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# Transport properties of $\text{MoS}_2/\text{V}_7(\text{Bz})_8$ and graphene/ $\text{V}_7(\text{Bz})_8$ vdW junctions tuned by bias and gate voltages†

Hong Yu,<sup>a</sup> Danting Li,<sup>a</sup> Yan Shang,<sup>a</sup> Lei Pei,<sup>a</sup> Guiling Zhang,<sup>a</sup>  <sup>\*,a</sup> Hong Yan<sup>\*,a</sup> and Long Wang<sup>b</sup>

The  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$  and graphene/ $\text{V}_7(\text{Bz})_8$  vdW junctions are designed and the transport properties of their four-terminal devices are comparatively investigated based on density functional theory (DFT) and the nonequilibrium Green's function (NEGF) technique. The  $\text{MoS}_2$  and graphene nanoribbons act as the source-to-drain channel and the spin-polarized one-dimensional (1D) benzene–V multidecker complex nanowire ( $\text{V}_7(\text{Bz})_8$ ) serves as the gate channel. Gate voltages applied on  $\text{V}_7(\text{Bz})_8$  exert different influences of electron transport on  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$  and graphene/ $\text{V}_7(\text{Bz})_8$ . In  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$ , the interplay of source and gate bias potentials could induce promising properties such as negative differential resistance (NDR) behavior, output/input current switching, and spin-polarized currents. In contrast, the gate bias plays an insignificant effect on the transport along graphene in graphene/ $\text{V}_7(\text{Bz})_8$ . This dissimilarity is attributed to the fact that the conductivity follows the sequence of  $\text{MoS}_2 < \text{V}_7(\text{Bz})_8 < \text{graphene}$ . These transport characteristics are examined by analyzing the conductivity, the currents, the local density of states (LDOS), and the transmission spectra. These results are valuable in designing multi-terminal nanoelectronic devices.

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## 1. Introduction

Two-dimensional (2D) materials have become the up-to-date focal point of research in past decades owing their prominent electronic, optical, and magnetism properties, which are used in wide fields from field effect transistors (FETs),<sup>1–3</sup> optoelectronics,<sup>4–7</sup> to spintronics.<sup>8–10</sup> As the brightest star in the 2D material family, graphene has attracted constant academic and individual research interest and enthusiasm.<sup>11–17</sup> Inspired by these pioneering works, various kinds of 2D materials, for example, 2D transition metal dichalcogenides (TMDs),<sup>18–22</sup> have sprung up and become research hotspots. Among TMDs,  $\text{MoS}_2$  nanosheets have drawn special attention due to their high stability, good semi-conductivity, large surface, *etc.*, which are desirable in nano electronic devices.<sup>23</sup> In recent years, numerous efforts have been devoted to fabricate van der Waals (vdW) heterojunctions of 2D materials by stacking alien components using various chemical techniques for engineering expected and improved properties toward practical

applications.<sup>24–31</sup> The carrier mobility of 2D materials can be retained from vdW interactions, since it does not destroy the bonding properties in intralayers. The interlayer coupling between the two vdW-stacked 2D layers can result in novel physical properties. Many theoretical works have focused on studying the transport properties of these 2D heterojunctions by constructing two-terminal devices. However, in experiments, electron transport is usually measured using a four-probe apparatus.<sup>32–37</sup> On the other hand, for most practical applications, multi-terminal electronic devices are needed. Therefore, it is highly necessary to theoretically investigate the transport properties of multi-terminal devices, which is still less studied to date.

Simultaneously, with the approaching physical limit of silicon based transistors, seeking emerging device architectures and new electronic materials has become an essential topic to meet the Moore's law. Recently, spintronic devices have received extensive attentions due to the combined merits of the electron transporting, the magnetic moment, and the electron spin. Therefore, introducing spin-polarized component into 2D materials to construct vdW heterojunctions is anticipated to be an effective strategy to explore novel functional materials with desirable properties.<sup>38–41</sup> However, as far as we know, studies on the spin-polarized 2D vdW heterojunctions are still scarce from both experimental and theoretical points of view. The one-dimensional (1D) benzene–V multidecker complexes nanowire,  $\text{V}_n(\text{Bz})_m$ , not only has been successfully synthesized from

<sup>a</sup>School of Materials Science and Chemical Engineering, Harbin University of Science and Technology, Harbin 150040, P. R. China. E-mail: guiling-002@163.com; yanhong204821@aliyun.com

<sup>b</sup>HeiLongjiang Construction Investment Group Co. Ltd, No. 523, Sanda Dongli Road, Harbin, 150040, P. R. China

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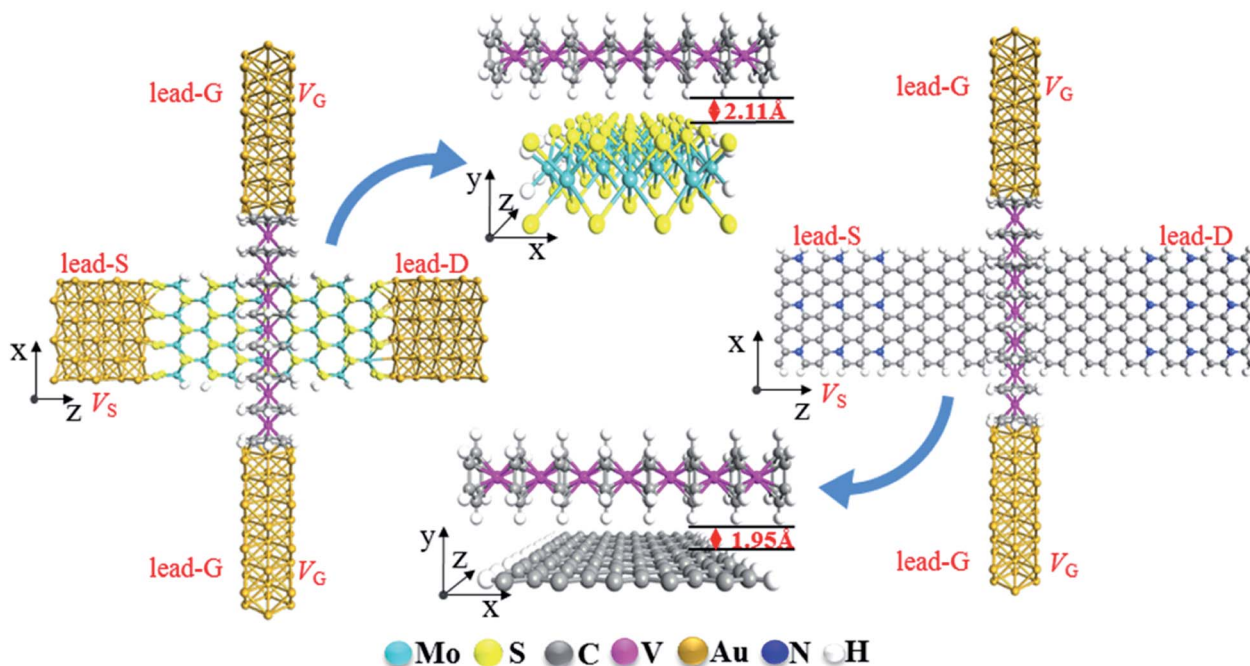


Fig. 1 Schematic plots of  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$  and  $\text{Graphene}/\text{V}_7(\text{Bz})_8$  four-terminal devices. Where lead-S, lead-D, and lead-G represent the source, drain, and gate leads, respectively.  $V_S$  and  $V_G$  represent the bias voltage applied on lead-S and the gate voltage applied on lead-G, respectively.

reaction of laser-vaporized metal atoms with  $\text{C}_6\text{H}_6$  in a He atmosphere, but also have been imaged or detected using scanning tunneling microscopy (STM), UV-vis and IR, electron paramagnetic resonance (EPR), time-of-flight mass spectroscopies, and photoionization spectroscopies.<sup>42–45</sup> Experimental works as well as theoretical studies have confirmed that the unpaired electrons on the V atoms are coupled ferromagnetically (FM), and also suggest that the ground state of  $\text{V}_n(\text{Bz})_m$  exhibits half-metallicity and spin filter effect.<sup>44,46</sup>

Further onward, stimulated by the splendid work on  $\text{MoS}_2$  transistors using a single-walled carbon nanotube as the gate electrodes to tune the transport properties,<sup>47</sup> we construct four-terminal devices for  $\text{MoS}_2$  and graphene heterojunctions:  $\text{MoS}_2$  and graphene are connected to the source and drain leads,

while the spin polarized  $\text{V}_n(\text{Bz})_m$  nanowire serves as gate terminals. Bias voltages and gate voltages are applied on the source lead and gate lead, respectively. The synergistic tuning effects of the bias and gate voltages on transport properties of  $\text{MoS}_2$  and graphene are investigated by employing the density functional theory (DFT) with non-equilibrium Green's function (NEGF) methodology.

## 2. Models and computational methods

Fig. 1 shows the structures of the four-terminal devices for  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$  and  $\text{graphene}/\text{V}_7(\text{Bz})_8$  vdW heterojunctions.  $\text{MoS}_2$

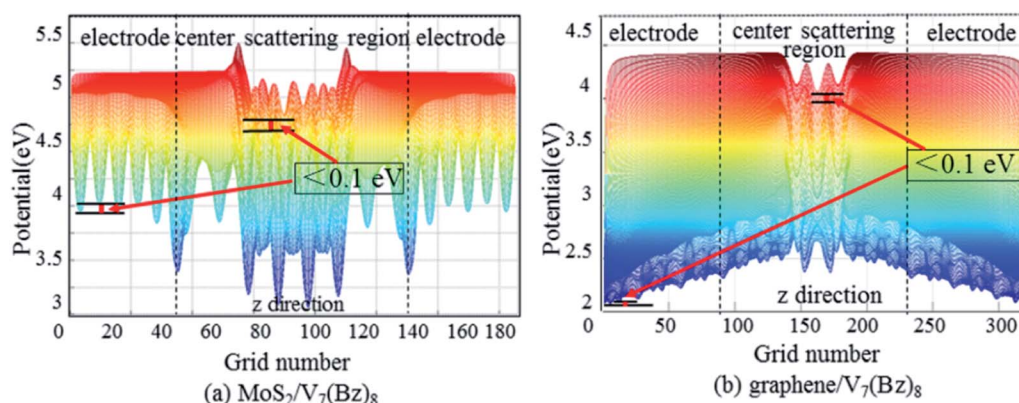


Fig. 2 The calculated potential distribution of the buffers near the leads in the two four-terminal devices, (a) is that in  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$  device and (b) is that in  $\text{graphene}/\text{V}_7(\text{Bz})_8$  device.

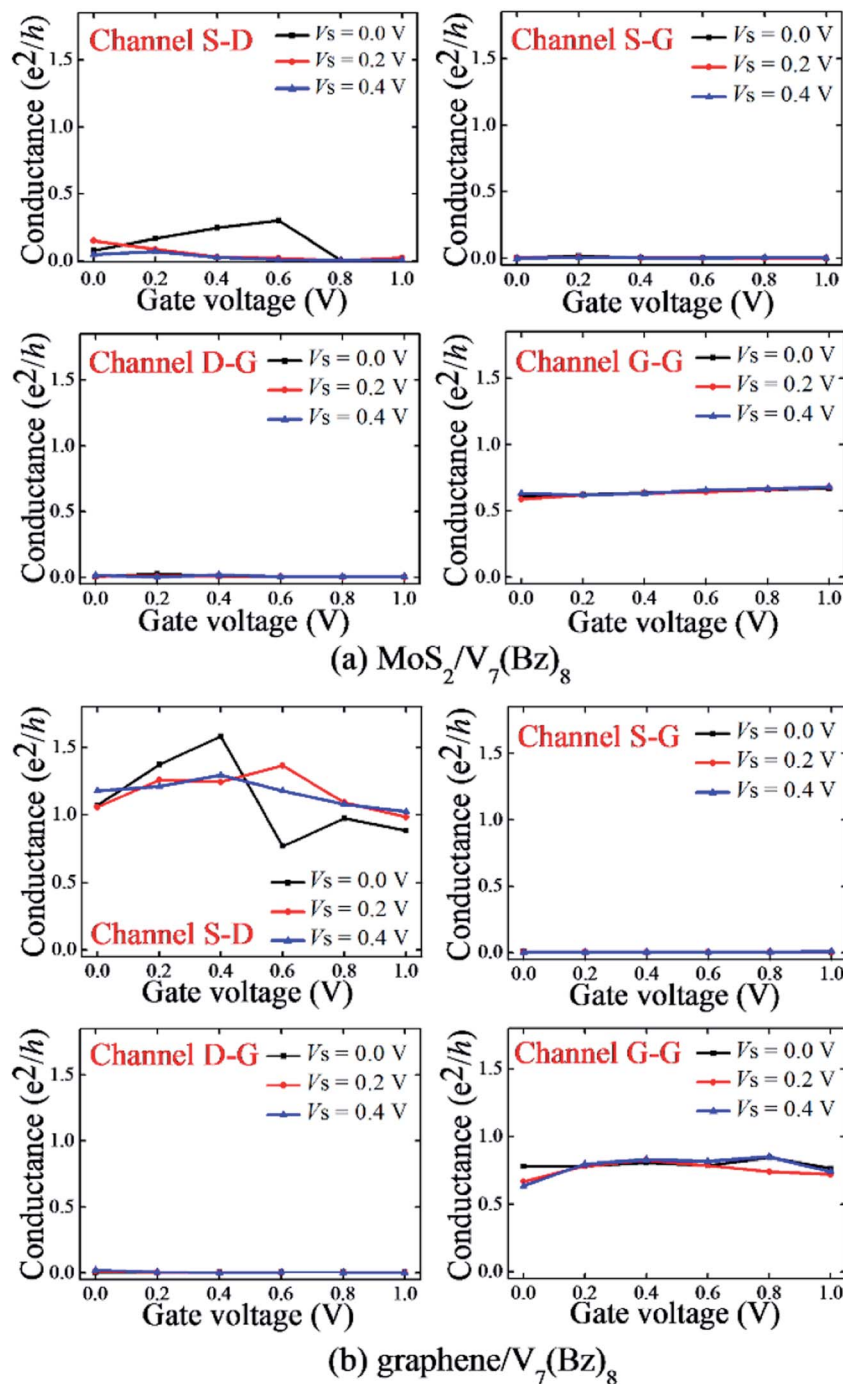


Fig. 3 Conductance for channel S-D, channel S-G, channel D-G, and channel G-G as a function of  $V_G$  at fixed values of  $V_S = 0.0, 0.2$ , or  $0.4$  V of the four-terminal  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$  and  $\text{graphene}/\text{V}_7(\text{Bz})_8$  devices. (a and b) Are for the  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$  and  $\text{graphene}/\text{V}_7(\text{Bz})_8$  devices, respectively.

and graphene offer the source-to-drain scatter region, and  $\text{V}_7(\text{Bz})_8$  nanowire provides the gate channel. For the sake of simplicity, the leads are denoted as lead-S, lead-D, and lead-G for the source, drain, and gate leads, respectively. The source-to-drain direction is defined as the  $z$  direction, and the gate to gate direction is referred as the  $x$  direction. The  $\text{MoS}_2$  and graphene are placed in the  $x \times z$  plane. The  $\text{MoS}_2$  scatter region contains 9-layered S and 8-layered Mo in the  $z$  direction, as well

as 4-layered S and 3-layered Mo in the  $x$  direction, which extends a  $11.33 \text{ \AA} \times 21.90 \text{ \AA}$  ( $x \times z$ ) plane. A similar size of graphene ribbon in the  $x$  and  $z$  direction ( $11.07 \text{ \AA} \times 19.89 \text{ \AA}$ ) is tailored as the scatter region, too. Such scatter region is long enough in the  $z$  direction to ignore the interaction between lead-S and lead-D. This can be confirmed by the small potential changes near the electrodes ( $<0.1$  eV) calculated for the two devices (Fig. 2). Every dangling bonds at edges are saturated by H atoms to stabilize



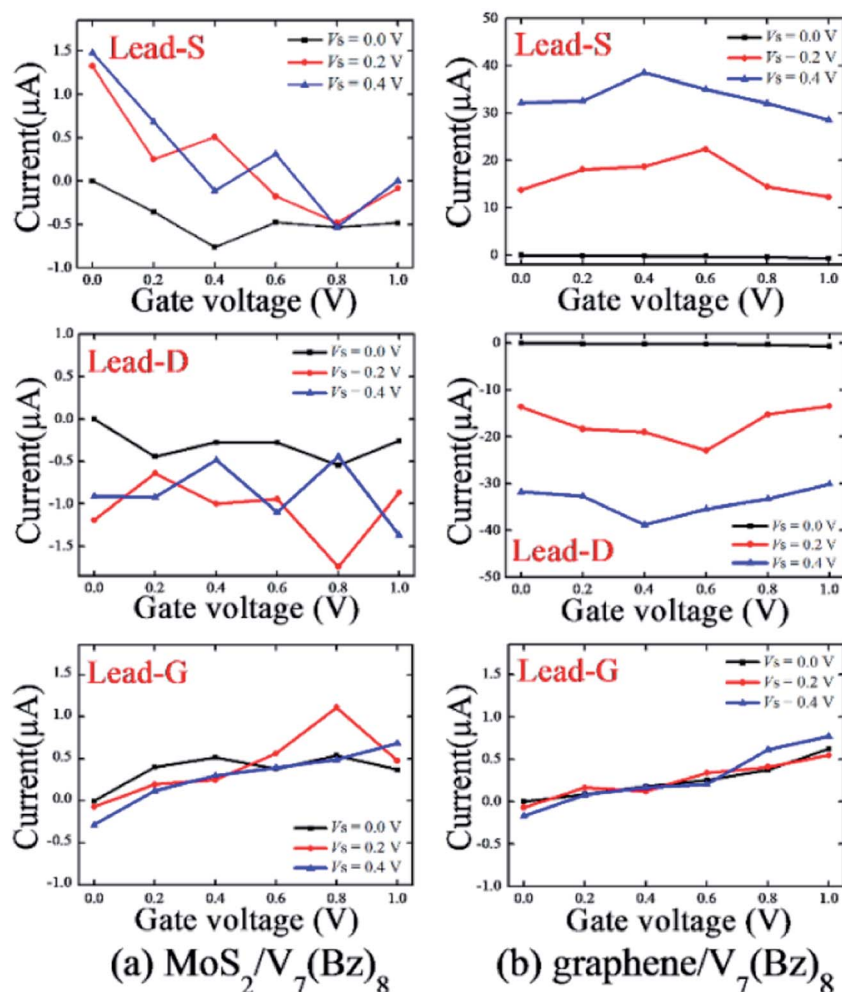


Fig. 4 Total currents for lead-S, lead-D, and lead-G as a function of  $V_G$  at fixed values of  $V_S = 0.0, 0.2$ , or  $0.4$  V of the four-terminal devices, where (a) shows those in  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$  devices, and (b) is in  $\text{graphene}/\text{V}_7(\text{Bz})_8$ .

the system. The  $\text{V}_7(\text{Bz})_8$  nanowire overlies the center of  $\text{MoS}_2$  and graphene surfaces with a vdW distance from them, and its longitudinal axis is parallel to the  $x$  direction and perpendicular to the  $z$  direction. Each lead-G in  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$  and  $\text{graphene}/\text{V}_7(\text{Bz})_8$  is modeled by a  $\text{Au}(100)-(3 \times 3)$  surface with 8 layers. The Au-C distance is set to be  $2.05 \text{ \AA}$  based on their covalent radius.<sup>48</sup>  $\text{MoS}_2$  is semi-conductor, so we use  $\text{Au}(100)-(6 \times 3)$  surfaces with 5 layers as the lead-S and lead-D for  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$ . As the S atom has good affinity with the gold surface, dithiolate derivatives have been used for the construction of metal/scatter region/metal devices in general.<sup>49–51</sup> Therefore, in the present work, we also use the S atom layer of  $\text{MoS}_2$  to link the Au electrodes. The S-Au distance was set as  $2.34 \text{ \AA}$  according to the reported literature.<sup>52</sup> To ensure that the bonds match well between graphene and leads, we elongate the graphene additionally by  $31.26 \text{ \AA}$  as the lead-S and lead-D. Since graphene nanoribbon is semi-conductor along the armchair direction, we use the N atoms to dope the graphene lead-S and lead-D since the N atoms can introduce additional  $\pi$  electrons to make the two leads become conductor.<sup>53,54</sup> A vacuum of  $20 \text{ \AA}$  in  $y$  direction is involved to eliminate the coupling between adjacent images.

The DFT method with generalized gradient approximation (GGA) implemented in ATK package is employed to optimize the four-terminal structures. For both  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$  and  $\text{graphene}/\text{V}_7(\text{Bz})_8$  four-terminal devices, all the leads are frozen and the scatter regions are optimized. The self-consistent total energies are converged to  $10^{-4} \text{ eV}$  and the forces are converged to  $0.05 \text{ eV \AA}^{-1}$ . The optimized cartesian coordinates of all the models are given in Tables S1 and S2 in the ESI.† The optimized distance is  $2.11 \text{ \AA}$  from  $\text{MoS}_2$  to  $\text{V}_7(\text{Bz})_8$  and  $1.95 \text{ \AA}$  from graphene to  $\text{V}_7(\text{Bz})_8$ , a typical vdW interaction. Then, the transport properties of the optimized devices are calculated using DFT with NEGF methodology within Nanodcal software package.<sup>55,56</sup> The linear combination of atomic orbitals (LCAO) is employed to expand physical parameters. The standard nonlocal norm-conserving pseudopotentials is used to describe the atomic core, and the double-zeta polarized (DZP) basis set is for valence electronic orbitals.<sup>57,58</sup> The exchange-correlation function is considered by the local density approximation (LDA).<sup>58,59</sup> The  $k$ -point for the central region is meshed as  $1 \times 1 \times 1$  ( $x \times y \times z$ ), and that is  $1 \times 1 \times 100$  for lead-S and lead-D,  $100 \times 1 \times 1$  for lead-G. It is worth noting that, in all calculations the spin of V



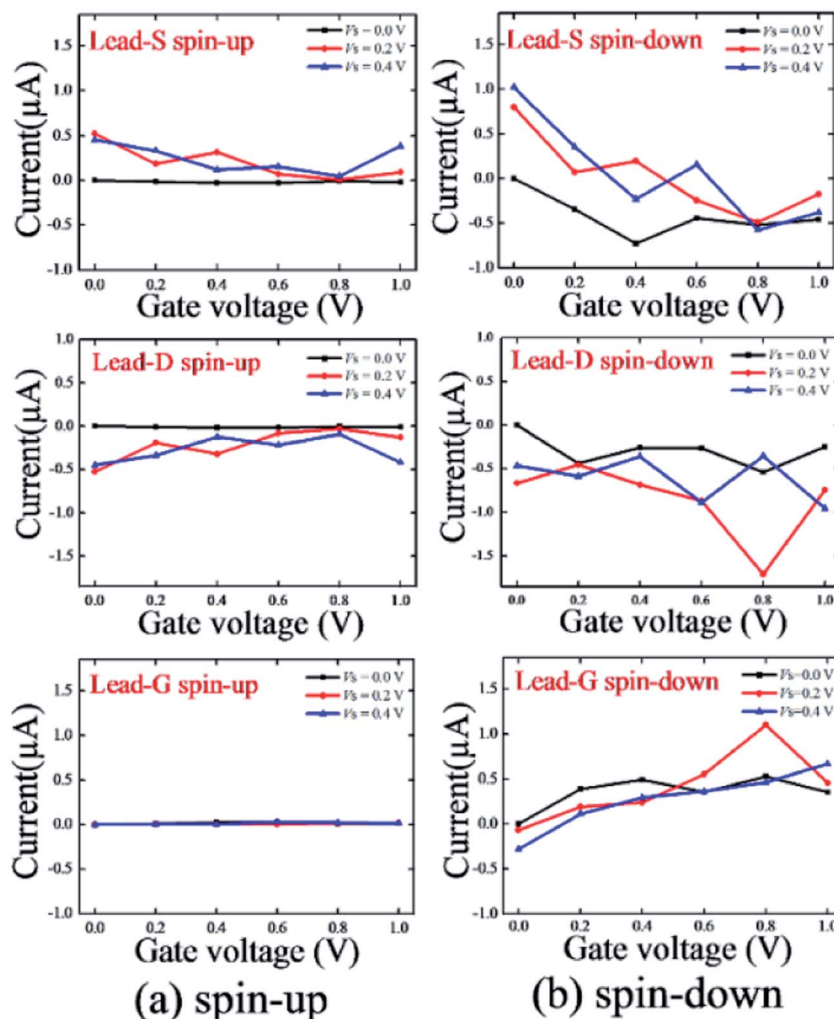


Fig. 5 Spin currents for lead-S, lead-D, and lead-G as a function of  $V_G$  at fixed values of  $V_S = 0.0, 0.2$ , or  $0.4$  V of the four-terminal  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$  device. (a) Depicts the spin-up states of  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$ , and (b) depicts those for spin-down state, respectively.

atoms is considered by using LAD +  $U$  scheme, and the U-J is set as  $3.0 \text{ eV}$ .<sup>60</sup> 160 Rydberg cutoff energy is applied.

The spin-dependent current in leads for multi-terminal system can be obtained using the Landauer-Büttiker formula (1):<sup>57,61,62</sup>

$$I_\sigma = \frac{e}{h} \sum_{\beta} \int_{-\infty}^{+\infty} dE T_{\alpha\beta}(E, V_\alpha, V_\beta) [f_\beta(E) - f_\alpha(E)] \quad (1)$$

where  $f_\alpha$  ( $f_\beta$ ) describes the Fermi distribution functions of lead- $\alpha$  (lead- $\beta$ ).  $V_\alpha$  ( $V_\beta$ ) is the bias voltages applied on lead- $\alpha$  (lead- $\beta$ ).  $\sigma \equiv \uparrow, \downarrow$  is the spin index. The transmission coefficient  $T_{\alpha\beta}$  between lead- $\alpha$  and lead- $\beta$  is dependent on  $E$ ,  $V_\alpha$ , and  $V_\beta$ . And the total current is  $I = I_\uparrow + I_\downarrow$ .

Considering the multi-terminals, the conductance between lead- $\alpha$  and lead- $\beta$  can be evaluated by formula (2):

$$G_{\alpha\beta} = \frac{e^2}{h} T_{\alpha\beta}(E) \quad (2)$$

### 3. Results and discussion

In this section, we mainly document on how the bias voltage ( $V_S$ ) and the gate voltage ( $V_G$ ) synergistically influence the transport properties of  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$  and graphene/ $\text{V}_7(\text{Bz})_8$  vdW junctions from the aspects of conductance, currents, LDOS, and transmission. The  $V_S$  is set as  $0.0, 0.2, 0.4 \text{ V}$ , respectively. And at each certain  $V_S$ , the  $V_G$  ranges from  $0.0$  to  $1.0 \text{ V}$  by a step of  $0.2 \text{ V}$ .

#### 3.1 Conductivity

To convenient description, the channels between any two terminals are denoted as channel S-D, channel S-G, channel D-G, and channel G-G.

Fig. 3 shows the calculated conductance of each channel of  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$  and graphene/ $\text{V}_7(\text{Bz})_8$  devices with the variation of  $V_S$  and  $V_G$ . For the sake of comparison, the scales of conductance in vertical-axis in Fig. 3 are the same for all the channels, and Fig. S1† with adapted vertical-axis scales is supplied as



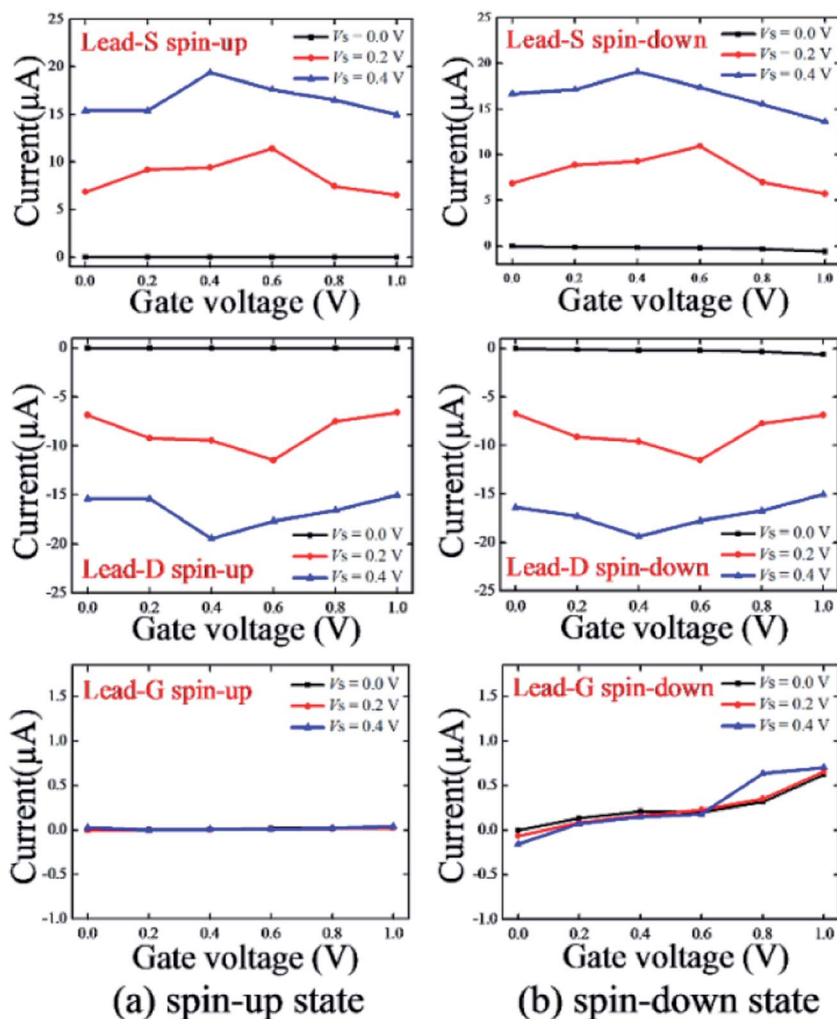


Fig. 6 Spin currents for lead-S, lead-D, and lead-G as a function of  $V_G$  at fixed values of  $V_S = 0.0, 0.2$ , or  $0.4$  V of the four-terminal graphene/ $V_7(Bz)_8$  device. (a) Depicts the spin-up states, and (b) depicts those for spin-down state, respectively.

ESI.† The conductance is bigger than  $0.8e^2/h$ , in the range of  $0.5\text{--}1.0e^2/h$ , and smaller than  $0.3e^2/h$  for graphene channel S-D,  $V_7(Bz)_8$  channel G-G, and  $MoS_2$  channel S-D, respectively. That is, the conductivity follows the sequence of  $MoS_2 < V_7(Bz)_8 < \text{graphene}$ . Therefore, using  $V_7(Bz)_8$  as gate may play different effects on tuning electron transporting through  $MoS_2$  and graphene. Since the  $V_7(Bz)_8$  nanowire attaches to  $MoS_2$  and graphene *via* vdW interaction, the conductance of channels S-G and D-G are very small ( $<0.03e^2/h$ ). The potentials from gate and from source couple in different degrees in  $MoS_2/V_7(Bz)_8$  and graphene/ $V_7(Bz)_8$  under various  $V_S$  and  $V_G$ . This integrates the change of resistance from source to drain, which makes the conductance goes up or down.

### 3.2 Current

To further understand the transport property of the four-terminal devices, we calculated the total currents,  $I_S$ ,  $I_D$ , and  $I_G$ , passing through lead-S, lead-D, and lead-G, respectively. Currents from lead to center region are input and defined as

positive values, while those from center region to lead are output and appointed as negative values.

Fig. 4 displays the total currents through the leads with the variation of  $V_S$  and  $V_G$  applied on  $MoS_2/V_7(Bz)_8$  and graphene/ $V_7(Bz)_8$  devices. Multi-channels in four-terminal devices could cause complicated transport properties aroused from the intricacy of channel couplings. Actually, the magnitude of each lead current comes from the synergistic action of all pathways. For example, though lead-S current is directly related to channels S-D and S-G, channels S-D and S-G are simultaneously influenced by channels D-G and G-G.

Now we consider the total currents,  $I_S$ ,  $I_D$ , and  $I_G$  of  $MoS_2/V_7(Bz)_8$  device (Fig. 4a). With free gate voltage, a small amount of current leaks off lead-G when applied  $V_S = 0.2$  and  $0.4$  V, while the lead-G current becomes input once  $V_G$  imposed, which could exert considerable effect on the source-to-drain transporting. The lead-D currents are always output owing to its low voltage potential compared with other three leads. At  $V_S = 0.0$  V, lead-S and lead-D give output currents with the absolute values of  $I_S < 1.0$   $\mu\text{A}$  and  $I_D < 0.5$   $\mu\text{A}$ , which are induced by the applied gate



**Table 1** Local density of states (LDOS) of the four-terminal MoS<sub>2</sub>/V<sub>7</sub>(Bz)<sub>8</sub> and graphene/V<sub>7</sub>(Bz)<sub>8</sub> devices at various V<sub>G</sub> exemplified by V<sub>S</sub> = 0.2 V

V <sub>G</sub> (V)	LDOS at V <sub>S</sub> = 0.2 V			
	MoS <sub>2</sub> /V <sub>7</sub> (Bz) <sub>8</sub> device		Graphene/V <sub>7</sub> (Bz) <sub>8</sub> device	
	Spin-up	Spin-down	Spin-up	Spin-down
0.0				
0.2				
0.4				
0.6				
0.8				
1.0				

voltages. Interestingly, regardless of any V<sub>S</sub>, the I<sub>S</sub>, I<sub>D</sub>, and I<sub>G</sub> curves present a pattern of up and down oscillation with the changing of V<sub>G</sub>, resulting in NDR peaks. Of course, the magnitude, the position, and the numbers of these NDR peaks are different for different leads at various V<sub>S</sub> and V<sub>G</sub>. In fact, conductivity, as a complicated phenomenon, relates to comprehensive factors such as scatter region structure, scatter region-lead interaction, couplings between channels, changes of resistance induced by voltages, *etc.* For example, the asymmetric character of MoS<sub>2</sub> structure not only induces different resistance for different channels, but also induces different left and right MoS<sub>2</sub>-lead interactions. Therefore, at zero V<sub>S</sub> bias, the I<sub>S</sub> and I<sub>D</sub> curves give different shapes. All these comprehensive factors may also be responsible for the dissimilarity of the NDR behaviors of I<sub>S</sub>, I<sub>D</sub>, and I<sub>G</sub>. When nonzero V<sub>S</sub> is applied, as the V<sub>G</sub> increases, the input I<sub>S</sub> current decreases as a whole and could switch to output

at certain V<sub>G</sub>, *e.g.*, at V<sub>G</sub> = 0.6 V for V<sub>S</sub> = 0.2 V and at V<sub>G</sub> = 0.4 and 0.8 V for V<sub>S</sub> = 0.4 V. Therefore, the currents across lead-S are rather dependent on the interplay of V<sub>S</sub> and V<sub>G</sub>. At V<sub>S</sub> = 0.2 and 0.4 V, the absolute values of I<sub>D</sub> rise comparing with V<sub>S</sub> = 0.0 V. This is a reasonable result from the higher potential injected into the channel S-D. The magnitude of the output currents of lead-D tends to be balanced around 1.0 μA. All these properties are much desirable for designing functional devices with fascinating characteristics such as NDR behavior and input/output switching.

As to the graphene/V<sub>7</sub>(Bz)<sub>8</sub> device, the total currents of I<sub>S</sub>, I<sub>D</sub>, and I<sub>G</sub> display a much different feature from MoS<sub>2</sub>/V<sub>7</sub>(Bz)<sub>8</sub> as shown from Fig. 4b. Intuitively, the gate V<sub>G</sub> plays a small role on the carries transport *via* the graphene plane. This result can be deduced from two phenomena: one is that the magnitudes of I<sub>S</sub> and I<sub>D</sub> (10–40 μA) are much higher than I<sub>G</sub> (<1.0 μA) under applied V<sub>G</sub>; another is that the input values of I<sub>S</sub> are almost



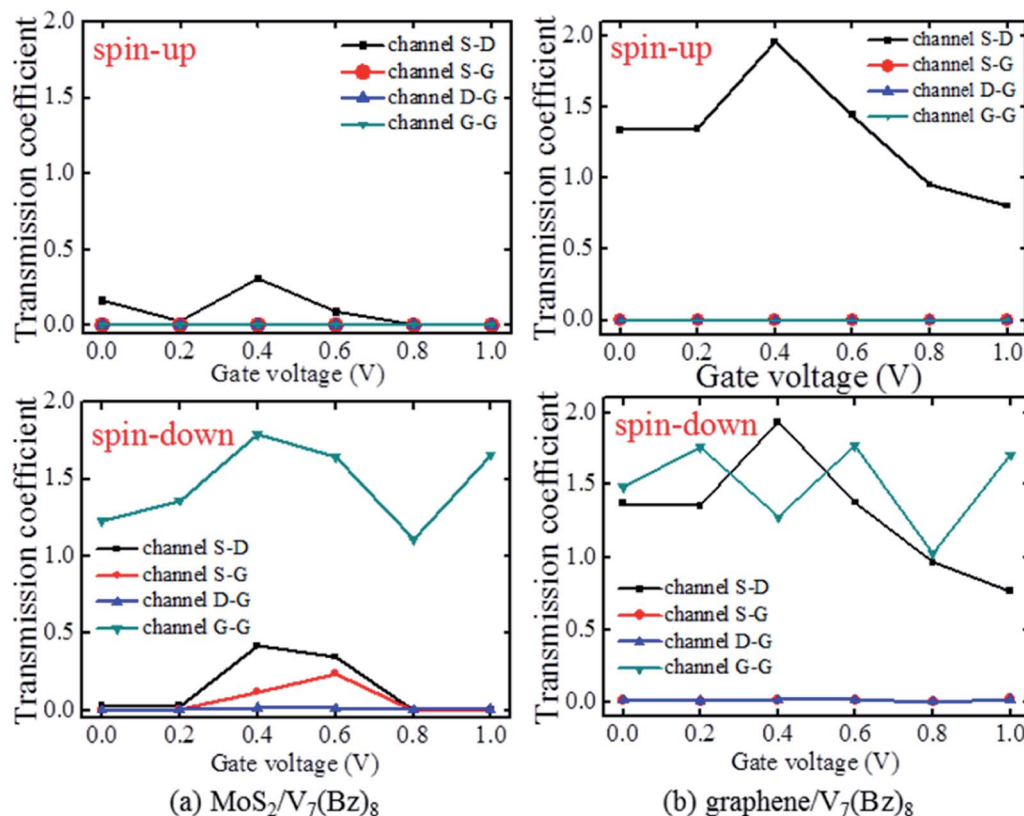


Fig. 7 Transmission coefficient of all channels in two devices with variational  $V_G$  from 0.0 to 1.0 V at  $V_S = 0$  V. (a) Shows spin-up and spin-down state in  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$  device, and (b) is that for  $\text{graphene}/\text{V}_7(\text{Bz})_8$  device.

equal to the output data of  $I_D$  at given  $V_S$  and  $V_G$ , indicating that the injected  $I_S$  almost arrives to the lead-D completely. No evident NDR behavior is observed for  $\text{graphene}/\text{V}_7(\text{Bz})_8$ . In addition, unlike in  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$  which there exists input/output switching of lead-S, the current of lead-S in  $\text{graphene}/\text{V}_7(\text{Bz})_8$  is always input within the considered gate bias. The reason for these different phenomena between  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$  and  $\text{graphene}/\text{V}_7(\text{Bz})_8$  lies in the fact that graphene has higher conductivity than  $\text{MoS}_2$ .

To investigate how the spin-polarized character of  $\text{V}_7(\text{Bz})_8$  influence the transport property of  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$  and  $\text{graphene}/\text{V}_7(\text{Bz})_8$ , we computed the spin-up currents,  $I_{S\uparrow}$ ,  $I_{D\uparrow}$ , and  $I_{G\uparrow}$  as well as the spin-down currents,  $I_{S\downarrow}$ ,  $I_{D\downarrow}$ , and  $I_{G\downarrow}$ , for lead-S, lead-D, and lead-G, respectively. The results are plotted in Fig. 5 and 6 as a function of  $V_G$  at a given  $V_S$ .

As expected, for  $\text{V}_7(\text{Bz})_8$  in both two devices, the  $I_{G\uparrow}$  remains nearly zero, while the  $I_{G\downarrow}$  possesses the same magnitude to the total  $I_G$ . This thoroughly inherits the intrinsic spin-polarized feature of pure  $\text{V}_7(\text{Bz})_8$ .<sup>60,63</sup> Quite significantly, the spin polarized feature of  $\text{V}_7(\text{Bz})_8$  induces different polarized transport property from  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$  to  $\text{graphene}/\text{V}_7(\text{Bz})_8$ . In the case of  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$ , spin-up and spin-down channels of lead-S and lead-G are split by the polarization property of  $\text{V}_7(\text{Bz})_8$ . At  $V_S = 0.0$  V, the spin-up paths of both lead-S and lead-D of  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$  are closed with almost zero current, and the total currents  $I_S$  and  $I_D$  entirely come from the spin-down  $I_{S\downarrow}$  and  $I_{D\downarrow}$ . This is mainly due to the fact that the spin-up channel of  $\text{V}_7(\text{Bz})_8$

is almost closed and hence cannot cause perturbation on the spin-up pathway of  $\text{MoS}_2$ . On the contrary, the spin-down state of  $\text{V}_7(\text{Bz})_8$  dominates electron transporting, enabling strong coupling with the spin-down channel of  $\text{MoS}_2$ . Under nonzero  $V_S$ , the currents through lead-S and lead-D in  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$  are polarized by a large extent. For the spin-up state, the currents of  $I_{S\uparrow}$  are all input and those of  $I_{D\uparrow}$  are all output. Both  $I_{S\uparrow}$  and  $I_{D\uparrow}$  curves of  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$  change smoothly in a small scale of 0.0–0.5  $\mu\text{A}$  along the variation of  $V_G$ . In striking contrast, the spin-down currents of  $I_{S\downarrow}$  and  $I_{D\downarrow}$  fluctuate significantly with the changing of  $V_G$ , with a large range of from 1.0 to  $-0.5$   $\mu\text{A}$  for  $I_{S\downarrow}$  and from 0.0 to  $-1.7$   $\mu\text{A}$  for  $I_{D\downarrow}$ . Observed carefully, one can find that the changing trends of  $I_{S\downarrow}$  and  $I_{D\downarrow}$  curves are roughly analogous to those of the total  $I_S$  and  $I_D$ . The NDR behavior and the input/output switch character of  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$  are mainly dominated by the spin-down state. Now turn to  $\text{graphene}/\text{V}_7(\text{Bz})_8$ , quite different from  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$ , no matter what  $V_S$  and  $V_G$  applied, the spin-up currents are almost equivalent to spin-down values for both lead-S and lead-D. Though the polarized characteristic of  $\text{V}_7(\text{Bz})_8$  is preserved in  $\text{graphene}/\text{V}_7(\text{Bz})_8$ , its conductivity is too much weaker than graphene that can not exert valid influence on channel S–D.

### 3.3 Local density of states

Table 1 displays the computed local density of states (LDOS) results of two devices under  $V_S = 0.2$  V, where the even LDOS



distribution suggests an effective transport channel. The scenario of the polarized transport behavior of  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$  leads becomes more obvious by analyzing the LDOS distributions. As representative example, it is clear that the LDOS of the spin-down state of  $\text{V}_7(\text{Bz})_8$  delocalizes stronger than the spin-up state, which could bring about coupling with the spin-down state of  $\text{MoS}_2$ , and consequently, lead to the spin-polarized splitting of  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$ . The electronic potential on the whole system changes along with the relative magnitude of  $V_s$  and  $V_g$ , as can be reflected from the LDOS distributions. Clearly, charge carries in the spin-up state of either channel S-D or channel G-G are blocked more seriously than the spin-down state. As for the  $\text{graphene}/\text{V}_7(\text{Bz})_8$  device, the spin-down passage of  $\text{V}_7(\text{Bz})_8$  is open with uniform LDOS distribution while the spin-up path is closed with almost no LDOS. Despite of the maintained polarized character of  $\text{V}_7(\text{Bz})_8$ , spin-up and spin-down channels along graphene has nearly identical LDOS spreading. This again demonstrates the unpolarized character of lead-S and lead-D currents.

### 3.4 Transmission coefficients and spectra

To further shed light on how the gate bias tune the transport property of  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$  and  $\text{graphene}/\text{V}_7(\text{Bz})_8$  devices, we computed spin polarized transmission coefficients (TC) of each

channel at the case of  $V_s = 0.0$  V and  $V_g$  from 0.0 to 1.0 V, and the results are shown in Fig. 7.

Evidently, in both  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$  and  $\text{graphene}/\text{V}_7(\text{Bz})_8$  devices, spin-down channel G-G composed of  $\text{V}_7(\text{Bz})_8$  has much higher TC than other channels, again demonstrating the preserved striking half-metallic character of  $\text{V}_7(\text{Bz})_8$ . One can find that, in  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$ , the main channel S-D consisting of  $\text{MoS}_2$  can not be largely influenced by the spin polarized character of the gate  $\text{V}_7(\text{Bz})_8$ , as can be clearly seen from the comparable spin-up and spin-down TC. However, the spin polarized  $\text{V}_7(\text{Bz})_8$  directly consists of channels S-G and D-G, and thereby, it could induce spin split of channels S-G and D-G and generate larger spin-down TC than spin-up TC. Therefore, the spin polarized transport character of  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$  mainly derives from the perturbation of the spin-down state of  $\text{V}_7(\text{Bz})_8$  upon the channels S-G and D-G. Regarding  $\text{graphene}/\text{V}_7(\text{Bz})_8$ , charge carriers can easily flow from lead-S to lead-D since channel S-D has large TC, while other channels are almost closed with nearly zero TC except the spin-down channel G-G. In addition, spin-up and spin-down states of channel S-D display the same TC curves, illustrating the ignoring polarization influence of  $\text{V}_7(\text{Bz})_8$  on graphene.

Furtherly, we calculated the transmission spectra (TS) of the main channel S-D for  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$  and  $\text{graphene}/\text{V}_7(\text{Bz})_8$

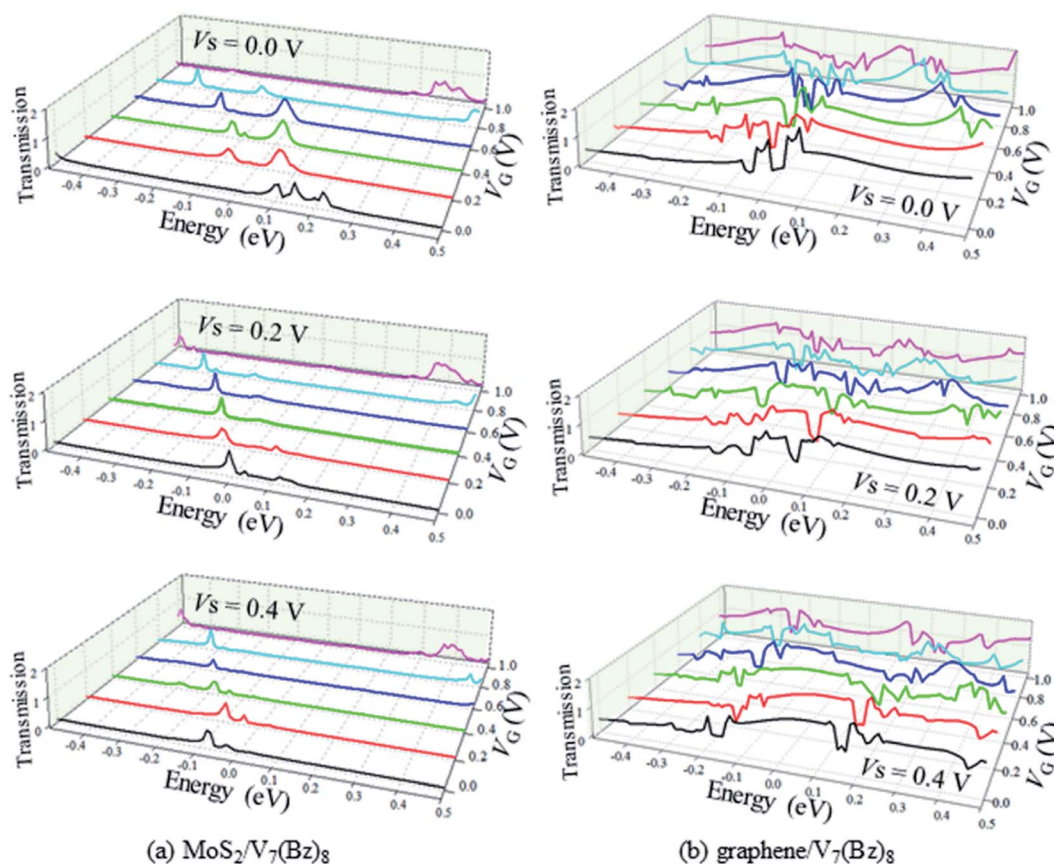


Fig. 8 Transmission spectra of channel S-D in  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$  and  $\text{graphene}/\text{V}_7(\text{Bz})_8$  device with variational  $V_g$  from 0.0 to 1.0 V at a certain  $V_s$ , where (a) depicts the case in  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$ , and (b) is that for  $\text{graphene}/\text{V}_7(\text{Bz})_8$ , respectively.



devices at each  $V_S$  and  $V_G$ , as plotted in Fig. 8. In the case of  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$ , at  $V_S = 0.0$  V and  $V_G = 0.0$  V, TS peak located above the  $E_F$  mainly arises from the native  $\text{MoS}_2$ . When adding  $V_G$  merely, an extra peak appears near the  $E_F$  owing to the gate carries injection into the channel S–D of  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$ . Once adding  $V_S$  together with  $V_G$ , only one peak exists below the  $E_F$  due to the suppression effect of the incoming  $V_G$  potential upon the  $V_S$ . Furthermore, with the increasing of  $V_G$ , such suppression effect in  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$  becomes more significant, as the TS peak moves farther away from the  $E_F$ . In contrast, regardless of  $V_S$  and  $V_G$ , lots of TS peaks of channel S–D appear around the  $E_F$  of graphene/ $\text{V}_7(\text{Bz})_8$ , again indicating the high conductivity of graphene.

## 4. Conclusion

The  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$  and graphene/ $\text{V}_7(\text{Bz})_8$  vdW junctions are designed and the transport properties of their four-terminal devices are comparatively investigated based on the DFT and NEGF techniques. The  $\text{MoS}_2$  and graphene nanoribbons act as the source-to-drain channel and the spin-polarized  $\text{V}_7(\text{Bz})_8$  nanowire serves as the gate channel. The transport characteristic is explored by investigating the conductance, currents, LDOS, and transmission spectra. Gate voltages applied on  $\text{V}_7(\text{Bz})_8$  exert different influences of electron transporting on  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$  and graphene/ $\text{V}_7(\text{Bz})_8$ . The interplay of source and gate bias potentials generates a pronounced influence on the transport property of  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$ . Evident NDR behavior, input/output current switches, as well as spin-polarized currents are found for  $\text{MoS}_2/\text{V}_7(\text{Bz})_8$ . In contrast, the gate bias plays insignificant effect on the transporting along graphene. These results are promising in designing multi-terminal nanoelectronic devices.

## Conflicts of interest

There are no conflicts of interest to declare.

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