# **RSC Advances**



## **PAPER**

View Article Online
View Journal | View Issue



Cite this: RSC Adv., 2022, 12, 34627

# Theoretical insights into interfacial stability and ionic transport of Li<sub>2</sub>OHBr solid electrolyte for all-solid-state batteries†

Bo Liu, \*\*D\*\*\*\* Piguang Liao, \*\* Xiaowen Shi, \*\* Yufeng Wen, \*\* Qingdong Gou, \*\* Meidong Yu, \*\* Shenlin Zhou\*\* and Xinyuan Sun\*\*

Li-rich antiperovskite materials are promising candidates as inorganic solid electrolytes (ISEs) for all-solid-state Li-ion batteries (ASSLIBs). However, the material faces several pressing issues for its application, concerning the phase stability and electrochemical stability of the synthesized material and the Li-ion transport mechanism in it. Herein, we performed first-principles computational studies on the phase stability, interfacial stability, defect chemistry, and electronic/ionic transport properties of Li<sub>2</sub>OHBr material. The calculation results show that the Li<sub>2</sub>OHBr is thermodynamically metastable at 0 K and can be synthesized experimentally. This material exhibits a wider intrinsic electrochemical stability window (0.80–3.15 V) compared with sulfide solid electrolytes. Moreover, the Li<sub>2</sub>OHBr displays significant chemical stability when in contact with typical cathode materials (LiCoO<sub>2</sub>, LiMn<sub>2</sub>O<sub>4</sub>, LiFePO<sub>4</sub>) and moisture. The dominant defects of Li<sub>2</sub>OHBr are predicted to be V<sub>Li</sub>- and Li<sub>i</sub>+, corresponding to lower Li-ion migration barriers of 0.38 and 0.49 eV, respectively, while the replacement of some of the OH<sup>-</sup> by F<sup>-</sup> is shown to be effective in decreasing migration barriers in Li<sub>2</sub>OHBr. These findings provide a theoretical framework for further designing high performance ISEs.

Received 1st November 2022 Accepted 25th November 2022

DOI: 10.1039/d2ra06921k

rsc.li/rsc-advances

## 1. Introduction

In recent years, inorganic solid electrolytes (ISEs) have received wide attention to replace organic liquid electrolytes currently used in commercial lithium-ion batteries.¹ Unlike organic liquid electrolytes, ISEs are nonflammable, nonvolatile, and have no liquid leakage problem, and are also expected to overcome the phenomenon of lithium dendrites, thus they have high safety performance.² Moreover, ISEs have the potential to improve interfacial stability, which could enable the application of a high-voltage cathode and even lithium metal anode.³ In terms of Li-ion transport properties, several ISEs have been reported, such as the Li<sub>10</sub>GeP<sub>2</sub>S<sub>12</sub>,⁴ Li<sub>7</sub>P<sub>3</sub>S<sub>11</sub>,⁵ and lithium-rich anti-perovskites (LRAP),⁶ with high Li⁺ conductivities comparable to or even surpassing those of traditional liquid electrolytes.

One promising class of ISEs is antiperovskites  $\text{Li}_{3-n}\text{OH}_n\text{X}$  (n = 0--1, X = Cl, Br) for ASSLIBS.<sup>7</sup> For example, Zhao *et al.* first experimental reported that  $\text{Li}_3\text{OCl}$  and  $\text{Li}_3\text{O}(\text{Cl}_{0.5}\text{Br}_{0.5})$  showed a high ion-conductivity (>10<sup>-3</sup> S cm<sup>-1</sup> at 300 K) with low activation energies (0.18–0.26 eV).<sup>8</sup> Subsequently, a good stability

and low  $Li^+$  vacancy migration barrier of  $Li_3OX$  (X = Cl, Br) was verified by first-principles calculations.9 Sugumar et al. reported the successful preparation of Li2OHBr by dry ball-milling of LiOH and LiBr at room temperature, which obtained high ion conductivity of  $1.1 \times 10^{-6} \, \mathrm{S \ cm^{-1}}$  with the activation energy of 0.54 eV.10 Recently, Yamamoto et al.11 reported that an ASSLIB composed of Li/Li<sub>2</sub>OHBr/Fe<sub>2</sub>(MoO<sub>4</sub>)<sub>3</sub> were fabricated by pressing at room temperature, which exhibited good chargedischarge performance and excellent cycle stability. Zhao and co-workers<sup>12</sup> proposed the Li<sub>2</sub>OHBr as a protective layer for the Li<sub>1.5</sub>Al<sub>0.5</sub>Ge<sub>1.5</sub>(PO<sub>4</sub>)<sub>3</sub> (LAGP) solid electrolyte to prevent the side reaction caused by direct contact between LAGP and Li metal anode. For practical applications, whether Li<sub>2</sub>OHBr material acts as a solid electrolyte or protective layer, it is crucial that the material shows good thermodynamic stability, electrochemical stability and fast Li-ion diffusion, which are the keys to ameliorating the electrochemical and rate performance of Li<sub>2</sub>-OHBr material. Moreover, Li<sub>2</sub>OHBr should possess the ability of moisture resistance and oxidation resistance, which will simplify the packaging of ASSLIBs in practice. However, indepth understanding of these important issues, has been hindered by the complicated synthesis and measurement conditions during the experiments. Therefore, it is critical that we explore the fundamental issues of the phase stability, electrochemical stability, chemical stability and electron/ion transport mechanism of Li<sub>2</sub>OHBr through reliable theoretical

<sup>&</sup>quot;College of Mathematics and Physics, Jinggangshan University, Ji'an, Jiangxi 343009, China. E-mail: liubo@jgsu.edu.cn

bScience and Technology Innovation Development Center, Ji'an, Jiangxi 343006, China † Electronic supplementary information (ESI) available. See DOI: https://doi.org/10.1039/d2ra06921k

approaches to elucidating the main behind physics mechanism.

In this work, we employ first-principles calculations to assess the phase stability, interfacial stability against electrode material, defect chemistry and electron/ion transport mechanism of anti-perovskites  $\text{Li}_2\text{OHBr}$ . We predict a wide electrochemical window and low chemical reactivity for  $\text{Li}_2\text{OHBr}$ , ensuring that this material is thermodynamically stable under high-voltage operation and in air. The analysis Li-ion transport mechanism shows the existence of low migration barriers involving charge carriers  $(V_{\text{Li}^-}, \text{Li}_i^+)$  in  $\text{Li}_2\text{OHBr}$ . We also study the effect of F<sup>-</sup> substitution of OH<sup>-</sup> on the Li<sup>+</sup> migration barriers in  $\text{Li}_2\text{OHBr}$ . The computational approach in this work can be extended to the design of other ISEs system.

# 2. Computational methods

All calculations are performed based on density functional theory (DFT) by using the projector augmented wave method, as implemented in the Vienna *ab initio* Simulation Package (VASP).<sup>13</sup> The generalized gradient approximation (GGA) with Perdew–Burke–Ernzerhof (PBE) is applied to treat the electronic exchange–correlation interactions.<sup>14</sup> The cutoff energy is set to 520 eV. The atomic force and energy convergence parameters are consistent with Materials Project (MP)<sup>15</sup> for all calculations. Based on DFT ground-state energies in the MP database, the phase stability and interfacial stability (including electrochemical and chemical stability) of Li<sub>2</sub>OHBr are evaluated using the same scheme in the previous work.<sup>16</sup>

The phase stability of  $\rm Li_2OHBr$  is assessed by computing the energy above convex hull, corresponding the decomposition energy to the thermodynamic phase equilibria. The electrochemical window of  $\rm Li_2OHBr$  is calculated by using the Li grand potential phase diagram. In this method, the grand potential  $\phi$  of the  $\rm Li_2OHBr$  is defined as:<sup>16</sup>

$$\phi[c, \mu_{Li}] = E[c] - n_{Li}[c]\mu_{Li}, \tag{1}$$

where  $\mu_{Li}$ ,  $n_{Li}[c]$ , and E[c] are the Li chemical potential, Li concentration and DFT energy of composition c, respectively.

The chemical stability of  $\text{Li}_2\text{OHBr/cathode}$  interfaces are determined by estimating the reaction between the  $\text{Li}_2\text{OHBr}$  and cathode with the lowest reaction energy  $(\Delta E)$ , <sup>17</sup> namely:

$$\Delta E(c_{\text{cathode}}, c_{\text{Li}_2\text{OHBr}}) = \min_{x \in [0, 1]} \left\{ E_{\text{pd}}[xc_{\text{Li}_2\text{OHBr}} + (1 - x)c_{\text{cathode}}] - xE[c_{\text{Li}_2\text{OHBr}}] - (1 - x)E[c_{\text{cathode}}] \right\},$$
(2)

where x is the mixing fraction of the cathode/coating compositions,  $E[c_{\text{cathode}}]$  and  $E[c_{\text{Li}_2\text{OHBr}}]$  are DFT energies of the cathode and Li<sub>2</sub>OHBr, respectively.  $E_{\text{pd}}[c]$  represents the DFT energy corresponding to the given composition C in the chemical reaction equilibrium.

The defect formation energy  $E_{\rm f}(i,q)$  with a defect i at charge state q is calculated via the following equation:<sup>18</sup>

$$E_{\rm f}(i, q) = E_{\rm tot}(i, q) - E_{\rm tot}({\rm Li_2OHBr, bulk}) - n\mu_i + q(\varepsilon_{\rm F} + E_{\rm V}), (3)$$

where  $E_{\text{tot}}(\text{Li}_2\text{OHBr}$ , bulk) and  $E_{\text{tot}}(i, q)$  are the DFT energies of Li<sub>2</sub>OHBr supercell without and with a defect, respectively. n is the number of Li or Li<sub>2</sub>O compound added (n > 0) or removed (n < 0) from the perfect supercell, and  $\mu_i$  is the chemical potential of Li (bcc structure) or Li<sub>2</sub>O (fcc structure) compound. The Fermi level  $(\varepsilon_F)$  is defined relative to the valence-band maximum  $(E_V)$  of perfect Li<sub>2</sub>OHBr and varies with defect concentration.

## 3. Results and discussion

#### 3.1 Structural and electronic property

The Li<sub>2</sub>OHBr orthorhombic structure with space group *Cmcm* is constructed, as shown in Fig. 1(a). Table 1 lists the optimized lattice parameters of Li<sub>2</sub>OHBr obtained by DFT calculations, together with the data from Howard *et al.*<sup>19</sup> The lattice parameters are a = 8.015 Å, b = 8.152 Å, and c = 7.944 Å, which are agree with the experiment data (within 2% error). To obtain the electronic properties of Li<sub>2</sub>OHBr, the electronic band structure of Li<sub>2</sub>OHBr is evaluated. Since the GGA-PBE functional generally

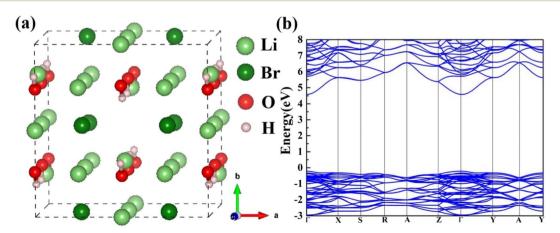


Fig. 1 (a) Atomic structure and (b) electronic band structure of Li<sub>2</sub>OHBr. The ball colors green, dark red, red, and gray indicate Li, Br, O, and H sites, respectively.

Table 1 Relaxed lattice parameters and atomic coordinates of  $\text{Li}_2$ -OHBr in the orthorhombic

System		Calc. (this work)			Exp. (ref. 19)		
Lattice parameters		a (Å) 8.015	b (Å) 8.152	c (Å) 7.944	a (Å) 8.010	b (Å) 8.030	c (Å) 7.880
		α (°)	$\beta$ (°)	$\gamma$ (°)	α (°)	$\beta$ (°)	$\gamma$ (°)
		90	90	90	90	90	90
Atom	Wyckoff	x	у	z	x	у	z
O	8f	0.000	0.744	0.514	0.000	0.741	0.513
H	8f	0.000	0.823	0.414	0.000	0.824	0.419
Cl	8g	0.749	0.489	0.250	0.747	0.487	0.250
Li	8d	0.250	0.250	0.000	0.250	0.250	0.000
Li	4b	0.000	0.500	0.000	0.000	0.500	0.000
Li	4c	0.000	0.195	0.250	0.000	0.197	0.250

underestimate the band gap, the advanced SCAN meta-GGA functional is used to provide a rigorous result.  $^{20}$  Fig. 1(b) shows that the direct band gap of Li<sub>2</sub>OHBr is 4.79 eV, and its VBM and CBM are both located at the  $\Gamma$  point of the first Brillouin zone. The large band gap indicates that Li<sub>2</sub>OHBr is electron insulator, which can effectively block electron leakage and prevent electrode corrosion. The band gap between the valence band maximum (VBM) and the conduction band minimum (CBM) provides an upper limit for the electrochemical window of solid electrolytes. When the chemical potential of the electrode/solid electrolyte is mismatched, *i.e.*, thermodynamically unstable, a chemical reaction between the two materials will occur spontaneously upon contact.

#### 3.2 Phase stability and interfacial stability

The feasibility and complexity level of the experimental synthesis of one given material could be evaluated by its phase stability. In general, for a particular component that does not have a stable phase, the corresponding system will decompose into a stable phase around its component coordinate points. The Li–H–O–Br quaternary phase diagram is constructed at 0 K by minimizing the formation energies of various compositions, as shown in Fig. 2(a). The phase diagram indicates that Li<sub>2</sub>OHBr is

energetically unstable duo to positive formation energy with respect to that form the  $\text{Li}_4\text{H}_3\text{BrO}_3$  and LiBr. However, the energy above hull ( $E_{\text{hull}}$ ) of  $\text{Li}_2\text{OHBr}$  is only 11.6 meV per atom, suggesting that  $\text{Li}_2\text{OHBr}$  is likely to be metastable at 0 K against possible decomposition products, and therefore may be stabilized by external conditions (such as pressure, temperature, and entropy). For example, the  $\text{Li}_2\text{OHBr}$  was synthesized from LiOH and LiBr starting materials by sintering method above 300 °C. In fact, successful synthesis of metastable phases has been widely reported for ISEs, such as  $\text{Li}_6\text{PS}_5\text{Cl}$  (21 meV per atom), Light-GeP<sub>2</sub>S<sub>12</sub> (25 meV per atom) and  $\text{Li}_7\text{P}_3\text{S}_{11}$  (27 meV per atom).

For the practical application of ASSLIBs, the ISEs should satisfy the conditions of good interfacial stability, including electrochemical stability and chemical stability.24 Using eqn (1), the phase equilibrium of Li<sub>2</sub>OHBr for a series of lithiation/ delithiation reactions is predicted to obtain the electrochemical stability window. The detailed lithiation/delithiation reactions with  $\mu_{Li}$  are listed in Table S1.† As shown in Fig. 2(b), Li<sub>2</sub>OHBr is oxidized to form Li<sub>4</sub>H<sub>3</sub>BrO<sub>3</sub> and Br when the oxidation voltage is higher than 3.15 V. Meanwhile, Li<sub>2</sub>OHBr is reductively decomposed into LiH, LiBr and Li2O starting from 0.80 V. The calculated electrochemical stability window range of Li<sub>2</sub>OHBr are 0.80-3.15 V vs. Li/Li<sup>+</sup>. Table 2 shows that Li<sub>2</sub>OHBr has much wider electrochemical window than that of reported sulfides and oxides solid electrolyte, such as Li<sub>10</sub>GeP<sub>2</sub>S<sub>12</sub> (1.71-2.14 V), Li<sub>7</sub>P<sub>3</sub>S<sub>11</sub> (2.28–2.31 V), Li<sub>6</sub>PS<sub>5</sub>Cl (1.71–2.01 V), Li<sub>2</sub>PO<sub>2</sub>N (0.68-2.63 V) and  $\text{Li}_7\text{La}_3\text{Zr}_2\text{O}_{12}$  (0.05-2.91). However, it should be pointed out that the dissociation of ISEs depends on kinetic factors, suggesting that the dissociation of the phases may be slowed down or interrupted under certain circumstances, such as slow electron/ion transport in dissociated phases. The above calculations assume complete thermodynamic equilibrium and no kinetic constraints in the reactions. Therefore, ISEs are expected to withstand a wider range of voltages in use than calculated, as measured by cyclic voltammetry (CV) using Li/ Li<sub>2</sub>OHBr/Au cell, the electrochemical potential window of Li<sub>2</sub>-OHBr was 1.7-3.5 V.11

The chemical stability between the Li<sub>2</sub>OHBr and various cathodes is calculated by using eqn (2). Three typical cathode

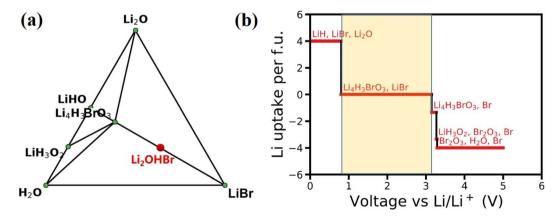


Fig. 2 (a)  $\text{Li}_{-}\text{H}_{-}\text{O}_{-}\text{Br}$  quaternary phase diagram. The metastable  $\text{Li}_{2}\text{OHBr}$  is marked in red font. (b) The voltage profiles and phase equilibria of  $\text{Li}_{2}\text{OHBr}$  at different potentials. The light yellow region indicates the electrochemical stability window range.

Table 2 Summary of electrochemical window (EW), equilibria phases at reduction and oxidation potentials of some typical solid-state electrolytes

Solid-state electrolytes	EW vs. Li/Na (V)	Equilibria phase at reduction potential	Equilibria phase at oxidation potential
Li <sub>2</sub> OHBr (this work)	0.80-3.15	LiH, LiBr, Li <sub>2</sub> O	Li <sub>4</sub> H <sub>3</sub> BrO <sub>3</sub> , Br
Li <sub>2</sub> OHCl (ref. 26)	0.82-3.15	LiH, LiCl, Li <sub>2</sub> O	LiH <sub>2</sub> ClO <sub>5</sub> , H <sub>2</sub> O, LiCl
Li <sub>3</sub> OCl (ref. 9)	0-2.55	Li <sub>3</sub> OCl	$\text{Li}_2\text{O}_2$ , LiCl
Na <sub>3</sub> OBr (ref. 27)	0-1.79	Na <sub>3</sub> OBr	Na <sub>2</sub> O <sub>2</sub> , NaBr
Li <sub>10</sub> GeP <sub>2</sub> S <sub>12</sub> (ref. 25)	1.71-2.14	Li <sub>4</sub> GeS <sub>4</sub> , Li <sub>2</sub> S, P	Li <sub>3</sub> PS <sub>4</sub> , GeS <sub>2</sub> , S
Li <sub>3</sub> PS <sub>4</sub> (ref. 25)	1.71-2.31	Li <sub>2</sub> S, P	P <sub>2</sub> S <sub>5</sub> , S
Na <sub>3</sub> PS <sub>4</sub> (ref. 27)	1.39-2.45	Na <sub>3</sub> P, Na <sub>2</sub> S	P <sub>2</sub> S <sub>5</sub> , S
Li <sub>7</sub> P <sub>3</sub> S <sub>11</sub> (ref. 25)	2.28-2.31	$\text{Li}_3\text{PS}_4$ , $\text{P}_4\text{S}_9$	$P_2S_5$ , S
Li <sub>6</sub> PS <sub>5</sub> Cl (ref. 25)	1.71-2.01	Li <sub>2</sub> S, LiCl, P	Li <sub>3</sub> PS <sub>4</sub> , LiCl, S
Li <sub>2</sub> PO <sub>2</sub> N (ref. 25)	0.68-2.63	Li <sub>3</sub> P, LiPN <sub>2</sub> , Li <sub>2</sub> O	$P_3N_5$ , $Li_4P_2O_7$ , $N_2$
$\text{Li}_7\text{La}_3\text{Zr}_2\text{O}_{12}$ (ref. 25)	0.05-2.91	$Zr_3O$ , $La_2O_3$ , $Li_2O$	$\mathrm{Li_2O_2}$ , $\mathrm{La_2O_3}$ , $\mathrm{Li_6Zr_2O_7}$

materials (e.g., layered LiCoO2, spinel LiMn2O4, olivine LiFePO<sub>4</sub>) are considered for fully-discharged and half-charged state. Fig. 3 shows the predicted reaction energies between the Li<sub>2</sub>OHBr and cathode materials, and the corresponding reaction products are listed in Table S2.† The mutual reaction energy  $\Delta E_{\min}$  of Li<sub>2</sub>OHBr with both fully-discharged cathodes are predicted to have low reaction energies (0 <  $|\Delta E_{\min}|$  < 50 meV per atom) in Fig. 3(a). The chemical reactivity sequence for a fully-discharged cathodes with the Li<sub>2</sub>OHBr is LiFePO<sub>4</sub> > LiMn<sub>2</sub>O<sub>4</sub> > LiCoO<sub>2</sub>, suggesting that LiCoO<sub>2</sub> cathode seems to have better interfacial compatibility with the Li<sub>2</sub>OHBr. An increased chemical reactivity of Li2OHBr with half-charged cathodes are observed from fully-discharged to half-charged cathode states, as shown in Fig. 3(b). It is worth noting that the Li<sub>2</sub>OHBr has either no reaction or a negligible driving force against LiCoO2 cathode, showing a significant thermodynamic chemical stability. Certainly, for all interfacial reactions  $\Delta E_{\min}$  < 0, a thermodynamically unstable interface is formed when Li<sub>2</sub>-OHBr is in contact with high-voltage cathodes, which may result in the formation of unwanted interfacial byproducts, thereby reducing the rate capacity and electrochemical performance of ASSLIBs. Therefore, further experiment techniques are awaited to assess potential interfacial products, such as X-ray photoelectron spectroscopy (XPS), X-ray absorption spectroscopy (XAS), and electron microscopy (EM).<sup>28</sup>

The air stability of ISEs is another issue involved in electrolyte handling and battery assembly in the development of ASSLIBs, where the electrolyte will inevitably be exposed to air and undergo structural changes if it is not chemically stable. For example, using the first-principles calculations, Zhang et al. The revealed the thermodynamic and kinetic mechanism in the reaction of  $\text{Li}_1\text{O}\text{GeP}_2\text{S}_{12}$  with  $\text{H}_2\text{O}$  in air to produce  $\text{H}_2\text{S}$  gas. Here, the stability of  $\text{Li}_2\text{O}\text{HBr}$  toward air is also studied using reaction energies  $\Delta E$  calculation for the reaction with moisture, the higher negative value indicates the material is strongly favorable to react with moisture. The estimation of driving force for  $\text{Li}_2\text{O}\text{HBr}$  when exposed to air  $\nu ia$  following reaction:

$$3Li_2OHBr + 2H_2O \rightarrow 2LiOH_2Br + Li_4(OH)_3Br,$$
 (4)

$$2\text{Li}_4(OH)_3Br + 3CO_2 \rightarrow 2\text{Li}OH_2Br + 3\text{Li}_2CO_3 + H_2O.$$
 (5)

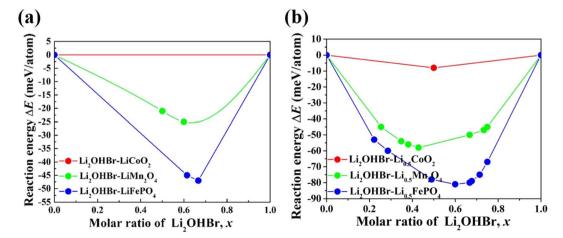


Fig. 3 Predicted reaction energies  $\Delta E$  between Li<sub>2</sub>OHBr and various cathode materials in both (a) fully-discharged and (b) half-charged states. The reaction energy  $\Delta E_{min}$  corresponds to the lowest mutual reaction energy at a ratio of x.

The estimated value of  $\Delta E$  is only -1 meV per atom when  $\mathrm{Li_2OHBr}$  reacts with  $\mathrm{H_2O}$ . While  $\mathrm{Li_4(OH)_3Br}$  forms as a hydrolysis intermediate and subsequently reacts favorably with  $\mathrm{CO_2}$  to produce  $\mathrm{LiOH_2Br}$ ,  $\mathrm{Li_2CO_3}$  and  $\mathrm{H_2O}$  ( $\Delta E = -118$  meV per atom). Therefore, it is suggested that the  $\mathrm{Li_2OHBr}$  solid electrolyte is stable in dry air. To understand the degradation mechanism of ISEs exposed to air, some experimental techniques, such as *in situ* scanning/transmission electron microscopy, neutron ray diffraction depth analysis and synchrotron X-ray imaging technologies, have been performed to track local nanoscale chemical evolution and structural information of interfacial phases.<sup>31</sup>

#### 3.3 Ionic transport mechanism

**3.3.1. Defect structure and formation energy.** Defects in ISEs have a remarkable effect on Li-ion diffusion, and the types of defects are determined by synthesis conditions (*e.g.*, chemical potential) affecting their formation energies. To elucidate the ion transport mechanism, it is necessary to understand which defects favor the diffusion of lithium ions in Li<sub>2</sub>OHBr. Herein, four defect types in Li<sub>2</sub>OHBr are considered, including lithium vacancy  $(V'_{Li})$ , lithium interstitial  $(Li_i)$ , lithium Frenkel defect pair  $(V'_{Li} - Li_i)$  and Li<sub>2</sub>O Schottky defect pair  $(2V'_{Li} - V'_{O})$ , as per our previous work. Among these defects, lithium Frenkel defect and Li<sub>2</sub>O Schottky defect pair in Li<sub>2</sub>OHBr are described as follows:

Lithium Frenkel defect:

$$Li_{i,i}^{\times} \rightarrow V_{i,i}^{\prime} + Li_{i,i}^{\prime},$$
 (6)

Li<sub>2</sub>O Schottky defect:

$$2Li_{Li}^{\times} + O_{O}^{\times} \rightarrow 2V_{Li}^{'} + V_{O}^{"} + Li_{2}O.$$
 (7)

According to the symmetry of  $\text{Li}_2\text{OHBr}$ , the possible defect configurations and formation energies are investigated to obtain the lowest energy configuration. The possible defect configurations include two different lithium vacancy defects  $(V'_{\text{Li}}\text{-I}\text{I} \text{ and } V'_{\text{Li}}\text{-II})$ , one lithium interstitial defect  $(\text{Li}_i^{\cdot})$ , two lithium Frenkel defect pair  $(V_{\text{near}} \text{ and } V_{\text{far}})$ , and three  $\text{Li}_2\text{O}$  Schottky defect pair  $(V_{\text{adjacent}}, V_{\text{separated-1}}, \text{and } V_{\text{separated-2}})$  in the ESI of Fig. S1 and S2.† The formation energies of four defect types  $(V'_{\text{Li}}, \text{Li}_i^{\cdot}, V'_{\text{Li}} - \text{Li}_i^{\cdot}, \text{ and } 2V'_{\text{Li}} - V'_{\text{O}})$  are calculated by eqn (3) in the neutral state, as listed in Table 3. By comparing the formation energies, it is found that the dominant defect configuration is  $V_{\text{near}}$ , and the corresponding defect formation energies is 0.34 eV. In contrast, single  $V'_{\text{Li}}$  and  $\text{Li}_i^{\cdot}$  show higher defect formation energies (3.93 eV and 1.29 eV), implying a lower concentration of lithium vacancy and lithium

interstitial defect in neutral Li<sub>2</sub>OHBr. In addition, the defect formation energies of lithium vacancy and interstitial at different charge states q as a function of Fermi level are also calculated in Fig. 4. The results show that the formation energies of Li<sub>i</sub><sup>+</sup> and V<sub>Li</sub><sup>-</sup> are lower than those of V'<sub>Li</sub>, Li<sup>-</sup><sub>i</sub>, V'<sub>Li</sub> – Li<sup>-</sup><sub>i</sub> and  $2V'_{Li} - V_{O}^{\bullet}$ , suggesting that the Li<sub>i</sub><sup>+</sup> and V<sub>Li</sub><sup>-</sup> may be the main defect types in Li<sub>2</sub>OHBr at room temperature. Therefore, charged defects (Li<sub>i</sub><sup>+</sup> and V<sub>Li</sub><sup>-</sup>) should act as charge carriers in Li<sub>2</sub>OHBr and will be used to calculate the migration barriers in the next section.

3.3.2. Li-ion migration via vacancy and interstitial mechanism. A high ionic conductivity at room temperature is one of most important indices for the practical application of ISEs in ASSBs. Herein, the Li-ion diffusion pathways and migration barriers of charge carriers (V<sub>Li</sub>-, Li<sub>i</sub>+) in Li<sub>2</sub>OHBr are studied using CI-NEB method. Fig. 5 shows the diffusion path of a single lithium vacancy  $(V_{Li^-})$  along the *ab*-plane and *c*-axis in Li<sub>2</sub>OHBr, corresponding to the migration barriers are 0.38 and 0.57 eV, respectively. For the direct lithium interstitial (Li<sub>i</sub><sup>+</sup>) migration in Li<sub>2</sub>OHBr, migration process from an interstitial site to its adjacent interstitial site needs to cross an energy barrier of 0.49 eV, as shown in Fig. 6. The energy barriers of V<sub>Li</sub>is slightly lower than that of Li<sub>i</sub><sup>+</sup>, suggesting that V<sub>Li</sub><sup>-</sup> and Li<sub>i</sub><sup>+</sup> will be generated simultaneously when the migrating Li leaves the lattice site. Therefore, we can conclude that the charge carriers (V<sub>Li</sub>- and Li<sub>i</sub>+) make main contribution to the ionic conductivity of Li2OHBr and that increasing the concentrations of V<sub>Li</sub>- and Li<sub>i</sub><sup>+</sup> defects are crucial to obtain high Li-ion conductivity.

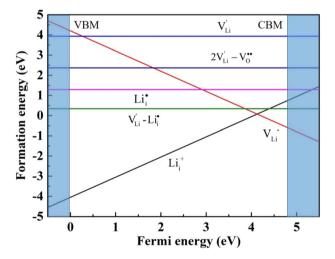


Fig. 4 Formation energies of lithium vacancy and interstitial at different charge states.

Table 3 Defect formation energies corresponding to various defect types and configurations in Li<sub>2</sub>OHBr

Defect type	$\mathbf{V}_{\mathrm{Li}}^{'}\text{-}\mathbf{I}$	$\mathbf{V}_{\mathrm{Li}}^{'}\text{-}\mathbf{II}$	$\mathrm{Li}^{\:\raisebox{3.5pt}{\text{\circle*{1.5}}}}_{\mathrm{i}}$	$V_{near}$	$V_{\rm far}$	$V_{\rm adjacent}$	$V_{\rm separated-1}$	V <sub>separated-2</sub>
$u_i$ (eV) $E_{\rm f}$ (eV)	-1.89 3.93	-1.89 $4.26$	-1.89 $1.29$	 0.34	— 0.62	-14.35 3.01	-14.35 $2.37$	-14.35 $2.75$

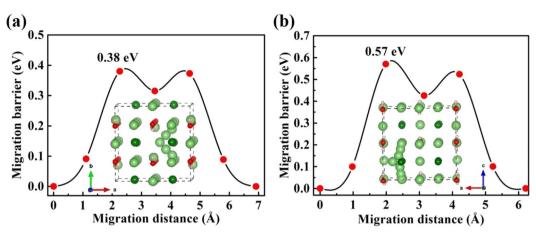


Fig. 5 Diffusion pathway and migration barriers of V<sub>Li</sub>- in Li<sub>2</sub>OHBr (a) along the ab-plane and (b) along the c-axis

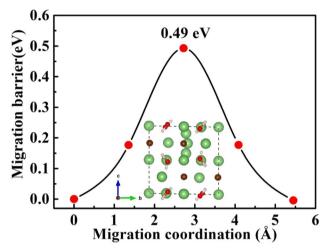


Fig. 6 Diffusion pathway and migration barriers of  $\mathrm{Li_{i}}^{+}$  in  $\mathrm{Li_{2}OHBr}$  along the c-axis.

Furthermore, the doping of Li<sub>2</sub>OHBr with halogen element (F or Cl) may be one of the key factors to enhance the ionic conductivity. In Fig. 7(a) and (b), we present the lithium vacancy

migration barriers obtained using CI-NEB calculations for F-doped Li<sub>2</sub>OHBr, namely Li<sub>2</sub>(OH)<sub>0.875</sub>F<sub>0.125</sub>Br. The results show that the migration barriers of Li<sub>2</sub>(OH)<sub>0.875</sub>F<sub>0.125</sub>Br is 0.37 and 0.33 eV along the *ab*-plane and *c*-axis, respectively, which is a lower value than that of pristine Li<sub>2</sub>OHBr. The main reason is that the substitution of F<sup>-</sup> for OH<sup>-</sup> increases the antiperovskite tolerance factor and favors a disordering of the OH<sup>-</sup> orientation for Li<sub>2</sub>OHBr, as previous reported by Li *et al.*<sup>21</sup> Therefore, the Li<sub>2</sub>OHBr can be doped by the substitution of OH<sup>-</sup> by F<sup>-</sup> is beneficial to reduce the migration barrier and improve the ionic conductivity.

# 4. Conclusions

In conclusion, the electronic properties, phase stability, interfacial stability, defect chemistry and Li-ions migration mechanisms of Li<sub>2</sub>OHBr have been systematically studied by the first-principles calculations. The calculations results indicate that Li<sub>2</sub>OHBr crystal structure is metastable by thermodynamics analysis. The electronic band structure shows that the Li<sub>2</sub>OHBr is an insulator with a wide direct band gap. The Li<sub>2</sub>OHBr has a wide electrochemical stability window that can be matched

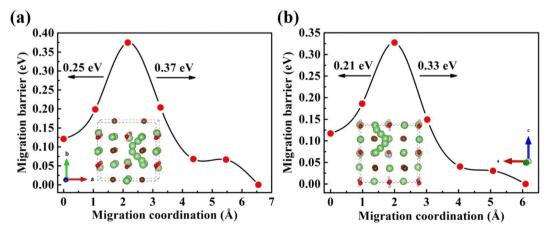


Fig. 7 Diffusion pathway and migration barriers of  $V_{Li^-}$  in F-doped Li<sub>2</sub>OHBr (a) along the *ab*-plane and (b) along the *c*-axis.

with the cathode materials. Moreover, the Li<sub>2</sub>OHBr also exhibits good chemical stability with typical cathode materials and in air. By comparing the defect formation energies in neutral and charged states, it is shown that charged Li<sub>i</sub><sup>+</sup> and V<sub>Li</sub><sup>-</sup> are the most dominant defect types in Li<sub>2</sub>OHBr. The Li<sub>2</sub>OHBr show the low migration barriers by using the CI-NEB calculation, while Fdoped Li<sub>2</sub>OHBr has a lower migration barrier compared with pristine Li<sub>2</sub>OHBr. This work provides insights into the thermodynamic and kinetic process of Li<sub>2</sub>OHBr and demonstrates the potential of computational methods in the efficient design of future ISEs.

## Conflicts of interest

There are no conflicts to declare.

# Acknowledgements

This work was supported by the National Natural Science Foundation of China (12004145), the Jiangxi Provincial Natural Science Foundation (20212BAB214032), the Innovation and Entrepreneurship Training Program for College Students in Jiangxi Province (202210419004), the Key Science and Technology Project of Ji'an City (20211015311), the Science and Technology Research Project of Jiangxi Provincial Department of Education (GJJ201030) and the PhD Start-up Fund of Natural Science Foundation of Jinggangshan University (JZB2013). All the calculations were supported by the high-performance computing platform of Shanghai University and Jinggangshan University.

## References

- 1 T. Famprikis, P. Canepa, J. A. Dawson, M. S. Islam and C. Masquelier, Nat. Mater., 2019, 18, 1278-1291.
- 2 X.-B. Cheng, R. Zhang, C.-Z. Zhao and Q. Zhang, Chem. Rev., 2017, 117, 10403-10473.
- 3 H. Xu, Y. Yu, Z. Wang and G. Shao, Energy Environ. Mater., 2019, 2, 234-250.
- 4 N. Kamaya, K. Homma, Y. Yamakawa, M. Hirayama, R. Kanno, M. Yonemura, T. Kamiyama, Y. Kato, S. Hama, K. Kawamoto and A. Mitsui, Nat. Mater., 2011, 10, 682-686.
- 5 H. Yamane, M. Shibata, Y. Shimane, T. Junke, Y. Seino, S. Adams, K. Minami, A. Hayashi and M. Tatsumisago, Solid State Ionics, 2007, 178, 1163-1167.
- 6 Z. Deng, D. Ni, D. Chen, Y. Bian, S. Li, Z. Wang and Y. Zhao, InfoMat, 2022, 4, e12252.
- 7 J. Zheng, B. Perry and Y. Wu, ACS Mater. Au, 2021, 1, 92–106.
- 8 Y. Zhao and L. L. Daemen, J. Am. Chem. Soc., 2012, 134, 15042-15047.
- 9 A. Emly, E. Kioupakis and A. Van der Ven, Chem. Mater., 2013, 25, 4663-4670.
- 10 M. K. Sugumar, T. Yamamoto, M. Motoyama and Y. Iriyama, Solid State Ionics, 2020, 349, 115298.

- 11 K. Yoshikawa, T. Yamamoto, M. K. Sugumar, M. Motoyama and Y. Iriyama, Energy Fuels, 2021, 35, 12581-12587.
- 12 L. Gao, R. Zhao, S. Han, S. Li, R. Zou and Y. Zhao, Batteries Supercaps, 2021, 4, 1491-1498.
- 13 G. Kresse, Phys. Rev. B: Condens. Matter Mater. Phys., 1996, 54, 11169-11186.
- 14 J. P. Perdew, K. Burke and M. Ernzerhof, Phys. Rev. Lett., 1996, 77, 3865-3868.
- 15 S. P. Ong, W. D. Richards, A. Jain, G. Hautier, M. Kocher, S. Cholia, D. Gunter, V. L. Chevrier, K. A. Persson and G. Ceder, Comput. Mater. Sci., 2013, 68, 314-319.
- 16 B. Liu, D. Wang, M. Avdeev, S. Shi, J. Yang and W. Zhang, ACS Sustainable Chem. Eng., 2020, 8, 948-957.
- 17 B. Liu, J. Liu, J. Yang, D. Wang, C. Ye, D. Wang, M. Avdeev, S. Shi, J. Yang and W. Zhang, J. Power Sources, 2020, 450, 227693.
- 18 S. B. Zhang and J. E. Northrup, Phys. Rev. Lett., 1991, 67, 2339-2342.
- 19 J. Howard and N. A. W. Holzwarth, Phys. Rev. B, 2019, 99,
- 20 J. Sun, A. Ruzsinszky and J. P. Perdew, Phys. Rev. Lett., 2015, **115**. 036402.
- 21 Y. Li, W. Zhou, S. Xin, S. Li, J. Zhu, X. Lü, Z. Cui, Q. Jia, J. Zhou, Y. Zhao and J. B. Goodenough, Angew. Chem., Int. Ed., 2016, 55, 9965-9968.
- 22 Z. Deng, Z. Zhu, I.-H. Chu and S. P. Ong, Chem. Mater., 2017, 29, 281-288.
- 23 I.-H. Chu, H. Nguyen, S. Hy, Y.-C. Lin, Z. Wang, Z. Xu, Z. Deng, Y. S. Meng and S. P. Ong, ACS Appl. Mater. Interfaces, 2016, 8, 7843-7853.
- 24 W. D. Richards, L. J. Miara, Y. Wang, J. C. Kim and G. Ceder, Chem. Mater., 2016, 28, 266-273.
- 25 Y. Zhu, X. He and Y. Mo, ACS Appl. Mater. Interfaces, 2015, 7, 23685-23693.
- 26 M. B. Effat, J. Liu, Z. Lu, T. H. Wan, A. Curcio and F. Ciucci, ACS Appl. Mater. Interfaces, 2020, 12, 55011-55022.
- 27 V. Lacivita, Y. Wang, S.-H. Bo and G. Ceder, J. Mater. Chem. A, 2019, 7, 8144-8155.
- 28 A. Banerjee, X. Wang, C. Fang, E. A. Wu and Y. S. Meng, Chem. Rev., 2020, 120, 6878-6933.
- 29 W. Xia, Y. Zhao, F. Zhao, K. Adair, R. Zhao, S. Li, R. Zou, Y. Zhao and X. Sun, Chem. Rev., 2022, 122, 3763-3819.
- 30 J. Zhang, L. Huang and X. Gu, Mater. Adv., 2022, 3, 3143-3150.
- 31 W. Li, J. Liang, M. Li, K. R. Adair, X. Li, Y. Hu, Q. Xiao, R. Feng, R. Li, L. Zhang, S. Lu, H. Huang, S. Zhao, T.-K. Sham and X. Sun, Chem. Mater., 2020, 32, 7019-7027.
- 32 S. Shi, P. Lu, Z. Liu, Y. Qi, L. G. Hector, H. Li and S. J. Harris, J. Am. Chem. Soc., 2012, 134, 15476-15487.
- 33 B. Liu, Q. Hu, T. Gao, P. Liao, Y. Wen, Z. Lu, J. Yang, S. Shi and W. Zhang, Journal of Materiomics, 2022, 8, 59-67.