

## PAPER

[View Article Online](#)  
[View Journal](#) | [View Issue](#)Cite this: *J. Mater. Chem. C*, 2021,  
9, 270Stability of organic thin-film transistors based  
on ultrathin films of dinaphtho[2,3-*b*:2',3'-*f*]thieno[3,2-*b*]thiophene (DNNT)<sup>†</sup>Rachana Acharya,<sup>a</sup> Darius Günder,<sup>c</sup> Tobias Breuer,<sup>c</sup> Guido Schmitz,<sup>b</sup>  
Hagen Klauk<sup>a</sup> and Gregor Witte<sup>c</sup>

Organic thin-film transistors (TFTs) based on ultrathin semiconductor films are potentially useful as highly sensitive physical, chemical or biological sensors and may also help in the development of a better understanding of the relations between structural and charge-transport characteristics of thin films of organic semiconductors. A particularly promising small-molecule organic semiconductor is dinaphtho[2,3-*b*:2',3'-*f*]thieno[3,2-*b*]thiophene (DNNT). However, it was recently reported that ultrathin DNNT films spontaneously undergo dramatic morphological changes within minutes after deposition that lead to the disaggregation of the initially closed (or at least connected) single-monolayer films into disconnected multilayer islands. Here, we investigate how this spontaneous structural reconfiguration affects the characteristics of TFTs based on ultrathin DNNT films and explore the extent to which it can be prevented by cryogenic cooling or *in situ* encapsulation. We fabricated inverted coplanar TFTs with a hybrid aluminum oxide/alkylphosphonic acid self-assembled monolayer gate dielectric and vacuum-deposited DNNT films with nominal thicknesses of 2.5 or 25 nm. Using atomic force microscopy (AFM) we monitored the spontaneous changes in the DNNT morphology in a quasi-continuous manner over a period of 12 hours after deposition. The charge-carrier mobility of the ultrathin DNNT TFTs was found to decrease rapidly, while the mobility of the TFTs with the thicker DNNT films is far more stable. We also found that the initial closed-monolayer morphology of the ultrathin DNNT films is preserved when the substrates are cooled to cryogenic temperatures immediately after the DNNT deposition, but that the morphological changes resume upon returning the substrates to room temperature. Furthermore, we fabricated TFTs in which the ultrathin DNNT films were encapsulated *in situ* with a vacuum-deposited film of polytetrafluoroethylene, C<sub>60</sub> or titanyl phthalocyanine immediately following the DNNT deposition and found that the encapsulation decelerates the structural reorganization of the ultrathin DNNT films and the concurrent degradation of the carrier mobility.

Received 23rd September 2020,  
Accepted 23rd November 2020

DOI: 10.1039/d0tc04554c

[rsc.li/materials-c](http://rsc.li/materials-c)

## Introduction

Organic thin-film transistors (TFTs) are being developed for flexible electronics applications, such as rollable displays,<sup>1,2</sup> conformable sensors,<sup>3,4</sup> and electronic skin.<sup>5–9</sup> For the fabrication of TFTs based on vacuum-deposited small-molecule semiconductors, the active semiconductor film is typically deposited as a polycrystalline film from the vapor phase and usually has a nominal thickness of a few tens of nanometers, or approximately 10 to 20 molecular

monolayers. However, since the gate-field-induced charge-carrier channel is usually confined to the first one to two molecular monolayers closest to the interface with the gate dielectric,<sup>10</sup> it is also possible to fabricate functional TFTs based on vacuum-deposited organic-semiconductor films with a thickness of only a few nanometers.

The ability to fabricate stable and reliable organic TFTs with such ultrathin semiconductor films could lead to an improved understanding of the charge-transport physics in organic semiconductors and improve the sensitivity of organic-TFT-based sensors,<sup>11,12</sup> while reducing fabrication costs by minimizing material consumption.

Dinaphtho[2,3-*b*:2',3'-*f*]thieno[3,2-*b*]thiophene (DNNT) is a particularly promising small-molecule semiconductor for the fabrication of organic-TFT-based active-matrix displays,<sup>13,14</sup> sensor arrays,<sup>15,16</sup> and integrated circuits.<sup>17–20</sup> DNNT TFTs are

<sup>a</sup> Max Planck Institute for Solid State Research, Heisenbergstr. 1, 70569 Stuttgart, Germany. E-mail: R.Acharya@fkf.mpg.de, H.Klauk@fkf.mpg.de<sup>b</sup> University of Stuttgart, Heisenbergstr. 3, 70569 Stuttgart, Germany<sup>c</sup> Philipps-Universität Marburg, Renthof 7, 35032 Marburg, Germany<sup>†</sup> Electronic supplementary information (ESI) available. See DOI: 10.1039/d0tc04554c

p-channel transistors that typically have charge-carrier mobilities up to about  $3 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ .<sup>5,21,22</sup> In terms of stability against oxidation, chemical decomposition, elevated temperature and formation of defect states, DNTT is one of the most stable organic semiconductors.<sup>5,23–26</sup> DNTT films deposited by thermal sublimation in vacuum onto electrically insulating surfaces typically grow in a Stranski–Krastanov (layer-plus-island) mode, much like films based on other small-molecule organic semiconductors with similar chemical structures,<sup>27–29</sup> with the molecules adopting an upright-standing molecular orientation both in the complete monolayer(s) that initially cover(s) the substrate surface and in the islands that are formed beyond a critical film thickness to minimize mechanical strain.<sup>30</sup> This upright-standing orientation is generally the preferred molecular arrangement for the purpose of fabricating organic TFTs, as it promotes  $\pi$ -orbital overlap and efficient charge transport in the lateral direction.<sup>31,32</sup>

However, Breuer *et al.* recently observed that DNTT films with a nominal thickness of one to two molecular monolayers spontaneously undergo a rapid structural reconfiguration at room temperature within minutes after deposition that proceeds by lateral and vertical mass transport and leads to a rearrangement of the initially closed (or at least connected) single-monolayer-thick DNTT films into disconnected islands with a thickness of several layers.<sup>33</sup> In analogy to the behavior of thin liquid films on solid surfaces,<sup>34</sup> this process is also referred to as post-deposition dewetting<sup>29,30</sup> or rapid roughening,<sup>36</sup> and it has been observed in ultrathin films based on a variety of organic materials, such as polystyrene,<sup>30</sup> diindenoperylene<sup>36</sup> and pentacene.<sup>37,38</sup> When a molecular monolayer of, for example, DNTT is deposited onto a solid surface, this monolayer will form during a non-equilibrium phase of the film-growth process and will thus be characterized by significant mechanical strain that results from the competition between adsorption and intermolecular forces.<sup>39</sup> If the molecule–substrate interactions are weaker than the intermolecular interactions, thermally activated molecular diffusion will cause this initially closed (or at least connected) molecular monolayer to break up, and a combination of lateral and vertical mass transport will lead to the formation of tall, disconnected islands. This can be considered similar to the Ostwald ripening process often observed during annealing processes, where large, stable droplets grow at the expense of smaller, less stable droplets.<sup>40</sup> While temporal morphological changes have been observed also for thick DNTT films (with a nominal thickness of 40 nm),<sup>33</sup> these occur on much larger time scales of weeks or months, due to the substantially larger mass transport involved in such thick films, and will not be discussed in this work.

Here we investigate how the spontaneously occurring structural reconfiguration of ultrathin DNTT films affects the electrical characteristics of TFTs based on such films and explore the extent to which it can be prevented by cryogenic cooling or *in situ* encapsulation. For this purpose, we fabricated inverted coplanar (bottom-gate, bottom-contact) TFTs based on DNTT films deposited by thermal sublimation in vacuum onto hybrid gate dielectrics consisting of a thin film of plasma-grown

aluminum oxide ( $\text{AlO}_x$ ) and a self-assembled monolayer (SAM) of *n*-tetradecylphosphonic acid. Such ultrathin  $\text{AlO}_x$ /SAM gate dielectrics are a popular choice for the fabrication of low-voltage DNTT TFTs on a variety of substrates<sup>18,41–44</sup> (although in these previous demonstrations, the DNTT films usually had a thickness of several tens of nanometers).

Using atomic force microscopy (AFM), we monitored the evolution of the morphology of the ultrathin DNTT films in a quasi-continuous manner. Due to the extremely small thickness of the films and the fact that the morphological changes occur within minutes after deposition, it was not possible to utilize more elaborate characterization techniques, such as X-ray diffraction (XRD) or Kelvin probe force microscopy (KPFM), to elucidate the crystal structure or the defect density of these ultrathin films, but the results from the quasi-continuous AFM analysis presented here provide a fairly complete picture of the dynamics of the morphological changes. To monitor the degradation of the carrier mobility resulting from the changes in thin-film morphology, we repeatedly measured the current–voltage characteristics of the ultrathin DNTT TFTs over a period of 20 days after device fabrication. We also found that the closed-monolayer morphology of the as-deposited ultrathin DNTT films is preserved when the films are cooled to cryogenic temperatures immediately after the DNTT deposition, but that the structural reconfiguration resumes when the substrate is brought back to room temperature.

In addition, we evaluated the extent to which the morphology of the ultrathin DNTT films can be stabilized by *in situ* encapsulation. Our choice of the encapsulation materials was dictated by the requirement to deposit the encapsulation immediately (within one to two minutes) following the completion of the deposition of the ultrathin DNTT film in the same vacuum-deposition system and without breaking the vacuum. This ruled out virtually all popular organic-TFT-encapsulation materials, such as  $\text{Al}_2\text{O}_3$ ,<sup>45</sup> Cytop<sup>46</sup> and parylene,<sup>47</sup> as none of these can be deposited in the vacuum-deposition system in which the DNTT films are deposited. The three materials we thus selected for the *in situ* encapsulation of the ultrathin DNTT films are polytetrafluoroethylene (PTFE), Buckminster-Fullerene ( $\text{C}_{60}$ ) and titanyl phthalocyanine ( $\text{TiOPc}$ ).

Our investigation is motivated by the need for a better understanding of the dynamics of the morphological changes that occur in ultrathin films of vacuum-deposited small-molecule semiconductors and of the effects that these morphology changes have on the long-term stability of ultrathin devices, such as highly sensitive chemical, physical or biological sensors. An important prerequisite for this investigation is our ability to unambiguously assign the changes observed in the electrical properties of the ultrathin semiconductor films to the mass transport and resulting changes in morphology, as opposed to any other processes, such as any chemical reactions, *e.g.*, the oxidation or decomposition of the molecules during exposure to air and humidity. This makes DNTT the perfect choice for this study, as DNTT is one of the most stable organic semiconductors with respect to long-term exposure to ambient air and elevated temperatures.<sup>5,23–26</sup> Owing to this excellent chemical stability,



the changes in the measured electrical properties of the ultrathin DNTT films can be unequivocally assigned to the morphological changes observed by AFM. The ability to fabricate ultrathin organic TFTs with sufficient stability may also be helpful in developing a better understanding of the charge-transport physics in organic semiconductors using surface-sensitive techniques.

## Results and discussion

### Determining the minimum nominal DNTT thickness required for functional TFTs

To estimate the smallest nominal DNTT thickness for which functional TFTs can be obtained, we fabricated four samples in which DNTT films with nominal thicknesses of 1.0, 1.5, 2.0 and 2.5 nm were deposited onto  $\text{AlO}_x/\text{SAM}$  dielectrics by sublimation in vacuum with a rate of about  $2 \text{ nm min}^{-1}$ . Immediately after the DNTT deposition, the substrates were removed from the vacuum system, and the morphology of the DNTT films was examined by AFM in ambient air at room temperature. Fig. 1 shows AFM images of the nominally 1.0, 1.5, 2.0 and 2.5 nm thick DNTT films, each obtained within 15 minutes after deposition. The AFM images indicate that a nominal thickness of 1.0 nm results in disconnected islands, and although the islands appear to be more and more connected for nominal thicknesses of 1.5 and 2.0 nm, a closed film is clearly identified only for a nominal thickness of 2.5 nm.

Based on these results, we chose 2.5 nm as the minimum nominal DNTT thickness for the experiments described below. Based on the fact that the DNTT molecules have a length of approximately 1.6 nm and stand approximately upright on the substrate surface (as confirmed by the height profile shown in Fig. S1, ESI<sup>†</sup>),<sup>31</sup> a nominal thickness of 2.5 nm corresponds to less than two complete molecular monolayers. As the nominal DNTT thickness is increased from 1.0 to 2.5 nm, the RMS surface roughness decreases from 1.86 nm to 0.69 nm (see Table S1, ESI<sup>†</sup>), indicating a more planar and uniform film morphology with increasing nominal thickness.

### Changes in the morphology of ultrathin DNTT films observed by AFM

Having determined the minimum nominal DNTT thickness required for functional TFTs to be 2.5 nm, we then monitored

the time-dependent changes in the morphology of these films. For this purpose, AFM images were recorded on a nominally 2.5 nm-thick DNTT film in a quasi-continuous manner by repeatedly scanning the same area on the surface of the DNTT film over a period of 12 hours, making it possible to observe the process of the spontaneous structural reconfiguration essentially in real time. A time-lapse video is available in the ESI,<sup>†</sup> and Fig. 2 shows six snapshots obtained in intervals of three hours, with the first image taken within 15 minutes after completion of the DNTT deposition. To rule out that the morphological changes seen in these AFM images were caused by the force of the AFM tip, an identical substrate was prepared on which AFM images were taken only at 15 minutes after completion of the DNTT deposition and again at 12 hours after deposition, and these images indicate the same time-dependent changes in the morphology of the DNTT films.

The AFM images confirm the spontaneous rearrangement of the ultrathin DNTT films, as reported by Breuer *et al.*<sup>33</sup> The maximum height of the DNTT film evolves from 3.18 nm (about 2 molecular monolayers) 15 minutes after deposition to 10.9 nm (about 6 to 7 molecular monolayers) 12 hours after deposition, while the root-mean-square surface roughness increases from 0.69 to 2.61 nm (see Table S1, ESI<sup>†</sup>). This increase in the maximum height and RMS surface roughness is accompanied by the appearance of tall, disconnected islands and by a substantial reduction of the surface coverage, *i.e.*, by the appearance of larger and larger areas of the dielectric not covered with DNTT.

The measurements in Fig. 2 were performed on a DNTT film with a nominal thickness of 2.5 nm deposited onto an  $\text{AlO}_x/n$ -tetradecylphosphonic acid SAM dielectric. Since morphological changes in ultrathin DNTT films were previously observed also on other dielectrics,<sup>33</sup> we also performed AFM measurements on nominally 2.5 nm-thick DNTT films deposited onto bare  $\text{AlO}_x$  (without SAM) and onto  $\text{AlO}_x$  covered with a SAM of either *n*-hexylphosphonic acid or *n*-octadecylphosphonic acid.<sup>48</sup> The results of these measurements are summarized in Fig. S2 (ESI<sup>†</sup>). Although there are some minor differences in the initial morphology of the ultrathin DNTT films on these surfaces, presumably due to differences in the surface energy and the packing density of the SAMs,<sup>48–51</sup> the morphological changes in the ultrathin DNTT films are quite similar, suggesting that they

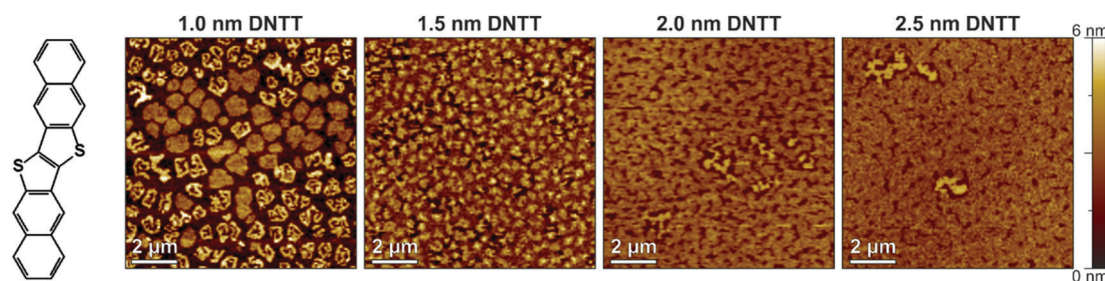


Fig. 1 Chemical structure of DNTT and thin-film morphology of DNTT films with nominal thicknesses of 1.0, 1.5, 2.0 and 2.5 nm deposited by thermal sublimation in vacuum onto a hybrid  $\text{AlO}_x/n$ -tetradecylphosphonic acid SAM dielectric, observed by AFM within 15 minutes after completion of the DNTT deposition.





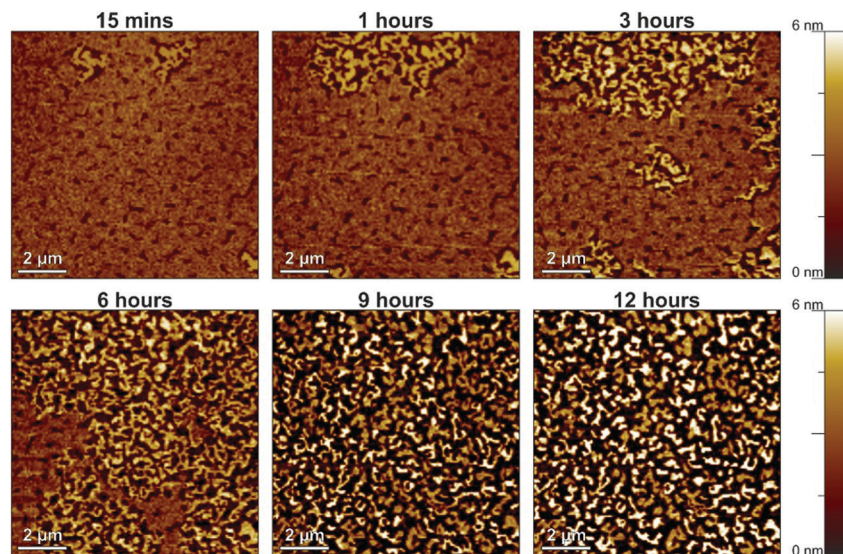


Fig. 2 Spontaneously occurring morphological changes in a DNTT film with a nominal thickness of 2.5 nm deposited by thermal sublimation in vacuum onto a hybrid  $\text{AlO}_x/n$ -tetradecylphosphonic acid SAM dielectric, observed by AFM over a period of 12 hours. A full-length video is available in the ESI† and results for DNTT films deposited onto bare  $\text{AlO}_x$  (without SAM) and onto  $\text{AlO}_x$  covered with a SAM of either *n*-hexylphosphonic acid or *n*-octadecylphosphonic acid are shown in Fig. S2 (ESI†).

are not significantly affected by the surface energy of the substrate. Results of X-ray diffraction measurements performed on DNTT films with a nominal thickness of 10 nm, summarized in Fig. S3 (ESI†), confirm that the DNTT molecules adopt an upright-standing molecular orientation.<sup>33</sup>

#### Influence of the changes in the DNTT thin-film morphology on the TFT characteristics

To investigate how the spontaneous morphological changes in the DNTT films described above affect the performance of TFTs in which such films serve as the semiconductor, we fabricated TFTs in the inverted coplanar (bottom-gate, bottom-contact) architecture. The TFTs were fabricated on oxidized silicon substrates to minimize the surface roughness. Aluminum gate electrodes with a root-mean-square surface roughness of less than 1 nm were deposited by thermal evaporation in vacuum.<sup>52</sup> The gate dielectric consists of plasma-grown  $\text{AlO}_x$  and an *n*-tetradecylphosphonic acid SAM.<sup>42</sup> Gold source/drain contacts were deposited by thermal evaporation in vacuum and functionalized with a monolayer of pentafluorobenzenethiol (PFBT) to minimize the contact resistance.<sup>53</sup> DNTT was deposited by vacuum sublimation with a rate of about  $2 \text{ nm min}^{-1}$ .<sup>44</sup> The deposition rate and the nominal thickness of the DNTT films were measured using a quartz-crystal microbalance positioned in close proximity to the substrate holder. TFTs with a nominal DNTT thickness of either 2.5 or 25 nm were fabricated. The current-voltage characteristics of the TFTs were measured on a probe station in ambient air at room temperature. The first measurement was performed within 15 minutes after the completion of the DNTT deposition, and the substrates were stored in ambient air at room temperature between measurements.

The results are summarized in Fig. 3. The initial TFT characteristics are quite similar, with effective charge-carrier

mobilities and threshold voltages of  $0.4 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$  and  $-1.5 \text{ V}$  for a nominal DNTT thickness of 2.5 nm (Fig. 3a and Fig. S4a, ESI†) and  $1 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$  and  $-1.7 \text{ V}$  for a nominal thickness of 25 nm (Fig. 3b and Fig. S4b, ESI†). This confirms that it is indeed possible to fabricate functional TFTs based on vacuum-deposited DNTT films with a nominal thickness of less than two molecular monolayers.

However, the evolution of the effective charge-carrier mobility is dramatically different for the two nominal DNTT thicknesses. While the effective charge-carrier mobility of the TFTs with the thick DNTT film decreases by less than 20% over a period of 20 days, the effective carrier mobility of the TFTs with the ultrathin DNTT film decreases far more rapidly and far more dramatically, from initially  $0.4$  to  $0.0007 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$  (a 99.9% decrease) within five days (Fig. 3c). The rapid decrease of the effective carrier mobility of the TFTs with the ultrathin DNTT film seen in Fig. 3c is consistent with the rapid reduction of the surface coverage of the semiconductor and the degree of percolation through the ultrathin DNTT film during the spontaneous rearrangement of the initially closed (or at least connected) ultrathin film into disconnected islands seen in Fig. 2. Although the morphological changes commence immediately after the completion of the DNTT deposition, their cumulative impact on the surface coverage and the degree of percolation appears to be most pronounced during the period between about one hour and about six hours after the DNTT deposition, which explains the observation that the drop in the effective charge-carrier mobility is most prominent during this period. The results confirm the severe impact of the spontaneous morphological changes of ultrathin DNTT films reported by Breuer *et al.*<sup>33</sup> on the effective charge-transport characteristics.

The thick DNTT films are far more stable than the ultrathin DNTT films, both in terms of the film morphology as observed



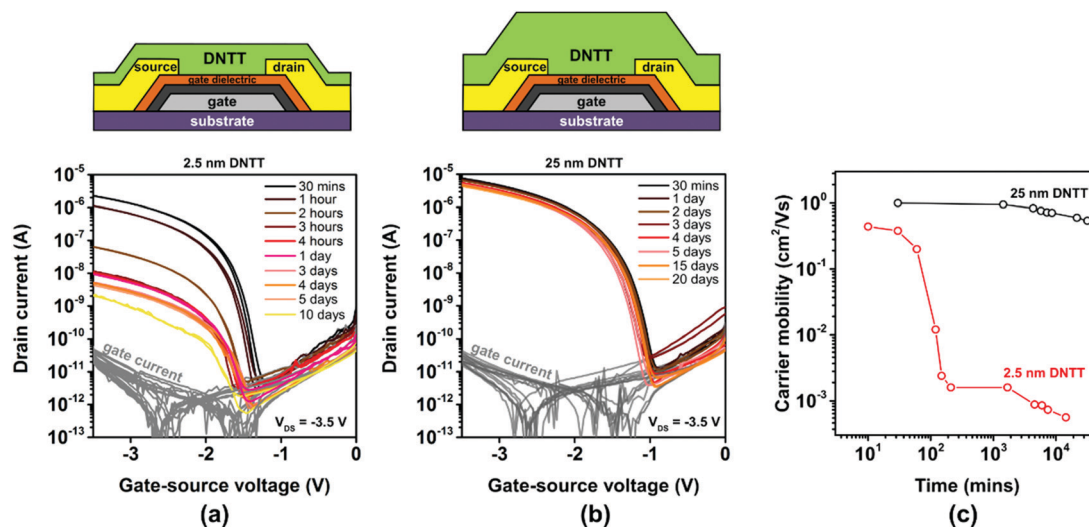


Fig. 3 Influence of the nominal thickness of vacuum-deposited DNTT films on the evolution of the electrical characteristics of TFTs based on these films. (a) Measured transfer curves of TFTs with an ultrathin DNTT film (nominal thickness of 2.5 nm) measured repeatedly over a period of 10 days after TFT fabrication. (b) Transfer curves of TFTs with a thick DNTT film (nominal thickness of 25 nm) measured over a period of 20 days after fabrication. (c) Evolution of the effective charge-carrier mobility extracted from the measured transfer curves. All TFTs have a channel length of 20  $\mu\text{m}$  and a channel width of 100  $\mu\text{m}$ .

by AFM (see Fig. S5, ESI<sup>†</sup>) and in terms of the electrical TFT characteristics (see Fig. 3b). This was already reported by Breuer *et al.*<sup>33</sup> and indicates that the spontaneous disaggregation of the first molecular monolayer into disconnected islands seen in Fig. 2 is effectively prevented when this first molecular monolayer is covered with a sufficient amount of additional material.

#### Effect of cryogenic cooling of ultrathin DNTT films

To confirm that the spontaneously occurring mass transport within the ultrathin DNTT films that leads to the disaggregation of the initially closed monolayer is indeed thermally activated, we cooled a substrate with an ultrathin DNTT film to a temperature of 193 K in dry nitrogen immediately after completion of the DNTT deposition and stored it at this temperature for a duration of 18 days. After allowing the substrate to warm up to room temperature, we immediately recorded an AFM image. An identical substrate with an identical ultrathin DNTT film served as a reference on which an AFM image was recorded immediately after the DNTT deposition and which was kept in ambient air at room temperature for a period of two weeks, after which another AFM image was taken. The results are summarized in Fig. 4.

Fig. 4a shows the AFM image recorded immediately after the DNTT deposition, indicating the initial closed-monolayer morphology (with a second monolayer seen as yet-disconnected islands). Fig. 4b confirms that this closed-monolayer morphology is perfectly preserved during storage at cryogenic temperature for 18 days. For comparison, Fig. 4c shows the break-up of the initially closed DNTT monolayer into tall, disconnected islands observed on the reference substrate after it had been stored at room temperature for two weeks (and also observed on the substrate that had been stored at cryogenic temperature within less than a day after it had been returned to room temperature). These observations strongly support the hypothesis

that these morphological changes in the ultrathin DNTT films are indeed thermally activated.<sup>35</sup>

#### In situ encapsulation of ultrathin DNTT films

As discussed in the Introduction, the morphological changes in the ultrathin DNTT films proceed by lateral and vertical mass transport that causes a rearrangement of the initially closed (or at least connected) single-monolayer-thick films into disconnected islands. While the results in Fig. 4 confirm that this mass transport is indeed thermally activated, the results in Fig. 3b show that it can be prevented by covering the ultrathin films with additional material (in the case of Fig. 3b with additional DNTT). In the following we explore the extent to which this can also be accomplished using other encapsulation materials. Our choice of these materials was dictated by the requirement to deposit the encapsulation immediately (within one to two minutes) following the completion of the deposition of the ultrathin DNTT film in the same vacuum-deposition system and without breaking the vacuum. This ruled

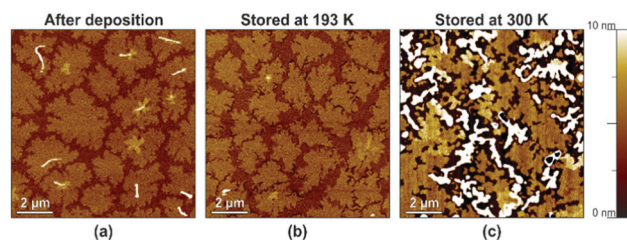


Fig. 4 Effect of cryogenic cooling on the spontaneously occurring morphological changes in ultrathin DNTT films, observed by AFM at room temperature. (a) DNTT morphology immediately after completion of the DNTT deposition. (b) DNTT morphology after the substrate had been stored at a temperature of 193 K for 18 days. (c) DNTT morphology on a reference substrate stored at room temperature for two weeks.



out virtually all popular organic-TFT-encapsulation materials, such as  $\text{Al}_2\text{O}_3$ ,<sup>45</sup> Cytop<sup>46</sup> and parylene,<sup>47</sup> as none of these can be deposited in the vacuum-deposition system in which the DNTT films are deposited. The three materials we selected for the *in situ* encapsulation of ultrathin DNTT films are polytetrafluoroethylene (PTFE),  $\text{C}_{60}$  and titanyl phthalocyanine (TiOPc).

### Encapsulation using vacuum-deposited polytetrafluoroethylene (PTFE)

Polytetrafluoroethylene (PTFE) is a synthetic fluoropolymer available as a powder that can be deposited by thermal sublimation in vacuum. PTFE is an excellent electrical insulator and has been used previously to encapsulate pentacene TFTs.<sup>54</sup> We deposited PTFE films with a nominal thickness of either 2.5 or 10 nm directly onto ultrathin DNTT films having a nominal thickness of 2.5 nm. The PTFE deposition was started within one to two minutes after completion of the DNTT deposition in the same vacuum-deposition system without breaking vacuum, by using one sublimation source for the DNTT deposition and another for the PTFE deposition. AFM images of a nominally 2.5 nm-thick PTFE film deposited onto a nominally 2.5 nm-thick DNTT film recorded within 15 minutes after the completion of the PTFE deposition and again after 72 hours are shown in Fig. S6b (ESI†). The AFM images indicate the occurrence of time-dependent morphological changes, but since it is not possible to distinguish the two materials in the AFM images, it is difficult to draw any useful conclusions from the images alone.

The results of the TFT measurements are summarized in Fig. 5. As can be seen, the initial TFT characteristics are not significantly affected by the PTFE encapsulation, but the degradation of the charge-carrier mobility is notably decelerated by the PTFE encapsulation, especially for the larger PTFE thickness (10 nm). However, even with the nominally 10 nm-thick PTFE encapsulation, the charge-carrier mobility of the TFTs still decreases by approximately one order of magnitude (from 0.8 to 0.08  $\text{cm}^2 \text{V}^{-1} \text{s}^{-1}$ )

within two months. Analysis of thin PTFE films deposited onto other surfaces suggests that PTFE films also undergo morphological changes (as seen in Fig. S7 and S8, ESI†),<sup>55,56</sup> and assuming these changes also occur in PTFE films deposited onto DNTT, this might explain the limited effectiveness of the PTFE encapsulation observed here. Nevertheless, as it is apparent that increasing the PTFE thickness from 2.5 to 10 nm increases the encapsulation effectiveness, we attempted the deposition of PTFE encapsulations with a nominal thickness of more than 10 nm, but this turned out to be quite challenging and prohibitively time-consuming in our deposition system, due to the unfavorable combination of the high vapor pressure of the PTFE powder and the limited throughput of the vacuum pump.

### Encapsulation using vacuum-deposited $\text{C}_{60}$

As an alternative to PTFE, we tested  $\text{C}_{60}$  for the encapsulation of ultrathin DNTT films. The vapor pressure of  $\text{C}_{60}$  is substantially smaller than that of PTFE, so that thick  $\text{C}_{60}$  films can be easily deposited by sublimation in vacuum. Other potential benefits of  $\text{C}_{60}$  for the purpose of encapsulating ultrathin DNTT films are the superior thermal and mechanical stability of  $\text{C}_{60}$  films<sup>57</sup> and the fact that  $\text{C}_{60}$  does not form undesirable Diels–Alder adducts with DNTT.<sup>58</sup> A substantial drawback of  $\text{C}_{60}$  for this purpose is the fact that it is an electron-transporting semiconductor and thus expected to adversely affect the current–voltage characteristics of the DNTT TFTs, especially their off-state behavior, due to the uncontrollable flow of electrons from the drain contact through the  $\text{C}_{60}$  film to the source contact. We deposited  $\text{C}_{60}$  films with a nominal thickness of either 2.5 or 25 nm directly onto ultrathin DNTT films having a nominal thickness of 2.5 nm.

The results are summarized in Fig. 6. As can be seen in Fig. 6b, the deposition of a nominally 25 nm-thick film of  $\text{C}_{60}$  indeed causes ambipolar transistor behavior, which is highly undesirable, as it increases the off-state leakage current of the

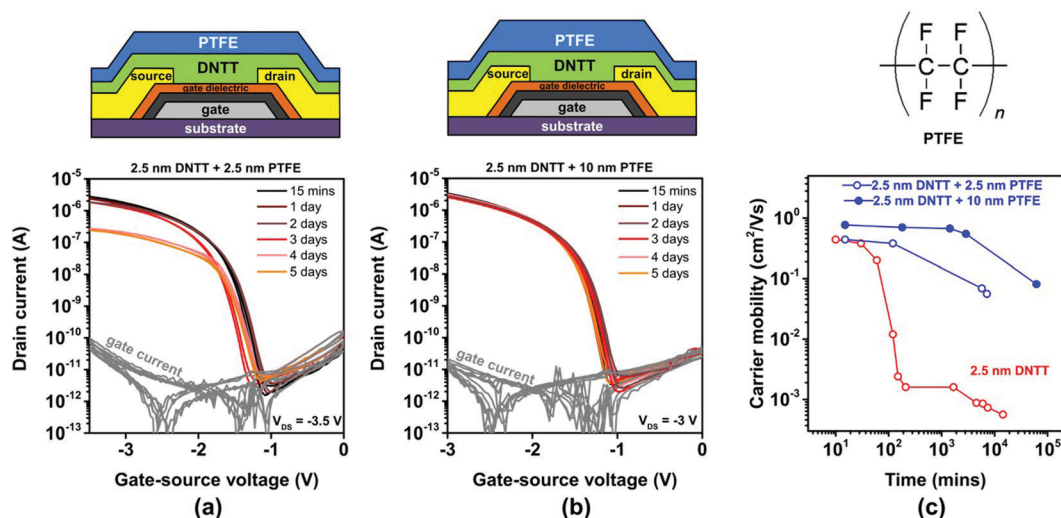


Fig. 5 Effect of *in situ* encapsulation of ultrathin, vacuum-deposited DNTT films (having a nominal thickness of 2.5 nm) with vacuum-deposited PTFE on the evolution of the TFT characteristics. Transfer curves of TFTs encapsulated with a nominally (a) 2.5 nm-thick and (b) 10 nm-thick PTFE film, measured repeatedly over a period of 5 days after TFT fabrication. (c) Evolution of the effective charge-carrier mobility extracted from the measured transfer curves.





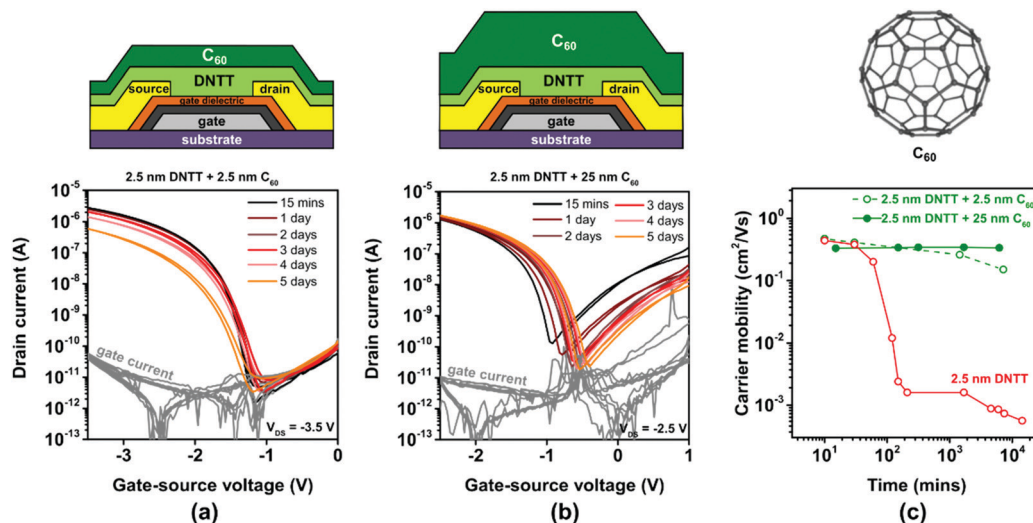


Fig. 6 Effect of *in situ* encapsulation of ultrathin, vacuum-deposited DNTT films (having a nominal thickness of 2.5 nm) with vacuum-deposited  $C_{60}$  on the evolution of the TFT characteristics. Transfer curves of TFTs encapsulated with a nominally (a) 2.5 nm-thick and (b) 25 nm-thick  $C_{60}$  film, measured repeatedly over a period of 5 days after TFT fabrication. (c) Evolution of the effective charge-carrier mobility extracted from the measured transfer curves.

TFTs and thus the static power consumption and the noise margins of logic circuits based on ambipolar TFTs. However, the degradation of the hole mobility in the ultrathin DNTT film is notably decelerated by the  $C_{60}$  encapsulation, which suggests that the morphology of the ultrathin DNTT film is indeed greatly stabilized by the  $C_{60}$  encapsulation.

A workable compromise between the beneficial effect of stabilizing the morphology of the ultrathin DNTT films and the drawback of increasing the leakage current is the use of a  $C_{60}$  encapsulation with a smaller thickness, as shown in Fig. 6a. With a nominally 2.5 nm-thick  $C_{60}$  encapsulation, the off-state drain current of the ultrathin-DNTT TFTs is approximately  $10^{-10}$  A, which is similar to that of ultrathin-DNTT TFTs without encapsulation (Fig. 3a) or with an electrically insulating encapsulation (Fig. 5a). However, the suppression of the degradation of the hole mobility in the ultrathin DNTT film is not quite as pronounced as in the case of the nominally 25 nm-thick  $C_{60}$  encapsulation (Fig. 6c). This is mainly due to the fact that the deposition of nominally 2.5 nm-thick  $C_{60}$  does not lead to the formation of a closed  $C_{60}$  film on the surface of the ultrathin DNTT film, but rather to edge decoration, as seen in Fig. S6c (ESI<sup>†</sup>) (and as reported previously for  $C_{60}$  deposited onto pentacene<sup>59</sup>).

### Encapsulation using vacuum-deposited titanyl phthalocyanine (TiOPc)

TiOPc is a hole-transporting small-molecule semiconductor.<sup>60</sup> TiOPc molecules usually adopt a flat-lying orientation when deposited onto metals,<sup>61</sup> but an upright-standing molecular orientation when deposited onto electrically insulating surfaces or organic-semiconductor films.<sup>62</sup> A potential benefit of TiOPc for the purpose of encapsulating ultrathin DNTT films is its excellent chemical robustness and thermal stability. A drawback of TiOPc for this purpose is that we cannot rule out that the TiOPc film will transport a hole current in parallel to the

ultrathin DNTT film, thus making it impossible to monitor the electrical behavior of the DNTT film exclusively. We deposited TiOPc films with a nominal thickness of either 2.5 or 25 nm directly onto ultrathin DNTT films.

The results are summarized in Fig. 7. The first observation is that the current-voltage characteristics of the ultrathin-DNTT TFTs with the TiOPc encapsulation (Fig. 7a and b) are exactly the same as those of the ultrathin-DNTT TFTs without encapsulation (Fig. 3a). This suggests that the TiOPc film does not participate in the charge transport, which is in contradiction to the results reported previously for TFTs based on a heterostructure of ultrathin dihexyl-sexithiophene and TiOPc, but this may be due to the fact that the carrier mobility in DNTT is significantly larger than that in both TiOPc and dihexyl-sexithiophene.

The second observation is that the TiOPc encapsulation slows down the degradation of the charge-carrier mobility of the TFTs, similar to encapsulation with PTFE and  $C_{60}$ . However, in stark contrast to PTFE (Fig. 5) and  $C_{60}$  (Fig. 6), this effect is virtually independent of the thickness of the TiOPc encapsulation. Nevertheless, the charge-carrier mobility of the TiOPc-encapsulated TFTs still decreases by approximately one order of magnitude within two weeks.

### Summary of encapsulation results

The results from the encapsulation experiments are summarized in Fig. 8. The transfer curves of the TFTs with and without the various encapsulation films recorded within 30 minutes after the last deposition and the evolution of the subthreshold swings extracted from these transfer curves are summarized in Fig. S9 (ESI<sup>†</sup>). The TFTs with the thick DNTT films (nominal thickness 25 nm) initially have an effective carrier mobility of  $1 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$  (Fig. 8a), which is similar to previous reports of bottom-gate, bottom contact DNTT TFTs.<sup>13</sup> The effective mobility of these TFTs degrades from 1 to  $0.8 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$  within 5 days and to  $0.5 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$  within 23 days (see also Fig. 3c).



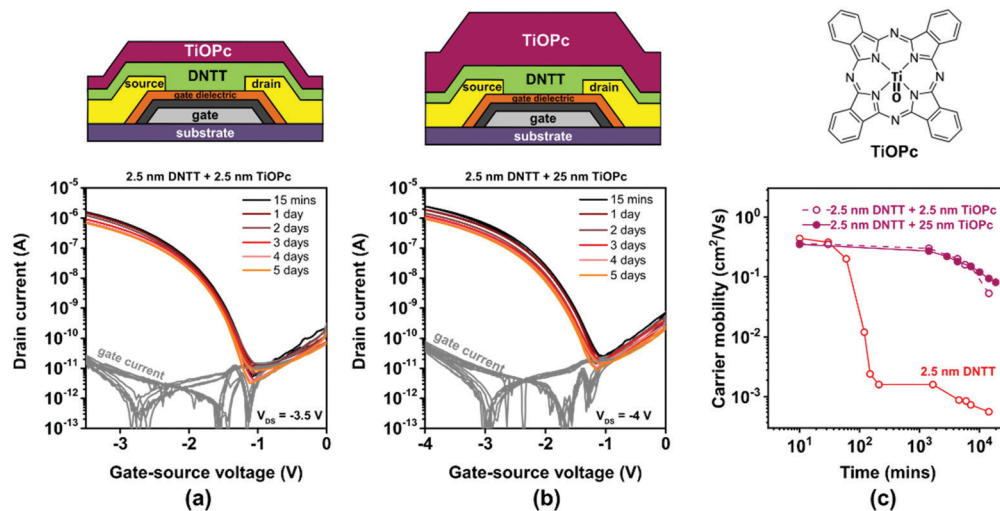


Fig. 7 Effect of *in situ* encapsulation of ultrathin, vacuum-deposited DNTT films (having a nominal thickness of 2.5 nm) with vacuum-deposited TiOPc on the evolution of the TFT characteristics. Transfer curves of TFTs encapsulated with a nominally (a) 2.5 nm-thick and (b) 25 nm-thick TiOPc film, measured repeatedly over a period of 5 days after TFT fabrication. (c) Evolution of the effective charge-carrier mobility extracted from the measured transfer curves.

This degradation in effective mobility is notably more rapid than the mobility degradation in top-contact DNTT TFTs,<sup>63</sup> but for the experiments described here, it was necessary to fabricate the TFTs in the bottom-contact architecture in order to minimize the delay between the completion of the DNTT deposition and the first measurement of the current-voltage characteristics. The long-term stability of bottom-contact DNTT TFTs has to our knowledge not been previously reported, and the reasons for the relatively rapid degradation of the effective mobility observed here is not known, but it may be related to a degradation of the interface between the source/drain contacts and the vacuum-deposited DNTT film.

The TFTs with the ultrathin DNTT films (nominal thickness 2.5 nm) initially have effective carrier mobilities between 0.3 and 0.4  $\text{cm}^2 \text{V}^{-1} \text{s}^{-1}$ , regardless of whether they were encapsulated or not and, if so, regardless of the encapsulation

material (PTFE,  $\text{C}_{60}$ , TiOPc) and the nominal thickness of the encapsulation film (Fig. 8a). The only exception is the TFT in which the ultrathin DNTT film was encapsulated with a nominally 10 nm-thick PTFE film, which has a somewhat larger initial mobility ( $0.7 \text{ cm}^2 \text{V}^{-1} \text{s}^{-1}$ ). The reason for this discrepancy is unknown, but it may be the result of unavoidable substrate-to-substrate variations in the fabrication process. The results confirm that it is indeed possible to fabricate functional TFTs based on vacuum-deposited DNTT films with a nominal thickness of less than two molecular monolayers.

Encapsulation of these films has a pronounced effect on their stability: while the effective mobility of unencapsulated ultrathin DNTT films decreases by two orders of magnitude within a day, the mobility of encapsulated films decreases by less than one order of magnitude within five days (Fig. 8b). The mechanism by which the encapsulation of the ultrathin DNTT

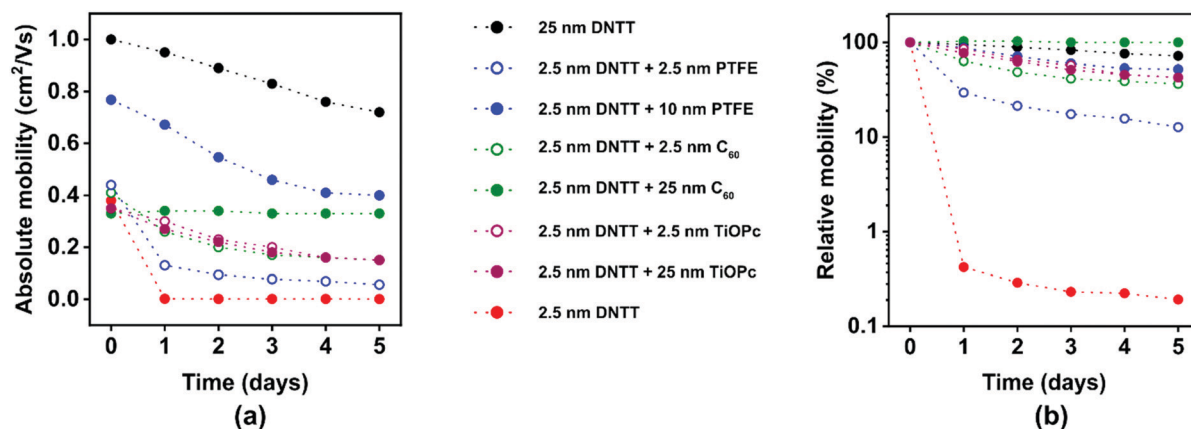


Fig. 8 Evolution of the effective charge-carrier mobilities of TFTs based on vacuum-deposited DNTT films with nominal thicknesses of 2.5 and 25 nm with and without encapsulation with PTFE,  $\text{C}_{60}$  or TiOPc with various nominal thicknesses. (a) Absolute value of the effective carrier mobility. (b) Effective carrier mobility normalized to the mobility measured immediately after TFT fabrication.





films decelerates the degradation of the carrier mobility in these films is not by protecting the DNTT molecules from ambient air (given the outstanding chemical stability of DNTT during long-term exposure to air and humidity<sup>24</sup>), but by stabilizing the fragile morphology of the ultrathin DNTT films by burying it under a thick film of additional mass. Although we are unable to inspect the morphology of the aged ultrathin DNTT films buried underneath the thick encapsulation films by AFM measurements, we hypothesize that the deceleration of the mobility degradation observed here is a direct result of preventing mass transport within the ultrathin DNTT films by the encapsulation film. This hypothesis is supported by the observation that for two of the three encapsulation materials, thicker encapsulation films provide better stability (Fig. 8b). The interesting question for how long the encapsulation can delay the break-up of the initially closed DNTT monolayer is beyond the scope of the present study.

The choice of the encapsulation material is dictated by a number of considerations. Since the spontaneous structural reconfiguration of the ultrathin DNTT films commences immediately after the DNTT deposition, it is necessary to start the deposition of the encapsulation film immediately after the completion of the DNTT deposition. This limits the choice of encapsulation materials to such materials that can be deposited in the same high-vacuum system in which the ultrathin DNTT films are deposited, without breaking vacuum and without time-consuming substrate transfers. If the purpose of the experiments is an investigation of the charge transport in ultrathin DNTT films, the choice of encapsulation materials should be limited to materials that do not participate in the charge transport. Of the three materials we have tested here, this obviously eliminates C<sub>60</sub> (and possibly also TiOPc, although Fig. 7 seems to indicate that the TiOPc films do not participate in charge transport). The observation that the stabilizing effect of the encapsulation increases with increasing thickness of the encapsulation film calls for materials that can be easily deposited as thick films. Unfortunately, this essentially eliminates PTFE, due to its unfavorably high vapor pressure.

Another aspect to be considered in the choice of the encapsulation material is the homogeneity and long-term morphological stability of the encapsulation film, which would likely benefit from choosing materials with a higher molecular weight and thus smaller diffusivity on the DNTT surface, and possibly also from performing the deposition of the encapsulation film at cryogenic temperatures to promote homogeneous coverage of the encapsulating material.<sup>59</sup>

## Conclusions

We have monitored the spontaneously occurring changes in the thin-film morphology of vacuum-deposited ultrathin DNTT films by AFM and analyzed how this spontaneous structural reconfiguration affects the electric characteristics of TFTs based on such films. Our results confirm that the observed transition from an initially closed (or at least connected) monolayer to tall,

isolated islands and the associated loss of percolation through the films cause a rapid and substantial degradation of the effective charge-carrier mobility of the TFTs. By showing that these morphological changes are prevented by keeping the substrates at cryogenic temperatures, we have confirmed that this structural reorganization is indeed thermally activated. We have also tested the extent to which encapsulation using three different materials, PTFE, C<sub>60</sub>, and TiOPc, is able to stabilize the morphology of the ultrathin DNTT films. Although we have not been able to stop the morphological changes entirely, we have found that encapsulation can slow the rate of degradation quite significantly. The ability to fabricate TFTs based on vacuum-deposited ultrathin organic–semiconductor films with sufficient long-term stability is potentially of interest for interface and charge-transport investigations.

## Experimental section

### Film characterization

Samples for film characterization were prepared on silicon substrates. Silicon was chosen as the substrate to provide the smallest possible surface roughness. Aluminum with a thickness of 30 nm and a root-mean-square surface roughness of less than 1 nm was deposited by thermal evaporation in vacuum at a rate of 2 nm s<sup>−1</sup>.<sup>52</sup> The aluminum surface was exposed to an RF oxygen plasma (oxygen flow rate: 30 sccm, oxygen partial pressure: 10 mTorr, RF power: 200 W, duration: 30 s) to increase the thickness of the native aluminum oxide (AlO<sub>x</sub>) to about 3.6 nm. On some substrates, ultrathin DNTT (Sigma-Aldrich) was deposited directly onto bare AlO<sub>x</sub>. On other substrates, a monolayer of *n*-tetradecylphosphonic acid, *n*-hexylphosphonic acid, or *n*-octadecylphosphonic acid (PCI Synthesis, Newburyport, MA, USA) was allowed to self-assemble on the AlO<sub>x</sub> surface by immersing the substrates into a 1 mM-solution of the phosphonic acid in 2-propanol for a duration of 12 hours, after which the substrates were rinsed with 2-propanol, dried with nitrogen and heated to a temperature of 80 °C for 10 minutes to stabilize the monolayer.<sup>64</sup> DNTT was deposited by thermal sublimation in vacuum at a rate of 2 nm min<sup>−1</sup> (without heating the substrate to avoid the unnecessary acceleration of the morphological changes). AFM images were acquired in tapping mode using a Bruker Nanoscope III MultiMode<sup>®</sup> or an Agilent SPM 5500 and cantilevers with resonance frequencies between 200 and 500 kHz and a typical tip radius of 8 nm. X-ray diffraction measurements were performed using a Bruker D8 Discovery diffractometer with monochromatized Cu K $\alpha$  radiation ( $\lambda$  = 1.542 Å) and a sensitive one-dimensional LynxEye silicon strip detector. All AFM and XRD measurements were performed in ambient air at room temperature.

### TFT fabrication and characterization

TFTs were fabricated in the inverted coplanar (bottom-gate, bottom-contact) architecture on silicon substrates coated with 100 nm-thick thermally grown silicon dioxide. This device architecture was chosen to minimize the delay between the



completion of the semiconductor deposition and the first measurement of the current–voltage characteristics. Aluminum gate electrodes with a thickness of 30 nm were deposited by thermal evaporation in vacuum at a rate of  $2 \text{ nm s}^{-1}$ . The gate dielectric consists of a thin, oxygen-plasma-grown film of  $\text{AlO}_x$  and a solution-processed *n*-tetradecylphosphonic acid SAM and has a capacitance of about  $0.7 \text{ } \mu\text{F cm}^{-2}$ .<sup>42</sup> For the source and drain contacts, titanium with a thickness of 0.2 nm (to improve adhesion) and gold with a thickness of 30 nm were sequentially deposited by thermal evaporation in vacuum. To minimize the contact resistance, the surface of the Au source/drain contacts was functionalized with a monolayer of pentafluorobenzethiol (PFBT).<sup>53</sup> DNTT was deposited by thermal sublimation in vacuum (base pressure  $10^{-6}$  mbar; substrate temperature  $20^\circ\text{C}$ ) at a rate of  $2 \text{ nm min}^{-1}$ . PTFE (Sigma-Aldrich, free-flowing powder, particle size  $<12 \text{ } \mu\text{m}$ ),  $\text{C}_{60}$  (Sigma-Aldrich) or TiOPc (Sigma-Aldrich) encapsulation films were deposited by thermal sublimation in vacuum within one to two minutes after completion of the DNTT deposition. Deposition rate and nominal thickness of all vacuum-deposited films were measured using a quartz-crystal microbalance positioned in close proximity to the substrate holder. Gate electrodes, source/drain contacts, organic semiconductor and encapsulation films were all patterned using polyimide shadow masks (CADiLAC Laser, Hilpoltstein, Germany). The TFTs have a channel length of  $20 \text{ } \mu\text{m}$  and a channel width of  $100 \text{ } \mu\text{m}$ . The electrical measurements were performed in ambient air at room temperature using a manual probe station and an Agilent 4156C Semiconductor Parameter Analyzer. The first measurement was performed within 15 minutes after the completion of the last vacuum deposition, and the substrates were stored in ambient air at room temperature between measurements.

## Conflicts of interest

There are no conflicts to declare.

## Acknowledgements

This work is partially funded by the German Research Foundation (DFG) under the Grants No. K. L. 2223/6-2 (SPP FFlexCom), K. L. 2223/7-1, INST 35/1429-1 (SFB 1249) and 223848855-SFB 1083 (TP A2). D. G. gratefully acknowledges financial support by the Cusanuswerk. Open Access funding provided by the Max Planck Society.

## Notes and references

- M. Noda, N. Kobayashi, M. Katsuhara, A. Yumoto, S. Ushikura, R. Yasuda, N. Hirai, G. Yukawa, I. Yagi, K. Nomoto and T. Urabe, *J. Soc. Inf. Disp.*, 2011, **19**, 316–322.
- B. Peng, X. Ren, Z. Wang, X. Wang, R. C. Roberts and P. K. L. Chan, *Sci. Rep.*, 2014, **4**, 6430.
- H. Fuketa, K. Yoshioka, Y. Shinozuka and K. Ishida, *IEEE Trans. Biomed. Eng.*, 2014, **8**, 824–833.
- Y. Chu, X. Wu, J. Lu, D. Liu, J. Du, G. Zhang and J. Huang, *Adv. Sci.*, 2016, **3**, 1500435.
- M. Kaltenbrunner, T. Sekitani, J. Reeder, T. Yokota, K. Kuribara, T. Tokuhara, M. Drack, R. Schwödiauer, I. Graz, S. Bauer-Gogonea, S. Bauer and T. Someya, *Nature*, 2013, **499**, 458–463.
- S. Lee, A. Reuveny, J. Reeder, S. Lee, H. Jin, Q. Liu, T. Yokota, T. Sekitani, T. Isoyama, Y. Abe, Z. Suo and T. Someya, *Nat. Nanotechnol.*, 2016, **11**, 472–478.
- W. Lee, D. Kim, J. Rivnay, N. Matsuhisa, T. Lonjaret, T. Yokota, H. Yawo, M. Sekino, G. G. Malliaras and T. Someya, *Adv. Mater.*, 2016, **28**, 9722–9728.
- X. Ren, K. Pei, B. Peng, Z. Zhang, Z. Wang, X. Wang and P. K. L. Chan, *Adv. Mater.*, 2016, **28**, 4832–4838.
- M. Kondo, M. Melzer, D. Karnaushenko, T. Uemura, S. Yoshimoto and M. Akiyama, *Sci. Adv.*, 2020, **6**, eaay6094.
- F. Dinelli, M. Murgia, P. Levy, M. Cavallini, F. Biscarini and D. M. De Leeuw, *Phys. Rev. Lett.*, 2004, **92**, 90–93.
- F. J. Zhang, C. A. Di, N. Berdunov, Y. Hu, Y. Hu, X. Gao, Q. Meng, H. Sirringhaus and D. Zhu, *Adv. Mater.*, 2013, **25**, 1401–1407.
- T. Someya, A. Dodabalapur, J. Huang, K. C. See and H. E. Katz, *Adv. Mater.*, 2010, **22**, 3799–3811.
- Y. Fujisaki, Y. Nakajima, T. Takei, H. Fukagawa, T. Yamamoto and H. Fujikake, *IEEE Trans. Electron Devices*, 2012, **59**, 3442–3449.
- K. Fukuda, T. Sekitani, U. Zschieschang, H. Klauk, K. Kuribara, T. Yokota, T. Sugino, K. Asaka, M. Ikeda, H. Kuwabara, T. Yamamoto, K. Takimiya, T. Fukushima, T. Aida, M. Takamiya, T. Sakurai and T. Someya, *Adv. Funct. Mater.*, 2011, **21**, 4019–4027.
- J. Milvich, T. Zaki, M. Aghamohammadi, R. Rödel, U. Kraft, H. Klauk and J. N. Burghartz, *Org. Electron.*, 2015, **20**, 63–68.
- S. Kim, H. Moon, H. Kwon, G. Lee, D. Yu, J. Choi, J. Park, S. J. Kim and S. Yoo, *Adv. Mater. Technol.*, 2019, **4**, 1800332.
- K. Ishida, T. C. Huang, K. Honda, Y. Shinozuka, H. Fuketa, T. Yokota, U. Zschieschang, H. Klauk, G. Tortissier, T. Sekitani, H. Toshiyoshi, M. Takamiya, T. Someya and T. Sakurai, *IEEE J. Solid-State Circuits*, 2013, **48**, 255–264.
- T. Sekitani, T. Yokota, K. Kuribara, M. Kaltenbrunner, T. Fukushima, Y. Inoue, M. Sekino, T. Isoyama, Y. Abe, H. Onodera and T. Someya, *Nat. Commun.*, 2016, **7**, 11425.
- A. K. M. Mahfuzul Islam, M. Hamamatsu, T. Yokota, S. Lee, W. Yukita, M. Takamiya, T. Someya and T. Sakurai, *IEEE J. Emerg. Sel. Top. Circuits Syst.*, 2017, **7**, 81–91.
- M. Elsobky, Y. Mahsereci, Z. Yu, H. Richter, J. N. Burghartz, J. Keck, H. Klauk and U. Zschieschang, *Electron. Lett.*, 2018, **54**, 338–340.
- X. Zhao, H. Li, Q. Tang, Y. Tong and Y. Liu, *J. Mater. Chem. C*, 2019, **7**, 3199–3205.
- U. Kraft, K. Takimiya, M. J. Kang, R. Rödel, F. Letzkus, J. N. Burghartz, E. Weber and H. Klauk, *Org. Electronics*, 2016, **35**, 33–40.
- E. D. Glowacki, S. Bauer, M. S. White, U. Monkowius, N. S. Sariciftci, M. Kaltenbrunner, L. Torsi, P. Mastorilli, J. Gsiorowski, G. P. Suranna, T. Sekitani, G. Romanazzi,



- T. Someya and M. Irimia-Vladu, *Adv. Mater.*, 2013, **25**, 1563–1569.
- 24 U. Zschieschang, F. Ante, D. Kälblein, T. Yamamoto, K. Takimiya, H. Kuwabara, M. Ikeda, T. Sekitani, T. Someya, J. B. Nimoth and H. Klauk, *Org. Electronics*, 2011, **12**, 1370–1375.
- 25 Z. Ding, G. Abbas, H. E. Assender, J. J. Morrison, S. G. Yeates, E. R. Patchett and D. M. Taylor, *ACS Appl. Mater. Interfaces*, 2014, **6**, 15224–15231.
- 26 S. Yogeve, R. Matsubara, M. Nakamura, U. Zschieschang, H. Klauk and Y. Rosenwaks, *Phys. Rev. Lett.*, 2013, **110**, 036803.
- 27 H. S. Lee, D. H. Kim, J. H. Cho, M. Hwang, Y. Jang, K. Cho, H. S. Lee, D. H. Kim, J. H. Cho, M. Hwang and Y. Jang, *J. Am. Chem. Soc.*, 2008, **130**, 10556–10564.
- 28 J. A. Venable, G. D. T. Spiller and M. Hanbucken, *Rep. Prog. Phys.*, 1984, **47**, 399–459.
- 29 M. C. Jung, M. R. Leyden, G. O. Nikiforov, M. V. Lee, H. K. Lee, T. J. Shin, K. Takimiya and Y. Qi, *ACS Appl. Mater. Interfaces*, 2015, **7**, 1833–1840.
- 30 R. Ruiz, D. Choudhary, B. Nickel, T. Toccoli, K. C. Chang, A. C. Mayer, P. Clancy, J. M. Blakely, R. L. Headrick, S. Iannotta and G. G. Malliaras, *Chem. Mater.*, 2004, **16**, 4497–4508.
- 31 T. Yamamoto and K. Takimiya, *J. Am. Chem. Soc.*, 2007, **129**, 2224–2225.
- 32 W. Xie, K. Willa, Y. Wu, R. Häusermann, K. Takimiya, B. Batlogg and C. D. Frisbie, *Adv. Mater.*, 2013, **25**, 3478–3484.
- 33 T. Breuer, A. Karthäuser, H. Klemm, F. Genuzio, G. Peschel, A. Fuhrich, T. Schmidt and G. Witte, *ACS Appl. Mater. Interfaces*, 2017, **9**, 8384–8392.
- 34 S. Herminghaus, K. Jacobs, K. Mecke, J. Bischof, A. Fery, M. Ibn-Elhaj and S. Schlagowski, *Science*, 1998, **282**, 916–919.
- 35 S. Rath and H. Port, *Chem. Phys. Lett.*, 2006, **421**, 152–156.
- 36 A. C. Dürr, F. Schreiber, K. A. Ritley, V. Kruppa, J. Krug, H. Dosch and B. Struth, *Phys. Rev. Lett.*, 2003, **90**, 016104.
- 37 D. Käfer, C. Wöll and G. Witte, *Appl. Phys. A: Mater. Sci. Process.*, 2009, **95**, 273–284.
- 38 A. Amassian, V. A. Pozdin, T. V. Desai, S. Hong, A. R. Woll, J. D. Ferguson, J. D. Brock, G. G. Malliaras and J. R. Engstrom, *J. Mater. Chem.*, 2009, **19**, 5580–5592.
- 39 A. A. Virkar, S. Mannsfeld, Z. Bao and N. Stingelin, *Adv. Mater.*, 2010, **22**, 3857–3875.
- 40 H. Chen, M. Li, Z. Lu, X. Wang, J. Yang, Z. Wang, F. Zhang, C. Gu, W. Zhang, Y. Sun, J. Sun, W. Zhu and X. Guo, *Nat. Commun.*, 2019, **10**, 3872.
- 41 S. Lee, Y. Inoue, D. Kim, A. Reuveny, K. Kuribara, T. Yokota, J. Reeder, M. Sekino, T. Sekitani, Y. Abe and T. Someya, *Nat. Commun.*, 2014, **5**, 5898.
- 42 U. Zschieschang and H. Klauk, *Org. Electronics*, 2015, **25**, 340–344.
- 43 U. Kraft, T. Zaki, F. Letzkus, J. N. Burghartz, E. Weber, B. Murmann and H. Klauk, *Adv. Electron. Mater.*, 2019, **5**, 1800453.
- 44 R. Acharya, B. Peng, P. K. L. Chan, G. Schmitz and H. Klauk, *ACS Appl. Mater. Interfaces*, 2019, **11**, 27104–27111.
- 45 D. Yu, Y. Q. Yang, Z. Chen, Y. Tao and Y. F. Liu, *Opt. Commun.*, 2016, **362**, 43–49.
- 46 J. M. Kim, J. Oh, K. M. Jung, K. C. Park, J. H. Jeon and Y. S. Kim, *Semicond. Sci. Technol.*, 2019, **34**, 075015.
- 47 D. Feili, M. Schuettler, T. Doerge, S. Kammer and T. Stieglitz, *Sens. Actuators, A*, 2005, **120**, 101–109.
- 48 O. Acton, G. G. Ting, P. J. Shamberger, F. S. Ohuchi, H. Ma and A. K. Y. Jen, *ACS Appl. Mater. Interfaces*, 2010, **2**, 511–520.
- 49 K. Fukuda, T. Hamamoto, T. Yokota, T. Sekitani, U. Zschieschang, H. Klauk and T. Someya, *Appl. Phys. Lett.*, 2009, **95**, 203301.
- 50 A. Jedaa, M. Burkhardt, U. Zschieschang, H. Klauk, D. Habich, G. Schmid and M. Halik, *Org. Electronics*, 2009, **10**, 1442–1447.
- 51 D. M. Spori, N. V. Venkataraman, S. G. P. Tosatti, F. Durmaz, N. D. Spencer and S. Zürcher, *Langmuir*, 2007, **23**, 8053–8060.
- 52 M. Geiger, R. Acharya, E. Reutter, T. Ferschke, U. Zschieschang, J. Weis, J. Pflaum, H. Klauk and R. T. Weitz, *Adv. Mater. Interfaces*, 2020, **7**, 1902145.
- 53 J. W. Borchert, B. Peng, F. Letzkus, J. N. Burghartz, P. K. L. Chan, K. Zojer, S. Ludwigs and H. Klauk, *Nat. Commun.*, 2019, **10**, 1119.
- 54 C. Pannemann, T. Diekmann, U. Hilleringmann, U. Schürmann, M. Scharnberg, V. Zaporozhchenko, R. Adelung and F. Faupel, *Mater. Sci. Pol.*, 2007, **25**, 95–101.
- 55 J. Fresnais, L. Benyahia and F. Poncin-Epaillard, *Surf. Interface Anal.*, 2006, **38**, 144–149.
- 56 N. Selway, V. Chan and J. R. Stokes, *Soft Matter*, 2017, **13**, 1702–1715.
- 57 R. C. Haddon, A. S. Perel, R. C. Morris, T. T. M. Palstra, A. F. Hebard and R. M. Fleming, *Appl. Phys. Lett.*, 1995, **67**, 121.
- 58 T. Breuer, T. Geiger, H. F. Bettinger and G. Witte, *J. Phys.: Condens. Matter*, 2019, **31**, 034003.
- 59 T. Breuer and G. Witte, *ACS Appl. Mater. Interfaces*, 2013, **5**, 9740–9745.
- 60 O. A. Melville, B. H. Lessard and T. P. Bender, *ACS Appl. Mater. Interfaces*, 2015, **7**, 13105–13118.
- 61 M. Kothe, F. Widdascheck and G. Witte, *J. Phys. Chem. C*, 2019, **123**, 6097–6106.
- 62 Y. Fujinaka, R. Ye, K. Ohta, K. Nishidate and M. Baba, *Mol. Cryst. Liq. Cryst.*, 2013, **580**, 110–116.
- 63 U. Zschieschang, J. W. Borchert, M. Geiger, F. Letzkus, J. N. Burghartz and H. Klauk, *Org. Electronics*, 2018, **61**, 65–69.
- 64 L. B. Goetting, T. Deng and G. M. Whitesides, *Langmuir*, 1999, **15**, 1182–1191.

