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## One-electron bonds in copper–aluminum and copper–gallium complexes†

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Odd-electron bonds have unique electronic structures and are often encountered as transiently stable, homonuclear species. In this study, a pair of copper complexes supported by Group 13 metalloligands,  $M[N((o-C_6H_4)NCH_2P^iPr_2)_3]$  ( $M = Al$  or  $Ga$ ), featuring two-center/one-electron (2c/1e)  $\sigma$ -bonds were synthesized by one-electron reduction of the corresponding  $Cu(I) \rightarrow M(III)$  counterparts. The copper bimetallic complexes were investigated by X-ray diffraction, cyclic voltammetry, electron paramagnetic spectroscopy, and density functional theory calculations. The combined experimental and theoretical data corroborate that the unpaired spin is delocalized across Cu, M, and ancillary atoms, and the singly occupied molecular orbital (SOMO) corresponds to a  $\sigma$ -(Cu–M) bond involving the Cu 4p<sub>z</sub> and M *ns*/*np*<sub>z</sub> atomic orbitals. Collectively, the data suggest the covalent nature of these interactions, which represent the first examples of odd-electron  $\sigma$ -bonds for the heavier Group 13 elements Al and Ga.

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Odd-electron  $\sigma$ -bonds, where the electrons are delocalized between two atoms, can occur as two-center/one-electron (2c/1e) or two-center/three-electron (2c/3e) interactions. Proposed by Pauling in 1931,<sup>1</sup> odd-electron  $\sigma$ -bonds have garnered attention because of their fundamental importance to chemical bonding and their relationship to radical species generated during oxidative stress in biological systems.<sup>2–14</sup> Examples of compounds exhibiting odd-electron bonding are typically homonuclear (like  $H_2^+$ ,  $He_2^+$ , and alkali metal dimers) and transiently stable, limiting them to spectroscopic characterization.<sup>1,11,15–18</sup>

The first solid-state structure of a formally one-electron  $\sigma$ -bond was a tetraphosphabenzene species (Fig. 1a) which was formed by the coupling of two diphosphirenyl radicals.<sup>19</sup> Following this discovery, the formation of discrete 2c/1e  $\sigma$ -bonds, where the odd-electron is delocalized between two homonuclear main group centers, was reported for B–B and then extended to P–P.<sup>8,17,20</sup> Of note, the first solid-state structure

of a B–B compound was reported in only 2014 (Fig. 1b).<sup>21</sup> Examples of 2c/1e  $\sigma$ -bonds between the heavier Group 13 congeners are even more lacking because of the greater propensity for their unpaired spins to couple, forming larger more stable clusters.<sup>8</sup> To our knowledge, there are only three structurally characterized examples of odd-electron bonds for the heavy Group 13 atoms,<sup>22</sup> and these examples are all homonuclear  $\pi$ -radicals (Fig. 1c).<sup>23–26</sup>

Heteronuclear odd-electron  $\sigma$ -bonds are also rare. The Cu(TPB) complex, where TPB is a trisphosphinoborane, is the single structural example of a 2c/1e bond between heteroatoms (Fig. 1d).<sup>27</sup> The authors described the bonding as Cu–B, where the unpaired electron is heavily polarized toward B. A theoretical study predicted that such a bond would also exist between Cu and Al, but no heavier analogues of Cu(TPB) have been

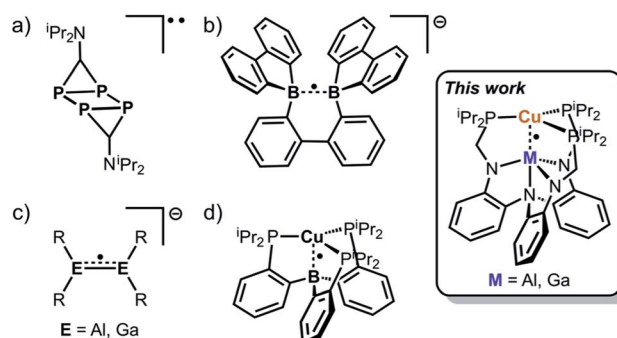


Fig. 1 Select examples of structurally characterized molecules (a–d) featuring odd-electron bonds.

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synthesized to date.<sup>28</sup> Furthermore, the heavier Group 13 elements by virtue of their lower electronegativity compared to B should facilitate greater covalent interactions with the Cu center.

Hence, we sought to target formally zerovalent Cu complexes supported by Al(III) or Ga(III) as an extension of the previously reported isoelectronic nickelate species and Cu(TPB).<sup>29</sup> Herein, we describe the synthesis, structure, spectroscopic characterization, and DFT calculations of cationic  $[\text{CuML}]^+$  complexes ( $\text{L} = [\text{N}((\text{o}-\text{C}_6\text{H}_4)\text{NCH}_2\text{P}^i\text{Pr}_2)_3]^{3-}$ ;  $\text{M} = \text{Al}$  and  $\text{Ga}$ ) as well as their one-electron reduced metalloradical counterparts that feature discrete  $2\text{c}/1\text{e}$  bonds.

## Results and discussion

### Synthesis and cyclic voltammetry

Initial synthetic attempts to produce cationic  $[\text{CuML}]^+$  complexes *via* metalation of ML with a copper halide followed by halide abstraction produced intractable mixtures. Hence, we sought an alternative Cu(I) precursor that eschews the halide and has labile ligands, such as  $[\text{Cu}(\text{COD})_2]\text{A}$  ( $\text{COD} = 1,5$ -cyclooctadiene and  $\text{A} = \text{counteranion}$ ). This cation has received little attention as a precursor to low-coordinate copper complexes, which may stem from the known salts ( $\text{A} = \text{ClO}_4^-$ ,  $\text{NO}_3^-$ ,  $\text{SbF}_6^-$ , *etc.*) lacking sufficient solubility in weakly coordinating solvents.<sup>30–32</sup> Because the  $\text{BAR}_4^{\text{F}}$  counteranion ( $\text{BAR}_4^{\text{F}} = \text{tetrakis}[3,5\text{-bis(trifluoromethyl)-phenyl}]\text{borate}$ ) confers high solubility in low dielectric solvents, we targeted and isolated  $[\text{Cu}(\text{COD})_2]\text{BAR}_4^{\text{F}}$  by protonating the pentameric  $(\text{CuMes})_5$  (ref. 33 and 34) ( $\text{Mes} = 2,4,6\text{-trimethylphenyl}$ ) with  $[(\text{H}(\text{OEt})_2)_2\text{BAR}_4^{\text{F}}]$  (ref. 35) in the presence of excess COD (see ESI for details, Fig. S1†). Gratifyingly,  $[\text{Cu}(\text{COD})_2]\text{BAR}_4^{\text{F}}$  reacted with the metalloligands ALL or GAL in PhF to provide  $[1]\text{BAR}_4^{\text{F}}$  or  $[2]\text{BAR}_4^{\text{F}}$  in good yields as yellow or orange powders, respectively (Scheme 1, Fig. S3–S22†).

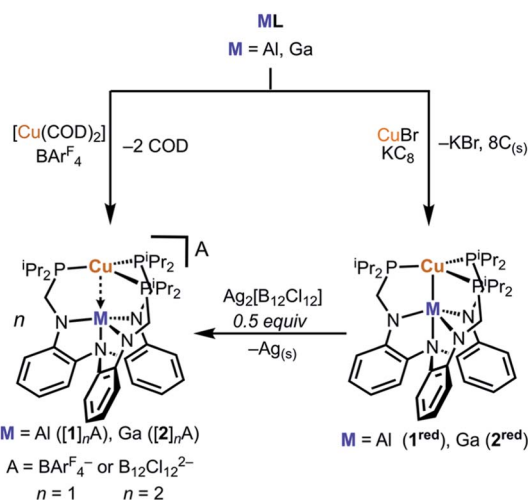
Cyclic voltammetry (CV) studies of both  $[1]\text{BAR}_4^{\text{F}}$  and  $[2]\text{BAR}_4^{\text{F}}$  suggested that  $\{\text{CuM}\}^{11}$  complexes are accessible at low

potentials (Fig. S31–36†). The CV of  $[1]\text{BAR}_4^{\text{F}}$  reveals a single reduction event at  $-2.03\text{ V}$  vs.  $[\text{FcCp}_2]^{0/+}$ , which corresponds to the  $[\text{CuALL}]^{+/0}$  redox couple. The CV of  $[2]\text{BAR}_4^{\text{F}}$  shows two single-electron reductions at  $-1.79\text{ V}$  and  $-2.42\text{ V}$  vs.  $[\text{FcCp}_2]^{0/+}$ , which correspond to the  $[\text{CuGAL}]^{+/0}$  and  $[\text{CuGAL}]^{0/-}$  redox couples, respectively. Of note,  $\text{Cu}(\text{TBP})$  was described to exhibit a similar CV profile to  $2^{\text{red}}$ .<sup>27</sup> Unfortunately, the second redox event for  $[1]\text{BAR}_4^{\text{F}}$  was not observed.<sup>36</sup> Scan rate studies indicate that the  $[\text{CuML}]^{+/0}$  redox couple is reversible for both bimetallics species, while the  $[\text{CuGAL}]^{0/-}$  redox process observed for  $[2]\text{BAR}_4^{\text{F}}$  is quasi-reversible (Fig. S32 and S34–S35†). For comparison, the Group 13 metalloligands and the monocopper complex,  $[\text{Cu}(\text{LH}_3)]\text{BAR}_4^{\text{F}}$ , exhibit no reversible redox processes in a similar potential window (Fig. S23–S30 and S37†), which suggests that the Cu–M unit may be necessary for accessing the reduced Cu states. Furthermore, the formal substitution of Al with Ga shifts the  $[\text{CuML}]^{+/0}$  redox couple by more than 200 mV. With this electrochemical insight, we sought to prepare the reduced  $\{\text{CuM}\}^{11}$  species. Chemical reduction of  $[1]\text{BAR}_4^{\text{F}}$  or  $[2]\text{BAR}_4^{\text{F}}$  was accomplished using  $\text{KC}_8$  to generate  $\text{CuALL}$  ( $1^{\text{red}}$ ) and  $\text{CuGAL}$  ( $2^{\text{red}}$ ), respectively. However, high-yielding syntheses of  $1^{\text{red}}$  and  $2^{\text{red}}$  were achieved by first metalating ML with CuBr to generate an *in situ* Cu(I) halide complex followed by addition of  $\text{KC}_8$  in THF at  $-78^\circ\text{C}$  (Scheme 1). The resultant reductions produced deeply colored maroon ( $1^{\text{red}}$ ) and dark red ( $2^{\text{red}}$ ) solutions, from which the products were purified by removing all volatiles and then extracting into benzene.

### X-ray crystallography

While  $[1]\text{BAR}_4^{\text{F}}$  and  $[2]\text{BAR}_4^{\text{F}}$  could be isolated, diffraction-quality crystals of these species were not readily obtained. To this end, we sought to incorporate the highly crystalline dodecachlorododecaborate dianion ( $[\text{B}_{12}\text{Cl}_{12}]^{2-}$ ). X-ray quality crystals of  $[1]_2[\text{B}_{12}\text{Cl}_{12}]$  and  $[2]_2[\text{B}_{12}\text{Cl}_{12}]$  were achieved by the oxidation of  $1^{\text{red}}$  and  $2^{\text{red}}$ , respectively, with  $\text{Ag}_2[\text{B}_{12}\text{Cl}_{12}]$ .<sup>37–39</sup> Crystals of the neutral copper species were grown from saturated  $\text{Et}_2\text{O}$  at  $-30^\circ\text{C}$  ( $1^{\text{red}}$ ) or toluene/pentane vapor diffusion at room temperature ( $2^{\text{red}}$ ). The solid-state structures of the four copper–Group 13 complexes are shown in Fig. 2, and the structural metrics are summarized in Table 1.

The structures of  $[1]_2[\text{B}_{12}\text{Cl}_{12}]$  and  $[2]_2[\text{B}_{12}\text{Cl}_{12}]$  reveal weak to nonexistent Cu(I)–M interactions as their Cu–M distances of 2.6239(8) and 2.5737(5) Å, respectively, are longer than the sum of the metals' covalent radii (*c.f.* 2.53 and 2.54 Å, respectively).<sup>40,41</sup> The formal shortness ratio ( $r$ ),<sup>42</sup> which is defined as the ratio of the metal–metal bond length to the sum of the single-bond covalent radii<sup>41</sup> was also calculated: 1.04 for  $[1]_2[\text{B}_{12}\text{Cl}_{12}]$  and 1.01 for  $[2]_2[\text{B}_{12}\text{Cl}_{12}]$ . The strength of the copper–Group 13 donor–acceptor interaction in  $[1]_2[\text{B}_{12}\text{Cl}_{12}]$  and  $[2]_2[\text{B}_{12}\text{Cl}_{12}]$  is much weaker compared to the corresponding isoelectronic neutral  $\text{Ni}(0)$ <sup>43,44</sup> and anionic  $\text{Co}(-)$ <sup>45</sup> bimetallics, which have  $r$  values  $<1$ , because the localized cationic change on the Cu(I) center should render it a weaker donor. The Cu center is nearly trigonal planar with minimal perturbation due to the Group 13 ion as both  $\sum(\angle\text{P}-\text{Cu}-\text{P})$  and  $\sum(\angle\text{N}_{\text{eq}}-\text{M}-\text{N}_{\text{eq}})$  approach  $360^\circ$ . For comparison, mononuclear  $[\text{Cu}(\text{LH}_3)]\text{BAR}_4^{\text{F}}$



Scheme 1 Synthesis of heterobimetallic copper–Group 13 complexes.

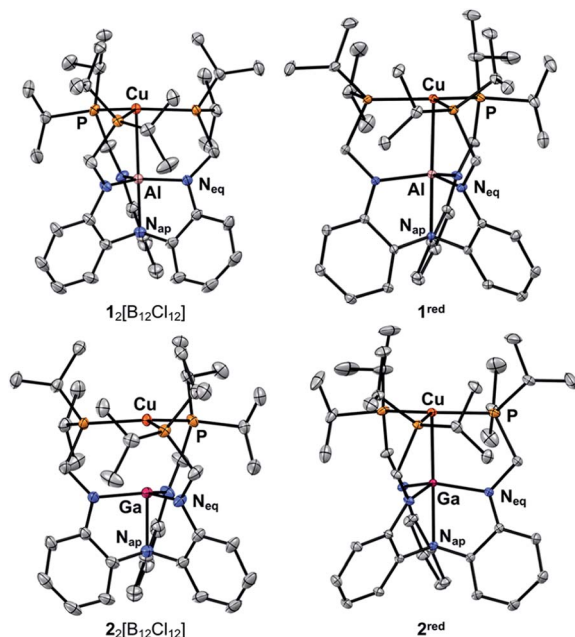


Fig. 2 Solid-state structures of  $[1]_2[B_{12}Cl_{12}]$ ,  $[2]_2[B_{12}Cl_{12}]$ ,  $1^{red}$ , and  $2^{red}$ . Thermal ellipsoids are depicted at 50% probability. Counteranions, co-crystallized solvent molecules, and hydrogen atoms are omitted for clarity.

Table 1 Structural metrics including bond lengths (Å) and angles (°) for complexes  $[1]_2[B_{12}Cl_{12}]$ ,  $[2]_2[B_{12}Cl_{12}]$ ,  $1^{red}$ , and  $2^{red}$

Parameter	$[1]_2[B_{12}Cl_{12}]$	$1^{red}$	$[2]_2[B_{12}Cl_{12}]$	$2^{red}$
M–Cu	2.6239(8)	2.5298(4)	2.5737(5)	2.4541(6)
$r^a$	1.04	1.00	1.01	0.97
Cu–P	2.2925(4) <sup>b</sup>	2.2487(7) <sup>c</sup>	2.2994(5) <sup>b</sup>	2.2698(9) <sup>c</sup>
M–N <sub>eq</sub>	1.85469(12) <sup>b</sup>	1.8915(1) <sup>c</sup>	1.9059(14) <sup>b</sup>	1.9574(9) <sup>c</sup>
M–N <sub>ap</sub>	2.000(2)	2.0836(1)	2.069(3)	2.1948(9)
Cu-to-P <sub>3</sub> -plane	0.188	0.118	0.180	0.136
M-to-N <sub>3</sub> -plane	0.07	0.246	0.142	0.355
$\Sigma(\angle P-Cu-P)$	358.02(1)	359.18(2)	358.19(2)	358.93(4)
$\Sigma(\angle N_{eq}-M-N_{eq})$	359.47(1)	355.01(1)	358.36(2)	350.36(4)

<sup>a</sup> Ratio of the Cu–Al/Ga bond length to the sum the Cu and Al/Ga Alvarez covalent radii.<sup>†</sup> <sup>b</sup> Trigonal space groups only display one value by symmetry. <sup>c</sup> Average of three unique values.

exhibits nearly perfect trigonal symmetry in the solid state ( $\Sigma(\angle P-Cu-P) = 359.96(2)^\circ$  and Cu-to-P<sub>3</sub>-plane = 0.027 Å, Table S2 and Fig. S2†). The average Cu–P bond length in  $[Cu(LH_3)]BARF_4$  of 2.3001(4) Å is similar to that of  $[1]_2[B_{12}Cl_{12}]$  and  $[2]_2[B_{12}Cl_{12}]$ , which indicates that all three complexes possess a Cu(I)  $d^{10}$  center.

After reduction to  $1^{red}$  and  $2^{red}$ , the Cu–M bond contracts significantly by 0.094 Å ( $r = 1.00$ ) and 0.120 Å ( $r = 0.97$ ), respectively, indicating an increase in the metal–metal bonding. Of relevance, bonding between copper and a Group 13 centers have been structurally characterized within clusters<sup>46</sup> and unsupported complexes,<sup>47–49</sup> although these prior examples do not feature odd-electron bonds. In comparison to  $1^{red}$  and  $2^{red}$ , the known Cu–Al<sup>49</sup> and Cu–Ga<sup>47–49</sup> complexes have lower  $r$  values

of 0.90–0.91. The greater  $r$  values for  $1^{red}$  and  $2^{red}$  are expected considering that their bonding involves only one electron. Another indication of a metal–metal interaction is the pyramidalization of the Group 13 center as indicated by the  $\Sigma(\angle N_{eq}-M-N_{eq})$  values of 355.01(1)° and 350.36(4)° for  $1^{red}$  and  $2^{red}$ , respectively, as both the M-to-(N<sub>3</sub>-plane) and M–N<sub>ap</sub> distances increase as the M center moves closer to Cu.<sup>50</sup> The local coordination environment around the Cu center also changes, with the slight lowering of the Cu center into the P<sub>3</sub>-plane and the contraction of the Cu–P bonds upon reduction to the  $\{Cu-M\}^{11}$  core. The Cu–P bond contraction is consistent with an increased Cu electron density that more strongly  $\pi$ -backbonds into the phosphine ligands. Moreover, the significant perturbation of the Group 13 centers suggests their direct participation in stabilizing the reduced  $\{Cu-M\}^{11}$  core, presumably *via* bonding interactions to the Cu center.

### EPR spectroscopy

To gain insight into the radical nature of the one-electron bonds, room-temperature EPR spectra were collected for solutions of  $1^{red}$  and  $2^{red}$  in a 1 : 1 solvent mixture of 2-MeTHF and toluene (Fig. 3). Complex  $1^{red}$  gave rise to a complex signal centered at  $g \sim 2.0$ , which corresponds to a highly overlapped manifold of hyperfine lines that can be decomposed into a leading 1 : 1 : 1 : 1 : 1 sextet due to hyperfine coupling (hfc) of the electronic spin  $S = 1/2$  with  $^{27}Al$  nuclei ( $I = 5/2$ , 100% natural abundance) and additional splitting of each of these hyperfine transitions into two 1 : 1 : 1 : 1 quartets due to additional coupling to the isotopes  $^{63}Cu$  and  $^{65}Cu$  on the copper site ( $I = 3/2$  both, with 69.2% and 30.8% natural abundance, respectively). The corresponding combined hyperfine transitions are assigned by stick spectrum shown on top of Fig. 3. The experimental spectrum was well modeled using an isotropic  $g$  value of 2.007 and three isotropic hyperfine values ( $A$ , in MHz): 269.7 ( $^{27}Al$ ), 171.4 ( $^{63}Cu$ ), and 182.6 ( $^{65}Cu$ ).

Complex  $2^{red}$  likewise exhibited a complex signal with an isotropic  $g$  value of 2.002, which was well modeled as a combined hyperfine splitting of 1 : 1 : 1 : 1 quartets into 1 : 1 : 1 : 1 quartets due to dominating hfc with  $^{69/71}Ga$  ( $I = 3/2$ ) and weaker coupling with  $^{63/65}Cu$  (also  $I = 3/2$ ). The two larger quartets arise from the distinct couplings to the two naturally abundant Ga isotopes with  $A(^{69}Ga) = 1199.3$  MHz (60.1%) and  $A(^{71}Ga) = 1524.3$  MHz (39.9%), while the smaller quartets result from the similar coupling to the two Cu isotopes:  $A(^{63}Cu) = 128.9$  MHz and  $A(^{65}Cu) = 137.3$  MHz. Furthermore, super-hyperfine coupling was observed at the superimposed  $m_I = -1/2$  transition for both the  $^{69}Ga$  and  $^{71}Ga$  quartets (inset of Fig. 3). A reasonable fitting was obtained by including the hyperfine interaction of three equivalent  $^{31}P$  nuclei ( $I = 1/2$ ,  $A = 29.4$  MHz), three equivalent  $^{14}N$  ( $I = 1$ ,  $A = 20.1$  MHz), and one unique  $^{14}N$  ( $A = 3.5$  MHz). The simulation is consistent with the ligand's three equivalent amido donors and one apical amine donor.<sup>29</sup>

The distribution of the unpaired spin density across the Cu and Group 13 valence orbitals in the  $\{Cu-M\}^{11}$  core can be evaluated by comparing the corresponding experimental  $A_{iso}$  values





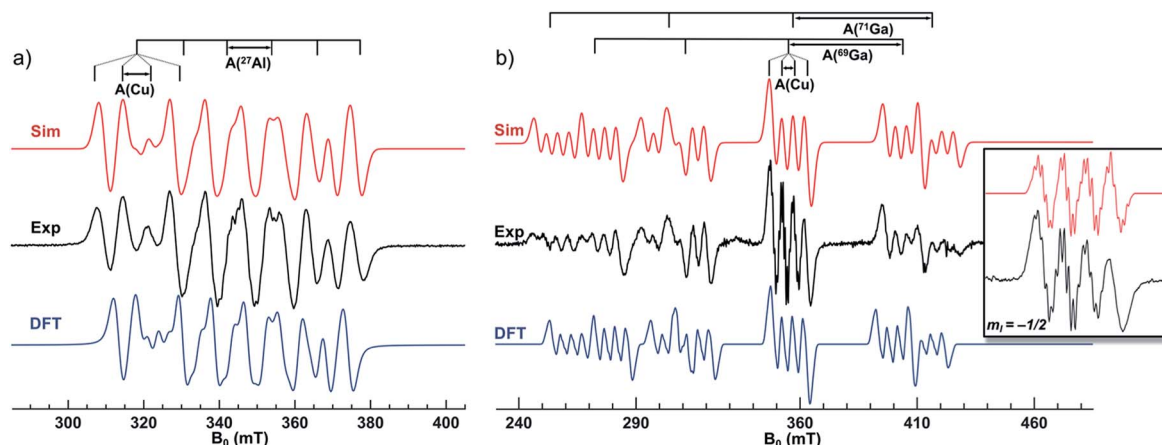


Fig. 3 X-band (9.64 GHz) EPR spectra of 2.5 mM solutions of (a)  $1^{\text{red}}$  and (b)  $2^{\text{red}}$  in a 1 : 1 mixture of 2-MeTHF/toluene at 298 K (time constant, 22 ms). The plot shows the simulated (red) and experimental (black) spectra, as well as a simulation with DFT-derived parameters using the PBE0 functional and def2-TZVP/def2/J basis sets (blue). Inset is a zoom of the  $^{63/65}\text{Cu}$  hyperfine pattern at the superimposed  $m_I = -1/2$   $^{69}\text{Ga}/^{71}\text{Ga}$ -transitions of  $2^{\text{red}}$ . Microwave power and modulation amplitudes used for  $1^{\text{red}}$  and  $2^{\text{red}}$  are 0.7 and 21.3 mW, and 0.98 and 0.50 mT, respectively. See ESI† for additional EPR spectra and details.

with tabulated reference values ( $a_{\text{iso}}^0$ ) for each specific elemental orbital.<sup>51</sup> Using the reference values for a localized valence  $s$ -based spin for Al (3911 MHz) and Ga (12,210 MHz), the  $^{27}\text{Al}$  3s and the  $^{69}\text{Ga}$  4s character in  $1^{\text{red}}$  and  $2^{\text{red}}$ , respectively, were found to be small at 6.9 and 9.8%, respectively.<sup>51</sup> A similar calculation was performed for the  $^{63}\text{Cu}$  4s orbital ( $a_{\text{iso}}^0 = 5995$  MHz) resulting in even smaller contributions of 2.9 and 2.2% in  $1^{\text{red}}$  and  $2^{\text{red}}$ , respectively. The frozen EPR spectra of  $1^{\text{red}}$  and  $2^{\text{red}}$  were also collected (Fig. S43 and S46†). However, only in the former case were we able to model the axial signal and thus extract the  $^{27}\text{Al}$  hyperfine tensor,  $\mathbf{A} = [229.6, 229.6, 325.4]$  MHz. The anisotropic component of the hfc tensor, defined as  $\mathbf{b} = [(A_x - A_{\text{iso}}), (A_y - A_{\text{iso}}), (A_z - A_{\text{iso}})] = [-32, -32, 64]$  MHz, was then compared to that of a fully localized spin in an Al  $3p_z$  orbital, where  $\mathbf{b}^0 = [-83, -83, 166]$  MHz. We found that the  $^{27}\text{Al}$   $3p_z$  character is much greater at 38%. Similarly, the  $^{63}\text{Cu}$   $4p_z$  contribution was found to be 19.9%.<sup>52</sup> The total estimated spin density between Cu and Al is *ca.* 67%, wherein the majority of the unpaired spin likely resides in the bonding axis between the two metals.

Only a handful of formal Cu(0) complexes are known in the literature, and they exhibit a range of isotropic  $A(^{63}\text{Cu})$  values. The complexes  $\text{Cu}(\text{PMe}_3)_3$  and  $\text{Cu}(\text{TPB})$  have  $A(^{63}\text{Cu})$  values of 120 and 191 MHz, respectively, which compare well to those measured for  $1^{\text{red}}$  and  $2^{\text{red}}$ .<sup>27,53,54</sup> For other literature complexes, the  $A(^{63}\text{Cu})$  values are significantly smaller (<30 MHz) because the unpaired spin is strongly delocalized over the ligand. For example,  $\text{Cu}(\text{Me}_2\text{-cAAC})_2$  (cAAC = cyclic (alkyl)(amino)carbene) displayed a small  $A(^{63}\text{Cu})$  coupling of 22.2 MHz, and the DFT-calculated spin density revealed that the majority of the spin density resides in the  $\text{C}_{\text{carbene}}\text{-N}$   $\pi$ -bond.<sup>55</sup> Similarly, the unpaired spin in  $\text{Cu}(\text{B}_2\text{P}_2)$  ( $\text{B}_2\text{P}_2$  = diphosphine-diboranthracene) was proposed to reside on the redox-active ligand core as no Cu hyperfine was discerned.<sup>56</sup>

Further comparison can be made with hfc to the Group 13 supporting ions with  $^{27}\text{Al}$  and  $^{69}\text{Ga}$  hyperfine values for  $1^{\text{red}}$  and  $2^{\text{red}}$  of 269.7 and 1199.3 MHz, respectively. Due to their highly

unstable nature, examples of mononuclear Al(II) and Ga(II) radicals are sparse. Both alane and gallane radical anions display isotropic hfc with their respective Group 13 centers with values of 432.2 MHz for  $[\text{AlH}_3]^{-}$  and 1179.5 MHz for  $[\text{GaH}_3]^{-}$ .<sup>57,58</sup> In contrast, the more bulky  $[\text{Al}(\text{SiMe}^t\text{Bu}_2)_3]^{-}$  and  $[\text{Ga}(\text{SiMe}^t\text{Bu}_2)_3]^{-}$  complexes exhibit much smaller isotropic hfc of 174.0 and 346.9 MHz, respectively.<sup>59</sup> Larger hfc in main group radicals has been shown to be directly related to increased  $s$ -character *via* pyramidalization.<sup>57,58,60,61</sup> The hyperfine values of  $1^{\text{red}}$  and  $2^{\text{red}}$  are closer to those of  $[\text{MH}_3]^{-}$  because the Group 13 ions are similarly pyramidalized<sup>57,58</sup> and correspondingly are different from  $[\text{Al}(\text{SiMe}^t\text{Bu}_2)_3]^{-}$ , which is more planar ( $\Sigma(\angle \text{Si-M-Si})$ : 358.4°).<sup>59</sup> Lastly, direct comparison to isoelectronic  $[\text{NiML}]^{-}$  analogues show similar hfc to  $^{27}\text{Al}$  (219.7 MHz) and  $^{69}\text{Ga}$  (1050 MHz).<sup>29</sup>

## Theoretical computations

To further elucidate the electronic structure of these species, theoretical calculations were performed using Kohn–Sham density functional theory (KS-DFT). Geometry optimizations were performed on the structures without truncation using the M06-L functional<sup>62</sup> with def2-series basis sets<sup>63</sup> and strong agreement was found between the optimized geometries and the X-ray structures (Table S3 and S4†). Based on the in depth study by Hedegård *et al.*,<sup>64</sup> the PBE0 (ref. 65 and 66) functional in conjunction with def2-TZVP and def2/J basis sets<sup>67</sup> were used to determine the EPR parameters with spin–orbit coupling included. The DFT-calculated EPR parameters also matched well to the experimental  $g$  and isotropic  $A$  values, as well as the anisotropic  $\mathbf{A}$  tensors, when available (Fig. 3, Tables S12 and S13†). The good agreement further supports the validity of the calculated electronic structures as models for the experimental species.

Examination of the molecular orbitals (MOs) of cationic species  $[\text{Cu}(\text{LH}_3)]^+$ ,  $[1]^+$ , and  $[2]^+$  revealed five doubly filled d-orbitals as expected for trigonal Cu(I) species. In both bimetallics, the lowest unoccupied molecular orbital (LUMO, Fig. 4)



was found to be a hybrid MO containing large contributions from the empty Cu  $4p_z$  and Group 13 ns and  $np_z$ -orbitals.<sup>68</sup> The neutral radical species  $1^{\text{red}}$  and  $2^{\text{red}}$  each contain a singly occupied molecular orbital (SOMO) that is similar to the LUMO of the cations, suggesting that one-electron reduction of the cations results in occupancy of the LUMO. The SOMO of  $1^{\text{red}}$  contains major contributions from Cu (8.4%), Al (33.8%), and P (35.2%). Similarly, the SOMO for  $2^{\text{red}}$  possesses contributions from Cu (12.7%), Ga (32.0%), and P (35.2%).

The contributions of the Group 13 atoms to the SOMOs have significant s-character (*ca.* 19% for each species), while the Cu center only possesses contributions from the  $4p_z$ . The  $s/p$ -mixing predicted to be present at both Al and Ga likely contributes to the pyramidalization of the Group 13 center. This result is consistent with the greater hfc values observed in the EPR spectra (*vide supra*). Notably, the P contribution nearly doubles upon reduction (from *ca.* 17% to 35% for both  $1^{\text{red}}$  and  $2^{\text{red}}$ ), further indicating the importance of the ancillary ligand in supporting the more electron-rich Cu center.

The Mulliken spin density of  $1^{\text{red}}$  and  $2^{\text{red}}$  revealed delocalization across Cu, Al/Ga, P, and N atoms as visualized through spin-density plots (Fig. 4, Table S11†). Comparison of spin densities between the two metal pockets  $\text{CuP}_3 : \text{MN}_3$  revealed a ratio of 1.0 : 1.2 suggesting the unpaired electron is almost equally delocalized across the two binding pockets. For  $1^{\text{red}}$ , spin densities were calculated to be 0.18 on Cu and 0.52 on Al,

with a moderate amount distributed throughout the ancillary atoms (0.09 per P and 0.01 per  $\text{N}_{\text{eq}}$ ), and these values are in good agreement to those estimated by EPR spectroscopy (*vide supra*). The calculated spin densities for  $2^{\text{red}}$  were found to be similar to that of  $1^{\text{red}}$  with 0.24 on Cu, 0.44 on Ga, and the remainder on the supporting donor atoms (0.08 per P and 0.03 per  $\text{N}_{\text{eq}}$ ).

By considering the SOMOs and noting the extent of spin delocalization in the spin density plots, the radical species  $1^{\text{red}}$  and  $2^{\text{red}}$  can therefore best be described as possessing  $\sigma$ -bonding interactions between the Cu ( $4p_z$ ) and Al/Ga ( $np_z/\text{ns}$ ) orbitals, as well as a  $\pi$ -backbonding interaction between Cu and the three P atoms.<sup>53,54,69</sup> The covalent nature of the  $\sigma(\text{Cu}-\text{M})$  bonds is supported by the significant contraction of the metal-metal bond lengths in the solid-state structures and the nearly equal spin delocalization between the two metal pockets, which matches well with the experimental EPR hfc values.

## Conclusion

In summary, two cationic copper-Group 13 complexes were synthesized and their subsequent reduction provided access to unusual  $17 e^- \{ \text{Cu}-\text{M} \}^{11}$  complexes. These radical species are characterized by short Cu-M bond lengths and exhibit significant hfc to Cu, M, and P atoms, collectively consistent with the highly covalent nature of SOMOs that involve both metal centers and the ancillary ligand donor atoms. These neutral radical  $\{ \text{Cu}-\text{M} \}^{11}$  species possess rarely observed  $2c/1e$   $\sigma$ -bonds and are the first report of any such bond for the heavier elements Al and Ga. Along with the previously reported  $\text{Cu}(\text{TPB})$ ,  $1^{\text{red}}$  and  $2^{\text{red}}$  complete a triad of copper-Group 13 complexes that demonstrate an increase in Cu bonding character and Cu-M covalency when traversing down the group in the order  $\text{B} < \text{Al} < \text{Ga}$ .<sup>27</sup> Investigation of the metalloradical reactivity of these bimetallics is ongoing.

## Data availability

Additional data for synthesis, characterization, and DFT calculations are available in the ESI.† The combined XYZ coordinates of DFT-optimized geometries have been uploaded as ESI.†

## Author contributions

BJG and MVV performed the experimental work supervised by CCL. TRS performed the theoretical calculations supervised by LG. MJD assisted with the EPR data collection, and EB assisted with the EPR analysis and interpretation. VGY provided X-ray support. The original draft was written by BJG with guidance from CCL, and all authors edited the manuscript.

## Conflicts of interest

There are no conflicts to declare.

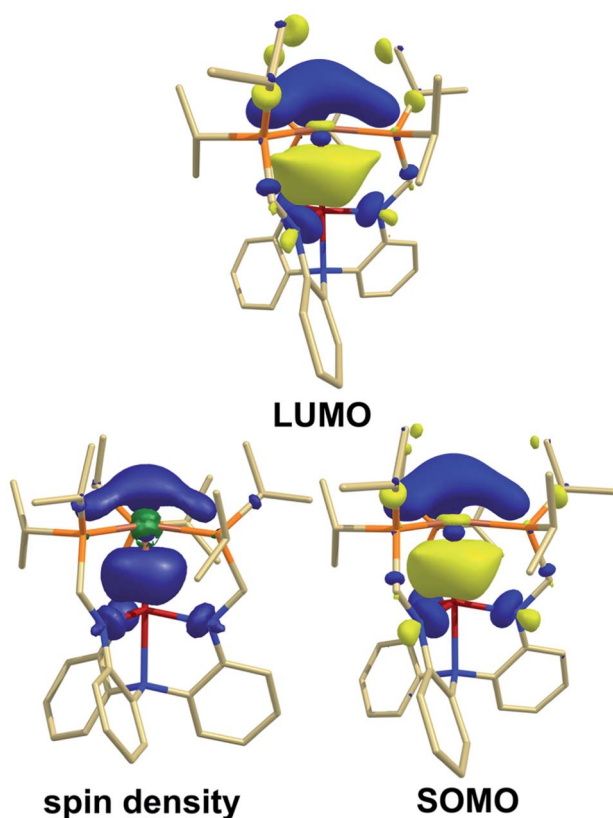


Fig. 4 Contour plots for the LUMO for **2** and the SOMO and spin density for  $2^{\text{red}}$ . The corresponding contour plots for  $1/1^{\text{red}}$  are similar (see ESI†).



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