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## Control of dominant conduction orbitals by peripheral substituents in paddle-wheel diruthenium alkynyl molecular junctions†

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Control of charge carriers that transport through the molecular junctions is essential for thermoelectric materials. In general, the charge carrier depends on the dominant conduction orbitals and is dominantly determined by the terminal anchor groups. The present study discloses the synthesis, physical properties in solution, and single-molecule conductance of paddle-wheel diruthenium complexes **1<sup>R</sup>** having diarylformamidinato supporting ligands (DARF: *p*-R-C<sub>6</sub>H<sub>4</sub>-NCHN-C<sub>6</sub>H<sub>4</sub>-R-*p*) and two axial thioanisylethynyl conducting anchor groups, revealing unique substituent effects with respect to the conduction orbitals. The complexes **1<sup>R</sup>** with a few different aryl substituents (R = OMe, H, Cl, and CF<sub>3</sub>) were fully characterized by spectroscopic and crystallographic analyses. The single-molecule conductance determined by the scanning tunneling microscope break junction (STM-BJ) technique was in the 10<sup>-5</sup> to 10<sup>-4</sup> G<sub>0</sub> region, and the order of conductance was **1<sup>OMe</sup>** > **1<sup>CF<sub>3</sub></sup>** ≫ **1<sup>H</sup>** ~ **1<sup>Cl</sup>**, which was not consistent with the Hammett substituent constants  $\sigma$  of R. Cyclic voltammetry revealed the narrow HOMO–LUMO gaps of **1<sup>R</sup>** originating from the diruthenium motif, as further supported by the DFT study. The DFT-NEGF analysis of this unique result revealed that the dominant conductance routes changed from HOMO conductance (for **1<sup>OMe</sup>**) to LUMO conductance (for **1<sup>CF<sub>3</sub></sup>**). The drastic change in the conductance properties originates from the intrinsic narrow HOMO–LUMO gaps.

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## Introduction

Molecular junctions, which are composed of two electrodes with a bridging molecule, are fundamental nanostructures for molecule-based devices. Control of the charge carrier transport through the molecular junction, *i.e.* either an electron or a hole, is an essential requirement for thermoelectric materials, where

both p- (hole transport) and n-type (electron transport) materials are necessary for a high performance series circuit.<sup>1,2</sup> Conduction orbitals close to the electrodes' Fermi energy levels determine the charge carrier and are mainly dominated by the terminal anchor groups, which connect to the electrodes.<sup>3,4</sup> The characteristics of the bridging linkers have a subtle effect with respect to the conduction of charge carriers because common organic molecules have a large separation between the HOMO and LUMO energies, which are far away from the Fermi energies of metal electrodes and, thus, structural modifications are often ineffective. As a rare example, Venkataraman and her coworkers have reported that the conductance pathway for dithiophene molecular wires with HOMO conducting thiomethyl anchor groups can be changed from the HOMO (hole) to the LUMO (electron) by insertion of tetra(thiophene-1,1-dioxide-2,5-diyls).<sup>5</sup> The oxidized thiophene part lowers the overall LUMO energy to change the dominant conduction orbitals, but it requires four repeating units. To date, however, control of the dominant conduction pathway *via* simple structural modifications, *i.e.* introduction of substituents, still remains elusive.

Chromium acetate (Cr<sub>2</sub>(OAc)<sub>4</sub>) synthesized in 1844 (ref. 6 and 7) and structurally characterized in 1951 was the first example of a dinuclear paddle-wheel metal complex (Fig. 1).<sup>8</sup> Since then, various kinds of paddle-wheel dimetal complexes

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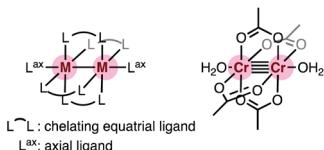
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**Fig. 1** Schematic representation of a paddle-wheel complex and the structure of chromium acetate with two axial water ligands

have been studied extensively, especially by Cotton,<sup>9,10</sup> because they exhibit intriguing electrochemical, magnetic, and optical properties derived from the metal–metal interaction.<sup>11,12</sup> Recent advances in single-molecule conductance measurements using break-junction (BJ) methods allow scientists to unveil the conductance properties of inorganic paddle-wheel complexes.<sup>13</sup> For example, Chen and Peng reported that the conductance of the pentanuclear paddle-wheel complexes having the oligo- $\alpha$ -pyridylamine ligands and isothiocyanate anchor groups (Fig. 2a)<sup>14</sup> was dependent on the embedded metal atoms (M = Ni, Co, and Cr) as well as the metal–metal bond order. The single-molecule conductance of the linearly arranged undeca-nickel complex with the tetra-naphthyridyltriamide ligand was also determined (Fig. 2b).<sup>15</sup> On the other hand, carrier transport through the equatorial paddle-wheel ligands in the dimolybdenum complex was also reported recently (Fig. 2c).<sup>16</sup>

Organometallic paddle-wheel diruthenium complexes with two axial alkynyl ligands constitute an interesting class of molecules for materials science due to the narrow HOMO-LUMO gaps and the rich redox properties. The first example of this class of complexes,  $\text{Ru}_2(\text{ap})_4(\text{C}\equiv\text{CPh})$ , ( $\text{ap}$  = 2-anilinopyridinate), was reported by Cotton in 1986.<sup>17</sup> Later on, Ren *et al.* extensively explored their application in molecular devices.<sup>18-22</sup>

To date, however, the single-molecule conductance measurement of organometallic paddle-wheel complexes remains to be studied. We describe, herein, the results of the single-molecule conductance study for a series of paddle-wheel diruthenium molecular wires **1<sup>R</sup>** with diarylformamidinato supporting ligands (DARF: *p*-R-C<sub>6</sub>H<sub>4</sub>-NCHN-C<sub>6</sub>H<sub>4</sub>-R-*p*) and two axial thioanisylethynyl ligands (Fig. 2 bottom; R = OMe, H, Cl, and CF<sub>3</sub>), in which the thiomethyl anchors show high affinity toward gold electrodes. The present study revealed a unique substituent effect of the aryl moieties in the supporting DARF ligands on the conductance, suggesting the dominant conduction band as well as the charge carrier conductance changed by the substituents.

## Results and discussion

## Synthesis and characterization

We prepared diruthenium dialkynyl complexes **1<sup>R</sup>** ( $R = p\text{-OMe}$ ,  $H$ ,  $p\text{-Cl}$ , and  $p\text{-CF}_3$ ) following the method previously reported by Ren *et al.*<sup>23–25</sup> Treatment of monochloro complexes **2<sup>R</sup>** with an excess amount of lithiated 4-ethynylthioanisole afforded **1<sup>R</sup>** in 17–49% yields. Conventional spectroscopic methods provide characteristic structural information of **1<sup>R</sup>**. The diruthenium dialkynyl complexes **1<sup>R</sup>** are diamagnetic, and the single sets of  $^1\text{H}$  NMR signals for the thioanisyl and DArF ligands indicate highly symmetrical structures of **1<sup>R</sup>** in solution (Fig. S1a, S2a, S4a, and S5a†). HR-ESI-TOF-MS spectra of **1<sup>R</sup>** contain molecular ion peaks corresponding to the deprotonated forms (Fig. S1e, S2f, S4h, and S5e†), and the  $\nu(\text{C}\equiv\text{C})$  vibrations appear in the range of 2092–2097  $\text{cm}^{-1}$  (Fig. S1f, S2g, S4i, and S5f†). Purple crystals were grown by slow evaporation of mixed  $\text{CH}_2\text{Cl}_2$ /hexane solutions containing **1<sup>H</sup>**, **1<sup>Cl</sup>**, and **1<sup>CF3</sup>**. Molecular

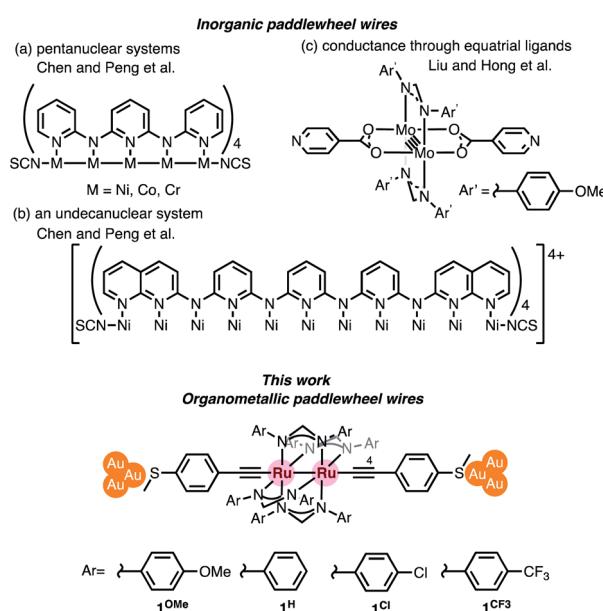
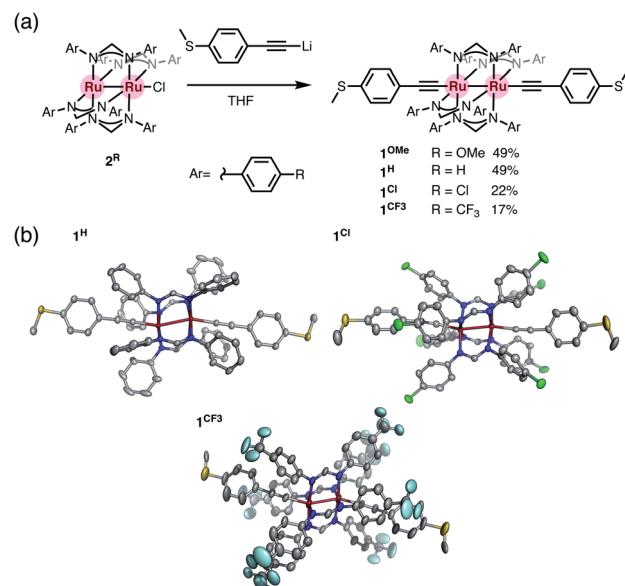


Fig. 2 Inorganic and organometallic molecular wires with paddle-wheel structural motifs



**Fig. 3** (a) Synthesis of  $\mathbf{1}^R$  ( $R = \text{H}$ ,  $\text{Cl}$ , and  $\text{CF}_3$ ). (b) Molecular structures of  $\mathbf{1}^{\text{H}}$ ,  $\mathbf{1}^{\text{Cl}}$ , and  $\mathbf{1}^{\text{CF}_3}$  with thermal ellipsoids at the 50% probability level. Hydrogen atoms and disordered parts are omitted for clarity.

structures of  $\mathbf{1}^{\mathbf{R}}$  ( $\mathbf{R} = \text{H, Cl, and CF}_3$ )<sup>†</sup> shown in Fig. 3 are very similar to each other regardless of the substituents. The  $\text{S}\cdots\text{S}$  distance is around 21 Å, and the Ru–Ru bond distance is *ca.* 2.55 Å ( $\mathbf{1}^{\mathbf{H}}$ : 2.5414(3) Å,  $\mathbf{1}^{\mathbf{Cl}}$ : 2.5468(3) Å, and  $\mathbf{1}^{\mathbf{CF}_3}$ : 2.5521(7) Å), while short Ru–C≡ bond lengths ( $d_{\text{Ru–C}} = 1.97\text{--}1.99$  Å) are observed. These structural features are typical for dialkynyl-Ru<sub>2</sub>(DArF)<sub>4</sub> complexes with no net formal  $\sigma$  bond between the two Ru centers with the  $\pi_4\delta_2\pi_4^*$  electronic configuration, where the electrons in the  $d_{z^2}$  orbitals are used for the Ru–C bonds.<sup>26</sup> The slightly bent Ru–Ru–C≡ linkages (154–166°) result from the second-order Jahn–Teller (SOJT) distortion (see below) to gain thermodynamic stability.<sup>25,27,28</sup>

### STM-BJ study

We carried out single-molecule conductance measurements by the scanning tunneling microscope break-junction (STM-BJ) method.<sup>29</sup> In this method, a gold STM tip repeatedly moves in and out on a gold electrode, and the molecular junction accidentally forms during this process when molecules are present in the solution. The single process gives a long-conductance trace. A tetraglyme solution of  $\mathbf{1}^{\mathbf{R}}$  ( $\mathbf{R} = \text{OMe, H, Cl, and CF}_3$ , 0.25 mM) was soaked into the gold substrate and the break-junction measurements were performed using a gold STM-tip with 100 mV bias voltage. Individual traces show the steps in the  $10^{-4}$  to  $10^{-5} G_0$  region (Fig. S6<sup>†</sup>). The 1D histogram constructed from 2000 successive traces reveals the peaks in the  $10^{-4}$  to  $10^{-5} G_0$  region (Fig. 4). The 2D histogram of  $\mathbf{1}^{\mathbf{H}}$  contains features extended up to  $\sim 1.5$  nm, which correspond to the molecular junction formation. Considering the snapback distance ( $\sim 0.5$  nm),<sup>30</sup> the sum of the lengths well agrees with the  $\text{S}\cdots\text{S}$  distance of  $\mathbf{1}^{\mathbf{H}}$ . On the other hand, the substituted analogues showed plateaus at shorter distances ( $\mathbf{1}^{\text{OMe}}$  and  $\mathbf{1}^{\text{Cl}}$ :  $\sim 1.3$  nm and  $\mathbf{1}^{\text{CF}_3}$ :  $\sim 0.6$  nm) probably because the bulkier DArF ligands prevent

the formation of molecular junctions and/or promote the breaking of molecular junctions. The single-molecule conductance was determined to be  $3.4 (\pm 1.0) \times 10^{-4} G_0$  ( $\mathbf{1}^{\text{OMe}}$ ),  $5.7 (\pm 0.7) \times 10^{-5} G_0$  ( $\mathbf{1}^{\mathbf{H}}$ ),  $5.6 (\pm 1.7) \times 10^{-5} G_0$  ( $\mathbf{1}^{\text{Cl}}$ ), and  $1.3 (\pm 0.2) \times 10^{-4} G_0$  ( $\mathbf{1}^{\text{CF}_3}$ ), with the trend of  $\mathbf{1}^{\text{OMe}} > \mathbf{1}^{\text{CF}_3} \gg \mathbf{1}^{\mathbf{H}} \sim \mathbf{1}^{\text{Cl}}$ . This trend is not in accord with the expectation that replacement with electron-withdrawing substituents (*cf.*  $\sigma$ :  $-0.27$  (OMe),  $0$  (H),  $0.23$  (Cl), and  $0.54$  (CF<sub>3</sub>)) deteriorates conductance because molecular wires with thioether anchor groups should prefer HOMO conductance caused by the energy alignment between the Au–sulfur lone pair bonds and the  $\pi$ -conjugated HOMO orbitals.<sup>3,31</sup>

### Cyclic voltammetry and DFT study

To estimate the frontier orbital energies, we carried out electrochemical measurements of  $\mathbf{1}^{\mathbf{R}}$  by cyclic voltammetry (Table 1, Fig. 5, [complex] = 1.0 mM in CH<sub>2</sub>Cl<sub>2</sub>, [NBu<sub>4</sub>][PF<sub>6</sub>] = 0.1 M). The paddle-wheel complexes  $\mathbf{1}^{\mathbf{R}}$  exhibit one 1e-oxidation and two successive 1e-reduction processes. The oxidation and reduction potentials cathodically shifted upon incorporation of the electron-withdrawing substituents ( $E_{1/2}^{-1}/E_{1/2}^{+2}$ ;  $\mathbf{1}^{\text{OMe}}$ : 0.12/–1.22,  $\mathbf{1}^{\mathbf{H}}$ : 0.19/–1.14,  $\mathbf{1}^{\text{Cl}}$ : 0.38/–0.86, and  $\mathbf{1}^{\text{CF}_3}$ : 0.51/–0.65 V (*vs.* FeCp<sub>2</sub>/FeCp<sub>2</sub><sup>+</sup>)). The potential shifts for the reduction processes caused by substitution ( $\Delta E_{1/2}^{-2} = E_{1/2}^{-2} \mathbf{1}^{\text{CF}_3} - E_{1/2}^{-2} \mathbf{1}^{\mathbf{H}} = 0.49$  V) are larger than those for the oxidation processes ( $\Delta E_{1/2}^{-1} = E_{1/2}^{-1} \mathbf{1}^{\text{CF}_3} - E_{1/2}^{-1} \mathbf{1}^{\mathbf{H}} = 0.32$  V).

To obtain further insights into the electronic structures of molecular wires  $\mathbf{1}^{\mathbf{R}}$ , we performed density functional theory (DFT) calculations at the B3LYP/LanL2DZ,6-31G(d) levels combined with the CPCM continuum solvent method (CH<sub>2</sub>Cl<sub>2</sub>) (Table 1, Fig. 6 and S9<sup>†</sup>). The optimized structures well reproduce the X-ray structures with the Ru<sub>2</sub>(DPhF)<sub>4</sub> core structures distorted from  $D_{4h}$  symmetry. The characteristics of the frontier

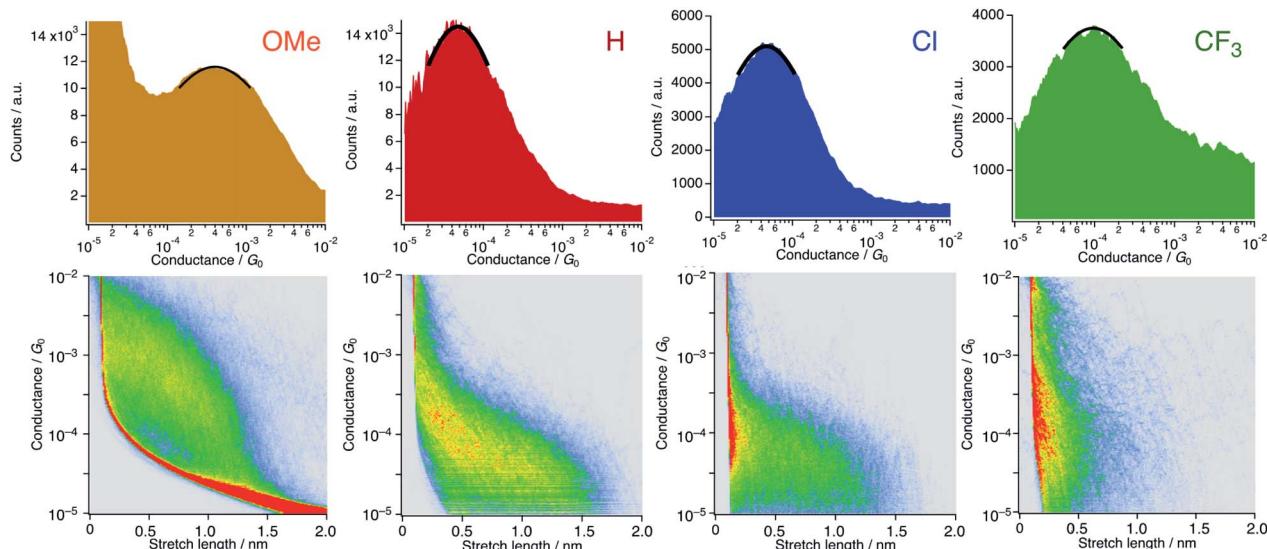


Fig. 4 1D and 2D histograms of  $\mathbf{1}^{\mathbf{R}}$  ( $\mathbf{X} = \text{OMe, H, Cl, and CF}_3$ ) constructed from 2000 successive traces.  $[\mathbf{1}^{\mathbf{R}}] = 0.25$  mM in tetraglyme. No data selection was performed.



Table 1 Summary of experimental and computational data for  $\mathbf{1}^R$  ( $R = \text{OMe}$ ,  $\text{H}$ ,  $\text{Cl}$ , and  $\text{CF}_3$ )

Complex	$E_{1/2}^1$ [V]	$E_{1/2}^2$ [V]	$E_{1/2}^3$ [V]	$\text{HOMO}^{\text{CV}a}$ [eV]	$\text{LUMO}^{\text{CV}b}$ [eV]	$\Delta^{\text{CV}}$ [eV]	$\text{HOMO}^{\text{DFT}}$ [eV]	$\text{LUMO}^{\text{DFT}}$ [eV]	$\Delta^{\text{DFT}}$ [eV]	$G_{\text{exp}}^c$ [ $10^{-5} G_0$ ]	$G_{\text{calc}}^c$ [ $10^{-5} G_0$ ]	$\Delta^{\text{NEGF}}$ [eV]
$\mathbf{1}^{\text{OMe}}$	0.12	-1.22	-2.20	-5.22	-3.88	1.34	-4.92	-3.39	1.53	34 $\pm$ 10	165	-0.13
$\mathbf{1}^{\text{H}}$	0.19	-1.14	-2.11	-5.29	-3.96	1.33	-5.06	-3.60	1.46	5.7 $\pm$ 0.7	1.03	-0.15
$\mathbf{1}^{\text{Cl}}$	0.38	-0.86	-1.85	-5.48	-4.24	1.24	-5.19	-3.84	1.35	5.6 $\pm$ 1.7	5.33	-0.16
$\mathbf{1}^{\text{CF}_3}$	0.51	-0.65	-1.66	-5.61	-4.46	1.15	-5.27	-4.00	1.27	13 $\pm$ 2.4	49.3	-0.10

<sup>a</sup>  $\text{HOMO}^{\text{CV}}$  (in eV) =  $-(E_{1/2}^1 \text{ (in V)} + 5.1)$ . <sup>b</sup>  $\text{LUMO}^{\text{CV}}$  (in eV) =  $-(E_{1/2}^2 \text{ (in V)} + 5.1)$ . <sup>c</sup> Calculated by the DFTNEGF study.

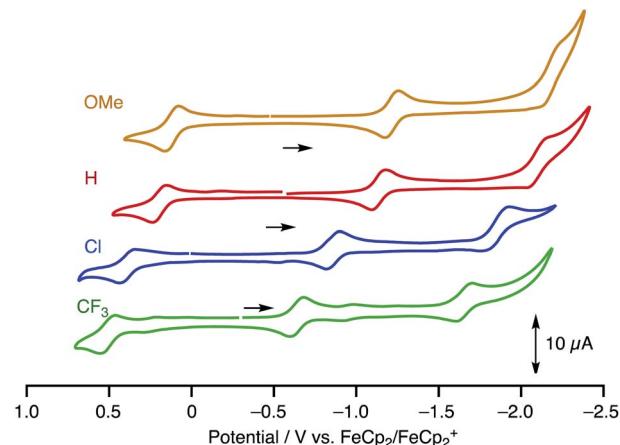


Fig. 5 CV charts of  $\mathbf{1}^R$  ( $R = \text{OMe}$ ,  $\text{H}$ ,  $\text{Cl}$  and  $\text{CF}_3$ ) ( $[\text{complex}] = 1.0 \text{ mM}$  in  $\text{CH}_2\text{Cl}_2$ ,  $[\text{NBu}_4]\text{[PF}_6]$  = 0.1 M, W.E. and C.E. Pt; R.E.  $\text{Ag}/\text{Ag}^+$ , at the scan rate of  $100 \text{ mV s}^{-1}$ ).

orbitals of  $\mathbf{1}^R$  are similar to each other. The highest occupied molecular orbitals (HOMOs), HOMO, HOMO-1 and HOMO-2, of  $\mathbf{1}^R$  are extended over the  $(\text{Ru}-\text{C}\equiv\text{C}-\text{C}_6\text{H}_4\text{SMe})_2$  linkages with the anti-bonding ligand  $\pi^*$  and ruthenium  $d\pi^*$  character, whereas the lowest unoccupied molecular orbitals (LUMOs) are localized on the  $\text{Ru}_2 \delta^*$  orbitals and the non-bonding N-C-N  $\pi$  orbitals of the DArF ligands. One notable feature of the HOMO and HOMO-1 is the overlap between the  $d\pi^*$  orbitals of the Ru cores due to the SOJT distortion as shown in Fig. 6.<sup>32</sup> As a result, the  $\text{Ru}_2$  complexes gain some  $\sigma$  bonding nature between the two Ru atoms, which may enhance the hole transport process through the HOMO orbital.

The HOMO and LUMO energies ( $\text{HOMO}^{\text{DFT}}/\text{LUMO}^{\text{DFT}}$ ) decrease from  $\mathbf{1}^{\text{OMe}}$  ( $-4.92/-3.39 \text{ eV}$ ) to  $\mathbf{1}^{\text{H}}$  ( $-5.06/-3.60 \text{ eV}$ ),  $\mathbf{1}^{\text{Cl}}$  ( $-5.19/-3.84 \text{ eV}$ ) and then to  $\mathbf{1}^{\text{CF}_3}$  ( $-5.27/-4.00 \text{ eV}$ ), and the values are in good agreement with those obtained from the CV analysis ( $\text{HOMO}^{\text{CV}}/\text{LUMO}^{\text{CV}}$ ). The LUMO energies are more sensitive to the substituents on the aryl substituents ( $R$ ) in accord with the CV results probably due to a subtle contribution

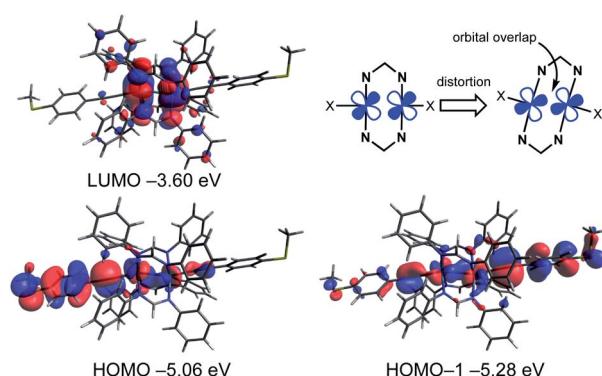


Fig. 6 Selected Kohn-Sham orbitals and their energies for  $\mathbf{1}^{\text{H}}$  obtained by the DFT calculations (B3LYP/LanL2DZ,6-31G(d) and CPCM( $\text{CH}_2\text{Cl}_2$ )) and the schematic representation of orbital overlap of Ru  $d\pi^*$  orbitals causing structural distortion.

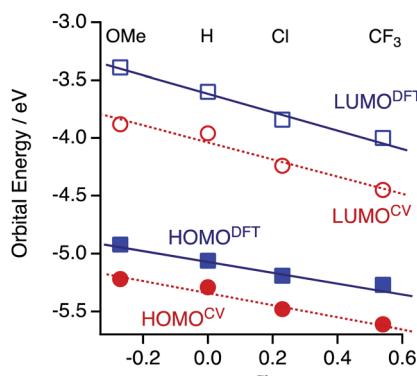


Fig. 7 Plots of the experimentally and theoretically obtained frontier orbital energies against the Hammett substituent constants  $\sigma$ .

of the aryl rings to the LUMO. As a result, the HOMO–LUMO gaps ( $\Delta$ ) decrease upon incorporation of the electron-withdrawing substituents ( $\Delta^{\text{DFT}}$ ;  $1^{\text{OMe}}$ ; 1.53,  $1^{\text{H}}$ ; 1.46,  $1^{\text{Cl}}$ ; 1.35, and  $1^{\text{CF}_3}$ ; 1.27 eV). It is noted that the plots of the HOMO and LUMO energies estimated from the CV and DFT studies against the Hammett substituent constants  $\sigma$  show a good linear relationship (Fig. 7), as reported for the paddle-wheel dimolybdenum complexes with the DArF ligands.<sup>33–35</sup> Thus, the relationship between the molecule and electrode coupling should be considered to understand the unusual conductance trend by substituents.

#### DFT-NEGF study

To gain insights into the substituent effect of the paddle-wheel diruthenium molecular wires and the electronic structures of the molecular junctions, we carried out the DFT-non equilibrium Green's function (NEGF) calculations by using the TranSIESTA code<sup>36</sup> and the level-broadening approach developed by Tada *et al.*,<sup>37</sup> both of which gave similar results. Here, we discuss the results obtained from the TranSIESTA code for  $\text{Au-1}^{\text{R}}\text{-Au}$  ( $\text{Au}$ : gold electrodes).<sup>38</sup> The trend of the theoretically estimated conductance agreed well with that obtained from the break-junction study ( $1^{\text{OMe}} > 1^{\text{CF}_3} \gg 1^{\text{H}} \sim 1^{\text{Cl}}$ ), showing the U-shaped plots (Table 1, Fig. 9a). The transmission spectra reveal many HOMO conduction orbitals within the range of 1.0 eV from the electrode's Fermi energy ( $E_F$ ) in contrast to the lower number of the LUMO conduction orbitals. The conduction orbitals shown in Fig. 8b are altered by the substituents. The dominant transmission eigenchannels of  $\text{Au-1}^{\text{OMe}}\text{-Au}$  and  $\text{Au-1}^{\text{H}}\text{-Au}$  are akin to the HOMOs of  $1^{\text{R}}$  without the electrodes (Fig. 6), whereas LUMO-like  $\delta^*$  orbitals appear for those of  $\text{Au-1}^{\text{Cl}}\text{-Au}$  and  $\text{Au-1}^{\text{CF}_3}\text{-Au}$ . Both the HOMO and LUMO conductance peaks ( $\text{HOMO}^{\text{NEGF}}$  and  $\text{LUMO}^{\text{NEGF}}$ ) shift to the negative side upon replacement of the substituents on the DArF ligands with the electron-withdrawing ones (Fig. 8a and 9b).<sup>39</sup> For  $1^{\text{OMe}}$ , the  $\text{HOMO}^{\text{NEGF}}$  (−0.13 eV) is closer to  $E_F$  than the  $\text{LUMO}^{\text{NEGF}}$  (0.18 eV), while the opposite trend is noted for  $1^{\text{CF}_3}$  ( $\text{HOMO}^{\text{NEGF}}$ : −0.10 and  $\text{LUMO}^{\text{NEGF}}$ : 0.04). These results clearly indicate that the main conduction orbital changes from the HOMO (for  $1^{\text{OMe}}$ ) to the LUMO (for  $1^{\text{CF}_3}$ ), leading to the conductance of  $1^{\text{CF}_3}$  being

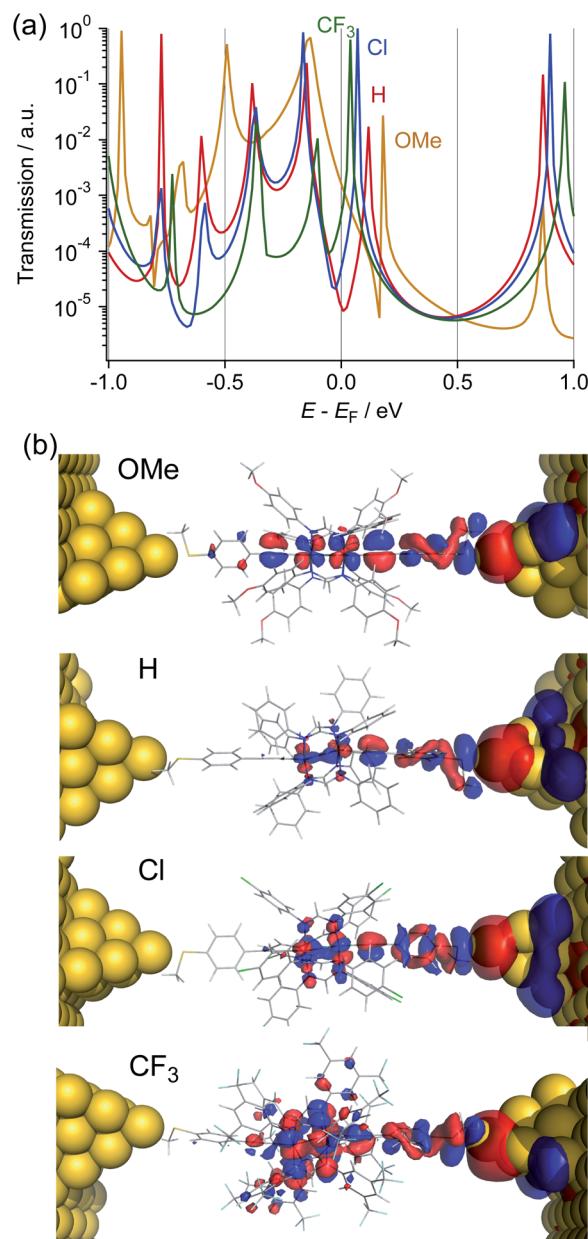


Fig. 8 (a) Transmission spectra and (b) the dominant transmission eigenchannel at the Fermi level of  $\text{Au-1}^{\text{R}}\text{-Au}$  ( $\text{R} = \text{OMe}, \text{H}, \text{Cl}$ , and  $\text{CF}_3$ ).

higher than that of  $1^{\text{H}}$  and  $1^{\text{Cl}}$ . It is also noted that even the LUMO having the  $\text{Ru}_2\delta^*$  orbitals could be an efficient carrier transport pathway, though there is negligible contribution of the thioanisylethynyl ligands in  $1^{\text{R}}$ . The effect of the substituents on the single-molecule conductance and transmission spectra reported so far was moderate or, sometimes, negligible because of their large HOMO–LUMO energy gaps.<sup>40–42</sup> The observed unique substituent effect should result from the fact that (1) the HOMO and LUMO conduction orbitals of  $1^{\text{R}}$  are located very close to  $E_F$  and, accordingly, the HOMO–LUMO gaps become very narrow, and (2) the LUMO energies being

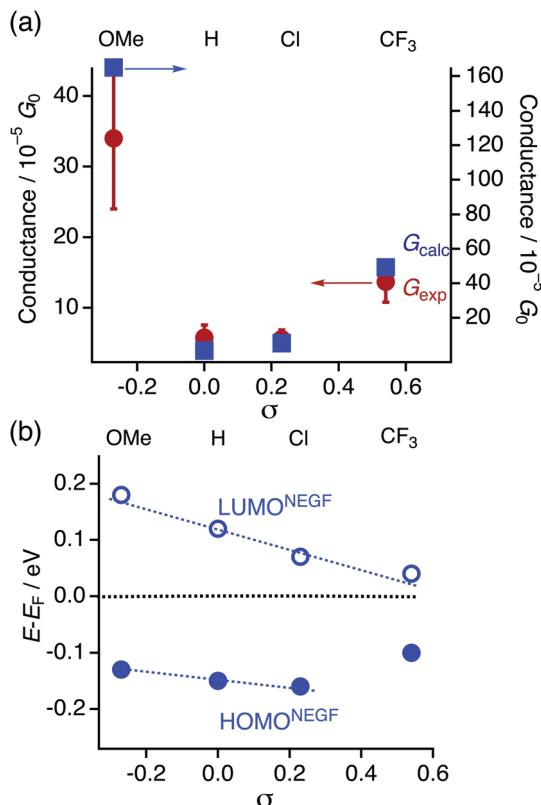


Fig. 9 Plots of (a) experimentally ( $G_{\text{exp}}$ ) and theoretically obtained conductance ( $G_{\text{calc}}$ ) and (b) the  $\text{HOMO}^{\text{NEGF}}$  and  $\text{LUMO}^{\text{NEGF}}$  against Hammett substituent constants  $\sigma$ .

more sensitive to the substituents causes the change of the conduction orbital from the HOMO to the LUMO.

## Conclusion

We have reported the synthesis, electrochemical properties, and single-molecule conductance of paddle-wheel diruthenium complexes **1<sup>R</sup>** ( $R = \text{OMe}$ ,  $\text{H}$ ,  $\text{Cl}$ , and  $\text{CF}_3$ ) having thio-anisylethynyl axial ligands. The unique trend of the single-molecule conductance caused by the replacement of the substituents on the DArF ligands with electron-withdrawing ones results from the change of the dominant conduction orbital from the HOMO (**1<sup>OMe</sup>**) to the LUMO (**1<sup>CF3</sup>**) as revealed by the DFT-NEGF study. The drastic change of the conduction orbitals is caused by the narrow HOMO–LUMO gaps of the paddle-wheel diruthenium cores. Thus, the present study reveals the importance of molecular wires embedded with metal–metal bonds with narrow HOMO–LUMO gaps, which could facilely modulate conduction orbitals. Because the control of the conduction orbitals is vital for thermoelectric materials, paddle-wheel diruthenium complexes turn out to be fascinating platforms for these purposes.

## Author contributions

SO, SC, YT, and MA conceived the project. SO, SC, and YT performed experiments. YT, TO, TT, and HT performed theoretical

calculations. SF and TN discussed the STM-BJ results. YT wrote the first draft and MA critically reviewed and edited the manuscript. All authors commented on drafts, and approved the final manuscript.

## Conflicts of interest

There are no conflicts to declare.

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