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Tuning phenoxyl-substituted diketopyrrolopyrroles from quinoidal to biradical ground states through (hetero-)aromatic linkers†

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Strongly fluorescent halochromic 2,6-di-*tert*-butyl-phenol-functionalised phenyl-, thienyl- and furyl-substituted diketopyrrolopyrrole (DPP) dyes were deprotonated and oxidised to give either phenylene-linked DPP1^{••} biradical ($y_0 = 0.75$) with a singlet open shell ground state and a thermally populated triplet state ($\Delta E_{ST} = 19$ meV; 1.8 kJ mol⁻¹; 0.43 kcal mol⁻¹) or thienylene/furylene-linked DPP2q and DPP3q compounds with closed shell quinoidal ground states. Accordingly, we identified the aromaticity of the conjugated (hetero-)aromatic bridge to be key for modulating the electronic character of these biradicaloid compounds and achieved a spin crossover from closed shell quinones DPP2q and DPP3q to open shell biradical DPP1^{••} as confirmed by optical and magnetic spectroscopic studies (UV/vis/NIR, NMR, EPR) as well as computational investigations (spin-flip TD-DFT calculations in combination with CASSCF(4,4) and harmonic oscillator model of aromaticity (HOMA) analysis). Spectroelectrochemical studies and comproportionation experiments further prove the reversible formation of mixed-valent radical anions for the DPP2q and DPP3q quinoidal compounds with absorption bands edging into the NIR spectral region.

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1. Introduction

Despite the long history,¹ open shell organic molecules have gained tremendous interest in recent years due to their unique electronic, magnetic and optical properties, which mainly arise from weakly coupled electron spins.² Biradicals are the smallest units to allow investigations of intramolecular spin interactions.³ They commonly feature small HOMO–LUMO energy gaps,⁴ enhanced second hyperpolarisability,⁵ redox amphoterism⁶ and large two-photon absorption cross-sections.⁷ While on the one hand these properties make biradicals promising candidates for organic spintronics,⁸ molecular magnetism,⁹ energy storage¹⁰ and electronic devices,¹¹ it is the unpaired electron spins, on the other hand, that drastically increase the reactivity and accelerate decomposition. Therefore, incorporation of biradicals into Kekulé-type quinoidal resonance structures is a common design strategy to stabilise these inherently highly reactive species. However, this brings up the question

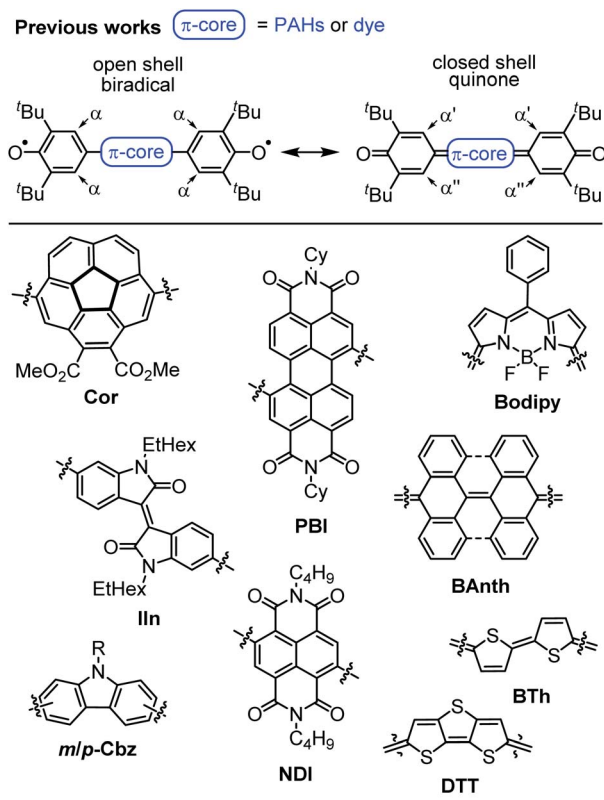
which factors favour biradicals and which ones quinodimethanes (\approx biradicaloids), as well as how to distinguish experimentally between both of these singlet states.² In general, quinodimethane-like compounds can be described by closed shell Kekulé structures, as open shell biradicals or as a superposition of both.^{12,13} Singlet biradicals are molecules with a singlet multiplicity of the lowest energy state but an open shell configuration and are described by the respective biradical character y_0 , which ranges from zero for closed shell to one for completely open shell compounds.² They are further characterised by the singlet triplet energy gap ΔE_{ST} , which typically lies between 10 and 500 meV ($=0.96$ – 48 kJ mol⁻¹; 0.23 – 11.5 kcal mol⁻¹).^{14–17} Within this material class, biradical(oid)s composed of π -systems substituted with two phenoxyl units are among the most outstanding representatives, due to the broad variety of accessible structures, which makes them an ideal model system for experimental and theoretical investigations (Scheme 1).^{18–25} In order to gain insight into the biradical character of these organic compounds, it is crucial to experimentally distinguish the biradical and quinone singlet state. As energy differences between these states tend to be very narrow, an even more careful choice of analytical methods is essential to uncover hidden discrepancies and draw solid conclusions. In general, biradicals show significantly broadened ¹H NMR spectra^{18,19} and/or pronounced EPR signals^{19,20} due to thermal triplet state population. Quinones, in contrast, are commonly EPR silent and characterised by sharp NMR resonance

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Scheme 1 (Top) Open and closed shell resonance structures of Kekulé type biradicals and biradicaloids and (bottom) examples of π -extended quinones and biradicals based on planar and bowl shaped polycyclic aromatic hydrocarbons, pigment chromophores and oligomeric heteroaromatics.^{18–25}

signals.^{21,22} Additionally, the (in)equivalency of phenyl protons α (Scheme 1, top) can be used as an indicator for the molecules' rigidity.^{20,22} Furthermore, in order to experimentally access the singlet–triplet energy gap (ΔE_{ST}), SQUID (superconducting quantum interference device) or variable temperature (VT) EPR can be utilised, of which the latter one commonly is performed in solution.^{19,22} In the solid state, further effects like enhanced intermolecular interactions have to be considered.²³ X-ray analysis is a powerful method to shine light on the electronic ground state, since it allows an experimental bond length determination. Accordingly, the appearance of a distinct bond length alternation (BLA) is commonly a clear hint towards closed shell quinones,^{20,22} whereas its absence indicates a dominating open shell biradical character.²⁴

Like for the experimental counterpart, also the theoretical description of biradicals requires sophisticated methods, since conventional (single reference) DFT is mainly suitable for closed shell compounds, monoradicals, decoupled pairs of local doublets or high spin states, but rather inappropriate to properly describe low spin (*i.e.* singlet) states in biradicals.^{22,26} Nevertheless, in case of very weakly interacting spin centres in biradicals, conventional DFT can lead to reasonable results in accordance with experimental findings (like for isoindigo derivative **IIIn**),¹⁹ but has to be treated with caution. Due to an

increasing number of compounds claimed to be open shell and a broad variety of methods applied – experimentally as well as theoretically – it is more than ever important to study molecules comprehensively because otherwise a comparison of investigated molecules like those collected in Scheme 1 becomes impossible. In this regard, only investigations applying the same theoretical (and ideally identical experimental) methods to all molecules allow relevant conclusions.²⁷ Unfortunately, an increasing number of recent publications still utilises isolated characterisation methods and (completely) neglects double checking of the gained results by comprehensive methods.

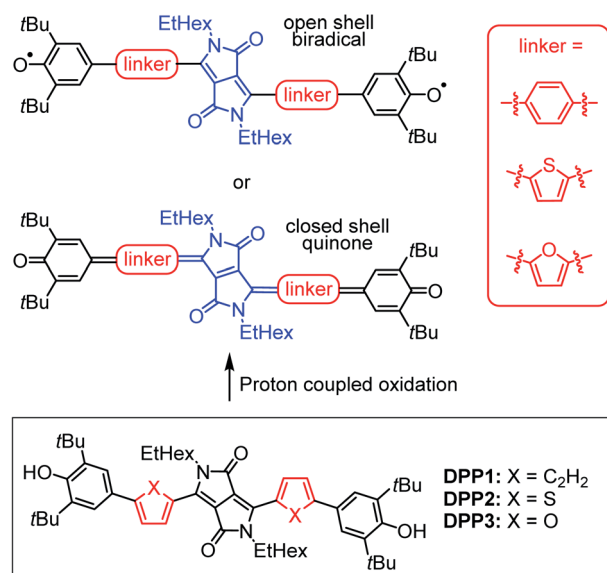
In the past, several factors influencing the electronic ground state and hence enabling a spin crossover from closed shell quinoidal to open shell biradical of π -conjugated compounds derived from quinodimethanes in general and twofold phenoxyl substituted chromophores in particular have been investigated. In this regard, the impact of the number of bridging phenyl units,^{28,29} the connecting π -core (Scheme 1, bottom),^{18–21,24,25} hetero atom effects, the number of attached (donor) substituents and the steric demand of *ortho*-positioned alkyl chains protecting the radical centres in such systems^{30–32} has been investigated. Furthermore, based on numerous longitudinally and laterally extended zethrenes,^{33–36} Wu and coworkers as well as Juriček and coworkers have demonstrated, that the size of the π -scaffold can significantly influence the biradical character of these polycyclic aromatic hydrocarbons (PAH). Likewise, Baumgarten, Feng and Müllen and in particular Haley and coworkers utilised structural isomerism and the number of Clar sextets^{37,38} in various (di)indenoacenes^{27,39–43} to rationally fine-tune the ΔE_{ST} of these compounds. However, a spin crossover has so far been mainly observed upon incremental elongation of an acene-like series.⁴⁴

In this work, we present the first detailed study on a spin crossover for this important class of twofold phenoxyl functionalised π -scaffolds. Toward this goal we chose diketopyrrolopyrrole (DPP) as a core unit and carried out a simple linker variation, while maintaining the molecular size (Scheme 2). DPPs are highly versatile chromophores with tunable optical and electronic properties^{45–49} and outstanding chemical stability which explains their wide application in coatings⁵⁰ and photoelectric devices.^{51–59} Furthermore, especially quinoidal DPPs and related oligo (hetero-)aromatic compounds have recently raised great attention as redox amphoteric dyes,^{60,61} in organic thin film transistors (OTFTs)^{62–65} or as near infrared (NIR) emitters.⁶⁶ As DPPs are known both in the “aromatic” and the quinoidal conjugation, they are hence predestinated to study questions related to the biradical/quinone form.

As demonstrated by a recent contribution from Zheng and co-workers³² reporting on a similar series of DPPs (including **DPP2** and **DPP3** and their oxidised derivatives with just different alkyl chains) we were not alone with this idea. However, in contrast to the conclusion of these authors that these two molecules upon oxidation exhibit an open shell ground state, our high level quantum chemical calculations and multifaceted experimental methods revealed just the opposite: the ground state of these compounds is fully closed shell, *i.e.* quinoidal, which could be proven among other techniques by a X-ray



This work: Introduction and variation of an aromatic linker



Scheme 2 (Top) Schematic illustration of DPP bridged biradicals and quinones, as well as (bottom) structures of synthesised derivatives **DPP1–3**.

crystal structure analysis for **DPP3q**. Most interestingly, for the oxidised species from the phenylene-spacer compound **DPP1** (that was missing in the study of Zheng), indeed the desired spin crossover into the biradicaloid ground state can be observed. Thus, our systematic study of these biradicals and quinones with optical and magnetic spectroscopic methods as well as computational investigations with spin-flip TD-DFT calculations in combination with CASSCF(4,4) and harmonic oscillator model of aromaticity (HOMA) analysis allow to rationalise how the aromaticity of the linker unit tunes these DPP dyes from quinoidal to biradicaloid ground states. Accordingly, this is a lucky case where independently acquired results from two different laboratories on partly identical compounds enable insights on the importance of experimental and theoretical methods for obtaining conclusive results in this important field of research.⁶⁷

2. Results and discussion

Diketopyrrolopyrrole derivatives **DPP1–3** were synthesised by Suzuki–Miyaura cross coupling of the respective literature known brominated DPP-precursors **1–3** with boronic ester **4** in 70%, 53% and 64% yield, respectively (Scheme S1a, ESI†). All new compounds were characterised by ¹H and ¹³C NMR spectroscopy as well as high resolution mass spectrometry. Additionally, the structure of bisphenol **DPP1** could be analysed by single crystal X-ray analysis (Fig. S4 and S5, ESI†).⁶⁸ The optical properties of **DPP1–3** in solution and in the solid state were investigated by UV/vis/NIR absorption and steady state fluorescence spectroscopy under ambient conditions (Fig. 1 and Table S3, ESI†).

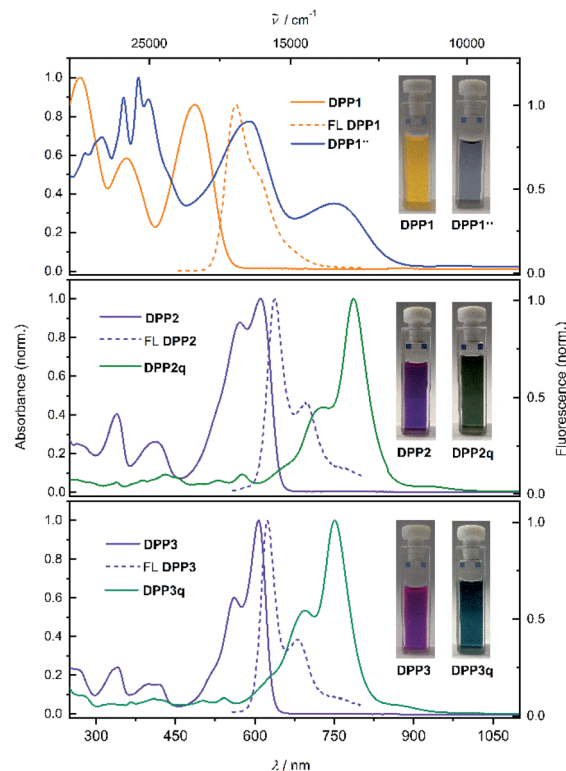


Fig. 1 Normalised UV/vis/NIR absorption spectra (CH_2Cl_2 , $c \approx 10 \mu\text{M}$, rt) of **DPP1** (orange solid line, top), **DPP2** (purple solid line, middle) and **DPP3** (purple solid line, bottom) as well as fluorescence spectra of **DPP1–3** (orange and purple dashed lines, from the top). Also shown are the absorption spectra of biradical **DPP1^{••}** (blue solid line, top), and quinones **DPP2q** (dark green solid line, middle) and **DPP3q** (green solid line, bottom) obtained from proton coupled oxidation of **DPP1–3**. Inset: photographs of the respective cuvettes.

The absorption spectrum of **DPP1** in CH_2Cl_2 (Fig. 1 top, orange solid line) is characterised by broad absorption in the UV and visible spectral range with maxima located at 269, 356 and 486 nm. These absorption maxima are only slightly bathochromically shifted compared to the parent chromophore Ph_2DPP ($\lambda_{\text{abs}} = 466 \text{ nm}$).⁶⁹ In contrast, heteroaromatic derivatives **DPP2** (Fig. 1 middle, purple solid line) and **DPP3** (Fig. 1 bottom, purple solid line) show absorption maxima at 610 nm (**DPP2**) and at 607 nm (**DPP3**) with well resolved vibronic progression, which are significantly redshifted compared to the unsubstituted parent compounds Fu_2DPP and Th_2DPP ($\lambda_{\text{abs}} = 535\text{--}550 \text{ nm}$).^{45,70} The bathochromic shift observed in the UV/vis spectra suggests a pronounced electron-donating effect from the terminal 4-hydroxyphenyl units to the DPP chromophore cores. **DPP1–3** show intense fluorescence at 565 nm for **DPP1** ($\Delta\tilde{\nu}_{\text{Stokes}} = 2580 \text{ cm}^{-1}$, $\Phi_{\text{FL}} = 49\%$, Fig. 1 top, orange dashed line), at 638 nm for **DPP2** ($\Delta\tilde{\nu}_{\text{Stokes}} = 719 \text{ cm}^{-1}$; $\Phi_{\text{FL}} = 38\%$, Fig. 1 middle, purple dashed line) and at 624 nm for **DPP3** ($\Delta\tilde{\nu}_{\text{Stokes}} = 449 \text{ cm}^{-1}$; $\Phi_{\text{FL}} = 48\%$, Fig. 1 bottom, purple dashed line). It is noteworthy that despite of the lowering of the band gap in this series of dyes, fluorescence quantum yields remain high and comparable to those reported in literature for the



respective unsubstituted chromophores Ph₂DPP, Th₂DPP or Fu₂DPP with much larger band gaps.^{45,69,70}

Dehydrogenation of bisphenols occurs easily using proton coupled oxidation which follows a concerted deprotonation-oxidation mechanism.^{18,19,25} Therefore, solutions of **DPP1**–**3** were treated with basic tetra-*n*-butylammonium fluoride (TBAF) in order to study the occurring spectral changes. Gradual changes in the respective absorption spectra were monitored by UV/vis/NIR spectroscopy (Fig. S7, ESI†). Upon stepwise addition of a solution of TBAF in CH₂Cl₂, the intensity of the absorption bands of **DPP1**, **DPP2** and **DPP3** decreases with concomitant rise of absorption bands at longer wavelengths at 704 nm for **DPP1**, at 813 nm for **DPP2** and at 789 nm for **DPP3**. The appearance of these bathochromically shifted absorption bands indicates the formation of the respective dianions **DPP1**^{2−}, **DPP2**^{2−} and **DPP3**^{2−} and can be explained by the strong charge transfer (CT) from the electron-rich phenoxide substituents to the electron-poor DPP core. Such CT bands are commonly observed in halochromic systems.^{18,19,25} Notably, the immediate addition of an equimolar amount of trifluoroacetic acid to freshly prepared solutions of the dianions leads to a complete recovery of the absorption spectral signatures of the corresponding phenols (Fig. S7, ESI†) and hence reveals the reversibility of the deprotonation process. Additionally, during deprotonation titration experiments of the heteroaromatic derivatives **DPP2** and **DPP3**, a concomitant rise of an additional NIR band at 1083 and 1063 nm, respectively, was observed, which does not appear upon titration of **DPP1** (Fig. S7, ESI†). The appearance of these bands indicates the formation of mixed valent radical anions **MV-DPP2** and **MV-DPP3** upon autoxidation. As these mixed valent species arise upon oxidation of dianions **DPP2**^{2−} and **DPP3**^{2−} under ambient conditions, we monitored UV/vis/NIR absorption spectral changes over time (Fig. S8, ESI†) to study the extent of this oxidation process. Within 17 h (**DPP2**^{2−}) and 4 h (**DPP3**^{2−}), the bands attributed to the presence of radical anions **MV-DPP2** and **MV-DPP3** rise continuously, followed by a subsequent decrease. Analogous time-dependent monitoring of **DPP1**^{2−} (Fig. S8a and b, ESI†) merely shows a decrease of CT band intensity upon oxidation, which is much slower than the processes observed for the heteroaromatic derivatives. It is worth to mention, that there is no hint for the formation of an analogous mixed valent species upon autoxidation of **DPP1**^{2−}. Taking the different extent and rates of the autoxidation processes into account, it can be concluded that the phenolate's sensitivity towards oxidation is decreasing in the order **DPP3**^{2−} > **DPP2**^{2−} > **DPP1**^{2−}. This trend was proven to hold true for neutral bisphenols **DPP1**–**3** as well by cyclic voltammetry, although the furane and thiophene derivatives feature a quite similar redox behaviour (Fig. S28 and Table S4, ESI†). Accordingly, the disparity between heterocyclic **DPP2/3** and **DPP2q/3q** (S ↔ O) caused by chalcogen effects is much lower than the differences (S/O ↔ Ph) to the aromatic phenyl system **DPP1**.

Our results indicate, that the first (aut-)oxidation step of **DPP2**^{2−} and **DPP3**^{2−} involves the formation of radical anions **MV-DPP2** and **MV-DPP3** by single electron transfer (SET). Apparently, a further defined oxidation process cannot be

achieved by using ambient oxygen. Therefore, we applied electrochemical oxidation in order to study the stepwise oxidation process and gain further insight into the redox properties. Accordingly, the electro-optical properties of dianions **DPP1**^{2−}, **DPP2**^{2−} and **DPP3**^{2−} were subsequently studied by spectroelectrochemistry (SEC, Fig. 2). Initial spectral changes observed for **DPP2**^{2−} and **DPP3**^{2−} mainly reproduce the transformations already monitored during their autoxidation (Fig. 2c, e and S8, ESI†). However, the NIR bands attributed to radical anions **MV-DPP2** and **MV-DPP3** vanish completely upon further raising the potential to 500 mV and 475 mV (*vs.* Pt pseudo reference electrode, Fig. 2d and f, red solid line), respectively, and simultaneously, the appearance of new absorption bands with maxima at 782 nm (**DPP2**) and 751 nm (**DPP3**) (Fig. 2d and f, blue solid line) can be observed, which resemble by spectral shape and vibronic structure those of **DPP2** and **DPP3**, but are significantly red shifted. The formed species represent quinones **DPP2q** and **DPP3q** (*vide infra*) with outstanding high molar extinction coefficients, *i.e.* tripled and doubled compared to **DPP2** and **DPP3**, respectively (Table S3, ESI†). Also for **DPP1**^{2−} a new optical signature (Fig. 2b, blue solid line) emerges upon electrochemical oxidation with increasing potential to 450 mV, which could not be observed during autoxidation experiments (Fig. S8, ESI†) and can be attributed to the formation of **DPP1**^{••} (*vide infra*). The resulting absorption spectrum is significantly broadened in the visible

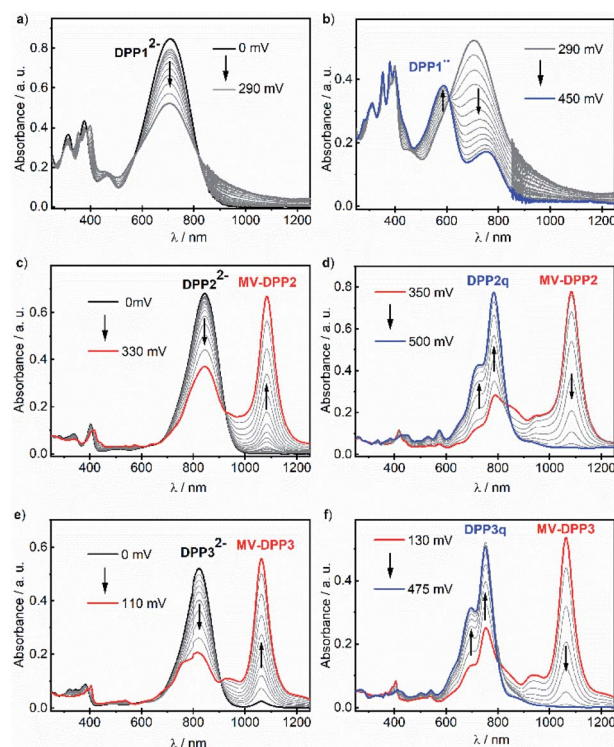


Fig. 2 UV/vis/NIR absorption spectral changes of (a and b) **DPP1**^{2−}, (c and d) **DPP2**^{2−} and (e and f) **DPP3**^{2−} (*c* ≈ 3 mM in CH₂Cl₂, rt, 0.2 M ^{*n*}Bu₄NPF₆) upon electrochemical oxidation in a spectroelectrochemical setup. Arrows indicate spectral changes with increasing positive potential.



region and shows pronounced panchromaticity with intense maxima at 380 and 588 nm and a band of lower intensity in the NIR region (749 nm). The appearance of several isosbestic points for all three dianions **DPP1**–**3**^{2−} indicates clearly defined oxidation processes. Our results indicate the formation of mixed valent radical anions **MV-DPP2** and **MV-DPP3** as well as twofold oxidised species **DPP1**^{••}, **DPP2q** and **DPP3q** by electrochemical oxidation of **DPP1**–**3**.

In order to enable in-depth studies, we attempted to obtain these species on a preparative scale as well by applying a suitable chemical oxidant. Accordingly, biradical **DPP1**^{••} and quinones **DPP2q** and **DPP3q** were synthesised almost quantitatively by oxidation of **DPP1** with potassium ferricyanide and by dehydrogenation of **DPP2** and **DPP3** using lead(IV) oxide (Scheme S1b, ESI†), which could not be applied to **DPP1** as it caused decomposition. All products feature the same UV/vis/NIR absorption spectral signatures as observed for the electrochemically generated ones (Fig. S9, ESI†). By selective reduction of **DPP2q** and **DPP3q** with (Me₂N)₂C=C(NMe₂)₂ (TDMAE), as well as by comproportionation of **DPP2**^{2−} or **DPP3**^{2−} with **DPP2q** or **DPP3q**, respectively, we could furthermore generate mixed valent compounds **MV-DPP2** and **MV-DPP3** (Schemes S8–S11 and Fig. S12, S13, ESI†), which, however, could not be isolated. Nevertheless, the result proves that comproportionation of (extended) quinones and bisphenoxides offers an efficient and convenient access to highly desirable radical anions⁷¹ in appropriately functionalised chromophores.

Whereas **DPP2q** and **DPP3q** could be easily isolated out of solution by evaporation of the solvent (CH₂Cl₂) and even be crystallised (**DPP3q**, Fig. 3), attempts to isolate **DPP1**^{••} failed and significant bleaching of the solution was observed instead even under inert conditions in degassed CCl₄ and more accelerated upon raising temperature to reflux. For this reason, isolation or recrystallization was not successful as well. Hence, to quantify the stability of **DPP1**^{••}, **DPP2q** and **DPP3q** in solution and in the solid state under ambient conditions, time-dependent UV/vis/NIR absorption spectra were recorded. The spectral signatures of **DPP2q** and **DPP3q** in solution and in the solid state remain unchanged over several days (Fig. S10, ESI†), and thus prove the high stability of these compounds. In contrast, **DPP1**^{••} decomposes within minutes in the solid state and within hours in solution as band intensities gradually decrease over time (Fig. S11, ESI†). By fitting the time-dependent data, a minimum half-life of 78 h in CCl₄ and 39 h in CH₂Cl₂ was obtained for **DPP1**^{••} at room temperature (Fig. S11 and Tables S1, S2, ESI†). Such a fast decomposition within the timescale of days compared to stable **DPP2q** and **DPP3q**, can be explained with a considerably higher reactivity of **DPP1**^{••}, presumably caused by the distinct biradical character. Life times in the range of hours to days are indeed typical for phenoxyl based biradicals.^{18,25}

Representatively for both quinones, the solid state structure of **DPP3q** could be determined by single crystal X-ray diffraction analysis (Fig. 3). Furane derivative **DPP3q** features an almost planar π -surface with negligible dihedral twist angles between 0.7(3)° and 1.7(3)° (Fig. 3a and b). In addition, a distinct bond length alternation over the whole chromophore can be observed

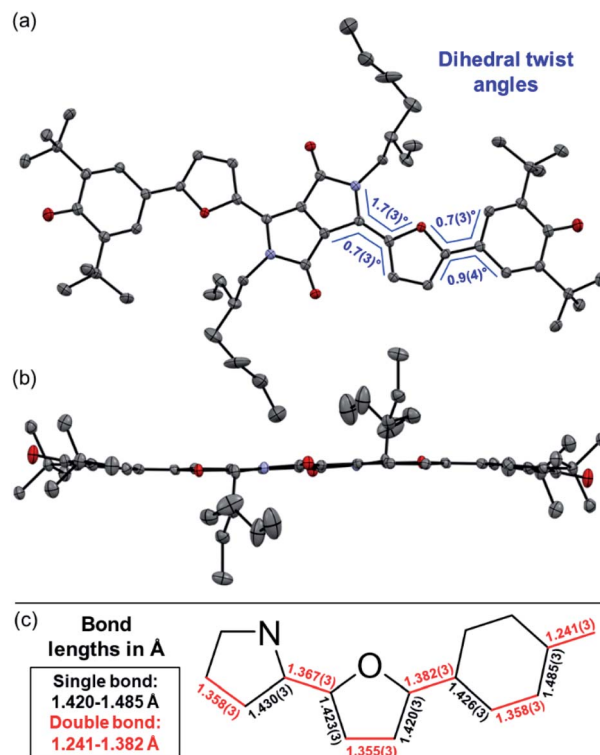


Fig. 3 Solid state molecular structure of **DPP3q** in (a) top view and (b) side view as well as (c) selected bond lengths determined by single crystal X-ray diffraction (ellipsoids set to 50% probability, carbon gray, nitrogen blue, oxygen red). H atoms and solvent (MeOH) molecules are omitted for clarity.

(Fig. 3c), which is well in line with the quinoidal character of **DPP3q** concluded by NMR spectroscopy (*vide infra*). The parallel displaced chromophores with an average π – π -distance of 3.19 Å form staircase like strands, which are oriented in a herringbone-type packing (Fig. S6, ESI†). As **DPP1**^{••} is stable for several hours in solution (CH₂Cl₂) and **DPP2q** and **DPP3q** show no significant decomposition, we were able to investigate the para- and diamagnetic properties of the compounds with electron paramagnetic resonance (EPR) and nuclear magnetic resonance (NMR) spectroscopy (Fig. 4).

DPP2q and **DPP3q** are virtually EPR silent (Fig. S31a and b, ESI†) in CH₂Cl₂ solution, but show a very weak signal in the solid state. However, temperature-dependent solid state EPR measurements in the range of 240 to 300 K did not reveal any significant signal intensity change (Fig. S31c and d, ESI†). Accordingly, no hints for thermal triplet state population were found, which proofs an insurmountable ΔE_{ST} and provides further evidence for a diamagnetic character. In contrast, **DPP1**^{••} features a pronounced EPR signal, which is centred at $g_{iso} = 2.0044$ with a peak-to-peak line width of 3.8 G (Fig. 4a). Temperature-dependent EPR spectroscopy of **DPP1**^{••} additionally allowed for a quantitative view on the energy difference between the singlet and triplet state. Double integration of EPR signals and data fitting according to the Bleaney–Bowers equation revealed a remarkably low ΔE_{ST} of 19 meV (1.8 kJ mol^{−1}; 0.43 kcal mol^{−1}, $2J = 147$ cm^{−1}) and hence



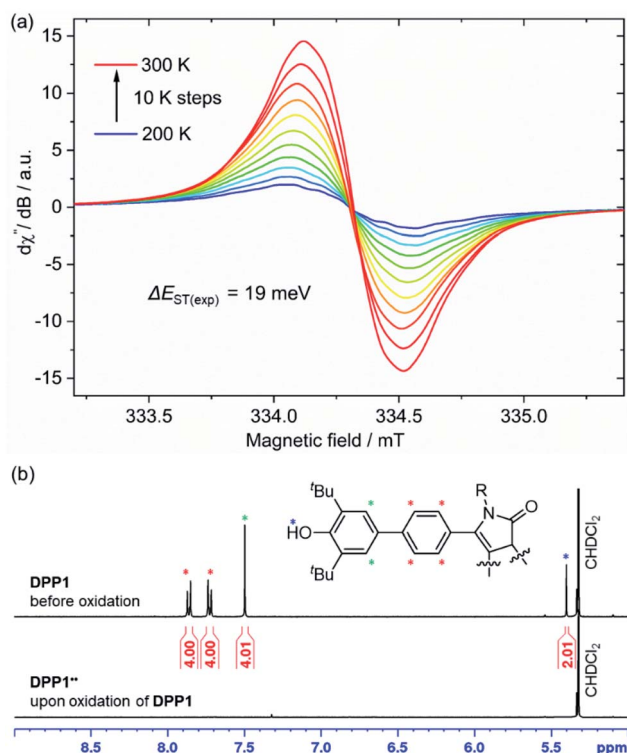


Fig. 4 (a) Temperature dependence of the X-band EPR signal of $\text{DPP1}^{\bullet\bullet}$ in CH_2Cl_2 ($c \approx 1 \text{ mM}$) and (b) aromatic region of the ^1H NMR spectrum (400 MHz) of DPP1 (top) and $\text{DPP1}^{\bullet\bullet}$ (bottom) in CD_2Cl_2 at rt.

indicates almost decoupled spin centres (Fig. S30, ESI†). As EPR signal integrals of $\text{DPP1}^{\bullet\bullet}$ increase with temperature, an open shell singlet ground state can be concluded for $\text{DPP1}^{\bullet\bullet}$. In compliance with the increasing amount of triplet species at higher temperature, also a pronounced broadening of resonance signals in the aromatic region of the respective ^1H NMR spectra is observed, thereby confirming the substantial paramagnetic character of $\text{DPP1}^{\bullet\bullet}$ at room temperature.^{18,19,25} Accordingly, the ^1H NMR spectrum of $\text{DPP1}^{\bullet\bullet}$ in CD_2Cl_2 at 298 K shows very broad signals (Fig. 4b), which significantly gain intensity upon decreasing the temperature to 180 K (Fig. S24, ESI†). This is in accordance with a biradical character in the ground state with a low energy difference to the NMR silent triplet state, which is thermally populated at higher temperature.

In contrast to $\text{DPP1}^{\bullet\bullet}$, sharp signals are detected for DPP2q and DPP3q in the aromatic region of the respective ^1H NMR spectra at room temperature. Notably, the resonance signals for the two heteroaryl protons are separated by 1.45 to 1.76 ppm (Fig. S16–S22, ESI†) with a significantly downfield shifted signal at 9.02 ppm (DPP2), 8.34 ppm (DPP3), 9.38 ppm (DPP2q) and 8.85 ppm (DPP3q). This shift can be explained by hydrogen bonding between a heteroaryl H atom and the carbonyl oxygen atom (Fig. S16–S22, ESI† hydrogen bonded proton highlighted in blue).^{61,62} Accordingly, a *N,S*- and *N,O*-*cis* configuration can be concluded for all heteroaromatic derivatives, which was also corroborated by X-ray structure analysis in the case of DPP3q

(Fig. 3). In addition, the resonance signal of the protons of the phenyl moiety in DPP2 (7.52 ppm) and DPP3 (7.63 ppm) with a singlet multiplicity splits up into two doublet signals for quinones DPP2q and DPP3q (Fig. S16–S22, ESI†). Thus, both protons are chemically not equivalent anymore, which can be explained by the formation of a double bond between the former phenyl and heteroaromatic unit.^{22,32} Therefore, the rotational freedom of the terminal phenyl unit is significantly hindered, as expected due to the quinoidal character of DPP2q and DPP3q . At an elevated temperature of 373 K, the resonance signals assigned to the heteroaryl protons of DPP3q remain sharp, whereas signals of DPP2q show a slight broadening (Fig. S25 and S26, ESI†). In particular, the resonance signals ascribed to the hydrogen bonded protons (*vide supra*) are well detectable, whereas the phenyl proton signal becomes broad in all heteroaromatic quinones.³² Accordingly, the different behavior of phenyl and heteroaryl proton signals upon heating can be explained by a rather rigid DPP core flanked by phenoxy groups with thermally enhanced rotational freedom rather than by a thermal triplet population. Summarising, it can be concluded that the ground state of DPP2q and DPP3q is dominated by a closed shell character with a large, thermally insurmountable singlet–triplet energy gap (ΔE_{ST}).

To further shine light onto the electronic ground state character of $\text{DPP1}^{\bullet\bullet}$, DPP2q and DPP3q , quantum chemical calculations have been performed. In general, biradicals can be defined as molecular systems with two electrons occupying two (almost) degenerate molecular frontier orbitals.² These orbital (near-)degeneracies result in wave functions, which are not dominated by a single configuration, but rather include several leading determinants. Since the accurate description of low-spin states in organic systems with partial open shell contributions requires at least two Slater determinants, biradicals usually cannot be described sufficiently using conventional density functional theory (DFT) or single-reference wavefunction-based methods,²⁶ but considerably more precisely by applying the spin-flip (SF) TD-DFT approach.⁷² This method uses a single-determinant high-spin triplet state as a reference, which is well represented by a single Slater determinant. From that configuration, the target manifold of low-spin states (that is: singlets and low-spin triplet) is generated in a single excitation by applying a linear spin-flipping excitation operator $\hat{R}_{M_s=-1}$:

$$\Psi_{M_s=0}^{s,t} = \hat{R}_{M_s=-1} \Psi_{M_s=1}^t \quad (1)$$

Optimization of the formally first excited state and subsequent analysis of its character allows determination of the electronic configuration in the relaxed ground state structure of the molecule. The geometry optimization of $\text{DPP1}^{\bullet\bullet}$, DPP2q and DPP3q has therefore been carried out in the framework of SF-TD-DFT employing the 50/50 functional (50% Hartree–Fock + 8% Slater⁷³ + 42% Becke⁷⁴ for exchange and 19% VWN⁷⁵ + 81% LYP⁷⁶ for correlation) along with the def2-SVP⁷⁷ basis set (Fig. S35, ESI†). Dihedral angles between the terminal phenoxy and bridging phenylene unit in $\text{DPP1}^{\bullet\bullet}$ are significantly reduced

to 11.14–14.73° compared to the parent bisphenole **DPP1** (30.2–32.2°, Fig. S4, ESI†), which can be explained by increasing conjugation in the dehydrogenated state.²⁴ In accordance with the single crystal X-ray results for **DPP3q**, the relaxed geometries of **DPP2q** and **DPP3q** exhibit negligible dihedral angles of only 0.5/0.6° and 0.2/0.4°, respectively. Furthermore, we observed a distinct, and continuous bond length alternation over the whole chromophore in **DPP2q** and **DPP3q**. In contrast, BLA is broken in **DPP1**^{••} with significantly elongated bonds e and q indicating a much higher biradical character (Fig. 5a).

As evidenced by ¹H NMR spectroscopy and X-ray analysis (*vide supra*), the *N,O*- and *N,S-cis* geometry – on which our calculations are based – is dominating in **DPP2q** and **DPP3q**. However, to complete our geometrical analysis, we performed optimisation of the *N,X-trans* geometries as well (Fig. S37, ESI†). The *trans* geometries were found to be 340 to 400 meV (=32.8–

38.6 kJ mol^{−1}; 7.84–9.22 kcal mol^{−1}) higher in energy, making the *cis* conformation a thermodynamic global minimum. Notably, the impact of the conformation on the bond lengths is very large. Accordingly, BLA is significantly reduced in the hypothetical *trans* geometry which is in good agreement with the respective values obtained by conventional DFT.³² For this reason, the real bond length alternation was underestimated in the earlier study by Zheng and co-workers, which is based on the *trans* geometry.³²

Additionally, analysis of the spin (*S*² = 0) and the natural orbital occupation number was used to determine the electronic ground state configuration and biradical character *y*₀.⁷⁸ Both together indicate that the energetically most favorable electronic configuration is governed by a closed shell occupation of orbitals for **DPP2q** and **DPP3q**, while **DPP1**^{••} is dominated by an open-shell singlet configuration and hence can be described as a singlet biradical.

In order to estimate the singlet-triplet energy gaps and quantify the singlet biradical character *y*₀, single point CASSCF(4,4)/def2-SVP^{75,77–82} calculations were employed for all three relaxed geometries, confirming the SF-TD-DFT results. The resulting energy gap is calculated to be 24 meV (2.3 kJ mol^{−1}, 0.54 kcal mol^{−1}) in case of **DPP1**^{••} and 1.22 eV (118 kJ mol^{−1}, 28.1 kcal mol^{−1})/1.31 eV (126 kJ mol^{−1}, 30.2 kcal mol^{−1}) for **DPP2q**/**DPP3q**, respectively (Fig. 5b). The calculated ΔE_{ST} of **DPP1**^{••} is in good accordance with the value obtained experimentally by temperature-dependent EPR spectroscopy. Accordingly, the value of *y*₀ = 0.75 obtained for **DPP1**^{••} clearly indicates the high biradical character of **DPP1**^{••}, while *y*₀ values of 0.004 and 0.003 validate the closed shell configuration of the quinones **DPP2q** and **DPP3q**, respectively. These findings further support the structures derived from ¹H NMR spectroscopy and X-ray analysis (*vide supra*). The low biradical character calculated for **DPP2q** and **DPP3q** is in contrast to the recent conclusions of Zheng and co-workers³² who reported a pronounced biradical character of *y*₀ = 0.64–0.65 for related dyes bearing just different alkyl chains. However, these values were calculated by employing conventional DFT calculations, which become less precise in case of stronger spin interactions and low spin states (*vide supra*),²⁶ as they are observed in **DPP2q** and **DPP3q**.

For the purpose of investigating the role of aromaticity as a key factor for the stabilization of the open shell configuration in **DPP1**^{••}, the “harmonic oscillator model of aromaticity” (HOMA)^{83,84} value was calculated for all three linkers. HOMA takes into account the deviation of bond length from an “optimal” value expected for a fully aromatic system.^{83,84}

$$\text{HOMA} = 1 - \sum_i \frac{\alpha_{ij}}{N} (R_{\text{Ref}} - R_{i,j})^2 \quad (2)$$

Here, *j* is the atom next to atom *i*, *N* denotes the total number of atoms, α and *R*_{Ref} are pre-calculated constants which are presented in the original article for each type of atom pair.⁸⁴ A perfectly aromatic compound thereby has a HOMA value of 1, whereas a non-aromatic compound has a value of 0 or below. For the phenyl linker in **DPP1**^{••} a value of 0.8 was obtained, which clearly indicates the presence of an aromatic benzene

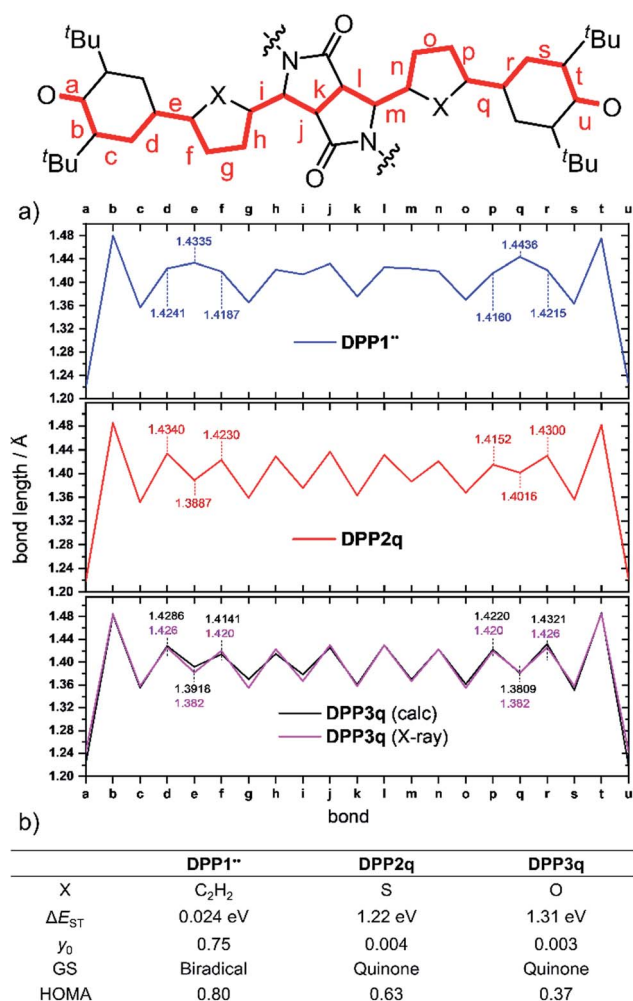


Fig. 5 (a) Bond lengths in Å between both oxygen sites in **DPP1**^{••}, **DPP2q** and **DPP3q** (blue, red and black solid line, from the top) as well as (b) calculated biradical/quinoidal properties of **DPP1**^{••}, **DPP2q** and **DPP3q**, ((SF) TD-DFT and CASSCF(4,4)/def2-SVP level of theory). Also shown are the bond lengths determined experimentally by X-ray diffraction (purple solid line). GS = ground state.



ring, thereby stabilising the biradical configuration, while in **DPP2q** and **DPP3q**, HOMA indices of 0.63 and 0.37, respectively, hint towards a much stronger quinoidal character (Fig. 5b). All theoretical findings are hence in agreement with the experimental results and our interpretation.

3. Conclusion

In conclusion, we reported a series of three with two 2,6-di-*tert*-butylphenoxy groups functionalised diketopyrrolopyrrole dyes **DPP1–3** that on first glance look very similar, but upon deprotonation and oxidation afford electronically very distinct compounds, *i.e.* biradical **DPP1^{••}** and quinones **DPP2q** and **DPP3q**. Our comprehensive optical and magnetic spectroscopic studies demonstrate the spin crossover from closed shell quinones **DPP2q** and **DPP3q** to a singlet open shell biradical **DPP1^{••}** by simple linker variation. Thus, the heteroaromatic thiophene or furane linkers between the DPP core and phenoxy substituent favour a closed shell ground state, thereby endowing **DPP2q** and **DPP3q** with bench stability and quinoidal structures as evidenced by X-ray diffraction (**DPP3q**) and UV/vis/NIR absorption, EPR and NMR spectroscopy as well as by quantum chemical calculations based on the spin-flip TD-DFT and CASSCF(4,4) level of theory. In contrast, the “isolating” strongly aromatic phenylene bridge endows derivative **DPP1^{••}** with biradical properties with a very small singlet–triplet-energy gap of 19 meV (1.8 kJ mol^{-1} ; $0.43 \text{ kcal mol}^{-1}$) and a large biradical character of $y_0 = 0.75$. As a consequence, **DPP1^{••}** loses stability and decomposes within the timescale of days in solution. The aromatic character of the bridging units in **DPP1^{••}**, **DPP2q** and **DPP3q** was investigated using HOMA index values, which show that the established order of aromaticity decrease (phenyl > thiophene > furane) applies for such biradicaloid systems with a central dye unit. Accordingly, we conclude that Clar’s sextet rule³⁸ also offers a design principle to derive open shell colourants with central dye and pigment units. In our example only **DPP1^{••}** with four benzenoid sextets prevailed in an open shell configuration as a biradical, whereas energy gains through non-benzenoid heteroaromatic furane or thiophene linkers were not large enough to counteract the transformation of the π -electron system into a fully conjugated quinoidal scaffold.

Conflicts of interest

There are no conflicts to declare.

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Notes and references

- 1 A. E. Tschitschibabin, *Ber. Dtsch. Chem. Ges.*, 1907, **40**, 1810–1819.
- 2 M. Abe, *Chem. Rev.*, 2013, **113**, 7011–7088.

- 3 M. Baumgarten, *Phys. Status Solidi B*, 2019, **256**, 1800642.
- 4 Z. Sun, Q. Ye, C. Chi and J. Wu, *Chem. Soc. Rev.*, 2012, **41**, 7857–7889.
- 5 K. Fukuda, J.-Y. Fujiyoshi, H. Matsui, T. Nagami, S. Takamuku, Y. Kitagawa, B. Champagne and M. Nakano, *Org. Chem. Front.*, 2017, **4**, 779–789.
- 6 G. E. Rudebusch, G. L. Espejo, J. L. Zafra, M. Peña-Alvarez, S. N. Spisak, K. Fukuda, Z. Wei, M. Nakano, M. A. Petrukhina, J. Casado and M. M. Haley, *J. Am. Chem. Soc.*, 2016, **138**, 12648–12654.
- 7 K. Kamada, K. Otha, T. Kubo, A. Shimizu, Y. Morita, K. Nakasuji, R. Kishi, S. Ohta, S.-I. Furukawa, H. Takahashi and M. Nakano, *Angew. Chem., Int. Ed.*, 2007, **46**, 3544–3546.
- 8 S. Shil, D. Bhattacharya, A. Misra and D. J. Klein, *Phys. Chem. Chem. Phys.*, 2015, **17**, 23378–23383.
- 9 I. Ratera and J. Veciana, *Chem. Soc. Rev.*, 2012, **41**, 303–349.
- 10 Y. Morita, S. Nishida, T. Murata, M. Moriguchi, A. Ueda, M. Satoh, K. Arifuku, K. Sato and T. Takui, *Nat. Mater.*, 2011, **10**, 947–951.
- 11 H. Koike, M. Chikamatsu, R. Azumi, J. Tsutsumi, K. Ogawa, W. Yamane, T. Nishiuchi, T. Kubo, T. Hasegawa and K. Kanai, *Adv. Funct. Mater.*, 2016, **26**, 277–283.
- 12 A. Konishi and T. Kubo, *Benzenoid Quinodimethanes in Physical organic chemistry of quinodimethanes*, ed. Y. Tobe and T. Kubo, Springer, Berlin Heidelberg New York, 2018, pp. 69–80.
- 13 J. Casado, *Top. Curr. Chem.*, 2017, **375**, 73.
- 14 W. Zeng, Z. Sun, T. S. Herng, T. P. Goncalves, T. Y. Gopalakrishna, K. W. Huang, J. Ding and J. Wu, *Angew. Chem., Int. Ed.*, 2016, **55**, 8615–8619.
- 15 S. H. Qiu, C. Q. Wang, S. Xie, X. B. Huang, L. L. Chen, Y. H. Zhao and Z. B. Zeng, *Chem. Commun.*, 2018, **54**, 11383–11386.
- 16 Z. B. Zeng, M. Ishida, J. L. Zafra, X. J. Zhu, Y. M. Sung, N. N. Bao, R. D. Webster, B. S. Lee, R. W. Li, W. D. Zeng, Y. Li, C. Y. Chi, J. T. L. Navarrete, J. Ding, J. Casado, D. Kim and J. Wu, *J. Am. Chem. Soc.*, 2013, **135**, 6363–6371.
- 17 Q. Wang, P. Hu, T. Tanaka, T. Y. Gopalakrishna, T. S. Herng, H. Phan, W. D. Zeng, J. Ding, A. Osuka, C. Y. Chi, J. Siegel and J. Wu, *Chem. Sci.*, 2018, **9**, 5100–5105.
- 18 D. Schmidt, M. Son, J. M. Lim, M. J. Lin, I. Krummenacher, H. Braunschweig, D. Kim and F. Würthner, *Angew. Chem., Int. Ed.*, 2015, **54**, 13980–13984.
- 19 R. Rausch, D. Schmidt, D. Bialas, I. Krummenacher, H. Braunschweig and F. Würthner, *Chem.–Eur. J.*, 2018, **24**, 3420–3424.
- 20 F. Miao, H. Phan and J. Wu, *Molecules*, 2019, **24**, 1446.
- 21 K. Zhang, K. W. Huang, J. L. Li, J. Luo, C. Y. Chi and J. Wu, *Org. Lett.*, 2009, **11**, 4854–4857.
- 22 F. Tampieri, L. Colella, A. Maghsoumi, J. Martí-Rujas, E. Parisini, M. Tommasini, C. Bertarelli and A. Barbon, *J. Phys. Chem. C*, 2016, **120**, 5732–5740.
- 23 G. Xue, X. Hu, H. Chen, L. Ge, W. Wang, J. Xiong, F. Miao and Y. Zheng, *Chem. Commun.*, 2020, **56**, 5143–5146.
- 24 A. Ueda, S. Nishida, K. Fukui, T. Ise, D. Shiomi, K. Sato, T. Takui, K. Nakasuji and Y. Morita, *Angew. Chem., Int. Ed.*, 2010, **49**, 1678–1682.



- 25 S. Lee, F. Miao, H. Phan, T. S. Herng, J. Ding, J. Wu and D. Kim, *ChemPhysChem*, 2017, **18**, 591–595.
- 26 N. Orms and A. I. Krylov, *Phys. Chem. Chem. Phys.*, 2018, **20**, 13127–13144.
- 27 G. E. Rudebusch, J. L. Zafra, K. Jorner, K. Fukuda, J. L. Marshall, I. Arrechea-Marcos, G. L. Espejo, R. P. Ortiz, C. J. Gomez-Garcia, L. N. Zakharov, M. Nakano, H. Ottosson, J. Casado and M. M. Haley, *Nat. Chem.*, 2016, **8**, 753–759.
- 28 R. West, J. A. Jorgenson, K. L. Stearley and J. C. Calabrese, *J. Chem. Soc., Chem. Commun.*, 1991, 1234–1235.
- 29 J. K. Zhou and A. Rieker, *J. Chem. Soc., Perkin Trans. 2*, 1997, 931–938.
- 30 E. V. Canesi, D. Fazzi, L. Colella, C. Bertarelli and C. Castiglioni, *J. Am. Chem. Soc.*, 2012, **134**, 19070–19083.
- 31 K. Takahashi, A. Gunji, K. Yanagi and M. Miki, *J. Org. Chem.*, 1996, **61**, 4784–4792.
- 32 W. Wang, L. Ge, G. Xue, F. Miao, P. Chen, H. Chen, Y. Lin, Y. Ni, J. Xiong, Y. Hu, J. Wu and Y. Zheng, *Chem. Commun.*, 2020, **56**, 1405–1408.
- 33 Y. Li, W. K. Heng, B. S. Lee, N. Aratani, J. L. Zafra, N. Bao, R. Lee, Y. M. Sung, Z. Sun, K. W. Huang, R. D. Webster, J. T. L. Navarrete, D. Kim, A. Osuka, J. Casado, J. Ding and J. Wu, *J. Am. Chem. Soc.*, 2012, **134**, 14913–14922.
- 34 P. Ravat, T. Solomek, D. Haussinger, O. Blacque and M. Juriček, *J. Am. Chem. Soc.*, 2018, **140**, 10839–10847.
- 35 P. Hu, S. Lee, K. H. Park, S. Das, T. S. Herng, T. P. Goncalves, K. W. Huang, J. Ding, D. Kim and J. Wu, *J. Org. Chem.*, 2016, **81**, 2911–2919.
- 36 W. D. Zeng, T. Y. Gopalakrishna, H. Phan, T. Tanaka, T. S. Herng, J. Ding, A. Osuka and J. Wu, *J. Am. Chem. Soc.*, 2018, **140**, 14054–14058.
- 37 J. Z. Liu, P. Ravat, M. Wagner, M. Baumgarten, X. L. Feng and K. Müllen, *Angew. Chem., Int. Ed.*, 2015, **54**, 12442–12446.
- 38 E. Clar, *The Aromatic Sextet*, Wiley, London, 1972.
- 39 J. J. Dressler, Z. Zhou, J. L. Marshall, R. Kishi, S. Takamuku, Z. Wei, S. N. Spisak, M. Nakano, M. A. Petrukhina and M. M. Haley, *Angew. Chem., Int. Ed.*, 2017, **56**, 15363–15367.
- 40 J. Ma, J. Z. Liu, M. Baumgarten, Y. B. Fu, Y. Z. Tan, K. S. Schellhammer, F. Ortmann, G. Cuniberti, H. Komber, R. Berger, K. Müllen and X. L. Feng, *Angew. Chem., Int. Ed.*, 2017, **56**, 3280–3284.
- 41 C. K. Frederickson, B. D. Rose and M. M. Haley, *Acc. Chem. Res.*, 2017, **50**, 977–987.
- 42 R. Nieman, N. J. Silva, A. J. A. Aquino, M. M. Haley and H. Lischka, *J. Org. Chem.*, 2020, **85**, 3664–3675.
- 43 J. J. Dressler, A. C. Valdivia, R. Kishi, G. E. Rudebusch, A. M. Ventura, B. E. Chastain, C. J. Gómez-García, L. N. Zakharov, M. Nakano, J. Casado and M. M. Haley, *Chem*, 2020, **6**, 1353–1368.
- 44 P. Hu and J. Wu, *Can. J. Chem.*, 2017, **95**, 223–233.
- 45 M. Grzybowski and D. T. Gryko, *Adv. Opt. Mater.*, 2015, **3**, 280–320.
- 46 M. Kaur and D. H. Choi, *Chem. Soc. Rev.*, 2015, **44**, 58–77.
- 47 C. Zhao, Y. Guo, Y. Zhang, N. Yan, S. You and W. Li, *J. Mater. Chem. A*, 2019, **7**, 10174–10199.
- 48 M. Pieczykolan, B. Sadowski and D. T. Gryko, *Angew. Chem., Int. Ed.*, 2020, **59**, 7528–7535.
- 49 H. Bürckstümmer, A. Weissenstein, D. Bialas and F. Würthner, *J. Org. Chem.*, 2011, **76**, 2426–2432.
- 50 H. Zollinger, *Colour Chemistry*, ed. P. M. Wallimann, Wiley-VCH, Weinheim, 3rd edn, 2003, pp. 342–344.
- 51 P. Data, A. Kurowska, S. Pluczyk, P. Zassowski, P. Pander, R. Jedrysiak, M. Czwartosz, L. Otulakowski, J. Suwinski, M. Lapkowski and A. P. Monkman, *J. Phys. Chem. C*, 2016, **120**, 2070–2078.
- 52 X. Gao and Y. Hu, *J. Mater. Chem. C*, 2014, **2**, 3099–3117.
- 53 Y. Zhang, C. Kim, J. Lin and T.-Q. Nguyen, *Adv. Funct. Mater.*, 2012, **22**, 97–105.
- 54 T. He, P. Leowanawat, C. Burschka, V. Stepanenko, M. Stolte and F. Würthner, *Adv. Mater.*, 2018, **30**, 1804032.
- 55 W. W. Li, K. H. Hendriks, M. M. Wienk and R. A. J. Janssen, *Acc. Chem. Res.*, 2016, **49**, 78–85.
- 56 V. Cuesta, M. Vartanian, P. Malhotra, S. Biswas, P. de la Cruz, G. D. Sharma and F. Langa, *J. Mater. Chem. A*, 2019, **7**, 11886–11894.
- 57 W. Yue, S.-L. Suraru, D. Bialas, M. Müller and F. Würthner, *Angew. Chem., Int. Ed.*, 2014, **53**, 6159–6162.
- 58 M. Grzybowski, E. Glodkowska-Mrowka, T. Stoklosa and D. T. Gryko, *Org. Lett.*, 2012, **14**, 2670–2673.
- 59 G. M. Fischer, E. Daltrozzi and A. Zumbusch, *Angew. Chem., Int. Ed.*, 2011, **50**, 1406–1409.
- 60 K. Takiyama and K. Kawabata, *J. Synth. Org. Chem., Jpn.*, 2018, **76**, 1176–1184.
- 61 S. Ray, S. Sharma, U. Salzner and S. Patil, *J. Phys. Chem. C*, 2017, **121**, 16088–16097.
- 62 Y. L. Qiao, Y. L. Guo, C. M. Yu, F. J. Zhang, W. Xu, Y. Q. Liu and D. B. Zhu, *J. Am. Chem. Soc.*, 2012, **134**, 4084–4087.
- 63 C. Wang, Y. P. Zang, Y. K. Qin, Q. Zhang, Y. H. Sun, C. A. Di, W. Xu and D. B. Zhu, *Chem.-Eur. J.*, 2014, **20**, 13755–13761.
- 64 C. Wang, Y. K. Qin, Y. H. Sun, Y. S. Guan, W. Xu and D. B. Zhu, *ACS Appl. Mater. Interfaces*, 2015, **7**, 15978–15987.
- 65 H. Z. Wu, Y. Wang, X. L. Qiao, D. L. Wang, X. D. Yang and H. X. Li, *Chem. Mater.*, 2018, **30**, 6992–6997.
- 66 L. B. Ren, F. Liu, X. X. Shen, C. Zhang, Y. P. Yi and X. Zhu, *J. Am. Chem. Soc.*, 2015, **137**, 11294–11302.
- 67 Parts of this work were presented at the 18th International Symposium on Novel Aromatic Compounds (ISNA-18) in July 2019 in Sapporo, Japan.
- 68 J. Mizuguchi, A. Grubenmann, G. Wooden and G. Rihs, *Acta Crystallogr., Sect. B: Struct. Sci.*, 1992, **48**, 696–700.
- 69 S. Celik, Y. Ergun and S. Alp, *J. Fluoresc.*, 2009, **19**, 829–835.
- 70 M. Murai, S. Y. Ku, N. D. Treat, M. J. Robb, M. L. Chabinyk and C. J. Hawker, *Chem. Sci.*, 2014, **5**, 3753–3760.
- 71 Y. Kumar, S. Kumar, D. Bansal and P. Mukhopadhyay, *Org. Lett.*, 2019, **21**, 2185–2188.
- 72 Y. Shao, A. I. Krylov and M. Head-Gordon, *J. Chem. Phys.*, 2003, **118**, 4807–4817.
- 73 J. C. Slater, *Quantum Theory of Molecules and Solids*, McGraw-Hill, New York, 1974.
- 74 A. D. Becke, *Phys. Rev. A: At., Mol., Opt. Phys.*, 1988, **38**, 3098–3100.



- 75 S. H. Vosko, L. Wilk and M. Nusair, *Can. J. Phys.*, 1980, **58**, 1200–1211.
- 76 C. Lee, W. Yang and R. G. Parr, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1988, **37**, 785–789.
- 77 F. Weigend and R. Ahlrichs, *Phys. Chem. Chem. Phys.*, 2005, **7**, 3297–3305.
- 78 D. Doehnert and J. Koutecky, *J. Am. Chem. Soc.*, 1980, **102**, 1789–1796.
- 79 B. O. Roos, *Adv. Chem. Phys.*, 1987, **69**, 399–445.
- 80 H.-J. Werner and P. J. Knowles, *J. Chem. Phys.*, 1985, **82**, 5053–5062.
- 81 T. Busch, A. D. Esposti and H. Werner, *J. Chem. Phys.*, 1991, **94**, 6708–6715.
- 82 P. J. Knowles and H.-J. Werner, *Chem. Phys. Lett.*, 1985, **115**, 259–267.
- 83 J. Kruszewski and T. M. Krygowski, *Tetrahedron Lett.*, 1972, **13**, 3839–3842.
- 84 T. M. Krygowski, *J. Chem. Inf. Comput. Sci.*, 1993, **33**, 70–78.

