


 Cite this: *RSC Adv.*, 2022, 12, 29022

Analysis of the growth and physicochemical properties of the newly developed stilbazolium derivative 4-*N,N*-dimethylamino-4-*N*-methyl stilbazolium 2-formyl benzene sulfonate (DSFS) single crystal: an effective material for nonlinear optical applications†

 Sekar Anand,^a Muthurakku Usha Rani,^{ID} *^a Sivaperuman Kalainathan^{ID} ^b and Ravi Shanker Babu^a

A novel ionic stilbazolium derivative single crystal of 4-*N,N*-dimethylamino-4-*N*-methyl stilbazolium 2-formyl benzene sulfonate (DSFS) was successfully cultivated with methanol as the solvent by using the slow evaporation technique. Structural confirmation was done using single-crystal X-ray diffraction (SCXRD), and the obtained results reveal that the DSFS crystal crystallized in a centrosymmetric pattern with $P\bar{1}$ space group. The presence of different vibrational modes in the DSFS crystal is affirmed via Fourier transform infrared (FTIR) spectroscopy. Linear optical constants such as cut-off wavelength, bandgap, extinction coefficient, Urbach energy, electrical conductivity and optical conductivity of the titular crystal were found from ultraviolet–visible–near infra-red spectroscopy (UV-Vis-NIR). The emission wavelength of the title crystal lies in the red region (620 nm), which was confirmed from the photoluminescence spectroscopy (PL). The surface nature of the grown ionic crystal was examined through the etching and atomic force microscopy (AFM) technique. With a He–Ne laser as the source, Z scan analysis was carried out to study the third-order nonlinear properties of the DSFS crystal.

 Received 24th August 2022
 Accepted 16th September 2022

DOI: 10.1039/d2ra05302k

rsc.li/rsc-advances

1 Introduction

In this technology-driven modern world, Nonlinear Optical (NLO) materials find their place due to their significant applications across various fields. The study of organic NLO crystals is a scholastic area of interest due to their ample applications in different domains, such as signal processing, electro-optic modulators, telecommunication, optical switches, deep UV light, high-density data storage and display technologies.^{1–8} The presence of delocalized π -electrons, donor and acceptor moieties, hydrogen bonds and weak Van der Waals forces induce high non-linearity in the organic crystals compared to its counterpart, inorganic materials.⁹ Also, tunability of the structure, high laser damage threshold, ultra-fast response and the significant electro-optic coefficient are the added advantages of the organic materials.¹⁰ Usually, organic crystals that possess

stable packing and high chromophore density tend to exhibit high non-linearity, high mechanical and photo-chemical stability and thermal properties.¹¹ Among the reported NLO crystals, 4-*N,N*-dimethylamino-4-*N*-methyl-stilbazolium tosylate (DAST) has been one of the yardstick crystals which has second harmonic generation (SHG) efficiency that is roughly 1000 times that of urea.^{12,13} One of the fundamental reasons behind this high nonlinearity among the DAST derivative crystals is the presence of the stilbazolium chromophore in the cation component. The presence of an ethylenic bridge (C=C) in the stilbazolium cation is one of the primary reasons behind the high NLO response of this derivative's materials, as this double bond acts as an interconnecting bridge between the donor and acceptor moiety. To form this carbon–carbon ethylenic bridge, many reaction strategies have been reported in the past, namely the Suzuki–Miyaura reaction, Horner–Wadsworth–Emmons (HWE) condensations and the Knoevenagel condensation. Among these methods, Knoevenagel condensation is used widely due to its simplicity.^{14–16} In the present study, the strategy of combining the stilbazolium (chromophore) cation with a new anion is adopted to develop a novel ionic organic single crystal. This approach of incorporating new counter-anions with an already existing stilbazolium cation has already been reported.¹⁷

^aSchool of Advanced Sciences, Vellore Institute of Technology, Vellore – 632 014, India. E-mail: usharani.m@vit.ac.in

^bCentre for Nanotechnology Research, Vellore Institute of Technology, Vellore – 632 014, India

† Electronic supplementary information (ESI) available. CCDC 2152094. For ESI and crystallographic data in CIF or other electronic format see <https://doi.org/10.1039/d2ra05302k>



This approach assists in creating a new NLO crystal with the desired physical and chemical properties.¹⁸ Previously reported results show that even a small change in the anion can induce a huge change in the non-linearity of the material by affecting the crystal stacking.¹⁹ The nature of the hydrogen bonds which connect the anion and cation decides the optoelectronic properties of the synthesized material.²⁰ The low toxicity level of the styryl pyridinium compound in the cation part makes it an appropriate material for applications like antibacterial drugs and disinfectants in hospital and food industries.^{21,22} The anion compound 2-formyl benzene sulfonate sodium salt is used to test the ability of fungal strains to change phenolic and non-phenolic precursors into stable and non-toxic dyes.²³ The findings from the extensive literature survey suggests that, so far, no crystal has been reported with 2-formyl benzene sulfonate as the anion with a stilbazolium derivative. So, in this article, we report the growth of DSFS crystal for the first time by employing the slow evaporation method. In addition, the synthesis and various physical and chemical properties such as structural, linear optical, surface analysis and nonlinear optical (NLO) properties of the grown crystal were inspected. The results are scrupulously discussed in the following sections. All the outcomes indicate that the grown DSFS crystal is a favourable candidate for NLO applications.

2 Materials and synthesis

All the required reagents were bought from Tokyo Chemical Limited (TCL). All the reagents were used without further purification. The overall synthesis process consists of three steps. Reactions involved in the synthesis process are explained in Fig. 1. Step 1: 1,4-dimethyl pyridinium iodide (1). Both the starting materials 4-picoline (2 ml) and methyl iodide (1.3 ml) were taken in an equimolar (20 mmol) ratio in an Erlenmeyer

bulb with acetone as the medium. The above mixture solution is placed in a reflux setup at about 60 °C for 2 hours. After 2 hours, the solution is left to cool down to the surrounding temperature and the resultant precipitate is filtered and dried in an oven. The final dried powder is 1,4-dimethyl pyridinium iodide (1). Step 2: 4-[2-(4-dimethylamino-phenyl)-vinyl]-1-methyl-pyridinium iodide (2). The final product of step 2 (2) is obtained *via* the Knoevenagel condensation reaction. Compound (1) (2.35 g) and 4-(dimethylamino)benzaldehyde (1.5 g) were dissolved in methanol solvent in an RB. Generally, weak bases like piperidine or ethylenediamine are used in the Knoevenagel condensation.²⁴ In the present study, piperidine is added to the solution mixture. Piperidine speeds up the reaction by providing its lone pair of electrons to complete the reaction. As a consequence of this condensation process, a carbon-carbon double bond is formed with water as a by-product.²⁴ Then, the prepared solution is refluxed for 12 hours at 60 °C using a modified Dean-Stark apparatus equipped with a heating mantle. At the end of the reaction, the Erlenmeyer bulb is allowed to cool down to room temperature. The unreacted initial reactants were removed from the final purple precipitate by washing it with diethyl ether solution. Then, the final powder product (2) is dried in an oven at 100 °C for 1 h. Methanol solvent is used for the recrystallization process to enhance the purity of the final powder. Step 3: 4-*N,N*-dimethylamino-4-*N*-methyl stilbazolium 2-formyl benzene sulfonate (3). The final metallization reaction occurs between compound (2) and 2-sodium formyl benzene sulfonate. Both the materials were dissolved separately in a hot aqueous solution and kept on the stirrer for 1 hour at 70 °C. Then both the solutions were mixed, and the mixture was left undisturbed for a whole day so that the residues could settle down. Finally, the obtained reddish-brown precipitate is kept in an oven at about 100 °C for 2 h to get rid of the water content. The final obtained powder is *N,N*-dimethylamino-4-*N'*-methylstilbazolium 2-formyl benzene sulfonate (DSFS). Fig. 1 portrays the reaction scheme involved in the formation of the DSFS material.

2.1 Solubility and crystal growth

The gravimetric method was chosen to find the solubility of the pure DSFS powder in methanol solvent using an ultra-cryostat (± 0.01 °C accuracy). The experiment was carried out for various temperatures, ranging from 25–45 °C with an interval of 5 °C (25, 30, 35, 40 and 45 °C). Saturated solutions were prepared at the aforementioned temperatures to study the solubility of the synthesized powder. Fig. 2 shows the solubility curve of the DSFS powder in methanol solvent. From the graph, it is clear that with an increase in temperature the solubility of the DSFS material also increases. This increase in solubility with respect to the rise in temperature is due to the rise in the average kinetic energy of the solution and solute molecules. This increase in kinetic energy allows the solvent molecules to break more solute molecules. Also, the rise in the solute molecule's kinetic energy destabilizes the solid state of the material, so more solute molecules dissolve readily at higher temperature.²⁵

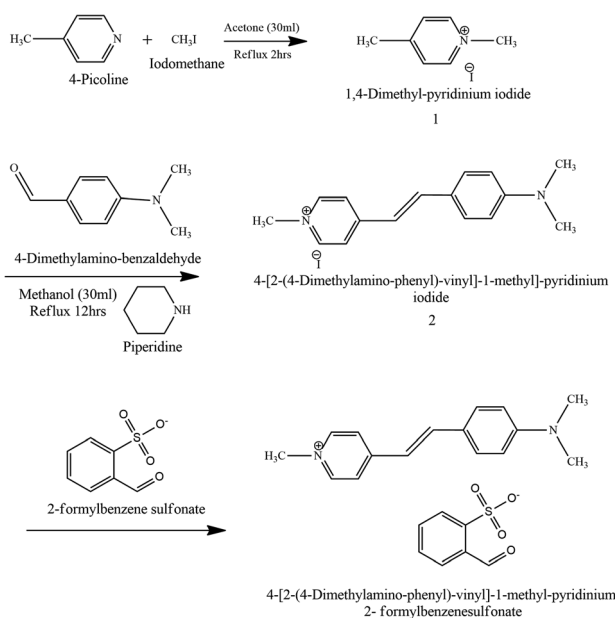


Fig. 1 Reaction scheme for the DSFS material.



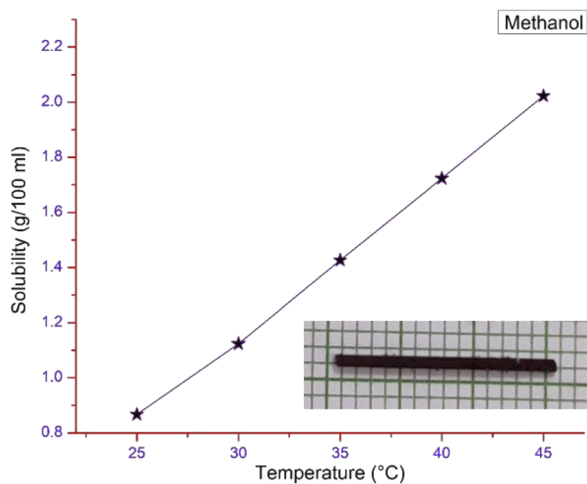


Fig. 2 Solubility curve and grown crystal image of DSFS.

The solubility study confirms that methanol is a suitable solvent for the crystal growth process. So, with methanol as the solvent, the crystal growth process is carried out by the slow evaporation method. Within the span of 20–25 days, a DSFS crystal with the dimensions of $30 \times 1 \times 1 \text{ mm}^3$ was obtained and the image of the grown title crystal is portrayed in the inset of Fig. 2.

3 Results and discussion

3.1 Structural studies

3.1.1 Single-crystal X-ray diffraction (SCXRD). To explore the structure of the title DSFS crystal, a single crystal of size $0.080 \times 0.140 \times 0.210 \text{ mm}^3$ was exposed to single crystal XRD. The molecular structure, interatomic interactions, packing and other structural properties of the grown crystal were interpreted through this study. Under the room temperature conditions, analysis was carried out using a D8 QUEST BRUKER. With the help of Apex 4 software, the obtained raw crystallographic data was solved and refined. Non-hydrogen atoms in the structure were positioned using the full-matrix least-squares refinement process by utilizing the SHELXL function.²⁶ Obtained crystallographic data shows that the DSFS crystallized in a centrosymmetric pattern with $(P\bar{1})$ space group with triclinic crystal system. Collected lattice constant values of the grown DSFS crystal are $a = 8.3812 \text{ \AA}$, $b = 9.8732 \text{ \AA}$, $c = 13.0649 \text{ \AA}$ and interfacial angle values are $\alpha = 82.404^\circ$, $\beta = 88.395^\circ$, $\gamma = 86.109^\circ$. Refined structural and crystallographic data of the title ion pair crystal are displayed in Table 1. Other important structural parameters, bond length and bond angle of selected bonds are depicted in Tables 1 and 2 in the ESI file.†

The title ion-pair crystal's ($\text{C}_{23}\text{H}_{24}\text{N}_2\text{O}_4\text{S}$) asymmetric unit comprises a positively charged stilbazolium cation and negatively charged sulfonate anion. The ORTEP diagram of the novel DSFS crystal is portrayed in Fig. 3. A partly viewed packing image of the novel DSFS crystal along the 'a' axis is depicted in Fig. 4. The cation's carbon-carbon (C=C) double bond indicates that it exhibits a *trans* or *E* configuration, with a torsional

angle of $\text{C}_1\text{-C}_6\text{=C}_7\text{-C}_8$ of 179.59° .²⁷ The torsional angle of the dimethyl group $\text{C}_{12}\text{-C}_{11}\text{-N}_2\text{-C}_{15}$ (5.83°) attached with the amino group $\text{C}_{10}\text{-C}_{11}\text{-N}_2\text{-C}_{16}$ (-9.70°) shows that the dimethyl moiety is slightly distorted from coplanarity with the benzene ring. The presence of the stilbazolium chromophore in the cation is the key to higher non-linearity. The main reason behind this efficiency is the deviation in the bond length among the ionic atoms and the presence of the D- π -A bridge, which in turn stimulates NLO susceptibility in the material.²⁴ NLO response of a chromophore also depends on the bond length alternation (BLA) degree. The BLA degree of the title cation for the bonds $\text{N}_1\text{-C}_{14}$ (1.48 \AA) and $\text{N}_2\text{-C}_{11}$ (1.37 \AA) is 0.11 \AA . The difference in the bond lengths of the dimethyl moiety $\text{N}_2\text{-C}_{15}$ (1.433 \AA) and $\text{N}_2\text{-C}_{16}$ (1.453 \AA) is 0.02 \AA . Similarly, the bond length difference between the bonds of $\text{N}_1\text{-C}_3$ (1.336 \AA) and $\text{N}_1\text{-C}_4$ (1.346 \AA) of the 4-picoline compound is 0.010 \AA . The ideal BLA value to exhibit higher non-linearity is approximately $\pm 0.05 \text{ \AA}$, whereas, in the case of the title cation, it is 0.11 \AA , which is slightly higher.²⁸ This BLA value indicates that the title stilbazolium cation compound also exhibits molecular nonlinearity. The geometrical description of the SO_3^- anion shows that the S^- unit adopts regular trihedral coordination with the three oxygen atoms through S-O bonds. The bond lengths and angles of the S-O and O-S-O interactions reside in the range of $1.4415\text{--}1.4416 \text{ \AA}$ and $113.09\text{--}114.12^\circ$. The three-dimensional network of the title ionic compound is made up of alternating perpendicular anions between the parallel cationic units. The CIF file of the grown crystal is deposited in the CCDC (CCDC no. 2152094).

3.1.2 Morphology. The morphology of the fully-grown DSFS crystal was predicted with the aid of the obtained single-crystal XRD data (CIF file) by using Winx Morph software.²⁹ The predicted morphology of the title compound with different crystal faces is illustrated in Fig. 5. From the morphology diagram, it is apparent that the growth rate is dominant along the 'a' axis.

3.1.3 Density measurement. Theoretical and experimental approaches were used to calculate the density of the novel DSFS crystal. Theoretically, the density of the new ionic crystal can be calculated using the single crystal XRD data with the help of the following equation:

$$\rho = \frac{MZ}{N_A V} \quad (1)$$

where all the values M (molecular weight - $424.50 \text{ g mol}^{-1}$), Z (number of molecules per unit cell - 2), and V (volume of the unit cell - 1068.96) can be directly obtained from the SCXRD results. N_A is Avogadro's number. The theoretically predicted density value for the DSFS crystal is 1.319 g cm^{-3} . Experimentally the density was measured using the floating technique.³⁰ With CCl_4 as the solvent, this measurement was carried out. Experimentally, the density value is derived from eqn (2)

$$\rho = \frac{m_0 \rho_{\text{solvent}}}{(m_0 - m')} \quad (2)$$

where m_0 and m' are the mass of the grown DSFS crystal in air and solvent (CCl_4) medium, ρ_{solvent} is the density of the CCl_4 solvent. The calculated density values of the theoretical



Table 1 Crystal data and structure refinement for DSFS

Identification code	DSFS	
Empirical formula	C ₂₃ H ₂₄ N ₂ O ₄ S	
Formula weight	424.50	
Temperature	300(2) K	
Wavelength	0.71073 Å	
Crystal system	Triclinic	
Space group	<i>P</i> $\bar{1}$	
Unit cell dimensions	<i>a</i> = 8.3812(2) Å <i>b</i> = 9.8732(6) Å <i>c</i> = 13.0649(11) Å	α = 82.404(3)° β = 88.395(3)° γ = 86.109(3)°
Volume	1068.96(15) Å ³	
Z	2	
Density (calculated)	1.319 g cm ⁻³	
Absorption coefficient	0.183 mm ⁻¹	
<i>F</i> (000)	448	
Crystal size	0.080 × 0.140 × 0.210 mm ³	
Theta range for data collection	1.57 to 28.38°	
Index ranges	-11 ≤ <i>h</i> ≤ 11, -13 ≤ <i>k</i> ≤ 13, -17 ≤ <i>l</i> ≤ 17	
Reflections collected	39 474	
Independent reflections	5323 [<i>R</i> (int) = 0.0516]	
Completeness to theta = 28.38°	100.0%	
Absorption correction	Multi-scan	
Max. and min. transmission	0.9850 and 0.9620	
Refinement method	Full-matrix least-squares on <i>F</i> ²	
Data/restraints/parameters	5323/0/274	
Goodness-of-fit on <i>F</i> ²	1.080	
Final <i>R</i> indices [<i>I</i> > 2σ(<i>I</i>)]	<i>R</i> 1 = 0.0487, <i>wR</i> 2 = 0.1256	
<i>R</i> indices (all data)	<i>R</i> 1 = 0.0798, <i>wR</i> 2 = 0.1556	
Largest diff. peak and hole	0.264 and -0.307 e Å ⁻³	

(1.319 g cm⁻³) and experimental (1.317 g cm⁻³) approaches agree with each other.

3.2 Fourier transform infra-red (FTIR) analysis

Structural confirmation of the title ionic crystal can be done with the help of functional groups and their corresponding vibrations by adopting FT-IR analysis. The KBr pellet technique is utilized to obtain the FTIR spectrum within the range of 3100 to 400 cm⁻¹, using a SHIMADZU instrument. The observed vibration spectrum of the synthesized material is displayed in Fig. 6. The title DSFS crystal consists of stilbazolium cation and formyl benzene sulfonate anion. The vibrational spectrum of the cation is attributed to the vinyl group, methyl group, pyridine group and phenyl ring. Similarly, in the case of the anion, the aldehyde and sulfonate groups' vibrations were analyzed.

3.2.1 Cation group absorption. The characteristic peak for this vinyl (C=C) stretching lies in the range of 1650 to 1600 cm⁻¹.³¹ The band at 1643 cm⁻¹ in the spectrum indicates the presence of vinyl stretching. Apart from the (C=C) stretching, stilbazolium derivatives show aliphatic C-H stretching in the vicinity of 3050 cm⁻¹. The peak centered at 3035 cm⁻¹ indicates the aliphatic C-H stretching. Generally, the methyl group's asymmetrical bending vibration and symmetrical bending vibration bands were positioned in and around 1450 and 1375 cm⁻¹, respectively. In Fig. 6, the peaks at 1473 and 1429 cm⁻¹ represent the asymmetrical bending vibrations. The peak observed at 1375 cm⁻¹ characterizes the symmetrical bending vibration of the methyl group. In the region between 1070–1010 cm⁻¹, methyl group rocking modes were observed.³² In the DSFS crystal, the peaks positioned at 1078 and 1014 cm⁻¹

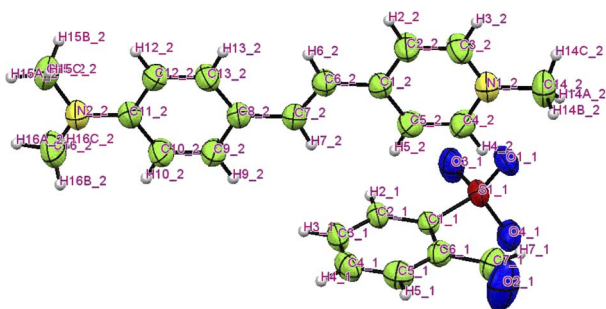


Fig. 3 ORTEP image of the DSFS crystal.

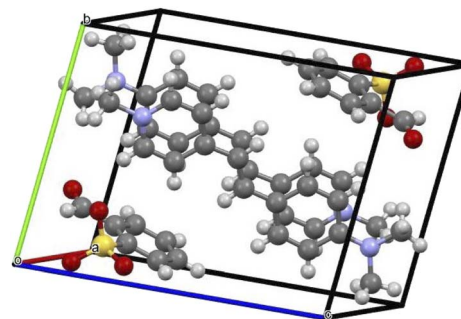


Fig. 4 Packing image of the DSFS crystal.



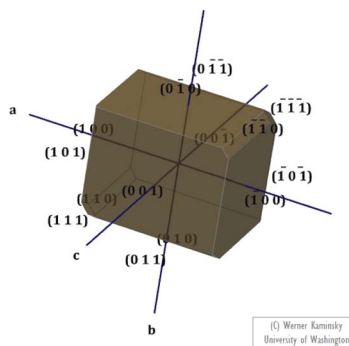


Fig. 5 Morphology of the DSFS crystal.

are attributed to the CH_3 rocking mode. The peak emitted at 1568 cm^{-1} is ascribed to aromatic C–C stretching. The peak positioned at 1527 cm^{-1} indicates ring (C=C) vibrations. Characteristic peaks for the aromatic C–H in-plane bending are observed at 1211 and 1078 cm^{-1} . Similarly, peaks observed at 781 and 721 cm^{-1} are assigned to the C–H out-of-plane bending vibrations.³¹ The peak centered at 887 cm^{-1} represents the 1,4-disubstituted aromatic C–H ring. Three peaks centered at 611 , 565 and 528 cm^{-1} in the spectrum indicate the phenyl ring mode (C–C–C) deformation vibration.

3.2.2 Anion group absorption. The anionic group of the DSFS crystal consists of SO_2 asymmetric and symmetric stretching, C–S stretching, SO_3 scissoring and SO_3 deformation vibrations. Peaks found at 1317 and 1014 cm^{-1} denote asymmetric and symmetric SO_2 stretching vibrations, respectively. The characteristic peak of the sulfonate group (S=O) stretching appeared at 1375 cm^{-1} . The sharp peak positioned at 565 cm^{-1} implies C–S stretching, SO_3 scissoring and SO_3 deformation vibrations were observed at 528 and 493 cm^{-1} , respectively. Finally, C=O vibrations of the aldehyde group in the anion group give the characteristic peak at 1687 cm^{-1} .

3.3 Linear optical studies

3.3.1 Ultraviolet-visible-near infra-red spectroscopy (UV-Vis-NIR). The response of the grown crystal in the Ultraviolet

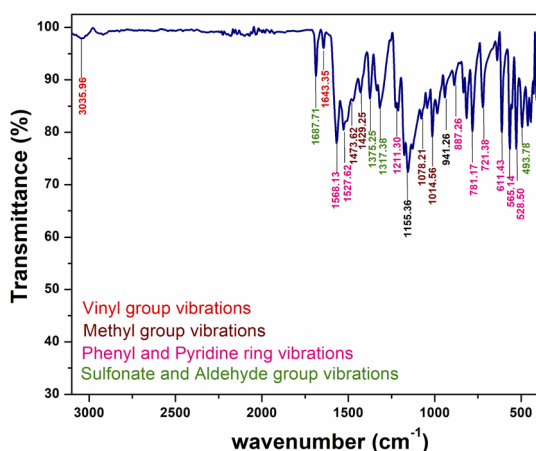


Fig. 6 FTIR spectrum of the grown DSFS crystal.

(200–400 nm), visible (400–800 nm) and near-infrared (800–1000) regions gives detail about the absorption edge and the orbital transitions.³³ To investigate the optical characteristics of the DSFS crystal in the range of 200–1000 nm, a 2 mm thick crystal is placed in an Analytik Jena Specord 210 Plus instrument. The absorbance spectrum of the ionic DSFS crystal is depicted in Fig. 7. From the plot, it is apparent that the DSFS material is opaque in the UV region (200–400 nm). The absorption edge of the DSFS crystal is 570 nm. DAST and DSFS are both crystals belonging to the same stilbazolium family. So, both the materials exhibit a similar kind of absorbance spectrum. When compared with DAST crystal ($\lambda_{\text{max}} = 550\text{ nm}$), DSFS crystal exhibits a slightly higher absorption edge of about 20 nm (redshift). It is because of the stilbazolium chromophore; in general the stilbazolium chromophore exhibits a redshift when it has higher molecular nonlinearity.³⁴ The bandgap of the grown crystal is obtained with the help of a standard Tauc plot relationship,³⁵ and the calculated E_g value is 2.25 eV. The graph between $(\alpha h\nu)^2$ and $h\nu$ is shown in Fig. 8.

In order to fabricate a potential optoelectronic device, knowledge about optical constants such as the coefficient of absorption (α), extinction coefficient (K), optical conductivity (σ_{opt}), and electrical conductivity (σ_{elc}) values are vital.

Fig. 9 represents the coefficient of the absorption spectrum of the DSFS crystal. Beyond 570 nm (cut-off wavelength), the absorption coefficient starts to decrease gradually. The extinction coefficient is the measure of light loss within a medium while propagating due to absorption and scattering phenomena.³⁶ From the absorption coefficient (α), the K value can be deduced by using the following relation:³⁷

$$K = \frac{\alpha\lambda}{4\pi} \quad (3)$$

Fig. 10 elucidates the variation of the extinction coefficient with respect to the photon energy. There is a steady decrease in the extinction coefficient value until 1.65 eV. Then it gradually increases and attains the maximum K value at 2.15 eV, which ratifies the bandgap region ($E_g = 2.25\text{ eV}$) of the title ionic

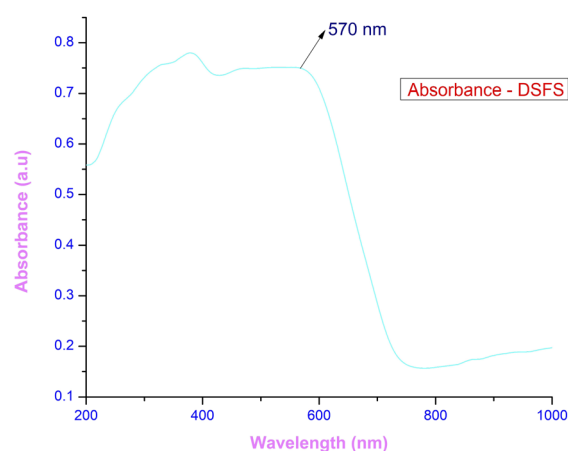


Fig. 7 Absorbance spectrum of the DSFS crystal.



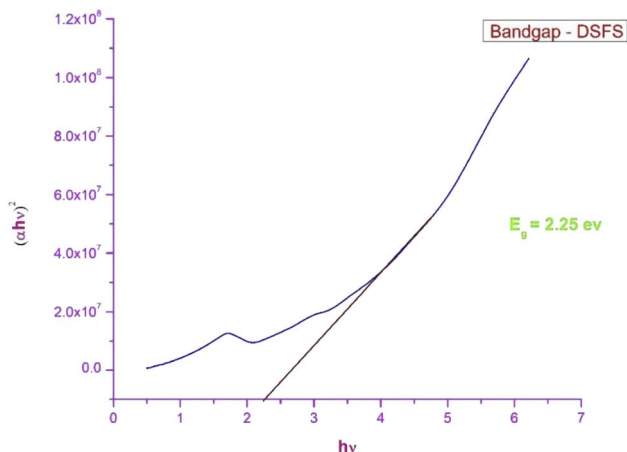


Fig. 8 Tauc plot of the DSFS crystal.

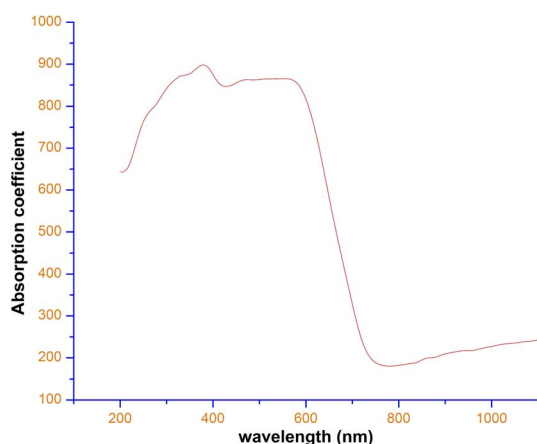


Fig. 9 Absorption coefficient vs. wavelength plot of DSFS crystal.

crystal. Along with that, the low K value signifies that the grown ionic crystal can be used for the construction of devices, which are ideal for information processing and computing technology.³⁸

3.3.2 Optical conductivity. Optical conductivity (σ_{opt}) is the measure of the response of the material to the incident light frequency and it can be derived from eqn (4):³⁹

$$\sigma = \frac{\alpha nc}{4\pi} \quad (4)$$

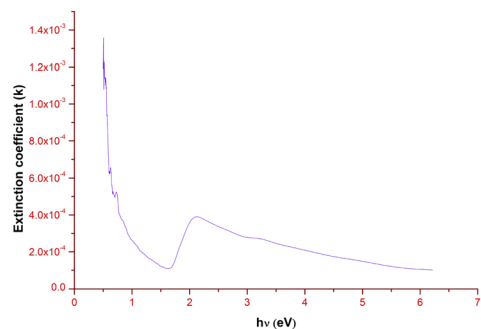


Fig. 10 Extinction coefficient vs. photon energy plot of DSFS crystal.

where α , c and n indicate the coefficient of absorption, light speed in air medium and refractive index, respectively. Fig. 11 indicates the disparity between the optical conductivity and its corresponding photon energy ($h\nu$). From the graph, it is obvious that with a rise in $h\nu$ value, there is a decrease in optical conductivity up to 1.6 eV. After that, there is a steep increase in the σ_{opt} in the bandgap region (1.61 to 2.21 eV). This increase in the optical conductivity in the band gap region is due to the rise in the optical absorption coefficient in this region.⁴⁰ The maximum (σ_{opt}) value is found to be $5 \times 10^{10} \text{ S}^{-1}$. This high value of optical conductivity (σ_{opt}) suggests that the synthesized crystal has a high photo response.⁴¹ Beyond 3.6 eV, the value of optical conductivity decreases due to a drop in α value.

With the help of σ_{opt} , the title crystal's electrical conductivity is calculated by using the following equation:

$$\sigma_{\text{elc}} = \frac{2\lambda\sigma_{\text{opt}}}{\alpha} \quad (5)$$

Fig. 11 portrays the disparity of electrical conductivity (σ_{elc}) with respect to photon energy. Fig. 11 shows that σ_{elc} decreases with an increase in $h\nu$. This drop in the electrical conductivity at higher photon energy indicates the semiconducting nature of the material.⁴²

The electrical susceptibility (χ_c) of the DSFS crystal is calculated at $\lambda = 570 \text{ nm}$, with the help of the already obtained optical constant value, from the following formula:⁴³

$$\chi_c = \frac{\mu^2 - k^2 - \epsilon_0}{4\pi} \quad (6)$$

where μ , ϵ_0 and k are the refractive index, vacuum permittivity and extinction coefficient, respectively. The χ_c value is found to be 0.4375 at 570 nm. By incorporating the calculated values of the coefficient of extinction and refractive index, real (ϵ_r) and imaginary (ϵ_i) dielectric constant values were elucidated from the below equation:

$$\epsilon_r = n^2 - K^2 \quad \text{and} \quad \epsilon_i = 2nk$$

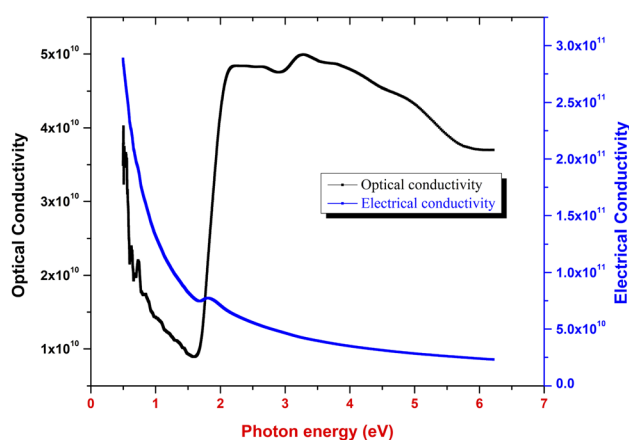


Fig. 11 Optical conductivity and electrical conductivity vs. photon energy plot.



Information about the grown material's electrical energy storage capacity and ohmic resistance can be obtained from the real and imaginary dielectric constants.⁴² With respect to the cut-off wavelength ($\lambda = 570$ nm), the calculated ϵ_r and ϵ_i values are 5.495 and 1.817×10^{-3} , respectively.

3.3.3 Urbach energy. Interaction of phonons in the bandgap region often ends up with the development of band tailing in the crystalline material because of the presence of disorders. From the following equation, it is possible to calculate the disorder in the system as a function of α , which corresponds to the transition between the localized states and the extended state of the conduction and valence bands, through a parameter known as the Urbach energy (E_u).^{41,44}

$$\alpha(h\nu) = \alpha_0 \exp\left(\frac{h\nu}{E_u}\right) \quad (7)$$

where h , ν , α_0 , and E_u represent Planck's constant, the incident frequency radiation, a constant and Urbach energy, respectively. The width of the elongated tail of the localized states in the forbidden electronic band gap (Urbach energy or band tail energy) was obtained by plotting a graph between the \ln of the absorption coefficient α vs. $h\nu$. The obtained slope value of the linear portion of the graph is 3.686. The Urbach energy can be calculated by taking the reciprocal of the obtained slope value and it is found to be 0.2712 eV. This low E_u value confirms that the novel DSFS ionic crystal possesses high crystallinity with minimal structural disorders. Fig. 12 depicts the Urbach energy graph.

From Table 2, we can conclude that the grown crystal possesses a wide transparency window in the Vis-NIR region, low extinction coefficient, high electrical and optical conductivity and few disorders (Urbach energy). These results indicate the grown crystal's appropriateness for optoelectronic, NLO and photonic applications.

3.4 Photoluminescence studies

Physical properties of the grown material, such as electronic transition, shallow and deep level defects and band gap, can be

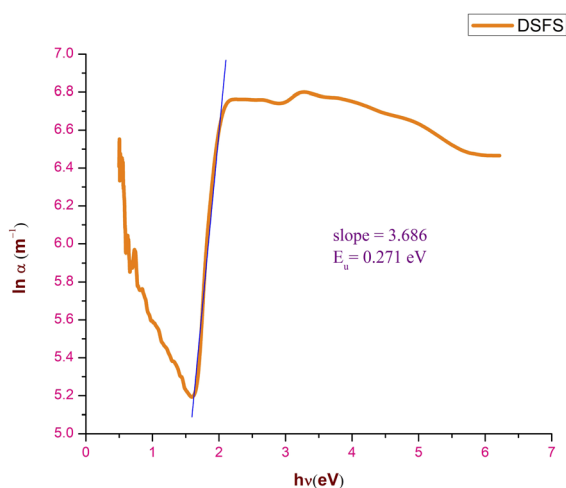


Fig. 12 Urbach energy plot of the DSFS crystal.

Table 2 Calculated optical parameters of the DSFS crystal

Optical constants	Calculated values
Cut-off wavelength (λ)	570 nm
Band gap (E_g)	2.25 eV
Extinction coefficient (K) @2.25 eV	3.755×10^{-4}
Optical conductivity (σ_{opt}) @2.25 eV	4.839×10^9
Electrical conductivity (σ_{elc}) @2.25 eV	6.176×10^{10}
Urbach energy (E_u)	0.271 eV

analyzed using photoluminescence studies.^{33,45} In general, compounds with aromatic rings and extended conjugated double bonds with a high degree of resonance stability exhibit a luminescence nature.¹² Upon irradiance of light, an electron-hole pair is generated. Due to very low relaxation time (τ), recombination of these excited electron-hole pairs takes place with the emission of photons.⁴⁶ Fig. 13 shows the output PL spectrum of the grown DSFS crystal obtained in the liquid state with methanol solvent using an F-7000 FL spectrophotometer, operated at room temperature. With 476 nm as the excitation wavelength, scanning was done between 500–800 nm. Only one broad peak centred at 620 nm (2.01 eV) confirms the absence of impurities. The stilbazolium chromophore is responsible for the emission of photons in the red colour region. Along with that, the existence of the electron donor and acceptor groups contributes to the red emission. The low quantum yield observed in the PL spectrum suggests that there is no direct electronic transition between the conduction and valence bands. Most probably, the electronic transition associated with the radiative recombination is between the trapped holes and electrons that occur between the localized energy states that existed within the bandgap region.³³ Various electronic transitions occurring within the bandgap region between different energy levels is responsible for the broadening of the spectrum. These transitions confirm that the shallow holes contribute more than deep holes toward the charge transfer process and also confirms the trapping of electrons.³³ The broad red emission peak indicates that the grown DSFS crystal is a promising

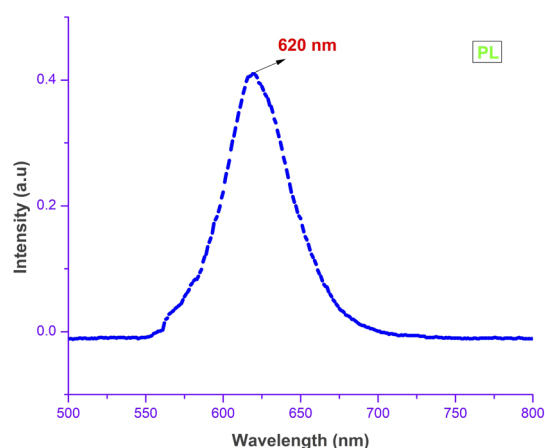


Fig. 13 PL spectrum of the DSFS crystal.



aspirant for various applications such as bandpass filters, optical switches in the near IR region and solid-state lasers. Also, stilbazolium salts can be used as fluorescence probe monitors in the biomedical research field.⁴⁷

3.5 Surface analysis

3.5.1 Etching studies. Quality and crystalline perfection are the essential factors when it comes to device fabrication. To assess the surface quality of the DSFS crystal, it is subjected to etching analysis. With the help of the developed etch pits, the growth mechanism involved in the crystal growth process and the defects can be analysed. A flat and optically good quality crystal is used for the chemical etching analysis with methanol + acetonitrile as the etchant. In the present study, etching was carried out at room temperature for two different etching times, 10 and 20 seconds. Initially, the DSFS crystal's grown surface is subjected to an optical microscope, and the output image is captured (Fig. 14a). The grown crystal is immersed in the etching solvent (10 s). Afterwards; the surface of the etched crystal is cleaned gently with a tissue to exclude the etchant from the surface. The exact process is repeated for 20 s, and the captured pictures of the surface are presented in (Fig. 14b and c). Rectangular-shaped etch pits were observed on the surface after 10 s (Fig. 14b). With the increased etching time (20 s), the etch pits are more precise and more visible in the dislocation site (Fig. 14c). It also ratifies that the mechanism involved in the crystal growth process is a 2D nucleation mechanism, and the synthesized crystal has few dislocations.⁴⁸

3.5.2 Atomic force microscopy analysis (AFM). A quantitative description of the surface topography is needed to comprehend the surface morphology of the grown ionic material. So, to study the surface morphology, the grown DSFS crystal is subjected to AFM analysis. In general, surface roughness parameters can be evaluated on the basis of five parameters, namely height or amplitude parameter, spatial parameter, hybrid parameter, functional parameter and feature parameter. Among these, the height parameter is the critical parameter which gives the necessary details about the surface topography

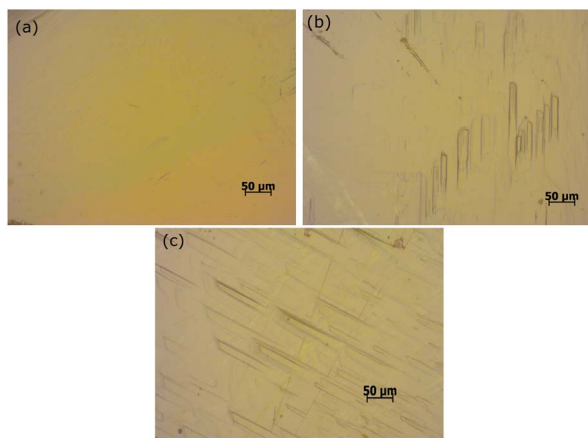


Fig. 14 Etch patterns of DSFS. (a) As grown surface; (b) 10 s and (c) 20 s.

of the material.⁴⁹ Fig. 15a and b show the 3D and 2D AFM images of the grown DSFS crystal. The root mean surface roughness value (S_q) and average roughness (S_a) value are found to be 32.975 nm and 26.824 nm. The difference between the average five maximum peaks and five minimum peaks is measured as the ten-point height (S_z), and the value of S_z is 140.297 nm. Based on the sign of the surface skewness (S_{sk}) moment, it can be concluded whether peaks or valleys are dominant on the surface. In the case of DSFS crystal, the S_{sk} value is negative (-0.0044), which clearly indicates that valleys are predominant and the surface is more planar.⁵⁰ The flatness of the crystal can be assessed through the surface kurtosis (S_{ku}) value. If the S_{ku} value is lower than 3, the surface is perfectly flat. If the value is higher than 3, then the surface possesses more peaks than valleys.⁵⁰ For the DSFS crystal, the S_{ku} value is 2.4297, which signifies the grown crystal surface is flat. It is obvious from the S_{sk} and S_{ku} values that the grown crystal is perfectly flat. All the measured surface roughness parameters are summarized in Table 3.

3.6 Nonlinear optical studies

3.6.1 Z scan measurement. The arrangement of chromophores in the crystal determines whether the material exhibits χ^2 or χ^3 effects. SCXRD studies affirmed that the title ionic crystal arranged in a centrosymmetric fashion with $P\bar{1}$ space group. In centrosymmetric crystals, due to the presence of an inversion centre, the resultant polarization has odd order susceptibility only, which means that even order susceptibility has null values. Due to the presence of odd order susceptibility the grown DSFS crystal exhibits third-order nonlinearity.⁵¹ Among the various technologies available to measure third-order nonlinearity, the Z scan technique is preferred over other methods due to its simplicity and reduced complexity.⁵² The basic principle behind the Z scan technique is that the sample under investigation is studied as a function of Z-position under the influence of a Gaussian beam to study the interaction between the sample and laser light. The four crucial parameters which help to determine the nonlinearity of the material can be accurately measured using the Z-scan technique: nonlinear refractive index (n_2), nonlinear absorption coefficient (β), third-order nonlinear optical susceptibility (χ^3) and second-order hyperpolarizability.⁵² In order to calculate the β and n_2 , two modes were adopted, namely open and closed

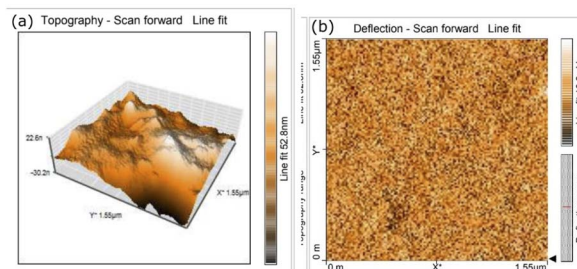


Fig. 15 AFM image of the grown DSFS crystal. (a) 3D image; (b) 2D image.



Table 3 Roughness parameters of DSFS crystal from AFM

Roughness parameters	Values
Root mean surface roughness (S_q)	32.975 nm
Average roughness (S_a)	26.824 nm
Ten-point height (S_z)	140.297 nm
Surface skewness (S_{sk})	-0.0044
Surface kurtosis (S_{ku})	2.4297

aperture modes. During the translation of the sample, under the Gaussian beam with respect to variation of transmittance, the intensity-dependent absorption of the material is studied *via* an open aperture and the nonlinear refractive index is measured with reference to the change in light intensity *via* the closed aperture method with a finite slit or aperture in front of the output detector.⁵³ With a helium–neon laser (632.8 nm) as the source, the Z-scan technique was carried out. With the aid of a Gaussian filter, the laser beam is converted to a Gaussian beam and the transformed beam is permitted to go through a convex lens with a focal distance of 30 mm to produce a focused beam with a beam waist diameter (W_0) of 12.05 μm . In order to check whether the subjected sample satisfies the thin sample approximation, the Rayleigh length value (Z_R) is calculated. For the Gaussian beam, the Rayleigh length ($Z_R = \pi W_0^2/\lambda$) is found to be 0.72 mm. The thickness (L) of the title crystal used for the analysis is 0.49 mm, which is less than the Rayleigh length; this clearly indicates that the sample used for measurement fulfilled the thin sample approximation ($Z_R \gg L$).⁵⁴ The crystal is translated through the Z-direction ($Z = +15$ to -15) with reference to the focus ($Z = 0$). The corresponding output reading is measured in the output power meter with regard to the position of the crystal.

Closed aperture. In order to evaluate the magnitude and sign of the nonlinear refractive index (n_2), a 2 mm aperture is placed in front of the photodetector. In the optical crystal, the illumination of high intense laser light tends to produce nonlinear thermal effects. The difference in the temperature of the crystal at the internally illuminated portion and outer region alters the material's refractive index. Thermal effects and the crystal's core structure are attributed to the material's self-focusing or defocusing nature.⁵⁵ The obtained closed aperture Z-scan graph is depicted in Fig. 16. The self-focusing behaviour ($+n_2$) of the grown ionic crystal is confirmed by the pre-focus valley followed by the post-focus peak in the closed aperture curve. From the graph, ΔT_{p-v} is calculated to figure out the nonlinear refractive index (n_2), with respect to the phase shift axis ($\Delta\Phi_0$). The difference in the normalized transmittance between the valley and peak can be derived from the following expression.

$$\Delta\Phi_0 = \frac{\Delta T_{(p-v)}}{0.406(1-S)^{0.25}} \quad (8)$$

'S' in eqn (8) denotes the linear transmittance through the aperture. By using the following expression, 'S' can be calculated:⁵⁶

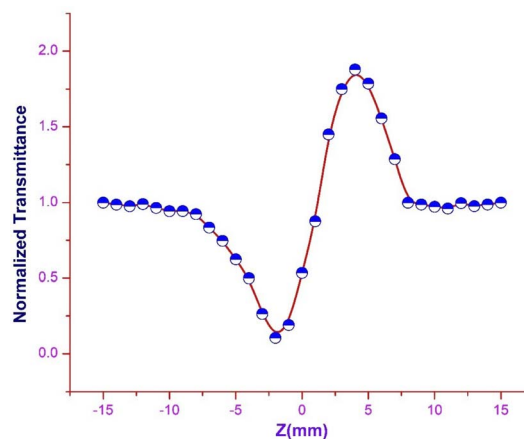


Fig. 16 Closed aperture trace of DSFS crystal.

$$S = 1 - \exp\left(\frac{-2r_a^2}{w_a^2}\right) \quad (9)$$

where r_a and w_a indicate the radius of the aperture and beam, respectively. The formula to calculate n_2 from the closed aperture graph is given as:⁵⁷

$$n_2 = \frac{\Delta\Phi_0}{Kl_0L_{\text{eff}}} \quad (10)$$

In the above eqn (10), K represents wavenumber ($K = 2\pi/\lambda$), L_{eff} indicates the effective thickness of the sample, and it is given by:

$$L_{\text{eff}} = [1 - \exp(\alpha L)]/\alpha \quad (11)$$

I_0 is the intensity of the incident beam at the focus. Where L and α represent the thickness and coefficient of absorption of the DSFS crystal.

Open aperture. The nonlinear coefficient of absorption of the title ionic crystal is evaluated through the open aperture method (Fig. 17). The valley formation at the focus ($Z = 0$) demonstrates reverse saturation absorption (RSA). This strong RSA effect is because of the higher absorption of molecules in the higher energy level state when compared with the ground state.⁵⁸ β is directly derived from the standard equation:⁵⁹

$$\beta = \frac{2\sqrt{2}\Delta T}{I_0L_{\text{eff}}} \quad (12)$$

ΔT represents the valley at the focus in the open aperture curve.

The real and imaginary parts of $\chi^{(3)}$ can be directly derived from the nonlinear refractive index and nonlinear coefficient of absorption with the following equations:⁵⁹

$$\text{Re}\chi^{(3)}(\text{esu}) = \frac{10^{-4}(\epsilon_0 c^2 n_0^2 n_2)}{\pi} \text{ cm}^2 \text{ W}^{-1} \quad (13)$$

$$\text{Im}\chi^{(3)}(\text{esu}) = \frac{10^{-2}(\epsilon_0 c^2 n_0^2 \lambda \beta)}{4\pi} \text{ cm W}^{-1} \quad (14)$$



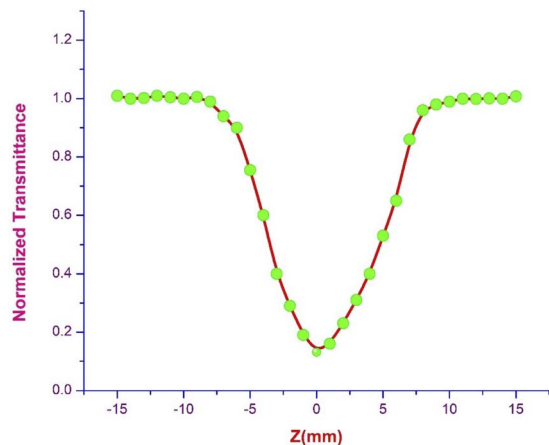


Fig. 17 Open aperture trace of DSFS crystal.

Table 4 Comparison of the $\chi^{(3)}$ value of the DSFS crystal with a few stilbazolium and organic crystals

Crystal name	$\chi^{(3)}$ value (esu)
DSFS (present work)	1.573×10^{-7}
VMST ³³	9.696×10^{-12}
DSMOS ⁶³	5.05×10^{-8}
KDP ⁶⁴	2.04×10^{-14}
4MBS ⁶⁵	11.04×10^{-7}

Coupling factor (ρ) is calculated for the obtained experimental results by adopting two methods: one directly by taking the ratio between the imaginary and real parts by using eqn (15):⁶⁰

$$\rho = \frac{\text{Im}(\chi^{(3)})}{\text{Re}(\chi^{(3)})} = \frac{\beta}{2kn_2} \quad (15)$$

The other method is by taking the ratio between ΔT and $\Delta\Phi_0$ of the open and closed aperture Z scan curves, respectively:

$$\rho = \frac{\Delta T}{\Delta\Phi_0} \quad (16)$$

The calculated values from eqn (15) and (16) are 0.23 and 0.17, respectively. Both the values lie well within the critical value of the coupling factor ($\rho_c = \frac{\sqrt{3}}{3} \approx 0.577$), which clearly indicates the validation of the obtained experimental data.⁶¹ The obtained value is less than 1/3 the value, which emphasises that the electronic polarization is responsible for the nonlinearity in the material.⁶² Both the n_2 and β are insensitive to the incident wavelength until it reaches a particular wavelength called the critical wavelength (λ_c). From the following formula, the λ_c of the DSFS crystal is calculated.

$$\lambda_c = \frac{4\pi n_2}{\sqrt{3}\beta} \quad (17)$$

The λ_c value of the DSFS crystal is 1530 nm. If the laser's operating wavelength is beyond 1530 nm, the transmittance

curve in the closed aperture method will disappear.⁶⁰ The calculated values of the nonlinear refractive index NLRI (n_2), nonlinear absorption coefficient NLAC (β) and third-order susceptibility ($\chi^{(3)}$) are found to be $4.6494 \times 10^{-11} \text{ m}^2 \text{ W}^{-1}$, $2.203 \times 10^{-4} \text{ m W}^{-1}$ and $1.573 \times 10^{-7} \text{ esu}$, respectively. The third-order susceptibility value of the title ionic crystal is compared with a few other stilbazolium derivatives and organic centrosymmetric crystals, and the comparison is tabulated in Table 4. These comparison results indicate that the title DSFS crystal possesses a comparatively large $\chi^{(3)}$ value.

The structure of the DSFS crystal is the main reason behind this high NLO response. As already discussed in the structural analysis, the presence of the BLA degree (0.11 Å) in the stilbazolium cation induces nonlinearity in the system. Along with that, the presence of a strong electron-donating group (dimethyl amino) and electron-withdrawing group (pyridinium) in the stilbazolium cation gives rise to the development of a conjugated structure where the delocalization of π electrons takes place through the styryl pyridinium (C=C) delocalization channel. Among various reported delocalized conjugated structures, styryl-pyridinium is considered to be one of the most efficient conjugated structures; it is due to the nitrogen and carbon atoms in the pyridine ring, which are in an sp^2 hybrid orbitals configuration and also a conjugated system is formed within the ring with one π orbital without a lone pair of electrons in the nitrogen atom.⁶⁶ This entire D- π -A setup tends to give larger electron density transfer within the cation structure, leading to the surge in the scale of the dipole moment, which in turn makes the molecules highly polarized and increases the non-linearity of the DSFS crystal.^{63,67} Apart from this, the interaction between the stilbazolium cation and the SO_3 ligand in the anion also induces non-linearity in the title crystal.³³ All the results attained from the Z scan analysis suggests that the ionic DSFS crystal can be used in optical limiting and switching applications.²⁷

4 Conclusion

The stilbazolium derivative single crystal DSFS with a new anion is fruitfully synthesized using the slow evaporation method. The structure of the grown ionic crystal is solved using SCXRD, and the CIF file is deposited in the CCDC database. The formation of the ethylenic bridge (C=C) is also confirmed *via* FTIR analysis. Results from the linear optical studies show that the bandgap and cut-off wavelength of the grown DSFS crystal are found to be 570 nm and 2.25 eV, respectively. The low value of the extinction coefficient (10^{-4}) and high value of the optical and electrical conductivity (10^{10}) indicates that the DSFS crystal is well suited for various optical applications. The S_{ku} value from AFM and etching analysis reveals the flatness and good optical quality of the title ionic crystal. Through the Z scan technique, the third-order susceptibility $\chi^{(3)}$ value of the DSFS crystal is determined to be $\chi^{(3)} = 1.573 \times 10^{-7} \text{ esu}$. The self-focusing effect along with the reverse saturation absorption nature of the title crystal suggests that the stilbazolium derivative crystal can be used for various applications such as remote sensing, photonic applications, optical limiting and optical switching. All the



characterization outcomes show that the DSFS crystal is suitable for various NLO applications.

Author contributions

S. A.: methodology, formal analysis, investigation, writing – original draft, and visualization. M. U. R.: conceptualization, visualization, supervision, review and editing. S. K.: conceptualization, review and editing. R. S. B.: conceptualization, review and editing.

Conflicts of interest

There are no conflicts to declare.

Acknowledgements

The authors would like to thank VIT – Vellore for providing the single crystal XRD measurement.

References

- M. Hasmuddin, M. Abdullah, P. Singh, M. Shkir, N. Vijayan and M. Wahab, *Opt. Laser Technol.*, 2015, **74**, 53–59.
- X. Wang, Y. Du, S. Ding, Q. Wang, G. Xiong, M. Xie, X. Shen and D. Pang, *J. Phys. Chem. B*, 2006, **110**, 1566–1570.
- R. Ragu, P. Mageshwari, M. Akilan and S. J. Das, *Appl. Phys. B: Lasers Opt.*, 2020, **126**, 1–10.
- T. H. Clara, R. Ragu, D. R. Jonathan and J. C. Prasana, *Opt. Mater.*, 2020, **109**, 110331.
- T. A. Hegde, A. Dutta, R. M. Jauhar, P. Karuppasamy, M. S. Pandian, M. Abith, T. S. Girisun and G. Vinitha, *Opt. Mater.*, 2020, **107**, 110033.
- S. Anand, S. Kalainathan, T. Amrutha, K. Senthil and M. Yamada, *J. Mater. Sci.: Mater. Electron.*, 2021, **32**, 11059–11071.
- C. Jin, F. Li, B. Cheng, H. Qiu, Z. Yang, S. Pan and M. Mutailipu, *Angew. Chem., Int. Ed.*, 2022, **61**, e202203984.
- M. Mutailipu, K. R. Poepelmeier and S. Pan, *Chem. Rev.*, 2021, **121**, 1130–1202.
- M. Faizan, M. Mehkoom, Z. Afroz, V. H. N. Rodrigues, S. Afzal and S. Ahmad, *J. Solid State Chem.*, 2021, **300**, 122255.
- M. Gomathi, C. Karnan, T. Sivanesan, J. C. Rhoda, S. Manivannan, V. Ragavendran, G. Vinitha and A. Prabakaran, *Chem. Phys. Lett.*, 2021, **776**, 138705.
- J. Yin, L. Li, Z. Yang, M. Jazbinsek, X. Tao, P. Günter and H. Yang, *Dyes Pigm.*, 2012, **94**, 120–126.
- K. Senthil, S. Kalainathan and A. R. Kumar, *CrystEngComm*, 2014, **16**, 9847–9856.
- S. Kannan, A. Sekar and K. Sivaperuman, *J. Mater. Chem. C*, 2020, **8**, 16668–16690.
- O.-P. Kwon, B. Ruiz, A. Choubey, L. Mutter, A. Schneider, M. Jazbinsek, V. Gramlich and P. Günter, *Chem. Mater.*, 2006, **18**, 4049–4054.
- E.-Y. Choi, M. Jazbinsek, J.-H. Jeong and O.-P. Kwon, *CrystEngComm*, 2012, **14**, 1045–1048.
- K. Senthil, S. Kalainathan and A. R. Kumar, *Spectrochim. Acta, Part A*, 2014, **124**, 603–610.
- Z. Yang, M. Jazbinsek, B. Ruiz, S. Aravazhi, V. Gramlich and P. Günter, *Chem. Mater.*, 2007, **19**, 3512–3518.
- A. A. Raj, S. J. Sundaram, R. Gunaseelan and P. Sagayaraj, *Spectrochim. Acta, Part A*, 2015, **149**, 957–964.
- A. Sekar, U. R. Muthurakku and K. Sivaperuman, *ChemistrySelect*, 2021, **6**, 10688–10716.
- S. Dharmalingam and G. N. Muthukrishnan, *J. Mol. Struct.*, 2022, **1250**, 131815.
- J. Pernak, J. Kalewska, H. Ksycińska and J. Cybulski, *Eur. J. Med. Chem.*, 2001, **36**, 899–907.
- K. Chanawanno, S. Chantrapromma, T. Anantapong, A. Kanjana-Opas and H.-K. Fun, *Eur. J. Med. Chem.*, 2010, **45**, 4199–4208.
- J. Polak and A. Jarosz-Wilkolazka, *Microb. Cell Fact.*, 2010, **9**, 51.
- M. K. Kumar, S. Sudhahar, G. Bhagavannarayana and R. M. Kumar, *Optik*, 2014, **125**, 5641–5646.
- J. A. M. Mani, *Investigation On Growth And Physicochemical Properties Of Nonlinear Optical Stilbazolium Family Crystals Dast Dapsh And Mbmsph*, PhD thesis, University of Madras, Chennai, 2017.
- G. Sheldrick, *Acta Crystallogr., Sect. A: Found. Adv.*, 2015, **71**, 3–8.
- N. Karuppanan and S. Kalainathan, *J. Phys. Chem. C*, 2018, **122**, 4572–4582.
- A. Teshome, M. D. H. Bhuiyan, G. J. Gainsford, M. Ashraf, I. Asselberghs, G. V. Williams, A. J. Kay and K. Clays, *Opt. Mater.*, 2011, **33**, 336–345.
- W. Kaminsky, *J. Appl. Crystallogr.*, 2005, **38**, 566–567.
- N. Vijayan, S. Rajasekaran, G. Bhagavannarayana, R. Ramesh Babu, R. Gopalakrishnan, M. Palanichamy and P. Ramasamy, *Cryst. Growth Des.*, 2006, **6**, 2441–2445.
- T. Vijayakumar, I. Hubert Joe, C. Reghunadhan Nair, M. Jazbinsek and V. Jayakumar, *J. Raman Spectrosc.*, 2009, **40**, 52–63.
- A. U. Rani, N. Sundaraganesan, M. Kurt, M. Cinar and M. Karabacak, *Spectrochim. Acta, Part A*, 2010, **75**, 1523–1529.
- M. K. Kumar, S. Sudhahar, P. Pandi, G. Bhagavannarayana and R. M. Kumar, *Opt. Mater.*, 2014, **36**, 988–995.
- M. Jazbinsek, *Appl. Sci.*, 2019, **9**, 882.
- S. Kalainathan, R. S. Babu, *et al.*, *J. Mater. Sci.: Mater. Electron.*, 2019, **30**, 10571–10578.
- K. Nivetha, W. Madhuri and S. Kalainathan, *J. Mater. Sci.: Mater. Electron.*, 2017, **28**, 8937–8949.
- J. Pernak, J. Kalewska, H. Ksycińska and J. Cybulski, *Eur. J. Med. Chem.*, 2001, **36**, 899–907.
- K. Pichan, S. P. Muthu and R. Perumalsamy, *J. Cryst. Growth*, 2017, **473**, 39–54.
- J. Polak and A. Jarosz-Wilkolazka, *Microb. Cell Fact.*, 2010, **9**, 51.
- R. Priya, S. Anitha, P. Mageshwari and R. Ragu, *J. Mater. Sci.: Mater. Electron.*, 2020, **31**, 21288–21302.
- A. Priyadarshini and S. Kalainathan, *J. Mater. Sci.: Mater. Electron.*, 2017, **28**, 7401–7412.



- 42 A. Priyadharshini and S. Kalainathan, *Opt. Mater.*, 2018, **78**, 35–43.
- 43 A. Priyadharshini and S. Kalainathan, *J. Phys. Chem. Solids*, 2018, **123**, 59–69.
- 44 R. Ragu, P. Mageshwari, M. Akilan and S. J. Das, *Appl. Phys. B: Lasers Opt.*, 2020, **126**, 1–10.
- 45 M. Rajalakshmi, R. Indirajith, P. Ramasamy and R. Gopalakrishnan, *Mol. Cryst. Liq. Cryst.*, 2011, **548**, 126–141.
- 46 T. R. Kumar, R. Jeyasekaran, S. R. Kumar, M. Vimalan and P. Sagayaraj, *Appl. Surf. Sci.*, 2010, **257**, 1684–1691.
- 47 A. U. Rani, N. Sundaraganesan, M. Kurt, M. Cinar and M. Karabacak, *Spectrochim. Acta, Part A*, 2010, **75**, 1523–1529.
- 48 V. Sangeetha, K. Gayathri, P. Krishnan, N. Sivakumar, N. Kanagathara and G. Anbalagan, *J. Cryst. Growth*, 2014, **389**, 30–38.
- 49 N. Saravanan, G. Palani, S. Shanmugan and V. Chithambaram, *Mater. Today: Proc.*, 2020, **30**, 115–122.
- 50 A. Sekar, U. R. Muthurakku and K. Sivaperuman, *ChemistrySelect*, 2021, **6**, 10688–10716.
- 51 K. Senthil, S. Kalainathan and A. R. Kumar, *CrystEngComm*, 2014, **16**, 9847–9856.
- 52 K. Senthil, S. Kalainathan, A. R. Kumar and P. Aravindan, *RSC Adv.*, 2014, **4**, 56112–56127.
- 53 S. Shalini, S. Shahil Kirupavathy, E. Jerusha and G. Vinitha, *J. Mater. Sci.: Mater. Electron.*, 2019, **30**, 6528–6536.
- 54 G. Sheldrick, *Acta Crystallogr., Sect. A: Found. Adv.*, 2015, **71**, 3–8.
- 55 K. Senthil, S. Kalainathan and A. R. Kumar, *Spectrochim. Acta, Part A*, 2014, **124**, 603–610.
- 56 S. Shettigar, G. Umesh, K. Chandrasekharan, B. Sarojini and B. Narayana, *Opt. Mater.*, 2008, **30**, 1297–1303.
- 57 A. Sundari and S. Manikandan, *J. Mater. Sci.: Mater. Electron.*, 2018, **29**, 558–567.
- 58 R. Surekha, P. Sagayaraj and K. Ambujam, *Opt. Mater.*, 2014, **36**, 945–949.
- 59 M. Suriya, M. Manimaran, B. M. Boaz and K. S. Murugesan, *J. Mater. Sci.: Mater. Electron.*, 2021, **32**, 11393–11417.
- 60 A. Teshome, M. D. H. Bhuiyan, G. J. Gainsford, M. Ashraf, I. Asselberghs, G. V. Williams, A. J. Kay and K. Clays, *Opt. Mater.*, 2011, **33**, 336–345.
- 61 B. Thirumalaiselvam, R. Kanagadurai, D. Jayaraman and V. Natarajan, *Opt. Mater.*, 2014, **37**, 74–79.
- 62 T. Vijayakumar, I. Hubert Joe, C. Reghunadhan Nair, M. Jazbinsek and V. Jayakumar, *J. Raman Spectrosc.*, 2009, **40**, 52–63.
- 63 N. Vijayan, S. Rajasekaran, G. Bhagavannarayana, R. Ramesh Babu, R. Gopalakrishnan, M. Palanichamy and P. Ramasamy, *Cryst. Growth Des.*, 2006, **6**, 2441–2445.
- 64 X. Wang, Y. Du, S. Ding, Q. Wang, G. Xiong, M. Xie, X. Shen and D. Pang, *J. Phys. Chem. B*, 2006, **110**, 1566–1570.
- 65 Z. Yang, M. Jazbinsek, B. Ruiz, S. Aravazhi, V. Gramlich and P. Günter, *Chem. Mater.*, 2007, **19**, 3512–3518.
- 66 J. Yin, L. Li, Z. Yang, M. Jazbinsek, X. Tao, P. Günter and H. Yang, *Dyes Pigm.*, 2012, **94**, 120–126.
- 67 I. M. Zahid, S. Kalaiyarasi, M. K. Kumar, T. Ganesh, V. Jaisankar and R. M. Kumar, *Mater. Sci.*, 2016, **34**, 811–818.

