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# Graphitization induced structural transformation of candle soot carbon into carbon nano-onion as a functional anode for metal-ion batteries†

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This work reports for the first time the structural transformation of glassy spherical candle soot carbon into graphitic polyhedral carbon nano-onions (multilayer fullerenes) by simple graphitization of candle soot carbon at varying temperatures from 1500 to 2400 °C. This structural transformation of candle soot carbon is thoroughly investigated by high-resolution transmission electron microscopy, X-ray diffraction, and Raman spectroscopy, with the aid of first-principles calculation based on density functional theory to understand the effect on physiochemical properties. Further, this understanding of structural transformation enables a promising pathway to tune the electrochemical properties of graphitized candle soot carbon to develop it as a functional anode for metal-ion (Li-ion as well K-ion) battery applications.

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## 1. Introduction

Fullerenes, especially C<sub>60</sub>, have attracted great interest in different fields of science and technology due to their appealing chemical, physical, biological, and electrochemical properties.<sup>1</sup> Nano-onion-like fullerenes, also called carbon nano-onions (CNOs), are important members of the fullerene family that have been widely studied and used in various applications, including energy storage materials, high-performance and high-temperature wear-resistant materials, superconductive materials, and biomaterials.<sup>2</sup> CNOs are zero-dimensional spherical or concentric shell structures of closed carbon shells, multi-layered quasi-spherical, and polyhedral-shaped shells.<sup>3</sup> A few methods have been reported to produce carbon onions, including arc-discharge,<sup>4</sup> electron irradiation,<sup>5</sup> chemical vapor deposition (CVD),<sup>6</sup> and nano-diamond annealing.<sup>7</sup> Recently, the graphitization of amorphous carbon into a polycrystalline graphitic structure has been reported. Liu *et al.*<sup>8</sup> proposed a strategy to enhance the graphitization and surface area of carbonaceous materials simultaneously, besides N, P co-doping. However, the quality of the obtained graphitized material was poor compared to other reported techniques, such as arc-discharge grown CNTs and mechanically exfoliated graphene.<sup>9</sup> Additionally, understanding of the graphitization of amorphous

carbon into a polycrystalline graphitic structure is still unclear and remains highly debatable.<sup>10</sup>

However, the cost-effective large-scale production of CNOs is still a challenging task. Therefore, it has limited their commercial availability. It has been reported<sup>11</sup> that quasi-spherical carbon nano-onions have an extrapolated price of \$3350 million per ton. The huge cost of carbon onions could be due to the high price of nano-diamond precursors (~\$2.4 million per ton) and the production process.<sup>11</sup> Therefore, the cost-effective mass production of carbon onions with uniform morphology and size is a bottleneck for their commercial viability and use in device applications.

Here, we report a facile method to produce spherical and polyhedral (hexagonal) carbon nano-onion structured materials by the cost-effective graphitization of low-cost candle soot carbon. Carbon soot is a relatively novel and not-so-well explored carbon nanomaterial among existing carbonaceous materials, consisting of small and ordered graphitic nano-structures and particle sizes ranging from 10 to 50 nm.<sup>12,13</sup> Here, we use the graphitization process to transform the amorphous nature of candle soot carbon to the spherical and polyhedral morphology of carbon nano-onions (multi-layer fullerenes) by improvising its lattice order and graphitic content. We comprehensively study the graphitization temperature dependence of the candle soot carbon micro-structure. Candle soot carbon is transformed into spherical carbon and graphitic polyhedral onions as the graphitization temperature increases (1500 to 2400 °C). This provides a toolkit to enhance its applications further. For example, in this work, we report the graphitization of candle soot carbon as a promising way to develop and tune its structural properties to enable its use as a

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functional anode material for energy storage applications, particularly for lithium and potassium-ion battery development.

## 2. Results and discussion

The candle soot carbon was prepared by a flame synthesis process, *i.e.*, direct combustion of a commercial candle in air, as shown in Fig. 1a. The calculated candle soot carbon deposition rate on a stainless-steel sheet (20 cm × 30 cm) is ~85 mg min<sup>-1</sup>. The physical and chemical properties of the candle soot carbon depend on the collection of soot from two different regions of the flame. For example, candle soot deposited from the tip of the flame is oxidized, conductive, and contains a lower amount of wax and organic compounds compared to the soot collected from the middle of the flame.<sup>14</sup> In this work, carbon candle soot was collected from the tip of the flame over a stainless-steel sheet, as shown in Fig. 1a. The low-resolution transmission electron micrograph (TEM) image (Fig. 1b) indicates that the carbon soot consists of interconnected nanoparticles, forming a fractal (branched) like a network of nanoparticles. At higher magnification (Fig. 1c), candle soot particles (~45 nm in size) have a short-range ordered graphitic layer with an interlayer spacing of 0.354 nm. Later, high-resolution transmission electron microscope (HRTEM) analysis was used to observe the structural change in candle soot carbon over/after each graphitization temperature. The HRTEM micrograph was also utilized to calculate the interlayer spacing of the microstructure and particle size.

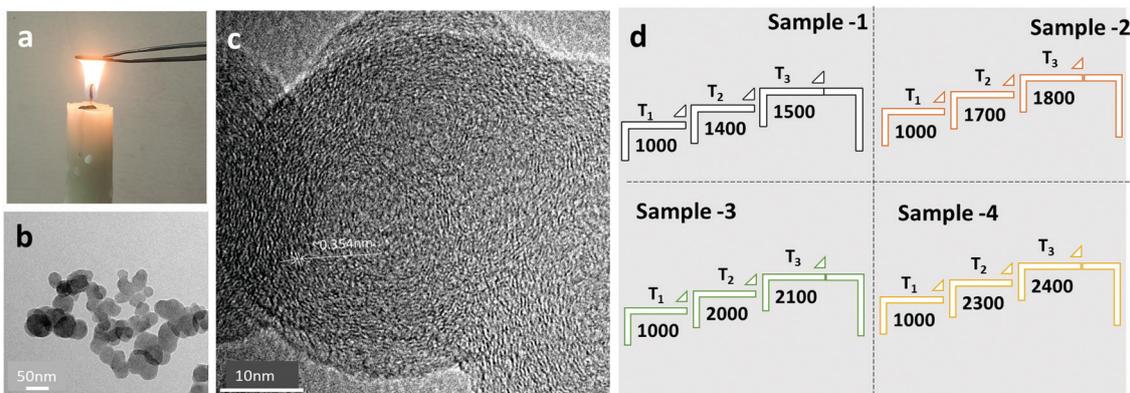
The collected candle soot carbon powder was cold-pressed into 12 mm diameter pellets with a uniaxial press. The resultant candle soot pellets were used for graphitization at different temperatures. Graphitization of candle soot pellets was carried out in a graphitization furnace (Make: Nabertherm GMBH) in an inert and high-purity nitrogen (99.99%) atmosphere. The pellet was placed into a tubular furnace with purging by 1.5 L min<sup>-1</sup> of N<sub>2</sub> gas flow for 15 min. After the initial purging, the N<sub>2</sub> gas flow was maintained at 0.5 L min<sup>-1</sup> throughout the graphitization process.

The graphitization of candle soot was carried out in a three-step process. (i) the temperature (T<sub>1</sub>) was raised at a ramp rate of 10 °C min<sup>-1</sup>, (ii) then, the temperature (T<sub>2</sub>) was reached at a ramp rate of 5 °C min<sup>-1</sup>, and (iii) finally, the graphitization temperature (T<sub>3</sub>) was raised at a ramp rate of 3 °C min<sup>-1</sup> and allowed to dwell there for 120 min. The furnace cooled down to room temperature in N<sub>2</sub> gas flow with a cooling rate of 30 °C min<sup>-1</sup>. The conditions for the entire graphitization process of candle soot are described in Fig. 1d.

After graphitization at 1500 °C, the morphology of the candle soot particles becomes spherical (Fig. 2a), from quasi-spherical (not symmetrical) observed in the case of bare candle soot (collected directly from the flame). We can observe there is a start to concentric graphite layers with an interlayer distance of 0.350 nm. Still, there are a few amorphous carbon regions with a graphitization trend forming an onion-like structure that also indicates the partial (initial) graphitization of the candle soot carbon. Also, the SAED pattern of the graphitized candle soot at 1500 °C has clear diffraction rings which correspond to the (002), (110), and (004) crystal planes of graphite.<sup>15</sup>

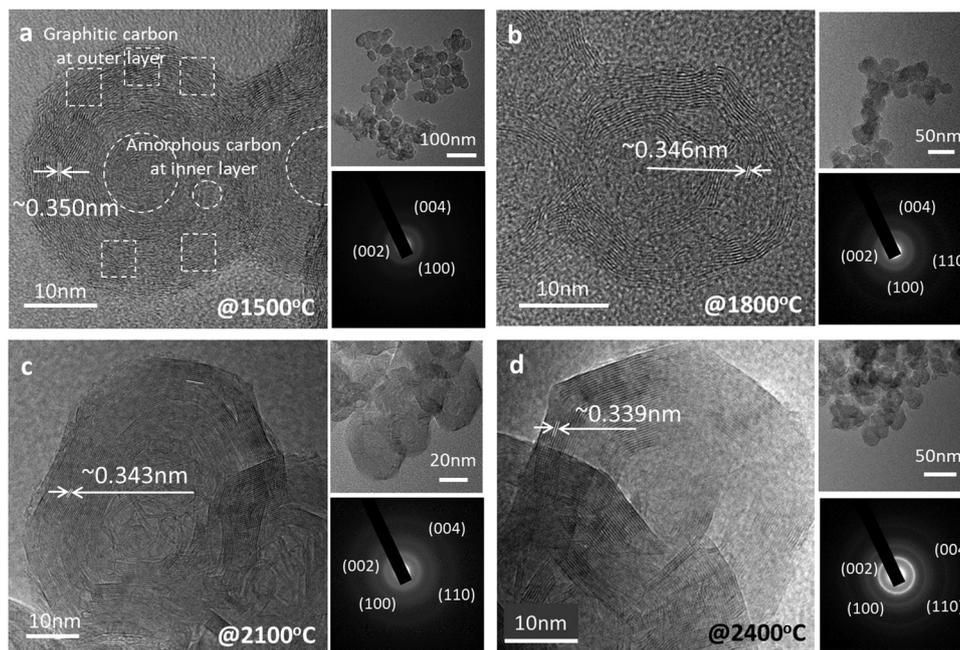
As the graphitization temperature increases to 1800 °C, the morphology of graphitized candle soot carbon becomes (partially) polyhedral with concentric graphite layers having an interlayer distance of 0.346 nm (Fig. 2b). These (partially) polyhedral carbon onions form more concentric graphite layers at the surface of the onions and fewer at the center. As the graphitization temperature increases further to 2100 °C (Fig. 2c), the candle soot carbon becomes polyhedral and irregular in shape. It has more graphitic content with lattice fringes of 0.343 nm. The SAED pattern also supports this observation and confirms that polyhedral carbon onion has more graphitic content.

However, at 2400 °C, the morphology of the carbon onion is nearly perfectly polyhedral (hexagonal) around 40 nm in diameter without any significant amount of amorphous carbon particles in its structure (Fig. 2d). The polyhedron has a nearly uniform graphitic structure from the surface of the onion towards the center. Also, the lattice fringes of the polyhedral



**Fig. 1** (a) Flame synthesis of candle soot carbon. (b) TEM morphology of candle soot particles after flame synthesis at lower magnification. (c) TEM morphology of candle soot particles at higher magnification. (d) Schematic showing the three-step graphitization of candle soot under different temperature conditions.

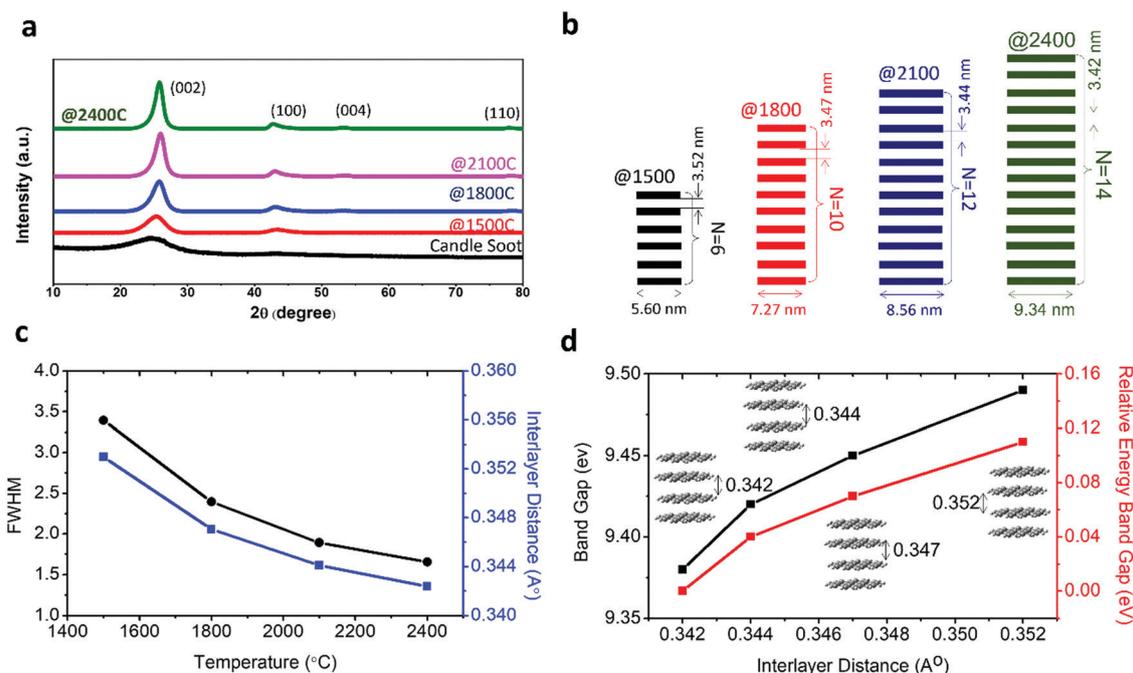




**Fig. 2** (a) TEM morphology of candle soot particles after graphitization at 1500 °C at lower and higher magnification with SAED pattern. (b) TEM morphology of candle soot particles after graphitization at 1800 °C at lower and higher magnification with SAED pattern. (c) TEM morphology of candle soot particles after graphitization at 2100 °C at lower and higher magnification with SAED pattern. (d) TEM morphology of candle soot particles after graphitization at 2400 °C at lower and higher magnification with SAED pattern.

onions are more distinct and straighter than those of the spherical onions of candle soot carbon with an interlayer distance of about 0.339 nm, close to that of the (002) graphite plane.

This indicates that candle soot carbon nanoparticles transformed to polyhedral onions have more ordered and concentric graphitic layers, as also shown in the SAED diffraction rings.



**Fig. 3** (a) X-ray diffraction of graphitized candle soot carbon after graphitization from 1500 °C to 2400 °C. (b) The schematic shows the calculated structural parameters, interlayer distance  $d_{(002)}$ , number of graphitic layers ( $N$ ),  $L_a$  and  $L_c$  of the graphitized candle soot carbon, with graphitization temperature. (c) Change in FWHM and interlayer distance with graphitization temperature. (d) Variation in bandgap energy with the change in interlayer distance in a carbon onion during graphitization.



The change in crystallinity during the graphitization of candle soot carbon at varying temperatures was also analyzed using powder X-ray diffractometry with a step size of about  $0.02^\circ$ , as shown in Fig. 3a. A broad and weak diffraction peak at about  $26.3^\circ$  confirms that bare candle soot carbon is amorphous (disordered) with an interlayer spacing ( $d_{002}$ ) of 0.355 nm. The interlayer spacing ( $d_{002}$ ) (eqn (S1), ESI†) was calculated from Bragg's law,<sup>16</sup> and crystallite height ( $L_c$ ) along the  $c$ -axis (eqn (S2), ESI†) and crystallite size ( $L_a$ ) along the  $a$ -axis (eqn (S3), ESI†) were evaluated by the Debye–Scherrer formula<sup>17</sup> (see Section S1 for more details, ESI†). The calculated structural parameters, interlayer distance  $d_{(002)}$ , number of graphitic layers ( $N$ ), and  $L_a$  and  $L_c$  of the graphitized candle soot carbon, with temperature are depicted in Fig. 3b.

Peaks at around  $25^\circ$ ,  $43^\circ$ ,  $53^\circ$  and  $78^\circ$  indicate the diffraction peaks corresponding to the (002) (100), (004) and (110) planes of graphitic carbon.<sup>15</sup> The (002) peak is attributed to the stacking structure of graphene layers along the  $c$ -axis (*i.e.*, the orientation of the aromatic ring carbon layers in a three-dimensional arrangement).<sup>18</sup> The (100) plane mainly represents the hexagonal plane structure of carbon along the  $a$ -axis.<sup>19</sup> The narrow and high intensities of the (002) and (100) peaks indicate the better orientation and the larger size of the aromatic layer slices, respectively.<sup>20</sup> After graphitization at  $1500^\circ\text{C}$ , the intensities of the (002) and (100) diffraction peaks increase; this indicates that the graphitic content in the candle soot carbon structure is enhanced with a closer interlayer spacing of 0.352 nm. The crystallite size ( $L_a$ ) and the crystallite height ( $L_c$ ) also increase, indicating the improvement in the stacking of

aromatic ring carbon layers along the  $c$ - and  $a$ -axes. As the graphitization temperature increases from  $1500^\circ\text{C}$  to  $2100^\circ\text{C}$ , the FWHM peak width decreases, indicating an increase in  $sp^2$ -bonded graphitic content in the graphitized carbon onion (Fig. 3c).

The calculated  $d$ -spacing of graphitized candle soot carbon from XRD analysis is also consistent with our observation from HRTEM (Fig. 2). With increasing temperature, the interlayer distance decreases and causes the two-dimensional microcrystallinity to grow gradually along the  $a$ - and  $c$ -axes. This is also supported by the calculated crystallite size ( $L_a$ ) and the crystallite height ( $L_c$ ) along the  $a$ - and  $c$ -axes, respectively. It is interesting to note that the (002) and (100) diffraction peaks become narrow and sharp, while other diffraction peaks of carbon corresponding to the (004) and (110) planes start evolving with increasing graphitization temperature (Fig. 3a). This indicates that the number of parallel graphitic layers increases, and the aromatic graphene layer becomes orderly with higher graphitization temperature.<sup>18</sup> At  $2400^\circ\text{C}$ , we achieved higher graphitization of candle soot carbon particles into polyhedral carbon onions with an interlayer spacing of 0.342 nm (comparable to pristine graphite), crystallite size ( $L_a$ ) of 4.91 nm, and crystallite height ( $L_c$ ) of 9.34 nm, with nearly 14 parallel graphitic layers along the  $c$ -axis.

To understand the effect of the change in interlayer distance of graphene layers in carbon onions during graphitization on the energy gap, we performed first-principles calculations based on the density functional theory (DFT) (see Section S2 for more details, ESI†). The variation in relative energy gap with interlayer distance is depicted in Fig. 3d. This result indicates

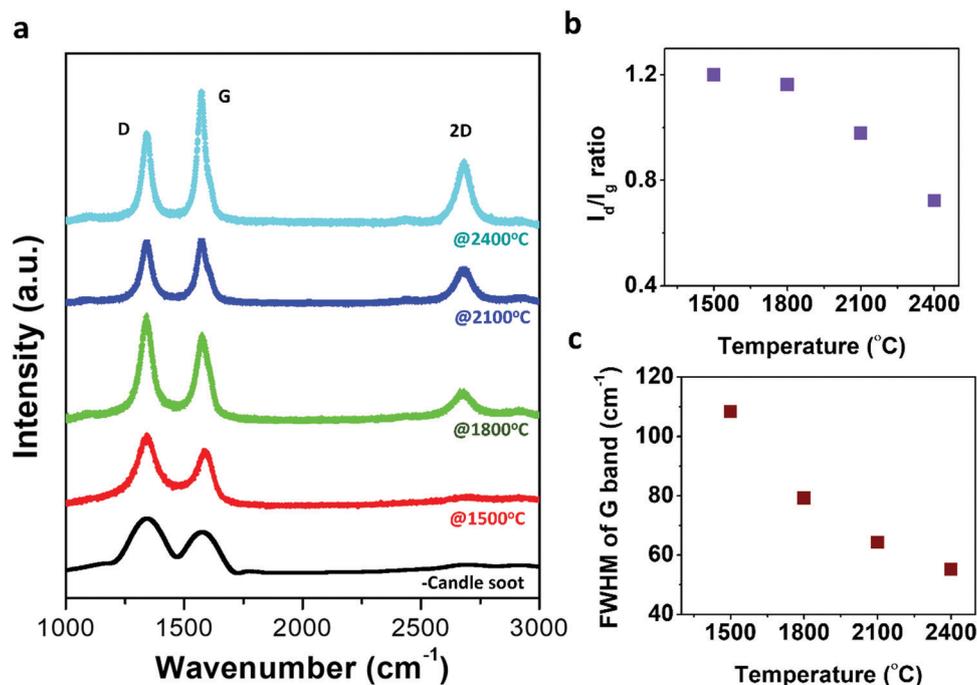


Fig. 4 (a) Raman spectra of collected candle soot and graphitized candle soot carbon after graphitization from  $1500^\circ\text{C}$  to  $2400^\circ\text{C}$ . (b) The change in  $I_d/I_g$  ratio with graphitization temperature. (c) Variation in FWHM of the G band with the structural transformation of candle soot carbon during graphitization.



that the relative energy gap increases with interlayer distance and can reach up to  $\sim 110$  meV at an interlayer distance of  $3.52 \text{ \AA}$  (it is achieved by graphitization at  $1500 \text{ }^\circ\text{C}$ ). Thus, we can consider that graphitization is also a promising and controllable way to tune the energy gap of candle soot carbon onions for developing different electronic and semiconductor devices such as p–n junctions, transistors, photodiodes, and lasers.<sup>21</sup>

Raman spectroscopy, a non-destructive characterization technique, was applied to understand the microstructural changes occurring during the graphitization process from candle soot carbon to polyhedral carbon onions and to investigate the ordering of the carbon structure during the transformation. Raman spectra were recorded at room temperature on candle soot carbon after graphitization from  $1500 \text{ }^\circ\text{C}$  to  $2400 \text{ }^\circ\text{C}$ , as shown in Fig. 4a.

Raman spectra show two modes; the first-order mode is in the ranges  $1000\text{--}1800 \text{ cm}^{-1}$  and the second-order is in the range  $2200\text{--}3100 \text{ cm}^{-1}$ . The first-order mode shows the G band and D band around  $1347 \text{ cm}^{-1}$  and  $1589 \text{ cm}^{-1}$ , respectively. The positions of the D band and G band mostly rely on the qualities of the carbonaceous material structure, such as vacancies, grain size, edge defects, dangling bonds, and other disorders.<sup>22</sup> The G band is related to the  $E_{2g}$  vibrations corresponding to crystalline  $sp^2$  domains and bond stretching of all pairs of  $sp^2$  atoms in both rings and chains (it is indicative of crystalline graphite content).<sup>23</sup> Whereas the D band is observed due to the radial  $A_{1g}$  breathing mode of vibration (breathing modes of  $sp^2$  atoms in the ring), which, in turn, signifies the presence of defects along with the disordered amorphous carbon content.<sup>23</sup> On increasing the graphitization temperature from  $1500 \text{ }^\circ\text{C}$  to  $2400 \text{ }^\circ\text{C}$ , the intensity of the G band monotonically increases with a reduction in the intensity of the D band, indicating an increase in the graphitic content and a reduction in defects with graphitization temperature. At  $2400 \text{ }^\circ\text{C}$ , the intensity of the G band dominated over the D band. The  $sp^2$  content was 56% more than the  $sp^3$  content (46%), indicating that polyhedral onions consist of more ordered graphitic  $sp^2$  planes. This observation can be confirmed with clear and straight lattice fringes of polyhedral candle soot carbon onions, as observed in the HRTEM micrograph. The  $sp^2$  and  $sp^3$  content were evaluated by calculating the relative area under the D and G bands. The relative intensity of the D to G bands ( $I_d/I_g$ ) is widely used to trace the defects present on carbon structures (*i.e.*, degree of graphitization).<sup>24,25</sup> Fig. 4b shows the ratio ( $I_d/I_g$ ) decreases with increasing graphitization temperature, indicating the polyhedral carbon onion structure has a more ordered structure and tends to be graphitic. Fig. 4c represents the change in full width at half maximum (FWHM) of the G band with graphitization temperature of the graphitized candle soot carbon material. FWHM of the G band mainly represents the disorder within the carbon  $sp^2$  sheets.<sup>26</sup> As the graphitization temperature increased, the FWHM decreased, indicating the lower disordered content within the carbon  $sp^2$  sheets. The in-plane crystallite thickness for the graphitized carbon structure is calculated using the Tuinstra and Koenig relation  $L_a = 4.4(I_g/I_d)$ .<sup>27</sup> The in-plane crystallite thickness for

candle soot carbon increases from  $3.6 \text{ nm}$  to  $6.0 \text{ nm}$ , with graphitization temperature ranging from  $1500 \text{ }^\circ\text{C}$  to  $2400 \text{ }^\circ\text{C}$ . It is also implied that the graphitized candle soot carbon onion has a greater number of  $sp^2$  domains with less disordered content compared to non-graphitized candle soot carbon.<sup>28</sup>

Additional peaks (including 2D and combined (D + G)) are observed in the range  $2500\text{--}2800 \text{ cm}^{-1}$ , which reflects the second-order two-phonon process associated with defects arising from all kinds of  $sp^2$  carbon materials.<sup>29</sup> It can be observed that the 2D band starts evolving and becomes sharp as the graphitization temperature increases. The 2D band in this graphitized carbon material exhibits a doublet structure, which means that the stacking of graphene layers is ordered along the hexagonal axis. It is mainly associated with the splitting of the  $\pi$  and  $\pi^*$  electronic states, as a result of interactions between the successive layer planes.<sup>23,30</sup> In the case of a one-layer graphene structure, the 2D band forms a singlet, whereas, for two-layer graphene, it exhibits a quadruplet structure.<sup>30</sup> Here, the Raman spectrum (2D band) of graphitized carbon onion at  $1500 \text{ }^\circ\text{C}$  is almost indistinguishable from that of a polyhedral carbon onion at  $2400 \text{ }^\circ\text{C}$ . This is due to the 2D band splitting merging into a doublet.<sup>30</sup>

At temperatures above  $1500 \text{ }^\circ\text{C}$ , the transformation of candle soot carbon nanoparticles to carbon nano-onions is initiated by the thermal energy associated with graphitization temperature. As the graphitization temperature increases, more ordered graphitic content develops with decreasing distance between the graphene layer and increasing crystallite size ( $L_c$ ), as shown in Fig. 5a. The calculated activation energy (please refer eqn (S4) and (S5), ESI†) for the transformation of candle soot carbon nanoparticles to polyhedral carbon nano-onions is  $8.74 \text{ kcal mol}^{-1}$  with a pre-exponential constant of  $0.002$  (see Section S3 for more details, ESI†). Here, the activation energy for the graphitization of candle soot carbon particles is comparably lower than for the graphitization of other precursor materials.<sup>31</sup> This may be due to the nano-effect of candle soot carbon particles, as the kinetics of nanomaterials is much faster than for bulk materials.<sup>32</sup>

The change in lattice orientation during graphitization can also be analyzed using the XRD pattern. Williamson and Hall reported that the observed XRD diffraction line broadening (FWHM) is due to crystallite size and micro-strain. Therefore, we used the diffraction patterns of graphitized candle soot carbon to evaluate the average crystallite size ( $d_{\text{XRD}}$ ) and microstrain ( $\epsilon$ ) at each graphitization temperature using the Williamson–Hall (W–H) relation (eqn (S6), ESI†)<sup>33,34</sup> (see Section S4 for more details, ESI†). The W–H curve for candle soot carbon with graphitization temperature (Fig. 5b) indicates that the strain present in the sample is compressive. This compressive strain in the material leads to the overall contraction of a particle with a reduction in interlayer spacing, which induces curvature across the carbon nano-onion and may form polyhedral (hexagonal) rings, as observed in the TEM images. As the graphitization temperature increases, the decrease in compressive strain may be due to the reduction in lattice imperfections, as observed in TEM, XRD, SAED, and Raman analyses.<sup>35</sup> Similarly, it has been



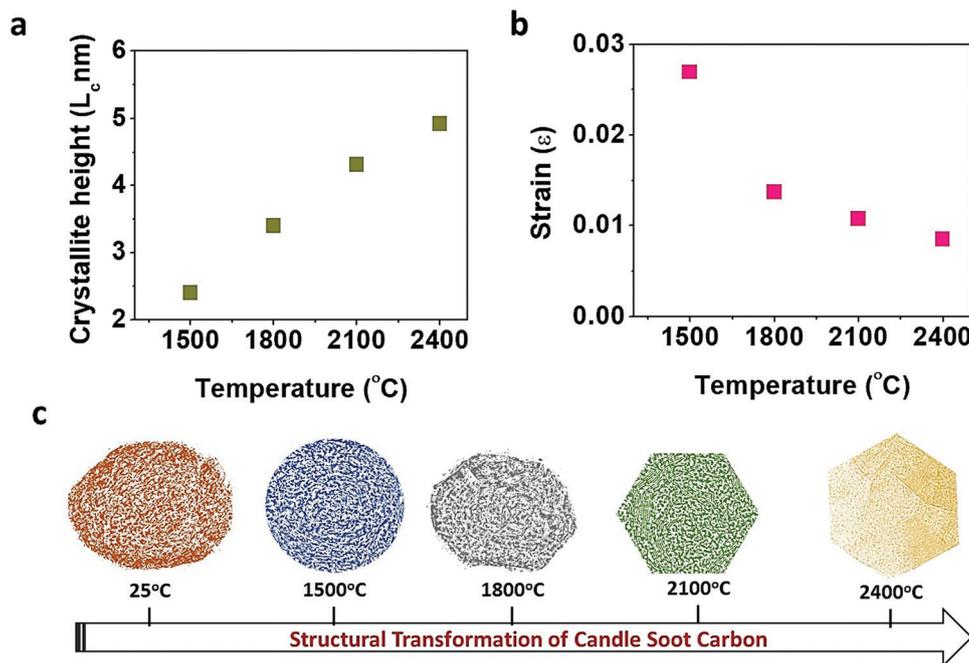


Fig. 5 (a) Variation of crystallite size ( $L_c$ ) with graphitization temperature. (b) The change in micro-strain with graphitization temperature. (c) Schematic showing the structural transformation of candle soot carbon during graphitization.

observed that when a carbon onion is exposed to electron irradiation it undergoes self-compression, leading to the formation of pentagonal and heptagonal rings that induce curvature in the network.<sup>36</sup>

The observed size of hexagonal carbon onions in the TEM micrograph is slightly smaller than the initial size of candle soot carbon nanoparticles. The decrease in the volume/size of the graphitized polyhedral carbon onion may be caused by the compressive stress and difference in the densities between the hard carbon nature of candle soot carbon (the density of hard carbon is  $\sim 1.45$ – $1.55$ ) and the more graphitic nature of the polyhedral carbon onion (the density of graphite is  $\sim 2.15$ – $2.25$ ).<sup>37</sup> The present study also clarifies that the polyhedral onions have a more ordered graphitic  $sp^2$  content and are energetically favorable at a higher temperature of graphitization (2400 °C) than spherical carbon onions obtained from graphitization of candle soot carbon at 1500 °C. A similar type of transformation and behavior is also observed in the case of CNTs.<sup>38</sup> It has been observed that normal nanotubes with circular shapes are energetically favorable at low temperatures. However, the configurational entropy accompanied by Stone–Wales defect creation at high temperatures makes polygonal nanotubes or nanofibers. These polygonal structures are thermodynamically more stable and contain less aromatic graphene interlayer spacing.<sup>38</sup> Here, spherical candle soot carbon onion is also energetically more stable at low temperature, as observed through TEM analysis. However, at high temperatures the polygonal (hexagonal) onion structure is more favorable with a reduction in graphene interlayer spacing, as observed in TEM and XRD patterns.

The TEM micrograph and Raman spectra show that the candle soot carbon and graphitized spherical carbon nano-onions consist

of small domains of graphitic  $sp^2$  sheets with defects in the structure. Therefore,  $\pi$ -electrons in candle soot carbon and spherical carbon onions are mostly localized in small domains, which cannot actively take part in the conduction of electrons.<sup>39</sup> However, the graphitization temperature increases the polyhedral carbon onion form with more graphitic  $sp^2$  sheets, reduces graphene interlayer spacing, and decreases the number of defects, resulting in the delocalization of  $\pi$  electrons.<sup>37</sup> Therefore, our studies along with other findings<sup>37,39</sup> strongly confirm that the structure of spherical onions is energetically favoured at low temperatures and far from the perfectly closed graphitic structure. In contrast, polyhedral carbon onions at high temperatures are thermodynamically more stable (Fig. 5c), accompanied by a highly ordered graphitic structure, with a reduction in the graphene interlayer spacing and more conduction  $\pi$  electrons within it. Further, we have measured the electrical conductivities of directly collected candle soot carbon (CCSC) and graphitized candle soot carbon at 2400 °C (GCSC). The electrical conductivities of CCSC and GCSC measured from the  $I$ - $V$  characteristic curves obtained by linear sweep voltammetry. The current response is recorded by varying the voltage from 1 to 4.0 V with a scan rate of  $0.01 \text{ V s}^{-1}$ . The current response varies linearly with the applied voltage between the two probes at a fixed distance. The slope of the linear fit of the curves is used to obtain the conductivity of the samples. The conductivities for CCSC and GCSC are calculated to be  $0.13$  and  $0.20 \text{ S cm}^{-1}$ . The electrical conductivity of GCSC was found to be the higher due to the high amount of ordered graphitic domains obtained during the graphitization process, as discussed in the previous section.

Further, we demonstrated graphitized candle soot carbon as an anode material for a metal-ion battery and how their



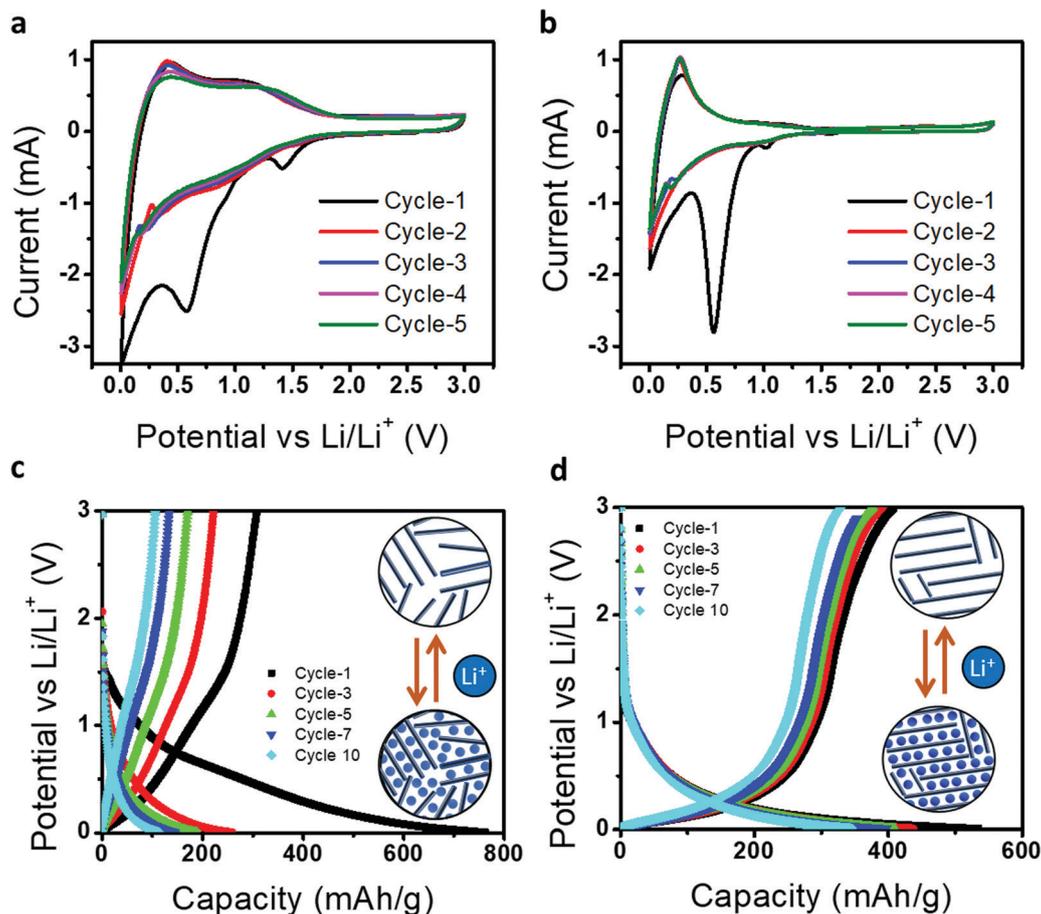


Fig. 6 (a) Cyclic voltammety curves (first 5 cycles) of collected candle soot carbon (CCSC) electrode for LIB. (b) Cyclic voltammety curves (first 5 cycles) of graphitized candle soot carbon (GCSC) electrode for LIB. (c) Charge–discharge curve of CCSC electrode for the first 10 cycles in LIB. (d) Charge–discharge curve of GCSC electrode for the first 10 cycles in LIB.

structural transformation engineered the tuning of the electrochemical performance. Fig. 6a and b represents the cyclic voltammety curves of directly collected candle soot carbon (CCSC) and graphitized candle soot carbon at 2400 °C (GCSC) at a scan rate of 0.1 mV s<sup>-1</sup> between 0.05 and 3.0 V. Galvanostatic charge–discharge was carried out at 50 mA g<sup>-1</sup> for all measurements. In the first discharge cycle for CCSC, there are two distinct reduction peaks at approximately 1.40 and 0.60 V (vs. Li<sup>+</sup>/Li) and for GCSC they are at 1.0 and 0.55 V, which is mainly attributed to electrolyte decomposition and the formation of a protective solid electrolyte interphase (SEI).<sup>40</sup> However, the peak (0.15–0.25 V) during the anodic scan indicates the removal of adsorbed/de-intercalation of Li-ions from the carbon structure.<sup>41</sup>

After the first cycle, GCSC shows stable CV curves in terms of shape and current density over the CCSC, indicating that the graphitization of CCSC at 2400 °C (GCSC) increases the graphitic content and efficient lithiation/de-lithiation over the GCSS. Fig. 6c and d show the electrochemical charge–discharge profiles of CCSC and GCSC, with a cut-off voltage of 0.005 to 3.0 V at a current density of 100 mA g<sup>-1</sup>, respectively. In the first cycle, CCSC offers high discharge capacity (762 mA h g<sup>-1</sup>) but suffers from low Coulombic efficiency (40.2%). Whereas GCSC shows an

initial discharge capacity of 534 mA h g<sup>-1</sup> with an improved Coulombic efficiency of 76.4%. Here CCSC offers a high initial discharge capacity (762 mA h g<sup>-1</sup>) but shows low capacity in the following cycles compared to GCSC because CCSC has a more disordered structure consisting of numerous structural defects (e.g., defects, edges, and residual heteroatoms) which provide plenty of storage sites with a wide distribution of binding energies for metal-ions.<sup>42</sup> The presence of a more disordered structure in CCSC results in a high initial discharge capacity with SEI formation, electrolyte decomposition and irreversible Li<sup>+</sup> capture. Furthermore, CCSC and GCSC have graphitic as well as disordered contents (from Raman analysis); therefore metal-ion storage in CCSC and GCSC materials possesses both sloping and plateau voltage regions, involving an intercalation mechanism, adsorption mechanism, and capacitive mechanism. However, CCSC achieved more capacity (discharge/charge curve) in the sloping voltage region because of the adsorption of ions on the surface of active sites (edge carbon fragments and heteroatoms such as oxygen and hydrogen) while the GCSC achieved the capacity in a low-voltage plateau region which is attributed to metal-ion intercalation/deintercalation into the graphitic layers.<sup>43</sup>

Further, during cycling of GCSC over CCSC, the discharge plateaus shift down (positively) and the charge plateaus shift up



(negatively), indicating that the structural transformation of candle soot carbon into carbon nano-onion lowers the polarization and increases the Li-ion cell reversibility.

Next, we demonstrated the use of CCSC and GCSC in a K-ion battery and their relative electrochemical performances. The graphite anode used in a potassium-ion battery has a limited theoretical capacity of  $279 \text{ mA h g}^{-1}$  with poor cyclic stability compared to an Li-ion battery;<sup>44</sup> therefore, it may be advantageous to use glassy carbon as an anode. Fig. 7a and b represent cyclic voltammetry (CV) curves for five cycles at a scan rate of  $0.1 \text{ mV s}^{-1}$  in a voltage range of 0.01–3.0 V (vs.  $\text{K}^+/\text{K}$ ).

For CCSC, the broad cathodic reduction peaks located at  $\sim 0.6$  and  $\sim 0.01$  V and anodic peaks located at 0.3 V correspond to the formation of a solid–electrolyte interface (SEI), intercalation, and de-intercalation of K-ions in CCSC, respectively.<sup>45</sup> Whereas GCSC has a cathodic peak located at  $\sim 0.78$  V and 0.65–0.7 V, and an anodic peak located at 0.68 V, representing the formation of SEI,<sup>46</sup> insertion of potassium ions, and extraction of potassium ions in GCSC,<sup>47</sup> respectively. Here, CCSC shows relatively broad cathodic and anodic profiles for the CV curves compared to GCSC. This indicates that CCSC

has high surface-driven storage of K-ions due to the presence of disordered carbon and multiple defects, as revealed in Fig. 2 and 3.

Further the K-ion storage performance of CCSC and GCSC was analyzed by galvanostatic charge–discharge, as shown in Fig. 7c and d. In the first cycle, CCSC shows a lower Coulombic efficiency (56.9%) than GCSC because CCSC contains many defective structures, leading to irreversible K-ion trapping and more SEI formation. We should also note that the initial Coulombic efficiency of CCSC is improved for the K-ion system (56.9%) compared to the Li-ion system (40.2%). This shows the advantage of the glassy nature of CCSC as an anode for a K-ion battery. However, as GCSC has a better degree of graphitization with limited defective structures, it still leads in terms of reversible and more efficient K-ion intercalation with higher Coulombic efficiency (76.2%). Further, it can be observed that charge–discharge curves of CCSC contain a highly sloping region with higher voltage, whereas GCSC shows a low sloping region at a lower voltage. This GCD behavior indicates that CCSC possesses surface-induced K-ion storage due to the presence of defects with amorphous carbon content. Whereas

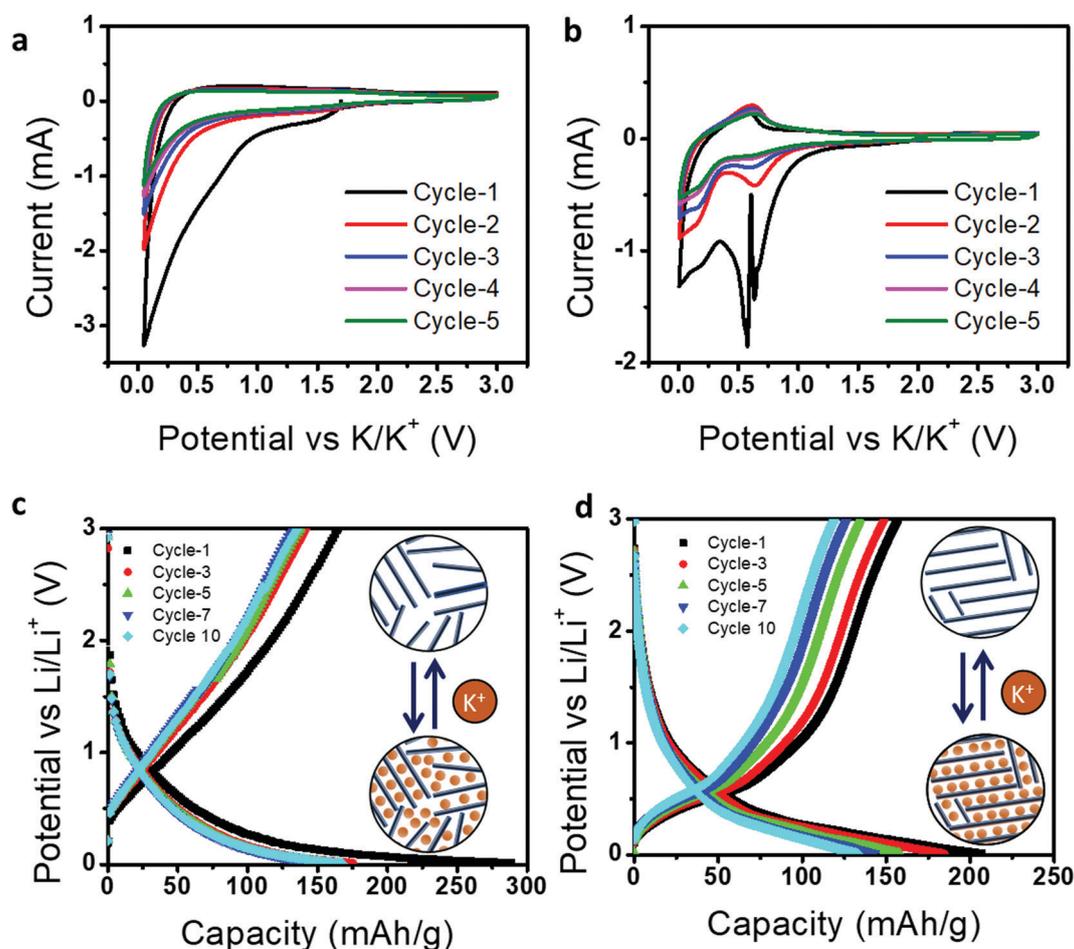


Fig. 7 (a) Cyclic voltammetry curves (first 5 cycles) of collected candle soot carbon (CCSC) electrode for KIB. (b) Cyclic voltammetry curves (first 5 cycles) of graphitized candle soot carbon (GCSC) electrode for KIB. (c) Charge–discharge curve of CCSC electrode for the first 10 cycles in KIB. (d) Charge–discharge curve of GCSC electrode for the first 10 cycles in KIB.



GCSC shows K-ion intercalation/de-intercalation due to the presence of a more graphitic structure. Thus, this study suggests that the electrochemical behaviors of energy storage devices are highly dependent on structural transformation and, therefore, by simply tuning the graphitization conditions of candle soot carbon materials, one can extend their utilization for wider electrochemical applications.

### 3. Conclusion

This work reports a facile method of preparing a spherical and polyhedral (hexagonal) carbon nano-onion structure by graphitization of candle soot carbon in the temperature range 1500 °C–2400 °C. The structural transformation of candle soot carbon with graphitization temperature was thoroughly investigated and it was observed that the structure of spherical onions is energetically favored at low graphitization temperature (1500 °C). However, the structure of polyhedral carbon onions is thermodynamically more stable at high graphitization temperature (2400 °C). Polyhedral carbon onions have a highly graphitic ordered structure, reduced graphene interlayer spacing, and greater crystallite size along the *a*- and *c*-axes compared to spherical onions. Thus, this work provides a mechanistic understanding and structural transformation of candle soot amorphous carbon into spherical and polyhedral graphitic carbon nano-onions. This study will also open new and unprecedented avenues for future research in graphitized candle soot carbon and tune its properties as a functional material, as described in this work as a case study for Li-ion and K-ion battery applications.

### Conflicts of interest

Authors have no conflicts of interest to disclose.

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### References

- H. Gergeroglu, S. Yildirim and M. F. Ebeoglugil, *SN Appl. Sci.*, 2020, **2**, 603.
- B. Xu, *New Carbon Mater.*, 2008, **23**, 289–301.
- A. Albu-Yaron, M. Levy, R. Tenne, R. Popovitz-Biro, M. Weidenbach, M. Bar-Sadan, L. Houben, A. N. Enyashin, G. Seifert, D. Feuermann, E. A. Katz and J. M. Gordon, *Angew. Chem., Int. Ed.*, 2011, **50**, 1810–1814.
- N. Sano, H. Wang, I. Alexandrou, M. Chhowalla, K. B. K. Teo, G. A. J. Amaratunga and K. Iimura, *J. Appl. Phys.*, 2002, **92**, 2783–2788.
- H. E. Troiani, A. Camacho-Bragado, V. Armendariz, J. L. Gardea Torresday and M. Jose Yacaman, *Chem. Mater.*, 2003, **15**, 1029–1031.
- C. Ruan, Z. Li, D. Zhang, X. Yuan, C. Liang, Y. Chang, H. Huang, L. Xu and M. Chen, *Carbon*, 2020, **161**, 622–628.
- L. Hawelek, A. Brodka, S. Tomita, J. C. Dore, V. Honkimäki and A. Burian, *Diamond Relat. Mater.*, 2011, **20**, 1333–1339.
- H. Liu, W. Zeng, Y. Yang, J. Chen, Y. Zhao and S. Mu, *J. Mater. Chem. A*, 2021, **9**, 1260–1268.
- A. Barreiro, F. Börrnert, S. M. Avdoshenko, B. Rellinghaus, G. Cuniberti, M. H. Rummeli and L. M. K. Vandersypen, *Sci. Rep.*, 2013, **3**, 1115.
- G. C. Loh and D. Baillargeat, *J. Appl. Phys.*, 2013, **114**, 33534.
- X. Liu, J. Ren, G. Licht, X. Wang and S. Licht, *Adv. Sustainable Syst.*, 2019, **3**, 1900056.
- S. Singh and N. Verma, *Int. J. Hydrogen Energy*, 2015, **40**, 1145–1153.
- S. Singh and N. Verma, *Int. J. Hydrogen Energy*, 2015, **40**, 5928–5938.
- C.-J. Liang, J.-D. Liao, A.-J. Li, C. Chen, H.-Y. Lin, X.-J. Wang and Y.-H. Xu, *Fuel*, 2014, **128**, 422–427.
- T. Ungár, J. Gubicza, G. Ribárik, C. Pantea and T. W. Zerda, *Carbon*, 2002, **40**, 929–937.
- K. T. Jeng, S. F. Lee, G. F. Tsai and C. H. Wang, *J. Power Sources*, 2004, **138**, 41–50.
- C. Ganesh Kumar, S. Pombala, Y. Poornachandra and S. V. Agarwal, Synthesis, characterization, and applications of nanobiomaterials for antimicrobial therapy, in *Nanobiomaterials in Antimicrobial Therapy*, ed. A. M. Grumezescu, William Andrew Publishing, 2016, pp. 103–152.
- T. Qiu, J.-G. Yang, X.-J. Bai and Y.-L. Wang, *RSC Adv.*, 2019, **9**, 12737–12746.
- J. L. Santos, M. A. Centeno and J. A. Odriozola, *J. Anal. Appl. Pyrolysis*, 2020, **148**, 104821.
- W. Li and Y. Zhu, *Energy Fuels*, 2014, **28**, 3645–3654.
- Y. Zhang, T.-T. Tang, C. Girit, Z. Hao, M. C. Martin, A. Zettl, M. F. Crommie, Y. R. Shen and F. Wang, *Nature*, 2009, **459**, 820–823.
- S. Stankovich, D. A. Dikin, R. D. Piner, K. A. Kohlhaas, A. Kleinhammes, Y. Jia, Y. Wu, S. T. Nguyen and R. S. Ruoff, *Carbon*, 2007, **45**, 1558–1565.
- L. Bokobza, J.-L. Bruneel and M. Couzi, *C*, 2015, **1**(1), 77–94.
- T. Palaniselvam, H. B. Aiyappa and S. Kurungot, *J. Mater. Chem.*, 2012, **22**, 23799–23805.
- V. M. Samoilov, D. B. Verbets, I. A. Bubnenkov, N. N. Stepanyova, A. V. Nikolaeva, E. A. Danilov, D. V. Ponomareva and E. I. Timoshchuk, *Inorg. Mater. Appl. Res.*, 2018, **9**, 890–899.
- D. S. Knight and W. B. White, *J. Mater. Res.*, 1989, **4**, 385–393.
- P. Puech, M. Kandara, G. Paredes, L. Moulin, E. Weiss-Hortala, A. Kundu, N. Ratel-Ramond, J.-M. Plewa, R. Pellenq and M. Monthieux, *C*, 2019, **5**(4), 69.
- G. Irmer and A. Dorner-Reisel, *Adv. Eng. Mater.*, 2005, **7**, 694–705.
- R. Pathak, K. Chen, A. Gurung, K. M. Reza, B. Bahrami, F. Wu, A. Chaudhary, N. Ghimire, B. Zhou, W.-H. Zhang, Y. Zhou and Q. Qiao, *Adv. Energy Mater.*, 2019, **9**, 1901486.
- A. C. Ferrari and D. M. Basko, *Nat. Nanotechnol.*, 2013, **8**, 235–246.



- 31 J. A. Newell, D. D. Edie and E. L. Fuller Jr., *J. Appl. Polym. Sci.*, 1996, **60**, 825–832.
- 32 Q.-S. Fu, Y.-Q. Xue, Z.-X. Cui and M.-F. Wang, *J. Nanomater.*, 2014, **2014**, 856489.
- 33 G. K. Williamson and W. H. Hall, *Acta Metall.*, 1953, **1**, 22–31.
- 34 I. F. Cruz, C. Freire, J. P. Araújo, C. Pereira and A. M. Pereira, in *Micro and Nano Technologies*, ed. A. A. El-Gendy, J. M. Barandiarán and R. L. B. T.-M. N. M. Hadimani, Elsevier, 2018, pp. 59–116.
- 35 A. K. Kole and P. Kumbhakar, *Results Phys.*, 2012, **2**, 150–155.
- 36 Y. Gan and F. Banhart, *Adv. Mater.*, 2008, **20**, 4751–4754.
- 37 S. Tomita, T. Sakurai, H. Ohta, M. Fujii and S. Hayashi, *J. Chem. Phys.*, 2001, **114**, 7477–7482.
- 38 M. Yoon, J. Howe, G. Tibbetts, G. Eres and Z. Zhang, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2007, **75**, 165402.
- 39 G. Treboux, P. Lapstun and K. Silverbrook, *Chem. Phys. Lett.*, 1999, **302**, 60–64.
- 40 X.-L. Wu, Q. Liu, Y.-G. Guo and W.-G. Song, *Electrochem. Commun.*, 2009, **11**, 1468–1471.
- 41 T. Wang, Y. Wang, G. Cheng, C. Ma, X. Liu, J. Wang, W. Qiao and L. Ling, *Energy Fuels*, 2020, **34**, 8911–8918.
- 42 Y. Huang, Y. Wang, P. Bai and Y. Xu, *ACS Appl. Mater. Interfaces*, 2021, **13**, 38441–38449.
- 43 R. Ding, Y. Huang, G. Li, Q. Liao, T. Wei, Y. Liu, Y. Huang and H. He, *Front. Chem.*, 2020, **8**, 1175.
- 44 Z. Tang, Y. Wang, Z. Zheng and X. Luo, *J. Mater. Chem. A*, 2021, **9**, 9191–9202.
- 45 G. Wang, X. Xiong, D. Xie, Z. Lin, J. Zheng, F. Zheng, Y. Li, Y. Liu, C. Yang and M. Liu, *J. Mater. Chem. A*, 2018, **6**, 24317–24323.
- 46 E. Zhang, X. Jia, B. Wang, J. Wang, X. Yu and B. Lu, *Adv. Sci.*, 2020, **7**, 2000470.
- 47 M. Wang, Y. Zhu, Y. Zhang, T. Yang, J. Duan and C. Wang, *Electrochim. Acta*, 2021, **368**, 137649.

