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## Tuning the optical band gap and electrical properties of NiO thin films by nitrogen doping: a joint experimental and theoretical study

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A joint experimental and theoretical study is presented to reveal the influence of nitrogen doping on the optical and electrical properties of NiO thin films. Nitrogen addition can significantly enhance the subgap absorption. The molecular state of nitrogen (N<sub>2</sub>) has been identified in these doped thin films by electron energy loss spectroscopy.

Transition metal oxides (TMOs) have attracted intensive studies in the fields of photodetectors and energy conversion during the last two decades. Among the various TMOs, nickel oxide (NiO) is of special interest, due to its wide band gap, p-type conduction, good stability, environmental friendliness and low-cost fabrication. NiO has been regarded as an interesting p-type semiconductor for photodetectors when assembled with n-type materials to form heterojunctions.<sup>1–4</sup> In addition, NiO seems to be a promising hole transport layer in perovskite solar cells, as it has enabled great progress in the conversion efficiency from 7.8% to more than 20% in a few years.<sup>5,6</sup> However, due to the poor electrical properties of intrinsic NiO, doping by metallic elements (such as Li) is generally required.<sup>7,8</sup> Nanocrystalline NiO deposited on nickel nanoparticles can be used as a humidity sensor.<sup>9</sup>

Doping with nonmetal elements (such as N) is another way to tune the functional properties of TMOs. For example, it has been evidenced that the incorporation of N into TiO<sub>2</sub> can shift the optical absorption edge from the UV into the visible-light region, which is suited for photocatalysts.<sup>10–12</sup> However, the chemical states of N in TiO<sub>2</sub> remain debatable. Theoretical calculations suggest that nitrogen substituents on the oxygen sites (N<sub>O</sub>) in TiO<sub>2</sub> due to their close radii,<sup>10</sup> whereas a recent experimental study reveals that certain amount of nitrogen is also present as molecular N<sub>2</sub>.<sup>13</sup> In the case of Cu<sub>2</sub>O, first-principles calculations have shown that molecular nitrogen substitution for Cu, *i.e.*, (N<sub>2</sub>)<sub>Cu</sub>, has a lower formation energy than N<sub>O</sub>.<sup>14</sup> Such (N<sub>2</sub>)<sub>Cu</sub> defects may act as acceptors and increase

the hole carrier concentration significantly.<sup>15</sup> This result disagrees with previous assumption suggesting that nitrogen is in the form of anion substituting oxygen in the Cu<sub>2</sub>O lattice.<sup>16,17</sup> However, less attention has been paid to nitrogen doping effect in NiO so far, although NiO shows great potential applications. A solely computational work using the HSE06 functional predicts that N may locate on the O site with the form of N<sub>O</sub> in NiO, but comparison of the formation energy to those of other N-related defects is yet to be reported.<sup>18</sup> An experimental investigation shows that the band gap of NiO can be narrowed by the incorporation of N.<sup>19</sup> However, the physical mechanisms of the nitrogen doping in NiO remain unknown.

In this work, a joint experimental and theoretical study has been performed to develop a comprehensive picture on the nitrogen doped NiO thin films. Experimentally, the tunable optical and electrical properties of the NiO thin films by nitrogen incorporation have been investigated, and the molecular state of nitrogen in doped thin films has been identified. Computationally, various N-related defects in doped NiO have been systematically calculated, and compared with the experimental results. This study may provide a new path to tune the functional properties of TMOs by nitrogen doping.

N-doped NiO thin films were deposited on glass and Si (100) substrates by reactive magnetron sputtering in Ar–O<sub>2</sub>–N<sub>2</sub> atmospheres. During the deposition, no intentional heating was applied to the substrates, and the self-established deposition temperature was close to room temperature. A pulsed-DC power supply (Pinnacle + Advanced Energy) was connected to the metallic nickel target (2 in. diameter and 1 mm thick with a purity of 99.995%). The current applied to the target, the frequency and the off-time were fixed to 0.3 A, 50 kHz and 4 μs, respectively. The distance between substrate and target was 75 mm. Ar and O<sub>2</sub> flow rates were fixed at 25 and 3 sccm, respectively. Various N<sub>2</sub> flow rates (0, 10, 30 and 40 sccm) were introduced into the sputtering chamber. The total sputtering

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pressure was adjusted to 1 Pa. The deposition time was fixed to 6 min, and the film thickness was found to be in the range of 230–400 nm, depending on the  $N_2$  flow rate. The structure of thin films has been checked by X-ray diffraction (XRD, Bruker D8 Advance with Cu  $K_\alpha$  radiation in Bragg Brentano configuration). Electrical properties were determined by resistivity measurements (Toyo Corp. Resitest 8300) using a van der Pauw configuration. The microstructure was investigated by transmission electron microscopy (TEM, JEOL ARM 200-Cold FEG microscope equipped with a GIF Quantum ER system). The nitrogen chemical state in N-doped NiO thin films was studied by electron energy loss spectroscopy (EELS) embedded in TEM. EELS analyses have been performed at an acceleration voltage of the electron beam of 200 kV with the energy resolution of 0.40 eV defined by the full width at half maximum (FWHM) of the zero-loss peak. For this purpose, N-doped NiO thin films with the thickness of around 60 nm were deposited on carbon grid.

First-principles calculations were performed using the projector augmented-wave (PAW) method,<sup>20</sup> as implemented in the VASP code.<sup>21,22</sup> PAW datasets with radial cutoffs of 1.22, 0.79 and 0.80 Å were employed for Ni, N, and O, respectively. Ni 3d and 4s, N 2s and 2p, and O 2s and 2p orbitals were considered as valence electrons. The Perdew–Burke–Ernzerhof generalized gradient approximation was used for all the calculations. The Hubbard  $U$  correction according to the Dudarev's formalism was applied to the localized Ni 3d orbitals with the effective  $U$  parameter of 5.3 eV.<sup>18,23,24</sup> The plane wave cut-off energy was set to 520 eV for all the calculations. Magnetic configuration on  $Ni^{2+}$  sublattice was set to the type-II antiferromagnetic (AFM) ordering following the experimental findings.<sup>18</sup> The lattice constant was optimized with the 4-atom rhombohedral primitive cell at this magnetic structure.  $7 \times 7 \times 7$   $k$ -point mesh was used for the reciprocal space sampling. The calculated lattice constant is 4.19 Å, in good agreement with 4.18 Å of our synthesized undoped NiO. The band structure shows an indirect type with a 3.2 eV bandgap.

The defect calculations were performed with 512-atom supercells under the fixed theoretical lattice constant and  $\Gamma$ -point mesh for the reciprocal space sampling. The residual forces on atoms were reduced to be less than 0.005 and 0.03 eV Å<sup>-1</sup> for the perfect-crystal and defect calculations, respectively. Defect formation energies were calculated following ref. 25 and energy corrections for finite-size errors were adopted accordingly. In this study, the O-rich condition is considered and therefore the Ni, O and N chemical potentials ( $\mu_{Ni}$ ,  $\mu_O$  and  $\mu_N$ ) are set to satisfy  $2\mu_O = E_{O_2}$ ,  $\mu_{Ni} + \mu_O = E_{NiO}$ , and  $\mu_N + 2\mu_O = E_{NO_2}$ , where  $E_{O_2}$ ,  $E_{NiO}$  and  $E_{NO_2}$  are respectively the total energies of an  $O_2$  molecule, rocksalt NiO, a  $NO_2$  molecule per formula unit. The Fermi level is considered to locate between the calculated valence band maximum and the conduction band minimum in the perfect crystal.

X-ray diffractograms of N-doped NiO thin films deposited with various  $N_2$  flow rates (0, 10, 30 and 40 sccm) demonstrate their single-phase characteristic without metallic Ni or nickel nitride ( $Ni_3N$ ), as only NiO diffraction peaks are detected in all the thin films (Fig. 1). The diffraction peaks have been assigned

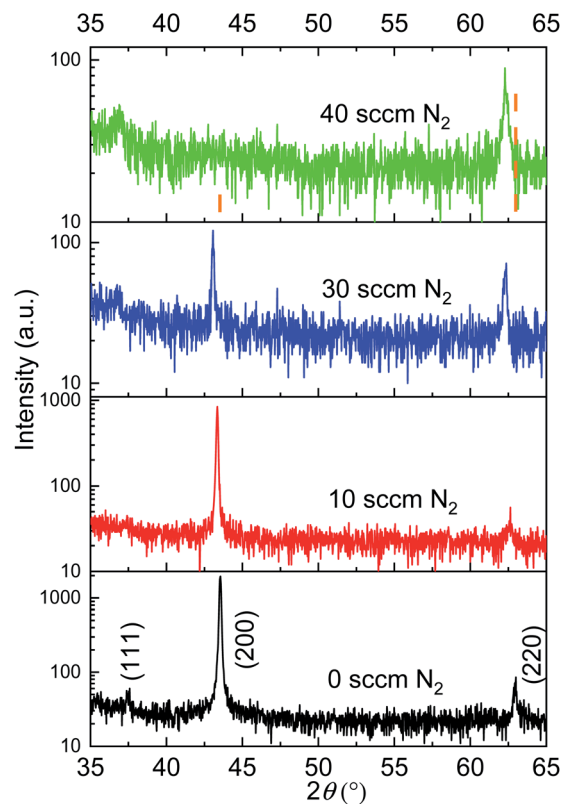


Fig. 1 X-ray diffractograms of nitrogen-doped NiO thin films.

according to the JCPDS 00-047-1049 file. As function of the nitrogen flow rate increase, the preferred orientation of NiO films changes from [200] to [220]. Although [111] and [200] are the most expected preferred orientation in rocksalt-like structures, the addition of nitrogen in the reactive may induce the growth of NiO films with [220] preferred orientation. Such preferred orientation has been already observed in DC sputtered NiO films.<sup>26</sup>

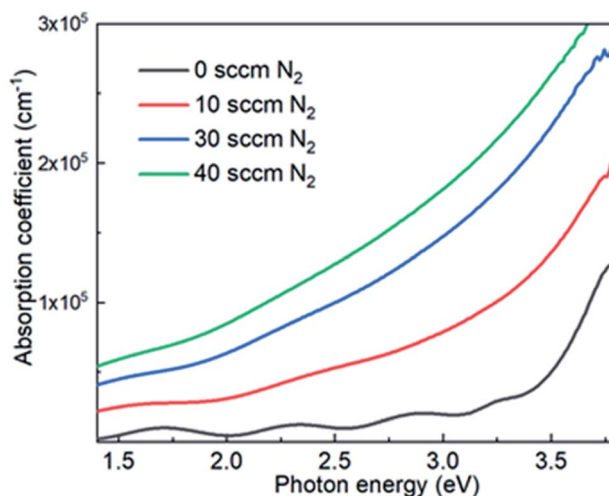


Fig. 2 Optical absorption coefficients of N-doped NiO thin films grown with various  $N_2$  flow rates (0, 10, 30 and 40 sccm).

The optical absorption coefficient ( $\alpha$ ) of N-doped NiO thin films grown with various  $N_2$  flow rates is shown in Fig. 2. It is seen that nitrogen doping in NiO can increase the optical absorption significantly in the photon energy between 1.4 and 3.8 eV. Such an enhancement of optical absorption by nitrogen doping is similar with that of nitrogen doped  $TiO_2$ .<sup>10,27</sup> Most of the published works based on transmittance measurements consider NiO as a direct semiconductor,<sup>8,19,28–30</sup> whereas theoretical calculations by various approaches including the present study predict an indirect band gap.<sup>31–33</sup> Furthermore, the literature does not mention the effect of nitrogen addition in NiO on the optical band gap value. Thus, the transmittance measurements have not been used to determine the optical band gap of N-doped NiO films using the well-known Tauc's method.

The electrical properties of N-doped NiO thin films at room temperature as a function of the  $N_2$  flow rate are plotted in Fig. 3. It demonstrates that the nitrogen incorporation can tune the electrical properties significantly. Taking the resistivity as an example, it is reduced from 69 200 (0 sccm  $N_2$ ) to 2268 (10 sccm  $N_2$ ), 273 (30 sccm  $N_2$ ) and 110  $\Omega$  cm (40 sccm  $N_2$ ). The mobility and carrier concentration have also been checked by AC Hall effect measurement. However, these two parameters of the undoped NiO (0 sccm  $N_2$ ) thin film could not be precisely determined probably due to its quite high resistivity indicating a probable low carrier concentration. In the case of other N-doped thin films (10, 30 and 40 sccm  $N_2$ ), the positive Hall coefficients imply the p-type conduction. As seen in Fig. 3, as increasing the  $N_2$  flow rate from 10 to 40 sccm, the mobility decreases from 3.4 to 0.14  $cm^2 V^{-1} s^{-1}$ . On the contrary, the hole concentration rises up from  $8.1 \times 10^{14}$  to  $4.0 \times 10^{17} cm^{-3}$ , indicating that nitrogen doping can increase the hole concentration remarkably. Since the increase of the hole concentration is higher than the mobility decrease, the addition of nitrogen induces a decrease of the film resistivity.

To reveal the mechanism behind such tunable optical and electrical properties by nitrogen incorporation, the nitrogen chemical states have been investigated by EELS, which is a useful technique to acquire the fruitful information of materials, such as chemical states, optical constants and density of electronic states.<sup>34–36</sup> The combination of EELS with TEM allows

the penetration depth of several tens to hundred nanometers (depending on the acceleration voltage), rather than the surface sensitivity in traditional photoelectron spectroscopy (such as X-ray photoemission spectroscopy (XPS)), which may reduce the signal disturbance by surface effects and/or contamination. In addition, high spatial resolution of EELS embedded in TEM yields higher reliability on the physical origins of spectrum for deep analyses. For instance, the EELS spectra from the inside of grains and grain boundaries can be collected precisely for a comparison. The energy loss near edge structure (ELNES) spectrum of N-K edge of the N-doped NiO thin film (40 sccm  $N_2$ ) is shown in Fig. 4a. In addition, the N-K edge ELNES spectra of molecular  $N_2$ ,<sup>37</sup> BN nanoparticles<sup>38</sup> and AlN thin film are also plotted alongside for comparison. As seen, a sharp peak with the full width at half maximum (FWHM) of 1 eV located at near 400 eV is clearly identified in N-doped NiO thin film. Beside, another broad peak with small intensity is observed in the energy range of 405–430 eV. These spectral features look quite similar to those of molecular  $N_2$  (see Fig. 4a). On the contrary, the ELNES shows significant discrepancies with that of BN and AlN, where the spectra show multi-peak features with larger widths at the onset of the N-K edge. Therefore, these comparisons imply that the N-K edge ELNES spectrum of the N-doped thin film exhibits the  $N_2$  molecular state.

The nitrogen chemical states located inside the grains and at the grain boundaries have been studied by the EELS line profile. Scanning transmission electron microscopy (STEM) image of the N-doped NiO thin film grown with 40 sccm  $N_2$  is shown in Fig. 4b, where the red line indicates the selected region for EELS line profile analyses. Fig. 4c exhibits the heatmap of ELNES intensity as a function of the energy loss and position for nine measurement points with an interval of around 1 nm. As seen in Fig. 4c, the ELNES spectra show the sharp peaks at around 400 eV irrelevant to the positions. It indicates that the detected nitrogen always predominantly exhibits the molecular state in the N-doped thin film, regardless whether in grains or grain boundaries. Therefore, the molecular state of nitrogen seems to be the intrinsic properties in our samples. The nitrogen content in the N-doped NiO thin film (40 sccm  $N_2$ ) has been estimated by energy-dispersive spectroscopy (EDS) embedded in TEM, which gives the N composition of around 3.6 at%.

Fig. 5a shows formation energies of various  $N_2$ -related defects. Several types of dimer  $N_2$  defects, including  $N_2$  on the nickel site ( $(N_2)_{Ni}$ ),  $N_2$  on the oxygen site ( $(N_2)_O$ ),  $N_2$  interstitial ( $(N_2)_i$ ), and  $N_2$  with multiple vacancies ( $N_2 + V_{Ni} + V_O$ ,  $N_2 + 2V_{Ni} + 2V_O$ , and  $N_2 + 4V_{Ni} + 4V_O$ ), have been taken into consideration. One can see that  $(N_2)_i$  shows very high formation energy, attributed to the huge size mismatch between the interstitial site and  $N_2$  molecule. Except for it,  $N_2$ -related defects are relatively low in formation energies with deep levels. Especially,  $N_2 + V_{Ni} + V_O$  acts as deep acceptor with a relatively low formation energy. This is qualitatively consistent with our experimental results: (i) the subgap optical absorption increases significantly with increasing the  $N_2$  flow rate (see Fig. 2), and (ii) the hole mobility decreases with increasing the  $N_2$  flow rate (see Fig. 3).

As shown in Fig. 5b, we also checked single N-related defects, including nitrogen on the nickel site ( $N_{Ni}$ ), nitrogen on the

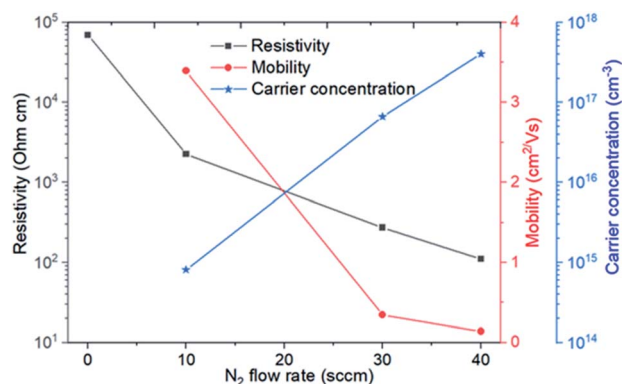


Fig. 3 Resistivity, mobility and carrier concentration of N-doped NiO thin films as a function of  $N_2$  flow rate.





## Author contributions

Y. W. carried out main parts of the experiments and wrote the first version of this paper. Y. W. and J. F. P. proposed the concept. S. B. and J. G. performed TEM analyses. Y. K., N. T. and F. O. completed the theoretical calculations. H. S. and B. D. contributed to the analysis and discussion for the results. All authors commented on the manuscript.

## Conflicts of interest

There are no conflicts to declare.

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