


 Cite this: *RSC Adv.*, 2022, 12, 13975

Transition metal decorated phthalocyanine as a potential host material for lithium polysulfides: a first-principles study†

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The shuttle effect caused by the soluble long-chain lithium polysulfides greatly hinders the practical application of lithium–sulfur (Li–S) batteries. Therefore, the introduction of suitable anchoring materials is more effective to mitigate this problem. Transition metal phthalocyanines (TMPC) are regarded as a new class of sulfur host materials. Here, 4d transition metal (Y, Zr, Nb, Mo, Tc, Ru, Rh, Pd, Ag, Cd) decorated phthalocyanines are designed and systematically researched for the performance analysis of anchoring S₈/LiPSs by first-principles calculations. The results reveal that the bonding strength of LiPSs can be well adjusted by introducing suitable 4d transition metals into the phthalocyanine structure. The electronic structure analysis indicates the formation of TM–S bonds between the TMPC substrate materials and the LiPSs, which is essential to weaken the Li–S bonds and hence slow down the shuttle effect of LiPSs. ZrPC and NbPC both exhibit excellent potential and thermal stability for facilitating the conversion of LiPSs, as well as a better promoting effect for the sulfur reduction reactions (SRR) with a reduced Gibbs free energy in the rate-determining step (*Li₂S₂ → *Li₂S) during the discharge reaction process. These findings in our work may encourage further experimental and theoretical research for anchoring LiPSs with TMPC as a host material.

Received 30th March 2022

Accepted 29th April 2022

DOI: 10.1039/d2ra02049a

rsc.li/rsc-advances

1. Introduction

Diverse electrical energy storage systems are closely related to human society, such as rechargeable batteries with competitive costs, high specific energy power and long cycling life.¹ These batteries play a leading role in the fields of portable electronic devices, electric vehicles, smart power grids, *etc.* Among the various batteries, lithium–sulfur (Li–S) batteries have attracted intense attention due to their exceptionally high theoretical capacity (1675 mA h g^{−1}), energy density (2600 W h kg^{−1}) and environmental friendliness.^{2,3} Compared with Li-ion batteries, Li–S batteries are expected to become one of the most promising alternatives for next-generation energy storage systems. Despite their great promise, the development of Li–S batteries is still hindered by some problems such as low active materials utilization, poor electrode stability, poor cycle life, and low Coulomb efficiency. The root cause of these matters lies in the low ion/electronic conductivity of elemental sulfur and

discharge products Li₂S/Li₂S₂ in the charge/discharge process, as well as the soluble long-chain polysulfides (Li₂S₄, Li₂S₆ and Li₂S₈) dissolved in the electrolyte solvent from the positive electrode (known as the “shuttle effect”).^{4–6} Therefore, solving these issues is particularly important for the design and the commercialization of Li–S batteries.

In the early days of Li–S batteries exploration, various two-dimensional (2D) materials with large special surface area and multiple adsorption sites, such as carbon nanotubes (CNTs),⁷ graphene sheets⁸ have been extensively used as the possible sulfur hosts to improve the conductivity of cathode and physically confine the shuttle effect of LiPSs. Among them, carbon-based materials can not only serve as a porous framework for adsorbing LiPSs, but also greatly improve the electronic conductivity of electrode materials. Therefore, the electrochemical performance of Li–S batteries has indeed been improved to a certain extent.⁹ However, until now, researchers have generally believed that the weak interaction between carbon-based materials and LiPSs is not sufficiently enough to prevent LiPSs from being transported into the electrolyte under high concentration gradients and internal electric fields.¹⁰ This phenomenon reveals that carbon-based materials alone cannot be used as a perfect matrix for promoting the overall performance of Li–S batteries. As an alternative, it has been proposed to anchor the LiPSs through relatively strong chemical interactions to enhance the capture capability¹¹ and strengthen the binding ability between the substrate materials and LiPSs

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 † Electronic supplementary information (ESI) available. See <https://doi.org/10.1039/d2ra02049a>

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species.¹² This method can effectively inhibit the dissolution of LiPSs into the electrolyte and is essential for improving the electrochemical performance of Li-S batteries. In the past few years, researchers have discovered that anchoring a single transition metal on a substrate material is an effective solution pattern, which can achieve a strong anchoring and catalytic effect on the basis of the Li anode and the S cathode.¹³ Moreover, it has been found that the interaction between substrate materials and adsorbents is mainly due to the transfer of charge between Li/S atoms in LiPSs and the substrate materials. For instance, Wu *et al.*¹⁴ investigated the InP₃ monolayer as a viable substrate material for anchoring LiPSs. Their research showed that during the adsorption process of LiPSs, electron accumulation mainly occurs between Li and P atoms, and electron depletion mainly occurs around S atoms. Zheng *et al.*¹⁵ discussed the use of non-metallic monolayer materials as anchoring materials in Li-S batteries for the adsorption of LiPSs. The remarkable performance of C₃N₄/C₂N is attributed to the bonding of Li-N/C-S and charge transfer. It can be seen that substrate materials with good physical and chemical properties are extremely important for the research of Li-S batteries.

Recently, materials based on phthalocyanine structure have received widespread discussion and research in the field of energy storage. Phthalocyanine (Pc) is a macrocyclic aromatic compound with a chemical formula of C₃₂H₁₆N₈.¹⁶ It has four isoindole subunits connected by a methine or azamethine bridge. The characteristics of this structure make it having excellent physical and chemical properties, flexible structure adjustability. Benefitting from the central cavity of tetranitrogen coordination, it can provide a perfect position to anchor a single transition metal atom.¹⁷ Therefore, it can form transition metal phthalocyanine (TMPC) structures with the doping of various transition metal ions, making them natural single-atom catalysts, which are widely used in various catalytic reactions, such as oxygen reduction reaction (ORR),¹⁸ nitrogen reduction reaction (NRR)¹⁹ and carbon dioxide reduction reaction (CO₂RR).²⁰ At present, some investigations have successfully synthesized TMPC materials with excellent conductivity, which can be directly applied to electrochemical catalysis reactions and can meet the needs of the reaction by flexibly adjusting the coordination of transition metal atoms. Cheviri *et al.*²¹ researched the adsorption of LiPSs on the OH-functionalized CoPc. Attributed to the existence of the Li-O bonds and the Li-N bonds, the adsorption of LiPSs was significantly stronger than that of carbon-based materials, indicating that CoPc is a good substrate material. Kim *et al.*²² adopted cobalt(II)-centered fluorinated phthalocyanine (F-Co(II)Pc) as the positive electrode of a Li-S battery, which effectively promoted the conversion reaction between LiPSs due to the existence of catalytic site on Co(II) active center. In addition, during the charging process, the strong combination of Co(II) and S atoms in Li₂S weakens the Li-S bonds, resulting in the lowest dissociation activation energy of Li₂S on Co(II)Pc. By means of coupling FePc and rGO, Ma *et al.*²³ designed a bi-functional anchored conversion composite material. This material can effectively capture the dissolved LiPSs, and the interconnected rGO substrate can also facilitate the transfer of

ions and electrons during the reaction, thus speeding up the redox reaction of the lithium-sulfur battery. Yu *et al.*²⁴ investigated the catalytic properties of different 3d transition metals anchored on the phthalocyanine structure and forecasted the potential application of TMPC as a battery cathode material in Li-S batteries. Besides, it is reported that the cavity of H₂Pc can be occupied by about 70 different metal ions with excellent chemical stability.^{25,26} In view of the large family of TMPC members and the transition metal atoms, a systematical investigation of the promising transition metal phthalocyanine materials is necessary for the development and application of TMPC-based host materials for LiPSs.

Inspired by the work of predecessors,²⁷⁻²⁹ herein, we systematically investigated the anchoring capability and catalytic performance of all 4d transition metal (TM = Y, Zr, Nb, Mo, Tc, Ru, Rh, Pd, Ag, Cd) phthalocyanines (TMPC) on a series of polysulfides S₈/LiPSs. By first-principles calculations, a comprehensive picture of the structures and electronic properties of TMPC-S₈/LiPSs complexes were provided. Overall, this work indicates that the binding ability of S₈/LiPSs can be well regulated by the appropriate 4d TMPC structures, and Zr/NbPc exhibits the excellent potential (the binding energy of ZrPc and NbPc is 5.18 eV and 4.63 eV) for inhibiting the shuttle effect in Li-S batteries. In addition, the results of *ab initio* molecular dynamics (AIMD) simulations show that most of the TMPC structures provide a theoretical basis for the feasibility of the experiment. Our simulations not only are helpful for the designation of TMPC-based hosting materials for LiPSs, but also can attract further experimental research.

2. Computational methods

2.1 Computational details

All spin-polarized density functional theory (DFT) calculations for all TMPC structures were performed by the Vienna *Ab initio* Simulation Package (VASP) code.^{30,31} The projector augmented wave (PAW) method was employed for describing the ion-electron interaction and the exchange correlation functional within the generalized gradient approximation (GGA) parameterized by Perdew, Burke, and Ernzerof (PBE) was used in all DFT calculations.^{32,33} In addition, the DFT-D3³⁴ correction with Grimme scheme was applied to describe van der Waals interaction between the adsorbates and adsorbents. The cutoff energy was set to 520 eV for the plane-wave basis. The convergence standards for the energy and the residual force were set to 10⁻⁵ eV and 0.02 eV Å⁻¹, respectively. The Brillouin zones were sampled with a 2 × 2 × 1 Monkhorst-Pack *k*-point grid for the H₂Pc and all the TMPC structures. The crystal orbital Hamilton populations (COHP)^{35,36} implemented in LOBSTER³⁷ was applied to describe the bonding/anti-bonding state between atoms.

2.2 Binding energy

The binding strength between the S₈/LiPSs and the H₂Pc/TMPC substrates was evaluated by the calculation of the binding energy (*E_b*) according to the following equation:^{38,39}



$$E_b = (E_{\text{substrate}} + E_{\text{adsorbate}}) - E_{\text{total}} \quad (1)$$

where E_{total} , $E_{\text{substrate}}$, and $E_{\text{adsorbate}}$ are the total energies of H₂Pc/TMPC-S₈/LiPSs complexes, H₂Pc/TMPC, and S₈/Li₂S_x ($x = 1, 2, 4, 6, 8$), respectively. The more positive the value of E_b , the stronger the interaction between the corresponding substrates and the adsorbates.

2.3 Charge density difference

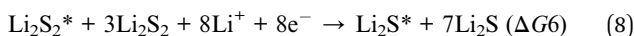
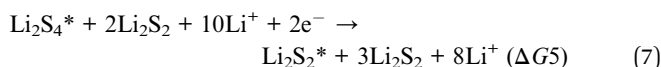
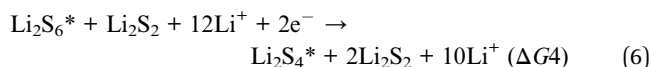
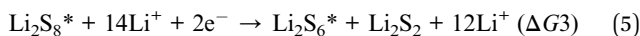
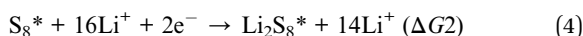
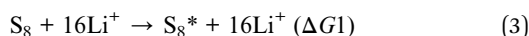
The charge density difference was defined as follows:⁴⁰

$$\Delta\rho = \rho_{\text{total}} - \rho_{\text{substrate}} - \rho_{\text{adsorbate}} \quad (2)$$

wherein ρ_{total} , $\rho_{\text{substrate}}$ and $\rho_{\text{adsorbate}}$ stand for the charge densities of H₂Pc/TMPC-S₈/LiPSs complexes, H₂Pc/TMPC, and S₈/LiPSs, respectively.

2.4 Free energy change of elementary reaction

The following eqn (3)–(8) described the overall discharge reactions or sulfur reduction reactions (SRRs) process in detail, where S₈/LiPSs molecules with “*” marked indicate that they were adsorbed on the TMPC/H₂Pc substrate materials.²⁴ The energy of a single Li ion and an electron (Li⁺ + e⁻) pair was treated as the energy of a crystalline Li atom.



The Gibbs free energy (ΔG) of the above SRRs process was calculated as follows:⁴¹

$$\Delta G = \Delta E_{\text{DFT}} + \Delta E_{\text{ZPE}} - T\Delta S - n_e U + \Delta G_{\text{pH}} \quad (9)$$

Here ΔE_{DFT} is the electronic energy difference before and after adsorption of reaction intermediates, ΔE_{ZPE} is the zero-point energy difference between the LiPSs and the adsorbed state, T is the room temperature (298.15 K), ΔS is the entropy change, n and U are the number of electrons transferred and the effect of an applied bias. ΔG_{pH} is the correction of pH, which can be obtained *via* $\Delta G_{\text{pH}} = 2.303 \times k_{\text{B}} T \times \text{pH}$, where k_{B} is the Boltzmann constant and the calculations presented here are for pH = 0. For the battery systems, the ΔE_{ZPE} and $T\Delta S$ can be ignored in the calculation, therefore, when $U = 0$ V, $\Delta G \approx \Delta E_{\text{DFT}}$. In order to gain further insight into the catalytic performance of the charging reaction (reverse reaction of SRRs), we carried out the transition state (TS) calculations for some structures. The climbing image nudged elastic band (CI-NEB)

method⁴² is used to calculate the dissociation energy barrier of Li₂S, and the cut-off energy, SCF convergence and max force are set to 520 eV, 1.0×10^{-7} eV per atom and 0.03 eV Å⁻¹, respectively.

2.5 *Ab initio* molecular dynamic (AIMD) simulation

Since the structure optimization process of the VASP calculation program is an optimization method using interatomic statics under standard 0 K conditions, it is not sufficient to understand the chemical reaction stability of the 4d TMPC structures in Li-S batteries. While density functional (DFT) *ab initio* molecular dynamics (AIMD) can reliably describe the time-evolution of solid-state systems directly embedded in real-world environments,⁴³ its computational strategy for simulated annealing is an upgrade from the standard 0 K DFT basic computation. Therefore, in order to check the thermodynamic stability of all TMPC structures, calculations of AIMD simulations were implemented. The optimized TMPC structures were annealed with 1000 steps to increase the temperature to the target temperature of 400 K (time step is 1 fs, temperature from 100 K to 400 K), and then 4 ps equilibrium simulation calculation was performed at 400 K with a time step of 1 fs. Temperature control is achieved through the Nosé thermostat model.

3. Results and discussions

3.1 Structural stability analysis

In this study, all 4d transition metal phthalocyanine (TMPC) structures were considered for the performance analysis for anchoring S₈/LiPSs. The optimized structures of H₂Pc/TMPC (TM = Y, Zr, Nb, Mo, Tc, Ru, Rh, Pd, Ag, Cd) and S₈/LiPSs are displayed in Fig. 1a and b. As illustrated in Fig. 1a, the non-metal-doped phthalocyanine (H₂Pc) structure composed of four isoindole groups and an inner porphyrine ring is a large conjugated system with 18 electrons and is a completely artificially synthesized compound. Theoretically, all transition metal (TM) ions can be used to replace the two hydrogen atoms situated in the center of H₂Pc structure, thereby forming a stable TMPC structure²⁵ (except AgPc with a positive formation energy and the details are described in Table S1†) for better anchoring S₈/LiPSs. Most of the optimized TMPC structures (except YPc, ZrPc, CdPc) can well maintain the initial planar structure of H₂Pc. The large radius of metal atoms Y, Zr and Cd result in a non-planar configuration and the introduced metal atoms are above the macrocyclic plane. Fig. 1b exhibits the fully relaxed S₈/Li₂S_x ($x = 1, 2, 4, 6, 8$) structures, which are consistent with the previous investigations.^{44,45}

3.2 Analysis of anchoring performance of S₈/LiPSs adsorbed on H₂Pc/TMPC substrates

Li-S battery is a rechargeable battery with elemental sulfur as the positive electrode and metal lithium as the negative electrode. Because of its wide range of material sources, high energy density, large theoretical capacity and so on, it has received extensive attention and research. However, the biggest technical obstacle of Li-S battery studies lies in the “shuttle effect” of



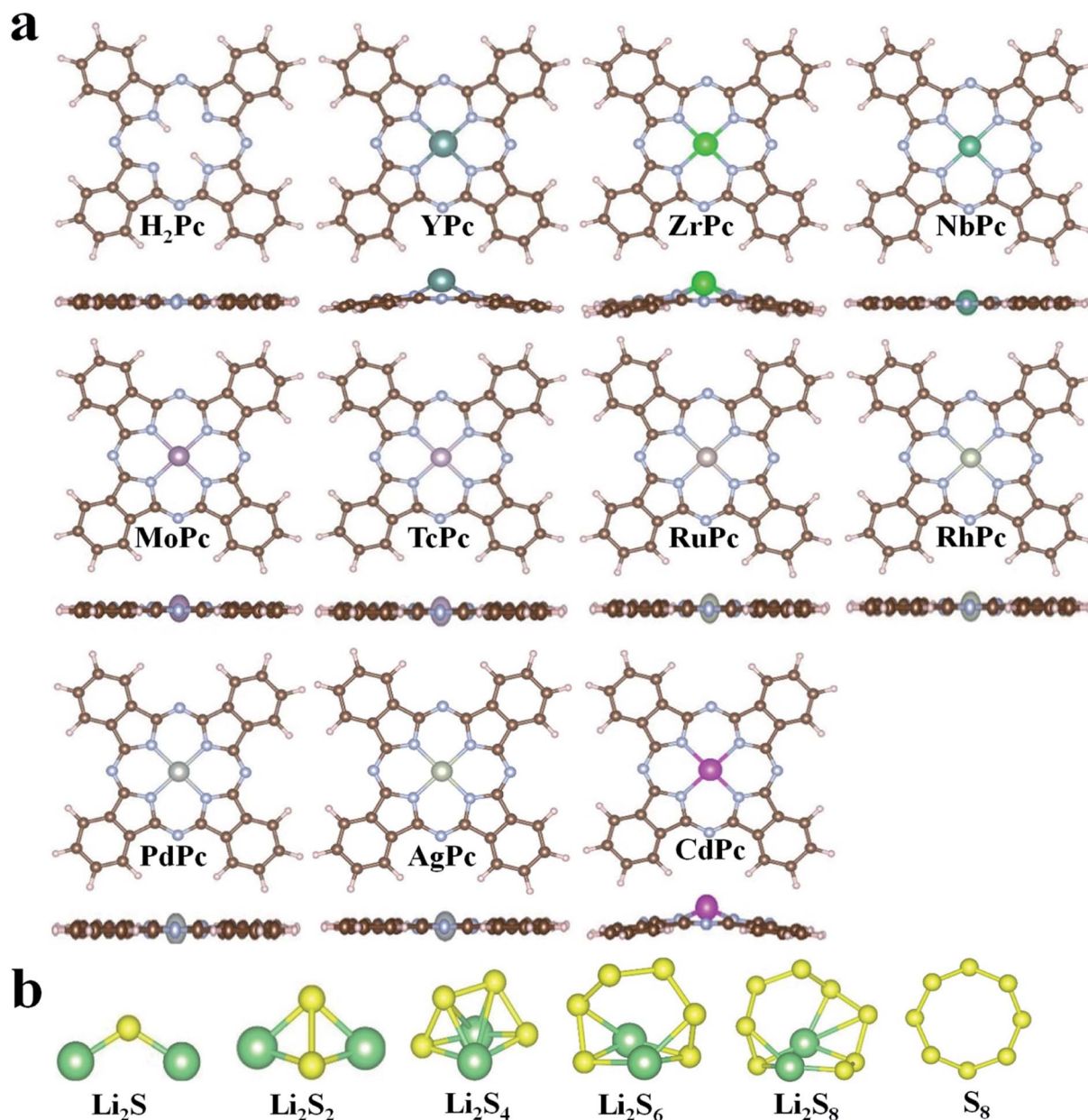


Fig. 1 (a) The schematic diagram of the optimized H_2Pc^{24} and TMPc structures. (b) The fully relaxed Li_2S , Li_2S_2 , Li_2S_4 , Li_2S_6 , Li_2S_8 and S_8 molecules. The N, C, H, S and Li are represented by silver, brown, lavender blush, yellow and green balls, respectively.

LiPSs⁴⁶ (especially Li_2S_4 , Li_2S_6 and Li_2S_8), which simultaneously reduces the performance of the cathode and anode. In the research of this article, we fully investigated the interaction between various H_2Pc /TMPc substrates with S_8 /LiPSs clusters. The binding energy results of H_2Pc /TMPc adsorbing S_8 /LiPSs are summarized and displayed in Fig. 2a. It can be intuitively seen that most of TMPc structures, especially ZrPc and NbPc, exhibit much better anchoring performance for S_8 /LiPSs than the initial phthalocyanine structure (H_2Pc), revealing that the shuttle effect of LiPSs can be effectively suppressed by the transition-metal adjusted Pc structure. The binding strength between most of the soluble LiPSs (*e.g.*, Li_2S_4 , Li_2S_6 and Li_2S_8) and H_2Pc /TMPc substrates is weaker than that of insoluble

LiPSs (*e.g.*, Li_2S and Li_2S_2), and the binding energy of S_8 is the minimum. The stable structures of S_8 /LiPSs species adsorbed on all the H_2Pc /TMPc substrate materials could be found in Fig. S1 and S2.† Fig. 2b is a schematic diagram of the optimized structure of Li_2S molecules adsorbed on the NbPc substrate. It can be seen that lithium-atom (Li_1/Li_2) are easier to bond with nitrogen-atom forming Li–N bonds, while TM (Nb) atoms tend to combine with sulfur-atom to form TM(Nb)–S bonds. The two newly formed bonds (Li–N and TM–S) are of vital importance in anchoring LiPSs and the interaction of TM–S bonds play a dominating role. In addition, the strength of the interaction between TMPc and Li_2S can also be reflected by the bond lengths of TM–S and Li–S bonds. The columns of different



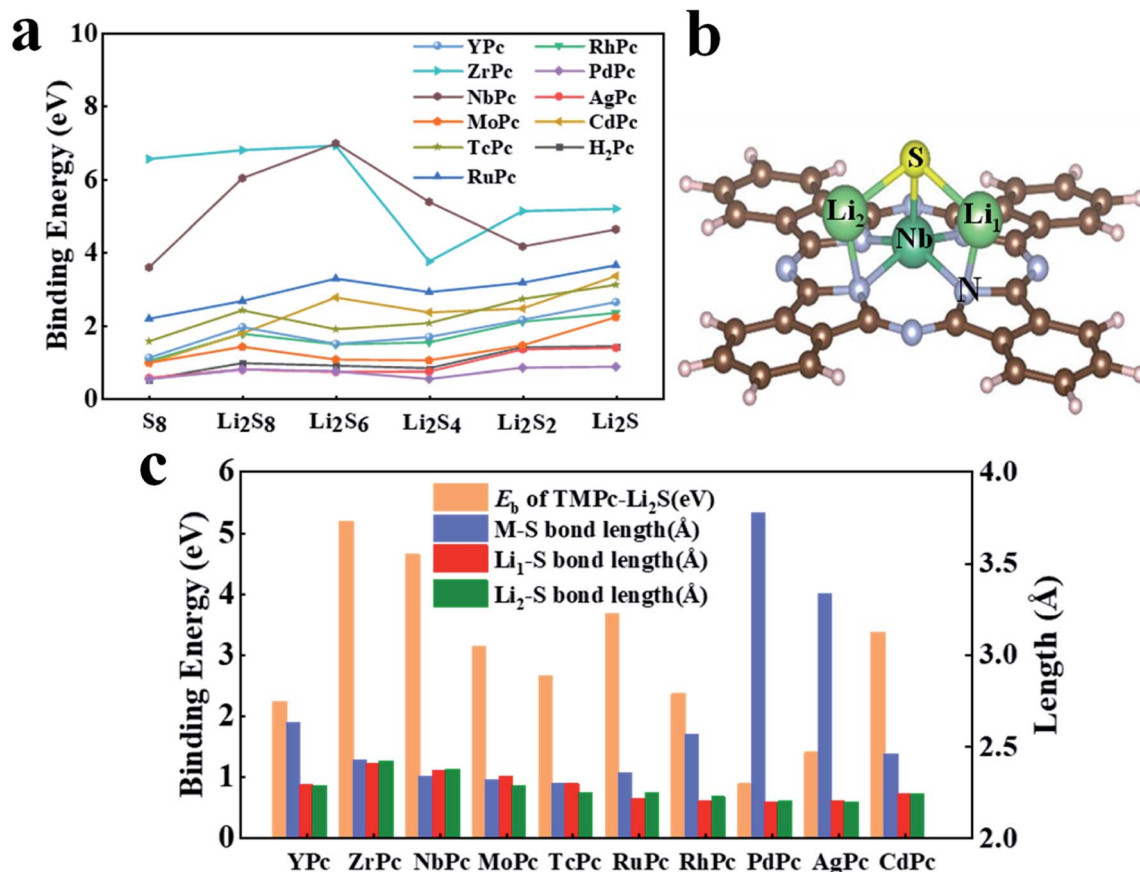


Fig. 2 (a) Binding energy of S₈/Li₂S_x (x = 1, 2, 4, 6, 8) adsorbed on different H₂Pc²⁴/TMPc surfaces; (b) the schematic diagram of NbPc-Li₂S binding system with atomic label (Li₁, Li₂, S, Nb and N atoms); (c) the binding energy, TM-S bond lengths and Li_{1/2}-S bond lengths histogram of different TMPc-Li₂S systems.

colors in Fig. 2c show the results of our statistics. The red, green, blue and orange columns represent the binding energy of TMPc adsorbing Li₂S, the bond lengths of TM-S, Li₁-S and Li₂-S (denoted as Li_{1/2}-S bonds), respectively. As shown in Fig. 2c, there is a correlation between the binding strength of TMPc adsorbing Li₂S and the TM-S bond lengths. The binding energies of PdPc, AgPc and YPc to Li₂S are relatively small among the TMPc structures, and the TM-S bond lengths after structure optimization is obviously larger than other structures. Among them, the Pd-S bond length of PdPc-Li₂S is the longest one (3.77 Å) and its binding energy is also the smallest (0.86 eV). Moreover, the Li_{1/2}-S bond lengths of all TMPc structures after adsorption of Li₂S molecules have varying degrees of elongation, and the longer Li_{1/2}-S bond corresponds to a stronger binding energy. Taking the ZrPc-Li₂S and PdPc-Li₂S systems as examples, the Li_{1/2}-S bond lengths of the Li₂S molecules before and after adsorption are both 2.09 Å, while the Li₁-S bond length in the two systems increased to 2.39 Å and 2.18 Å, respectively. Furthermore, the bond length of Li₂-S increased to 2.41 Å and 2.19 Å, and the corresponding binding energies are 5.18 eV and 0.86 eV. In short, in all the considered TMPc structures tuned by the 4d transition metal ions, ZrPc and NbPc have excellent anchoring properties for S₈/LiPSSs, which shows

that the two structures can effectively suppress the shuttle effect in Li-S batteries.

3.3 Electronic structure analysis of TM (Zr, Nb, Ru, Pd)Pc anchoring S₈/LiPSSs

In order to further understand the potential binding mechanism between S₈/LiPSSs and TMPc substrates, taking ZrPc, NbPc, RuPc and PdPc systems as examples, we studied the charge density difference before and after adsorption. As shown in Fig. 3, the electron accumulation is displayed in the yellow area, while the electron consumption is displayed in the cyan area, respectively. The charge density difference diagrams clearly exhibit the differentiation in charge transfer between each structure. It is worth noting that the charge transfer between PdPc and S₈/LiPSSs is significantly less than that between Zr(Nb, Ru)Pc and S₈/LiPSSs, which is highly consistent with the results of binding energy. Through comparative analysis, the charge transfer of ZrPc-Li₂S, NbPc-Li₂S and RuPc-Li₂S systems mainly occurs in the TM(Zr, Nb, Ru)-S bonds and Li-N bonds. However, for the PdPc-Li₂S system, the charge mainly accumulates between Li atom and PdPc substrate rather than the Pd/S atoms (Fig. 3d). So far, the above results show that Zr-S, Nb-S and Ru-S bonds are much stronger than the Pd-S bond.



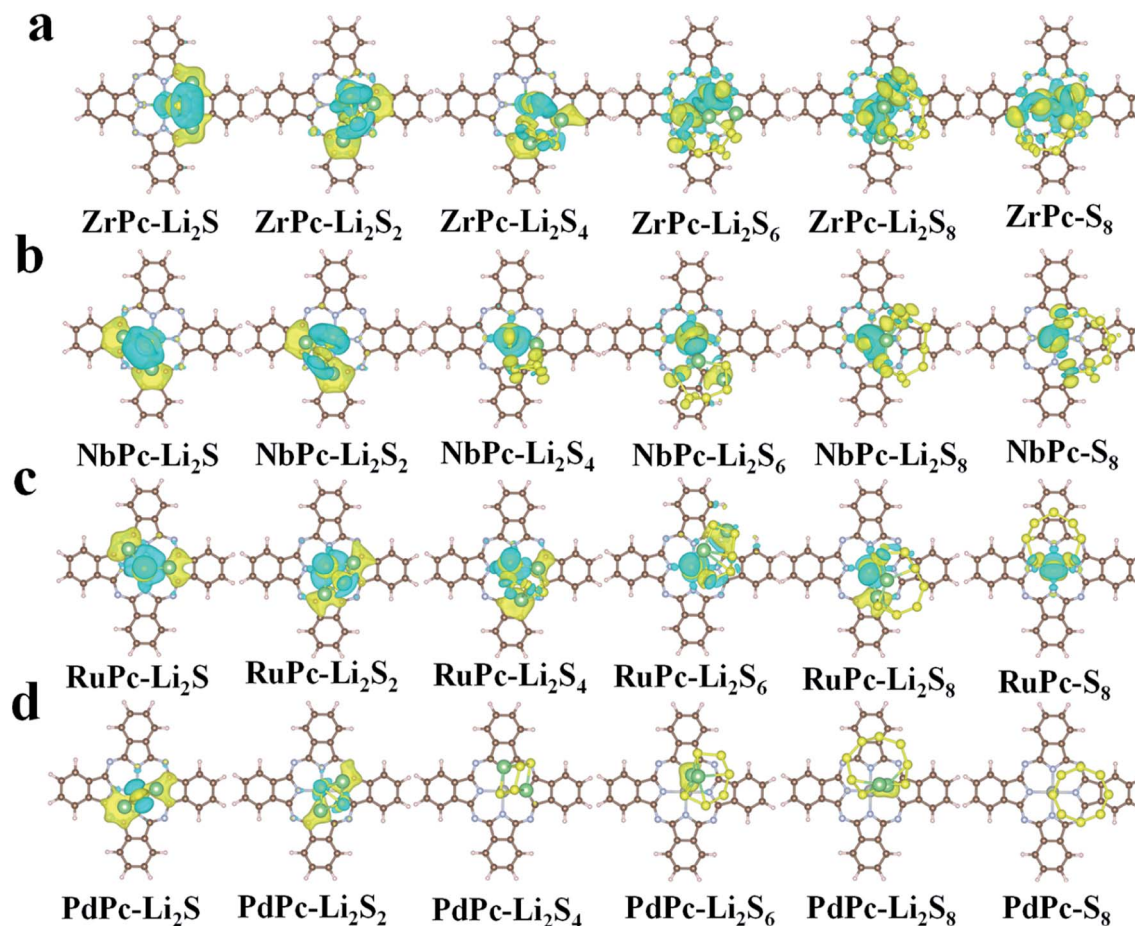


Fig. 3 Charge density differences of S_8 /LiPSs adsorbed on (a) ZrPc, (b) PdPc, (c) RuPc and (d) PdPc with an iso-surface value set to $0.002 e \text{ \AA}^{-3}$. The yellow area and cyan area in the figure indicate charge accumulation and depletion, respectively.

In addition, the binding energy of PdPc– Li_2S is similar to that of H_2Pc-Li_2S (Fig. 2a), indicating that the interaction of lithium atoms and nitrogen atoms is not sufficient to anchor S_8 /LiPSs clusters. In general, a strong TM–S bond is essential to alleviate the shuttle effect of LiPSs.

To reveal the mechanisms of TMPc substrates anchoring S_8 /LiPSs, we calculated the projected density of states (PDOS) of different TM (Zr, Nb, Ru and Pd)Pc before and after the adsorption of Li_2S and the corresponding energy of the d-band center (ϵ_d), which are presented in Fig. 4. Focusing on the TM–S and Li_1-N (Fig. 2b) bonds, the PDOS of the 4d orbitals of Zr, Nb, Ru and Pd (Zr_d, Nb_d, Ru_d and Pd_d) before adsorption and the 2p orbitals of N-atom (N_p) in the Zr(Nb, Ru, Pd)Pc structures are drawn. As shown at the bottom of the PDOS diagrams in Fig. 4a, ZrPc, NbPc and RuPc show a large overlap between the d orbital and p orbital in a wide energy range, which means a stronger bonding between Zr(Nb, Ru) and N atoms. The middle part in Fig. 4a show that the adsorption of Li_2S induces a strong hybridization of the 2p orbital (S_p) and the d orbital of Zr/Nb, revealing a stronger interaction between Zr–S and Nb–S. Meanwhile, the PDOSs of S_p and Zr/Nb_d split into two parts, that is the bonding state and the anti-bonding state, and are distributed on both sides of the Fermi level. Compared with

PdPc and Li_2S alone, the PDOS changes of Pd_d and S_p in PdPc– Li_2S system are very small, indicating a much weaker interaction between Pd and S atom. The analysis from the PDOS plots reveals that interactions between the TM_d and S_p orbitals engender the strong TM–S interactions, such as the Zr–S and Nb–S bonds. In addition, the energy of the d-band center can be approximately employed to describe the adsorption energy of the adsorbate. And the closer the value of the d-band center is to the Fermi level, the stronger the adsorption effect on S_8 /LiPSs.⁴⁷ Therefore, we calculated the energy of the d-band center in the ZrPc– Li_2S , NbPc– Li_2S , RuPc– Li_2S , and PdPc– Li_2S systems. As shown by the purple solid line in Fig. 4, the corresponding energy values of the d-band center (ϵ_d) are 0.33 eV (Zr_d), -0.61 eV (Nb_d), -1.22 eV (Ru_d) and -3.51 eV (Pd_d), respectively. Obviously, the energy of the d-band center of the ZrPc– Li_2S and NbPc– Li_2S systems is relatively closer to the Fermi level, indicating that the effect of Zr/NbPc anchoring LiPSs is indeed much stronger than other TMPc structures.

The COHP analysis is applied to explain the change of TM–S bonding strength. Hence COHP analysis of the TM–S bonds for ZrPc– Li_2S , NbPc– Li_2S , RuPc– Li_2S and PdPc– Li_2S systems are calculated and plotted in Fig. 4b. The Zr/Nb–S bonds have more electronic states near the Fermi level than the Ru/Pd–S bonds,



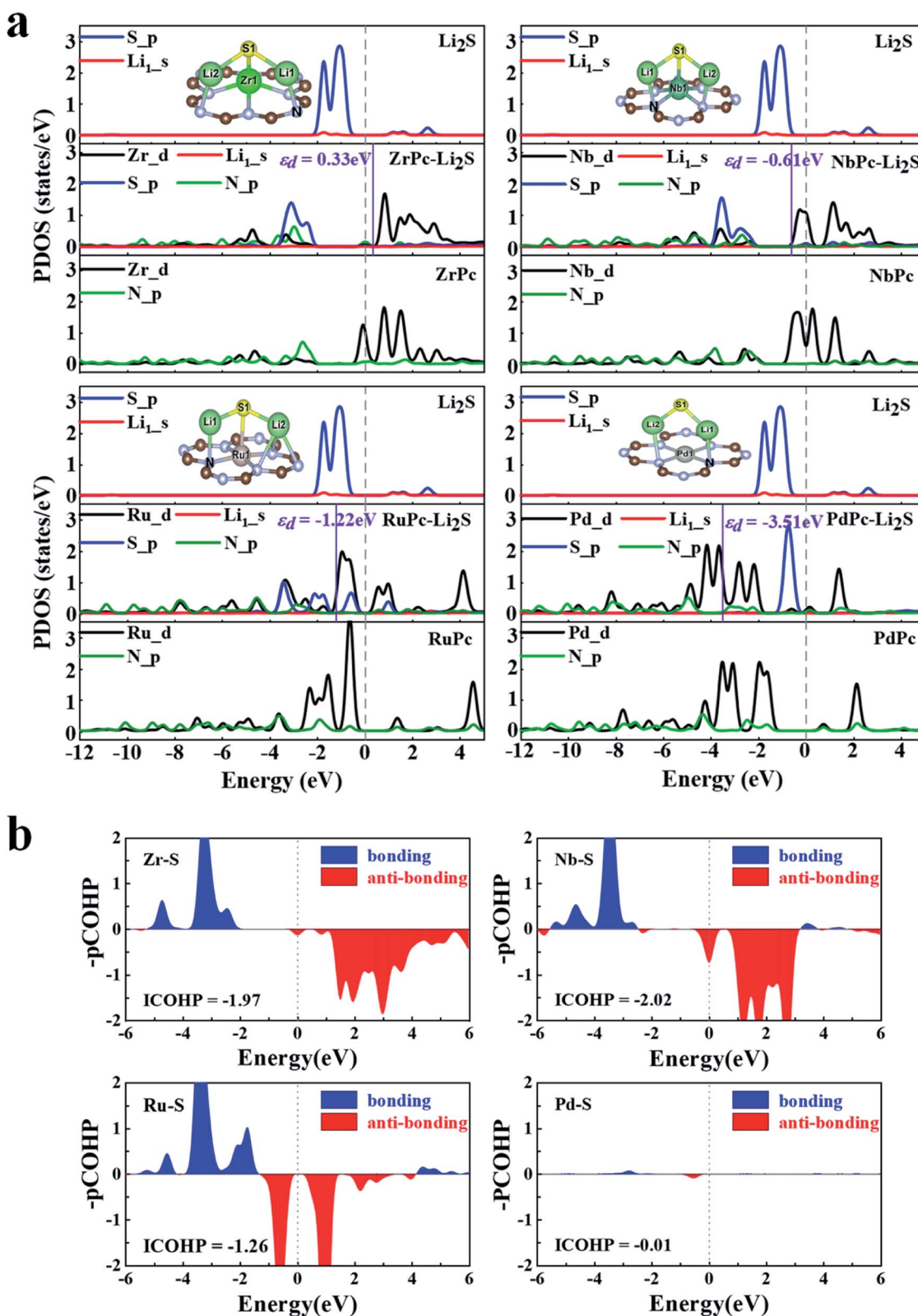


Fig. 4 (a) The partial density of states (PDOS) plots and d-band central energies (ϵ_d) for ZrPc-Li₂S, NbPc-Li₂S, RuPc-Li₂S and PdPc-Li₂S binding systems, where vertical dashed and solid lines represent the Fermi level and the energy of the d-band center, respectively. (b) The projected crystal Hamiltonian population (pCOHP) and corresponding integrated value (ICOHP) of Zr-S bond in ZrPc-Li₂S, Nb-S bond in NbPc-Li₂S, Ru-S bond in RuPc-Li₂S and Pd-S bond in PdPc-Li₂S.



as shown in Fig. 4b. The above findings indicate that the Zr/Nb–S bonds in Zr/NbPc–Li₂S systems have a much stronger interaction than the Ru/Pd–S bonds of Ru/PdPc–Li₂S systems. In order to obtain a quantitative and in-depth understanding of the interaction between TM–S bonds, the integral value of COHP (ICOHP) is further calculated. A more negative value of the ICOHP indicates that the stronger interaction of the TM–S bonds. The ICOHP results of different TM–S bonds in the four systems are –1.97 (Zr–S), –2.02 (Nb–S), –1.26 (Ru–S) and –0.01 (Pd–S), which is in good agreement with the trend of their binding energies. Among them, the ICOHP values of Zr/Nb–S are very close, which further verifies the conclusion that both ZrPc and NbPc all have excellent anchoring effect on S₈/LiPSSs.

3.4 Catalytic performance analysis

Through the previous analysis of H₂Pc/TMPc structures anchoring S₈/LiPSSs, it can be found that ZrPc and NbPc can well inhibit the shuttle effect, and are promising candidates for SRRs electrocatalysts. In addition, the rapid conversion of LiPSSs during charging and discharging of lithium–sulfur batteries is also an important factor affecting its overall performance. The discharge process is mainly the conversion of soluble Li₂S₄, Li₂S₆ and Li₂S₈ to insoluble or insulating Li₂S/Li₂S₂, making it difficult for the reaction to proceed quickly and sustainably. The charging process is a reversible reaction of the discharging process and the first reaction is the dissociation of Li₂S (Li₂S → LiS + Li⁺ + e[–]), which is the key step to produce Li–S batteries with high capacity and high rate. In view of this, the free energy diagram in the discharge reaction process (S₈* → Li₂S*) and the

dissociation energy barrier of Li₂S* in the charging reaction process are further simulated and analyzed for estimating catalytic performance of TMPcs.

The entire charge and discharge reaction process involve the transfer of 16 electrons. The SRRs from S₈* to Li₂S* are divided into six basic reactions (eqn (3)–(8)). In the discharge process, S₈* is first converted to Li₂S₈*, and then gradually reduced to Li₂S₆*, Li₂S₄*, Li₂S₂* and Li₂S*.⁴⁸ Fig. 5a presents the free energy diagrams of ZrPc, NbPc, RuPc, PdPc and H₂Pc during SRRs. Obviously, the conversion from S₈* to Li₂S₈* has the most negative ΔG, demonstrating that this reaction is the easiest to occur. Only the last step (Li₂S₂* is reduced to Li₂S*) has the most positive Gibbs free energy (ΔG), indicating that the rate-determining step of the entire SRRs process is the last step. Among them, the ΔGs of ZrPc, NbPc, RuPc, PdPc and H₂Pc for the rate-determining step are 3.19 eV, 2.94 eV, 3.32 eV, 3.83 eV and 3.86 eV, respectively. Our simulations indicate that ZrPc and NbPc have a smaller rate-limiting step ΔG, which are excellent to promote the kinetics of SRRs.

In addition, Fig. 5b exhibits the results of Li₂S dissociated on the surface of ZrPc and NbPc. The ZrPc/NbPc–Li₂S systems have a close (0.71 eV and 0.57 eV) and low Li₂S decomposition barrier (initial, transient and final state are plotted in the insets of Fig. 5b). On the contrary, the PdPc/H₂Pc–Li₂S systems cannot obtain the stable configuration after the dissociation of Li₂S* (as exhibited in Fig. S3†). The LiS and Li radicals in the dissociated state are fully optimized and then recombined to form Li₂S, indicating that they have almost no catalytic activity for the decomposition of Li₂S. It is worth mentioning that the

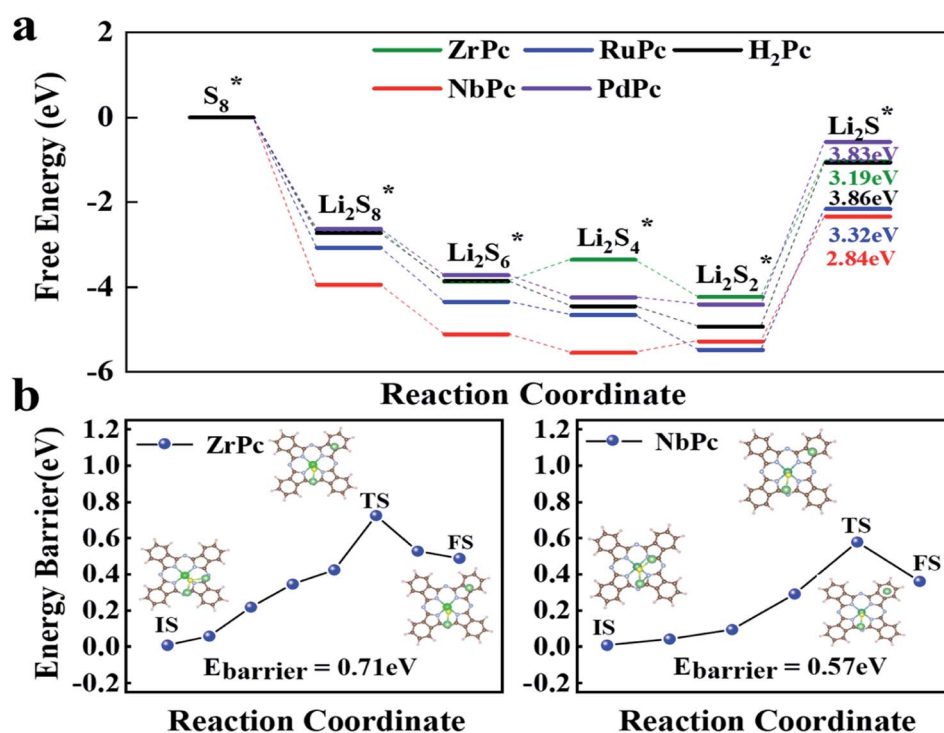


Fig. 5 (a) Relative free energy profiles for the discharging process from S₈ to Li₂S on the ZrPc, NbPc, RuPc, PdPc and H₂Pc surfaces. (b) The decomposition free energy barriers of the Li₂S on ZrPc and NbPc surfaces, and the corresponding initial, transition, and final structures are shown in the illustration.



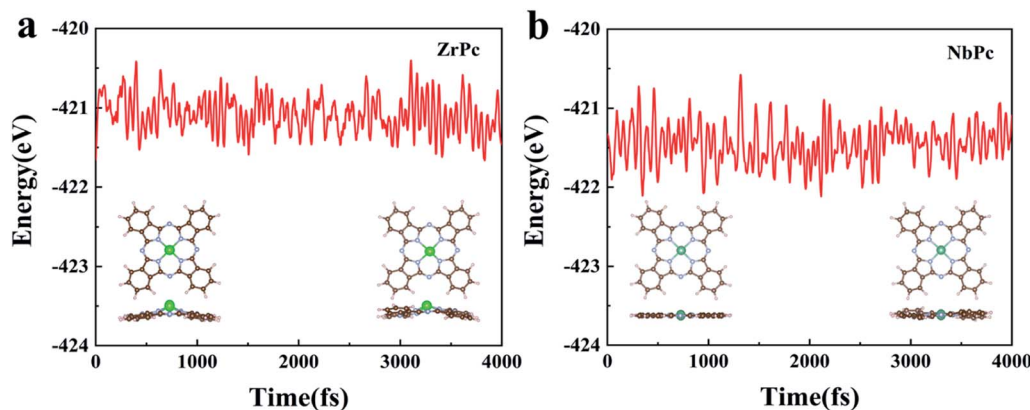


Fig. 6 (a) and (b) are the results of AIMD simulations of ZrPc and NbPc structures, respectively.

dissociation performance of the base material for Li_2S has a certain correlation with the strength of anchoring Li_2S^* . This is due to the fact that the Li–S bonds in the $\text{H}_2\text{Pc}/\text{TMPc}-\text{Li}_2\text{S}$ system are elongated and weakened to different degrees. The ZrPc/NbPc– Li_2S systems have a strong Zr/Nb–S bonds and hence possess a lower Li_2S^* dissociation energy barrier attributed to the strong interaction of the Zr/Nb–S bonds weakening the $\text{Li}_{1/2}\text{-S}$ bonds (according to Table S2,† the bond lengths of Li–S in pristine Li_2S , ZrPc– Li_2S , NbPc– Li_2S , PdPc– Li_2S and $\text{H}_2\text{Pc}-\text{Li}_2\text{S}$ are ~ 2.09 Å, ~ 2.41 Å, ~ 2.37 Å, ~ 2.19 Å, ~ 2.19 Å, respectively), so as to better promote the dissociation of Li_2S . Consequently, ZrPc and NbPc also exhibit excellent catalytic performance during the charging reaction.

For practical usage of TMPc catalysts, the thermodynamic stability of ZrPc and NbPc are evaluated using AIMD simulations at 400 K for 4 ps with 1 fs interval. The results of final structures and the total energy change are shown in Fig. 6. The AIMD simulation results of the remaining TMPc (TM = Y, Mo, Tc, Ru, Rh, Pd, Ag, Cd) structures are given in Fig. S4.† Most of TMPc structures are thermodynamically stable, which is basically consistent with the results of their formation energy. For ZrPc and NbPc, our results show that the total energy oscillate within a small range, and the geometric morphology are also preserved well without significant distortion and bond breaking, guaranteeing their relatively high stability. Based on this, it is of great interest that ZrPc and NbPc can serve as sulfur host materials in LiPSs with high experimental feasibility under different reaction environments.

4. Conclusions

In this work, density functional theory (DFT) calculations were used to explore the anchoring and catalytic performance of H_2Pc and TMPc-based electrocatalysts for adsorbing S_8/LiPSs clusters. The results ensure higher adsorption ability of TMPc (especially ZrPc and NbPc) for anchoring the S_8/LiPSs clusters than H_2Pc , as well as two bonding patterns (TM–S and Li–N bonding). Moreover, TM–S bonds are more vital to anchor LiPSs than Li–N bonds, due to the strong interaction between the 4d orbital of TM and the 2p orbital of S. The small decomposition

barriers and low rate-limiting step ΔG indicate that ZrPc and NbPc have excellent anchoring properties for S_8/LiPSs and can better inhibit the “shuttle effect” when used as the sulfur hosts in Li–S batteries. The AIMD simulation results also further predict the experimental feasibility of the TMPc substrate materials. Our research gave atomic insights for explaining the interaction between $\text{H}_2\text{Pc}/\text{TMPc}$ substrates and S_8/LiPSs molecules, and provided a feasible strategy for the rational design of better sulfur cathode hosts.

Conflicts of interest

The authors declare no conflicts of interest.

Acknowledgements

This work was financially supported by the National Natural Science Foundation of China (Grant No. 22168036), Department of Education of Tibet Autonomous Region ([2021]1-21 Theoretical research and design of two dimensional electrocatalytic materials).

Notes and references

- 1 M. Zhao, B. Q. Li, H. J. Peng, H. Yuan, J. Y. Wei and J. Q. Huang, *Angew. Chem., Int. Ed.*, 2020, **59**, 12636–12652.
- 2 J. X. Zhao, Y. G. Yang, R. S. Katiyar and Z. F. Chen, *J. Mater. Chem. A*, 2016, **4**, 6124–6130.
- 3 H. Yuan, H. J. Peng, J. Q. Huang and Q. Zhang, *Adv. Mater. Interfaces*, 2019, **6**, 1802046.
- 4 T. Li, X. Bai, U. Gulzar, Y. J. Bai, C. Capiglia, W. Deng, X. F. Zhou, Z. P. Liu, Z. F. Feng and R. P. Zaccaria, *Adv. Funct. Mater.*, 2019, **29**, 1901730.
- 5 Y. Hu, W. Chen, T. Y. Lei, Y. Jiao, J. W. Huang, A. J. Hu, C. H. Gong, C. Y. Yan, X. F. Wang and J. Xiong, *Adv. Energy Mater.*, 2020, **10**, 2000082.
- 6 M. Zhang, W. Chen, L. X. Xue, Y. Jiao, T. Y. Lei, J. W. Chu, J. W. Huang, C. H. Gong, C. Y. Yan, Y. C. Yan, Y. Hu, X. F. Wang and J. Xiong, *Adv. Energy Mater.*, 2020, **10**, 1903008.



- 7 A. A. Razzaq, Y. Z. Yao, R. Shah, P. W. Qi, L. X. Miao, M. Z. Chen, X. H. Zhao, Y. Peng and Z. Deng, *Energy Storage Mater.*, 2019, **16**, 194–202.
- 8 P. C. Shi, Y. Wang, X. Liang, Y. Sun, S. Cheng, C. H. Chen and H. F. Xiang, *ACS Sustainable Chem. Eng.*, 2018, **6**, 9661–9670.
- 9 Y. Yang, G. Y. Zheng and Y. Cui, *Chem. Soc. Rev.*, 2013, **42**, 3018–3032.
- 10 Z. Q. Ye, Y. Jiang, J. Qian, W. L. Li, T. Feng, L. Li, F. Wu and R. J. Chen, *Nano Energy*, 2019, **64**, 103965.
- 11 R. Jayan and M. M. Islam, *J. Phys. Chem. C*, 2020, **124**, 27323–27332.
- 12 J. Lei, T. Liu, J. J. Chen, M. S. Zheng, Q. Zhang, B. W. Mao and Q. F. Dong, *Chem*, 2020, **6**, 2533–2557.
- 13 J. F. Wu, S. Y. Ding, S. H. Ye and C. Lai, *J. Energy Chem.*, 2020, **42**, 27–33.
- 14 W. X. Wu, Y. M. Zhang, Y. H. Guo, J. X. Bai, C. H. Zhang, Z. F. Chen, Y. X. Liu and B. B. Xiao, *Appl. Surf. Sci.*, 2020, **526**, 146717.
- 15 Y. P. Zheng, H. H. Li, H. Y. Yuan, H. H. Fan, W. L. Li and J. P. Zhang, *Appl. Surf. Sci.*, 2018, **434**, 596–603.
- 16 D. Gounden, N. Nombona and W. E. van Zyl, *Coord. Chem. Rev.*, 2020, **420**, 213359.
- 17 I. E. Brumboiu, S. Haldar, J. Luder, O. Eriksson, H. C. Herper, B. Brena and B. Sanyal, *J. Chem. Theory Comput.*, 2016, **12**, 1772–1785.
- 18 Y. N. Zhou, G. P. Gao, W. Chu and L. W. Wang, *Nanoscale Adv.*, 2020, **2**, 710–716.
- 19 S. Q. Liu, Y. W. Liu, X. P. Gao, Y. J. Tan, Z. M. Shen and M. H. Fan, *Appl. Surf. Sci.*, 2020, **500**, 144032.
- 20 X. H. Wan, Z. F. Zhang, H. Niu, Y. H. Yin, C. G. Kuai, J. Wang, C. Shao and Y. Z. Guo, *J. Phys. Chem. Lett.*, 2021, **12**, 6111–6118.
- 21 M. Cheviri and S. Lakshmipathi, *Chem. Phys. Lett.*, 2020, **739**, 136942.
- 22 J. Kim, H. Shin, D. J. Yoo, S. Kang, S. Y. Chung, K. Char and J. W. Choi, *Adv. Funct. Mater.*, 2021, **11**, 2106679.
- 23 J. S. Ma, M. P. Yu, J. W. Zhu, W. T. Li, W. Gong and H. Qiu, *Ionics*, 2021, **27**, 3007–3016.
- 24 B. Yu, Q. He and Y. Zhao, *Appl. Surf. Sci.*, 2021, **558**, 149928.
- 25 M. S. Liao and S. Scheiner, *J. Chem. Phys.*, 2001, **114**, 9780–9791.
- 26 R. Zeis, T. Siegrist and C. Kloc, *Appl. Phys. Lett.*, 2005, **86**, 022103.
- 27 R. B. dos Santos, R. Rivelino, F. de Brito Mota, G. K. Gueorguiev and A. Kakanakova-Georgieva, *J. Phys. D: Appl. Phys.*, 2015, **48**, 295104.
- 28 A. Kakanakova-Georgieva, G. K. Gueorguiev, R. Yakimova and E. Janzén, *J. Appl. Phys.*, 2004, **96**, 5293–5297.
- 29 R. B. Dos Santos, R. Rivelino, B. Mota Fde, A. Kakanakova-Georgieva and G. K. Gueorguiev, *Dalton Trans.*, 2015, **44**, 3356–3366.
- 30 B. Hammer, L. B. Hansen and J. K. Norskov, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1999, **59**, 7413–7421.
- 31 J. Hafner, *J. Comput. Chem.*, 2008, **29**, 2044–2078.
- 32 Y. M. Juan, E. Kaxiras and R. G. Gordon, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1995, **51**, 9521–9525.
- 33 J. P. Perdew, K. Burke and M. Ernzerhof, *Phys. Rev. Lett.*, 1997, **78**, 1396.
- 34 S. Grimme, J. Antony, S. Ehrlich and H. Krieg, *J. Chem. Phys.*, 2010, **132**, 154104.
- 35 V. L. Deringer, A. L. Tchougreff and R. Dronskowski, *J. Phys. Chem. A*, 2011, **115**, 5461–5466.
- 36 R. Dronskowski and P. E. Blochl, *J. Phys. Chem.*, 1993, **97**, 8617–8624.
- 37 S. Maintz, V. L. Deringer, A. L. Tchougreff and R. Dronskowski, *J. Comput. Chem.*, 2016, **37**, 1030–1035.
- 38 Y. M. Zhao and J. X. Zhao, *Appl. Surf. Sci.*, 2017, **412**, 591–598.
- 39 L. H. Song, M. G. Zhang and J. Guo, *Chem. Phys. Lett.*, 2021, **777**, 138711.
- 40 Z. L. Yi, F. Y. Su, L. Huo, G. Y. Cui, C. L. Zhang, P. D. Han, N. Dong and C. M. Chen, *Appl. Surf. Sci.*, 2020, **503**, 144446.
- 41 Z. Y. Gao, H. Y. Huang, S. P. Xu, L. L. Li, G. Yan, M. L. Zhao, W. J. Yang and X. J. Zhao, *Mol. Catal.*, 2020, **493**, 111091.
- 42 G. Henkelman, B. P. Uberuaga and H. Jónsson, *J. Chem. Phys.*, 2000, **113**, 9901–9904.
- 43 D. G. Sangiovanni, G. K. Gueorguiev and A. Kakanakova-Georgieva, *Phys. Chem. Chem. Phys.*, 2018, **20**, 17751–17761.
- 44 L. J. Wang, T. R. Zhang, S. Q. Yang, F. Y. Cheng, J. Liang and J. Chen, *J. Energy Chem.*, 2013, **22**, 72–77.
- 45 Y. M. Zhao, L. Yang, J. X. Zhao, Q. H. Cai and P. Jin, *Phys. Chem. Chem. Phys.*, 2017, **19**, 18208–18216.
- 46 H. Ci, M. Wang, Z. Sun, C. Wei, J. Cai, C. Lu, G. Cui, Z. Liu and J. Sun, *J. Energy Chem.*, 2021, 474–482.
- 47 Y. Zhou, Z. Z. Zhou, R. X. Shen, R. G. Ma, Q. Liu, G. Z. Cao and J. C. Wang, *Energy Storage Mater.*, 2018, **13**, 189–198.
- 48 R. S. Assary, L. A. Curtiss and J. S. Moore, *J. Phys. Chem. C*, 2014, **118**, 11545–11558.

