



Cite this: *RSC Adv.*, 2022, **12**, 13749

Janus transition metal dichalcogenides in combination with MoS₂ for high-efficiency photovoltaic applications: a DFT study†

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Exotic features of two-dimensional materials have been demonstrated, making them particularly appealing for both photocatalytic and photovoltaic applications. van der Waals corrected density functional theory calculations were performed on AAII-Se MoS₂, AAII-Te MoS₂, and AAII-Se WS₂ heterostructures in this study. Our findings reveal that the heterostructures have high stability due to the tiny lattice mismatch and binding energy, which is extremely favorable for epitaxial growth of these heterostructures. According to the electronic band gap calculation, AAII-Se MoS₂ and AAII-Se WS₂ are semiconducting materials, while AAII-Te MoS₂ has metallic properties. Interestingly, all three heterostructures have type II band gap alignment, which is advantageous for photovoltaic and photocatalytic applications. Furthermore, it was discovered that AAII-Se MoS₂ and AAII-Se WS₂ heterostructures exhibit high power conversion efficiency of up to 12.15% and 9.37%, respectively. Based on these intriguing features, the two heterostructures are excellent prospects for photovoltaic applications. The heterostructures have no appropriate band edge sites for overall water splitting at pH = 0, but they are good for the oxygen evolution process. It is feasible to alter the position of the band edges using strain resulting in improved overall water splitting by the heterostructures.

Received 5th February 2022
 Accepted 1st May 2022

DOI: 10.1039/d2ra00775d
rsc.li/rsc-advances

Introduction

The discovery of 2D materials has accelerated^{1,2} since the discovery of graphene in 2004,³ including that of molybdenum disulphide (MoS₂)^{4,5} graphitic carbon nitride (g-C₃N₄)^{6,7} phosphorene,^{8,9} and MXenes.^{10–12} The piqued interest of many researchers is because of their exceptional electronic and mechanical properties as well as possible application in nano-devices. In addition, a number of derivatives of 2D materials are emerging, such as spliced 3D compounds, Janus 2D materials and 2D van der Waals heterojunctions.^{13–17} 2D van der Waals (vdW) heterojunctions are gaining significant attention because it combines the advantages of different from monolayers. These 2D heterostructures combines two or more different layered materials, which could result in novel properties and potential applications.^{18,19} The relative weak interlayer coupling of

heterostructures allows for most of the intrinsic properties from the individual monolayers to remain intact.²⁰ Thus, the opportunities to combine the intrinsic different properties of various monolayers is an underlying advantage of heterostructures.²⁰ Also, 2D vdW heterostructure could overcome the shortcomings of single materials, such as low quantum efficiency, high charge recombination, and serious chemical back-reactions.^{21,22}

The world's finite energy supplies of non-renewable and need to decarbonize to meet the growing need for energy calls for the need to rapidly develop renewable energy. As a result, photocatalytic water splitting, and solar cell technology have been seen as two suitable approaches to providing clean energy. Several studies have been carried out, which focused on finding an efficient photocatalyst^{23–27} and solar cell^{28–30} materials, such as the development of MoS₂ monolayer as a potential photocatalyst for water splitting^{31–33} and solar cell³⁴ applications. The photogenerated electron–hole pairs in MoS₂, on the other hand, remain in the same spatial locations, resulting in a high rate of recombination. Photogenerated electron–hole pairs in the MoS₂ monolayer, remain in the same spatial locations resulting in a high rate of recombination. vdW heterostructures can overcome this issue while also providing improved overall characteristics. There are three types of band alignment for semiconductor heterojunction, including straddling gap (type I), staggered gap (type II) and broken gap (type III).³⁵ Because the valence-band maximum (VBM) and conduction-band minimum (CBM) reside in distinct layers, type-II (staggered)³⁶

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† Electronic supplementary information (ESI) available. See <https://doi.org/10.1039/d2ra00775d>.



heterostructures appear promise for achieving efficient charge carrier separation. As a result, when photogenerated electron-hole couples break at the interface, the electrons are moved to one layer and the holes are sent to another. This property makes type-II heterostructures ideal for use in photovoltaic and photocatalytic devices. Wang, S. *et al.*³⁷ investigated photocatalytic devices based on MoS₂/ZnO van der Waals heterostructure and Bernardi *et al.*³⁸ investigated photovoltaic devices based on the MoS₂/WS₂ bilayer. Both studies discovered that the heterostructures exhibit type-II band edge alignment. From their investigations, the band edge positions satisfy the energy levels required for the water splitting reaction in the photocatalytic application, and as a photovoltaic device, the power-conversion efficiency of up to 1% and a far higher power density than existing ultrathin solar cells was observed. The findings showed that creating hetero-structured materials based on MoS₂ have potential application for photocatalytic water splitting and photovoltaic.

Recently Janus 2D materials particularly Janus transition metal dichalcogenides (TMD) material, are gaining increasing attention due to their distinct properties, which are different from the traditional 2D materials.³¹ Unlike typical TMDs, Janus TMD (JTMD) monolayers lack both in-plane and out-of-plane symmetries, giving them an extra degree of freedom to regulate their properties.^{32,39,40} Using chemical vapor deposition technique, a JTMD material (*i.e.*, MoS₂) was recently produced experimentally by selective selenization of the top atomic layer in a single-layer MoS₂.^{41,42} JTMD's fascinating features imply that they are crucial 2D materials. Furthermore, unlike traditional vdW heterostructures, the intrinsic interlayer polarization from Janus material will couple with the interlayer built-in polarization field, providing an additional degree of freedom to modulate the physical/chemical properties of the heterostructure, resulting in novel features and potential applications.²⁰

We used *ab initio* calculations to determine the structural and electrical properties of the MoS₂/JTMD (JTMD = MoS₂, MoSTe, MoSeTe, WS₂, and WSe₂) bilayer in this study. The calculated binding energies of the bilayer system indicates that configurations AAII-Se MoS₂, AAII-Te MoSTe and AAII-Se WS₂ is stable. The calculated electronic band gap for the two heterostructure is semiconducting and one shows metallic property. The AAII-Se MoS₂, AAII-Te MoSTe and AAII-Se WS₂ vdW heterostructures forms a typical type-II band alignment, which lead to effective separation of electron-hole pairs. Moreover, the AAII-Se MoS₂ and AAII-Se WS₂ vdW heterostructure also exhibits very good power conversion efficiency with a value of (7.289–12.15%) and (5.62–9.37%) respectively suggesting that they are a good candidate for photovoltaic applications. The heterostructures might be a suitable material for photocatalytic water splitting if strain is applied on them.

Computational method

Density functional theory (DFT)⁴³ with the projected augmented wave (PAW) and generalized gradient approximation (GGA)⁴⁴ with the PBE variant as implemented in the Quantum ESPRESSO and CASTEP codes were used to perform first-

principles computations on various MoS₂ and Janus TMDs heterostructures. The plane wave basis set's kinetic energy cutoff is set to 540 eV and the Brillouin zone is sampled using a $12 \times 12 \times 1$ Monkhorst-Pack *k*-point grid. All atomic positions were completely relaxed until the force was less than 0.01 eV Å⁻¹, and the overall energy convergence criteria was set to 10^{-5} eV. To preclude contact between the neighboring slabs, a 25 Å vacuum was presented in the *z*-direction. Grimme's DFT-D2 method has been widely used to provide a reliable description of all types of vdW heterostructures.^{45,46} As a result, throughout the calculations, the Grimme DFT-D2 dispersion correction method is used. Next, the HSE06 hybrid functional was used in the current investigation to perform valence band alignment and band structure computation.^{47,48} Binding energy (E_b) of the heterostructures is calculated using,

$$E_b = E_{\text{heterostructure}} - (E_{\text{monolayer1}} + E_{\text{monolayer2}}) \quad (1)$$

where E_b is binding energy, $E_{\text{heterostructure}}$ is the total energy of the heterostructures, and $E_{\text{monolayer1}}$ and $E_{\text{monolayer2}}$ are the total energy of the JTMDs, and MoS₂ monolayers, respectively.

Results and discussions

(a) Structure properties and energetics

In this study, the monolayers and heterostructures of MoS₂ TMD and five Janus TMDs (MoS₂, MoSTe, MoSeTe, WS₂, and WSe₂) was considered. One monolayer of MoS₂ and a variation of the five different Janus TMDs was used to construct the heterostructures. The calculated optimum lattice parameter for MoS₂ is 3.19 Å, MoS₂ is 3.26 Å, MoSTe is 3.36 Å, MoSeTe is 3.42 Å, WS₂ is 3.26 Å, and WSe₂ is 3.45 Å as presented in Table 1. These calculated lattice parameters agree with the reported values in ref. 49–55. The lattice mismatch between MoS₂ and the five Janus TMDs (MoS₂, MoSTe, MoSeTe, WS₂, WSe₂) was estimated, as shown in Table 1. Three of the five heterostructures considered in the current study have a lattice mismatch $\leq 5\%$ as presented in Table 1. Experimentally,⁵⁶ lattice mismatch less than or equal to 5 have been shown to be appropriate for generation of vdW heterostructures. The three heterostructures with appropriate lattice mismatch are MoS₂/MoS₂, MoS₂/MoSTe, and MoS₂/WS₂. These structures are further considered as candidate materials for photovoltaic and photocatalytic applications.

As presented in Fig. 1, eight different stacking patterns for the heterostructures were created. In Fig. 1, different schematic

Table 1 Optimized lattice parameter and lattice mismatch

MXY	Current work lattice parameter a (Å)	Lattice mismatch (%)
MoS ₂	3.19	
MoS ₂	3.26	2.147
MoSTe	3.36	5.059
MoSeTe	3.42	6.725
WS ₂	3.26	2.147
WSe ₂	3.45	7.536



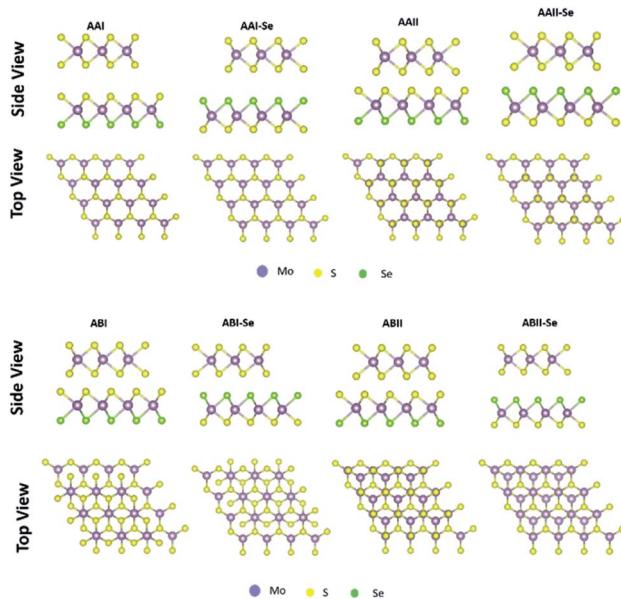


Fig. 1 Stacking patterns of $\text{MoS}_2/\text{MoSSe}$ heterostructure with different stacking orientations. AA and AB refer to the stacking pattern, whether the two hexagonal lattices are aligned on top of one another or shifted. I and II refer to whether the Mo is aligned below the metal or the chalcogen atom of JTMDs. Se implies that the Se chalcogen atom is facing the MoS_2 monolayer if not mentioned the other chalcogen of the JTMD is facing the MoS_2 monolayer.

representation of the $\text{MoS}_2/\text{MoSSe}$ heterostructure studied in this work is presented. The stacking pattern referred to as AAI have the Mo atom transition metal of one layer placed on the top of metal atom of second layer, while chalcogen atom of one layer is on top of the chalcogen of the second layer, AAII-Se stacking is similar to AAI stacking except that the chalcogen facing the MoS_2 is Se instead of S, stacking AAII has the Mo atom from the MoS_2 monolayer on top of the chalcogen atom of the Janus TMDs, while the S from the MoS_2 is on top of the metal atom of the Janus TMDs, stacking AAII-Se is similar to the AAII pattern except that the Se is facing the MoS_2 instead of S, ABI stacking similar to AAI stacking but shifted in the horizontal direction, the Mo from MoS_2 is on top of the metal atom of the Janus TMDs but the chalcogen of the MoS_2 is not on the top of the chalcogen of the Janus TMDs, ABI-Se stacking is similar to ABI stacking except the Se is facing the MoS_2 instead of S facing MoS_2 , ABII stacking is similar to AAII stacking but shifted in the horizontal direction, the chalcogen from the MoS_2 is on top of the metal atom of the Janus TMDs but Mo is not on top of the chalcogen of the Janus TMDs, ABII-Se stacking the same configuration as ABII stacking except the Se is facing the MoS_2 instead of S facing MoS_2 .

To determine the stability of the hetero-structures and the most favourable stacking pattern for each TMD-JTMD configuration, we calculated the binding energy for all the stacking patterns. This consist of 24 different structures and the results are presented in Table 2.

The most stable stacking pattern of the TMD-JTMD, which is the configuration with the most negative binding energy is used

Table 2 Binding energies (E_b in eV) and stacking pattern of different configuration of $\text{MoS}_2/\text{Janus TMDCs}$ heterostructure

Heterostructure	Stacking pattern	Binding energy E_b (eV)
$\text{MoS}_2/\text{MoSSe}$	AAI	-0.006
	AAI-Se	-0.006
	AAII	-0.009
	AAII-Se	-0.011
	ABI	-0.008
	ABI-Se	-0.010
	ABII	-0.009
	ABII-Se	-0.011
	AAI	0.001
	AAI-Te	-0.002
	AAII	-0.003
	AAII-Te	-0.009
	ABI	-0.003
	ABI-Te	-0.006
	ABII	-0.004
	ABII-Te	-0.008
$\text{MoS}_2/\text{MoSTe}$	AAI	-0.008
	AAI-Se	-0.009
	AAII	-0.013
	AAII-Se	-0.015
	ABI	-0.012
	ABI-Se	-0.013
	ABII	-0.013
	ABII-Se	-0.015
	AAI	-0.008
	AAI-Te	-0.002
	AAII	-0.003
	AAII-Te	-0.009
MoS_2/WSSe	AAI	-0.008
	AAI-Se	-0.009
	AAII	-0.013
	AAII-Se	-0.015
	ABI	-0.012
	ABI-Se	-0.013
	ABII	-0.013
	ABII-Se	-0.015
	AAI	-0.008
	AAI-Te	-0.002
	AAII	-0.003
	AAII-Te	-0.009

for further studies. As shown in Table 2, the stable stacking pattern are the AAII-Se MoSSe, AAII-Te MoSTe, and AAII-Se WSSe heterostructures. The electronic properties, photocatalytic water splitting, and photovoltaic solar cells performance are presented and discussed below.

(b) Electronic properties

The electronic structure calculation of the investigated materials is summarized below in Tables 3 and 4 for the monolayers and heterostructures respectively. Table 3 shows the optimized lattice constant a (Å), band gap (calculated using PBE and HSE06 in eV), work function (φ in eV), valence and conduction band edge relative to the vacuum (EVB & ECB) for the three considered heterostructures. The calculated band gaps are indirect (MoSSe and MoSTe) and direct (MoS_2 and WSSe) using

Table 3 Optimized lattice constant a (Å), band gap (calculated using PBE and HSE06 in eV), work function (φ in eV), valence and conduction band edge relative to the vacuum (EVB & ECB) for the monolayers

Monolayer	MoS_2	MoSSe	MoSTe	WSSe
a (This study)	3.19	3.26	3.260	3.26
a (Expt.)	3.159 [ref. 69]	—	—	—
E_g -PBE	1.74 (dir)	1.65 (dir)	1.16 (ind)	1.77 (dir)
E_g -HSE06	2.278	2.028	1.571	2.195
E_g -Expt.	1.9 eV [ref. 68]	1.68 eV [ref. 70]	—	—
φ	6.273	5.955	5.289	5.728
EVB	-6.273	-5.955	-5.289	-5.728
ECB	-3.995	-3.927	-3.718	-3.533



Table 4 Optimized lattice constant a (Å), band gap (calculated using PBE and HSE06 in eV), work function (φ in eV), valence and conduction band edge relative to the vacuum (EVB & ECB) for the Heterostructures

Heterostructure	AAII-Se MoSSe	AAII-Te MoSTe	AAII-Se WSSe
a	3.220	3.270	3.220
E_g -PBE	1.026 (ind)	0.303 (ind)	0.838 (dir)
E_g -HSE06	1.133	0.000	0.928
φ	5.904	5.435	5.646
EVB	-6.33	-5.90	-6.07
ECB	-3.73	-4.47	-3.67

PBE functional. The calculated band gap using the HSE06 functional is consistently larger than those determined using the PBE functional for all the considered monolayers.

To the best of our knowledge, two types of Janus layered material have been synthesized thus far, which are the MoSSe and WSSe monolayers. As such, a direct comparison of the evaluated findings with the experimental investigation is not feasible. However, the theoretical method applied yields consistent qualitative results with experimental data. The calculated lattice parameters 3.19, 3.26, 3.26 Å and band gap 2.278, 2.03, 2.2 eV for the pristine MoS₂, MoSSe, WSSe, are consistent with available experimental data as presented in Table 3. MoSSe has a lattice constant that falls between MoS₂ (3.159 Å) and MoSe₂ (3.30 Å),^{71,72} which is appropriate given MoSSe's composition.

Table 3 shows the optimized lattice constant a (Å), band gap (calculated using PBE and HSE06 in eV), work function (φ in eV), valence and conduction band edge relative to the vacuum (EVB & ECB) for the three considered heterostructures. The calculated band gaps are indirect (AAII-Se MoSSe and AAII-Te MoSTe) and direct (AAII-Se WSSe) using PBE functional, however with HSE06 we found that the AAII-Se MoSSe and AAII-Se WSSe are semi-conducting whereas the AAII-Te MoSTe is metallic. The AAII-Se MoSSe and AAII-Se WSSe are semiconductor with a band gap value of 1.33 eV and 0.928 eV respectively. These two configurations are further considered as possible photocatalyst and photovoltaic material. The band structure diagram using HSE06 hybrid functional are presented as Fig. 2 for these heterostructures.

(c) Photocatalytic water splitting

Photocatalytic water splitting has received a lot of interest as an energy-efficient and environmentally friendly way to create

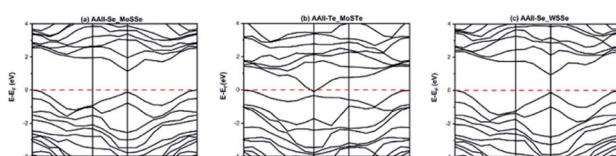


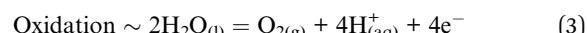
Fig. 2 Band Structure of AAII-Se MoSSe, AAII-Te MoSTe, and AAII-Se WSSe.

hydrogen.^{57,58} The overall photocatalytic water splitting process can proceed as presented in eqn (2):



$$\Delta E^0 = 1.23 \text{ V}$$

where ΔE^0 is the equilibrium potential at standard conditions. The oxidation and reduction half-cell reactions are involved in the photoelectrochemical (PCE) path to water splitting and is presented below:



$$\Delta E^0 = 1.23 \text{ V vs. SHE}$$

and



$$\Delta E^0 = 0.00 \text{ V vs. SHE}$$

An efficient photocatalyst should have the conduction band (CB) potential greater than the H^+/H_2 potential for the reduction reaction (eqn (4)) to occur. While the valence band (VB) energy level should be greater than the OH^-/O_2 potential for the oxidation reaction (eqn (3)) to occur. Thus, the bandgap criterion for photocatalysis corresponds to the minimal potential difference, which is 1.23 eV.⁵⁹ The conventional water redox potentials in terms of vacuum are 4.44 eV for reduction (H^+/H_2) and 5.67 eV for oxidation ($\text{H}_2\text{O}/\text{O}_2$).⁶⁰

The band edge position of the heterostructures of AAII-Se MoSSe and AAII-Se WSSe for water splitting is shown in Fig. 2.

As demonstrated in Fig. 3a, the band edge positions of the AAII-Se MoSSe and AAII-Se WSSe heterostructures do not straddle the water redox potentials at pH = 0. This means that these heterostructures fail to meet the thermodynamic criteria of an efficient photocatalyst for splitting water into hydrogen and oxygen molecules. However, AAII-Se MoSSe is a good photocatalyst for the oxygen evolution reaction as shown in the plot, and AAII-Se WSSe is as well if the band edge is slightly modified. As we know, strain engineering is a viable method for tuning the band gap and band-edge positions, as demonstrated by several research for various materials. In future studies, we would demonstrate the effect of strain on the conduction band edge to make this heterostructure possible for overall water splitting.

Overall, the good catalytic performance obtained can be attributed to the following: (a) several intrinsic features of the heterostructure remain unchanged because of the weak van der Waals interlayer interaction. As a result, the good catalytic performance of MoS₂ is transferred to some of these



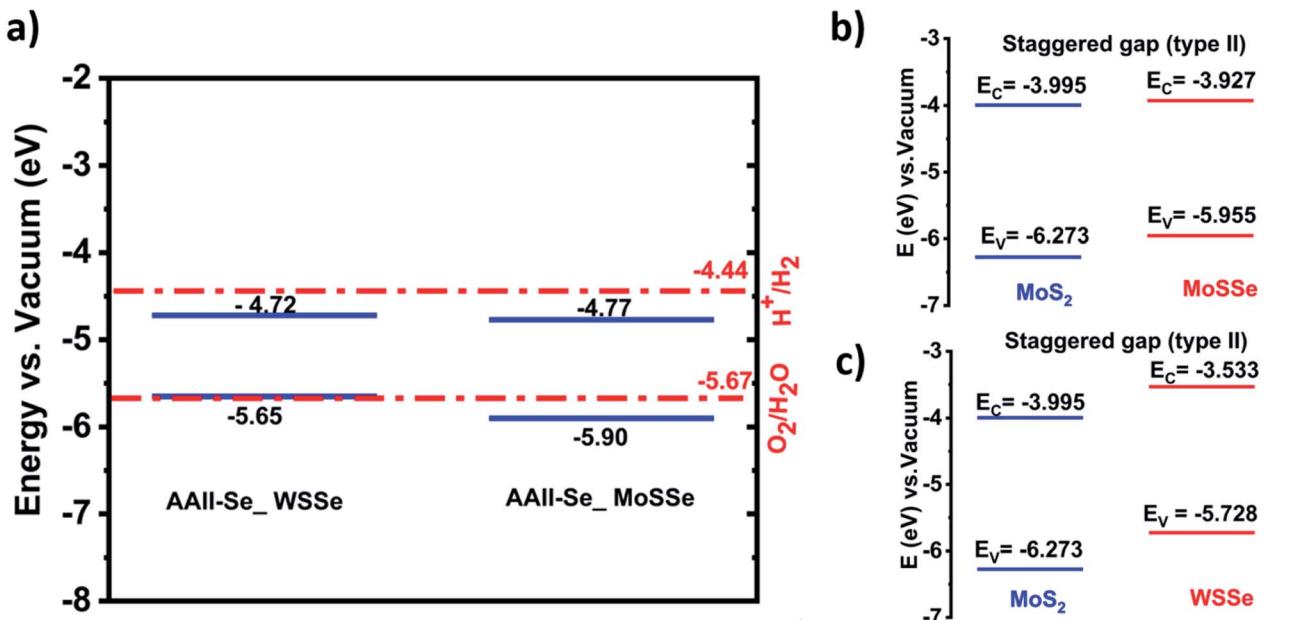


Fig. 3 (a) Band edge positions of the two heterojunctions for photocatalytic water splitting. The redox potential of water splitting at pH = 0 is shown by the red dashed line. (b) Schematic representation of type II band alignment AAII-Se MoSSe heterostructure. (c) Schematic representation of type II band alignment AAII-Se WSSe heterostructure.

heterostructures. (b) The heterostructure overcomes a notable drawback of a single MoS₂ monolayer, which is rapid charge recombination resulting in improved catalytic performance. The redox potential *versus* light absorption in the semiconductor is an irreconcilable contradiction for photocatalytic water-splitting. A small band gap is beneficial for the high light absorption, in contrast a large bandgap is essential for a high redox potential for water splitting. The balance is hard to be attain in single monolayers but can be addressed in vdW heterostructures. Thus, these vdW heterostructures combines layers with different properties such as lattice parameter, electronic band gaps, *etc.* resulting in exciting new characteristics. The combination of MoS₂ and Janus monolayers shows some of these attributes as discussed and evaluated in the current study.

Considering, the band gap alignment, we obtained that both structures exhibit type II alignment, as illustrated in Fig. 3b and c. This allows for efficient charge carrier separation and increases the photocatalytic water splitting activity. As indicated in the introduction, the short lifetimes of photogenerated electrons that easily recombine with holes jeopardize MoS₂'s suitability as a photocatalyst. This issue is overcome, however, by forming a type-II heterostructure with MoSSe and WSSe.

(d) Photovoltaic solar cells

Heterostructures based on atomically thin semiconductors are a promising developing technology for achieving ultrathin and lightweight photovoltaic solar cells on flexible substrates.⁶¹ The stable two heterostructure configurations considered have reasonable band gap and exhibit type-II band alignment, which is critical for photovoltaic energy conversion⁶² making them a good candidate for photovoltaic applications. Thus, the

performance of these heterostructure for photovoltaic applications is evaluated. To investigate the solar cell storage or conversion performance of these anticipated vdW heterostructured materials, we used the following equation⁶³ to calculate the maximum power conversion efficiency (PCE) (η). η which is expressed below is the most essential metric describing a photovoltaic device.

$$\eta = \frac{FF \cdot V_{oc} \cdot J_{sc}}{P_{sun}} \quad (5)$$

Using

$$eV_{oc} = E_g - E_{loss} \quad (6)$$

$$J_{sc} = e \int_{E_g}^{\infty} \frac{S(E)}{E} dE \quad (7)$$

$$P_{sun} = \int_0^{\infty} S(E) dE \quad (8)$$

where FF denotes the fill factor, V_{oc} denotes the maximum open-circuit voltage ($E_g - E_{loss}$), J_{sc} denotes the maximum short-circuit current density, and $S(E)$ is calculated using the NREL AM1.5 dataset. The bandgap energy is denoted by E_g . According to previous reports, the fill factors in 2D heterostructure photovoltaic structures are typically in the range of 0.3–0.5,⁶¹ so we take this value, and energy losses E_{loss} can be assumed to be 0.3.⁶⁴ Because a previous experimental research discovered a fill factor of 0.57 for a MoS₂/p-Si heterostructure.⁶⁵ We use this as our fill factor because it is the only experimental article that the authors are aware of that reports a fill factor for a TMD. Using the equation presented above, the computed PCE values for

AAII-Se MoSSe and AAII-Se WSSe are (7.289–12.15%) and (5.62–9.37%) respectively, which are much better than the results for MoS₂/p-Si (5.23%)⁶⁶ WS₂–MoS₂ (2.56%).⁶⁷

Conclusions

Using first-principles calculations, we studied the structural and electrical characteristics of the MoS₂/Janus TMDs van der Waals heterostructure. The most energetically stable stacking patterns of AAII-Se MoSSe and AAII-Se WSSe exhibit type II band gap alignment as well as indirect and direct band gap alignment. At pH = 0, the heterostructures have no acceptable band edge positions for overall water splitting but are good for the oxygen evolution reaction. Many investigations for various materials have demonstrated that strain engineering is a viable method for tuning the band gap and band-edge positions therefore there is a chance vary the position of the band edges and make the heterostructures active for overall water splitting applications. Furthermore, the power conversion efficiency of both heterostructures is comparable to and greater than the value published for similar 2D van der Waals heterostructures. Overall, the type-II heterostructure and high-power conversion efficiency imply that the vdW heterostructure we studied has a lot of promise for usage in photovoltaic devices and, with minor modifications, photocatalytic devices.

Conflicts of interest

There are no conflicts to declare.

Acknowledgements

This work was substantially sponsored by funding from Addis Ababa University through a thematic project (Grant No. RD/LT-095/2019). We also gratefully acknowledge Dr Tekalign T. Debela, Postdoctoral Associate at the University of Oregon, for providing insight and expertise that considerably aided the research. The authors thank the Centre for High Performance Computing (CHPC) in Cape Town, South Africa, for the computational resources used in this study. KOO acknowledges the HySA Infrastructure Centre of Competence, Faculty of Engineering, North-West University for their financial support.

References

- 1 G. A. Asres, A. Dombovari, T. Sipola, R. Puskás, A. Kukovecz, Z. Kónya, A. Popov, J. F. Lin, G. S. Lorite, M. Mohl, G. Toth, A. Lloyd Spetz and K. Kordas, *Sci. Rep.*, 2016, **6**, 1–7.
- 2 K. O. Obodo, C. N. M. Ouma, J. T. Obodo and M. Braun, *Phys. Chem. Chem. Phys.*, 2017, **19**, 19050–19057.
- 3 K. S. Novoselov, D. Jiang, F. Schedin, T. J. Booth, V. V Khotkevich, S. V Morozov and A. K. Geim, *Proc. Natl. Acad. Sci. U. S. A.*, 2005, **102**, 10451.
- 4 K. F. Mak, C. Lee, J. Hone, J. Shan and T. F. Heinz, *Phys. Rev. Lett.*, 2010, **105**, 2–5.
- 5 B. Radisavljevic, A. Radenovic, J. Brivio, V. Giacometti and A. Kis, *Nat. Nanotechnol.*, 2011, **6**, 147–150.
- 6 X. Wang, K. Maeda, A. Thomas, K. Takanabe, G. Xin, J. M. Carlsson, K. Domen and M. Antonietti, *Nat. Mater.*, 2009, **8**, 76–80.
- 7 X. Wang, K. Maeda, X. Chen, K. Takanabe, K. Domen, Y. Hou, X. Fu and M. Antonietti, *J. Am. Chem. Soc.*, 2009, **131**, 1680–1681.
- 8 M. Zhang, Q. Wu, F. Zhang, L. Chen, X. Jin, Y. Hu, Z. Zheng and H. Zhang, *Adv. Opt. Mater.*, 2019, **7**, 1–18.
- 9 S. Guo, Y. Zhang, Y. Ge, S. Zhang, H. Zeng and H. Zhang, *Adv. Mater.*, 2019, **31**, 1–19.
- 10 K. O. Obodo, C. N. M. Ouma, P. M. Modisha and D. Bessarabov, *Appl. Surf. Sci.*, 2020, **529**, 147186.
- 11 I. C. Onyia, S. O. Ezeonu, D. Bessarabov and K. O. Obodo, *Comput. Mater. Sci.*, 2021, **197**, 110613.
- 12 X. Jiang, A. V. Kuklin, A. Baev, Y. Ge, H. Ågren, H. Zhang and P. N. Prasad, *Phys. Rep.*, 2020, **848**, 1–58.
- 13 Z. Lin, B. R. Carvalho, E. Kahn, R. T. Lv, R. Rao, H. Terrones, M. A. Pimenta and M. Terrones, *2D Mater.*, 2016, **3**, 022002.
- 14 N. Lu, H. Guo, L. Li, J. Dai, L. Wang, W.-N. Mei, X. Wu and X. C. Zeng, *Nanoscale*, 2014, **6**, 2879–2886.
- 15 K. Koåmider and J. Fernández-Rossier, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2013, **87**, 2–5.
- 16 A. Vargas, F. Liu, C. Lane, D. Rubin, I. Bilgin, Z. Hennighausen, M. Decapua, A. Bansil and S. Kar, *Sci. Adv.*, 2017, **3**(7), DOI: [10.1126/sciadv.1601741](https://doi.org/10.1126/sciadv.1601741).
- 17 C. Tan, J. Chen, X. J. Wu and H. Zhang, *Nat. Rev. Mater.*, 2018, **3**, 1–13.
- 18 Y. Zhang, C. K. Lim, Z. Dai, G. Yu, J. W. Haus, H. Zhang and P. N. Prasad, *Phys. Rep.*, 2019, **795**, 1–51.
- 19 J. Low, J. Yu, M. Jaroniec, S. Wageh and A. A. Al-Ghamdi, *Adv. Mater.*, 2017, **29**, 1–20.
- 20 L. Ju, M. Bie, X. Zhang, *et al.*, *Front. Phys.*, 2017, **16**, 13201.
- 21 Y. Boyjoo, H. Sun, J. Liu, V. K. Pareek and S. Wang, *Chem. Eng. J.*, 2017, **310**, 537–559.
- 22 J. Zhang, G. Xiao, F. X. Xiao and B. Liu, *Mater. Chem. Front.*, 2017, **1**, 231–250.
- 23 C. Jiang, K. Y. Lee, C. M. A. Parlett, M. K. Bayazit, C. C. Lau, Q. Ruan, S. J. A. Moniz, A. F. Lee and J. Tang, *Appl. Catal. A*, 2016, **521**, 133–139.
- 24 Z. Jiang, Z. Y. Zhang, W. Shangguan, M. A. Isaacs, L. J. Durndell, C. M. A. Parlett and A. F. Lee, *Catal. Sci. Technol.*, 2016, **6**, 81–88.
- 25 S. Kumar, C. M. A. Parlett, M. A. Isaacs, D. V. Jowett, R. E. Douthwaite, M. C. R. Cockett and A. F. Lee, *Appl. Catal. B*, 2016, **189**, 226–232.
- 26 E. Benavente, F. Durán, C. Sotomayor-Torres and G. González, *J. Phys. Chem. Solids*, 2018, **113**, 119–124.
- 27 R. Kumar, D. Das and A. K. Singh, *J. Catal.*, 2018, **359**, 143–150.
- 28 E. Singh and H. S. Nalwa, *J. Nanosci. Nanotechnol.*, 2015, **15**, 6237–6278.
- 29 E. Singh and H. S. Nalwa, *RSC Adv.*, 2015, **5**, 73575–73600.
- 30 E. Singh and H. S. Nalwa, *Sci. Adv. Mater.*, 2015, **7**, 1863–1912.
- 31 R. Li, Y. Cheng and W. Huang, *Small*, 2018, **14**, 1802091.



32 Y. C. Cheng, Z. Y. Zhu, W. B. Mi, Z. B. Guo and U. Schwingenschlögl, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2013, **87**, 2.

33 K. O. Obodo, C. N. M. Ouma, J. T. Obodo, M. Braun and D. Bessarabov, *Comput. Condens. Matter*, 2019, **21**, e00419.

34 S. Wi, H. Kim, M. Chen, H. Nam, L. J. Guo, E. Meyhofer and X. Liang, *ACS Nano*, 2014, **8**, 5270–5281.

35 J. Zhang, M. Zhang, R. Q. Sun and X. Wang, *Angew. Chem., Int. Ed.*, 2012, **51**, 10145–10149.

36 V. O. Özçelik, J. G. Azadani, C. Yang, S. J. Koester and T. Low, *Phys. Rev. B*, 2016, **94**, 035125.

37 S. Wang, C. Ren, H. Tian, J. Yu and M. Sun, *Phys. Chem. Chem. Phys.*, 2018, **20**, 13394–13399.

38 M. Bernardi, M. Palummo and J. C. Grossman, *Nano Lett.*, 2013, **13**, 3664–3670.

39 F. Li, W. Wei, P. Zhao, B. Huang and Y. Dai, *J. Phys. Chem. Lett.*, 2017, **8**(23), 5959–5965, DOI: [10.1021/acs.jpclett.7b02841](https://doi.org/10.1021/acs.jpclett.7b02841).

40 Y. Ji, M. Yang, H. Lin, T. Hou, L. Wang, Y. Li and S. T. Lee, *J. Phys. Chem. C*, 2018, **122**, 3123–3129.

41 A. Y. Lu, H. Zhu, J. Xiao, C. P. Chuu, Y. Han, M. H. Chiu, C. C. Cheng, C. W. Yang, K. H. Wei, Y. Yang, Y. Wang, D. Sokaras, D. Nordlund, P. Yang, D. A. Muller, M. Y. Chou, X. Zhang and L. J. Li, *Nat. Nanotechnol.*, 2017, **12**, 744–749.

42 J. Zhang, S. Jia, I. Kholmanov, L. Dong, D. Er, W. Chen, H. Guo, Z. Jin, V. B. Shenoy, L. Shi and J. Lou, *ACS Nano*, 2017, **11**, 8192–8198.

43 P. E. Blöchl, *Phys. Rev. B*, 1994, **50**, 17953.

44 J. P. Perdew, K. Burke and M. Ernzerhof, *Phys. Rev. Lett.*, 1996, **77**(4), 3865.

45 B. Ghosh, S. Nahas, S. Bhowmick and A. Agarwal, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2015, **91**, 1–6.

46 B. J. Wang, X. H. Li, R. Zhao, X. L. Cai, W. Y. Yu, W. Bin Li, Z. S. Liu, L. W. Zhang and S. H. Ke, *J. Mater. Chem. A*, 2018, **6**, 8923–8929.

47 J. Heyd, G. E. Scuseria and M. Ernzerhof, *J. Chem. Phys.*, 2003, **118**, 8207–8215.

48 A. V. Krukau, O. A. Vydrov, A. F. Izmaylov and G. E. Scuseria, *J. Chem. Phys.*, 2006, **125**, 224106.

49 A. Miralrio, E. Rangel Cortes and M. Castro, *Appl. Surf. Sci.*, 2018, **455**, 758–770.

50 M. Farmanbar and G. Brocks, *Phys. Rev. B*, 2016, **93**, 1–12.

51 S. Ahmad, I. Ahmad, N. Van and B. Amin, *RSC Adv.*, 2020, **10**, 38114–38119.

52 M. Yagmurcukardes, C. Sevik and F. M. Peeters, *Phys. Rev. B*, 2019, **100**, 1–8.

53 X. Yang, D. Singh, Z. Xu, Z. Wang and R. Ahuja, *J. Mater. Chem. C*, 2019, **7**, 12312–12320.

54 H. G. Abbas, T. T. Debela, J. R. Hahn and H. S. Kang, *J. Phys. Chem. C*, 2021, **125**, 7458–7465.

55 T. V. Vu, N. V. Hieu, H. V. Phuc, N. N. Hieu, H. D. Bui, M. Idrees, B. Amin and C. V. Nguyen, *Appl. Surf. Sci.*, 2020, **507**, 145036.

56 T. A. Alrebdi and B. Amin, *Chem. Phys.*, 2021, **549**, 111252.

57 K. Maeda, K. Teramura, D. Lu, T. Takata, N. Saito, Y. Inoue and K. Domen, *Nature*, 2006, **440**, 295.

58 S. Y. Reece, J. A. Hamel, K. Sung, T. D. Jarvi, A. J. Esswein, J. J. H. Pijpers and D. G. Nocera, *Science*, 2011, **334**, 645–648.

59 E. L. Miller, *Energy Environ. Sci.*, 2015, **8**, 2809–2810.

60 J. Kang, S. Tongay, J. Zhou, J. Li and J. Wu, *Appl. Phys. Lett.*, 2013, **102**, 22–25.

61 M. M. Furchi, F. Höller, L. Dobusch, D. K. Polyushkin, S. Schuler and T. Mueller, *npj 2D Mater. Appl.*, 2018, **2**, 1–7.

62 J. Kang, H. Sahin and F. M. Peeters, *J. Phys. Chem. C*, 2015, **119**, 9580–9586.

63 M. R. Filip, C. Verdi and F. Giustino, *J. Phys. Chem. C*, 2015, **119**, 25209–25219.

64 H. J. Snaith, *Adv. Funct. Mater.*, 2010, **20**, 13–19.

65 S. K. Pradhan, B. Xiao and A. K. Pradhan, *Sol. Energy Mater. Sol. Cells*, 2016, **144**, 117–127.

66 M. L. Tsai, S. H. Su, J. K. Chang, D. S. Tsai, C. H. Chen, C. I. Wu, L. J. Li, L. J. Chen and J. H. He, *ACS Nano*, 2014, **8**, 8317–8322.

67 M. L. Tsai, M. Y. Li, J. R. D. Retamal, K. T. Lam, Y. C. Lin, K. Suenaga, L. J. Chen, G. Liang, L. J. Li and J. H. He, *Adv. Mater.*, 2017, **29**, 1–7.

68 K. F. Mak, C. Lee, J. Hone, J. Shan and T. F. Heinz, *Phys. Rev. Lett.*, 2010, **105**(13), 136805.

69 N. Bandaru, R. S. Kumar, D. Snead, O. Tschauner, J. Baker, D. Antonio, S. Luo, T. Hartmann, Y. Zhao and R. Venkat, *J. Phys. Chem. C*, 2014, **118**(6), 3230–3235.

70 Y. N. Wen, M. G. Xia and S. L. Zhang, *Comput. Mater. Sci.*, 2018, **152**, 20–27.

71 L. Kou, A. Du, C. Chen and T. Frauenheim, *Nanoscale*, 2014, **6**(10), 5156–5161.

72 Q. Yue, Z. Shao, S. Chang and J. Li, *Nanoscale Res. Lett.*, 2013, **8**(1), 1–7.

