

Cite this: *Mater. Adv.*, 2022,  
3, 6343

# Pulsed laser deposition for conformal growth of MoS<sub>2</sub> on GaN nanorods for highly efficient self-powered photodetection†

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Nanostructure-based photodetectors (PDs) have recently attracted intensive attention due to their immense potential in high performance next-generation photonic devices. Herein, we report a self-powered, ultrafast and broadband PD based on conformally grown MoS<sub>2</sub> on GaN nanorods (NRs). The higher aspect-ratio and extensively exposed edges of conformal 2D materials provide a large interface for better optical absorption and light harvesting. The device has been realized by the deposition of MoS<sub>2</sub> by pulsed laser deposition on GaN NRs/Si grown via molecular beam epitaxy. The MoS<sub>2</sub>/GaN/Si-based 1D–2D conformal PD exhibits a high photoresponse in a broad range of wavelengths (300–1000 nm) in the self-powered mode. The maximum responsivity of the PD is found to be ~14.22 A W<sup>-1</sup> with ultrafast rise and fall times of 38.8 and 8.2 μs, respectively. Our findings demonstrate that conformal PDs based on the MoS<sub>2</sub>/GaN NR heterojunction hold promising potential for applications in the field of optoelectronics.

Received 23rd May 2022,  
Accepted 23rd June 2022

DOI: 10.1039/d2ma00577h

rsc.li/materials-advances

## 1 Introduction

Over the last few years, photodetectors (PDs) have become an integral component of most of the semiconductor-based systems as they find use in various fields such as defense and missile technology, communication systems, environmental and industrial monitoring, and so on.<sup>1,2</sup> A critical parameter associated with PDs is their energy consumption.<sup>3</sup> They generally require an external energy source for their operation and this adds to the overall cost of the device system.<sup>3</sup> Thus, the present research on PDs emphasizes self-powered PDs and photovoltaics,<sup>4–7</sup> and various strategies have been employed to achieve self-powered detection. These include utilizing devices based on p–n junctions, the photothermoelectric effect, asymmetric contacts in linear configurations, heterojunctions having a large difference in the carrier concentrations of the constituent semiconductors, and so on. Recent advancements in the study of two-dimensional (2D) materials, especially

molybdenum disulfide (MoS<sub>2</sub>),<sup>8–10</sup> have shown promising results for self-powered optoelectronic systems because of their outstanding properties such as large carrier mobility, high light absorption properties, high current on/off ratio, the absence of dangling bonds at the surface, *etc.*<sup>11–17</sup> Moreover, heterostructures of MoS<sub>2</sub> with the established conventional three-dimensional (3D) semiconductors such as gallium nitride (GaN) offer the utilization of both the unique properties of MoS<sub>2</sub> as well as the matured technology of III-nitrides.<sup>9,18</sup> GaN is a very important wide band gap semiconductor which finds its applications in high speed and power devices, and is currently used in defense technology, wireless communication systems, radars, ultraviolet PDs, and electronic warfare systems.<sup>19</sup> Hence, integration of MoS<sub>2</sub> with GaN can be very interesting because of the complimentary nature of the electrical and optical properties of these two material systems, in terms of their band gaps and electronic mobilities. In addition, such integrations can open up a window to boost the device's responsivity as well as wavelength detection range, to achieve a broadband spectral response.

To date, intensive research has been done on the growth of MoS<sub>2</sub> on planar GaN thin films for photodetection. Huang *et al.*<sup>20</sup> have shown an n-MoS<sub>2</sub>/p-GaN heterostructure, exhibiting a responsivity of ~105 A W<sup>-1</sup> at 532 nm under applied reverse bias. Jain *et al.*<sup>21</sup> have reported on the current transport in MoS<sub>2</sub> grown on GaN and AlGaN thin films. Zhang *et al.*<sup>22</sup> have reported the design of a two-dimensional MoS<sub>2</sub>/GaN van der Waals (vdW) heterostructure PD consisting of few-layer p-type MoS<sub>2</sub> and very

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† Electronic supplementary information (ESI) available: Schematic of the device, Raman spectra and glancing angle XRD of MoS<sub>2</sub>, morphology of MoS<sub>2</sub>/GaN NRs, variation of sensitivity and gain with power density of the incident light for the MoS<sub>2</sub>/GaN NRs/Si based device, band gap estimation of GaN, energy band diagrams of MoS<sub>2</sub>/GaN/Si and MoS<sub>2</sub>/Si heterojunctions, responsivity as a function of wavelength for the MoS<sub>2</sub>/GaN NRs/Si-based device in two different configurations: top–top contact and top–bottom contact, along with a comparison with the GaN NRs/Si-based device. See DOI: <https://doi.org/10.1039/d2ma00577h>



thin n-type GaN flakes. Zhuo *et al.*<sup>9</sup> have demonstrated a deep ultraviolet (UV) PD based on an MoS<sub>2</sub>/GaN heterojunction, showing a responsivity of 187 mA W<sup>-1</sup> in the self-biased mode. Recently, Singh *et al.*<sup>23</sup> have reported an MoS<sub>2</sub>/GaN/Si-based PD showing a broadband photoresponse, with an ultrahigh responsivity in the near infrared region (23.81 A W<sup>-1</sup> at 995 nm) under self-powered conditions. However, the responsivity in the UV-visible region is very low with the reversal of charge transport. Overall, it can be seen that most of these devices suffer from either narrowband detection or low responsivities under self-powered conditions. To overcome such limitations, considerable study is being done on one-dimensional (1D) nanostructures such as nanowires and nanorods (NRs) because of their exceptional chemical and physical properties, such as higher surface area and so on.<sup>24–26</sup> As of now, only a few reports exist where MoS<sub>2</sub> has been grown with an alternative structure, that is, its vertical growth or growth along 1D nanostructures.<sup>27,28</sup> Cai *et al.*<sup>29</sup> have reported the photovoltaic conversion and the interfacial charge transfer mechanism in C<sub>60</sub>/MoS<sub>2</sub> 0D–2D mixed-dimensional vdW heterostructures based on atomic-bond-relaxation consideration and first-principles calculations. Compared to monolayer MoS<sub>2</sub>, it is observed that the C<sub>60</sub>/MoS<sub>2</sub> heterostructures show a type II band alignment along with lower interface reflectivity, thereby leading to an improvement in the power conversion efficiency. Jariwala *et al.*<sup>30</sup> have presented a concise and critical review of emerging mixed-dimensional vdW heterostructure devices. Liu *et al.*<sup>31</sup> have also reviewed the progress, challenges as well as the opportunities of vdW materials integration in their perspective. In the case of NRs, the channel width is very low, which leads to easier charge transport to metal electrodes, resulting in ultrafast photodetection. Moreover, due to the higher aspect-ratio and extensively exposed surfaces of the NRs, heterostructures of 1D GaN with MoS<sub>2</sub> can lead to better device performance, offering great potential in various nanoscale electronic and optoelectronic devices. Thus, it is very crucial to explore devices based on the conformal growth of MoS<sub>2</sub> over GaN NRs, which can be promising for photodetection.

In this article, we demonstrate the fabrication of a MoS<sub>2</sub>/GaN NRs/Si 1D–2D heterostructure and report the photoreponse behavior of the device based on this heterostructure. Interestingly, the device exhibits broadband photodetection (300–1000 nm) at 0 V with a maximum responsivity of ~14.22 A W<sup>-1</sup> at a wavelength of 900 nm for low light intensity (0.075 mW cm<sup>-2</sup>). The electrical characteristics of the device show an excellent photoresponse, with an ultrafast response/recovery time of 38.8/8.2 μs. The mechanism of the self-powered photoresponse is explained on the basis of the energy band diagrams of the heterojunction. The present work shows the importance of conformal heterostructures for the fabrication of highly efficient PDs.

## 2 Materials and methods

### Synthesis and characterization of the MoS<sub>2</sub>/GaN NRs/Si heterostructure

GaN NRs were grown on Si(111) using Omicron Nanotechnology molecular beam epitaxy (MBE) system. Prior to deposition, the Si

substrate (1 cm<sup>2</sup> in area) was cleaned using a standard cleaning procedure reported elsewhere.<sup>32</sup> The substrate was further cleaned thermally inside the MBE chamber at 1000 °C for 30 min. After thermal cleaning, GaN NRs were grown at a substrate temperature of 860 °C for 3 h. Ultra-high purity gallium was evaporated from a standard Knudsen effusion cell and the temperature and beam equivalent pressure were maintained at 960 °C and 1.22 × 10<sup>-7</sup> mbar, respectively. During the growth process, the nitrogen pressure was kept at 1.2 × 10<sup>-4</sup> mbar with a flow rate of 2 sccm, whereas the forward RF-power of the plasma source was maintained at 450 W.

MoS<sub>2</sub> was further deposited by pulsed laser deposition (PLD) on the MBE-grown GaN NRs/Si using a physical mask of 1000 μm diameter. The laser energy was kept at 200 mJ, resulting in a laser fluence of ~1 J cm<sup>-2</sup>. The number of laser shots was fixed to 1000. The details regarding the preparation of the MoS<sub>2</sub> target as well as the other PLD parameters for synthesis of MoS<sub>2</sub> have been given in our previous report.<sup>33</sup> For structural examination of the prepared heterojunction, μ-Raman spectroscopy and high-resolution X-ray photoelectron spectroscopy (HRXPS) were carried out using a LabRAM HR (532 nm laser) and a Thermo Scientific K-Alpha XPS (monochromatic X-ray beam: Al target), respectively. Scanning electron microscopy (SEM) was done using an Ultra55 FE-SEM Karl Zeiss EDS. Cross-sectional transmission electron microscopy (TEM) and energy dispersive spectroscopy (EDS) were performed using a 300 kV Titan Themis from FEI. TEM lamella preparation was done using a PIPS II: Precision Ion Polishing System from Gatan. High resolution X-ray diffraction (HRXRD) measurements were done using a diffractometer (manufactured by Bruker (Discover D8)) equipped with a Cu-Kα target of monochromatic wavelength (λ = 1.5404 Å) operating at 40 kV/40 mA.

### Photodetection and electrical measurements

The electrical and photodetection measurements of the device were performed as described elsewhere.<sup>33</sup> Al (100 nm) and Cr/Au (10 nm/100 nm) circular metal electrodes (diameter – 600 μm each) were deposited by thermal evaporation at the bottom (Si) of the device and by e-beam evaporation at the top (MoS<sub>2</sub>) of the device, respectively. Hence, the effective area of the device becomes 0.502 mm<sup>2</sup>. The schematic of the MoS<sub>2</sub>/GaN NRs/Si-based device is illustrated in Fig. S1 of the ESI.†

## 3 Results and discussion

To confirm the formation of the 2H phase of MoS<sub>2</sub>, the μ-Raman spectrum and glancing angle HRXRD pattern are shown in Fig. S2 of the ESI.† The characteristic Raman modes of MoS<sub>2</sub> (A<sub>1g</sub> and E<sub>12g</sub>) observed at 407.52 and 381.68 cm<sup>-1</sup>,<sup>10,33</sup> along with the XRD peak<sup>28</sup> at ~14.1°, confirm the presence of the 2H phase of MoS<sub>2</sub>. To get an insight into the growth morphology of MoS<sub>2</sub> on GaN NRs, scanning electron microscopy (SEM) and transmission electron microscopy (TEM) were done. Fig. 1(a and b) show the top surface and cross-sectional morphology of the GaN NRs



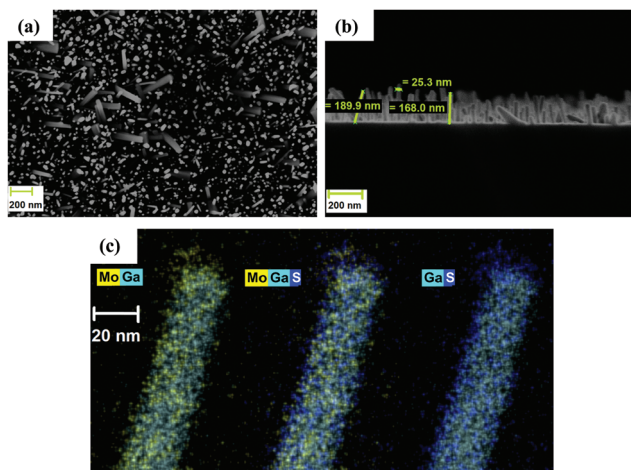


Fig. 1 SEM images of GaN NRs on Si showing the (a) top surface and (b) cross-sectional view. (c) Cross-sectional TEM-EDS micrographs of MoS<sub>2</sub>/GaN NRs, indicating conformal growth of MoS<sub>2</sub> around the GaN NRs.

acquired through SEM. It is evident from this figure that the growth density of the NRs is medium with the Si substrate exposed in between NRs, and the average length of the NRs is  $\sim 180$  nm. The bird's-eye view SEM image of MoS<sub>2</sub>/GaN NRs is

shown in Fig. S3 of the ESI.† Fig. 1(c) shows the cross-sectional TEM-energy dispersive spectroscopy (EDS) micrographs of the MoS<sub>2</sub>/GaN NRs/Si heterostructure. Overall, these results confirm the conformal growth of MoS<sub>2</sub> on the GaN NRs.

For the performance evaluation of the MoS<sub>2</sub>/GaN NRs/Si-based PD, spectral response studies were conducted in the wavelength range of 300–1000 nm ( $0.075 \text{ mW cm}^{-2}$ ) under self-powered conditions and the results are shown in Fig. 2(a). The response is observed in the entire wavelength range, and the maximum response is obtained at 900 nm. Responsivity starts to increase from the UV region to the visible region, and becomes maximum in the NIR region. Due to the non-uniform and thin coating of MoS<sub>2</sub> on GaN NRs as well as Si, a considerable amount of UV and visible light is absorbed by GaN and Si. Hence, a significant photoresponse is obtained in the UV-visible region, in contrast to the PD based on a thicker MoS<sub>2</sub> film on GaN thin film/Si as demonstrated in our previous report.<sup>23</sup> In the NIR region, the response becomes maximum around 900 nm as the absorbed photons lead to band-to-band transitions in MoS<sub>2</sub> and Si, following which the photon energy becomes insufficient to excite the electrons present in the valence band of MoS<sub>2</sub> and Si. Thus, the responsivity follows the behavior as observed in Fig. 2(a).  $I$ - $V$  characteristics in the

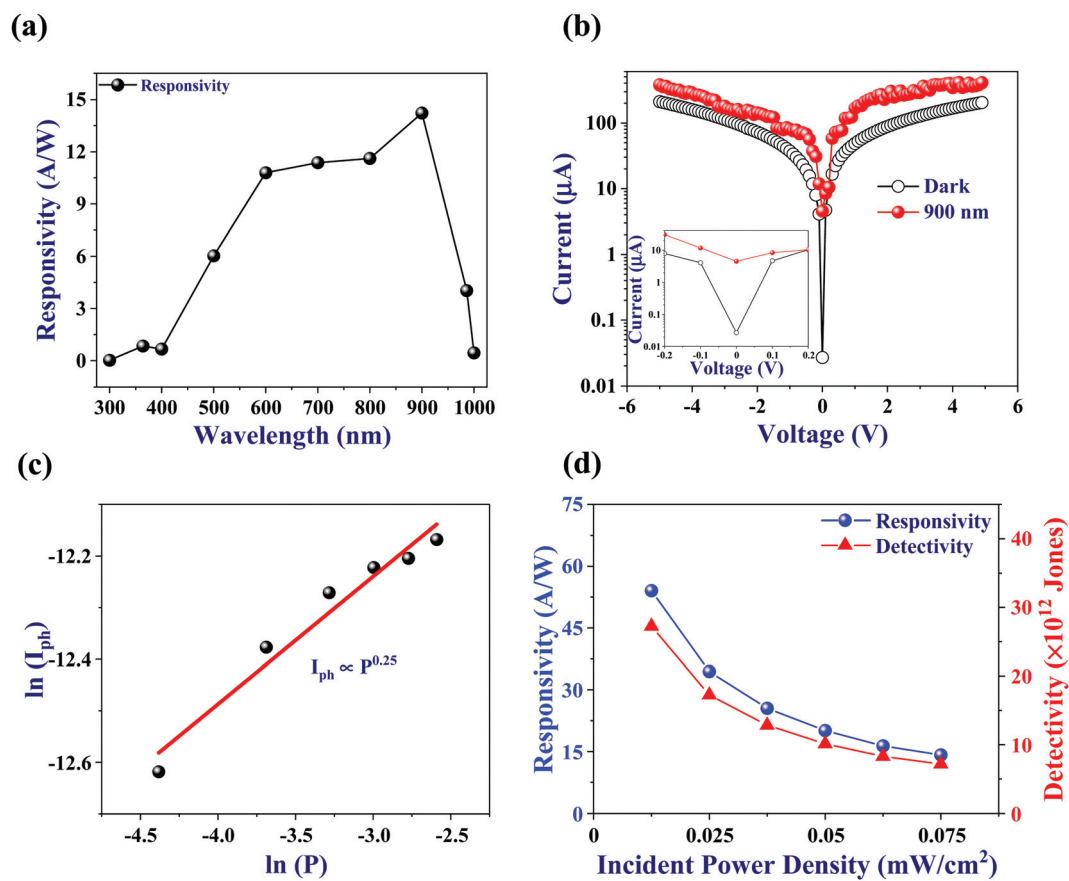


Fig. 2 (a) Spectral response, (b) dark and illuminated  $I$ - $V$  characteristics of the device (the inset shows a zoomed view of the  $I$ - $V$  characteristics near zero bias), (c) variation in photocurrent with different power intensities of irradiated light, and (d) variation in responsivity and detectivity with different power densities of irradiated light for the MoS<sub>2</sub>/GaN NRs/Si-based device.



dark and under light illumination ( $\lambda = 900$  nm) are presented in Fig. 2(b). The maximum photocurrent is seen at zero bias (inset of Fig. 2(b)), which shows the self-powered behavior of the PD. The variation in the photocurrent with light intensity was studied and the  $\log(\text{photocurrent})$  vs.  $\log(\text{power density})$  plot is shown in Fig. 2(c). Linear fitting of the curve shows a nonlinear dependence of photocurrent on the power density of the incident light ( $I_{\text{ph}} \propto P^\theta$ ,  $\theta = 0.25$ ). This non-unity value of  $\theta$  suggests that complex processes are involved in the generation of charge carriers *i.e.*, trapping and recombination.<sup>34,35</sup> Fig. 2(d) depicts the variation of the responsivity and detectivity with illumination power density. Responsivity decreases with an increase in the power density of the incident light source, which can be related to the enhanced recombination or scattering rate of photogenerated charge carriers, because their concentration increases at higher power densities. Just like responsivity, detectivity follows a similar trend and decreases with increasing power density, as the two quantities are directly proportional to each other, and the dark current is the same for different incident power densities. The variation of sensitivity and gain with incident power density is shown in Fig. S4 (ESI<sup>†</sup>). The temporal response of the device is measured at 0 V (Fig. 3(a)) and a reproducible and stable photoresponse is observed. From the  $I$ - $t$  plot, the time interval between the two sampling points comes out to be  $\sim 100$  ms, which is limited by the experimental setup. Hence, to estimate the rise and fall time constants with precision and accuracy, an ultrafast time measurement setup has been used (Fig. 3(b)), and the details of

this setup have been described elsewhere.<sup>33</sup> Fig. 3(c–e) depict the transient response of the detector at different chopper frequencies (1, 5 and 10 kHz), and the response/recovery times of the device were estimated by fitting a single response–recovery cycle with given first-order exponential equations. The estimated response/recovery times are as follows: 1 kHz – 164.9/93.8  $\mu\text{s}$ , 5 kHz – 57.2/16.9  $\mu\text{s}$ , and 10 kHz – 38.8/8.2  $\mu\text{s}$ .

It is seen that the response time of the PD is slower than the recovery time. As observed from the power law, the response of a semiconductor to incident photons is a complex process. Therefore, during the photogeneration of carriers, traps and surface states might be playing a significant role, leading to a slower response. In contrast, when the illumination is switched off, a fast charge carrier recombination leads to a quick recovery. A similar trend has been seen in previously reported NR-based PDs.<sup>36</sup>

The probable working mechanism of the  $\text{MoS}_2/\text{GaN}$  NRs/Si-based PD is explained with the help of the energy band diagrams of the heterojunction. The band gaps of  $\text{MoS}_2$  and Si have been taken to be  $\sim 1.34$  and  $\sim 1.11$  eV, respectively, from our previous reports.<sup>10,33</sup> The band gap of GaN has been estimated from diffuse reflectance spectroscopy, and is found to be  $\sim 3.37$  eV (Fig. S5 of the ESI<sup>†</sup>). In this case, two different heterojunctions,  $\text{MoS}_2/\text{GaN}/\text{Si}$  and  $\text{MoS}_2/\text{Si}$ , come into the picture as  $\text{MoS}_2$  is also directly deposited on Si. Using Hall measurements for  $\text{MoS}_2/\text{Si}$  and the Si substrate, the carrier concentrations ( $N_D$ ) of  $\text{MoS}_2$  and Si have been estimated to be  $\sim 4.3 \times 10^{16}$  and  $\sim 3.1 \times 10^{12}$   $\text{cm}^{-3}$ , respectively. For

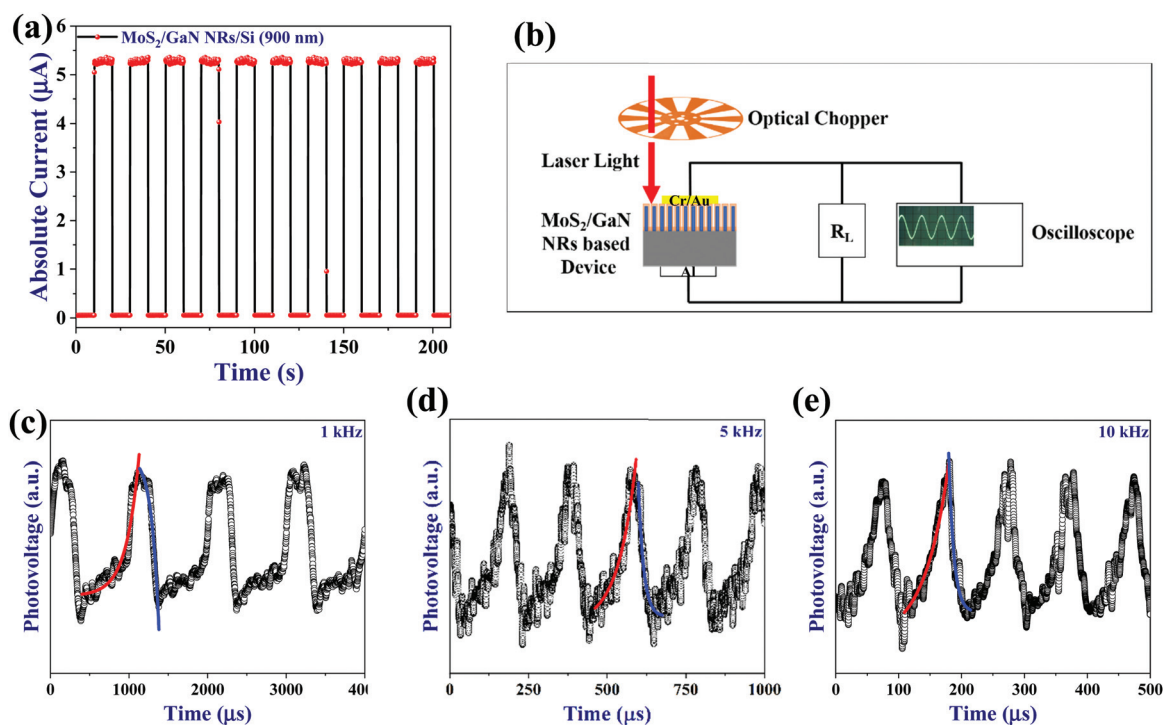


Fig. 3 (a) Temporal response of the  $\text{MoS}_2/\text{GaN}$  NRs/Si device at 0 V under 900 nm light illumination, (b) schematic of the ultrafast time measurement setup. Temporal response of the  $\text{MoS}_2/\text{GaN}$  NRs/Si device at a frequency of (c) 1 kHz, (d) 5 kHz and (e) 10 kHz measured with the ultrafast measurement setup.



estimating the carrier concentration of GaN, we have taken a GaN thin film sample and the carrier concentration is observed to be  $\sim 1.8 \times 10^{19} \text{ cm}^{-3}$ . All three samples (MoS<sub>2</sub>, GaN and Si) show unintentional n-type behavior, and the values of  $E_C - E_F$  are calculated for all three semiconductors using eqn (1):<sup>37</sup>

$$E_C - E_F = k_B T \ln \left( \frac{N_C}{N_D} \right) \quad (1)$$

where  $k_B$ ,  $T$ ,  $E_C$ ,  $E_F$  and  $N_C$  are the Boltzmann constant, temperature, conduction band minimum, Fermi level and the number of states in the conduction band, respectively. Moreover, to determine the band alignments at these two heterojunctions, HRXPS measurements were conducted. HRXPS is the most reliable, direct and convenient method for the evaluation of the valence band offset (VBO) and conduction band offset

(CBO) of a given heterojunction interface. The VBO ( $\Delta E_V$ ) and CBO ( $\Delta E_C$ ) values for the MoS<sub>2</sub>/GaN heterointerface have been estimated using the equations<sup>38</sup> given below:

$$\Delta E_V = \Delta E_{\text{MoS}_2/\text{GaN}}^{\text{MoS}_2} + \Delta E_{\text{GaN}2p_{3/2}-\text{MoS}_2}^{\text{MoS}_2/\text{GaN}} - \Delta E_{\text{GaN}2p_{3/2}-\text{VBM}}^{\text{GaN}} \quad (2)$$

$$\Delta E_C = E_G^{\text{MoS}_2} + \Delta E_V - E_G^{\text{GaN}} \quad (3)$$

where the first term ( $\Delta E_{\text{MoS}_2/\text{GaN}}^{\text{MoS}_2}$ ) of eqn (2) corresponds to the difference in the energy of the Mo core level with respect to its valence band maximum (VBM) and is found to be 228.44 eV (Fig. 4(a)). The subsequent term ( $\Delta E_{\text{GaN}2p_{3/2}-\text{MoS}_2}^{\text{MoS}_2/\text{GaN}}$ ) of eqn (2) refers to the energy difference of the Ga and Mo core levels, and

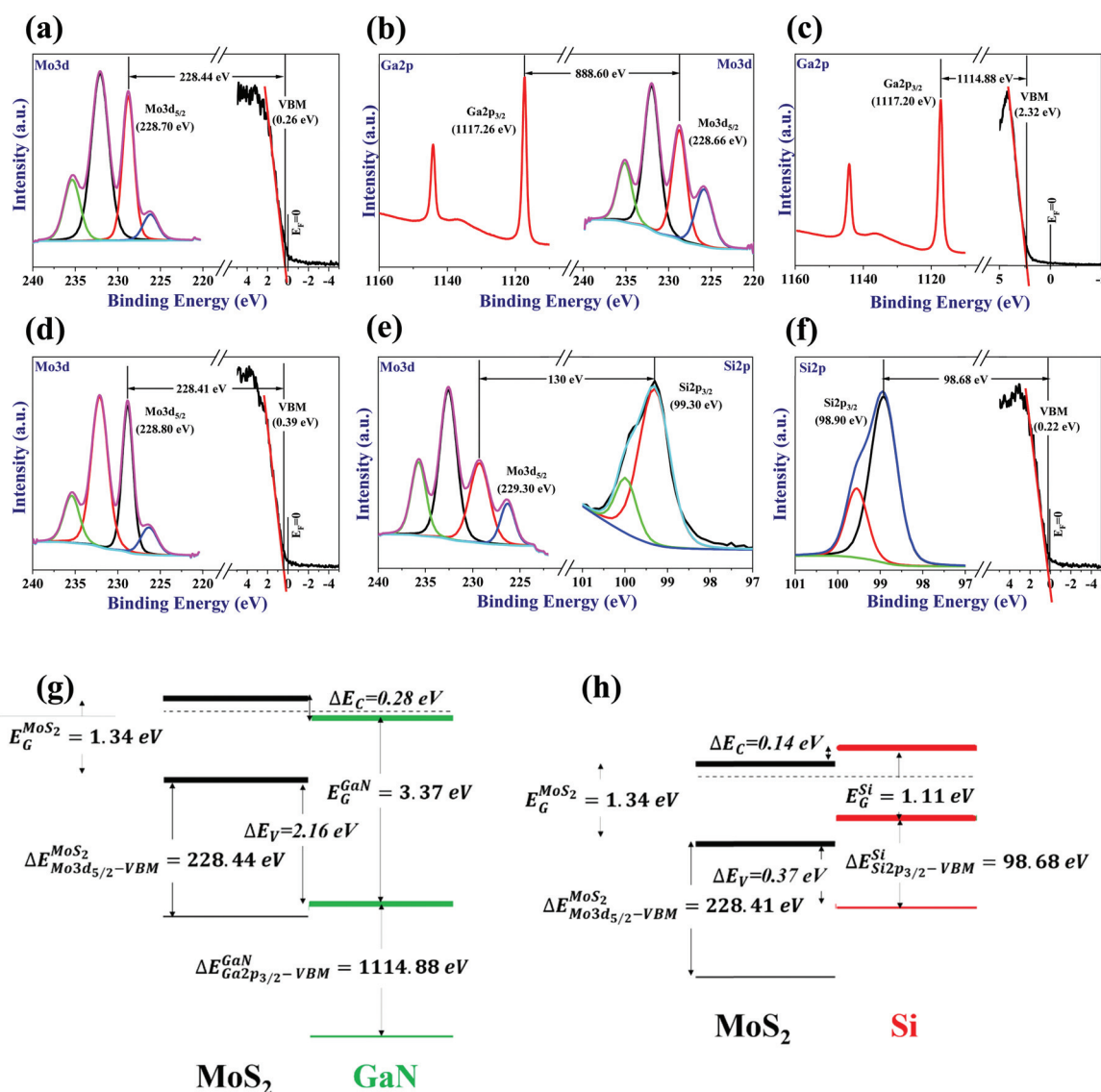


Fig. 4 For MoS<sub>2</sub>/GaN, (a) Mo3d core level and valence band spectra (VBS) acquired from MoS<sub>2</sub>, (b) Ga2p and Mo3d core-level spectra of the MoS<sub>2</sub>/GaN heterojunction, and (c) Ga2p core level and VBS of the GaN epilayer. For MoS<sub>2</sub>/Si, (d) Mo3d core level and VBS acquired from MoS<sub>2</sub>, (e) Mo3d and Si2p core-level spectra of the MoS<sub>2</sub>/Si heterojunction, and (f) Si2p core level and VBS of the Si substrate. Band alignment at (g) MoS<sub>2</sub>/GaN and (h) MoS<sub>2</sub>/Si heterojunctions.



is estimated to be 888.60 eV (Fig. 4(b)). The last term ( $\Delta E_{\text{GaN}2p_{3/2}\text{-VBM}}^{\text{GaN}}$ ) of eqn (2) denotes the energy of the core level of Ga relative to its respective VBM, which is observed to be 1114.88 eV (Fig. 4(c)). Hence, from eqn (2),  $\Delta E_V$  is calculated to be 2.16 eV. In eqn (3),  $E_G^{\text{GaN}}$  and  $E_G^{\text{MoS}_2}$  represent the band gaps of GaN and MoS<sub>2</sub>, respectively. Thus, from eqn (3),  $\Delta E_C$  is found to be 0.28 eV.

Similarly, the band offset parameters of MoS<sub>2</sub>/Si have been determined by substituting the appropriate terms in eqn (2) and (3) with the corresponding values, as shown in Fig. 4(d–f). The VBO and CBO values for the MoS<sub>2</sub>/Si heterojunction are observed to be –0.37 and –0.14 eV, respectively. Here, the negative sign indicates that the band edges of MoS<sub>2</sub> are below those of Si. Based on these band offset values as well as the values of  $E_C - E_F$  obtained from eqn (1), the band alignments at the MoS<sub>2</sub>/GaN and MoS<sub>2</sub>/Si heterointerfaces are shown in Fig. 4(g and h). The corresponding energy band diagrams of the MoS<sub>2</sub>/GaN/Si and MoS<sub>2</sub>/Si heterojunctions are shown in Fig. S6 (ESI<sup>†</sup>) and are in agreement with the band alignments evaluated through HRXPS measurements.

When the device is illuminated, electrons and holes are generated at both the MoS<sub>2</sub> as well as the Si side. Due to the overall built-in potential developed in the MoS<sub>2</sub>/GaN NRs/Si heterojunction, electrons in the Si side move into the conduction band of MoS<sub>2</sub>. Moreover, these electrons are then collected by the electrode at the MoS<sub>2</sub> side, whereas the holes remain in the valence bands of MoS<sub>2</sub> and Si as depicted in Fig. 5(a). On the other hand, for the MoS<sub>2</sub>/Si heterojunction, the built-in potential is directed from the MoS<sub>2</sub> towards the Si side. Thus, the electrons present in the conduction band of MoS<sub>2</sub> and Si are collected by the electrode at the MoS<sub>2</sub> side, whereas the photogenerated holes move towards the Si side and are collected by the metal contact on Si because of the built-in potential as shown in Fig. 5(b). In addition, the remnant photogenerated holes in the valence bands of MoS<sub>2</sub> and Si in the MoS<sub>2</sub>/GaN/Si heterojunction move into the valence band of Si (as MoS<sub>2</sub> is continuously distributed all over GaN NRs and Si) and are collected by the electrode at the Si side.

Responsivity, detectivity, gain and sensitivity are some important figures of merit used to assess the performance of

a detector. We have calculated these figures of merit for a light illumination intensity of 0.075 mW cm<sup>-2</sup> at 0 V ( $\lambda = 900$  nm). The responsivity ( $R_\lambda$ ) is the ratio of the photocurrent generated to the incident power density and is calculated by using the following equation:<sup>39</sup>

$$R_\lambda = \frac{I_{\text{ph}}}{P_\lambda \times A} \quad (4)$$

where  $I_{\text{ph}}$ ,  $P_\lambda$  and  $A$  are the photocurrent, incident power density, and effective area of the PD.

The specific detectivity ( $D^*$ ) can be determined by the relation:<sup>39</sup>

$$D^* = \frac{R_\lambda}{(2e \times J_d)^{1/2}} \quad (5)$$

where  $J_d$ ,  $R_\lambda$ , and  $e$  are the dark current density, calculated responsivity, and electronic charge ( $1.6 \times 10^{-19}$  C), respectively.

Sensitivity ( $S$ ) is the ratio of the photocurrent to the dark current of the PD. It is determined using the following equation:<sup>39</sup>

$$S = \frac{I_{\text{ph}}}{I_{\text{dark}}} \quad (6)$$

The gain ( $G$ ) of a PD,<sup>39</sup> which can be defined as the number of photogenerated electrons detected per incident photon, assuming that the external quantum efficiency ( $\eta$ ) is 100%, is calculated using the following equation:

$$G = \frac{hc \times R_\lambda}{ne\lambda} \quad (7)$$

From these results, the responsivity, detectivity, sensitivity and gain are estimated to be 14.22 A W<sup>-1</sup>,  $2.72 \times 10^{13}$  Jones,  $\sim 1.17 \times 10^2$ , and 74.49, respectively, in the self-biased mode (light intensity = 0.075 mW cm<sup>-2</sup>). The gain achieved for the device is greater than unity, and this can be ascribed to the strong built-in potential present at the interface, which results in effective carrier separation and hole trapping due to the adsorption/desorption of oxygen species. Due to the sulfur vacancies in MoS<sub>2</sub>, atmospheric oxygen gets adsorbed, capturing the free electrons. When the junction is illuminated with light, electron–hole pairs are created and the photogenerated

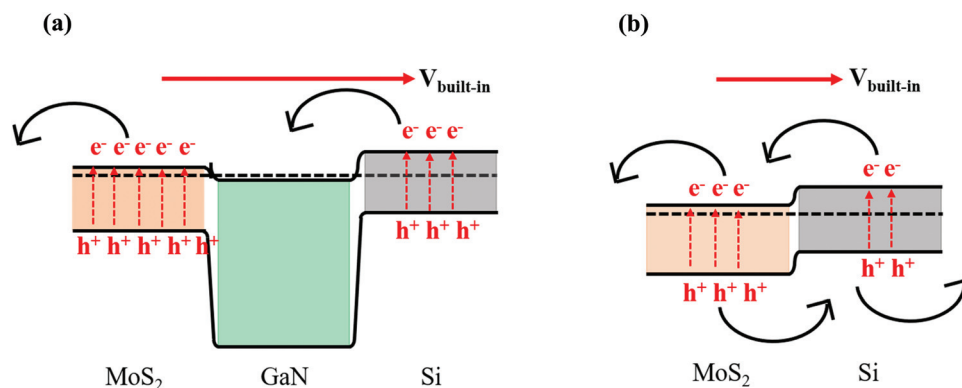


Fig. 5 Carrier transport in the self-powered mode for (a) MoS<sub>2</sub>/GaN NRs/Si and (b) MoS<sub>2</sub>/Si heterojunctions depicted by energy band diagrams.



**Table 1** Performance parameters of the MoS<sub>2</sub>/GaN NRs/Si based PD (this work) in comparison with other reported PDs (working wavelength in the NIR region)

Photodetector	Bias (V)	Responsivity (mA W <sup>-1</sup> )	Detectivity (Jones)	Rise/fall time (ms)
MoS <sub>2</sub> /GaN <sup>42</sup>	20	~5 × 10 <sup>-1</sup> @ 900 nm	—	—/—
GeSe <sup>43</sup>	2	63 @ 900 nm (1.5 W m <sup>-2</sup> )	8 × 10 <sup>8</sup>	—
PbS quantum dots/graphene <sup>44</sup>	4	6.81 × 10 <sup>3</sup> @ 900 nm (27.5 μW)	—	300/3.2 × 10 <sup>3</sup>
MoTe <sub>2</sub> /monolayer graphene <sup>45</sup>	V <sub>D</sub> = 1	9.71 × 10 <sup>5</sup> @ 1064 nm (202 μW)	1.55 × 10 <sup>11</sup>	78/375
MoTe <sub>2</sub> <sup>46</sup>	V <sub>D</sub> = 10, V <sub>G</sub> = 10	24 @ 1060 nm (20 nW)	1.30 × 10 <sup>9</sup>	1.6/1.3
Black phosphorus <sup>47</sup>	V <sub>D</sub> = -1, V <sub>G</sub> = -15	~1 × 10 <sup>6</sup> @ 900 nm (5 mW cm <sup>-2</sup> )	—	~5/—
WSe <sub>2</sub> /SnSe <sub>2</sub> <sup>48</sup>	V <sub>D</sub> = -3, V <sub>G</sub> = -100	3.96 × 10 <sup>5</sup> @ 980 nm (100 mW cm <sup>-2</sup> )	4.40 × 10 <sup>10</sup>	16/45
MoS <sub>2</sub> /Si <sup>49</sup>	0	300 @ 808 nm (10 mW cm <sup>-2</sup> )	~10 <sup>13</sup>	0.003/0.040
MoS <sub>2</sub> /GaN <sup>50</sup>	1	8 × 10 <sup>4</sup> @ 850 nm (2 mW cm <sup>-2</sup> )	—	~2.4 × 10 <sup>3</sup> /~9 × 10 <sup>3</sup>
MoS <sub>2</sub> /GaN NRs/Si (this work)	0	1.42 × 10 <sup>4</sup> @ 900 nm (0.075 mW cm <sup>-2</sup> )	2.72 × 10 <sup>13</sup>	0.039/0.008

holes discharge the negatively charged oxygen ions from MoS<sub>2</sub>. Overall, the adsorption and desorption of the oxygen species on the MoS<sub>2</sub> surface play a crucial role in the trapping of holes and prolong the mean lifetime of the holes. When the mean lifetime of the holes becomes larger than the transit time of the electrons, a high internal gain can be achieved.<sup>33,40,41</sup> To reveal the influence and the importance of MoS<sub>2</sub>, spectral response measurements of the MoS<sub>2</sub>/GaN NRs/Si-based device (in the top–top contact configuration, *i.e.* by taking both the contacts from the top) as well as the GaN NRs/Si-based device have been carried out in the self-powered mode for comparison. It has been observed that the maximum responsivity in the case of the top–top contact configuration is ~1.83 A W<sup>-1</sup> (at a wavelength of 986 nm), whereas, for the case of the GaN NRs/Si-based PD, the maximum responsivity is found to be ~150 mA W<sup>-1</sup> (at a 364 nm wavelength). Moreover, the response in this case is higher in the UV region as compared to that in the visible/NIR region due to the absence of MoS<sub>2</sub>. These results are presented in Fig. S7 of the ESI.†

The performance parameters of our device are summarized in Table 1 and compared with those of several reported PDs from the literature (working wavelength in the NIR region). It can be clearly deduced from this comparison that our hybrid NRs-based device shows excellent performance exhibiting a self-biased, broadband and ultrafast photoresponse. These results show that devices based on the 1D–2D conformal architecture can be a route towards the advancement of futuristic electronics. We have also compared our MoS<sub>2</sub>/GaN NRs/Si-based PD with a previously reported MoS<sub>2</sub>/GaN film/Si-based device.<sup>23</sup> It may be noted that the device shows a broadband photoresponse along with a much-pronounced response in the visible region as GaN NRs also contribute to the photocurrent, in contrast to the narrow range detection obtained in the case of the GaN film-based PD.<sup>23</sup> Furthermore, the response is maximum at 900 nm instead of 995 nm.<sup>23</sup> We believe the differences are due to the very low thickness of MoS<sub>2</sub> on GaN NRs. Moreover, due to the lower channel width in the case of GaN NRs, the response/recovery is much faster as compared to the GaN film-based device. In addition, the phenomenon of polarity inversion is missing in the present case, and the most plausible reason for this could be the low thickness as well as the non-uniform distribution of MoS<sub>2</sub> over GaN NRs, which may result in the development of thermoelectric potential

which is insufficient to separate electrons and holes. Furthermore, the difference in the observed photodetection behavior may also be due to the different growth orientations of GaN in the case of MoS<sub>2</sub>/GaN film/Si and MoS<sub>2</sub>/GaN NRs/Si-based PDs. Thus, a thorough study is required for better understanding of these different behaviors.

## 4 Conclusions

In conclusion, we have reported the conformal growth of MoS<sub>2</sub> on GaN NRs and a broadband, self-biased and ultrafast PD based on the MoS<sub>2</sub>/GaN NRs/Si heterostructure is demonstrated. The PD shows a high photoresponsivity of 14.22 A W<sup>-1</sup> in the self-powered mode for low intensity light illumination along with an ultrafast photoresponse with a rise/fall time of 38.8/8.2 μs. We believe that this study could provide a very important step towards engineering core–shell nanostructured devices in the field of low power optoelectronics.

## Conflicts of interest

The authors declare no conflicts of interest.

## Acknowledgements

The authors thank Dr. Arun Malla Chowdhury, Dr. Basanta Roul and Ms. Malti Kumari for useful discussions and help in characterization. The authors acknowledge Prof. Kausik Majumdar and Mr Nithin Abraham from the Department of Electrical Communication Engineering, Indian Institute of Science, Bangalore, India for their assistance in ultrafast time measurements. The authors are thankful to Mr. Ambresh M. for his help in TEM and its sample preparation. The authors also thank Mr. Debashis Tripathy from the Department of Inorganic and Physical Chemistry, Indian Institute of Science, Bangalore, India for assistance in HRXPS measurements. D. K. S. is thankful to the Council of Scientific and Industrial Research (CSIR), Government of India, New Delhi and Indian Institute of Science, Bangalore, India for providing a senior research fellowship and research associateship, respectively. S. B. K. acknowledges the INSA senior scientist fellowship.



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