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# Thermoelectric performance enhancement of eco-friendly $Cu_2Se$ through incorporating $CB_4$

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We have prepared $Cu_2Se$ + $x$ wt% $CB_4$ composites with $x = 0$, 0.1, 0.3, 0.5, and 0.7 by a hydrothermal method and hot-pressing technique. The structural and compositional analysis indicates that pure phase $Cu_2Se$ powders were synthesized and the densified layered bulk samples were obtained. Electrical properties testing showed that the sample with $x = 0.5$ has the high power factor of 0.886 mW m$^{-1}$ K$^{-2}$ due to its high Seebeck coefficient. Meanwhile, the thermal conductivity was suppressed to 0.6 W m$^{-1}$ K$^{-1}$ at 773 K. As a result, the final optimized $ZT$ value of 1.46 at 773 K was achieved. These results suggest that $CB_4$ could be an alternative inclusion to improve effectively the thermoelectric performance of $Cu_2Se$.

## 1. Introduction

With the wide demand for environment-friendly energy, thermoelectric (TE) materials have received significant attention during the past two decades.[1–6] TE materials can realize the conversion directly between thermal and electrical energy.[7–9] High-efficiency TE materials require a high dimensionless figure of merit ($ZT$), $ZT = S^2\sigma T/\kappa$, where $S$ is the Seebeck coefficient, $\sigma$ is the electrical conductivity, $T$ is the absolute temperature, and $\kappa$ is the total thermal conductivity respectively.[10–13] Numerous exploratory efforts tried to optimize these parameters, *i.e.*, enhance the power factor (PF = $S^2\sigma$) and reduce the thermal conductivity simultaneously, to achieve high TE performance.[14,15] However, these parameters are strongly coupled together, making it challenging to improve the $ZT$s.[16,17]

Chalcogenide TE materials have attracted widespread interests and have been intensively studied because of their intrinsically low lattice thermal conductivity ($\kappa_L$).[18–21] Very recently, Chung and co-workers reported[22] the ultralow $\kappa_L$ of 0.07 W m$^{-1}$ K$^{-1}$ obtained in polycrystal SnSe and thereby got a high $ZT$ of 3.1, which breaks thermoelectric performance limits.[23] Cu-based chalcogenide also received much attention for environmental-friendly TE application and $Cu_2Se$ is one of the strongest candidates.[24,25] To improve the TE performance of $Cu_2Se$, numerous strategies have been developed, such as nano-engineering,[26,27] solid solution alloying,[28,29] element doping,[30–32] and introducing highly dispersed nanoparticles.[33–36] Among them, adding nanoparticles to $Cu_2Se$ matrix has been demonstrated to be an effective and prevailing approach to improve TE performance. The nanoparticles in composites, such as $TiO_2$,[37] graphene,[38] nano-boron,[39] carbon dots[40,41] and CNTs,[42] can effectively scatter intermediate frequency phonons and hence suppress $\kappa_L$.

In this work, we try to introduce boron carbide ($CB_4$) as a second phase in the $Cu_2Se$ matrix because nano-$CB_4$ has good thermal stability, high temperature electrical conductivity and relatively large Seebeck coefficient.[43] For preparing $Cu_2Se$ samples with different contents of $CB_4$, the hydrothermal method and hot-pressing technique are employed. The results show that $CB_4$ inclusions do reduce $\kappa$ significantly and meanwhile enhance PF slightly over a wide range of temperature. As a result, the highest $ZT$ reaches 1.46 at 773 K, which is much higher than the value of undoped sample.

## 2. Experimental section

$Cu_2Se$ nanopowders were prepared by hydrothermal method and $CB_4$ nanopowders were added during the hydrothermal process. We first mixed 20 mmol copper chloride dehydrate ($CuCl_2 \cdot 2H_2O$, 99.9%), 10 mmol selenium dioxide ($SeO_2$, 99.9%) and a recommended amount of $CB_4$ (99%) into about 60 milliliters deionized water, and then magnetic stirred continuously for 12 h. The amount of $CB_4$ is determined by the mass ratio of $Cu_2Se$ and $CB_4$. In this work, we consider five $Cu_2Se$ + $x$ wt% $CB_4$ ($x = 0$, 0.1, 0.3, 0.5, 0.7) samples, which are named as S0, S1, S3, S5 and S7 respectively for the convenience of description. After the components were completely dissolved, we added 20 milliliters hydrazine hydrate ($N_2H_4 \cdot H_2O$) into the mixture solution. The solution was then loaded into a Teflon-lined stainless steel autoclave and sealed. The autoclave was heated at 453 K for 24 h and cooled down to room temperature naturally. The product was collected by centrifuging and washed by deionized water and anhydrous ethanol for several times. We finally dried the product in vacuum at 333 K for at least 12 h. For getting the final samples, the synthesized powders were sintered into pallets with a diameter of 20 mm by

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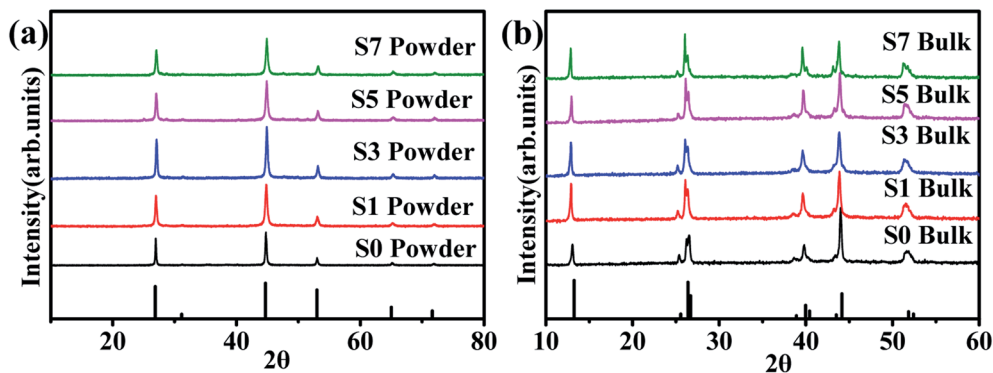
**Fig. 1** (a) XRD patterns of the $\beta$-$Cu_2Se$ (ICDD file No. 00-006-0680) synthesized nanopowders. (b) XRD patterns of $\alpha$-$Cu_2Se$ (ICDD file No. 00-019-0401) bulk samples.

hot-pressing method at 973 K and under a pressure of 40 MPa in the vacuum for 30 minutes.

The phase compositions of the prepared samples were identified by X-ray diffraction measurements (XRD, PANalytical, Netherlands) with Cu K$\alpha$ radiation. The fracture surface microstructure was observed by field emission scanning electron microscopy (SEM, TM4000) and Energy-dispersive X-ray spectroscopy (EDS, X-MAX$^N$20, England). Seebeck coefficients
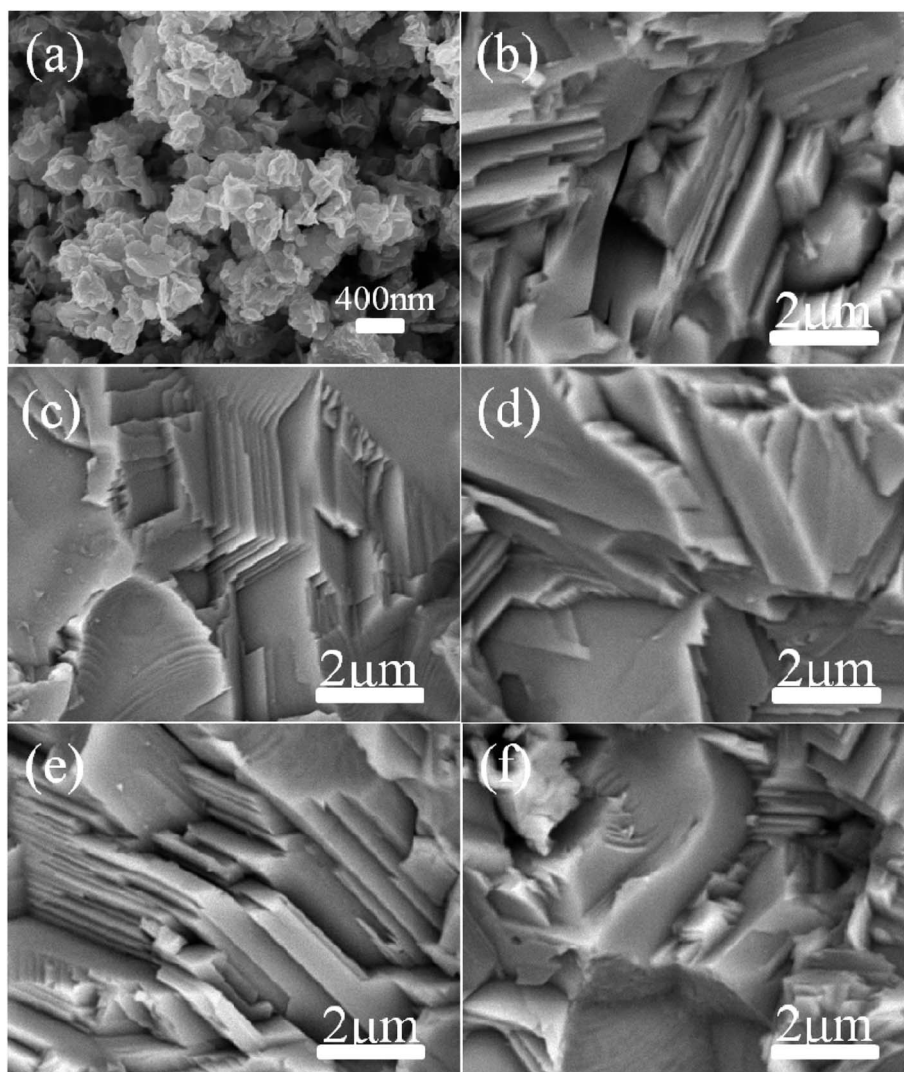


**Fig. 2** SEM images for (a) $Cu_2Se$ nanopowders, the samples of (b) S0, (c) S1, (d) S3, (e) S5 and (f) S7.

and electrical conductivity simultaneously were measured by using a CTA-3 system (Cryoall, China) from 323 K to 773 K in a helium atmosphere. $\kappa$ was calculated according to the formula $\kappa = \rho D C_p$, where $\rho$ is the density of the samples, $D$ is the thermal diffusivity, and $C_p$ is the heat capacity respectively. $D$ was measured using the laser flash method (DFL-1600, Netzsch) on a square piece specimen with 10 mm in side and 1.5 mm in thickness. $C_p$ was measured by differential scanning calorimetry (DSCQ-2000, TA). The density of the samples $\rho$ was obtained from Archimedes' method.

## 3.   Results and discussion

Fig. 1(a) shows the XRD patterns of the synthesized $Cu_2Se$ + $x$ wt% $CB_4$ ($x = 0, 0.1, 0.3, 0.5, 0.7$) nanopowders. The main diffraction peaks are all in consistent with the ICDD file No. 00-006-0680, indicating β-phase $Cu_2Se$ is prepared successfully.[40,44–46] No $CB_4$ phase is to be detected, which might be due to introducing only a small amount into the main β-phase

$Cu_2Se$. Fig. 1(b) shows the XRD patterns of the sintered bulk samples. The results match with the ICDD file No. 00-019-0401 very well, indicating that the as-prepared β-$Cu_2Se$ nanostructures has transferred to pure α-$Cu_2Se$ after sintering.[40,46,47]

The typical morphology of the synthesized $Cu_2Se$ precursor nanopowders is shown in Fig. 2(a). Nanoplates with a lateral size of 100–500 nm, often aggregating together, can be clearly observed. Fig. 2(b–f) display the cross-sectional microstructures of the hot-pressed $Cu_2Se$ + $x$ wt% $CB_4$ bulk samples. Similar lamellar structure is observed for each sample. Such structure is conducive to reducing the thermal conductivity due to the increased scattering of phonons between layered interfaces.[48] A few pores on the surface can be spotted which might be due to Se volatilization during the sintering process and hence the corresponding density is slightly lower than the pure $Cu_2Se$. The samples S0, S1, S3, S5, and S7 have densities of 5.75, 5.65, 5.38, 5.29, 5.20 g cm$^{-3}$ respectively (theoretical density of $Cu_2Se$ is 6.74 g cm$^{-3}$). Fig. 3 shows the surface EDS mapping of the sample S5. These maps suggest that copper, selenium, carbon

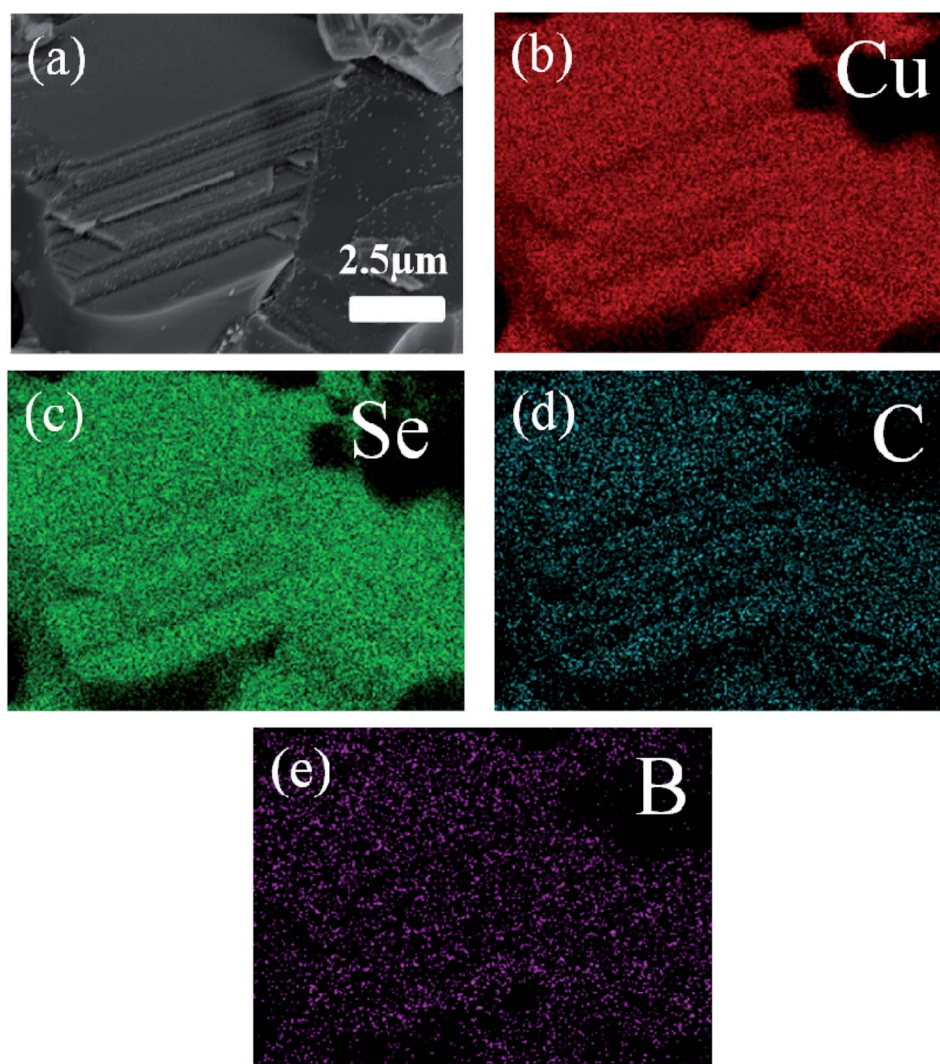

Fig. 3   (a) SEM image of the typical morphology of the S5. (b)–(e) The corresponding EDS maps.

and boron signals overlap, showing $CB_4$ is uniformly distributed in the sample.

The temperature dependence of the electrical conductivity $\sigma$ of our samples is shown in Fig. 4(a). Within the temperature range of 400–773 K, the electrical conductivity $\sigma$ of all samples decreases monotonously with temperature increasing, which shows the typical behaviors of metallic-like characteristic. However, we can clearly find that the electrical conductivity $\sigma$ varies abnormally around 400 K, which have been proved to be due to the phase transition from $\alpha$-$Cu_2Se$ to $\beta$-$Cu_2Se$.[49] In addition, the electrical conductivity $\sigma$ decreases with the increasing of $CB_4$ components at each testing temperature. For example, compared to the sample S0, the electrical conductivity $\sigma$ is about 3 times smaller around 400 K and 2 times smaller at 773 K for the sample S7. This trend in electrical conductivity $\sigma$ is very different from the report of $Cu_2Se$/SiC composite,[27] but similar to the report of $Cu_2S$/SiC composite.[17] For clarifying the underlying mechanism of the electrical transport behaviors, the room temperature carrier concentration ($n$) and carrier mobility ($\mu$) are measured, as shown in Table 1. Clearly, there is a reduction in carrier concentration and an increase in carrier mobility with increasing of the $CB_4$ composite content. For the superionic conductor $Cu_2Se$, unstable copper ions show a liquid-like behavior and hence cause the unrepeatable electrical transport properties. The carrier concentration $n$ contains two parts: the majority carrier ($n_{hole}$) and the ion carrier ($n_{ion}$), and therefore, the electrical conductivity $\sigma = \sigma_{hole} + \sigma_{ion}$. In our

work, the introduction of $CB_4$ might effectively block the migration of copper ions, which causes the reduction in $n_{ion}$ and hence in the whole carrier concentration $n$. In addition, the migration restrictions of copper ions can also effectively enhance the stability of $Cu_2Se$ chemical structure, which is confirmed by the cyclic testing shown in Fig. 4. The increase in mobility $\mu$ might be due to the grain growth during the process of hot pressing, as can be clearly seen in Fig. 2. Because of the difficulty for measuring the transport of ions, more related discussions do not be present here.

In our work, the samples S5 and S7 show a quite low electrical conductivity $\sigma$. This fact would not be very good for optimizing the thermoelectric properties of $Cu_2Se$. However, thankfully, the reduction in electrical conductivity $\sigma$ is often accompanied by enhancement of the Seebeck coefficient. As we can see from Fig. 4(b), the samples S5 and S7 exhibit a larger

**Table 1** Carrier concentration ($n$), mobility ($\mu$) at room temperature

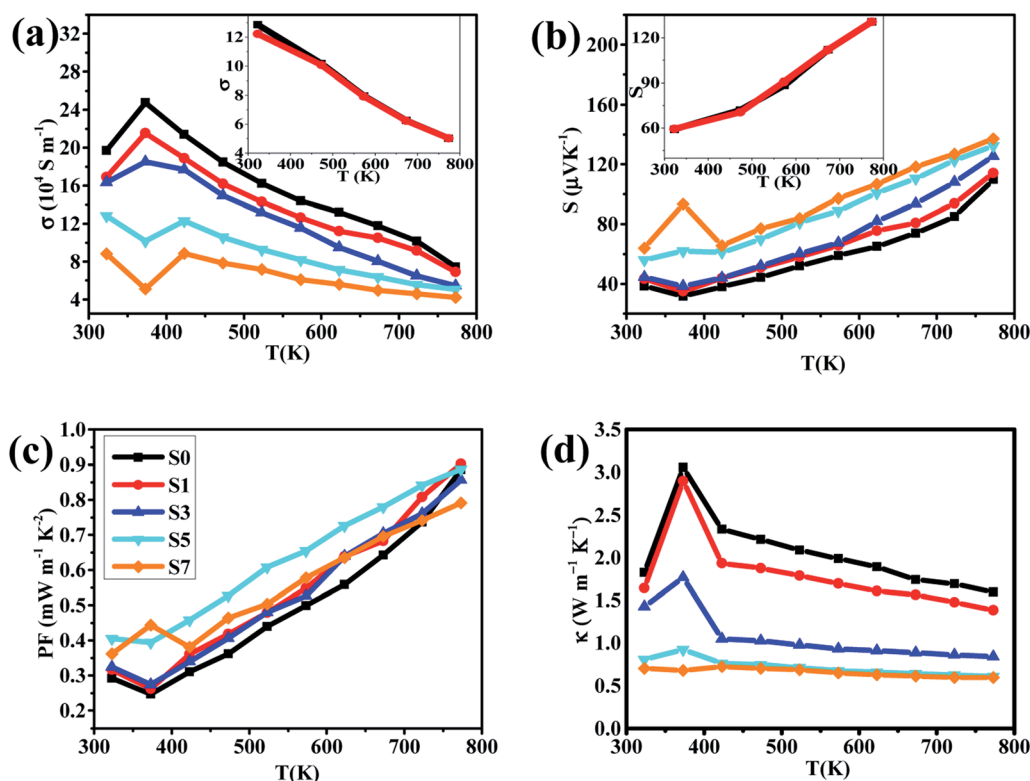| Sample | $n_{hole} + n_{ion}$ ($10^{19}$ cm$^{-3}$) | $\mu$ (cm$^2$ V$^{-1}$ s$^{-1}$) |
|---|---|---|
| S0 | 7.9 | 73.6 |
| S1 | 6.6 | 101 |
| S3 | 5.5 | 104 |
| S5 | 5.3 | 112 |
| S7 | 3.6 | 119 |



**Fig. 4** Temperature dependence of (a) electrical conductivities, (b) Seebeck coefficients, (c) power factors and (d) thermal conductivities for the samples of S0, S1, S3, S5 and S7. Insets of (a) and (b) are the heating (red line) and cooling (black line) curves for electrical conductivity and Seebeck coefficient of sample S5 respectively.

Seebeck coefficient than the others within the entire testing temperature range and the optimal value reaches 137 μV K$^{-1}$ at 773 K for the sample S7. Additionally, the positive Seebeck coefficient for all samples suggests the *p*-type electrical transport in Cu$_2$Se. The calculated PFs of the samples are shown in Fig. 4(c). Very clearly, PFs curves of the samples S0, S1, S3 and S5 show a valley around 400 K due to the phase transition from β-Cu$_2$Se to α-Cu$_2$Se,[50,51] while the sample S7 has the abnormal high PF because of its very high *S* and low σ. However, we cannot evaluate precisely these parameters around 400 K because they are very sensitive to temperature change during the process of phase transformation. This might be the reason that some previous studies did not show the related data close to 400 K.

As *T* >400 K, all PFs increase rapidly as temperature increases. However, the sample S7 does not show the superior PF value, which might be due to its very low electrical conductivity. In contrast, the sample S5 has the optimal PF as *T* >400 K (PF$_{max}$ = 0.886 mW m$^{-1}$ K$^{-2}$ at 773 K), indicating that a moderate content of CB$_4$ composites may result in a better electrical performance of Cu$_2$Se.

Another important parameter to determine the TE properties is the thermal conductivity κ. The thermal conductivity can be divided into two parts: κ = κ$_L$ + κ$_{carrier}$, where κ$_L$ comes from heat phonons travelling through the crystal lattice and κ$_{carrier}$ arises from heat carriers' movement. Generally, at high temperature, κ$_L$ dominates in Cu$_2$Se matrix.[39] The measured thermal conductivity is shown in Fig. 4(d). Clearly, all of the thermal conductivity κ decrease with temperature increasing, which is a reasonable trend because of the increasing in phonon scattering and the reduction in carrier concentration. For a fixed temperature, as CB$_4$ component increases from 0.0 to 0.7%, the thermal conductivity decreases very rapidly. To take a step further, carrier thermal conductivity (κ$_{carrier}$) is calculated by using the Wiedemann–Franz relationship, κ$_{carrier}$ = *L*σT, where *L* is the Lorenz number. When the detailed band structure and scattering mechanism are not known, the Lorenz number cannot be viewed as a constant but can be calculated by *L* = 1.5 + exp[−|*S*|/116], where *L* is in unit of 10$^{-8}$ W Ω K$^{-2}$ and *S* (Seebeck coefficient) in unit of μV K$^{-1}$.[52] The calculated *L* and κ$_{carrier}$ are showed in Fig. 5. However, it seems that the measured total

thermal conductivity of S5 and S7 are lower than the calculated values. The similar phenomenon happened to Cu$_2$Se/SiC system[27] and Cu–S system.[17,53] It might be due to the different thermal transport properties between electron/hole carriers and ion carriers. The equation κ$_{carrier}$ = *L*σT may not be suitable for Cu$_2$Se/CB$_4$ composite. The ion carriers have a non-negligible contribution to total thermal transport, so the modified equation can be taken as κ$_{carrier}$ = *L*(σ − σ$_{ion}$)T + κ$_{ion}$,[27,30,53] where κ$_{ion}$ is the ionic thermal conductivity, σ$_{ion}$ is ionic electrical conductivity. However, σ$_{ion}$ and κ$_{ion}$ cannot be measured precisely now and hence the carrier thermal conductivity and lattice thermal conductivity are not offered here.

Considering the total thermal conductivity, the samples S5 and S7 remain very low values and very close to each other as *T* >400 K. The lowest κ value is 0.59 W m$^{-1}$ K$^{-1}$ at 773 K for S7 sample. This result indicates that the higher concentration of CB$_4$ inclusions could not improve further the thermal conductivity. Moreover, as shown in Fig. 4(c), the PF of sample S7 is also smaller than sample S5. Therefore, in this work, the S5 sample would show the optimal thermoelectric performance. This can be confirmed positively by the relationship between ZT values and temperature, as shown in Fig. 6. Sample S5 exhibits

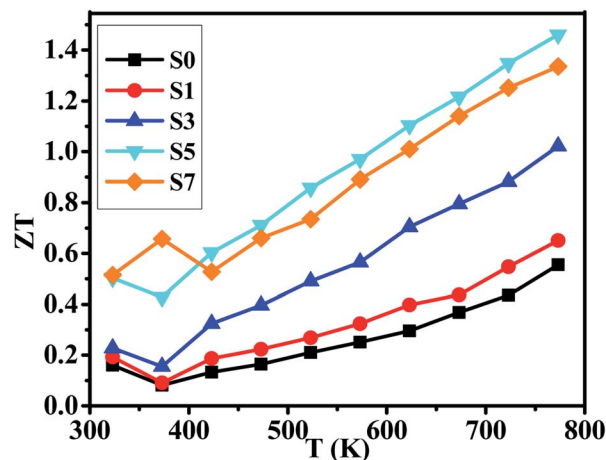**Fig. 6** Temperature dependence of *ZT*s for samples of S0, S1, S3, S5 and S7.
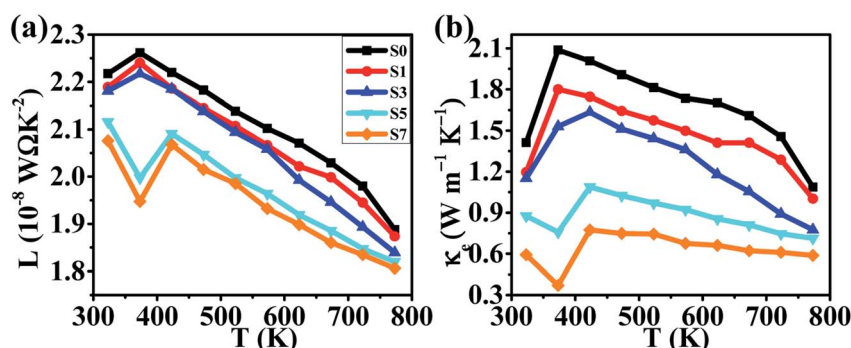
**Fig. 5** Temperature dependence of (a) the calculated Lorenz number, (b) the calculated thermal conductivities for samples of S0, S1, S3, S5 and S7.

the largest $ZT$ than the others as $T > 400$ K and the largest value is of 1.46 at 773 K, which is much larger than that of undoped sample S0 ($ZT_{max} = 0.55$ at 773 K). The average $ZT$ value of S5 ($ZT_{average} = 1.034$) is also larger than one within the temperature range of 400–773 K. These results suggest that $Cu_2Se$ compositing a certain amount of $CB_4$ is very promising for TE application.

## 4. Conclusions

We have prepared the pure phase $Cu_2Se + x$ wt% $CB_4$ ($x = 0$, 0.1, 0.3, 0.5, 0.7) nanopowders by hydrothermal method and have sintered the corresponding bulk samples by hot-pressing technique. With the content of $CB_4$ increasing, the electrical conductivities of our samples decrease very rapidly, while the Seebeck coefficients steadily increase. The final PFs show that the sample S5 possesses the optimal value within the testing temperature range. The thermal conductivity testing confirms that the $CB_4$ nano-inclusions can drastically reduce the $\kappa$ of our samples. The lowest $\kappa$ value is 0.59 W m$^{-1}$ K$^{-1}$ at 773 K for S7 sample. However, the $\kappa$ of sample S5 is very close to that of sample S7 as $T > 400$ K. Finally, the optimized $ZT$ value of 1.46 at 773 K is achieved for S5 sample, which is about 165% larger than the undoped sample.

## Conflicts of interest

The authors have no conflicts to declare.

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