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## Enhanced excitonic features in an anisotropic $\text{ReS}_2/\text{WSe}_2$ heterostructure†

Arslan Usman, <sup>a,f</sup> M. Adel Aly, <sup>a,b</sup> Hilary Masenda, <sup>a,c</sup> Joshua J. P. Thompson, <sup>a</sup> Surani M. Gunasekera, <sup>d</sup> Marcin Mucha-Kruczyński, <sup>d,e</sup> Samuel Brem, <sup>a</sup> Ermin Malic<sup>a</sup> and Martin Koch<sup>a</sup>

Two-dimensional (2D) semiconductors have opened new horizons for future optoelectronic applications through efficient light–matter and many-body interactions at quantum level. Anisotropic 2D materials like rhenium disulphide ( $\text{ReS}_2$ ) present a new class of materials with polarized excitonic resonances. Here, we demonstrate a  $\text{WSe}_2/\text{ReS}_2$  heterostructure which exhibits a significant photoluminescence quenching at room temperature as well as at low temperatures. This indicates an efficient charge transfer due to the electron–hole exchange interaction. The band alignment of two materials suggests that electrons optically injected into  $\text{WSe}_2$  are transferred to  $\text{ReS}_2$ . Polarization resolved luminescence measurements reveal two additional polarization-sensitive exciton peaks in  $\text{ReS}_2$  in addition to the two conventional exciton resonances  $X_1$  and  $X_2$ . Furthermore, for  $\text{ReS}_2$  we observe two charged excitons (trions) with binding energies of 18 meV and 15 meV, respectively. The bi-excitons of  $\text{WSe}_2$  become polarization sensitive and inherit polarizing properties from the underlying  $\text{ReS}_2$  layers, which act as patterned substrates for top layer. Overall, our findings provide a better understanding of optical signatures in 2D anisotropic materials.

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## 1. Introduction

Two-dimensional (2D) materials have gained immense research interest since the inception of graphene.<sup>1</sup> The 2D family consists of a wide variety of layered materials which include metals, insulators, and semiconductors. The layered semiconductors hold strong excitonic features<sup>2–6</sup> that can be tuned by external means by probing the physics of spin and valley polarization. Among all layered semiconductors, transition metal dichalcogenides (TMDs) exhibit exceptional optical, electrical, and topological properties when scaled down to monolayers. Most monolayers possess a direct bandgap in contrast to their bulk crystals. These ultrathin layers are promising materials for flexible, transparent, and robust electronic and optoelectronic devices such as photodetectors,<sup>7,8</sup> solar cells<sup>9,10</sup> and polarization sensitive

switches.<sup>11</sup> Large exciton binding energies ( $>100$  meV)<sup>12</sup> enabling excitons to overcome thermal dissociation at room-temperature are a consequence of enhanced quantum confinement and reduced dielectric screening.<sup>13</sup> Another field of research that is emerging rapidly in the last few years is the formation of heterostructures<sup>14–18</sup> by stacking monolayers of two different TMDs. These structures possess an atomically sharp and nearly defect-free interface.<sup>15</sup> The stack represents a hybrid material arranged in a desired sequence.<sup>19,20</sup> In such heterostructures interlayer excitons are formed because of the band alignment and an efficient charge transfer across the heterostructure interface.<sup>21,22</sup> Such charge transfer excitons have characteristic features like long lifetimes and spin-valley polarization.<sup>23</sup> Thus, the domain of two-dimensional materials and van der Waals heterostructures opens new avenues leading to a fast-growing research area in semiconductor technology. Typically, monolayers absorb about 10% of the incoming radiation in the visible range.

The absorption gradually increases with an increasing number of layers and easily reaches  $\sim 30\%$  for thicker flakes. Following this concept, heterostructures with more than two monolayers show a considerable photon absorption. Consequently, the increased concentration of photo generated charge carriers leads to enhanced quantum efficiency that is crucial for photovoltaic and optoelectronic applications.

Rhenium disulphide ( $\text{ReS}_2$ ) is an anisotropic layered material from the VII-B group which has a layer independent

<sup>a</sup>Department of Physics and Materials Sciences Centre, Philipps-Universität Marburg, Marburg 35032, Germany. E-mail: arslan.usman@physik.uni-marburg.de

<sup>b</sup>Department of Physics, Faculty of Science, Ain Shams University, Cairo 11566, Egypt

<sup>c</sup>School of Physics, University of the Witwatersrand, 2050 Johannesburg, South Africa

<sup>d</sup>Department of Physics, University of Bath, Claverton Down, Bath BA2 7AY, UK

<sup>e</sup>Centre for Nanoscience and Nanotechnology, University of Bath, Claverton Down, Bath BA2 7AY, UK

<sup>f</sup>Department of Physics, COMSATS University Islamabad-Lahore-Campus, Pakistan

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bandgap due to its weak interlayer coupling and low crystal symmetry. It has a distorted 1T structure with Re-Re metallic atoms in a zig-zag chain-like geometry which is maintained *via* covalent bonds. This lack of ordering is the overriding reason for low interlayer coupling of  $\text{ReS}_2$ .<sup>24</sup> Importantly, with respect to its *b*-axis (metallic atomic chains) it shows a polarization sensitive optical excitation and emission. The fabrication of  $\text{ReS}_2$  based heterostructure and its integration in optoelectronic applications<sup>25–27</sup> has not yet been explored extensively.

Here, we study the photoluminescence from a 2D layered heterostructure composed of a few-layers (FL) of rhenium disulphide covered with a monolayer of  $\text{WSe}_2$ . The low-temperature luminescence spectrum shows numerous excitonic peaks with a distinct polarization dependence which includes charged trions and polarized bi-excitons that have not been reported yet. We have also investigated the exciton polarizability using polarization resolved PL. From the pronounced luminescence we conclude on an efficient charge transfer due to the distinct type II band alignment.<sup>28</sup> The efficient charge transfer between  $\text{ReS}_2$  and  $\text{WSe}_2$  could be a promising basis for the application of such heterostructures in optoelectronic devices.

## 2. Methods and characterizations

A monolayer (ML) of  $\text{WSe}_2$  and a few layer (FL) flake of  $\text{ReS}_2$  were exfoliated from bulk crystals (2D Semiconductors, Inc.) on polydimethylsiloxane (PDMS) gel. A flake (few layers) of hexagonal boron nitride (hBN) was transferred using the standard viscoelastic stamping technique onto a 300 nm  $\text{SiO}_2/\text{Si}$  substrate. The FL- $\text{ReS}_2$  flake was then transferred on top of the hBN which acts as substrate. Secondly, a glass slide containing a monolayer of  $\text{WSe}_2$  on PDMS is well aligned and transferred on top of the  $\text{ReS}_2$  flake ensuring a good overlay for the heterostructure. The optical micrograph of the heterostructure is illustrated in Fig. 1(a), while panel (b) illustrates a schematic of the prepared heterostructure. The entire transfer process is done under a microscope using a manual controlled  $xyz\theta$  stage with a 20× objective. During the transfer process before moving to the next step, each material on PDMS gel was

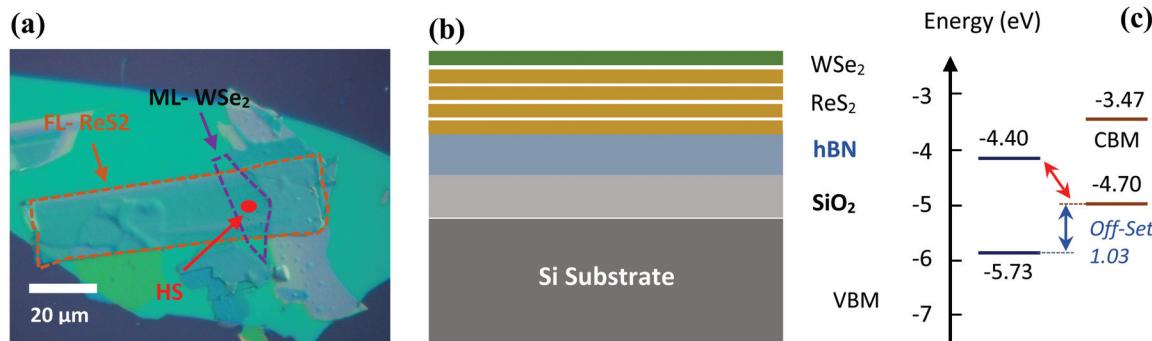
characterized using Raman and photoluminescence (PL) spectroscopy. The heterostructure is also confirmed using atomic force microscopy, Fig. S2 (ESI).<sup>†</sup> Raman and PL spectra from each individual monolayer as well as from the overlap region confirms the proper material stacking and consequently reveals good interface quality. The crystallographic orientation was estimated by identifying the *b*-axis, as  $\text{ReS}_2$  is most likely to be cleaved along that direction (ESI, Fig. S2<sup>†</sup>).

## 3. Results and discussion

### 3.1 Photoluminescence

In this work, a steady state  $\mu$ -PL setup is used to obtain the spectra of the two individual materials as well as their heterostructure. In the setup, which comprises a 40× objective, the sample is excited with a CW laser ( $\lambda = 532$  nm). The emission signal is collected by the same objective and is analyzed by monochromator with a liquid nitrogen-cooled CCD detector (Princeton Instruments). The obtained spectra are shown in Fig. 2.  $\text{WSe}_2$  is a direct gap semiconductor at monolayer thickness with an emission peak at 1.64 eV as shown in Fig. 2(a), whereas few-layer  $\text{ReS}_2$  has a direct bandgap at 1.55 eV. The orange curve in the PL spectrum has one dominant peak at 1.64 eV, which we attribute to the neutral exciton ( $\text{X}^0$ ) of  $\text{WSe}_2$ . In addition to that, the peak has a tail on the low energy side which suggests the presence of a weak additional peak redshifted by  $\sim$ 30 meV which originates from the negatively charged exciton ( $\text{X}^-$ ). This peak has already been observed for  $\text{WSe}_2$  in previous studies.<sup>29–32</sup>

The blue curve represents the emission from few-layer  $\text{ReS}_2$ . The broad peak at 1.56 eV with a shoulder at 1.52 eV represents the two neutral excitons  $\text{X}_2$  and  $\text{X}_1$ , respectively. These two characteristic polarization sensitive excitons are not resolvable at room temperature. There is an additional peak on the high-energy side at 1.63 eV, which is attributed to an indirect emission from dark states.<sup>33</sup> In the heterostructure when the  $\text{WSe}_2$  monolayer is stacked over  $\text{ReS}_2$  the reduced dielectric screening causes a redshift of 5 meV of the neutral exciton ( $\text{X}^0$ ) peak. The luminescence spectrum of the heterostructure is



**Fig. 1**  $\text{WSe}_2/\text{ReS}_2$  vertical heterostructure, (a). Optical micrograph with individual and overlay regions, (b). Schematic of vertical heterostructure, (c). Energy level diagram.

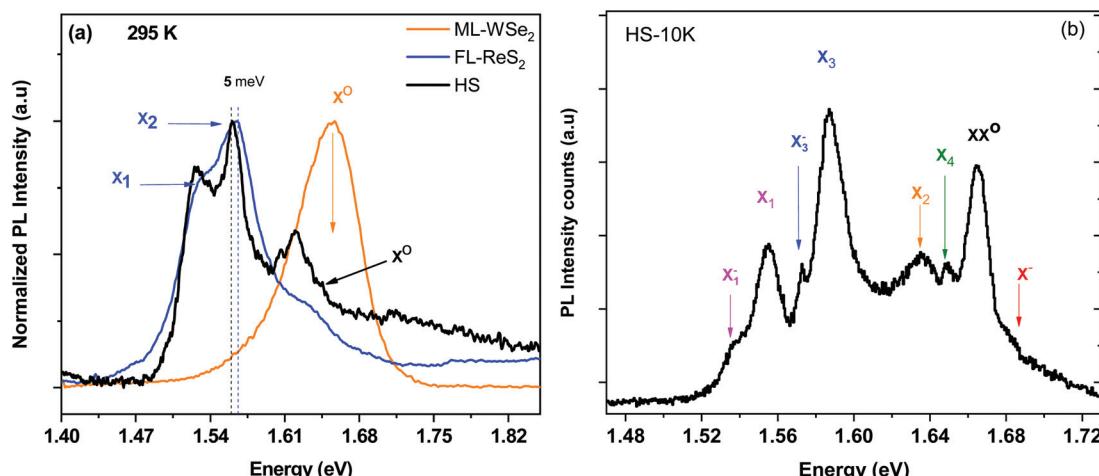


Fig. 2 Photoluminescence spectra of monolayer WSe<sub>2</sub>, FL-ReS<sub>2</sub> & their heterostructure at (a). 295 K (b). 10 K.

shown as a black curve. It shows an enhanced luminescence from WSe<sub>2</sub> bi-excitons (XX<sup>o</sup>). This originates from a localization of charge carriers in WSe<sub>2</sub> arising from directional screening of the underlaying ReS<sub>2</sub>. The origin of the peak has been concluded from its respective position. It is also confirmed by a power dependent measurement (see Fig. 3). The ReS<sub>2</sub> acts as a patterned substrate that influences the isotropic nature of WSe<sub>2</sub>.

Therefore, the charge species act as polarized excitons and trions. The confinement due to anisotropic screening field pronounced the mixed states (bright and semi-dark) in biexcitons due to electron-electron intervalley scattering as reported earlier in.<sup>34</sup> It is further observed in this heterostructure that the PL intensity of the neutral exciton (X<sup>o</sup>) of WSe<sub>2</sub> is significantly quenched. This quenching effect arises from the efficient charge transfer across the heterojunction (electrons to ReS<sub>2</sub>) which in turn originates from the type-II band alignment, schematically shown in Fig. 1(c). The type II band alignment for this material system was already predicted by DFT

calculations.<sup>35,36</sup> It can be further elucidated from the spectra that the enhanced XX<sup>o</sup> is due to the change of the electron mobility under local confinement. The band off-set in the conduction band minimum (CBM) is 0.93 eV and is, hence, larger than the thermal vibrational energy of ~0.02 meV (phonon mode) at room temperature. The band off-set in the valence band (VBM) is ~1.03 eV. The emission from the interlayer charge transfer exciton should appear at an energy which is beyond our detector's limit. At the room temperature, the excitonic properties of single-layer and few-layer TMDCs are largely influenced by phonon interactions. To obtain sharp and well-resolved excitonic features we perform PL measurements at low temperature (~10 K). The excitation light was circularly polarized to obtain uniform excitation in all possible directions. The PL spectra shown in Fig. 2(b) reveal eight emission peaks corresponding to different excitonic features of the heterostructure. The spectra exhibit two higher energy peaks at 1.66 eV and 1.68 eV attributed to the bi-exciton (XX<sup>o</sup>) and the charged exciton (X<sup>-</sup>) of WSe<sub>2</sub>, respectively.

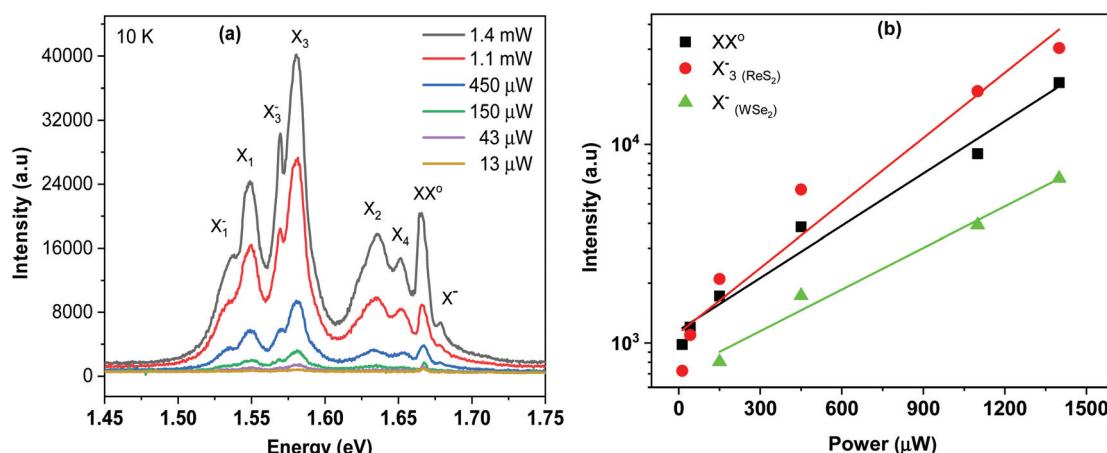


Fig. 3 (a) Power dependent PL of the heterostructure at 10 K. (b) Semi-log plot of the observed intensities as a function of the excitation power.



At low temperatures, the localized states thermalize more easily than excitons and trions and a significant amount of charge carriers are trapped in these localized states. This has also been reported earlier.<sup>37</sup> At the same time the directional screening from the underlying substrate also forces these carriers to remain in these localized states. There are four exciton peaks at 1.55, 1.58, 1.63 and 1.64 eV related to  $\text{ReS}_2$ , designated as  $X_1$ ,  $X_3$ ,  $X_2$  and  $X_4$ , respectively. The additional peaks  $X_3$  and  $X_4$  originate from the splitting of excitonic states due to the electron-hole exchange interactions.<sup>38</sup> Structural anisotropy resulting from the broken rotational symmetry also contributes to this spin-orbital splitting.

Here, we report for the first time the observation of negatively charged trions of  $\text{ReS}_2$  ( $X_1^-$  and  $X_3^-$ ) in an ungated structure, shown in Fig. 2(b). These newly observed trion peaks are referred to as  $X_1^-$  and  $X_3^-$  and appear at 1.536 eV and 1.572 eV, respectively. They are negatively charged due to the charge transfer-induced electron accumulation in  $\text{ReS}_2$ . The trion peaks are located approximately 18 meV and 15 meV below their respective neutral exciton peaks corresponding to the trion binding energies. These values agree well with previous observations.<sup>39</sup> The efficient transfer of charge carriers across the heterostructure is also the reason for the suppression of the neutral exciton ( $X^0$ ) in  $\text{WSe}_2$ . The emission intensity of the  $\text{WSe}_2$  bi-exciton is much higher than that of both the trion ( $X^-$ ) and the neutral exciton ( $X^0$ ), which we attribute to the directional dielectric-screening-induced exciton localization arising from the underlaying  $\text{ReS}_2$ . This peculiarity has not been observed in previous studies on an isolated  $\text{WSe}_2$  monolayer.<sup>40–42</sup> The enhanced emission is due to the change of screening to a directional environment in this heterostructure.

To clarify the nature of the individual PL peaks power-dependent photoluminescence measurements are performed. The spectra in Fig. 3 along with a plot of the corresponding intensities confirms the assignment of the peaks to biexciton  $XX^0$  and trion  $X^-$  of  $\text{WSe}_2$  which shows almost linear behaviour with reference to excitation power. The charged trion ( $X_3^-$ ) of  $\text{ReS}_2$  shows nearly quadratic trend. These bi-excitons evolve from two identical neutral excitons with a Bohr radius exactly equal to that of a single exciton ( $\sim 1$  nm) and separated by  $\sim 3$ –4 nm.<sup>43</sup> Under suitable conditions, these four-body quasi-particles would thus be a strong channel for light emission. Such many body excitonic correlated states favor new quantum properties such as exciton condensates. It could be interesting to produce coherent photons by utilizing these correlated excitons.<sup>44</sup> Furthermore, low-temperature PL spectra show two further excitons, designated as  $X_2$  and  $X_4$ , that are close to each other (separated by  $\sim 12$  meV). They have a weaker oscillator strength and therefore disappear at higher temperatures.

### 3.2 Polarization-resolved Photoluminescence

The  $\text{ReS}_2$  crystal is anisotropic and therefore the PL shows a pronounced polarization dependence. The two neutral polarized excitons in  $\text{ReS}_2$ , namely  $X_1$  and  $X_2$  have been reported

earlier in a single layer and in few-layer samples.<sup>45–47</sup> Since, we have a heterostructure we observe here the additional excitons  $X_3$  and  $X_4$ . In the following, we investigate their anisotropic nature using an additional polarizer in front of the detector. Fig. 4(a) shows the polarized PL emission as a function of the polarizer angle. Another representation is the 2D contour plot shown in Fig. 4(b). In this plot the dependence of exciton and trion resonances can clearly be seen for a complete  $360^\circ$  turn of the analyser.

We also study the degree of polarization of the light from the heterostructure by luminescence measurements employing right-handed and left-handed circular polarized excitation and kept the detection angle fixed, so that the orthogonal helicities are obtained as  $(\sigma^+)$  and  $(\sigma^-)$ . Fig. 4(c) shows that the excitons  $X_1$  and  $X_3$  are more sensitive to  $(\sigma^+)$  and  $(\sigma^-)$  polarizations, respectively. The exciton and trion of same orientation are very much consistent with one type of circular polarization. One possible reason would be partial excitation of electron-hole pairs in valleys of dark states  $X_D$  of localized charge carriers in the respective exciton  $X_1$  or  $X_3$ . The variation in the degree of polarization can also be attributed to a variable oscillator strength and its electron-phonon coupling strength. The comparative low emission signal from  $X_2$  and  $X_4$  excitons is expected because of lower oscillator strength compared to  $X_1$  and  $X_3$  excitons as revealed in the temperature dependent PL spectra. The polar plots for the prominent excitonic features of the heterostructure are presented in Fig. 5 as a function of detection angle  $\theta$ . The exciton  $X_1$ ( $X_3$ ) as well as the negative trion  $X_1^-$ ( $X_3^-$ ) emission exhibits a double-lobed shape, oriented at  $63^\circ$  with respect to each other, corresponding to the angle between  $a$  and  $b$ -axis of  $\text{ReS}_2$ . This polarization resolved peaks emerge from the splitting of singlet and triplet states; the broken rotational symmetry and structural anisotropy are the reasons for such splitting. The excitonic signal stems primarily from non-degenerate direct excitons.<sup>45</sup> The corresponding trions also exhibit a strongly linearly polarized emission of the same orientation. This behaviour suggests the spread of the anisotropic trion wavefunction, which is primarily elongated along and perpendicular to the  $b$ -axis for the  $X_1^-$ ( $X_3^-$ ) trion.

The  $\text{WSe}_2$  localized excitons also show a linear polarization as shown in Fig. 5(c), these excitonic species are well aligned with the underlying directional screening field which is consistent with previous results.<sup>47</sup>

To illustrate the anisotropic nature of the excitonic wavefunction calculations have been performed (*cf.* ESI†), the results of which are presented in Fig. 6. In an isotropic material, the wave function of the 1s exciton is circular, whereas in this case it is distinctly elliptical. Unlike conventional TMDs (Mo, W), the rhenium TMDs possess inversion symmetry, which leads to a spin degeneracy of the valence and conduction bands. Previous works have suggested that this gives rise to the strong exchange splitting of the 1s-exciton peak of the  $\text{ReS}_2$  monolayer. While this calculation is beyond the scope of our analytical methods, we can estimate the exchange splitting by referring to the literature.<sup>48</sup>



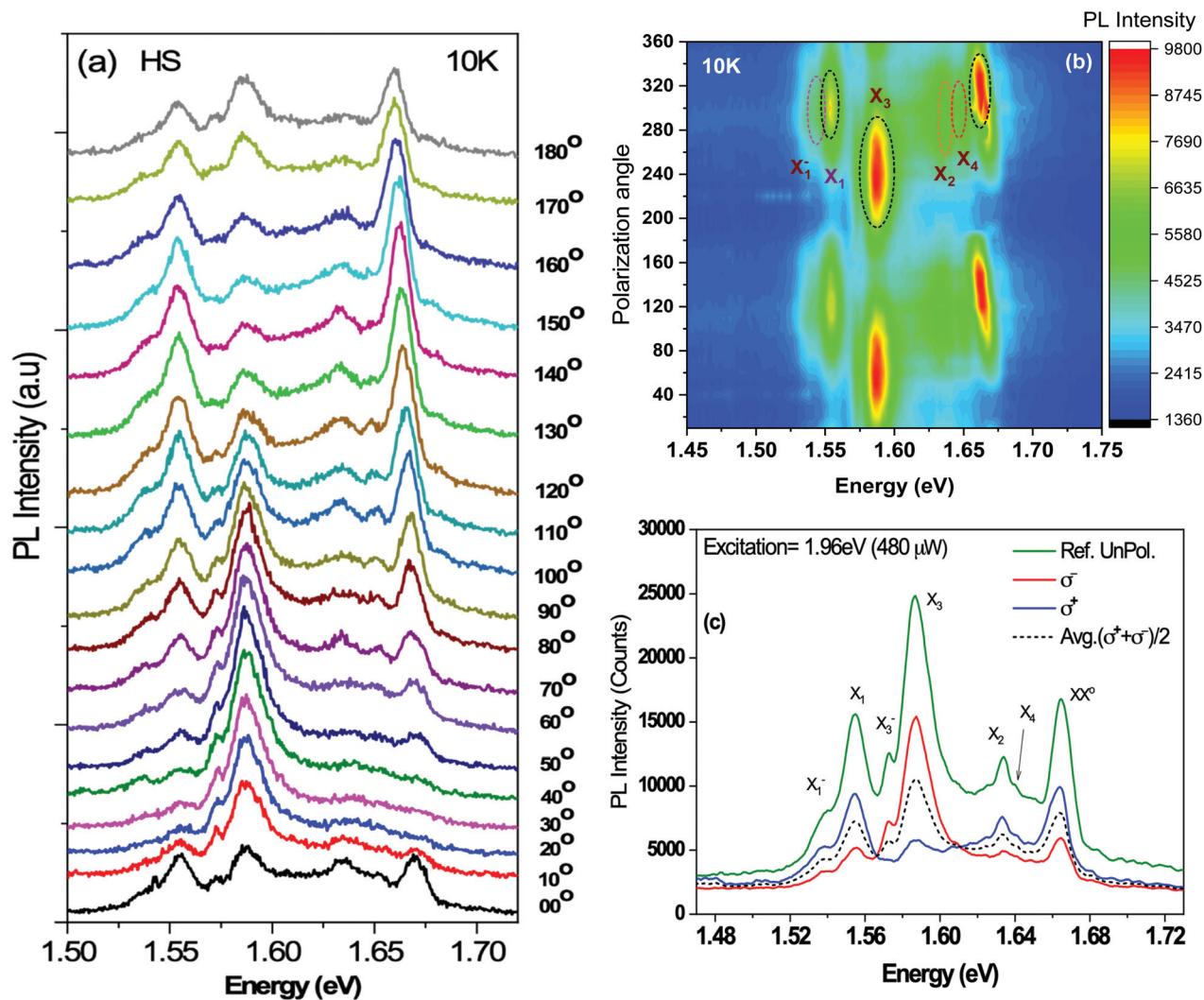


Fig. 4 (a). Polarization-resolved PL spectra of the WSe<sub>2</sub>/ReS<sub>2</sub> heterostructure at 10 K. The detection angle varies between 0° to 180° (b). 2D contour plot of PL spectra with polarization angle ( $\theta$ ). (c). The right- and left-handed circular polarization PL spectra.

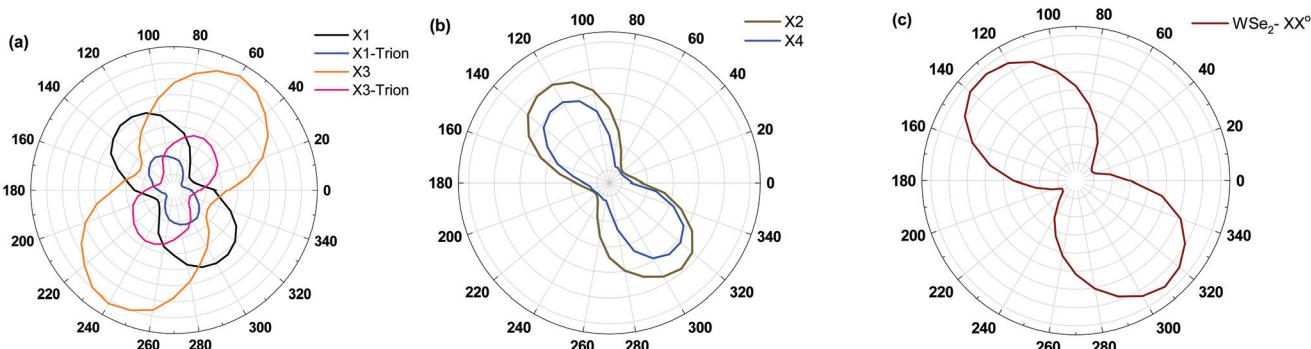


Fig. 5 Polar plot of integrated PL intensities for the WSe<sub>2</sub>/ReS<sub>2</sub> heterostructure at 10 K. (a) Intensity of the PL emission from the exciton and trion of ReS<sub>2</sub> (X<sub>1</sub> and X<sub>3</sub>, respectively) as a function of the analyzer angle. (b) Intensity of the PL emission from the X<sub>2</sub> and X<sub>4</sub> excitons of ReS<sub>2</sub>. (c) Intensity of the PL emission from bi-excitons in WSe<sub>2</sub> (XX<sup>0</sup>).



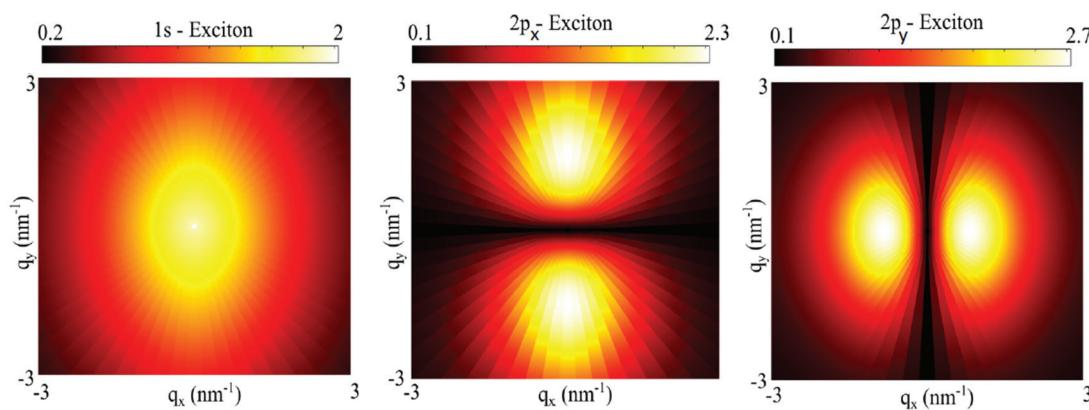


Fig. 6 Square modulus of excitonic wave functions  $|\varphi_{q_x, q_y}^n|^2$  for  $n = 0, 1, 2$ , for 1s and 2p excitons.

### 3.3 Temperature dependent Photoluminescence

In order to study the role of phonons, temperature dependent photoluminescence measurements of few layer  $\text{ReS}_2$ , monolayer  $\text{WSe}_2$  and their heterostructure were performed. The results are presented in Fig. 7. The FL- $\text{ReS}_2$  shows the commonly observed exciton  $X_1$  at an energy of 1.59 eV with a shoulder at 1.55 eV. In addition, a broad peak is observed at 1.66 eV with an anomalous variation of the intensity as the

temperature is increased. This behaviour is attributed to low-lying dark states  $X_D$ <sup>49</sup> recently reported in transition metal dichalcogenides.<sup>50-53</sup> As can be seen from Fig. 7(a), the 10 K PL spectrum of the  $\text{WSe}_2$  monolayer shows well resolved characteristic features of exciton, trion and bi-exciton at 1.72 eV, 1.70 eV and 1.68 eV, respectively.

It can be elucidated from the heterostructure PL spectra that the emission arising from bi-excitons of  $\text{WSe}_2$  has been substantially enhanced as compared to the monolayer. The

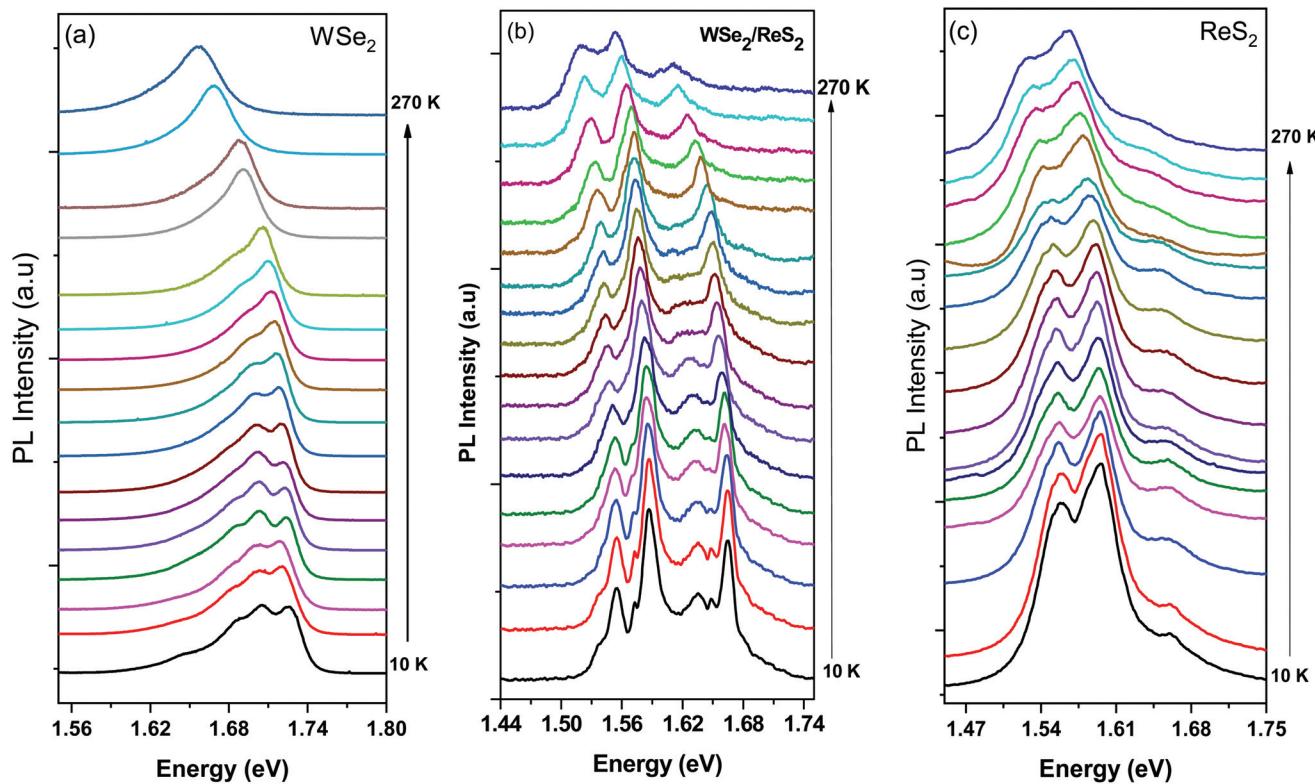


Fig. 7 PL spectra measured at different temperatures with a fixed angle of the analyzer, (a).  $\text{WSe}_2$  (b). Heterostructure of  $\text{WSe}_2/\text{ReS}_2$  (c). Few layers of  $\text{ReS}_2$ .



ReS<sub>2</sub> acts as substrate for WSe<sub>2</sub> monolayer and its screening field confines the carrier within WSe<sub>2</sub>, as a result of this confinement the phenomenon of mixed states (bright and semi-dark) is more pronounced in biexcitons due to electron–electron intervalley scattering as reported earlier.<sup>34</sup> The lifetime of biexcitons is bit longer than neutral exciton and trion<sup>54</sup> which are also responsible for higher quantum yield. Therefore, the confinement due to screening will significantly influence more on biexciton in this heterostructure. Furthermore, it is redshifted by 15 meV. This could be due to further confinement of localized charges because of the screening field underneath. At low temperature these localized states are easy to thermalize than excitons and trions and significant amount of charge carriers are captured in those localized states.<sup>37</sup>

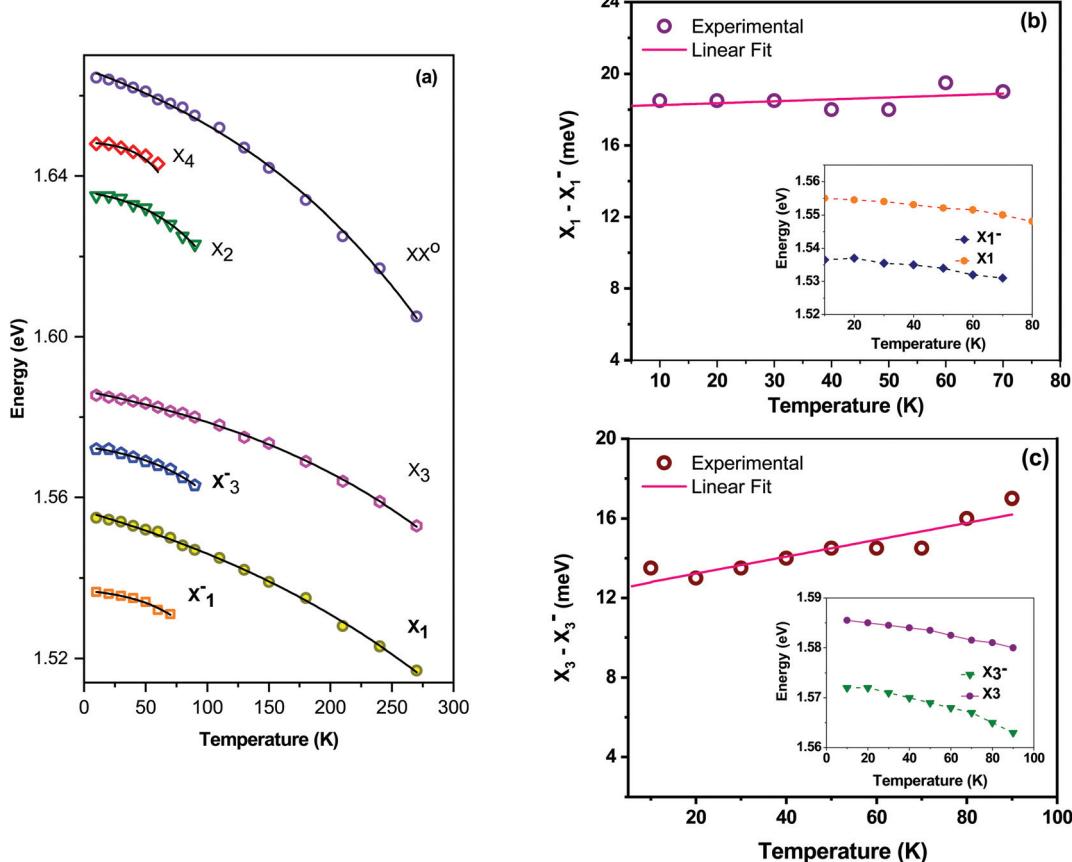
The X<sub>1</sub> exciton which is observed at high temperature splits into two excitons with a narrow linewidth, namely X<sub>1</sub> and X<sub>3</sub> at low temperature. In addition, one observes corresponding trions as X<sub>1</sub><sup>−</sup> and X<sub>3</sub><sup>−</sup>. The increased electron injection from WSe<sub>2</sub> to ReS<sub>2</sub> because of charge transfer as well as the encapsulation of ReS<sub>2</sub> layers in hBN and WSe<sub>2</sub> are the prominent reasons for the emergence of these trions and the narrowing of exciton linewidths. Furthermore, one observes two additional peaks at 1.633 eV and 1.65 eV which we attribute to the X<sub>2</sub> and X<sub>4</sub> excitons of ReS<sub>2</sub>, respectively. These two peaks

are only resolvable in the temperature range between 10 K and 70 K. At higher temperatures these excitons are thermally populated by absorbing phonons and scattered out of the light cone, as explained by the rate equation model.<sup>33</sup> Such phonon modes have been recently observed for ReS<sub>2</sub>.<sup>55</sup> Hence, for the heterostructure we observe a temperature dependent redshift of the peak positions of all exciton species.

We observed Varshni shift for energies with increasing temperature which has been traditionally applied for the bandgap shrinkage in semiconductors;<sup>56</sup>

$$E_g(T) = E_g(0) - S\hbar\omega \left[ \coth\left(\frac{\hbar\omega}{2kT}\right) - 1 \right]$$

Here,  $E_g(0)$  is the resonance energy of the exciton,  $\hbar\omega$  is the average phonon energy and  $S$  is the electron–phonon coupling constant.  $E_g(0)$ ,  $S$  and  $\hbar\omega$  are the fitted parameters. Fig. 8(a) reveals that the energy separation between the peaks X<sub>1</sub> and X<sub>3</sub> of ReS<sub>2</sub> is almost constant with temperature. In contrast, the separation between the peaks X<sub>2</sub> and X<sub>4</sub> is converging with temperature. Fig. 8(b and c) shows the temperature dependent peak positions of the trion X<sub>1</sub><sup>−</sup> and X<sub>3</sub><sup>−</sup> and their corresponding neutral excitons. The insets show the original peak positions, the main figures show the energy difference between exciton and trion. The trion X<sub>3</sub><sup>−</sup> shows a larger energy



**Fig. 8** (a). Temperature dependent shift of the excitonic peak positions of the heterostructure. Energy separation between exciton and trion for (b). X<sub>1</sub> (c). X<sub>3</sub>.



shift than  $X_1^-$  because it lies in a plane having different dipole moment which is more aligned towards *b*-axis ( $\sim 25^\circ$ ). The linearly polarized emission of the trion in a quantum confined quasi-1D system like this suggests an anisotropic trion wavefunction<sup>57</sup> as discussed earlier. This also enhances the Coulomb interaction between charge carriers.<sup>58-60</sup>

### 3.4 Exciton binding energy in fractional dimension

To get an insight into the exciton and trion binding energies for FL-ReS<sub>2</sub> with variable dielectric constant, a theoretical analysis is implemented. Numerous techniques are available to estimate the exciton and trion binding energies in quasi-2D materials:<sup>60-62</sup> *e.g.* the diffusion Monte Carlo method<sup>63</sup> and the boundary matching matrix method.<sup>60,63</sup> Following these approaches, we can obtain an estimate for the binding energies. The binding energies of exciton ( $E_{\text{ex}}$ ) and trion ( $E_{\text{T}}$ ) were calculated using the following equations.<sup>65</sup>

$$E_{\text{ex}} = \frac{4}{(D-1)^2} \times R_y^*$$

$$E_{\text{T}} = \left( -\frac{4}{(D-1)^2} - c_0 - \sum_{i=1}^4 c_i D^{-i} e^{-D} \right) \times R_y^*$$

Here  $D_{\text{ex}}$  and  $D_{\text{T}}$  are the effective dimension of the exciton and trion, respectively.  $\epsilon$  is the dielectric coefficient and  $R_y^*$  is the effective Rydberg.<sup>64</sup>

$$R_y^* = \frac{13\,600}{\epsilon^2} \frac{\mu}{m_e} \text{ meV}$$

The effective dimension of the trion is slightly smaller than that of the exciton. To keep the calculation simple, we approximate  $D = D_{\text{ex}} = D_{\text{T}}$ . The coefficients  $c_{i=1-4}$ <sup>65</sup> are based on a Monte Carlo study of charged excitons and are described in the ESI (Table S1†). Employing these equations and parameters, we estimated exciton and trion binding energies as a function of fractional dimension for a unique dielectric constant. The result of this calculation is presented in Fig. 9. Based on two dielectric constants of ReS<sub>2</sub> along and perpendicular to the *b*-axis<sup>64,66</sup> the exciton and trion wavefunction is elongated in two directions. The excitons which are along *b*-axis have higher energy as compared to excitons which are confined in a perpendicular direction. A similar behaviour can also be seen for the corresponding trions. The obtained exciton binding energy in ReS<sub>2</sub> is slightly higher than reported for other TMDs.<sup>67,68</sup> This also supports its quasi-1D nature and agrees with the already reported binding energy.<sup>69</sup> It can also be seen that the binding energies of exciton and trion increase with the decrease of the fractional dimension  $D$ .

To draw a comparison of our results obtained on samples with substrates to the exciton binding energy which one would expect for an isolated layer of ReS<sub>2</sub>, we theoretically calculated the binding energies of excitonic states of a free-standing monolayer using the Wannier equation. The calculated binding energy for the 1s exciton is 859 meV (see ESI Fig. S3†).

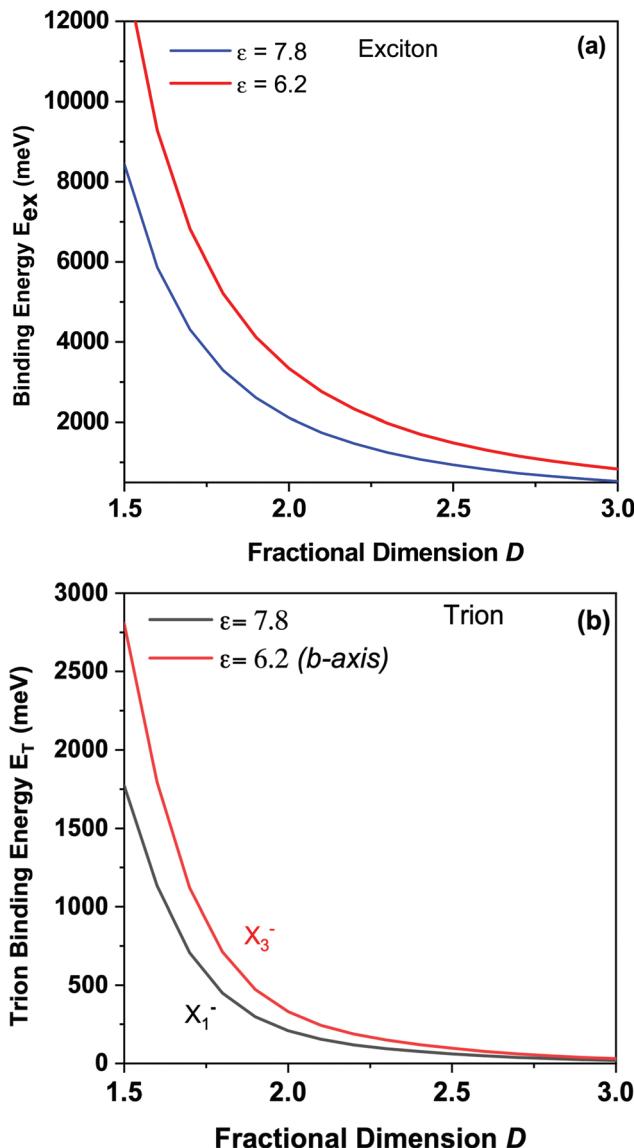


Fig. 9 The calculated binding energy for ReS<sub>2</sub> as a function of effective dimension  $D$  for (a). Exciton (b). Trion.

## 4. Conclusions

We have successfully demonstrated the charge transfer exciton in an ungated WSe<sub>2</sub>/ReS<sub>2</sub> heterostructure. We have also observed trions which have thus far only been observed experimentally in a gated structure. Two additional peaks are evident in the luminescence spectrum of ReS<sub>2</sub>, they are accompanied by two luminescence maxima which we attribute to spin degenerate excitons of ReS<sub>2</sub>. Polarization-resolved spectra reveal distinct orientations of these excitons and of trions which are elongated in the same orientation as the excitons. The distorted 1T structure of anisotropic ReS<sub>2</sub> favours the directional screening due to the existence of variable dielectric constants. The high exciton and trion binding energies also confirm their strongly confined quasi-1D nature. Moreover,



ReS<sub>2</sub> acts as an optically patterned substrate for the top WSe<sub>2</sub> monolayer. Consequently, the WSe<sub>2</sub> excitons show a polarized behaviour which differs from their intrinsic properties. Our observations provide new insight into the many-body physics in optically excited states of anisotropic van der Waals systems which are promising materials for optoelectronic devices, including polarization sensitive switches and spin manipulated devices.

## Conflicts of interest

There are no conflicts to declare.

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