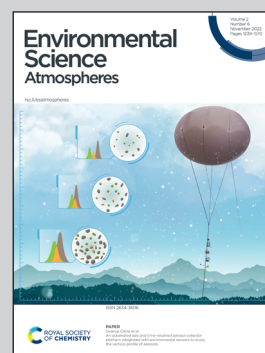


Showcasing research from Professor Fárník's laboratory,  
Department of Dynamics of Molecules and Clusters,  
J. Heyrovský Institute of Physical Chemistry, CAS, Prague,  
Czech Republic.

Molecular-level insight into uptake of dimethylamine  
on hydrated nitric acid clusters

Clusters in molecular beams can represent laboratory model systems for atmospheric aerosols. In this study, hydrated nitric acid clusters pick up individual dimethylamine molecules. Mass spectrometry of positively and negatively charged clusters accompanied by ab initio calculations reveals details about acid-base reactions and ionization processes in these small mixed clusters. Such research can provide a detailed molecular-level insight into the new particle formation processes in the atmosphere.

### As featured in:



See Jozef Lengyel, Eva Pluhařová,  
Michal Fárník *et al.*,  
*Environ. Sci.: Atmos.*, 2022, 2, 1292.

## PAPER

View Article Online  
View Journal | View Issue



Cite this: *Environ. Sci.: Atmos.*, 2022, 2, 1292

## Molecular-level insight into uptake of dimethylamine on hydrated nitric acid clusters†

Andriy Pysanenko,<sup>a</sup> Karolína Fárníková,<sup>a</sup> Jozef Lengyel,<sup>b</sup> Eva Pluhařová<sup>a</sup> and Michal Fárník<sup>a</sup>

Mixed nitric acid/water clusters with dimethylamine (DMA) represent a suitable model system for understanding acid–base chemistry in atmospherically relevant clusters. We investigate these clusters in a detailed molecular-beam experiment accompanied by *ab initio* calculations. The  $(\text{HNO}_3)_M(\text{H}_2\text{O})_N$  clusters are produced by supersonic expansion into vacuum and doped by DMA molecules in a pickup process. Two complementary mass spectrometry approaches are employed to analyze the resulting  $(\text{DMA})_k(\text{HNO}_3)_M(\text{H}_2\text{O})_N$  clusters: (i) electron impact ionization at 70 eV to form positive cluster ions and (ii) low-energy electron attachment at 0–10 eV to form negative clusters. The positive ion spectra contain mainly protonated  $(\text{DMA})_k(\text{HNO}_3)_m\text{H}^+$  clusters with  $k = m + 1$ , whereas the negative ones are dominated by  $(\text{DMA})_k(\text{HNO}_3)_m\text{NO}_3^-$  with  $m \geq k$ , followed by  $(\text{DMA})_k(\text{HNO}_3)_m\text{NO}_2^-$  ( $m > k$ ) ions with low abundances. These observations are rationalized by our calculations, which exhibit the tendency of the mixed clusters to maximize the number  $\text{DMA} \cdot \text{H}^+$  and  $\text{NO}_3^-$  ions in the clusters. In the neutral clusters, this is fulfilled for 1 : 1 ratio of DMA and  $\text{HNO}_3$ , while the positively charged  $(\text{DMA})_k(\text{HNO}_3)_m\text{H}^+$  clusters satisfy this condition for  $k = m + 1$ . The protonated clusters always contain the  $\text{DMA} \cdot \text{H}^+$  moiety. For the negatively charged cluster ions, thermochemistry supports the prevailing formation of  $\text{NO}_3^-$  and  $m \geq k$  ion composition. Furthermore, the  $\text{NO}_3^-$ -containing cluster ions can form when an electron attaches to the protonated moiety of the  $\text{DMA} \cdot \text{H}^+ \cdots \text{NO}_3^-$  ion pair in the cluster, which leads to H atom evaporation. From the gas phase  $\text{HNO}_3$  molecule, where  $\text{NO}_2^-$  is formed exclusively upon an electron attachment, the tendency to form  $\text{NO}_3^-$  increases to hydrated  $\text{HNO}_3$  clusters, where both  $\text{NO}_2^-$  and  $\text{NO}_3^-$  ions are generated in approximately equal abundances, to the DMA doped clusters, where  $\text{NO}_3^-$  strongly prevails  $\text{NO}_2^-$ .

Received 27th July 2022  
Accepted 19th August 2022

DOI: 10.1039/d2ea00094f

rsc.li/esatmospheres

### Environmental significance

New particle formation (NPF) in atmospheric aerosol generation is still insufficiently understood. It is well-established that acid–base chemistry plays a key role in NPF, yet, it is difficult to obtain a truly molecular-level understanding from field and aerosol chamber investigations. Studies of isolated clusters in molecular beams can provide such picture. Nitric acid and dimethylamine (DMA) are ubiquitous throughout the atmosphere and are involved in NPF. We dope hydrated nitric acid clusters with DMA in a molecular beam experiment complemented by theoretical calculations, and provide a molecular-level insight into these species and their chemistry. Discussed relationship between neutral cluster structure and positively and negatively charged cluster ions is relevant for implementation of mass spectrometry in aerosol research.

## 1 Introduction

Aerosol particles play a pivotal role in the atmosphere: they influence the Earth's energy balance by scattering and absorption of solar and terrestrial radiation and thus determine the

global climate, provide environment for heterogeneous chemical reactions and multiphase processes and control the chemistry and composition of the atmosphere.<sup>1–4</sup> Moreover, aerosol particles have pronounced effects on human health.<sup>2,5</sup> Despite their significance and huge effort in their investigation, aerosols still belong to the least known component of our atmosphere.

Significant fraction of atmospheric aerosol particles originates from molecular clusters formed from single molecules in the gas phase *via* gas-to-particle conversion known as new particle formation (NPF). They grow to larger sizes and eventually become stable against evaporation. However, the detection of the smallest particles in the nanometer size range directly in the atmosphere is difficult.<sup>6</sup> Nevertheless, these sizes,

<sup>a</sup>J. Heyrovský Institute of Physical Chemistry, v.v.i., The Czech Academy of Sciences, Dolejškova 2155/3, Prague 8 182 23, Czech Republic. E-mail: michal.farnik@jh-inst.cas.cz; eva.pluharova@jh-inst.cas.cz

<sup>b</sup>Chair of Physical Chemistry, TUM School of Natural Sciences, Technical University of Munich, Garching 85748, Germany. E-mail: jozef.lengyel@tum.de

† Electronic supplementary information (ESI) available: Experimental details, additional mass spectra and their discussion, further computational results, cluster structures and coordinates. See <https://doi.org/10.1039/d2ea00094f>



where the transition from the gas phase clusters to stable particles occurs, are important.<sup>7</sup> The laboratory alternative to investigate the clusters in the NPF process is offered by various aerosol chamber experiments such as CESAM in Paris<sup>8</sup> or CLOUD in CERN<sup>9</sup> and others,<sup>10</sup> or various aerosol flow-tube experiments, see, *e.g.*, ref. 11–13. A different perspective is offered by experiments with clusters in molecular beams in vacuum. They can provide a bottom-up approach and truly molecular-level insight.

It has been well established that the key components in the NPF process are acids, especially sulfuric acid, which have to be stabilized in the clusters by bases such as ammonia and the simple alkyl amines abundant in the atmosphere.<sup>2,14–17</sup> Other species, such as highly oxygenated organic molecules (HOMs) might stabilize the clusters as well,<sup>18–21</sup> and water vapour is indispensable in the NPF process too.<sup>2,14,15</sup> However, dehydrated clusters are often detected in laboratory experiments,<sup>17,22–24</sup> thus the exact role of water still has to be better understood.<sup>15</sup>

Though nitric acid ( $\text{HNO}_3$ ) is mainly considered in polar stratospheric cloud particles,<sup>3,25</sup> it is ubiquitous throughout the entire atmosphere and its involvement in the particle growth in the atmosphere has been studied.<sup>26–29</sup> Recent work has indicated a rapid NPF and growth by nitric acid and ammonia condensation.<sup>9</sup> These species are more abundant in polluted areas and thus nitric acid can be responsible for urban smog. Although, there are no indications that clusters of nitric acid  $\text{HNO}_3$  with dimethylamine  $(\text{CH}_3)_2\text{NH}$  (DMA) are directly involved in NPF in atmosphere, they represent an excellent model system for understanding detailed acid–base chemistry in the NPF process. Nanoparticle formation and growth from the nitric acid  $\text{HNO}_3$  and dimethylamine  $(\text{CH}_3)_2\text{NH}$  (DMA) was investigated before in a flow tube experiment accompanied by theoretical modelling.<sup>13</sup> The authors deduced acid : base ratio close to 1 : 1 in all their relative humidity conditions, which was different from the previous work on sulfuric acid/base particles, namely  $\text{H}_2\text{SO}_4$ –DMA.<sup>12,30,31</sup> Here we approach similar system, DMA– $\text{HNO}_3$ – $\text{H}_2\text{O}$  clusters, from a complementary experimental and theoretical perspective.

Molecular beam studies of clusters with atmospheric relevance started from the early cluster experiments in 1970s (*e.g.*, see recent reviews<sup>32,33</sup>). The early experiments with the hydrated  $\text{HNO}_3/\text{H}_2\text{O}$  clusters performed in Castleman's group,<sup>34</sup> revealed the acid dissociation and  $\text{H}_3\text{O}^+\cdots\text{NO}_3^-$  ion pair generation in these clusters. These early studies were extended by our experiments with Na-doping of the  $(\text{HNO}_3)_M(\text{H}_2\text{O})_N$  clusters,<sup>35</sup> and electron attachment to these clusters.<sup>36,37</sup> Later on, we have investigated also uptake of various atmospherically relevant molecules by these clusters.<sup>33,38,39</sup>

Here, we dope the preexisting hydrated nitric acid clusters with the dimethylamine molecules: the small well defined  $(\text{HNO}_3)_M(\text{H}_2\text{O})_N$  clusters fly isolated in a molecular beam in vacuum and adsorb individual DMA molecules in a pickup chamber.<sup>33</sup> Due to the finite size of the clusters, this method allows us observing acid–base reactions and water evaporation directly in the small  $(\text{DMA})_K(\text{HNO}_3)_M(\text{H}_2\text{O})_N$  clusters. The mass spectrometry of the positively as well as negatively charged clusters accompanied by *ab initio* quantum chemical

calculations of these clusters offers an unprecedented detailed insight into the reaction and ionization processes in these clusters. This truly molecular-level insight provides a basis for detailed understanding of the NPF processes involving nitric acid and dimethylamine.

## 2 Experimental and theoretical methods

### 2.1 Experiment

The experiments have been performed with our versatile cluster beam apparatus CLUB described previously, *e.g.*, in our recent reviews<sup>32,33</sup> and references therein. The apparatus scheme is provided in accompanying ESI.† Briefly, the  $(\text{HNO}_3)_M(\text{H}_2\text{O})_N$  clusters are produced by a continuous supersonic expansion of nitric acid vapour from  $\text{HNO}_3$ –water solution (65% Lach-Ner) in buffer gas helium (99.9999%, Messer). The nitric acid solution was evaporated in a temperature controlled stainless steel reservoir placed directly in the source vacuum chamber with an attached conical nozzle (diameter 90  $\mu\text{m}$ , length 2 mm, full opening angle 30 deg.). The nozzle was heated separately to a temperature usually 5–10 °C higher than the reservoir to prevent condensation in the nozzle. The clustering conditions were controlled by the reservoir temperature  $T_R$  and He stagnation pressure  $P_0$ , as exploited in our previous work.<sup>35–37</sup> The present results were obtained at  $T_R = 65^\circ\text{C}$  and  $P_0 = 3$  bar. Previous analysis<sup>35</sup> revealed that under similar conditions neutral  $(\text{HNO}_3)_M(\text{H}_2\text{O})_N$  clusters were generated with  $M \approx 1\text{--}6$  and  $N \approx 1\text{--}15$ .

The cluster beam was skimmed and passed to a pickup chamber, where dimethylamine (99.8% Sigma-Aldrich) was introduced *via* a needle valve with a controlled pressure ranging from  $p_c = 2 \times 10^{-4}$  mbar to  $2 \times 10^{-3}$  mbar. Mass spectra at different DMA pressures were recorded. It ought to be stressed that the individual clusters fly isolated without any interaction in high vacuum, and even at the highest pickup pressure  $p_c = 2 \times 10^{-3}$  mbar they undergo only several individual collisions with isolated molecules (note that this pickup pressure is still almost 6 orders of magnitude below the atmospheric pressure).

After flying another  $\approx 1.5$  m through two differentially pumped chambers in ultrahigh vacuum ( $\leq 10^{-7}$  mbar) the clusters entered the vacuum chamber with a perpendicularly mounted reflectron time-of-flight mass spectrometer (RTOF), which was operated either in the positive or negative ion mode.

In the positive mode, the clusters were ionized by a pulsed electron gun at a frequency of 10 kHz at 70 eV electron energy. A typical electron current was 10  $\mu\text{A}$ . The ionization pulse width was 5  $\mu\text{s}$ . After 0.5  $\mu\text{s}$  delay, to avoid any free electron effects, the ions were extracted by a 4 kV pulse and further accelerated to 8 keV into the time-of-flight region. After  $\approx 95$  cm flight path, the ions were detected by an MCP detector.

In the negative ion mode, the voltages on the extraction plates were inverted. The clusters were ionized at low electron energies  $E_e$  0–14 eV, and the spectra were recorded in the electron energy steps of 0.2 eV and integrated. The electron pulse was 5  $\mu\text{s}$  long, and after 0.5  $\mu\text{s}$  delay a 3  $\mu\text{s}$  extraction pulse was applied to extract the





negative ions. A typical electron current at the present experiment was 0.2  $\mu\text{A}$  at 5 eV. The individual mass spectra recorded at different electron energies allow for the evaluation of the electron energy dependence of the negative ion yield for individual mass peaks.<sup>37</sup> The energy spectra were normalized to the electron current. Since the current is very low at electron energies close to zero, below  $E_e \approx 1.5$  eV the spectra can be subject to a large error.<sup>37</sup>

## 2.2 Theory

In order to provide detailed interpretation of collected mass spectra, we characterized structures of dimethylamine with hydrated nitric acid clusters using *ab initio* calculations and *ab initio* molecular dynamics simulations. We focused on the following clusters:  $\text{DMA}_K(\text{HNO}_3)_M(\text{H}_2\text{O})_N$  ( $K = 1-3$ ;  $M = 1-3$ ;  $N = 2-6$ ),  $\text{DMA}_K\text{HNO}_3(\text{H}_2\text{O})_n\text{H}^+$  ( $k = 1-3$ ;  $n = 2-6$ ) and negatively charged  $\text{DMA}_K(\text{HNO}_3)_m(\text{H}_2\text{O})_n\text{NO}_3^-$  and  $\text{DMA}_K(\text{HNO}_3)_m(\text{H}_2\text{O})_n\text{NO}_2^-$  ( $k = 1-3$ ;  $m = 1, 2$ ;  $n = 0-4$ ) which will be used to interpret the mass spectra.

To sample the configurational space of the clusters, we used Born–Oppenheimer molecular dynamics<sup>40</sup> using the CP2K program package.<sup>41</sup> For treatment of the electronic structure we applied revPBE<sup>42</sup> functional with the D3 dispersion correction.<sup>43</sup> The core electrons were described by the Goedecker–Tetter–Hutter norm-conserving pseudopotentials,<sup>44</sup> with Kohn–Sham orbitals being expanded in TZV2P MOLOPT basis set.<sup>45</sup> For the auxiliary plane-basis set a cutoff of 400 Ry was used. The equations of motion were integrated with a time step of 0.5 fs. To keep the temperature constant a CSVR thermostat<sup>46</sup> was used with a coupling constant of 50 fs. We generated at least one 20 ps trajectory at 300 K for each of the systems, specifically one for clusters containing one  $\text{HNO}_3$  unit, two for dehydrated clusters with two  $\text{HNO}_3$  units and three for larger clusters. Periodic boundary conditions were applied. Box length of the cubic simulation cell was 14 Å.

To calculate the most stable structure of each system, we optimized structures at first at the PM6 (ref. 47) level of theory which was then followed up by the M06-2X/aug-cc-pVDZ<sup>48,49</sup> level of theory. The DFT functional has been chosen in order to be consistent with previous studies of systems containing hydrated nitric acid<sup>36,39</sup> where it was benchmarked against CCSD(T). M06-2X has been shown to provide reliable results for atmospherically relevant clusters.<sup>50</sup> We gathered these structures from snapshots of the previously calculated trajectories. Snapshots were spaced by 1 ps. We performed frequency analysis to confirm the local minimum. The energy span of the resulting structures was 20–90  $\text{kJ mol}^{-1}$ , the structure with the lowest energy was then used to evaluate the binding energy of either dimethylamine or water to the cluster at M06-2X/aug-cc-pVDZ level. The binding energy includes zero-point vibrational energy. We performed all quantum chemical calculations in the Gaussian 16.<sup>51</sup>

## 3 Experimental results

### 3.1 Positive spectra

(Fig. 1) summarizes the evolution of the positive ion mass spectra with the increasing pressure of DMA in the pickup chamber. Top

panel (a) shows the spectrum of pure  $(\text{HNO}_3)_M(\text{H}_2\text{O})_N$  clusters without any pickup. It exhibits the protonated mass peak series  $(\text{HNO}_3)_m(\text{H}_2\text{O})_n\text{H}^+$ , which are labeled by different symbols and colors for series with different number of nitric acid molecules  $m = 0-4$ . Some peaks are labeled by their acid-to-water  $m : n$  ratio and the corresponding  $m/z$  for easier orientation. The peculiar shape of the spectrum has been observed in previous experiments<sup>34,35,38,39</sup> and it was explained by the acid dissociation and  $\text{H}_3\text{O}^+ \cdots \text{NO}_3^-$  ion pair generation.<sup>34,35</sup>

Panels below show the spectra evolution with increasing pressure of DMA in the pickup chamber (b)  $8.6 \times 10^{-4}$  mbar and (c)  $1.1 \times 10^{-3}$  mbar. Additional mass spectra measured at different  $p_c$  pressures are shown in ESI.† The spectra were measured under the same conditions except for the pickup gas pressure  $p_c$ . Thus, the differences between panels in the relative signal intensities reflect the decreasing cluster beam intensity with increasing DMA pressure due to the cluster scattering out of the beam. The relative increase of peak intensities at the higher mass side of the spectrum with respect to the lower mass side is simply the consequence of the easier scattering of smaller (lighter) clusters out of the beam.

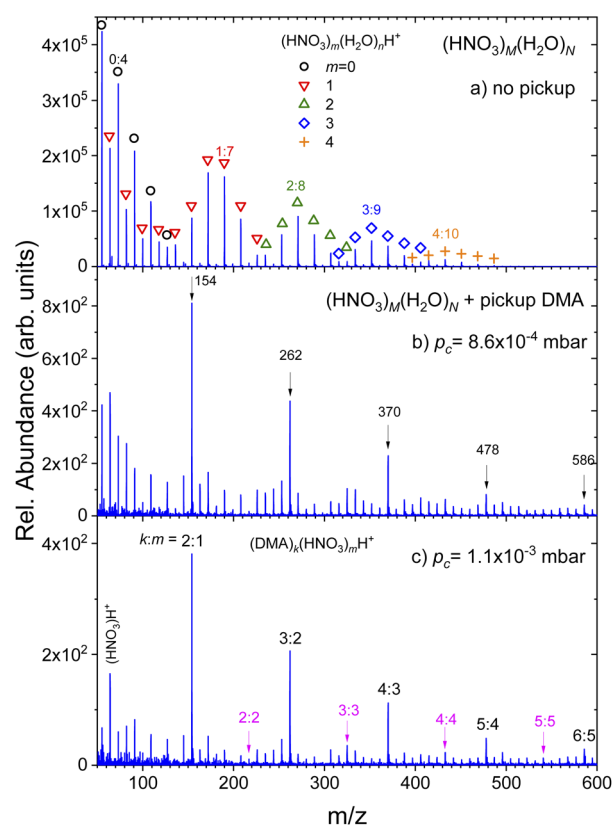


Fig. 1 Mass spectra of  $(\text{HNO}_3)_M(\text{H}_2\text{O})_N$  clusters after the pickup of DMA molecules. The top spectrum (a) shows the pure clusters without any pickup. Individual protonated mass peak series  $(\text{HNO}_3)_m(\text{H}_2\text{O})_n\text{H}^+$  are labeled by symbols and some peaks are labeled by  $m : n$  ratio for easier orientation. Panels below show the spectra evolution with increasing pressure of DMA in the pickup chamber: (b)  $8.6 \times 10^{-4}$  mbar and (c)  $1.1 \times 10^{-3}$  mbar. The most pronounced peaks are labeled by their  $m/z$  in (b), and assigned  $k : m$  of  $(\text{DMA})_k(\text{HNO}_3)_m\text{H}^+$  peaks in (c).



The important difference between the spectra are the pronounced peaks that clearly stand out as the DMA pressure increases. They are labeled by their mass-to-charge ratio ( $m/z$ ) in panel (b). Although there are mass coincidences with the original  $(\text{HNO}_3)_m(\text{H}_2\text{O})_n\text{H}^+$  masses, these peaks must be due to the picked up DMA molecules. Therefore, they are assigned to  $(\text{DMA})_k(\text{HNO}_3)_m\text{H}^+$  ions with  $k = m + 1$ , as indicated in the last panel (c) where these peaks are labeled by  $k : m$ . Upon careful inspection, further mass peaks corresponding to the  $(\text{DMA})_k(\text{HNO}_3)_m\text{H}^+$  series with  $k = m$  could be unambiguously identified as indicated also in panel (c). However, this series is very weak compared to the  $k = m + 1$  series. No hydrated species  $(\text{DMA})_k(\text{HNO}_3)_m(\text{H}_2\text{O})_n\text{H}^+$  could be identified unambiguously. They would partly coincide with the original  $(\text{HNO}_3)_m(\text{H}_2\text{O})_n\text{H}^+$  series, however, we have not observed any intensity anomalies which would point to the generation of the  $(\text{DMA})_k(\text{HNO}_3)_m(\text{H}_2\text{O})_n\text{H}^+$  mass peaks upon the DMA uptake.

### 3.2 Negative spectra

(Fig. 2) shows the negative ion mass spectra after the electron attachment. Top panel (a) shows the spectrum of the pure  $(\text{HNO}_3)_M(\text{H}_2\text{O})_N$  clusters. The expansion conditions and the clusters were the same as in (Fig. 1a). The major ion series  $(\text{HNO}_3)_m(\text{H}_2\text{O})_n\text{NO}_2^-$  (red circles) and  $(\text{HNO}_3)_m(\text{H}_2\text{O})_n\text{NO}_3^-$  (green triangles) are indicated. Similar spectra were published and discussed previously.<sup>36,37</sup> We briefly recap: The dissociative electron attachment (DEA) to an isolated  $\text{HNO}_3$  molecule yielded

essentially only  $\text{NO}_2^-$  ions with the relative abundance of 96.5% and the rest were  $\text{OH}^-$  ions. The  $\text{NO}_3^-$  ions were observed with only negligible 0.03% abundance.<sup>36</sup> On the other hand, the electron attachment to the hydrated  $\text{HNO}_3$  in the clusters changed the populations in favor of the  $\text{NO}_3^-$  containing cluster ions. The  $(\text{HNO}_3)_m(\text{H}_2\text{O})_n\text{NO}_3^-$  ions with different  $m$  and  $n$  dominated the spectrum with integrated intensity of 56% over the  $(\text{HNO}_3)_m(\text{H}_2\text{O})_n\text{NO}_2^-$  ions with 32%. The rest corresponded to  $\text{HNO}_3^-$  and  $\text{OH}^-$  containing cluster ions with 8% and 2%, respectively.<sup>36,37</sup> In our current spectrum in (Fig. 2a), the ratio of integrated intensities of the ion series containing  $\text{NO}_2^-$  and  $\text{NO}_3^-$  is  $1.0 \pm 0.1$ , which is slightly different from the previous spectra<sup>36,37</sup> due to somewhat different expansion conditions as discussed in ESI.† Nevertheless, our main focus here is the qualitative change of the spectra with the addition of the DMA molecules.

The negative ion spectrum after the pickup is shown in the bottom panel (Fig. 2b). The change in the spectrum with pickup pressures was relatively abrupt and appeared just at higher  $p_c$ . Therefore only the spectrum at  $p_c = 8.4 \times 10^{-4}$  mbar is shown in (Fig. 2b), which corresponds to the positive ion spectrum in (Fig. 1b). Due to the scattering of the clusters from the beam, the ion abundances are very low at higher  $p_c$  pressures.

Clearly, new major series appear in the spectrum, which could be unambiguously assigned to  $(\text{DMA})_k(\text{HNO}_3)_m\text{NO}_3^-$  ions (triangles up). The series distinguished by different colors and connected by dashed lines always start with  $k = m$  ion followed by ions with more  $\text{HNO}_3$  than DMA molecules  $m > k$ . Series with  $k = 1-3$  and  $m$  up to  $k + 4$  and  $k + 3$  are indicated in (Fig. 2b). New series with DMA molecules and  $\text{NO}_2^-$  ions,  $(\text{DMA})_k(\text{HNO}_3)_m\text{NO}_2^-$ , could be identified in the spectrum too, however, at very low intensities. The series indicated by the orange diamonds corresponds to the  $k = 1$  case. The first member of this series corresponds to  $m = k + 1$ . A very weak  $k = 2$  series could be found upon a closer inspection (not indicated to avoid congestion).

Traces of the series from the top spectrum can be found in (Fig. 2b) too, and are labeled by the same symbols as in the top panel (circles and triangles down). We note that the relative intensities of the original  $\text{NO}_2^-$  and  $\text{NO}_3^-$  series change upon the DMA pickup: while in the pure cluster spectrum (a), the  $\text{NO}_2^-$  containing peaks (circles) are higher than the corresponding  $\text{NO}_3^-$  containing peaks (triangles down), in the bottom spectrum with DMA pickup (b) the ratio is reversed. The total integrated intensities of the  $\text{NO}_3^-$  ion containing series (including the ones with DMA) constitute over 80% of the mass spectrum, i.e., the ratio of the cluster ion series containing  $\text{NO}_3^- : \text{NO}_2^-$  is about  $4.3 \pm 0.5$  (the larger error bars are due to the lower intensity and ambiguity in the assignment of some less abundant peaks due to mass coincidences).

The positive ion mass spectra indicated no clear evidence for any mass peaks containing DMA,  $\text{HNO}_3$  and water. Since we start with strongly hydrated clusters that pick up DMA molecules, the question arises, if the water evaporates upon the DMA pickup and subsequent acid-base reaction, or if it is the result of the cluster ionization process eventually followed by ion-molecule reactions. Therefore, we seek for evidence of hydrated mixed

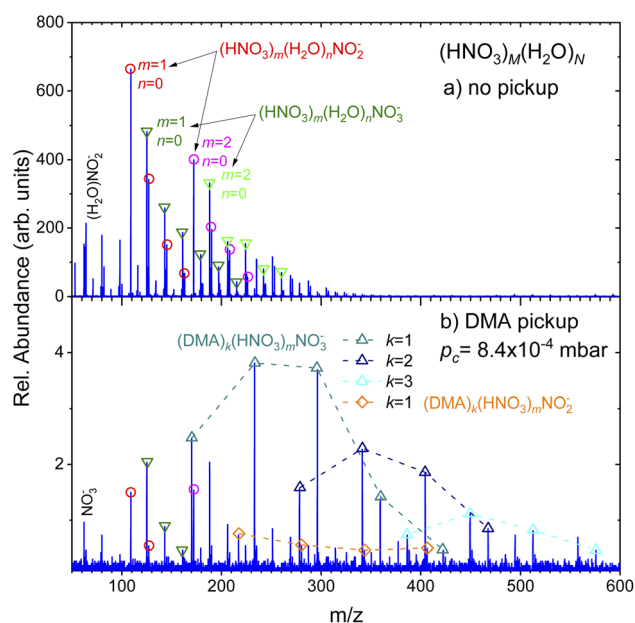


Fig. 2 Top panel (a) shows the negative ion mass spectra after the electron attachment to the pure  $(\text{HNO}_3)_M(\text{H}_2\text{O})_N$  clusters. Major ion series  $(\text{HNO}_3)_m(\text{H}_2\text{O})_n\text{NO}_2^-$  (circles) and  $(\text{HNO}_3)_m(\text{H}_2\text{O})_n\text{NO}_3^-$  (triangles down) are indicated. Bottom spectrum (b) shows the same clusters after pickup of DMA molecules (pressure in the pickup cell  $p_c = 8.4 \times 10^{-4}$  mbar). Here, the major series correspond to  $(\text{DMA})_k(\text{HNO}_3)_m\text{NO}_3^-$  (triangles up). A minor  $(\text{DMA})_k(\text{HNO}_3)_m\text{NO}_2^-$  series is indicated (diamonds).

clusters upon the electron attachment. The attachment might be a softer ionization leading less to water evaporation than the 70 eV electron impact ionization. There is a series of 3 peaks following the first  $(\text{DMA})_k(\text{HNO}_3)_m\text{NO}_3^-$ ,  $k = m = 1$ , which are displaced by  $m/z = 18$ . They could correspond to the hydrated cluster ions  $(\text{DMA})_k(\text{HNO}_3)_m(\text{H}_2\text{O})_n\text{NO}_3^-$ ,  $n = 1-3$ . Due to mass coincidences we discuss these peaks in more detail in ESI†. Nevertheless, we can safely conclude that the abundance of hydrated cluster ions is strongly suppressed with respect to the non-hydrated ones.

**3.2.1 Electron energy dependence.** Recording the negative ion mass spectra at the electron energies 0–14 eV in steps of 0.2 eV allows evaluation of the electron energy dependent ion yield for individual mass peaks. Examples of these spectra are shown in (Fig. 3) for a few selected cluster ions containing DMA molecules:  $(\text{DMA})_1(\text{HNO}_3)_1\text{NO}_3^-$ ,  $m/z = 170$  (green triangles up);  $(\text{DMA})_1(\text{HNO}_3)_2\text{NO}_3^-$ ,  $m/z = 233$  (red triangles down);  $(\text{DMA})_2(\text{HNO}_3)_3\text{NO}_3^-$ ,  $m/z = 341$  (blue opened diamonds). The spectra obtained for all other mass peaks with a sufficient signal-to-noise ratio are very similar.

For comparison, the grey line shows the  $E_e$ -dependence of  $(\text{HNO}_3)\text{NO}_3^-$ ,  $m/z = 125$  mass peak in the pure spectrum without pickup (note, that its intensity was scaled down by more than two orders of magnitude). This spectrum is identical with the  $E_e$ -dependencies of electron attachment to the pure  $(\text{HNO}_3)_M(\text{H}_2\text{O})_N$  clusters measured and analyzed in detail previously.<sup>37</sup> Thus, Fig. 3 demonstrates that the energy spectra of the DMA-containing clusters exhibit essentially the same character as the spectra of the pure clusters without the DMA pickup, and the electron attachment proceeds *via* the same mechanism to the pure  $(\text{HNO}_3)_M(\text{H}_2\text{O})_N$  as well as to the  $(\text{DMA})_K(\text{HNO}_3)_M(\text{H}_2\text{O})_N$  clusters with the adsorbed DMA molecules. Due to the relatively low signal-to-noise ratio of the negative ion signal and mass coincidences, no further

unambiguous conclusions could be drawn from the energy dependencies, and therefore we present somewhat extended discussion only in ESI† for interested reader.

## 4 Computational results

The calculations allow to disentangle the pickup of DMA on the hydrated acid clusters and subsequent positive or negative ionization.

### 4.1 Pickup of DMA on neutral clusters

We start with pickup of DMA on  $(\text{HNO}_3)(\text{H}_2\text{O})_5$  which corresponds to one of the more populated clusters in our experiment and is known to have acidically dissociated structure.<sup>36</sup> Fig. 4a and b show the binding energies of  $\text{H}_2\text{O}$  and DMA molecules, respectively, as a function of hydration with  $n$ -water molecules. Note that, throughout the whole manuscript, we use enthalpies rather than free energies<sup>52</sup> to allow for easier comparison with the experiment where the clusters have presumably a low temperature, though its exact value is challenging to determine experimentally. We use “binding strength” or “binding energy” without the negative sign, *i.e.*, the enthalpy change associated with the processes listed in captions of Figs. 4, 5 and 7. The

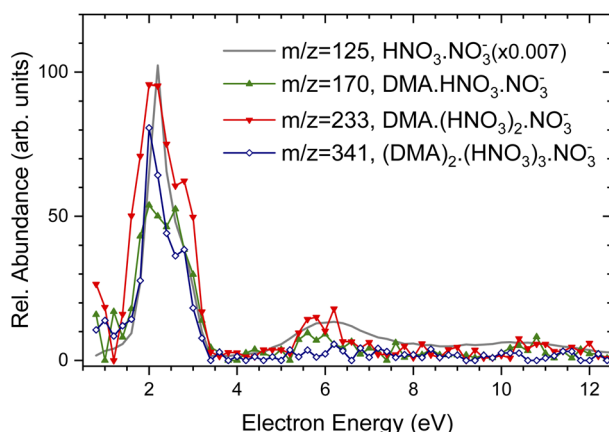


Fig. 3 Negative ion yield dependence on the electron energy for different mass peaks:  $(\text{DMA})_1(\text{HNO}_3)_1\text{NO}_3^-$ ,  $m/z = 170$ , green triangles up;  $(\text{DMA})_1(\text{HNO}_3)_2\text{NO}_3^-$ ,  $m/z = 233$ , red triangles down;  $(\text{DMA})_2(\text{HNO}_3)_3\text{NO}_3^-$ ,  $m/z = 341$ , blue opened diamonds. For comparison, the grey line shows the  $E_e$  dependence of  $(\text{HNO}_3)\text{NO}_3^-$ ,  $m/z = 125$  mass peak in the pure spectrum without pickup (intensity multiplied by a factor of 0.007).

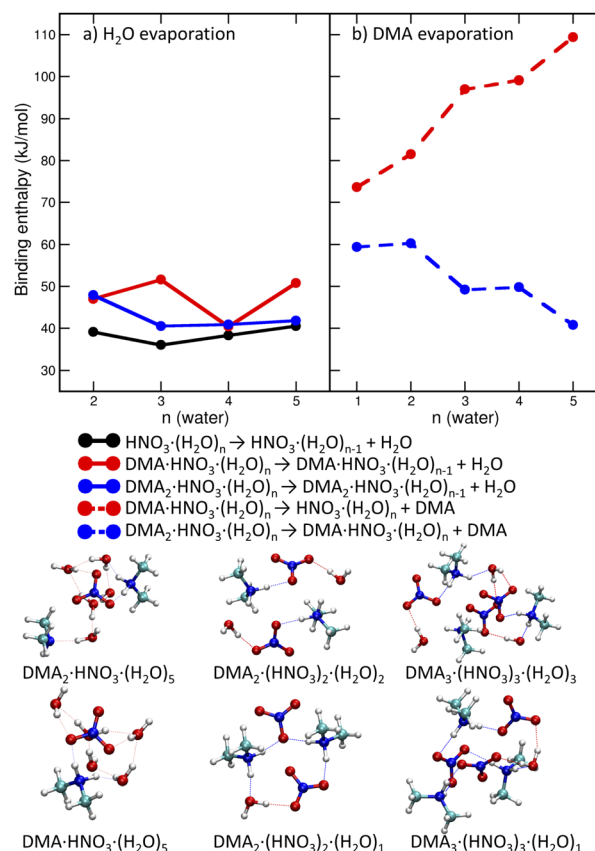


Fig. 4 Binding enthalpies of water molecules (a, solid lines) and DMA molecules (b, dashed lines) in  $\text{DMA}_K(\text{HNO}_3)(\text{H}_2\text{O})_N$  clusters as a function of number of water molecules. The bottom panel shows structures of selected neutral clusters.





attachment of a DMA molecule to the  $(\text{HNO}_3)(\text{H}_2\text{O})_n$  cluster is followed by an energy release of about  $100 \text{ kJ mol}^{-1}$  for a structure containing  $n = 3$  to 5 water molecules (Fig. 4b). This is due to the proton transfer from  $\text{H}_3\text{O}^+$  to DMA in the clusters with ion pair already generated,  $n = 4$  and 5. For  $n = 3$  the proton transfers directly from the  $\text{HNO}_3$  molecule. The binding strength of the first attached DMA increases with the cluster size because of the favorable hydration of the charged species. The trend for the second (neutral) DMA is the opposite and the binding energy of the second DMA to the clusters that are more populated in our experiment  $(\text{HNO}_3)(\text{H}_2\text{O})_n$ ,  $n = 3$ –5, is about  $45 \text{ kJ mol}^{-1}$ .

The left panel (Fig. 4a) illustrates that binding of a water molecule is about  $45 \text{ kJ mol}^{-1}$  nearly independently of the clusters size. Therefore, the pickup of DMA is sufficient to evaporate two water molecules. The pickup of the second DMA molecule is possible, but it releases less energy, because the molecule remains neutral and has a binding strength comparable to water.

DMA is a stronger base than  $\text{NO}_3^-$  or water, thus clusters with  $n \geq 1$  contain the  $\text{DMA} \cdot \text{H}^+ \cdots \text{NO}_3^-$  ion pair. This feature is preserved in larger structures ( $\text{DMA} : \text{HNO}_3$  2 : 2 or 3 : 3) which possess acidically dissociated structures even without any water (see Fig. 4 bottom).

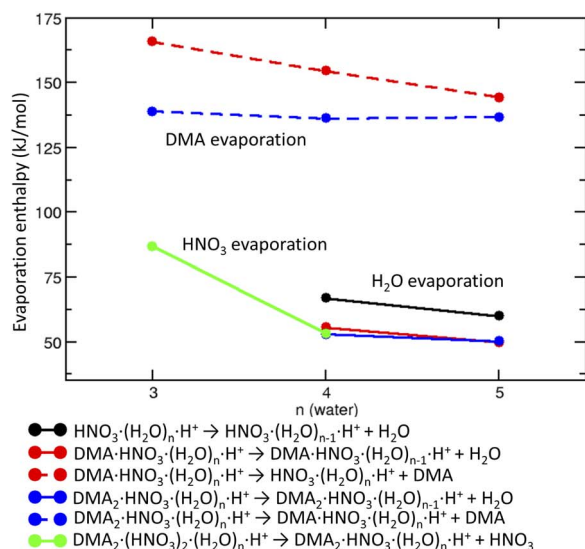


Fig. 5 Binding enthalpies of water molecules (solid lines), DMA molecules (dashed lines) and  $\text{HNO}_3$  (green line) in  $\text{DMA}_k(\text{HNO}_3)_M(\text{H}_2\text{O})_n\text{H}^+$  clusters as a function of number of water molecules (top). The bottom shows structures of selected protonated clusters.

## 4.2 Positively charged clusters

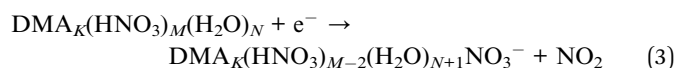
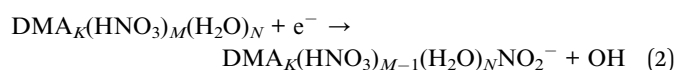
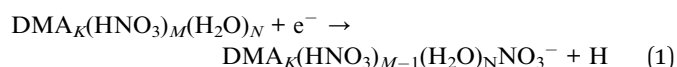
The second analyzed step was the ionization of clusters. The positive ionization of  $\text{DMA}_k(\text{HNO}_3)_M(\text{H}_2\text{O})_N$  by 70 eV electrons creates protonated clusters  $\text{DMA}_k(\text{HNO}_3)_M(\text{H}_2\text{O})_N\text{H}^+$ . The structures always contain protonated  $\text{DMA} \cdot \text{H}^+$ . In the cluster ions  $\text{DMA}(\text{HNO}_3)(\text{H}_2\text{O})_n\text{H}^+$  with less than  $n \leq 4$  water molecules, the neutral  $\text{HNO}_3$  is formed (see Fig. 5 bottom, right). More hydrated clusters  $\text{DMA}(\text{HNO}_3)(\text{H}_2\text{O})_n\text{H}^+$  ( $n \geq 5$ ) possess structures with a  $\text{DMA} \cdot \text{H}^+ \cdots \text{NO}_3^-$  ion pair and the additional proton resides on a water molecule forming  $\text{H}_3\text{O}^+$ . In case of  $(\text{DMA})_2(\text{HNO}_3)_2(\text{H}_2\text{O})_n\text{H}^+$ , structure containing  $\text{H}_3\text{O}^+$  is preferred for  $n \geq 3$  and neutral  $\text{HNO}_3$  is more likely in case of  $n \leq 2$  (see ESI†).

The experimental mass spectra contain predominantly the series  $(\text{DMA})_k(\text{HNO}_3)_m\text{H}^+$  ions with  $k = m + 1$  and small signals corresponding to  $k = m$ . This can be explained by the binding energy of each of the molecules (Fig. 5 top). The most likely molecule to evaporate upon the ionization is water with binding energy of around  $55 \text{ kJ mol}^{-1}$  in protonated clusters of different composition. The binding of water to the protonated clusters containing DMA (red and blue circles, full lines) is slightly weaker than to those without DMA (black circles). The overall positive charge of the cluster strengthens water binding by about  $10 \text{ kJ mol}^{-1}$  in comparison with the neutral structure. The second weakest bound moiety is nitric acid with about  $85 \text{ kJ mol}^{-1}$  cost for evaporation from  $\text{DMA}_2(\text{HNO}_3)_2(\text{H}_2\text{O})_3\text{H}^+$  (green line). The most tightly bound species is DMA (about  $150 \text{ kJ mol}^{-1}$ ), because it is protonated in all clusters. The binding energy of the first DMA is by about  $25 \text{ kJ mol}^{-1}$  higher than of the second one for  $n = 3$ . This difference decreases with increasing number of water molecules. Thus, evaporation of  $\text{HNO}_3$  from  $\text{DMA}_2(\text{HNO}_3)_2(\text{H}_2\text{O})_n\text{H}^+$  giving  $\text{DMA}_2(\text{HNO}_3)(\text{H}_2\text{O})_n\text{H}^+$  is plausible and corroborates with the dominant  $k = m + 1$  series in the mass spectra.

The tight binding of both DMA molecules in the positively charged clusters corresponds to the fact that in order to evaporate them, the extra proton has to be accepted by a weaker base at first. From structural point of view, the clusters maximizing the number of  $\text{DMA} \cdot \text{H}^+$  and  $\text{NO}_3^-$  ions are the optimal ones. This is fulfilled for  $K = M$  in the neutral clusters and for  $k = m + 1$  in the protonated ones (Fig. 5 left).

## 4.3 Negatively charged clusters

After a low energy electron attachment, the following chemical reactions releasing H, OH and  $\text{NO}_2$  radicals can occur:



Based on the reaction enthalpies (Fig. 6), the predominant pathway is the release of  $\text{NO}_2$  radical ( $\sim -130 \text{ kJ mol}^{-1}$ ) and formation of  $\text{NO}_3^-$ -containing ions. The reaction enthalpy does not depend significantly on the amount of DMA,  $\text{HNO}_3$  and  $\text{H}_2\text{O}$  in the clusters. The second, but less likely possibility, is the formation of OH radical and generation of  $\text{NO}_2^-$ -containing ions which is slightly exothermic ( $\sim -20 \text{ kJ mol}^{-1}$ ) in our clusters. The hydrogen radical is unlikely to form, because the reactions are endothermic ( $\sim 40 \text{ kJ mol}^{-1}$ ). Thus, the energetics suggests that the clusters containing  $\text{NO}_3^-$  will be more populated than those with  $\text{NO}_2^-$ .

If just the reactions triggered by electron attachment to the  $\text{DMA}_K(\text{HNO}_3)_M(\text{H}_2\text{O})_N$  clusters containing the same number of DMA and  $\text{HNO}_3$  molecules  $K = M$  were considered, the resulting structures should contain more DMA than  $\text{HNO}_3$ . This is not the case, since we observe  $\text{DMA}_k(\text{HNO}_3)_m\text{NO}_3^-$  cluster ions with  $m \geq k$ . Therefore, evaporation of molecules from the clusters generated immediately after the electron attachment has to be quantified too. The binding strengths of DMA and  $\text{H}_2\text{O}$  to the negatively charged clusters are similar to those in neutral clusters. The binding energy of a water molecule is about  $50 \text{ kJ mol}^{-1}$  (see ESI).<sup>†</sup> Binding of DMA in the intermediate clusters is on average by a few  $\text{kJ mol}^{-1}$  weaker (Fig. 7 black, blue and magenta lines). Clusters containing  $\text{NO}_3^-$  or  $\text{NO}_2^-$  do not differ in this aspect. It is likely that water molecules and DMA molecules which cannot be protonated evaporate leading to the cluster composition with  $m = k$  (see labels of the black, blue and magenta lines).

In contrast to that,  $\text{HNO}_3$  is tightly bound (at least  $150 \text{ kJ mol}^{-1}$ ) in the  $\text{DMA}_k(\text{HNO}_3)_m(\text{H}_2\text{O})_n\text{NO}_3^-$  or  $\text{DMA}_k(\text{HNO}_3)_m(\text{H}_2\text{O})_n\text{NO}_2^-$   $m = k$  clusters (Fig. 7 red). DMA

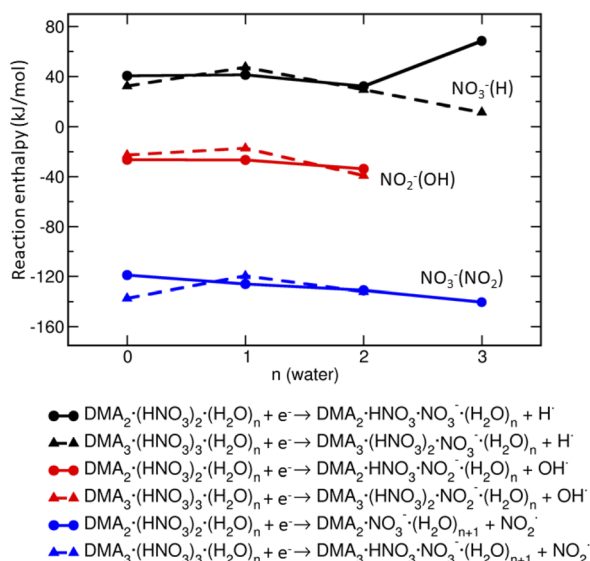


Fig. 6 Reactions enthalpies of potential processes occurring after electron attachment to  $\text{DMA}_K(\text{HNO}_3)_M(\text{H}_2\text{O})_N$ : formation of H radical with  $\text{NO}_3^-$  remaining in the cluster (black), OH radical with  $\text{NO}_2^-$  remaining in the cluster (red), and  $\text{NO}_2$  radical with  $\text{NO}_3^-$  remaining in the cluster (blue). Triangles correspond to the  $\text{DMA}_3(\text{HNO}_3)_3(\text{H}_2\text{O})_N$ , circles to the  $\text{DMA}_2(\text{HNO}_3)_2(\text{H}_2\text{O})_N$  clusters.

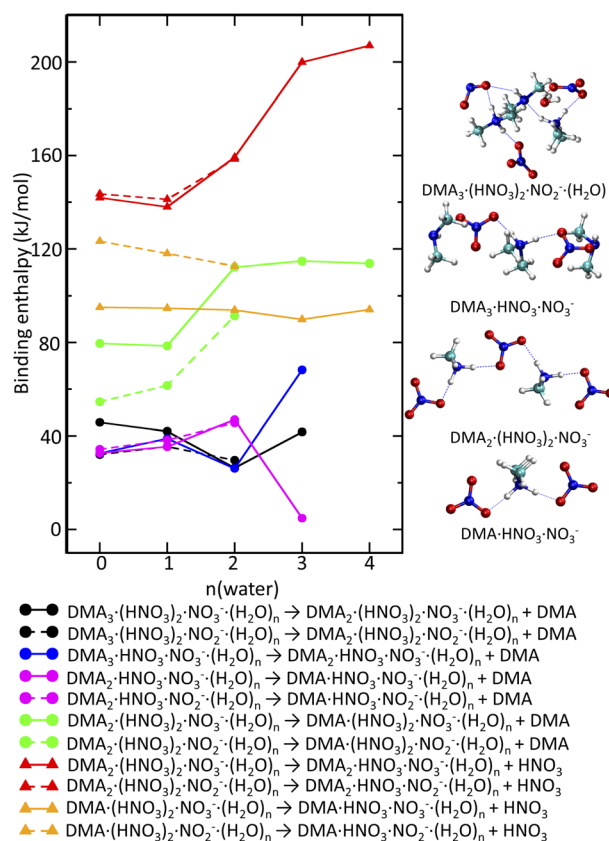


Fig. 7 Binding enthalpies of DMA (circles) and  $\text{HNO}_3$  (triangles) in  $\text{DMA}_K(\text{HNO}_3)_M(\text{H}_2\text{O})_N\text{NO}_3^-$  (solid lines) or  $\text{DMA}_K(\text{HNO}_3)_M(\text{H}_2\text{O})_N\text{NO}_2^-$  (dashed lines) clusters as a function of number of water molecules. Different colours correspond to different cluster composition. The right panel shows structures of selected negative clusters.

binding to the  $m = k$  cluster (Fig. 7 green) is also stronger than to the above mentioned  $m \leq k$  clusters, but the interaction is weaker than in case of  $\text{HNO}_3$ . The corresponding values increase from  $80 \text{ kJ mol}^{-1}$  for less hydrated clusters containing  $\text{NO}_3^-$  to  $120 \text{ kJ mol}^{-1}$  more hydrated clusters. The lower values for analogous process in  $\text{DMA}_k(\text{HNO}_3)_m(\text{H}_2\text{O})_n\text{NO}_2^-$  can be rationalized by the presence of  $\text{HONO}$  (not  $\text{NO}_2^-$ ) (see ESI).<sup>†</sup> Thus, the evaporation of DMA from the  $m = k$  clusters leading to the experimentally observed  $m \geq k$  series is more likely than the hypothetical  $\text{HNO}_3$  evaporation. Even structures with  $m \geq k$  bind  $\text{HNO}_3$  quite strongly, even though not as much as in the  $m = k$  case. The evaporation costs more than  $90 \text{ kJ mol}^{-1}$  (Fig. 7 orange).

Fig. 7 (right) shows some representative structures. The intermediate  $\text{DMA}_3(\text{HNO}_3)_3\text{NO}_3^-$  formed after electron attachment to the neutral  $\text{DMA}_3(\text{HNO}_3)_3$  clusters has one  $\text{DMA} \cdot \text{H}^+$  symmetrically surrounded by  $\text{NO}_3^-$  and two neutral DMA molecules. The facile evaporation of the later leads to  $\text{DMA}(\text{HNO}_3)_3\text{NO}_3^-$  which is in fact  $\text{DMA} \cdot \text{H}^+(\text{NO}_3^-)_2$ . The tendency to maximize the number of protonated  $\text{DMA} \cdot \text{H}^+$  and anionic units is preserved in negatively charged clusters too. It occurs for  $k = m$  in  $\text{DMA}_k(\text{HNO}_3)_m\text{NO}_3^-$ . In addition, the evaporation of



DMA from the  $m = k$  clusters is more facile than evaporation of  $\text{HNO}_3$ . The resulting  $m \geq k$  structures observed in the experimental spectrum are still relatively stable with respect to the  $\text{HNO}_3$  evaporation.

## 5 Discussion

We start with the composition of the clusters focusing on the  $\text{DMA}:\text{HNO}_3$  ratio detected in the mass spectra under different types of ionization. The positive ion mass spectra in Fig. 1 show that upon pickup of DMA molecules on the hydrated  $(\text{HNO}_3)_M(\text{H}_2\text{O})_N$  clusters and subsequent electron ionization, protonated clusters  $(\text{DMA})_k(\text{HNO}_3)_m\text{H}^+$  are generated with one more DMA molecule than  $\text{HNO}_3$ ,  $k = m + 1$ . Minor series with  $k = m$  is also discernible. On the other hand, electron attachment to the same neutral clusters yields mainly  $(\text{DMA})_k(\text{HNO}_3)_m\text{NO}_3^-$  series with  $m \geq k$ . Analogous series with  $\text{NO}_2^-$  and  $m > k$  are also present at very small abundances.

The calculations suggest that the optimal structures are those maximizing the number of  $\text{DMA}\cdot\text{H}^+$  and  $\text{NO}_3^-$  ions. This occurs for  $K = M$  in neutral clusters, for  $k = m + 1$  in protonated ones and for  $k = m$  in  $(\text{DMA})_k(\text{HNO}_3)_m\text{NO}_3^-$ . Fig. 7 rationalizes the experimental observation that the major peaks in the negative spectra correspond to  $m > k$  (specifically,  $m = k + 1$  and  $m = k + 2$ ), since additional DMA molecules can evaporate from the negatively charged clusters, while the evaporation of  $\text{HNO}_3$  is less likely, as discussed above. In addition, we have to explain why the clusters with  $m \geq k + 2$  are still observed experimentally. The experimental spectrum reflects the stability of the ions as well as the population of the original neutral species. If we start from neutral clusters with more  $\text{HNO}_3$  than DMA molecules  $M > K$ , it is easier to end up with the  $m > k$  negatively charged ionic clusters. It ought to be mentioned that the neutral  $(\text{DMA})_K(\text{HNO}_3)_M(\text{H}_2\text{O})_N$  clusters are generated by the pickup of  $K$  molecules of DMA on  $(\text{HNO}_3)_M(\text{H}_2\text{O})_N$ , and thus the initial  $K:M$  composition can be different from the optimal  $K = M$  neutral cluster structure, and the clusters with  $M > K$  can be more abundant. Upon the positive ionization, there is enough energy delivered in the ionization process to reach the optimal  $k = m + 1$  structure of the protonated  $(\text{DMA})_k(\text{HNO}_3)_m\text{H}^+$  cluster ions, while the clusters with  $m \geq k + 2$  are still observed upon the electron attachment.

Second point to discuss is the dehydration of clusters. We start the experiment with highly hydrated clusters before the DMA pickup. In the positive  $(\text{HNO}_3)_m(\text{H}_2\text{O})_n\text{H}^+$  spectrum of pure clusters without DMA pickup, the average cluster ion contains six to three  $\text{H}_2\text{O}$  molecules per one  $\text{HNO}_3$  molecule. In addition, at least OH and probably also some water molecules are evaporated upon the electron ionization at 70 eV. Therefore, the number of  $\text{H}_2\text{O}$  molecules in the neutral cluster is even larger than in the observed cluster ion fragment,  $N > n$ . Upon the DMA pickup mostly cluster ions without any water are observed in the mass spectrum. Does the water evaporate from the neutral cluster due to the acid–base reaction, or upon the subsequent ionization? The comparison of the binding energies suggest that the uptake of one DMA and subsequent acid–base reaction is sufficient to evaporate about two water molecules. In

addition, the average collision energy of DMA molecules with clusters in our experiments corresponds to about  $58 \text{ kJ mol}^{-1}$ , which could lead to another single water molecule evaporation. However, this is not quite sufficient for the complete dehydration of the clusters, thus, the remaining water molecules are evaporated by both, positive and negative ionization processes.

Finally, we concentrate on the effect of cluster medium observed in the negative ion spectra. Let us remind that the DEA of an isolated  $\text{HNO}_3$  molecule yielded essentially only  $\text{NO}_2^-$  ions with the relative abundance of 96.5% with almost no traces of  $\text{NO}_3^-$ .<sup>36</sup> Going from the gas phase  $\text{HNO}_3$  to the hydrated  $\text{HNO}_3$  clusters, our previous investigations<sup>36,37</sup> showed that the relative population of  $\text{NO}_3^-$  ions upon the electron attachment substantially increased. In the present experiment, the ratio of integrated intensities of the ion series containing  $\text{NO}_3^-$  and  $\text{NO}_2^-$  is  $1.0 \pm 0.1$ . Upon further addition of the DMA molecule, the propensity for the  $\text{NO}_3^-$  ion based clusters is even stronger with the ratio of the cluster ion series containing  $\text{NO}_3^-:\text{NO}_2^-$  of about  $4.3 \pm 0.5$ .

The change of the  $\text{NO}_3^-:\text{NO}_2^-$  ratio upon the uptake of DMA can be discussed in terms of a simple mechanism of electron attachment to an ion pair formed in cluster. If a neutral cluster contains an ion pair – either  $\text{H}_3\text{O}^+\cdots\text{NO}_3^-$  or  $\text{DMA}\cdot\text{H}^+\cdots\text{NO}_3^-$  – the electron is likely to attach to the protonated species leading to the hydrogen atom ejection from the cluster and the  $\text{NO}_3^-$ -containing cluster ion formation. For  $(\text{HNO}_3)_2(\text{H}_2\text{O})_N$  clusters with single  $\text{HNO}_3$  molecule, the  $\text{H}_3\text{O}^+\cdots\text{NO}_3^-$  ion pair formation occurs at  $n \geq 4$ .<sup>34–36,52,53</sup> Our present calculations indicate that the  $\text{DMA}\cdot\text{H}^+\cdots\text{NO}_3^-$  ion pair is already formed in clusters with a single water molecule because of DMA higher basicity. With addition of second  $\text{HNO}_3$  molecule, the number of water molecules required to form an ion pair is reduced even further. For example, with two  $\text{HNO}_3$  molecules only three water molecules are needed for the ion pair to be energetically preferred in  $(\text{HNO}_3)_2(\text{H}_2\text{O})_3$  cluster, and the presence of DMA in the cluster induces spontaneous proton transfer even without any water. Thus, the relative abundance of ion-pair containing clusters increases upon the DMA addition. Since they lead to the  $\text{NO}_3^-$ -containing cluster ions upon the electron attachment, the relative abundance of these ions increases upon the DMA addition, as observed experimentally.

However, this simple mechanism cannot explain the observation of the  $\text{NO}_2^-$  containing cluster ions, which are present in the spectra too, albeit at low intensities. Therefore, we present the argumentation based on the reaction energetics summarized in Fig. 6: the formation of  $\text{NO}_3^-$  is more exothermic than  $\text{NO}_2^-$ , but  $\text{NO}_2^-$ -containing clusters can be formed as well.

## 6 Conclusions

We have investigated the formation of atmospherically relevant  $(\text{DMA})_K(\text{HNO}_3)_M(\text{H}_2\text{O})_N$  clusters by uptake of DMA on  $(\text{HNO}_3)_M(\text{H}_2\text{O})_N$  and subsequent mass spectrometry analysis using both positive and negative ionization. Our experimental results aided by theoretical calculations indicate that:

- Optimal neutral cluster composition corresponds to DMA and  $\text{HNO}_3$  in 1 : 1 ratio. This is rationalized by the tendency to



maximize the number  $\text{DMA}\cdot\text{H}^+$  and  $\text{NO}_3^-$  ions in the clusters and by binding energies of cluster constituents.

- The dominant series in the positive mass spectra is  $k = m + 1$  which is nicely supported by calculated structures where the additional  $\text{H}^+$  resides on the extra DMA and prevents its evaporation.

- Upon an electron attachment, a formation of  $\text{NO}_3^-$  is more likely than  $\text{NO}_2^-$  which is due to the presence of the  $\text{DMA}\cdot\text{H}^+\cdots\text{NO}_3^-$  ion pair structures and reaction energetics (shown in Fig. 6). In addition, calculations give  $k = m$  as the optimal composition. The observation  $m > k$  is due to the fact that the negatively charged cluster evaporate DMA more easily than  $\text{HNO}_3$  according to the calculated evaporation enthalpies, and also due to the experimental neutral cluster composition.

An interesting point is the effect of cluster environment on the product ions observed in the electron attachment experiments. The population of  $\text{NO}_3^-$  and  $\text{NO}_2^-$  ions depends strongly on the medium. The dissociative electron attachment (DEA) to an isolated  $\text{HNO}_3$  molecule yields essentially exclusively  $\text{NO}_2^-$  ions, while the  $\text{NO}_3^-$  ion abundance increases with hydration: in the present pure  $\text{HNO}_3/\text{H}_2\text{O}$  clusters the ratio of integrated intensities of the ion series containing  $\text{NO}_3^-$  and  $\text{NO}_2^-$  ions is  $1.0 \pm 0.1$ . With addition of the strong base DMA, this ratio increased to about  $4.3 \pm 0.5$ . The propensity for  $\text{NO}_3^-$ -containing cluster ions can be explained by a simple model of electron attachment to the protonated species in the neutral clusters containing  $\text{H}_3\text{O}^+\cdots\text{NO}_3^-$  or  $\text{DMA}\cdot\text{H}^+\cdots\text{NO}_3^-$  ion pairs, which are generated due to the proton transfer in the acid–base reaction. Nevertheless, to justify the generation of both  $\text{NO}_3^-$  and  $\text{NO}_2^-$ -containing cluster ions, we calculate the reaction energetics, which explains the observed tendencies qualitatively.

We also demonstrate that the observed ionic species strongly depend on the ionization procedure. While the positive ion spectrum suggests prevalence of DMA molecules in the clusters, the negative spectrum clearly points to more  $\text{HNO}_3$  molecules in clusters. In addition, the observed major cluster ions do not contain any water, while the spectra of pure clusters without DMA pickup clearly exhibit a high level of hydration in the neutral clusters. Therefore, extreme caution has to be exercised when conclusions about neutral cluster composition are made from mass spectra, which should be considered in any implementation of mass spectrometry in aerosol research.

## Author contributions

A. Pysanenko: performed experiment and data evaluation; K. Fárniková: theoretical calculations and visualization; J. Lengyel: proposed idea for the experiment, helped data acquisition, interpretation and writing; E. Pluhařová: supervised and performed theory, interpretation and writing; M. Fárnik: supervised and designed experiment, interpretation, original draft and writing.

## Conflicts of interest

There are no conflicts to declare.

## Acknowledgements

The authors acknowledge support of the Czech Science Foundation (GAČR) project no.: 21-07062S. J. Lengyel acknowledges support of Deutsche Forschungsgemeinschaft (DFG, German Research Foundation; Project LE 4583/1-1). Computational resources were supplied by the project “e-Infrastruktura CZ” (e-INFRA CZ LM2018140) supported by the Ministry of Education, Youth and Sports of the Czech Republic.

## Notes and references

- 1 B. J. Finlayson-Pitts and J. N. Pitts, *Chemistry of the upper and lower atmosphere*, Academic Press, San Diego, 2000.
- 2 U. Pöschl, *Angew. Chem., Int. Ed.*, 2005, **44**, 7520–7540.
- 3 C. E. Kolb and D. R. Worsnop, *Annu. Rev. Phys. Chem.*, 2012, **63**, 471–491.
- 4 J. H. Seinfeld and S. N. Pandis, *Atmospheric Chemistry and Physics: From Air Pollution to Climate Change*, John Wiley & Sons Inc., Hoboken, NJ, 2016, p. 3rd ed.
- 5 B. J. Finlayson-Pitts and J. N. Pitts, *Science*, 1997, **276**, 1045–1051.
- 6 F. Bianchi, J. Tröstl, H. Junninen, C. Frege, S. Henne, C. R. Hoyle, U. Molteni, E. Herrmann, A. Adamov, N. Bukowiecki, X. Chen, J. Duplissy, M. Gysel, M. Hutterli, J. Kangasluoma, J. Kontkanen, A. Kürten, H. E. Manninen, S. Münch, O. Peräkylä, T. Petäjä, L. Rondo, C. Williamson, E. Weingartner, J. Curtius, D. R. Worsnop, M. Kulmala, J. Dommen and U. Baltensperger, *Science*, 2016, **352**, 1109–1112.
- 7 M. Kulmala, J. Kontkanen, H. Junninen, K. Lehtipalo, H. E. Manninen, T. Nieminen, T. Petäjä, M. Sipilä, S. Schobesberger, P. Rantala, A. Franchin, T. Jokinen, E. Järvinen, M. Äijälä, J. Kangasluoma, J. Hakala, P. P. Aalto, P. Paasonen, J. Mikkilä, J. Vanhanen, J. Aalto, H. Hakola, U. Makkonen, T. Ruuskanen, R. L. Mauldin, J. Duplissy, H. Vehkamäki, J. Bäck, A. Kortelainen, I. Riipinen, T. Kurtén, M. V. Johnston, J. N. Smith, M. Ehn, T. F. Mentel, K. E. J. Lehtinen, A. Laaksonen, V.-M. Kerminen and D. R. Worsnop, *Science*, 2013, **339**, 943–946.
- 8 J. Wang, J. F. Doussin, S. Perrier, E. Perraudin, Y. Katrib, E. Pangui and B. Picquet-Varraut, *Atmos. Meas. Tech.*, 2011, **4**, 2465–2494.
- 9 M. Wang, W. Kong, R. Marten, X.-C. He, D. Chen, J. Pfeifer, A. Heitto, J. Kontkanen, L. Dada, A. Kürten, T. Yli-Juuti, H. E. Manninen, S. Amanatidis, A. Amorim, R. Baalbaki, A. Baccarini, D. M. Bell, B. Bertozzi, S. Bräkling, S. Brilke, L. Caudillo Murillo, R. Chiu, B. Chu, L.-P. De Menezes, J. Duplissy, H. Finkenzeller, L. Gonzalez Carracedo, M. Granzin, R. Guida, A. Hansel, V. Hofbauer, J. Krechmer, K. Lehtipalo, H. Lamkaddam, M. Lampimäki, C. Ping Lee, V. Makhmutov, G. Marie, S. Mathot, R. L. Mauldin, B. Mentler, T. Müller, A. Onnela, E. Partoll, T. Petäjä, M. Philippov, V. Pospisilova, A. Ranjithkumar, M. Rissanen, B. Rörup, W. Scholz, J. Shen, M. Simon, M. Sipilä, G. Steiner, D. Stolzenburg, Y. Jun Tham,



- A. Tomé, A. C. Wagner, D. S. Wang, Y. Wang, S. K. Weber, P. M. Winkler, P. J. Wlasits, Y. Wu, M. Xiao, Q. Ye, M. Zauner-Wieczorek, X. Zhou, R. Volkamer, I. Riipinen, J. Dommen, J. Curtius, U. Baltensperger, M. Kulmala, D. R. Worsnop, J. Kirkby, J. H. Seinfeld, I. El-Haddad, R. C. Flagan and N. M. Donahue, *Nature*, 2020, **581**, 184–189.
- 10 H. Herrmann, T. Schaefer, A. Tilgner, S. A. Styler, C. Weller, M. Teich and T. Otto, *Chem. Rev.*, 2015, **115**, 4259–4334.
  - 11 Y. Huang, M. M. Coggon, R. Zhao, H. Lingell, M. U. Bauer, C. Flagan and J. H. Seinfeld, *Atmos. Meas. Tech.*, 2017, **10**, 839–867.
  - 12 H. Chen, S. Chee, M. J. Lawler, K. C. Barsanti, B. M. Wong and J. N. Smith, *Aerosol Sci. Technol.*, 2018, **4**, 1120–1133.
  - 13 S. Chee, N. Myllys, K. C. Barsanti, B. M. Wong and J. N. Smith, *J. Phys. Chem. A*, 2019, **123**, 5640–5648.
  - 14 R. Zhang, A. Khalizov, L. Wang, M. Hu and W. Xu, *Chem. Rev.*, 2012, **112**, 1957–2011.
  - 15 J. Elm, *J. Phys. Chem. A*, 2021, **125**, 895–902.
  - 16 T. Olenius, O. Kupiainen-Määttä, I. K. Ortega, T. Kurtén and H. Vehkamäki, *J. Chem. Phys.*, 2013, **139**, 084312.
  - 17 A. Kürten, T. Jokinen, M. Simona, M. Sipilä, N. Sarnel, H. Junninen, A. Adamov, J. Almeida, A. Amorim, F. Bianchi, M. Breitenlechner, J. Dommen, N. M. Donahue, J. Duplissy, S. Ehrhart, R. C. Flagan, A. Franchin, J. Hakala, A. Hansel, M. Heinritzi, M. Hutterli, J. Kangasluoma, J. Kirkby, A. Laaksonen, K. Lehtipalo, M. Leiminger, V. Makhmutov, S. Mathot, A. Onnela, T. Petäjä, A. P. Praplan, F. Riccobono, M. P. Rissanen, L. Rondo, S. Schobesberger, J. H. Seinfeld, G. Steiner, A. Tomé, J. Tröstl, P. M. Winkler, C. Williamson, D. Wimmer, P. Ye, U. Baltensperger, K. S. Carslaw, M. Kulmala, D. R. Worsnop and J. Curtius, *Proc. Natl. Acad. Sci. U. S. A.*, 2014, **111**, 15019–15024.
  - 18 J. Elm, T. Kurtén, M. Bilde and K. V. Mikkelsen, *J. Phys. Chem. A*, 2014, **118**, 7892–7900.
  - 19 J. Elm, N. Myllys, T. Olenius, R. Halonen, T. Kurtén and H. Vehkamäki, *Phys. Chem. Chem. Phys.*, 2017, **19**, 4877–4886.
  - 20 J. Elm, N. Myllys and T. Kurtén, *J. Phys. Chem. A*, 2017, **121**, 4578–4587.
  - 21 F. Bianchi, T. Kurtén, M. Riva, C. Mohr, M. P. Rissanen, P. Roldin, T. Berndt, J. D. Crounse, P. O. Wennberg, T. F. Mentel, J. Wildt, H. Junninen, T. Jokinen, M. Kulmala, D. R. Worsnop, J. A. Thornton, N. Donahue, H. G. Kjaergaard and M. Ehn, *Chem. Rev.*, 2019, **119**, 3472–3509.
  - 22 J. Kirkby, J. Curtius, J. Almeida, E. Dunne, J. Duplissy, S. Ehrhart, A. Franchin, S. Gagne, L. Ickes, A. Kürten, A. Kupc, A. Metzger, F. Riccobono, L. Rondo, S. Schobesberger, G. Tsagkogeorgas, D. Wimmer, A. Amorim, F. Bianchi, M. Breitenlechner, A. David, J. Dommen, A. Downard, M. Ehn, R. C. Flagan, S. Haider, A. Hansel, D. Hauser, W. Jud, H. Junninen, F. Kreissl, A. Kvashin, A. Laaksonen, K. Lehtipalo, J. Lima, E. R. Lovejoy, V. Makhmutov, S. Mathot, J. Mikkilä, P. Minginette, S. Mogo, T. Nieminen, A. Onnela, P. Pereira, T. Petäjä, R. Schnitzhofer, J. H. Seinfeld, M. Sipilä, Y. Stozhkov, F. Stratmann, A. Tomé, J. Vanhanen, Y. Viisanen, A. Vrtala, P. E. Wagner, H. Walther, E. Weingartner, H. Wex, P. M. Winkler, K. S. Carslaw, D. R. Worsnop, U. Baltensperger and M. Kulmala, *Nature*, 2011, **476**, 429–433.
  - 23 K. Lehtipalo, L. Rondo, J. Kontkanen, S. Schobesberger, T. Jokinen, N. Sarnela, A. Kürten, S. Ehrhart, A. Franchin, T. Nieminen, F. Riccobono, M. Sipilä, T. Yli-Juuti, J. Duplissy, A. Adamov, L. Ahlm, J. Almeida, A. Amorim, F. Bianchi, M. Breitenlechner, J. Dommen, A. J. Downard, E. M. Dunne, R. C. Flagan, R. Guida, J. Hakala, A. Hansel, W. Jud, J. Kangasluoma, V.-M. Kerminen, H. Keskinen, J. Kim, J. Kirkby, A. Kupc, O. Kupiainen-Määttä, A. Laaksonen, M. J. Lawler, M. Leiminger, S. Mathot, T. Olenius, I. K. Ortega, A. Onnela, T. Petäjä, A. Praplan, M. P. Rissanen, T. Ruuskanen, F. D. Santos, S. Schallhart, R. Schnitzhofer, M. Simon, J. N. Smith, J. Tröstl, G. Tsagkogeorgas, A. Tomé, P. Vaattovaara, H. Vehkamäki, A. E. Vrtala, P. E. Wagner, C. Williamson, D. Wimmer, P. M. Winkler, A. Virtanen, N. M. Donahue, K. S. Carslaw, U. Baltensperger, I. Riipinen, J. Curtius, D. R. Worsnop and M. Kulmala, *Nature*, 2016, **7**, 11594.
  - 24 J. Lengyel, A. Pysanenko and M. Fárnik, *Atmos. Chem. Phys.*, 2017, **17**, 14171–14180.
  - 25 T. Peter, *Annu. Rev. Phys. Chem.*, 1997, **48**, 785–822.
  - 26 G. Huang, X. Zhou, G. Deng, H. Qiao and K. Civerolo, *Atmos. Environ.*, 2002, **36**, 2225–2235.
  - 27 S. M. Murphy, A. Sorooshian, J. H. Kroll, N. L. Ng, P. Chhabra, C. Tong, J. D. Surratt, E. Knipping, R. C. Flagan and J. H. Seinfeld, *Atmos. Chem. Phys.*, 2007, **7**, 2313–2337.
  - 28 A. Sorooshian, S. M. Murphy, S. Hersey, H. Gates, L. T. Padro, A. Nenes, F. J. Brechtel, H. Jonsson, R. C. Flagan and J. H. Seinfeld, *Atmos. Chem. Phys.*, 2008, **8**, 5489–5520.
  - 29 B. R. Bzdek, D. P. Ridge and M. V. Johnston, *Atmos. Chem. Phys.*, 2010, **10**, 3495–3503.
  - 30 J. Kim, L. Ahlm, T. Yli-Juuti, M. Lawler, H. Keskinen, J. Tröstl, S. Schobesberger, J. Duplissy, A. Amorim, F. Bianchi, N. M. Donahue, R. C. Flagan, J. Hakala, M. Heinritzi, T. Jokinen, A. Kürten, A. Laaksonen, K. Lehtipalo, P. Miettinen, T. Petäjä, M. P. Rissanen, L. Rondo, K. Sengupta, M. Simon, A. Tomé, C. Williamson, D. Wimmer, P. M. Winkler, S. Ehrhart, P. Ye, J. Kirkby, J. Curtius, U. Baltensperger, M. Kulmala, K. E. J. Lehtinen, J. N. Smith, I. Riipinen and A. Virtanen, *Atmos. Chem. Phys.*, 2016, **16**, 293–304.
  - 31 M. J. Lawler, P. M. Winkler, J. Kim, L. Ahlm, J. Tröstl, A. P. Praplan, S. Schobesberger, A. Kürten, J. Kirkby, F. Bianchi, J. Duplissy, A. Hansel, T. Jokinen, H. Keskinen, K. Lehtipalo, M. Leiminger, T. Petäjä, M. Rissanen, L. Rondo, M. Simon, M. Sipilä, C. Williamson, D. Wimmer, I. Riipinen, A. Virtanen and J. N. Smith, *Atmos. Chem. Phys.*, 2016, **16**, 13601–13618.
  - 32 M. Fárnik and J. Lengyel, *Mass Spectrom. Rev.*, 2018, **37**, 630–651.





- 33 M. Fárník, J. Fedor, J. Kočišek, J. Lengyel, E. Pluhařová, V. Poterya and A. Pysanenko, *Phys. Chem. Chem. Phys.*, 2021, **23**, 3195–3213.
- 34 B. D. Kay, V. Hermann and A. W. Castleman, *Chem. Phys. Lett.*, 1981, **80**, 469–474.
- 35 J. Lengyel, A. Pysanenko, J. Kočišek, V. Poterya, C. Pradzynski, T. Zeuch, P. Slaviček and M. Fárník, *J. Phys. Chem. Lett.*, 2012, **3**, 3096–3109.
- 36 J. Lengyel, M. Ončák, J. Fedor, J. Kočišek, A. Pysanenko, M. K. Beyer and M. Fárník, *Phys. Chem. Chem. Phys.*, 2017, **19**, 11753–11758.
- 37 J. Lengyel, J. Fedor and M. Fárník, *Phys. Chem. Chem. Phys.*, 2019, **21**, 8691–8697.
- 38 A. Pysanenko, J. Lengyel and M. Fárník, *J. Chem. Phys.*, 2017, **148**, 154301.
- 39 J. Lengyel, A. Pysanenko, K. Fárníková, E. Pluhařová and M. Fárník, *J. Phys. Chem. Lett.*, 2020, **11**, 2101–2105.
- 40 J. VandeVondele, M. Krack, F. Mohamed, M. Parrinello, T. Chassaing and J. Hutter, *Comput. Phys. Commun.*, 2005, **167**, 103–128.
- 41 T. D. Kühne, M. Iannuzzi, M. D. Ben, V. V. Rybkin, P. Seewald, F. Stein, T. Laino, R. Z. Khaliullin, O. Schütt, F. Schiffmann, D. Golze, J. Wilhelm, S. Chulkov, M. H. Bani-Hashemian, V. Weber, U. Borštnik, M. Taillefumier, A. S. Jakobovits, A. Lazzaro, H. Pabst, T. Müller, R. Schade, M. Guidon, S. Andermatt, N. Holmberg, G. K. Schenter, A. Hehn, A. Bussy, F. Belleflamme, G. Tabacchi, A. Glöß, M. Lass, I. Bethune, C. J. Mundy, C. Plessl, M. Watkins, J. VandeVondele, M. Krack and J. Hutter, *J. Chem. Phys.*, 2020, **152**, 194103.
- 42 Y. Zhang and W. Yang, *J. Chem. Phys.*, 1998, **109**, 2604.
- 43 S. Grimme, J. Antony, S. Ehrlich and H. Krieg, *J. Chem. Phys.*, 2010, **132**, 154104.
- 44 S. Goedecker, M. Teter and J. Hutter, *Phys. Rev. B*, 1996, **54**, 1703–1710.
- 45 J. VandeVondele and J. Hutter, *J. Chem. Phys.*, 2007, **127**, 114105.
- 46 G. Bussi, D. Donadio and M. Parrinello, *J. Chem. Phys.*, 2007, **126**, 014101.
- 47 J. J. P. Stewart, *J. Mol. Model.*, 2007, **13**, 1173–1213.
- 48 Y. Zhao and D. G. Truhlar, *Theor. Chem. Acc.*, 2008, **120**, 215–241.
- 49 T. Dunning, *J. Chem. Phys.*, 1989, **90**, 1007–1023.
- 50 J. Elm, J. Kubečka, V. Besel, M. J. Jääskeläinen, R. Halonen, T. Kurtén and H. Vehkamäki, *J. Aerosol Sci.*, 2020, **149**, 105621.
- 51 M. J. Frisch, G. W. Trucks, H. B. Schlegel, G. E. Scuseria, M. A. Robb, J. R. Cheeseman, G. Scalmani, V. Barone, G. A. Petersson, H. Nakatsuji, X. Li, M. Caricato, A. V. Marenich, J. Bloino, B. G. Janesko, R. Gomperts, B. Mennucci, H. P. Hratchian, J. V. Ortiz, A. F. Izmaylov, J. L. Sonnenberg, D. Williams-Young, F. Ding, F. Lipparini, F. Egidi, J. Goings, B. Peng, A. Petrone, T. Henderson, D. Ranasinghe, V. G. Zakrzewski, J. Gao, N. Rega, G. Zheng, W. Liang, M. Hada, M. Ehara, K. Toyota, R. Fukuda, J. Hasegawa, M. Ishida, T. Nakajima, Y. Honda, O. Kitao, H. Nakai, T. Vreven, K. Throssell, J. A. Montgomery Jr, J. E. Peralta, F. Ogliaro, M. J. Bearpark, J. J. Heyd, E. N. Brothers, K. N. Kudin, V. N. Staroverov, T. A. Keith, R. Kobayashi, J. Normand, K. Raghavachari, A. P. Rendell, J. C. Burant, S. S. Iyengar, J. Tomasi, M. Cossi, J. M. Millam, M. Klene, C. Adamo, R. Cammi, J. W. Ochterski, R. L. Martin, K. Morokuma, O. Farkas, J. B. Foresman and D. J. Fox, *Gaussian ~16 Revision C.01*, 2016, Gaussian Inc. Wallingford CT.
- 52 P. R. McCurdy, W. P. Hess and S. S. Xantheas, *J. Phys. Chem. A*, 2002, **106**, 7628–7635.
- 53 K. R. Leopold, *Annu. Rev. Phys. Chem.*, 2011, **62**, 327–349.

