1. Introduction

Rare-earth doped calcium aluminate phosphors have been intensively studied due to their excellent afterglow properties, among which Dy$^{3+}$ doped calcium aluminate (CaAl$_2$O$_4$:Dy$^{3+}$) is a well-known inorganic phosphor which emits white photoluminescence (PL) under ultraviolet excitation. Apart from its white PL, CaAl$_2$O$_4$:Dy$^{3+}$ exhibits a white afterglow even after the removal of ultraviolet excitation. As reported by Liu et al. in 2005, Dy$^{3+}$ works as the luminescence center of afterglow for CaAl$_2$O$_4$:Dy$^{3+}$ and the white afterglow of solid state reaction derived CaAl$_2$O$_4$:Dy$^{3+}$ lasts 32 min at the optimal doping concentration of 2 at%.[1] For afterglow materials, the processes of charged carrier excitation, migration, trapping, detrapping, and radiative recombination are critically important to understand their afterglow properties.[7,8] For example, long afterglow can be achieved at room temperature only when traps have an appropriate activation energy somewhere around 0.65 eV whereas shallow traps ($E \leq 0.4$ eV) and deep traps ($E > 2$ eV) are not favorable because they can be emptied either easily or with difficulty at room temperature.[7] Up to date, only a single report exists on the afterglow of CaAl$_2$O$_4$:Dy$^{3+}$, leaving the processes of charged carrier excitation, migration, trapping and detrapping not fully revealed. The lack of such knowledge hampers the comprehensive understanding of the afterglow mechanism of CaAl$_2$O$_4$:Dy$^{3+}$.

With respect to solid state reaction derived CaAl$_2$O$_4$:Dy$^{3+}$, sol–gel combustion derived CaAl$_2$O$_4$:Dy$^{3+}$ is abundant in oxygen and calcium vacancies. In CaAl$_2$O$_4$:Dy$^{3+}$, oxygen vacancies are potential electron traps because they are positively charged whereas calcium vacancies are potential hole traps because they are negatively charged.[8,9] Therefore, sol–gel combustion derived CaAl$_2$O$_4$:Dy$^{3+}$ phosphors are suitable for exploring the afterglow mechanisms of CaAl$_2$O$_4$:Dy$^{3+}$. In this paper, we report the PL and afterglow properties of sol–gel combustion derived CaAl$_2$O$_4$:Dy$^{3+}$ with the doping concentrations varying from 0.1 to 5.0 mol%. A picture on the PL and afterglow mechanisms of CaAl$_2$O$_4$:Dy$^{3+}$ is given to reveal the processes of charged carriers’ excitation, migration, trapping, detrapping, and recombination in CaAl$_2$O$_4$:Dy$^{3+}$.

2. Experimental section

2.1 Sol–gel combustion synthesis of CaAl$_2$O$_4$:Dy$^{3+}$

CaAl$_2$O$_4$:Dy$^{3+}$ phosphors with the doping concentrations varying from 0.1 to 5.0 mol% were synthesized via the sol–gel combustion with urea as fuel.[10–13] All reagents were in analytical grade and provided by Sinopharm Chemical Reagents Co., Ltd (Shanghai, China). The purity of dysprosium oxide (Dy$_2$O$_3$) was 99.99%. Under stirring with a magnetic bar, Ca(NO$_3$)$_2$·4H$_2$O (0.2 mol), Al(NO$_3$)$_3$·9H$_2$O (0.4 mol), urea (6.0 mol), H$_3$BO$_3$ (0.02 mol) and stoichiometric amount of Dy$_2$O$_3$ were dissolved in deionized water (600 mL) to form a transparent solution. A homogeneous solution was obtained after the mixture was
stirred vigorously for 60 min in a glass beaker. After having been aged at room temperature for two weeks, the solution was ready for the sol–gel combustion. Urea and boric acid functioned as the fuel and flux, respectively. Alumina crucibles, each with the volume capacity of 50 mL, were employed as the reaction containers. Half filled with the aged solution, the solution-containing alumina crucible was transferred into an air-filled box furnace for self-propagating combustion. The temperature in the furnace was preset at 780 °C. After the solution-containing crucible was transferred into the furnace, the temperature in the furnace was gradually dropped to about 706 °C. Heated at such high temperatures, water in the solution got boiling and the starting materials (fuels and metal nitrates) in the crucible were partially decomposed until the organic fuels were automatically ignited to initiate the exothermic reactions. The sol–gel combustion yielded voluminous gases and bright flames. In this work, the synthesis was initiated by point-heating of a small part of the mixture in the crucible, and the ignition was started at around 750 °C. Once started, a wave of exothermic reactions swept through the remaining material in the crucible. This combustion synthesis lasted for about 40 s, which was exceptionally fast when compared to high-temperature solid state reactions. During the sol–gel combustion, a large amount of energy was released from the exothermic reactions, which in turn raised the temperature in the furnace up to 830 °C. Measured with an infrared thermometer, the temperature in the flame was up to 1300 °C. After the fire was extinguished, the crucible was taken out of the furnace immediately. The total holding time of the crucible in the furnace was about 4 min. White powders were resulted after the sol–gel combustion. According to molar ratio of Dy³⁺ to Ca²⁺ in the starting materials, a series of CaAl₂O₄: Dy³⁺ phosphors were obtained with the nominal doping concentration of Dy³⁺ varying in the range of 0.1–5.0 mol%. No further calcination was applied to the sol–gel combustion derived CaAl₂O₄: Dy³⁺.

2.2 Solid state synthesis of CaAl₂O₄: Dy³⁺ phosphors
For comparison, solid state reaction derived CaAl₂O₄: Dy³⁺ phosphors were synthesized with CaCO₃ (0.02 mol), Al₂O₃ (0.02 mol), H₃BO₃ (0.002 mol), Dy₂O₃ as raw materials. Stoichiometric mixtures of the raw materials were homogeneously mixed and ground. Subsequently the mixture was transferred into alumina crucibles and sintered at 900 °C for 4 h in an air-filled electric crucible furnace. The sintered products were ground again in an agate mortar, then the powder products were calcined at 1350 °C for 10 h in air. According to the molar ratio of Dy³⁺ to Ca²⁺ in the starting materials, the doping concentration of Dy³⁺ in the solid state reaction derived CaAl₂O₄: Dy³⁺ varied in the range 0.1–10.0 mol%.

2.3 Phase, morphology and elemental composition of CaAl₂O₄: Dy³⁺ phosphors
X-Ray diffraction (XRD) profiles of CaAl₂O₄: Dy³⁺ phosphors were recorded on X-ray diffractometer (D/max 2500 PC, Rigaku Corporation, Akishima, Japan) using Cu Kα radiation (λ = 0.15405 nm). The scanning electron microscope (SEM) (model S-4800, Hitachi, Tokyo, Japan) was employed to analyze the morphology of the synthesized products. The SEM was coupled with a silicon drifted detector as the X-ray analyzer to record the energy dispersive X-ray (EDX) spectrum of the synthesized products. The micrographs of CaAl₂O₄: Dy³⁺ nanocrystals were recorded on a transmission electron microscope (TEM) (model JEM-2100, Japan Electronics Corp, Akishima, Japan). Samples for TEM analysis were prepared by suspending the particles in ethanol under the excitation of ultrasonification and then drying a drop of the suspension on a carbon-coated copper grid.

2.4 PL excitation and emission spectra of CaAl₂O₄: Dy³⁺
The PL excitation spectrum of CaAl₂O₄: Dy³⁺ was measured with the fluorescence spectrometer F-7000 (Hitachi, Japan). The spectrophotometer (Tianjin Gangdong Ltd., Tianjin, China) was used to acquire the steady-state PL spectra of CaAl₂O₄: Dy³⁺. The excitation source of the PL spectrum was provided by a helium-cadmium laser (Kimmon Electric Co. Ltd., Tokyo, Japan). The emission wavelength and the output power of the laser radiation were 325 nm and 13 mW, respectively.

2.5 Afterglow spectra and thermoluminescence (TL) glow curves of CaAl₂O₄: Dy³⁺
The afterglow spectra of CaAl₂O₄: Dy³⁺ were recorded with the PL spectrophotometer (Tianjin Gangdong Ltd, Tianjin, China) immediately after the ultraviolet irradiation of a high-pressure mercury lamp was blocked off. The output power of the high-pressure mercury lamp was 175 W. The irradiation duration of the high-pressure lamp irradiation was 3 min. Afterglow decay curve was taken by focusing the afterglows into the entrance slit of the spectrometer. The TL glow curves of CaAl₂O₄: Dy³⁺ were measured on a TL meter constructed according to the scheme given by Yamashita et al.⁴⁴ The phosphor was placed on an electrically heated plate, the temperature of the plate was controlled with a program in computer. As the temperature of the plate was raised linearly with time, the light output of the phosphor was recorded using a photomultiplier in a photon-counting mode. A bialkali photomultiplier tube (model H10425, Hamamatsu, Japan) was used in the TL meter for the luminescence detection, it covers the 350–600 nm wavelength range. Prior to the TL measurements, CaAl₂O₄: Dy³⁺ phosphors were exposed to the 254 nm irradiation of a low-pressure mercury lamp for 5 min. The output power of the low-pressure mercury lamp was 32 W. The TL signals of CaAl₂O₄: Dy³⁺ phosphors were recorded when the CaAl₂O₄: Dy³⁺ phosphors were heated from 10 to 200 °C at a rate of 2 °C s⁻¹.

3. Results and discussions
3.1 Phase and morphology of CaAl₂O₄: Dy³⁺
Fig. 1 represents the XRD profiles of CaAl₂O₄: Dy³⁺ with different doping concentrations. The XRD data of monoclinic CaAl₂O₄ registered in Joint Committee on Powder Diffraction Standards (JCPDS) card no. 23-1306 are depicted at the bottom of the figure for comparison. The scale of the intensity of powder X-ray
Figure 1 shows the XRD curves of CaAl₂O₄:Dy³⁺ with doping concentrations varying from 0.1 to 5.0 mol%. The standard diffraction data of CaAl₂O₄ (JCPDS 23-1036) are displayed at the bottom of this figure for comparison.

Diffraction is counts of photons received by the detector. For the purpose of clarity, the XRD curves in Fig. 1 are evenly spaced after having been shifted upwards one by one, so the scale of the intensity of powder XRD in Fig. 1 is arbitrary units. As can be seen in Fig. 1, each XRD profile of CaAl₂O₄:Dy³⁺ exhibits distinct peaks at 2θ = 16.01, 18.99, 21.98, 23.97, 30.08, 31.35, 35.63, 37.46, 41.11, 44.76, and 47.24°. According to the diffraction data registered in JCPDS card no. 23-1036, these peaks can be attributed to the X-ray diffractions from crystallographic planes (111), (112), (020), (211), (123), (015), (006), (313), (232), (040) and (226) of monoclinic CaAl₂O₄, respectively. Among these peaks, the strongest one is located at 30.98° and indexed as (123). Actually, this peak is contributed jointly by two crystallographic planes (220) and (123) of monoclinic CaAl₂O₄, which are located at 2θ = 30.11 and 30.05°, respectively. Being closely packed and nearly equal in diffraction intensity, the two peaks are not distinguishable in the XRD profiles. This is the reason why some researchers indexed the peak at around 30.8° as (220)."
charged carriers. Till date, several rare earth doped CaAl2O4 materials are studied and available in literature. For example, Liu et al. prepared CaAl2O4:Dy3+ via solid state reaction at 1350 °C. Due to the features of solid state reaction, the solid state reaction derived CaAl2O4:Dy3+ phosphors are different from the sol–gel combustion derived ones in the aspects of population density of intrinsic defects. Especially in the sol–gel combustion method, addition of H3BO3 lowers reaction temperature, accelerates a diffusion of raw materials, and eventually influences the afterglow properties. Detailed discussions are given by Takeuchi and Kishine in the case of Eu2+ and Dy3+ codoped SrAl2O4.

3.2 EDX spectrum of CaAl2O4:Dy3+
Fig. 3 depicts the EDX spectrum of sol–gel combustion derived CaAl2O4:Dy3+ (5.0 mol%). As can be seen in Fig. 3, the characteristic X-ray emissions of O(Ka), Al(Ka) and Ca(Ka2), and Ca(Kb1,3) are located at 0.525, 1.486, 3.693, and 4.013 keV, respectively. Moreover, the characteristic emissions of Dy(La1) and Dy(Lb) can be identified in the EDX spectrum at 6.495 and 7.248 keV, respectively. The X-ray emissions of Au(Ma1) at 2.122 keV and Au(La1) at 9.713 keV are due to Au sputtering for the convenience of SEM characterization. Apparently, the EDX spectrum of CaAl2O4:Dy3+ confirms the presence of elements Ca, Al, O and Dy in the sample.

3.3 XPS spectrum of CaAl2O4:Dy3+
The oxidation states of chemical elements in CaAl2O4:Dy3+ are investigated with the XPS. Fig. 4 represents the XPS survey scan (a) and high-resolution XPS spectra of Ca 2p (b), Al 2p (c), O 1s (d), Dy 3d (e), and Dy 4d (f) in sol–gel combustion derived CaAl2O4:Dy3+. The nominal doping concentration of Dy3+ is 2.0 mol%. As can be seen in Fig. 4(b–d), the binding energies of Ca 2p3/2, Ca 2p1/2, Al 2p1/2, and O 1s are located at approximately 347.3, 350.9, 74.1, and 531.8 eV, respectively. Fig. 4(e) shows that the binding energies of Dy 3d5/2 and Dy 3d3/2 are located at 1297.6 and 1335.1 eV, respectively. Fig. 4(f) shows that the binding energies of Dy 4d5/2 and Dy 4d3/2 are located at around 154.0 and 157.3 eV, respectively. The data in Fig. 4 point out the presence of Ca, Al, O and Dy3+ in the phosphor.

3.4 PL excitation spectrum of CaAl2O4:Dy3+
Fig. 5 shows the PL excitation spectrum of sol–gel combustion derived CaAl2O4:Dy3+ (2.0 mol%). The emission wavelength is
fixed at 574 nm for the PL excitation measurement. As can be seen in Fig. 5, there are 6 peaks in the PL excitation spectrum of CaAl$_2$O$_4$:Dy$^{3+}$ (2.0 mol%), which are located at 296, 326, 389, 425, and 452 nm, respectively. The excitation peaks can be assigned to the 4f–4f transitions of Dy$^{3+}$. The strongest absorption of CaAl$_2$O$_4$:Dy$^{3+}$ is located at 350 nm, which is due to the $^6$H$_{15/2} \rightarrow ^4$D$_{7/2}$, $^6$H$_{15/2} \rightarrow ^4$F$_{5/2}$, $^6$H$_{15/2} \rightarrow ^4$F$_{7/2}$, $^6$H$_{15/2} \rightarrow ^4$I$_{13/2}$, $^6$H$_{15/2} \rightarrow ^4$G$_{11/2}$ and $^6$H$_{15/2} \rightarrow ^4$I$_{15/2}$ transitions of Dy$^{3+}$, respectively.26–31 The PL excitation spectrum indicates that a light source with the emission wavelength of 350 nm excites CaAl$_2$O$_4$:Dy$^{3+}$ most efficiently. As a contrast, a light source with the emission wavelength of 326 nm can excite CaAl$_2$O$_4$:Dy$^{3+}$ but in a much less efficient manner.

3.5 PL spectra of CaAl$_2$O$_4$:Dy$^{3+}$ at different doping concentrations

Fig. 6 represents the PL spectra of sol–gel combustion derived CaAl$_2$O$_4$:Dy$^{3+}$ with doping concentrations in the range of 0.1–5.0 mol%. When the doping concentration is as low as 0.1 mol%, the PL spectrum of CaAl$_2$O$_4$:Dy$^{3+}$ consists of a broad PL band peaking at about 400 nm whereas the characteristic emissions of Dy$^{3+}$ are hardly discernible, as shown by the PL spectrum (a). The broadband emissions can be attributed to the intrinsic defect emissions in CaAl$_2$O$_4$, which are namely the oxygen and calcium vacancies.3,4 The PL properties of undoped CaAl$_2$O$_4$ and their origins were discussed in details in our previous work.15,16 Actually, similar broadband emissions are recorded for undoped CaAl$_2$O$_4$,15,16 Tb$^{3+}$ doped CaAl$_2$O$_4$,3,4

Fig. 4 XPS survey scan (a) and high-resolution XPS spectra of Ca 2p (b), Al 2p (c), O 1s (d), Dy 3d (e), and Dy 4d (f) in sol–gel combustion derived CaAl$_2$O$_4$:Dy$^{3+}$ (2.0 mol%).
undoped SrAl2O4,8 and Dy3+ doped BaAl2O4.28 The characteristic emissions of Dy3+ activator are represented by two narrowband emissions peaking at around 482 and 574 nm, respectively, which are due to the 4F9/2 → 6H15/2 and the 4F9/2 → 6H13/2 transitions of Dy3+ activator.2-4 As documented in the literature, the transition 4F9/2 → 6H15/2 (∆L = 2 and ∆J = 3) is magnetic dipole allowed whereas the transition 4F9/2 → 6H13/2 is identified as a hypersensitive electric dipole transition (∆L = 2 and ∆J = 2) of Dy3+. The reason why the characteristic emissions of Dy3+ activator are hardly discernible in the PL spectrum of CaAl2O4:Dy3+ [0.1 mol%] rests on the fact that the characteristic emissions of Dy3+ are too weak when compared to the strong emissions from the intrinsic defects in the host. As the doping concentration is increased to 0.2 and 0.4 mol%, the narrowband emissions of Dy3+ at 574 nm become discernible in the PL spectrum of CaAl2O4:Dy3+ [0.2 and 0.4 mol%], as shown by the PL spectra (b) and (c). As the doping concentration is elevated further to 0.6 mol%, the narrowband emissions of Dy3+ can be identified clearly in the PL spectrum of CaAl2O4:Dy3+, as shown by the PL spectrum (d). It is apparent in Fig. 6 that the intensity of the characteristic emissions of Dy3+ increases monotonically with the doping concentration. Consequently, the PL spectrum of CaAl2O4:Dy3+ indicates that the intrinsic defects in CaAl2O4 and the doping species Dy3+ act as two independent sets of luminescence center of PL in CaAl2O4:Dy3+.

The insight reason for the change in the PL intensity with doping of Dy3+ is due to the changes in the population density of the luminescence center of PL. The substitution of Dy3+ for Ca2+ in the lattice of CaAl2O4 promotes the production of two kinds of luminescence center of PL, as shown in eqn (1). The first kind of luminescence center of PL is the Dy3+ in Ca2+ site while the second kind of luminescence center of PL is the oxygen vacancy.

As the doping concentration of Dy3+ increases, the population densities of the two kinds of luminescence center of PL are increased, which in turn lead to the enhancement in the PL intensities. It is noted that the PL intensity of each spectrum in Fig. 6 is normalized at 400 nm.

$$\text{Dy}^{3+} + \text{Ca}^{2+} + \text{O}^{2-}_0 \rightarrow \text{Dy}^{3+} + \text{V}^{\omega}_0 + \text{Ca}^{2+} + \text{O}^{2-} \quad (1)$$

To check the influences of excitation wavelength on the emissions of CaAl2O4:Dy3+, we measured the PL spectra of CaAl2O4:Dy3+ (5 mol%) under different excitation wavelengths. The emission spectra of CaAl2O4:Dy3+ under the excitations of 326 nm and 350 nm are shown in Fig. S3.† As can be seen in Fig. S3, the structures of the two PL spectra are nearly identical to each other regardless the variation in the excitation wavelength. However, the emission intensity of CaAl2O4:Dy3+ is apparently sensitive to the excitation wavelength. For example, the characteristic emissions of Dy3+ are stronger under the excitation of 350 nm than those under the excitation of 326 nm. One of the reasons rests on the fact that Dy3+ exhibits stronger absorption at 350 nm than at 326 nm, as endorsed by Fig. 5.

Due to the combined contributions of defect emissions from the host and the characteristic emissions of the dopant, the color coordinates of the emissions of CaAl2O4:Dy3+ varies with the doping concentration. Color coordinates of luminescent materials, which can be calculated from their PL spectral data, are important parameters to quantitatively describe the emission color for luminescent materials.25-27 The CIE chromaticity coordinates of CaAl2O4:Dy3+ are given in Table S1† for different doping concentrations. As shown in Table S1,† the chromaticity coordinates of CaAl2O4:Dy3+ change with the doping concentration. For example, the chromaticity coordinates of CaAl2O4:Dy3+ (0.1 mol%) are (0.156, 0.094), those of CaAl2O4:Dy3+ (5.0 mol%) are (0.187, 0.176). The emission photos and the CIE chromaticity diagram of CaAl2O4:Dy3+ are shown in Fig. S4.† It is clear that the PL color of CaAl2O4:Dy3+ keeps blue when the doping concentration varies in the range of 0.1–5.0 mol%. As referred to the PL spectrum in Fig. 6, the defect-related emissions (blue) of the host are a major constituent in the PL spectrum whereas the characteristic emissions of Dy3+ are a minor constituent. That is the reason why the solid-gel combustion derived CaAl2O4:Dy3+ exhibits blue colored emission in spite of the variation of the doping concentration in the range of 0.1–5.0 mol%.

The assignment of the broadband emissions in Fig. 6 to intrinsic defects in CaAl2O4 gains further support from the PL spectrum of solid state reaction derived CaAl2O4:Dy3+. The solid state reaction route involves chemical decomposition and reactions at much high temperatures (often from 1000 to 1500 °C) to produce a new solid composition. At such high temperatures, oxygen atoms migrate into the crystal lattice to repair the defects with the result of decreased population density of oxygen vacancies in CaAl2O4. Therefore, the broadband emissions should be weakened if the density of oxygen vacancies in CaAl2O4 is reduced at high temperature. Fig. 7 depicts the normalized PL spectra of solid state reaction derived
CaAl$_2$O$_4$:Dy$^{3+}$ at 1350 °C for 10 h. Obviously, each PL spectrum of the solid state reaction derived CaAl$_2$O$_4$:Dy$^{3+}$ consists of the characteristic emissions of Dy$^{3+}$. The most striking feature in Fig. 7 is that the broadband emissions extending from 350 to 450 nm are weakened for the solid state reaction derived CaAl$_2$O$_4$:Dy$^{3+}$. The chromaticity coordinates of the solid state reaction derived CaAl$_2$O$_4$:Dy$^{3+}$ (10 mol%) are (0.270, 0.365), as shown in Fig. S5.

### 3.6 Afterglow spectrum of CaAl$_2$O$_4$:Dy$^{3+}$

The emission spectrum of a high-pressure mercury lamp contains a pronounced spectral line at 365.4 nm, which is close to an absorption peak of CaAl$_2$O$_4$:Dy$^{3+}$. After exposure to the illumination of a high-pressure mercury lamp (175 W) for 3 min, CaAl$_2$O$_4$:Dy$^{3+}$ exhibits intense afterglow after the ultraviolet excitation is blocked off. Fig. S6† depicts the afterglow photos of CaAl$_2$O$_4$:Dy$^{3+}$ (0.8 mol%) taken at different times after the irradiation of the high-pressure mercury lamp is blocked off. As shown in Fig. S6,† the white afterglow of CaAl$_2$O$_4$:Dy$^{3+}$ (0.8 mol%) can last more than 60 min. The initial luminance of the afterglow is found to depend on the doping concentration of Dy$^{3+}$. Fig. 8(a) shows the plot of the integrated afterglow intensity of sol–gel combustion derived CaAl$_2$O$_4$:Dy$^{3+}$ versus the dopant concentration of Dy$^{3+}$. As can be seen in Fig. 8(a), the optimal doping concentration is around 0.8 mol%. It is noted that the afterglow gets quenched when the doping concentration is high (i.e., 5 mol%). Just like the case of PL quenching at high doping concentration, the non-radiative interaction between dopants is one of the reason of the concentration induced afterglow quenching. However, the traps generated by the dopant in the lattice of CaAl$_2$O$_4$ should be the key factor to be responsible for the afterglow quenching. As described in eqn (1), doping CaAl$_2$O$_4$ with Dy$^{3+}$ results in oxygen vacancies. These positively charged oxygen vacancies can act as electron traps. The increase in the population density as well as the change in the trap depth of these the positively charged oxygen vacancies generate significant effects on the afterglow duration of the phosphor. For example, the afterglow duration is very short when the trap depth is shallow ($E < 0.6$ eV), and no afterglow can be observed when the trap depth is too deep ($E > 2.0$ eV). To observe afterglow at room temperature, the traps should have an appropriate activation energy somewhere between these two extremes, a trap depth around 0.65 eV is considered to be optimal. 7 Fig. 8(b) depicts the afterglow spectrum of the sol–gel combustion derived CaAl$_2$O$_4$:Dy$^{3+}$ at the doping concentration of 0.8 mol%. As can be seen in Fig. 8(b), the afterglow spectrum consists of two narrow emission bands of Dy$^{3+}$, which is distinctly different from the broadband afterglow of CaAl$_2$O$_4$:Eu$^{2+}$.16,34,52,53 Thus Fig. 8(b) demonstrates that Dy$^{3+}$ acts as the luminescence center of afterglow for CaAl$_2$O$_4$:Dy$^{3+}$. The CIE chromaticity coordinates of the afterglow of CaAl$_2$O$_4$:Dy$^{3+}$ (0.8 mol%) are calculated to be (0.265, 0.305), as shown in Fig. S7.† It is clear that the afterglow color of CaAl$_2$O$_4$:Dy$^{3+}$ is close to white at the doping level 0.8 mol%. For
3.7 Afterglow decay profile of CaAl₂O₄:Dy³⁺

The duration of afterglow is a straightforward and standardized parameter to evaluate the properties of an afterglow material, whereby 0.32 mcd m⁻² is often used as threshold for defining the duration of an afterglow. Fig. 9 depicts the afterglow decay profile of sol–gel combustion derived CaAl₂O₄:Dy³⁺ (0.8 mol%). The phosphor was exposed to illumination of a high-pressure mercury lamp irradiation (175 W) for 3 min before the measurement of afterglow decay curve. As shown by the raw data in Fig. 9, the afterglow duration of CaAl₂O₄:Dy³⁺ is determined to be about 115 min. It is found that the decay curves in Fig. 9 could be best fitted to the tri-exponential function according to eqn (2)

\[ I(t) = I_1 \exp(-t/\tau_1) + I_2 \exp(-t/\tau_2) + I_3 \exp(-t/\tau_3), \]

where \( I(t) \) is the afterglow intensity at time \( t \) after blocking the laser excitation, \( I_i \) is the prefactor of the exponential component whose lifetime decay constant is \( \tau_i \) (\( i = 1–3 \)). The red solid line in Fig. 9 represents the fit of the experimental signals by eqn (2). The fitting parameters are tabulated in the figure. Clearly, the CaAl₂O₄:Dy³⁺ has three largely different decay components with constants of \( \tau_1 = 2.51 \text{ min}, \tau_2 = 11.31 \text{ min and } \tau_3 = 89.29 \text{ min.} \) We can see that \( \tau_2 \) and \( \tau_3 \) are much longer than \( \tau_1 \), suggesting the presence of deeper traps in the phosphor. The luminance reading of the first data point in Fig. 9 is 8.3 mcd m⁻² at 3.67 min. In order to derive the luminance at \( t = 0 \), it is necessary to extrapolate the data, yielding the value of the luminance to be 22.35 mcd m⁻² at \( t = 0 \). It is worth of noting that the luminance of the phosphor at the initial moment is different from the luminance of the first component at the initial moment (i.e., \( I_1 = 16.56 \text{ mcd m}^{-2}. \)). Actually, the luminance of the phosphor at the initial moment is the summation of the luminance of the three components at the initial moment, that is, the sum of \( I_1, I_2 \) and \( I_3 \). Moreover, the afterglow duration of the phosphor is different from the longest lifetime decay constant \( \tau_3 \) (i.e., 89.29 min). According to eqn (1), the afterglow duration of the phosphor is defined by the variable \( t \) when \( I(t) \) reaches the threshold 0.32 mcd m⁻². The luminance decreases to 0.3205 mcd m⁻² when \( t = 111 \text{ min.} \) Consequently, the afterglow duration derived from eqn (2) (about 111 min) is very close the afterglow duration derived from raw data (about 115 min).

Being distinctly different from PL decay, the afterglow phenomenon is a particular case of thermostimulated luminescence and is a defect dependent phenomenon. Additional notes can be found in the ESI.† The afterglow duration of CaAl₂O₄:Dy³⁺ heavily depends on the wavelength of ultraviolet excitation. The afterglow of CaAl₂O₄:Dy³⁺ (0.8 mol%) lasts for about 60 s when the 325 nm excitation of the HeCd laser (13 mW) is switched off. Conversely, the afterglow of CaAl₂O₄:Dy³⁺ (0.8 mol%) lasts for only 30 s when the 254 nm irradiation of a low-pressure mercury lamp (32 W) is turned off. Moreover, the afterglow duration of CaAl₂O₄:Dy³⁺ depends on the doping concentration when the wavelength of ultraviolet excitation is fixed. For example, after exposed to the 254 nm irradiation of a low-pressure mercury lamp (32 W) for 5 min, the afterglows of CaAl₂O₄:Dy³⁺ last for about 5, 8, 14, 19, 30, 60, 120, 75, 10 and 0 s when the doping concentrations are 0.1, 0.2, 0.4, 0.6, 0.8, 1.0, 2.0, 3.0, 4.0 and 5.0 mol%, respectively.

3.8 Possible PL and afterglow mechanisms of CaAl₂O₄:Dy³⁺

CaAl₂O₄ is an insulator with a bandgap of around 6.7 eV. Oxygen and calcium vacancies are intrinsic defects in CaAl₂O₄:Dy³⁺. On one hand, these intrinsic vacancies act as luminescence center of PL in CaAl₂O₄:Dy³⁺. The energy levels of oxygen vacancies in CaAl₂O₄ are proposed to work as electron traps because oxygen vacancies have lower energy levels than the luminescence center of PL with the result of a broadband PL spectrum peaking at about 400 nm. On the other hand, these intrinsic vacancies work as traps for charged carriers. For example, oxygen vacancies in CaAl₂O₄ are proposed to work as electron traps because they are positively charged whereas calcium vacancies are potential hole traps because they are positively charged. Dopant Dy³⁺ ions are incorporated into the host, a series of defect energy levels are introduced into the bandgap of CaAl₂O₄. The lowest energy level of the excited state of Dy³⁺ is known as \( ^{4}F_{9/2} \) while the energy levels of the ground state of Dy³⁺ are denoted as \( ^{4}H_{J} (J = 15/2–5/2) \). As evidenced by the characteristic emissions in Fig. 6, 7 and S3,‡ dopant Dy³⁺ ions act as luminescent center of PL in CaAl₂O₄:Dy³⁺ to yield narrowband emissions peaking at 482 and 574 nm, respectively. Furthermore, this extrinsic defect can work as electron trap because it is positively charged.

Fig. 10 schematically illustrates the PL and afterglow mechanisms of CaAl₂O₄:Dy³⁺. The energy levels of oxygen vacancies and dopant Dy³⁺ are depicted in the bandgap of CaAl₂O₄. As shown in Fig. 10, oxygen vacancies and Dy³⁺ activator are two independent sets of photon absorbers in CaAl₂O₄:Dy³⁺. Upon the ultraviolet excitation at 325 nm (3.82 eV), a portion of the incident photons are absorbed by the host due to the presence of a large number of oxygen vacancies in CaAl₂O₄ (process ①).
After non-radiative relaxations, the hot electrons are captured by either the oxygen vacancies (process ③) or the electron traps (process ⑤). The subsequent radiative recombination of electrons captured at oxygen vacancies with holes in the valence band yields the broadband emissions peaking at about 400 nm (process ④). Some of the excited photons absorbed by Dy$^{3+}$ activator in CaAl$_2$O$_4$ (process ⑥) produce the characteristic emissions of Dy$^{3+}$ (process ⑦): non-radiative relaxation of hot electrons at excited states of Dy$^{3+}$ to electron traps in CaAl$_2$O$_4$. Process ⑧: thermal release of electrons from electron traps to oxygen vacancies in CaAl$_2$O$_4$. Process ⑨: thermal release of electrons from electron traps to the excited states of Dy$^{3+}$.

Fig. 10 Schematic illustration on the PL and afterglow mechanisms of CaAl$_2$O$_4$:Dy$^{3+}$. Process ①: host absorption of incident photons due to the presence of intrinsic defects in CaAl$_2$O$_4$. Process ②: non-radiative relaxation of hot electrons to oxygen vacancies (luminescence center) in CaAl$_2$O$_4$. Process ③: non-radiative relaxation of hot electrons to electron traps in CaAl$_2$O$_4$. Process ④: radiative recombination of electrons captured at oxygen vacancy with holes in the valence band to yield the broadband emissions peaking at about 400 nm. Process ⑤: absorption of incident photons due to the transition $^6H_{15/2} \rightarrow ^3P_{3/2}$ of Dy$^{3+}$ (around 325 nm) in CaAl$_2$O$_4$. Process ⑥: electron transitions $^4F_{9/2} \rightarrow ^4H_{15/2}$ and $^4F_{9/2} \rightarrow ^4H_{13/2}$ of Dy$^{3+}$ to produce the characteristic emissions of Dy$^{3+}$. Process ⑦: non-radiative relaxation of hot electrons at excited states of Dy$^{3+}$ to electron traps in CaAl$_2$O$_4$. Process ⑧: thermal release of electrons from electron traps to oxygen vacancies in CaAl$_2$O$_4$. Process ⑨: thermal release of electrons from electron traps to the excited states of Dy$^{3+}$.

Once the ultraviolet excitation is ceased, processes ①–⑨ are stopped immediately. Under thermal activation, electrons can be released from the electron traps via processes ② and ⑤. Subsequently, afterglow with characteristic emissions of Dy$^{3+}$ is resulted when the electrons released from the excited states of Dy$^{3+}$ recombine radiatively with holes via process ③. Theoretically speaking, afterglow with broadband emissions peaking at about 400 nm should be observed when thermally detrapped electrons recombine radiatively with holes via process ③. In practice, such a broadband afterglow is negligible because its intensity is several orders of magnitude weaker than that of Dy$^{3+}$ related afterglows. Therefore, the afterglow spectrum of the CaAl$_2$O$_4$:Dy$^{3+}$ consists of two narrowband emissions peaking at 482 and 574 nm, respectively.

In the light of the PL and afterglow mechanisms of CaAl$_2$O$_4$:Dy$^{3+}$ in Fig. 10, electrons in the electron traps contribute to the PL under photoexication, too. Upon the photoexication, some electrons in the electron traps are detrapped via processes ② and ⑤. Broadband emissions peaking at around 400 nm can be expected when the detrapped electrons recombine radiatively with holes via the radiative process ③, and narrowband emissions of Dy$^{3+}$ can be resulted when the detrapped electrons recombine radiatively with holes via the radiative process ⑨. In most cases, however, such contributions are negligible because they are many times weaker than the emissions resulted directly from the ultraviolet photoexication. Only when the afterglow is sufficiently strong, such contributions to the PL are discernible. For example, the PL spectrum of Dy$^{3+}$ doped SrAl$_2$O$_4$ is the result of superposition of the broadband emissions of the host peaking at about 400 nm and another broadband emissions peaking at about 520 nm.

3.9 TL glow curves of CaAl$_2$O$_4$:Dy$^{3+}$

TL is an important tool to determine the activation energies (i.e., trap depths) of trapping levels in crystals. TL represents the TL glow curves of sol–gel combustion derived CaAl$_2$O$_4$:Dy$^{3+}$ with different doping concentrations. Each CaAl$_2$O$_4$:Dy$^{3+}$ phosphor was exposed to ultraviolet light of 254 nm for 5 min before the TL measurements. The temperature rising rate was 2 K min$^{-1}$. It can be seen that each TL glow curve in Fig. 11 exhibits an extremely broad and asymmetric profile, suggesting the presence of multiple trap levels in CaAl$_2$O$_4$:Dy$^{3+}$. Moreover, both the profile and the peak temperature of the TL glow curve are sensitive to the doping concentration: (i) the TL glow curve consists of a primary peak at around 350 K and a secondary peak at about 310 K when the doping concentration increases from 0.1 to 0.4 mol%; (ii) the TL glow curve has only one peak, which gradually shifts to higher temperature (from 350 to 375 K) as the doping concentration increases from 0.6 to 0.8 mol%; and (iii) the TL glow curve has only one peak, which gradually shifts to lower temperature (from 375 to 350 K) as the doping concentration increases further from 1.0 to 5.0 mol%. The vertical dash line in Fig. 11 marks the position of 350 K. With the vertical line as a guideline, the evolution of the peak temperature with the doping concentration can be identified clearly. Interestingly, the peak...
temperature of the TL glow curve of CaAl₂O₄:Dy³⁺ (0.8 mol%) is the highest (around 375 K) among the 10 phosphors under test, which coincides with the best afterglow performance of CaAl₂O₄:Dy³⁺ at optimal doping concentration of 0.8 mol%.

To understand the TL behavior of CaAl₂O₄:Dy³⁺, it is necessary to deconvolute each TL glow curve and evaluate the trapping parameters. The TL curve based on the general order function is given by the following equation:

\[
I(T) = n_0 \exp\left(\frac{-E}{kT}\right) \times \left[1 + \left(\frac{b - 1}{\beta}\right)s \int_{T_0}^{T} \exp\left(\frac{-E}{kT}\right) dT\right]^{-\frac{1}{b(\beta - 1)}},
\]

where \(I\) is the TL intensity at temperature \(T\), \(s\) is the pre-exponential factor with the unit of \(\text{s}^{-1}\), \(n_0\) is the concentration of trapped charges at time \(t = 0\), \(E\) is the trap depth, \(k\) is the Boltzmann constant, \(b\) is the order of kinetics, and \(\beta\) is the heating rate. In order to determine the kinetic parameters of the multiple traps, computerized glow curve deconvolution of the TL glow curves is carried out with general order kinetics using a computer program given by Chung et al.\(^4\)

Fig. 12 depicts the computerized glow curve deconvolution of the TL glow curve of CaAl₂O₄:Dy³⁺ at the doping concentration of 0.8 mol%. It is found that this TL glow curve can be described satisfactorily by using the general order kinetics to model 5 traps in CaAl₂O₄:Dy³⁺. The figure-of-merit (FOM) of the deconvolution is 2.071%. The kinetic parameters and the electron lifetime at room temperature (\(\tau_{300}\)) are summarized in Table 1 for each trap in the CaAl₂O₄:Dy³⁺ (0.8 mol%). As can be seen in Table 1, the calculated \(E\) values are 0.3919, 0.8219, 1.1466, 0.8403 and 1.0352 eV for the five traps in CaAl₂O₄:Dy³⁺ (0.8 mol%). Apparently, most of these traps are suitable for long-time afterglow.

In addition to the TL glow curve shown in Fig. 12, we have also deconvoluted the TL glow curves of CaAl₂O₄:Dy³⁺ with the doping concentrations of 0.1, 0.2, 0.4, 0.6, 1.0, 2.0, 3.0, 4.0 and 5.0 mol%, respectively. For the sake of brevity, the computerized glow curve deconvolutions of the TL glow curves of CaAl₂O₄:Dy³⁺ are shown in ESI as Fig. S9–S17.\(^†\) The kinetic parameters and the electron lifetime at room temperature are listed in Table S2\(^†\) for each trap in the CaAl₂O₄:Dy³⁺. Apparently, the parameters of electron traps in CaAl₂O₄:Dy³⁺ can be effectively tuned via the control of doping concentration, which in turn can be exploited to modify the brightness and duration of the afterglow of CaAl₂O₄:Dy³⁺.

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**Table 1** Kinetic parameters of the computerized glow curve deconvolution of the TL glow curve of CaAl₂O₄:Dy³⁺ (0.8 mol%). \(T_m\) represents the peak temperature, \(E\) is the trap-depth, \(s\) is the frequency factor, \(b\) is the order of kinetics, and \(\tau_{300}\) is the room temperature electron lifetime in the trap.

<table>
<thead>
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<th>Trap number</th>
<th>(T_m) (K)</th>
<th>(E) (eV)</th>
<th>(s) (s(^{-1}))</th>
<th>(b)</th>
<th>(\tau_{300}) (h)</th>
<th>FOM (%)</th>
</tr>
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<tr>
<td>1</td>
<td>339.25</td>
<td>0.3919</td>
<td>9.234 × 10(^5)</td>
<td>1.000</td>
<td>1.93 × 10(^6)</td>
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<tr>
<td>2</td>
<td>365.55</td>
<td>0.8219</td>
<td>4.810 × 10(^8)</td>
<td>1.482</td>
<td>2.11 × 10(^7)</td>
<td></td>
</tr>
<tr>
<td>3</td>
<td>386.45</td>
<td>1.1466</td>
<td>2.256 × 10(^8)</td>
<td>2.000</td>
<td>1.02 × 10(^6)</td>
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</tr>
<tr>
<td>4</td>
<td>418.95</td>
<td>0.8403</td>
<td>2.336 × 10(^7)</td>
<td>2.000</td>
<td>4.71 × 10(^5)</td>
<td></td>
</tr>
<tr>
<td>5</td>
<td>472.85</td>
<td>1.0352</td>
<td>1.683 × 10(^8)</td>
<td>2.000</td>
<td>1.59 × 10(^6)</td>
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4. Conclusions

With doping concentration varying in the range of 0.1–5.0 mol%, a series of CaAl2O4:Dy3+ phosphors have been synthesized via sol–gel combustion technique. The PL and afterglow mechanisms of sol–gel combustion derived CaAl2O4:Dy3+ are explored by means of XRD, SEM, TEM, EDX, XPS, PL, afterglow spectroscopy and TL dosimetry. The PL spectrum of CaAl2O4:Dy3+, which consists of a broad PL band peaking at about 400 nm and the characteristic emissions of Dy3+, indicates that oxygen vacancies and Dy3+ work as two independent sets of luminescence center of PL for CaAl2O4:Dy3+. As a contrast, the afterglow spectrum of CaAl2O4:Dy3+ consists of the characteristic emissions of Dy3+ only, verifying that dopant Dy3+ acts as the luminescence center of afterglow for CaAl2O4:Dy3+. The afterglow duration of CaAl2O4:Dy3+ is found to depend on the doping concentration, and 115 min-long afterglow is recorded for CaAl2O4:Dy3+ at the optimal doping concentration of 0.8 mol%. The profile and peak temperature of the TL glow curve are sensitive to the doping concentration. The TL glow curve of CaAl2O4:Dy3+ (0.8 mol%) exhibits a maximal peak temperature at around 375 K. A picture on the PL and afterglow mechanisms is given to reveal the processes of excitation, migration, trapping, detrapping, and radiative recombination of charged carriers in CaAl2O4:Dy3+.

Conflicts of interest

There are no conflicts to declare.

Acknowledgements

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