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Introduction

Semiconductor quantum dots are excellent building blocks for designing light-harvesting assemblies.^{1,2} The ability to chemically modify the surface with a functionalized ligand or to couple with another semiconductor particle offers a variety of ways to harvest visible photons.^{3–5} Since the 1990s, metal chalcogenide quantum dots (QDs), CdSe in particular, have served as the prototypical compound to elucidate excited state and charge transfer properties.^{6–8} In recent years, another quantum dot system, *viz.*, perovskite nanocrystals (CsPbX_3 , X = Cl, Br, I), has emerged as a model semiconductor QD system to probe light induced optoelectronic and photocatalytic properties.^{9–13} We have recently elucidated the photocatalytic aspects of CsPbBr_3 QDs by probing the interfacial electron transfer to methyl viologen^{14,15} and ferrocenium cation.¹⁶ To date, the use of these perovskite quantum dots in photocatalysis has been limited only to a few nonpolar solvents.^{17–21} The weakly-binding organic ligand shell around perovskite nanocrystals does not provide sufficient stability in polar solvents. In order to expand the scope of the perovskite nanocrystals to a wide range of

photocatalytic applications (*e.g.*, solar hydrogen production or CO_2 reduction), it is important to provide protection against chemical transformation in the presence of a redox couple or in a polar medium.

One simple approach to achieve stability in polar solvents is to cap the semiconductor nanocrystals with a thin inorganic shell. Design of such heterostructures has been successfully employed for binary and ternary semiconductors like CdSe/ZnS,^{22,23} InP/ZnS,²⁴ and AgInS₂/ZnS.²⁵ The heterostructure with type I or type II band alignment offers strategies to enhance emission of the core QD or improve charge separation within the heterostructure. Although a few reports exist to-date of capping CsPbBr_3 QDs with SiO_2 ,²⁶ CdS,²⁷ or ZnS^{28,29} shells, none of these heterostructures have shown a major leap in achieving improved performance with a long-term stability in polar solvents. Ambiguity still exists whether the added material forms a continuous shell around the perovskite core or forms smaller discontinuous islands on the surface.²⁶ Given the difficulty in imaging the thin inorganic shell around perovskite nanocrystals, because of the image contrast, one employs stability tests in a polar medium or its resistance to halide exchange to confirm surface modification.^{30–33}

Designing perovskite heterostructures with metal chalcogenide shells can have several distinct advantages: (i) providing stability towards increased polarity of the solvent, (ii) remediating surface defects by directly interacting with the vacancies, and (iii) allowing for type I or quasi-type II band alignment to promote increased charge recombination in the core (increased emission yield) or improved charge separation.^{34–36}

^aRadiation Laboratory, University of Notre Dame, Notre Dame, Indiana 46556, USA

^bDepartment of Chemistry and Biochemistry, University of Notre Dame, Notre Dame, Indiana 46556, USA

^cDepartment of Chemical and Biomolecular Engineering, University of Notre Dame, Notre Dame, Indiana 46556, USA. E-mail: pkamat@nd.edu

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In this context, a CsPbBr_3 – CdS heterostructure offers an attractive means to tune the band energies, as their conduction bands are nearly isoenergetic (E_{CB} of CsPbBr_3 and quantized $\text{CdS} \approx -0.8$ V *versus* NHE)^{37–39} and facilitate charge separation. In addition, cubic CsPbBr_3 nanocrystals have a lattice constant (a) of 5.85 Å (ref. 40) while that of CdS (zinc blende structure) is 5.83 Å.⁴¹ The similarity of the two values signifies the possibility of having a less strained interface with reduced defect states.⁴² We have now successfully prepared CsPbBr_3 – CdS heterostructures in a two-step method, and the optical properties of these structures are discussed.

Results and discussion

CsPbBr_3 – CdS heterostructure

CsPbBr_3 quantum dots (QDs) dispersed in octadecene (ODE) were prepared using a previously reported procedure.⁴³ These QDs were then treated with cadmium diethyldithiocarbamate ($\text{Cd}(\text{DDTC})_2$) at 110 °C to obtain CdS capped CsPbBr_3 QDs. Experimental details on the synthesis of CsPbBr_3 and CsPbBr_3 – CdS QDs are presented in the ESI (Fig. S1†). The transmission electron microscopy (TEM) images of the two nanocrystals are shown in Fig. 1A and B. These cubic particles are similar in size showing particles of lengths 8–9 nm (see Fig. S2† for size distribution analysis). This shows that CdS capping in the heterostructure is relatively thin compared to the CsPbBr_3 core and it is difficult to identify with the image contrast in TEM images. This observation is consistent with earlier work which reports difficulty in characterizing the shell in a CsPbBr_3 heterostructure using TEM analysis.^{44,45} These studies have attributed the imaging difficulty to low electron density contrast of the shell. However, other techniques such as elemental analysis can be useful to overcome these limitations.

We succeeded in establishing the presence of CdS in the CsPbBr_3 – CdS heterostructures through elemental analysis with TEM energy dispersive X-ray (EDX) spectroscopy. Fig. 1C–G which present elemental mapping, confirm the presence of Cd and S along with Cs, Pb, and Br. The elemental ratio (Fig. 1H) also suggests a relatively low concentration of CdS as compared to Cs, Pb and Br in the heterostructure. We can conclude that any CdS in the heterostructure is of the order of a monolayer. However, there is also the possibility of forming small clusters in and around the CsPbBr_3 QDs.

Evidence of surface modification with CdS was also seen through the change in the surface charge. Zeta potential measurements indicated that CsPbBr_3 – CdS QDs suspended in toluene carry more negative surface charge (−37.4 mV) than pristine CsPbBr_3 nanocrystals in toluene (−15.8 mV). This increased surface charge of CsPbBr_3 – CdS QDs enabled us to carry out electrophoretic deposition of a film under the influence of a DC field (see ESI† for details; Fig. S3†). The increased surface negativity and ability to be deposited as a film under applied bias indicates a modified surface around CsPbBr_3 . Similar electrophoretic deposition was also possible when CsPbBr_3 nanocrystals were coated with a PbSO_4 –oleate shell.⁴⁶ It should be noted that pristine CsPbBr_3 QDs suspended in

toluene cannot be deposited as film using electrophoresis as it does not carry sufficient surface charge.

In addition to CsPbBr_3 and CsPbBr_3 – CdS heterostructures, we also synthesized CdS QDs with the same ligands following a similar experimental procedure (*i.e.* without CsPbBr_3). Fig. 2A shows the absorption spectra of CsPbBr_3 , CdS and CsPbBr_3 – CdS QDs in toluene. The CdS and CsPbBr_3 QDs exhibit characteristic excitonic peaks at 434 and 518 nm, respectively. Although the absorption spectrum of the CsPbBr_3 – CdS heterostructure shows two peaks that overlap with the absorption of individual QDs, the peak around 434 nm may also arise from the deposition of small size CdS particles on the surface of CsPbBr_3 nanocrystals. Such a decoration of CdS particles, if any, would give rise to a CdS excitonic peak in the heterostructures. The excitonic peak at 518 nm (CsPbBr_3) remains unaffected after heterostructure formation, thus ruling out any interference of exchange of metal ions. Similarly, the tail absorption at longer wavelengths arises from the scattering effects, similar to what is seen in PbSO_4 –oleate capped CsPbBr_3 .⁴⁶

The emission spectra of these three nanostructures are shown in Fig. 2B. Whereas CdS QDs remain the least emissive, both CsPbBr_3 and CsPbBr_3 – CdS QDs are highly emissive. Additionally, the emission features of CsPbBr_3 QDs remain unchanged following the deposition of CdS : the emission maximum (521 nm) and the full width at half maximum (18 nm) of CsPbBr_3 are unaffected after CdS modification. These results confirm that there is no substitution of cations during the heterostructures synthesis, and thus the emission characteristics of the parent CsPbBr_3 QDs are retained in the heterostructure. If there was any substitution of Pb^{2+} with Cd^{2+} we would expect a blue shift in the absorption and emission maxima of the perovskite QD.⁴⁷ The excitation spectra recorded at different emission wavelengths confirm the origin of the emission to arise from the CsPbBr_3 (Fig. S4 in the ESI†).

Another interesting aspect of the capping with CdS is the enhancement in emission yield. The emission quantum yields as determined using an integrating sphere were 39% and 60% for CsPbBr_3 and CsPbBr_3 – CdS QDs, respectively (see ESI† for details). Passivation of surface defects by CdS is expected to suppress nonradiative processes and thus lead to increased emission yield. For example, capping of CdSe with CdS has resulted in the significant enhancement of emission yield.^{48–51} The flow of charge carriers from the CdS shell to CdSe core in these studies was established through emission and excitation spectral measurements. Control experiments were carried out to check whether introduction of Cd^{2+} ions alone can induce similar changes in the emission properties. Fig. S5† shows a decrease in emission yield and lifetime when CsPbBr_3 was treated with cadmium acetylacetone ($\text{Cd}(\text{acac})_2$) instead of $\text{Cd}(\text{DDTC})_2$. Additionally, treatment with Cd^{2+} did not change the absorption of the QDs. This further confirms that the observed optical properties are due to the presence of CdS in the heterostructure.

We employed time-resolved emission measurements to monitor the excited state behavior of CsPbBr_3 before and after CdS deposition. The emission decay at 520 nm was monitored for CsPbBr_3 and CsPbBr_3 – CdS samples (Fig. 2C). Each trace was



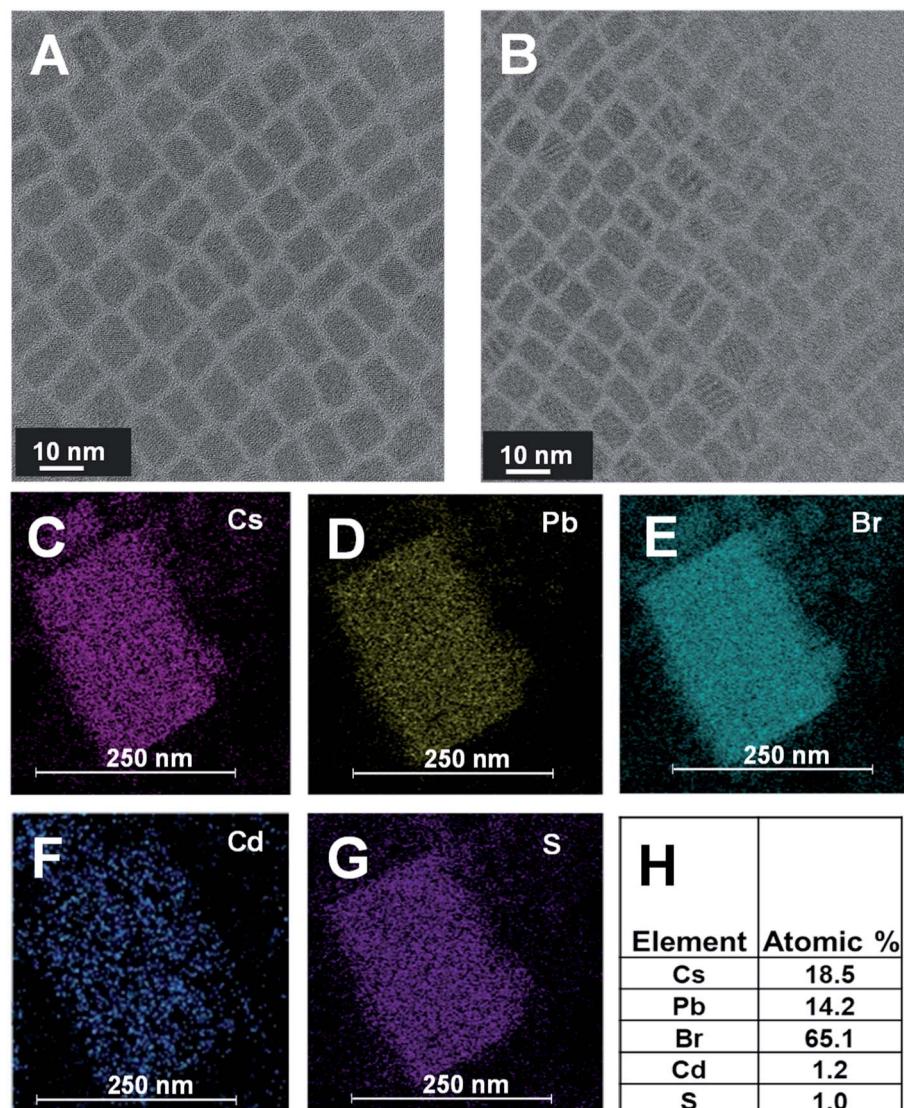


Fig. 1 TEM image for (A) pristine CsPbBr_3 and (B) CsPbBr_3 –CdS. (C–G) Elemental maps showing presence of Cs, Pb, Br, Cd, and S, respectively. (H) Atomic percentage of different elements as measured by EDAX.

analyzed using a biexponential decay fit and the fitting parameters are presented in Table S1.† Of interest is the increase in emission lifetime of CsPbBr_3 upon capping with CdS. A nearly seven-fold increase in average lifetime of CsPbBr_3 –CdS QDs ($\tau_{\text{ave}} = 46.9$ ns) was observed over that of pristine CsPbBr_3 QDs ($\tau_{\text{ave}} = 7.0$ ns). This shows that CdS deposition facilitates long-lived charge separation in CsPbBr_3 –CdS. In addition to surface passivation, we can also expect the formation of a quasi-type II heterojunction as shown in the scheme (Fig. 2D). Whereas direct charge carrier recombination is dominant in pristine CsPbBr_3 , the nearly isoenergetic conduction bands of CsPbBr_3 and CdS can facilitate delocalization of electrons across the two semiconductors, thus improving charge separation. Since the CdS layer is relatively thin, its contribution to the emission is expected to be small. The excitation spectra (Fig. S4†) rule out the contribution from CdS to overall emission. The observed increase in lifetime

parallels the emission yield enhancement seen in the CsPbBr_3 –CdS heterostructure.

Stability in polar environment

Attaining long term stability of CsPbBr_3 nanocrystals in polar solvents remains a challenge. CsPbBr_3 nanocrystals undergo rapid degradation in polar solvents, which has hampered their applications in photocatalysis. Even the addition of a small amount of polar solvents such as ethanol or water can induce chemical transformation/precipitation of CsPbBr_3 QDs and thus a loss of photoactivity.^{9,52,53} Recently, it was reported that ZnS-capped CsPbBr_3 QDs were stable in a toluene : water biphasic mixture. The contact with water was made by periodic shaking since the two solvents are immiscible.²⁸ We also conducted a similar stability test of CsPbBr_3 and CsPbBr_3 –CdS nanocrystals using a biphasic mixture of toluene and water with periodic shaking. The emission spectra of CsPbBr_3 and



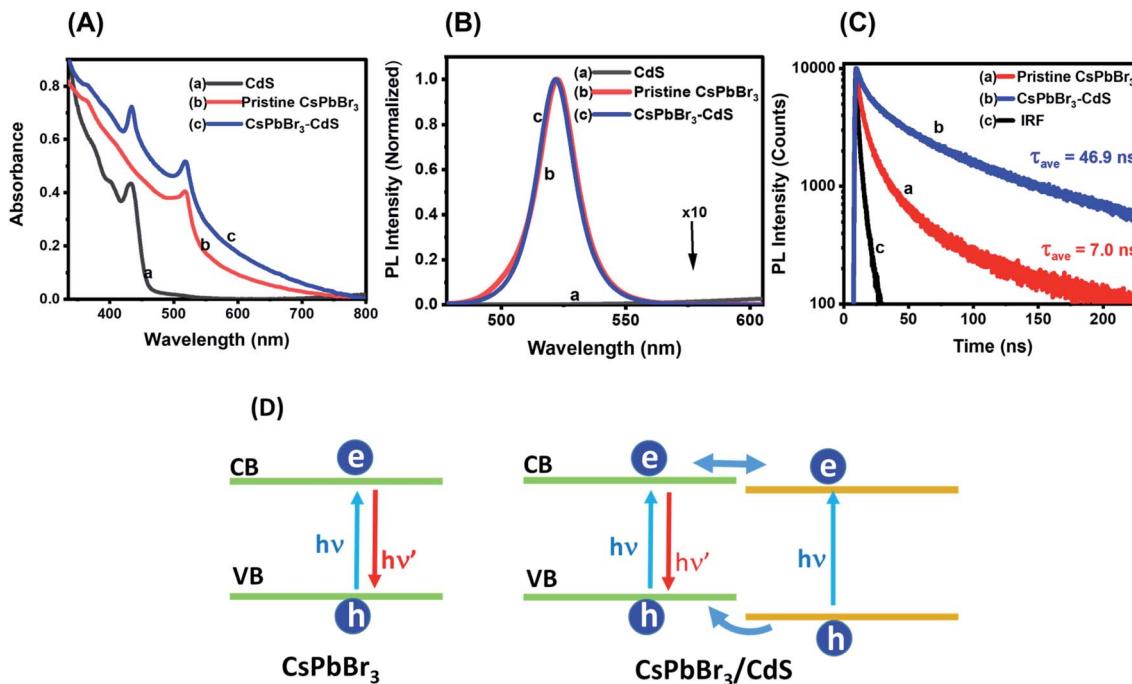


Fig. 2 (A) Absorption spectra and (B) corresponding photoluminescence (PL) spectra of (a) CdS, (b) pristine CsPbBr₃ and (c) CsPbBr₃-CdS QD suspensions in toluene. (C) PL decay traces monitored at 520 nm using 370 nm excitation: (a) pristine CsPbBr₃ and (b) CsPbBr₃-CdS nanocrystal suspensions in toluene, and (c) instrument response (IRF). The kinetic analysis is presented in the ESI.† (D) Schematic diagram illustrating charge separation and charge recombination in CsPbBr₃ and CsPbBr₃-CdS nanocrystals.

CsPbBr₃-CdS nanocrystals and the PL intensity variation recorded during 30 hour period are shown in Fig. S6.† CsPbBr₃ nanocrystals became non emissive after 20 hours of exposure in biphasic mixture. On the other hand, CsPbBr₃-CdS nanocrystals, after an initial drop in photoluminescence, maintained more than 40% emission even after 30 hours.

The biphasic solvent mixture approach does not represent an increase in the overall polarity of the medium. Ideally, an inorganic shell should prevent direct contact of CsPbBr₃ with a polar environment and maintain its photostability. We checked the stability of CsPbBr₃ and CsPbBr₃-CdS by introducing a miscible polar solvent (ethanol) to a toluene suspension of the QDs and monitoring the absorption and emission spectra over time. Fig. 3A and B show the absorption and emission spectra recorded following addition of ethanol (15% v/v) to toluene solution over a period of 60 minutes. The absorption of pristine CsPbBr₃ QDs shows enhanced absorbance with time due to scattering effects caused by turbidity as the ligands from QD surface become detached in the polar medium.⁵⁴ A ~85% decrease in the CsPbBr₃ emission yield is seen immediately after the addition of ethanol to the toluene solution. In addition, upon exposure to ethanol we also see a change in the absorption of CsPbBr₃ due to particle aggregation. These results confirm the susceptibility of CsPbBr₃ QDs to polar environment (Fig. 3C). On the other hand, CsPbBr₃-CdS QDs exhibit only a small decrease (~15%) in emission with a relatively small change in the absorption during 60 min of exposure in toluene/ethanol mixed solvent. The CdS deposition

provides the necessary protection for CsPbBr₃, and thus decreases its susceptibility to ethanol-induced degradation.

Excited state interactions with an electron acceptor

Since CsPbBr₃-CdS QDs were stable in toluene/ethanol mixed solvent, we were able to probe the excited state interactions and interfacial electron transfer with a cationic electron acceptor, ethyl viologen, EV²⁺. Emission spectra recorded at different concentrations of EV²⁺ are shown in Fig. 4A. The quenching of photoluminescence confirmed the excited state interaction between CsPbBr₃-CdS QDs and EV²⁺. As a control, we tested the solubility of EV²⁺ in the toluene/ethanol mixed solvent separately by recording absorption spectra and confirming the probe molecules are fully soluble at the concentrations employed in this study (Fig. S7†). Earlier studies have shown direct complexation between CsPbBr₃ and methyl viologen and elucidated the role of surface bound ligands in dictating the complexation constant.^{14,15} Here, we were able to quench the emission at micromolar concentrations of EV²⁺, thus indicating a complex formation in the ground state between CsPbBr₃-CdS and EV²⁺.⁵¹ The equilibrium of the bound and unbound EV²⁺ molecules (reaction (1)) can be expressed in terms of the apparent association constant, K_{app} and emission yields (expression (2)).⁵⁵ The observed quantum yield ($\phi_{(obs)}$) takes into account the emission arising from EV²⁺-bound (ϕ_f) and pristine (ϕ_f^0) CsPbBr₃ QDs. With increasing concentration of EV²⁺, more CsPbBr₃-CdS QDs bind to viologen and thus exhibit a decrease in the emission yield.



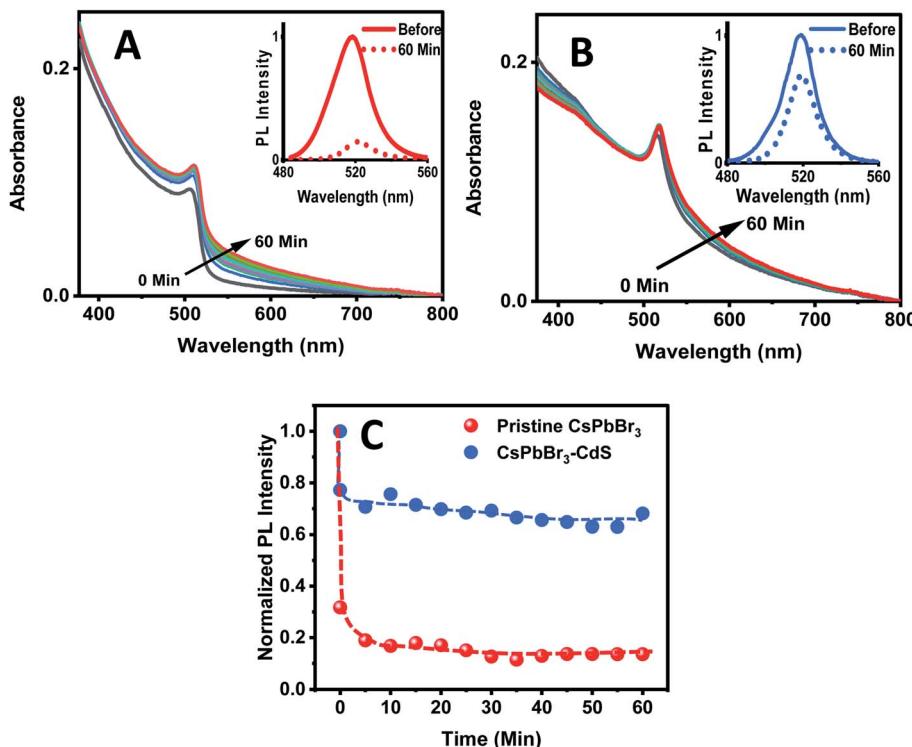


Fig. 3 Absorption spectra of (A) CsPbBr₃ and (B) CsPbBr₃-CdS QDs in toluene recorded at different time intervals after addition of ethanol (15% v/v). Spectra were recorded at a time interval of 5 min. Insets: PL spectra of QDs before and 60 min after ethanol addition (excitation: 370 nm). (C) Changes in the normalized PL intensity of the CsPbBr₃ and CsPbBr₃-CdS QDs monitored at 520 nm as a function of time after addition of ethanol (15% v/v) to the toluene suspension.



$$\frac{1}{(\phi_f^0 - \phi_f(\text{obs}))} = \frac{1}{(\phi_f^0 - \phi_f')} + \frac{1}{K_{\text{app}}(\phi_f^0 - \phi_f')[\text{EV}^{2+}]} \quad (2)$$

The photoluminescence quenching data was analyzed using expression (2). The emission intensity of the QDs at the emission maximum (which is proportional to quantum yield, $(\phi_f(\text{obs}))$) was monitored at different concentration of EV^{2+} .^{55,56} The linear dependence of the double reciprocal plot $(1/(\phi_f^0 - \phi_f(\text{obs})))$ versus $1/[\text{EV}^{2+}]$ in Fig. 4B confirms the validity of the association between CsPbBr₃-CdS and EV^{2+} . The apparent association constant K_{app} determined from the slope and intercept of the plot in Fig. 4B was $7.0 \times 10^4 \text{ M}^{-1}$. This complexation constant is 1–2 orders of magnitude smaller than the one observed for uncapped CsPbBr₃ and viologen ($0.8\text{--}7.0 \times 10^6 \text{ M}^{-1}$).¹⁵ The decrease in K_{app} further indicates that the presence of CdS reduces the surface interactions with the viologen. The K_{app} value we obtain in this study is in line with literature values of CdS interacting with viologens.⁵⁷

To further establish the excited state interactions, we monitored the photoluminescence lifetime of the CsPbBr₃-CdS QDs at different EV^{2+} concentrations. Time-resolved luminescence decay traces were recorded in toluene : ethanol (85 : 15% v/v) using an excitation source at 370 nm. The lifetimes were

fitted to a biexponential kinetic expression (expression 3), and the fitting parameters are given in Table S2.†

$$y = A_1 e^{-t/\tau_1} + A_2 e^{-t/\tau_2} \quad (3)$$

The average lifetime (τ_{ave}) decreased with increasing concentration of EV^{2+} in accordance with the photoluminescence quenching seen in Fig. 4A. The decrease in average lifetime from 42.4 ns to 18.1 ns upon addition of 12 μM EV^{2+} is indicative of a competing excited state deactivation pathway involving electron transfer from the CsPbBr₃-CdS QDs to EV^{2+} . Since this electron transfer is likely to occur within the time resolution ($\sim 1 \text{ ns}$) of our photoluminescence lifetime set up, we employed femtosecond transient absorption spectroscopy to resolve the electron transfer process.

Transient absorption spectra were recorded following laser pulse excitation at 400 nm (16 $\mu\text{J cm}^{-2}$). The transient absorption spectra of a representative CsPbBr₃-CdS QD sample containing 0 and 4 μM EV^{2+} are shown in Fig. 5A and B respectively. The negative absorption (exciton bleach) feature centered at $\sim 527 \text{ nm}$ corresponds to the charge separated state within the NCs.^{58–60} The charge separation which occurs within the laser pulse is seen in spectrum 'a' recorded with a probe delay of 1 ps. As electrons and holes recombine, a recovery in the bleached absorption is seen. The bleach recovery at 527 nm for the samples containing different amounts of EV^{2+} are presented in Fig. 5C (see Fig. S8† for longer time scale kinetics). With

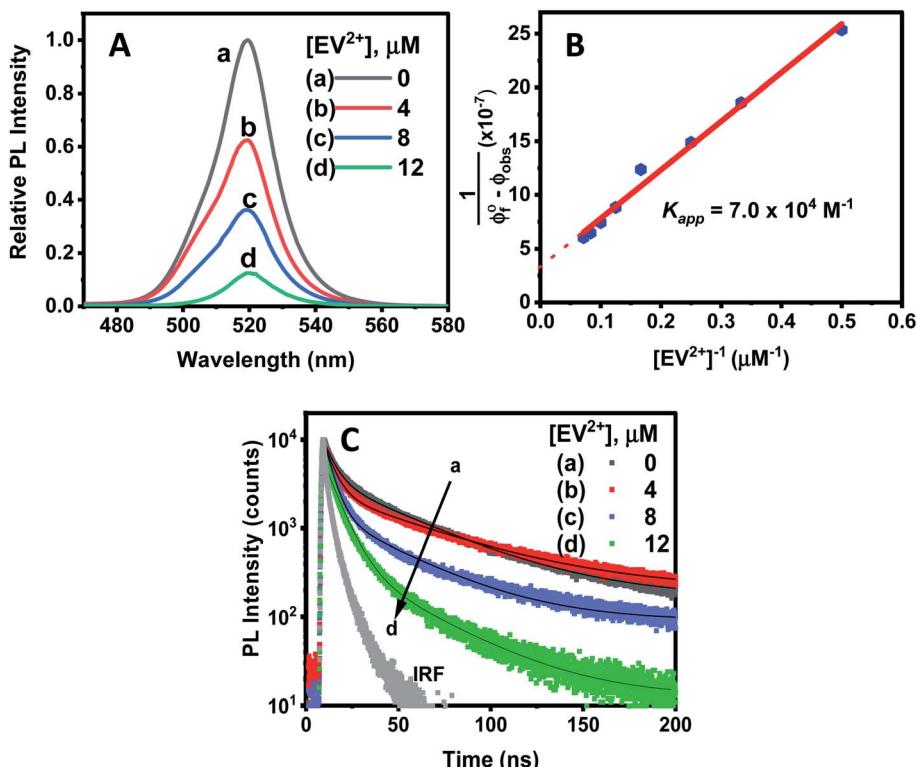
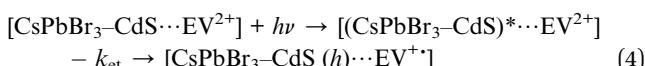


Fig. 4 (A) PL quenching of $\text{CsPbBr}_3\text{-CdS}$ QDs ($\approx 16 \text{ nM}$) in toluene : ethanol (90 : 10% v/v) upon additions of increasing concentration of ethyl viologen (EV^{2+}) dissolved in ethanol. (B) Double reciprocal plot analysis of emission quenching of $\text{CsPbBr}_3\text{-CdS}$ in toluene : ethanol. The slope (3.3×10^{-7}) and intercept ($4.3 \times 10^{-12} \text{ M}^{-1}$) was used to determine the apparent association constant, K_{app} , $7.0 (\pm 0.8) \times 10^4 \text{ M}^{-1}$. (C) Photoluminescence decay at different EV^{2+} concentrations, recorded with 370 nm excitation and monitored at $\lambda_{\text{max}} = 520 \text{ nm}$.

increasing EV^{2+} concentration we observe a quick recovery of the bleach feature, thus confirming the presence of an additional deactivation pathway for the photogenerated electrons, *viz.*, electron transfer to EV^{2+} (reaction (4)).



The bleach recovery was analyzed using a biexponential kinetic fit³⁵ and the fitting parameters (a_1, τ_1) and (a_2, τ_2) corresponding to fast and slow components are presented in Table S3.[†] While the fast component varied in the range 18.6–6.6 ps the long component varied in the range of 129.6–71.7 ps.

If we assign the decrease in the fast component to the electron transfer pathway from $\text{CsPbBr}_3\text{-CdS}$ to EV^{2+} , we can obtain the rate constant for electron transfer through the expression (5).

$$k_{\text{et}} = \frac{1}{\tau_1(\text{CsPbBr}_3\text{-CdS}-\text{EV}^{2+})} - \frac{1}{\tau_1(\text{CsPbBr}_3\text{-CdS})} \quad (5)$$

If we substitute the fast time components (τ_1) of $\text{CsPbBr}_3\text{-CdS}$ and that of $\text{CsPbBr}_3\text{-CdS}$ with 32 μM EV^{2+} in expression 5 (Table S3[†]), we obtain rate constant (k_{et}) of the electron transfer in the range of $3.8\text{--}9.8 \times 10^{10} \text{ s}^{-1}$ for three different EV^{2+}

concentrations (4–32 μM) or an average rate constant of $6.5 \times 10^{10} \text{ s}^{-1}$. This rate constant for electron transfer is lower than the one obtained for electron transfer between oleic acid/oleylamine capped CsPbBr_3 QDs and viologen ($k_{\text{et}} = 3.6 \times 10^{11} \text{ s}^{-1}$).¹⁵ As discussed in the emission quenching experiments, the surface interactions play a role in dictating the kinetics of interfacial electron transfer.

It is evident that modification of CsPbBr_3 with CdS slows down the electron transfer rate. Schemes 1A and B illustrate the two scenarios for achieving electron transfer, *viz.*, without and with CdS modification. As shown earlier, when methyl viologen is directly bound to CsPbBr_3 , the charge separation is extended, with electrons residing in the viologen moiety and holes residing within the CsPbBr_3 QDs.¹⁴ This extended charge separation as a bound pair was manifested as a long-lived transient bleach component. In contrast, the electron transfer with $\text{CsPbBr}_3\text{-CdS}$ is mediated through the CdS layer as expected by the quasi-type II band alignment (Scheme 1C). The reduced ethyl viologen (EV^{2+}) is no longer directly bound to the CsPbBr_3 , but instead is now linked to the CdS layer. Similar CdS-mediated electron transfer has been observed in CdSe–CdS heterostructures.⁵¹ The bleaching recovery accelerates with increasing viologen concentrations as the electrons are depleted from the CsPbBr_3 core, mediated through CdS. The distinct difference between the two bleach recovery kinetics observed with pristine CsPbBr_3 and $\text{CsPbBr}_3\text{-CdS}$ QDs further highlights



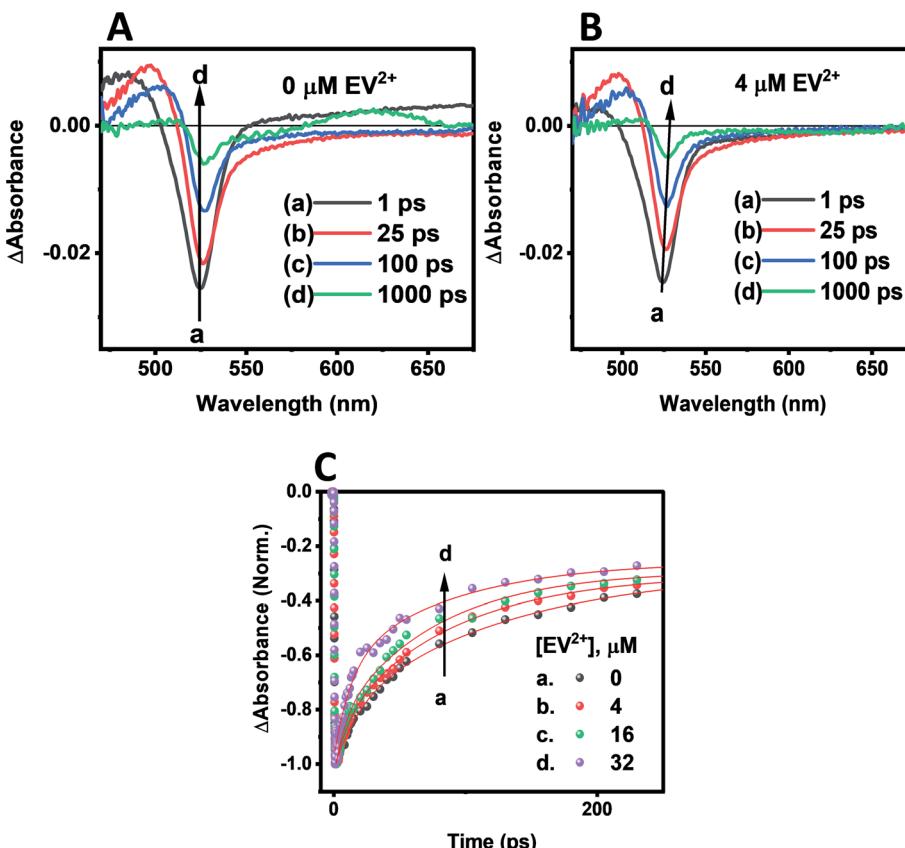


Fig. 5 Transient absorption spectra of CsPbBr_3 – CdS suspension in toluene (≈ 16 nM): (A) without EV^{2+} and (B) with $4 \mu\text{M}$ EV^{2+} . The difference absorption spectra were recorded following a 400 nm laser pulse ($16 \mu\text{J cm}^{-2}$) excitation. (C) Kinetic traces of the CsPbBr_3 – CdS bleach recovery at 527 nm in the absence and increasing concentrations of EV^{2+} .

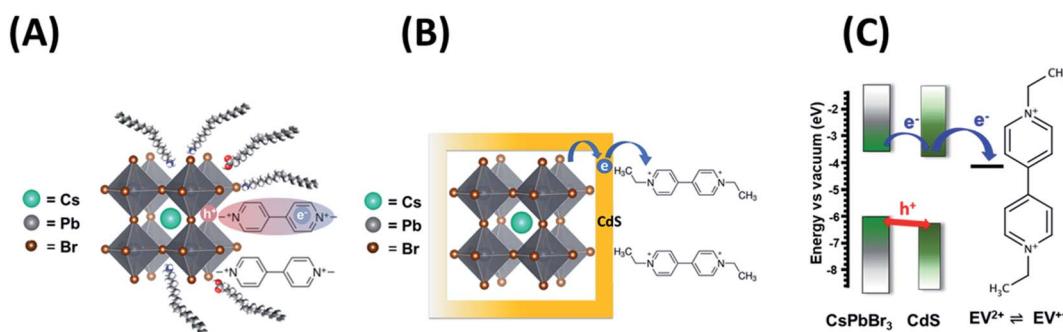
the electron mediation of CdS in promoting interfacial electron transfer in the heterostructure (Scheme 1C).

Steady state photolysis and Fermi level equilibration

Although there have been several studies that demonstrate the photocatalytic properties through product identification,^{17,18,36,61} direct spectroscopic identification of intermediates or electron transfer products is rather limited. If indeed the CsPbBr_3 – CdS heterostructure is responsible for photocatalytic

reduction, we should be able to observe the buildup of stable viologen radical under continuous photoirradiation.^{62,63}

Deaerated CsPbBr_3 – CdS nanocrystal suspensions in toluene : ethanol ($85 : 15$ v/v%) mixed solvent containing $100 \mu\text{M}$ of EV^{2+} were subjected to steady-state visible light illumination (>400 nm; 200 mW cm^{-2}). Absorbance spectra were recorded periodically during the steady state photolysis experiment. Fig. 6A shows a representative difference absorbance spectrum, which shows distinct peaks at 405 nm and 608 nm,



Scheme 1 Photoinduced electron transfer between CsPbBr_3 and viologen (A) without and (B) with mediation through a CdS layer. (C) Energetic diagram showing the flow of charge carriers and CdS mediated reduction of viologen.



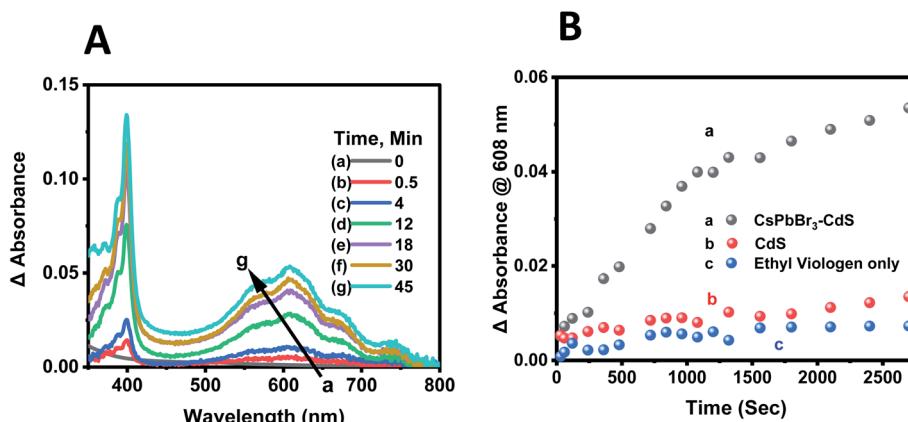


Fig. 6 (A) Difference absorbance spectra recorded during steady state photolysis corresponding to $\text{EV}^{+•}$ formation. The $[\text{CsPbBr}_3\text{-CdS} + \text{EV}^{2+}]$ sample before excitation was used as a reference. Absorption spectra were recorded at different light exposure times. (B) Evolution of $\text{EV}^{+•}$ formation monitored via the growth of absorbance at 608 nm with $\text{CsPbBr}_3\text{-CdS}$. Control experiments with CdS photocatalyst and in the absence of any photocatalyst are also shown. The initial concentration of EV^{2+} was 100 μM . Experiments were carried out in deaerated toluene/ethanol 85 : 15 v/v% with visible light excitation (Xe lamp, 400 nm long pass filter, 200 mW cm^{-2}).

corresponding to the absorbance of $\text{EV}^{+•}$.⁶⁴ Fig. 6B shows the growth of the 608 nm absorption over time, which attains a plateau after about 45 minutes. The steady state concentration of $\text{EV}^{+•}$ increased with increasing concentration of EV^{2+} . From the extinction coefficient of $\text{EV}^{+•}$ at 608 nm ($\epsilon = 1.4 \times 10^4 \text{ M}^{-1} \text{ cm}^{-1}$)^{63,64} we can determine the steady state concentration of the electron transfer product. The quantum yield (QY) of $\text{EV}^{+•}$ formation ($\Phi_{\text{EV}^{+•}}$) was determined using potassium ferrioxalate actinometry.^{51,65} Details of the actinometry experiments are presented in the ESI.† In the present case we obtain a maximum quantum efficiency ($\Phi_{\text{EV}^{+•}}$) of 8.4% for the electron transfer. Although the initial electron transfer yield as monitored from the transient absorption could be as high as 73%,¹⁵ the back-electron transfer during the steady state irradiation makes the net electron transfer yield lower than the value obtained immediately after laser pulse excitation.

The steady state concentration of the reduction product $\text{EV}^{+•}$ is dictated by the forward and back electron transfer processes. Because of the use of ethanol as a hole scavenger, the back electron transfer rate constant (k_{bet}) is significantly lower in the present experiments. The redox couple in contact with a semiconductor surface undergoes Fermi level equilibration that is dependent on the position of the conduction band and the potential of the redox couple. The equilibrium concentration of the reduced and oxidized species can be used to determine the flat band potential of the semiconductor. The steady state concentration of $[\text{EV}^{+•}]_{\text{ss}}$ is indicative of charge equilibration between the $\text{CsPbBr}_3\text{-CdS}$ nanocrystals and $\text{EV}^{2+}/\text{EV}^{+•}$ couple. We employed the Nernst equation (expression (6)) to obtain the flat band potential of the $\text{CsPbBr}_3\text{-CdS}$ heterostructure.

$$E_{\text{CB}} = E_{\text{FB}} = E^{\circ} - \left(\frac{0.059}{n} \right) \log \frac{[\text{EV}^{+•}]_{\text{ss}}}{[\text{EV}^{2+}]_0 - [\text{EV}^{+•}]_{\text{ss}}} \quad (6)$$

By substituting the redox potential of the $\text{EV}^{2+}/\text{EV}^{+•}$ couple, $E^{\circ} = -0.449 \text{ V vs. NHE}$,⁶⁴ and the steady state concentration values of the $\text{EV}^{+•}$ (3.65 μM) and EV^{2+} ($[\text{EV}^{2+}] = [\text{EV}^{2+}]_0 - [\text{EV}^{+•}]_{\text{ss}}$

= 96.35 μM), we obtain $E_{\text{FB}} = -0.365 \text{ V vs. NHE}$. It is evident that the close lying conduction band potential of the $\text{CsPbBr}_3\text{-CdS}$ heterostructure limits the one electron transfer to EV^{2+} (Scheme 1C). It should be noted that the flat band potential of the $\text{CsPbBr}_3\text{-CdS}$ heterostructure obtained in this study is based on the charge equilibration between the semiconductor QD and redox couple and may differ from the values obtained from theoretical estimates or bulk material. The estimate of flat band potential provides an estimate of the energetics of the $\text{CsPbBr}_3\text{-CdS}$ heterostructure suspended in the solvent medium to execute photocatalytic processes.

Conclusions

The design of $\text{CsPbBr}_3\text{-CdS}$ heterostructure offers stabilization of perovskite nanocrystals for photocatalytic applications in polar medium. The salient feature of the $\text{CsPbBr}_3\text{-CdS}$ heterostructure is realized through its stability in mixed solvents, remediation of surface defects, and increased emission yield. The stability of the CsPbBr_3 structure in mixed solvents with increased polarity has allowed us to accumulate electron transfer product (reduced viologen) under steady state irradiation conditions with a quantum efficiency of 8.4%. The relatively high electron transfer efficiency observed in the present study shows how a heterostructure design of perovskite nanocrystals plays a crucial role in dictating the photocatalytic properties of stable perovskite nanocrystals.

Data availability

The data is available within the main text and ESI.†

Author contributions

AK: conceptualization, synthesis, measurements, data analysis, validation, visualization, writing first draft and revision. JTD: data analysis, discussions, revision. JC: conceptualization of



material synthesis, preliminary investigation; revision. PVK: conceptualization, method development, funding acquisition, resources, supervision, discussions, writing.

Conflicts of interest

There are no conflicts to declare.

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