transferred from the MB molecule to MIL-101(Fe) and MIL-101(FeNi), respectively. While for trimetallic MIL-101(FeNiTi), the amount of charge transferred from MB increases to 1.59e. These results suggest that the charge transfer from the MB molecule is most efficiently promoted over trimetallic MIL-101(FeNiTi) with the completely HS state of Fe^{3+} .

To explore the relationship between MIL-101(FeNiTi) containing high-spin Fe^{3+} and CT efficiency, the interaction among different transition metals was further discussed in light of the electronic characteristics of mixed-valent metal ions in MIL-101(FeNiTi). For MIL-101(Fe) with LS Fe^{3+} , the dominant interaction between the fully occupied π -symmetry (t_{2g}) d-orbitals of Fe^{3+} and the bridging oxygen (O^{2-} , transition metal ions are connected through the bridging oxygen bridges) is electron–electron repulsion. For MIL-101(FeNi), Ni^{2+} with paired electrons in π -symmetry (t_{2g}) d-orbitals interacts with O^{2-} *via* electron–electron repulsion. However, because of the electronegativity difference between Fe^{3+} (1.651) and Ni^{2+} (1.367),^{58,59} the electrons would transfer from Ni to Fe sites (Fig. S14[†]), which also agrees well with the XANES and XPS results. The electron transfer would lead to electron delocalization in MIL-101(FeNi).⁶⁰ As for MIL-101(FeNiTi) (Fig. 6d), electrons in Fe^{3+} would be repulsed by O^{2-} , and the paired electrons in O^{2-} interact with Ti^{4+} *via* strong π -donation, resulting in partial electron transfer from Fe^{3+} to Ti^{4+} . Meanwhile, *via* O^{2-} as an intermediary, electrons in Ni^{2+} would transfer to d-orbitals of Ti^{4+} powered by strong π -donation. In this case, the electron transfer among mixed-valence metal ions with different filling degree of e_g orbitals finally leads to electron delocalization in MIL-101(FeNiTi). It is suggested that the electron delocalization caused by the coupling between HS Fe, Ni and Ti ions would improve the conductivity of the SERS substrate,⁶¹ which is verified by the highest photocurrent and the minimum electrochemical impedance (Fig. S15[†]). The



enhanced conductivity is supposed to promote CT.^{62,63} Therefore, benefitting from the formation of high-spin Fe³⁺, the strong bonding and promoted charge transfer between molecules and trimetallic MIL-101(FeNiTi) both contribute to the remarkable SERS activity.

Conclusions

In summary, the low-to-high spin conversion of Fe³⁺ in MIL-101 was tuned through the introduction of aliovalent Ni²⁺ and Ti⁴⁺ as a result of distortion of the octahedral symmetry. Compared with MIL-101(Fe) and bimetallic MIL-101(FeNi), the trimetallic MIL-101(FeNiTi) with completely HS Fe³⁺ exhibited a more distinct SERS activity with EF up to 6.1×10^6 . Meanwhile, a prominent detection stability was revealed during long-term preservation under ambient conditions. The combination of electronic, magnetic and electrochemical analyses with DFT calculations demonstrated that the prominent SERS activity is attributed to the enhanced bonding strength and promoted charge transfer between molecules and MIL-101 with HS Fe³⁺. To our knowledge, to date, this is the first report on the in-depth investigation of the relationship between the spin state of transition metal-based substrates and their SERS activity. This development would open up a new avenue to rational design of distinguished non-noble metal-based SERS substrates.

Data availability

The datasets generated and/or analyzed during the present study are available from the corresponding author upon reasonable request.

Author contributions

Xinlu Zheng designed the experiments, analyzed data, performed theoretical calculations and wrote the manuscript. Xiao Wu and Letian Zhang contributed to sample fabrication and data acquisition. Jianjian Kang, Man Zhou and Yang Zhong assisted in the data interpretation. Jinlong Zhang contributed to the supervision and the theoretical understanding of the paper. Lingzhi Wang contributed to the supervision, the theoretical understanding of the paper and revising the paper.

Conflicts of interest

There are no conflicts to declare.

Acknowledgements

This study was supported by the National Key R&D Program of China (2021YFC2103500), the National Natural Science Foundation of China (No. 21972040, 21822603, and 21773062), Shanghai Municipal Science and Technology (No. 21ZR1417900, 2018SHZDZX03, 17520711500, and 20DZ2250400), the Innovation Program of Shanghai Municipal Education Commission, the Program of Introducing Talents of

Discipline to Universities (No. B16017 and B20031), and the Fundamental Research Funds for the Central Universities.

Notes and references

- 1 T. Chen, H. Wang, G. Chen, Y. Wang, Y. Feng, W. S. Teo, T. Wu and H. Chen, *ACS Nano*, 2010, **4**, 3087–3094.
- 2 Y. Wang, C. Zhao, J. Wang, X. Luo, L. Xie, S. Zhan, J. Kim, X. Wang, X. Liu and Y. Ying, *Sci. Adv.*, 2021, **7**, eabe4553.
- 3 Y. Zeng, J.-Q. Ren, A.-G. Shen and J.-M. Hu, *J. Am. Chem. Soc.*, 2018, **140**, 10649–10652.
- 4 K. Liu, Y. Bai, L. Zhang, Z. Yang, Q. Fan, H. Zheng, Y. Yin and C. Gao, *Nano Lett.*, 2016, **16**, 3675–3681.
- 5 Z. Ye, C. Li, M. Celentano, M. Lindley, T. O'Reilly, A. J. Greer, Y. Huang, C. Hardacre, S. J. Haigh and Y. Xu, *JACS Au*, 2022, **2**, 178–187.
- 6 Z. Yin, Y. Wang, C. Song, L. Zheng, N. Ma, X. Liu, S. Li, L. Lin, M. Li and Y. Xu, *J. Am. Chem. Soc.*, 2018, **140**, 864–867.
- 7 A. Li, X. Wang, L. Guo and S. Li, *J. Phys. Chem. C*, 2019, **123**, 25394–25401.
- 8 X. Wang, G. Ma, A. Li, J. Yu, Z. Yang, J. Lin, A. Li, X. Han and L. Guo, *Chem. Sci.*, 2018, **9**, 4009–4015.
- 9 K. Yuan, Q. Mei, X. Guo, Y. Xu, D. Yang, B. J. Sánchez, B. Sheng, C. Liu, Z. Hu and G. Yu, *Chem. Sci.*, 2018, **9**, 8781–8795.
- 10 W. Lee, B.-H. Kang, H. Yang, M. Park, J. H. Kwak, T. Chung, Y. Jeong, B. K. Kim and K.-H. Jeong, *Nat. Commun.*, 2021, **12**, 1–10.
- 11 X. Qin, Y. Si, D. Wang, Z. Wu, J. Li and Y. Yin, *Anal. Chem.*, 2019, **91**, 4529–4536.
- 12 S. Yang, J. Yao, Y. Quan, M. Hu, R. Su, M. Gao, D. Han and J. Yang, *Light: Sci. Appl.*, 2020, **9**, 1–7.
- 13 L. Yang, Y. Peng, Y. Yang, J. Liu, H. Huang, B. Yu, J. Zhao, Y. Lu, Z. Huang and Z. Li, *Adv. Sci.*, 2019, **6**, 1900310–1900319.
- 14 H. Sun, S. Cong, Z. Zheng, Z. Wang, Z. Chen and Z. Zhao, *J. Am. Chem. Soc.*, 2019, **141**, 870–878.
- 15 E. Feng, T. Zheng, X. He, J. Chen and Y. Tian, *Sci. Adv.*, 2018, **4**, eaau3494.
- 16 B. Yang, S. Jin, S. Guo, Y. Park, L. Chen, B. Zhao and Y. M. Jung, *ACS Omega*, 2019, **4**, 20101–20108.
- 17 J. Lin, W. Ren, A. Li, C. Yao, T. Chen, X. Ma, X. Wang and A. Wu, *ACS Appl. Mater. Interfaces*, 2020, **12**, 4204–4211.
- 18 J. Lin, J. Yu, O. Akakuru, X. Wang, B. Yuan, T. Chen, L. Guo and A. Wu, *Chem. Sci.*, 2020, **11**, 9414–9420.
- 19 V. Rajput, R. K. Gupta and J. Prakash, *J. Mater. Chem. C*, 2022, **10**, 73–95.
- 20 D. Qi, L. Lu, L. Wang and J. Zhang, *J. Am. Chem. Soc.*, 2014, **136**, 9886–9889.
- 21 X. Zheng, W. Zhang, J. Zhang and L. Wang, *Appl. Surf. Sci.*, 2021, **541**, 148422–148430.
- 22 X. Tan, L. Wang, C. Cheng, X. Yan, B. Shen and J. Zhang, *Chem. Commun.*, 2016, **52**, 2893–2896.
- 23 J. Lin, Y. Shang, X. Li, J. Yu, X. Wang and L. Guo, *Adv. Mater.*, 2017, **29**, 1604797–1604806.
- 24 S. Cong, Y. Yuan, Z. Chen, J. Hou, M. Yang, Y. Su, Y. Zhang, L. Li, Q. Li and F. Geng, *Nat. Commun.*, 2015, **6**, 1–7.



- 25 C. Gu, D. Li, S. Zeng, T. Jiang, X. Shen and H. Zhang, *Nanoscale*, 2021, **13**, 5620–5651.
- 26 G. Shen, L. Pan, R. Zhang, S. Sun, F. Hou, X. Zhang and J. J. Zou, *Adv. Mater.*, 2020, **32**, 1905988–1905997.
- 27 S. Chen, Z. Kang, X. Hu, X. Zhang, H. Wang, J. Xie, X. Zheng, W. Yan, B. Pan and Y. Xie, *Adv. Mater.*, 2017, **29**, 1701687–1701694.
- 28 R. Gao, J. Wang, Z.-F. Huang, R. Zhang, W. Wang, L. Pan, J. Zhang, W. Zhu, X. Zhang and C. Shi, *Nat. Energy*, 2021, **6**, 614–623.
- 29 J. Jiang, F. Sun, S. Zhou, W. Hu, H. Zhang, J. Dong, Z. Jiang, J. Zhao, J. Li and W. Yan, *Nat. Commun.*, 2018, **9**, 1–12.
- 30 J. Huang, J. Chen, T. Yao, J. He, S. Jiang, Z. Sun, Q. Liu, W. Cheng, F. Hu and Y. Jiang, *Angew. Chem.*, 2015, **127**, 8846–8851.
- 31 Z. Lu, H. Wang, D. Kong, K. Yan, P.-C. Hsu, G. Zheng, H. Yao, Z. Liang, X. Sun and Y. Cui, *Nat. Commun.*, 2014, **5**, 1–7.
- 32 K. M. Taylor-Pashow, J. Della Rocca, Z. Xie, S. Tran and W. Lin, *J. Am. Chem. Soc.*, 2009, **131**, 14261–14263.
- 33 F. Guo, S. Yang, Y. Liu, P. Wang, J. Huang and W.-Y. Sun, *ACS Catal.*, 2019, **9**, 8464–8470.
- 34 P. Xiong, H. Zhang, G. Li, C. Liao and G. Jiang, *Sci. Total Environ.*, 2021, **797**, 149179–149186.
- 35 U. I. Kramm, L. Ni and S. Wagner, *Adv. Mater.*, 2019, **31**, 1805623–1805629.
- 36 M. L. Aubrey, B. M. Wiers, S. C. Andrews, T. Sakurai, S. E. Reyes-Lillo, S. M. Hamed, C.-J. Yu, L. E. Darago, J. A. Mason and J.-O. Baeg, *Nat. Mater.*, 2018, **17**, 625–632.
- 37 M. Sathiya, G. Rousse, K. Ramesha, C. Laisa, H. Vezin, M. T. Sougrati, M.-L. Doublet, D. Foix, D. Gonbeau and W. Walker, *Nat. Mater.*, 2013, **12**, 827–835.
- 38 S. Zhou, X. Miao, X. Zhao, C. Ma, Y. Qiu, Z. Hu, J. Zhao, L. Shi and J. Zeng, *Nat. Commun.*, 2016, **7**, 1–7.
- 39 L. Zhao, Y.-S. Meng, Q. Liu, O. Sato, Q. Shi, H. Oshio and T. Liu, *Nat. Chem.*, 2021, 1–7.
- 40 Z. Li, Z. Zhuang, F. Lv, H. Zhu, L. Zhou, M. Luo, J. Zhu, Z. Lang, S. Feng and W. Chen, *Adv. Mater.*, 2018, **30**, 1803220–1803227.
- 41 X. T. Wang, T. Ouyang, L. Wang, J. H. Zhong, T. Ma and Z. Q. Liu, *Angew. Chem.*, 2019, **131**, 13425–13430.
- 42 G. Song, W. Gong, S. Cong and Z. Zhao, *Angew. Chem., Int. Ed.*, 2021, **60**, 5505–5511.
- 43 T. E. Westre, P. Kennepohl, J. G. DeWitt, B. Hedman, K. O. Hodgson and E. I. Solomon, *J. Am. Chem. Soc.*, 1997, **119**, 6297–6314.
- 44 J. Kim, S. J. Kim, E. Jung, D. H. Mok, V. K. Paidi, J. Lee, H. S. Lee, Y. Jeoun, W. Ko and H. Shin, *Adv. Func. Mater.*, 2022, 2110857–2110865.
- 45 J.-Y. Zhang, Y. Yan, B. Mei, R. Qi, T. He, Z. Wang, W. Fang, S. Zaman, Y. Su and S. Ding, *Energy Environ. Sci.*, 2021, **14**, 365–373.
- 46 K. Fan, Y. Ji, H. Zou, J. Zhang, B. Zhu, H. Chen, Q. Daniel, Y. Luo, J. Yu and L. Sun, *Angew. Chem., Int. Ed.*, 2017, **56**, 3289–3293.
- 47 K. Fan, H. Chen, Y. Ji, H. Huang, P. M. Claesson, Q. Daniel, B. Philippe, H. Rensmo, F. Li and Y. Luo, *Nat. Commun.*, 2016, **7**, 1–9.
- 48 K. Chen, M. Cao, G. Ni, S. Chen, H. Liao, L. Zhu, H. Li, J. Fu, J. Hu and E. Cortés, *Appl. Catal., B*, 2022, 121093–121102.
- 49 W. Zhao, K. Zhang, L. Wu, Q. Wang, D. Shang and Q. Zhong, *J. Colloid Interf. Sci.*, 2021, **581**, 76–83.
- 50 Y. Tong, J. Wu, P. Chen, H. Liu, W. Chu, C. Wu and Y. Xie, *J. Am. Chem. Soc.*, 2018, **140**, 11165–11169.
- 51 A. Vojvodic and J. K. Nørskov, *Science*, 2011, **334**, 1355–1356.
- 52 R. G. Burns and R. G. Burns, *Mineralogical applications of crystal field theory*, Cambridge University Press, 1993.
- 53 J. C. Rienstra-Kiracofe, G. S. Tschumper, H. F. Schaefer, S. Nandi and G. B. Ellison, *Chem. Rev.*, 2002, **102**, 231–282.
- 54 S. Bhattacharjee, U. V. Waghmare and S.-C. Lee, *Sci. Rep.*, 2016, **6**, 1–10.
- 55 L. Tao, K. Chen, Z. Chen, C. Cong, C. Qiu, J. Chen, X. Wang, H. Chen, T. Yu and W. Xie, *J. Am. Chem. Soc.*, 2018, **140**, 8696–8704.
- 56 J. Seo, Y. Kim, J. Lee, E. Son, M.-H. Jung, Y.-M. Kim, H. Y. Jeong, G. Lee and H. Park, *J. Mater. Chem. A*, 2022, **10**, 13298–13304.
- 57 F. L. Hirshfeld, *Theor. Chim. Acta*, 1977, **44**, 129–138.
- 58 M. Ying, R. Tang, W. Yang, W. Liang, G. Yang, H. Pan, X. Liao and J. Huang, *ACS Appl. Nano Mater.*, 2021, **4**, 1967–1975.
- 59 H. Liu, M. Zha, Z. Liu, J. Tian, G. Hu and L. Feng, *Chem. Commun.*, 2020, **56**, 7889–7892.
- 60 H. Li, J. Wang, R. Qi, Y. Hu, J. Zhang, H. Zhao, J. Zhang and Y. Zhao, *Appl. Catal., B*, 2021, **285**, 119778–1197787.
- 61 L. Yang and M. Dincă, *Angew. Chem., Int. Ed.*, 2021, **60**, 23784–23789.
- 62 R. Wang, C. Chen, Y. Zheng, H. Wang, J.-W. Liu and S.-H. Yu, *Mater. Chem. Front.*, 2020, **4**, 2881–2903.
- 63 M. Pradhan, A. K. Sinha and T. Pal, *Chem.–Eur. J.*, 2014, **20**, 9111–9119.

