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## N-type and p-type molecular doping on monolayer MoS<sub>2</sub><sup>†</sup>

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Monolayer MoS<sub>2</sub> has attracted much attention due to its high on/off current ratio, transparency, and suitability for optoelectronic devices. Surface doping by molecular adsorption has proven to be an effective method to facilitate the usage of MoS<sub>2</sub>. However, there are no works available to systematically clarify the effects of the adsorption of F<sub>4</sub>TCNQ, PTCDA, and tetracene on the electronic and optical properties of the material. Therefore, this work elucidated the problem by using density functional theory calculations. We found that the adsorption of F<sub>4</sub>TCNQ and PTCDA turns MoS<sub>2</sub> into a p-type semiconductor, while the tetracene converts MoS<sub>2</sub> into an n-type semiconductor. The occurrence of a new energy level in the conduction band for F<sub>4</sub>TCNQ and PTCDA and the valence band for tetracene reduces the bandgap of the monolayer MoS<sub>2</sub>. Besides, the MoS<sub>2</sub>/F<sub>4</sub>TCNQ and MoS<sub>2</sub>/PTCDA systems exhibit an auxiliary optical peak at the long wavelengths of 950 and 850 nm, respectively. Contrastingly, the MoS<sub>2</sub>/tetracene modifies the optical spectrum of the monolayer MoS<sub>2</sub> only in the ultraviolet region. The findings are in good agreement with the experiments.

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### 1. Introduction

The monolayer MoS<sub>2</sub> is a semiconducting transition metal dichalcogenide with an optical bandgap of about 1.8 eV. This material has proven to possess high on/off current ratio and transparency, making it suitable for electronic and optoelectronic applications.<sup>1,2</sup> The properties of monolayer MoS<sub>2</sub> are adjustable by modifying the crystal structure, doping with different elements, and the adsorption of organic molecules.<sup>3–6</sup> The application of organic molecular adsorption is attractive because of advantages such as mechanical flexibility, easy fabrication, light weight, and low cost.<sup>7,8</sup> Organic molecular adsorption has been demonstrated as a viable method to optimize material properties,<sup>6</sup> also to protect MoS<sub>2</sub> from oxidation under ambient conditions.<sup>9</sup> Therefore, it is essential to understand the behavior of the organic adsorption on the monolayer MoS<sub>2</sub> to optimize its performance for optoelectronic uses.

Here, we studied the adsorption of three organic molecules, *i.e.*, 2,3,5,6-tetrafluoro-7,7,8,8-tetracyanoquinodimethane

(F<sub>4</sub>TCNQ), perylene-3,4,9,10-tetracarboxylic dianhydride (PTCDA), and tetracene. Where, F<sub>4</sub>TCNQ is one of the most effective p-type dopants due to its strong ability of electron acceptance.<sup>10–13</sup> The experiment showed that F<sub>4</sub>TCNQ could modify the electrical and optical properties of MoS<sub>2</sub>.<sup>11</sup> Hu *et al.* confirmed that the deposition of F<sub>4</sub>TCNQ on the monolayer MoS<sub>2</sub> allows convenient control of the radioactive exciton recombination for laser and LED applications.<sup>12</sup> The charge transfer within the F<sub>4</sub>TCNQ–MoS<sub>2</sub> interface can tune electrical and gas sensing properties.<sup>13</sup> PTCDA and tetracene are popular organic semiconductors that have attracted many research interests in the past decades.<sup>8,14–19</sup> Tetracene was used for field-effect transistors.<sup>20,21</sup> The study of Habib showed the outstanding luminescent properties of the MoS<sub>2</sub>/PTCDA system.<sup>17</sup> Also, Wang *et al.* have demonstrated the novel MoS<sub>2</sub>/PTCDA hybrid heterojunction synapse transistor with both electrical and optical modulation and efficient gate tunability.<sup>18</sup> Besides, the MoS<sub>2</sub>/tetracene is a good candidate for the ambipolar field-effect transistor.<sup>21</sup>

The literature review showed that most of the available research is experiments, and there is no theoretical study available to elucidate the interaction between the MoS<sub>2</sub> monolayer and the mentioned organic molecules. Therefore, this work is going to clarify it by using the density functional theory (DFT) method and analyzing the electronic structure and optical properties of MoS<sub>2</sub> without and with the adsorption of F<sub>4</sub>TCNQ, PTCDA, and tetracene. Furthermore, this work also considers the influences of pressure on the properties.

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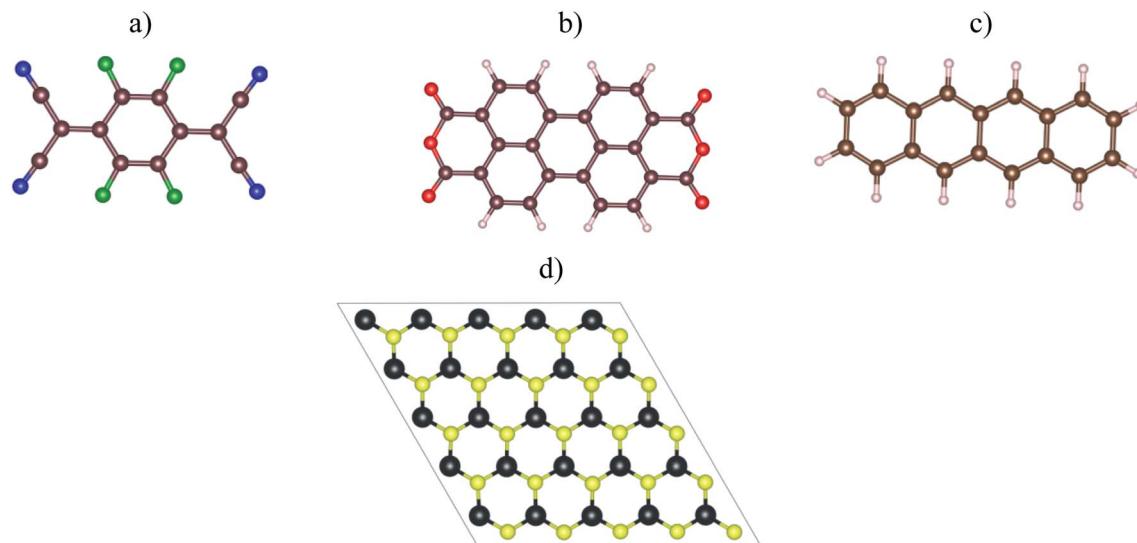


Fig. 1 The top view of organic molecules: (a)  $\text{F}_4\text{TCNQ}$ , (b) PTCDA, (c) tetracene, and (d) the monolayer  $\text{MoS}_2$ . Blue (N), green (F), brown (C), red (O), white (H), yellow (S), and black (Mo).

## 2. Computational method

Fig. 1a–c visualized the molecular structure of  $\text{F}_4\text{TCNQ}$  ( $\text{C}_{12}\text{F}_4\text{N}_4$ ), PTCDA ( $\text{C}_{24}\text{H}_8\text{O}_6$ ), and tetracene ( $\text{C}_{18}\text{H}_{12}$ ), respectively. The monolayer  $\text{MoS}_2$ , shown in Fig. 1d, was described by a slab model with the  $5 \times 5$  unit cell and a vacuum space of  $27 \text{ \AA}$  along the normal direction of the  $\text{MoS}_2$  surface.

The optimized geometry structure and total energy of the monolayer  $\text{MoS}_2$  with organic molecular adsorption were obtained with the aid of the Vienna *ab initio* simulation package (VASP).<sup>22–24</sup> The projector-augmented-wave method was used to describe the valence electron–core ion interaction.<sup>25,26</sup> The

electron exchange–correlation interactions were treated by using the PBE-GGA approximation of the Perdew–Burke–Ernzerhof,<sup>27,28</sup> which has shown to be a reasonable approximation for the bandgap of the monolayer  $\text{MoS}_2$ .<sup>29</sup> The plane-wave cutoff energy was 600 eV. The  $k$ -point mesh was sampled at  $3 \times 3 \times 1$  in the first Brillouin zone by Monkhorst–Pack technique.<sup>30</sup> All atomic positions of  $\text{MoS}_2$  and  $\text{MoS}_2$ /organic-molecule systems are fully relaxed during the geometry optimization until the interatomic forces meeting the upper criterion of  $0.001 \text{ eV \AA}^{-1}$ .

The pressure applying to the system is defined by

$$P = \frac{E - E_0}{V - V_0}, \quad (1)$$

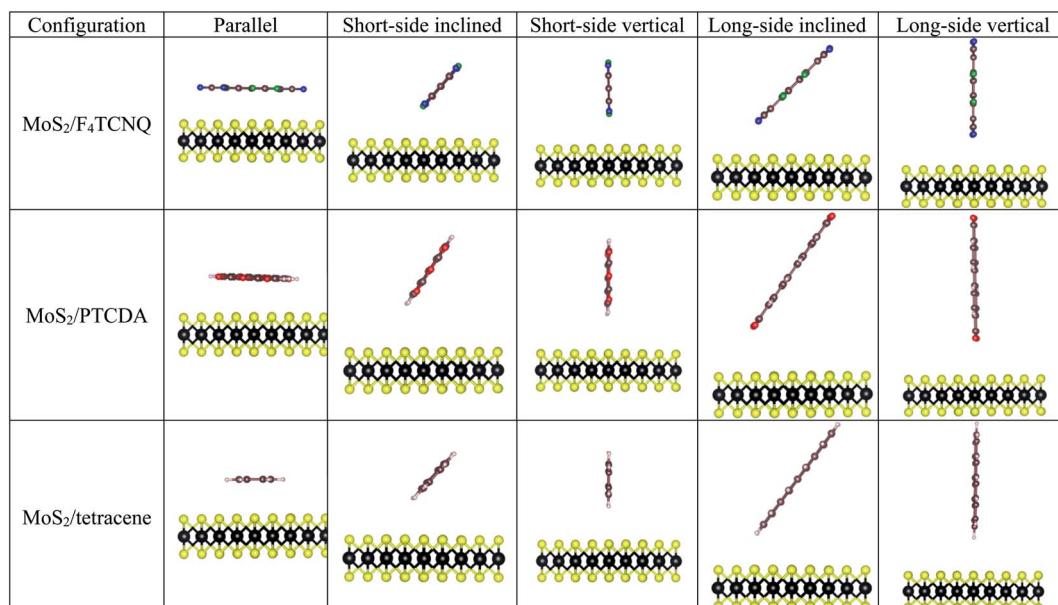


Fig. 2 The side view of the possible adsorption configurations of  $\text{F}_4\text{TCNQ}$ , PTCDA, and tetracene on the monolayer  $\text{MoS}_2$ .



**Table 1** Adsorption energy  $E_a$  (eV) for different configurations on the monolayer  $\text{MoS}_2$  without strain

Adsorption configuration	$\text{F}_4\text{TCNQ}$	PTCDA	Tetracene
Parallel	−0.009	0.019	0.039
Short-side inclined	<b>−0.055</b>	<b>−0.071</b>	<b>−0.046</b>
Short-side vertical	−0.038	−0.036	−0.030
Long-side inclined	−0.027	−0.031	−0.022
Long-side vertical	−0.023	−0.17	0.137

**Table 2** Adsorption energy  $E_a$  and the shortest distance  $d$  from the organic molecule to the substrate surface

$a^a$ (Å)	MoS <sub>2</sub> /F <sub>4</sub> TCNQ		MoS <sub>2</sub> /PTCDA		MoS <sub>2</sub> /tetracene	
	$E_a$ (eV)	$d$ (Å)	$E_a$ (eV)	$d$ (Å)	$E_a$ (eV)	$d$ (Å)
2.84	0.034	3.44	0.095	3.48	0.105	3.49
3.00	0.066	3.41	0.098	3.46	0.092	3.41
3.18	−0.009	3.47	0.019	3.53	0.039	3.49
3.32	−0.090	3.58	−0.019	3.60	−0.018	3.58
3.48	−0.130	3.60	−0.074	3.64	−0.128	3.61
3.64	−0.145	3.77	−0.080	3.79	−0.230	3.72

<sup>a</sup> 2.84 and 3.00 Å: for compressive pressure; 3.18 Å: without pressure; 3.32, 3.48, and 3.64 Å: for tensile strain.

where, the energy and volume of the strained (unstrained) system are denoted as  $E$  and  $V$  ( $E_0$  and  $V_0$ ), respectively. The volume of the unit cell of the monolayer  $\text{MoS}_2$  is  $V = \frac{\sqrt{3}}{2}a^2c$  ( $a$  is the lattice constant and  $c$  is the thickness of the slab). The pressure is modified by changing the lattice constant  $a = b$ .

The adsorption energy of  $\text{F}_4\text{TCNQ}$ , PTCDA, and tetracene on the monolayer  $\text{MoS}_2$  is calculated by

$$E_a = E_{\text{sub+om}} - (E_{\text{sub}} + E_{\text{om}}), \quad (2)$$

where,  $E_{\text{sub+om}}$ ,  $E_{\text{sub}}$ , and  $E_{\text{om}}$  are the total energy of the  $\text{MoS}_2$ /organic molecule system, the isolated monolayer  $\text{MoS}_2$ , and the isolated organic molecule, respectively.

The optical property of the systems was determined through the imaginary part  $\epsilon(\omega)$  of dielectric constant, which is calculated by the summation over empty states as follows:<sup>31,32</sup>

$$\epsilon(\omega) = \frac{2\pi e^2}{V\varepsilon_0} \sum_{\mathbf{k},\mathbf{v},\mathbf{c}} \left| \langle \xi_{\mathbf{k}}^{\mathbf{c}} | \hat{\mathbf{v}} \cdot \mathbf{r} | \xi_{\mathbf{k}}^{\mathbf{v}} \rangle \right|^2 \delta[\hbar\omega - (E_{\mathbf{k}}^{\mathbf{v}} - E_{\mathbf{k}}^{\mathbf{c}})], \quad (3)$$

here,  $\omega$  is the frequency of emission,  $\varepsilon_0$  is the dielectric constant of free space. The valence and conduction band wave functions are  $\xi_{\mathbf{k}}^{\mathbf{v}}$  and  $\xi_{\mathbf{k}}^{\mathbf{c}}$ , respectively. The unit cell volume is  $V$ , and the polarization vector of the electric field of the emission is  $\hat{\mathbf{v}}$ .

### 3. Results and discussion

#### 3.1. Structural properties

The optimized structural parameters for the monolayer  $\text{MoS}_2$  were obtained in the previous publication,<sup>33</sup> where the optimized lattice constant was found to be  $a = b = 3.18 \text{ \AA}$ . To explore the adsorption properties of  $\text{F}_4\text{TCNQ}$ , PTCDA, and tetracene on the monolayer  $\text{MoS}_2$ , the geometry of the organic molecules with different configurations was optimized at several adsorption sites on the surface of  $\text{MoS}_2$  as shown in Fig. 2. The most favorable adsorption of the  $\text{F}_4\text{TCNQ}$ , PTCDA, and tetracene molecules on the  $\text{MoS}_2$  substrate was obtained for the short-side inclined configuration based on the most negative adsorption energy listed in Table 1.

Although the adsorption energy of the short-side inclined and parallel configurations are different, the electronic density of states (DOS) and the dielectric function are almost the same, as shown in Fig. S1 and S2 in the ESI.† Therefore, to make a comparison with the previous publication on polyethylenimine,<sup>33</sup> we will perform a detailed analysis of the geometrical structure, the DOS, and the dielectric function for the parallel configuration of  $\text{F}_4\text{TCNQ}$ , PTCDA, and tetracene. Furthermore, we also found that the extension of the unit cell size to  $6 \times 6$  and the inclusion of van der Waals correction (using vdW-DF functional) generated more negative adsorption energy (Tables S1 and S2†). However, the electronic structure of the  $\text{MoS}_2/\text{F}_4\text{TCNQ}$ ,  $\text{MoS}_2/\text{PTCDA}$ , and  $\text{MoS}_2/\text{tetracene}$  systems remain the same as that of the unit cell  $5 \times 5$  without van der Waals correction, as shown in Fig. S3.† Thus, it is acceptable to



Fig. 4 The index of atoms for the bond lengths and angles.

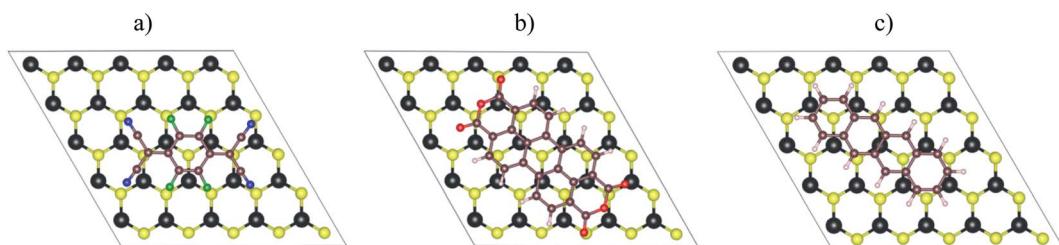
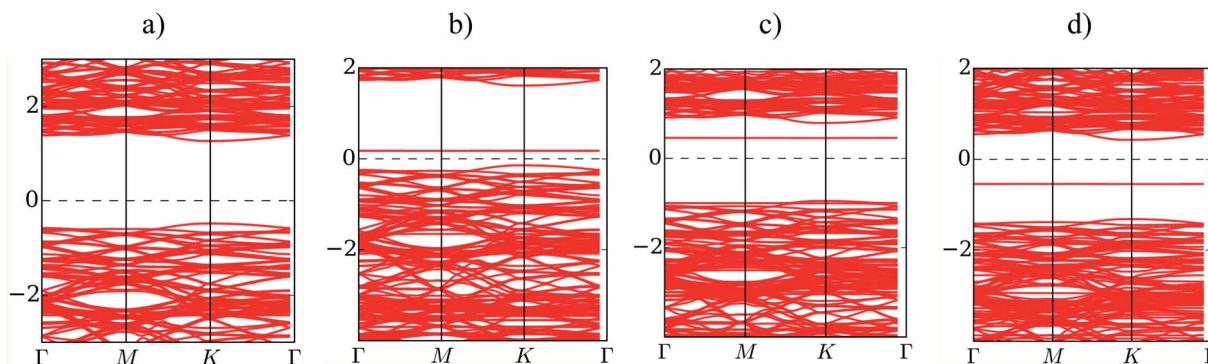


Fig. 3 The top view of the parallel adsorption configuration of the organic molecules on the monolayer  $\text{MoS}_2$ : (a)  $\text{F}_4\text{TCNQ}$ , (b) PTCDA, and (c) tetracene.



**Table 3** Structural parameters of the MoS<sub>2</sub> substrate before and after the adsorption of the organic molecules. The index of atoms is described in Fig. 4

Parameters	MoS <sub>2</sub>	MoS <sub>2</sub> /F <sub>4</sub> TCNQ	MoS <sub>2</sub> /PTCDA	MoS <sub>2</sub> /tetracene
Bonds (Å)	S <sub>1</sub> –S <sub>2</sub>	3.183	3.160	3.162
	S <sub>1</sub> –S <sub>3</sub>	3.126	3.144	3.146
	Mo <sub>1</sub> –Mo <sub>2</sub>	3.187	3.183	3.180
	Mo <sub>2</sub> –S <sub>1</sub>	2.415	2.415	2.415
Angles (°)	Mo <sub>1</sub> S <sub>1</sub> Mo <sub>2</sub>	82.30	82.46	82.34
	S <sub>1</sub> Mo <sub>2</sub> S <sub>2</sub>	82.30	81.76	81.88
	S <sub>1</sub> Mo <sub>1</sub> S <sub>3</sub>	80.54	80.90	81.02



**Fig. 5** The electronic band structure of (a) the monolayer MoS<sub>2</sub>, (b) MoS<sub>2</sub>/F<sub>4</sub>TCNQ, (c) MoS<sub>2</sub>/PTCDA, and (d) MoS<sub>2</sub>/tetracene at zero pressure.

select the unit cell  $5 \times 5$  and the PBE-GGA approximation in the present study.

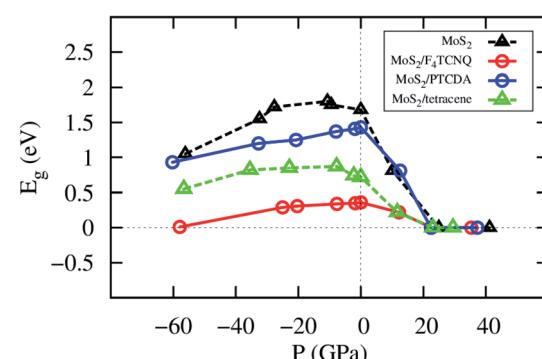
Table 2 listed the adsorption energy and the vertical nearest distance to the MoS<sub>2</sub> surface of the F<sub>4</sub>TCNQ, PTCDA, and tetracene molecules in the parallel configuration (Fig. 3). With the application of the tensile strain to the system, the organic molecules become stabilized with more negative adsorption energy relative to that at the lattice constant of 3.18 Å (without strain). The compressive pressure makes the molecules less stable with more positive adsorption energy. Table 2 also shows that the vertical nearest distance to the surface increases with the increase of the lattice constant. This result seems contradicted to the tendency of the more stable configuration of the adsorption upon modifying the tensile strain. Although the vertical nearest distance increases, the electrostatically attractive force between the adsorbates and MoS<sub>2</sub> increases at a faster rate.<sup>33</sup> Therefore, the adsorption energy becomes more negative. Besides, the interaction of the molecules with the surface is physisorption.

The geometrical parameters such as the bond length and bond angle of the MoS<sub>2</sub> substrate before and after the adsorption of the organic molecules are exhibited in Table 3. We found that the structural parameters did not change significantly after the adsorption of the molecules, which is consistent with the physical adsorption found in this work.

### 3.2. Electronic properties

Fig. 5 presents the band structure of MoS<sub>2</sub> before and after the adsorption of F<sub>4</sub>TCNQ, PTCDA, and tetracene at zero pressure.

The bandgap is determined by the difference between the conduction band minimum and the valence band maximum. We find that the isolated substrate shows a direct bandgap of 1.68 eV at the K point, as listed in Table 4. The calculated bandgap of 1.68 eV comes from the electronic contribution only, while the experimental value of 1.80 eV has included the exciton binding energy.<sup>2,3,34</sup> Upon the adsorption of the organic molecules, the new energy state occurs in the bandgap region of the MoS<sub>2</sub> substrate. The new energy level occurring above the Fermi level for the adsorption of F<sub>4</sub>TCNQ and PTCDA implies that the substrate became the p-type semiconductor. Whilst, it occurring below the Fermi level for the adsorption of tetracene indicates that the monolayer MoS<sub>2</sub> became the n-type semiconductor. The existence of the new energy level shrinks the bandgap of the MoS<sub>2</sub> to 0.36, 1.43, and 0.72 eV for the



**Fig. 6** The bandgap versus pressure of the monolayer MoS<sub>2</sub> with and without the adsorption of the organic molecules.



Table 4 The bandgap  $E_g$  of the monolayer MoS<sub>2</sub> with and without the adsorption of the organic molecules

$a$ (Å)	MoS <sub>2</sub> <sup>a</sup>		MoS <sub>2</sub> /F <sub>4</sub> TCNQ		MoS <sub>2</sub> /PTCDA		MoS <sub>2</sub> /tetracene	
	$P$ (GPa)	$E_g$ (eV)	$P$ (GPa)	$E_g$ (eV)	$P$ (GPa)	$E_g$ (eV)	$P$ (GPa)	$E_g$ (eV)
2.84	-56.18	1.05	-57.93	0.01	-60.27	0.93	-56.61	0.55
3.00	-32.44	1.55	-25.03	0.29	-32.70	1.20	-35.46	0.82
3.04	-27.65	1.72	-20.31	0.31	-20.74	1.25	-22.88	0.85
3.12	-10.69	1.80	-7.64	0.34	-7.93	1.37	-7.78	0.87
3.16	-9.46	1.75	-1.79	0.35	-1.91	1.41	-2.24	0.74
3.18	0	1.68	0	0.36	0	1.43	0	0.72
3.32	10.13	0.81	12.21	0.22	12.57	0.81	11.68	0.22
3.48	24.95	0	22.59	0	22.41	0	23.03	0
3.64	41.14	0	35.36	0	37.26	0	29.44	0

<sup>a</sup> In ref. 33.

adsorption of F<sub>4</sub>TCNQ, PTCDA, and tetracene, respectively. However, the bandgap of the MoS<sub>2</sub> remains the direct bandgap nature after the adsorption of the organic molecules.

The change of the lattice constant modifies the pressure applied to the system. The details of the bandgap and the pressure are described in Table 4 and visualized in Fig. 6. We found that the bandgap behavior of MoS<sub>2</sub>/tetracene is similar to

that of the monolayer MoS<sub>2</sub>, which exhibits a concave-downward curve with the maximum bandgap of 0.87 and 1.80 eV found at the pressure of -7.78 and -10.69 GPa for the MoS<sub>2</sub>/tetracene and monolayer MoS<sub>2</sub>, respectively. The concave-downward curve for the bandgap *versus* the pressure for the MoS<sub>2</sub> system was found to be in good agreement with experiment,<sup>35</sup> while that for the MoS<sub>2</sub>/tetracene is similar to that of an n-type semiconductor, the MoS<sub>2</sub>/PEI interface, in the previous study.<sup>33,36</sup> There is a little difference in the behavior of the bandgap *versus* the pressure for MoS<sub>2</sub>/F<sub>4</sub>TCNQ and MoS<sub>2</sub>/PTCDA compared to that of the monolayer MoS<sub>2</sub>. That is the bandgap somehow linearly increases to a maximum value and then decreases with more positive pressures. The maximum value is 0.36 and 1.43 eV at the zero pressure for MoS<sub>2</sub>/F<sub>4</sub>TCNQ and MoS<sub>2</sub>/PTCDA, respectively. We have to impress that MoS<sub>2</sub>/F<sub>4</sub>TCNQ and MoS<sub>2</sub>/PTCDA are the p-type semiconductors, which are different from the n-type of the MoS<sub>2</sub>/tetracene. The monolayer MoS<sub>2</sub>, MoS<sub>2</sub>/F<sub>4</sub>TCNQ, MoS<sub>2</sub>/PTCDA, and MoS<sub>2</sub>/tetracene were found to be completely converted to the metal with the bandgap of 0 eV at the pressure of 24.95, 22.59, 22.41, and 23.03 GPa, respectively.

In a complex structure, the Bader point charge of an atom was determined in two steps: (1) calculating the charge by the Bader partition technique, and (2) subtracting the charge obtained in step (1) to that of the neutral atom. In the pseudo-potential method of the DFT, one considers the contribution of valence electrons only. Table 5 presents the minus-plus signs as the loss and gain of charge, respectively. We found that F<sub>4</sub>TCNQ and PTCDA accept while MoS<sub>2</sub> donates the charge. Contrastingly, the tetracene donates while MoS<sub>2</sub> gains the charge. These results are in a good correlation with the p-type doping nature of F<sub>4</sub>TCNQ and PTCDA and the n-type doping of tetracene on the MoS<sub>2</sub> monolayer.

The electronic density of states (DOS) in Fig. 7a-c shows that the adsorption of F<sub>4</sub>TCNQ, PTCDA, and tetracene contributes to the appearance of the new state coming from the p<sub>z</sub> orbital of the organic molecules. As shown in Fig. 7d-f, the charge clouds of the atoms of the organic molecules clearly show the shape of p<sub>z</sub> orbitals along the vertical direction, the vector **c**. The charge accumulation was found to more dominant on the F<sub>4</sub>TCNQ and

Table 5 The point charge (in the unit of e) of the MoS<sub>2</sub>/F<sub>4</sub>TCNQ, MoS<sub>2</sub>/PTCDA, and MoS<sub>2</sub>/tetracene systems

MoS <sub>2</sub> /F <sub>4</sub> TCNQ		
$P$ (GPa)	MoS <sub>2</sub>	F <sub>4</sub> TCNQ
-57.93	-0.226	+0.226
-25.03	-0.098	+0.098
0	-0.073	+0.073
12.21	-0.144	+0.144
22.59	-0.231	+0.231
MoS <sub>2</sub> /PTCDA		
$P$ (GPa)	MoS <sub>2</sub>	PTCDA
23.03	-0.015	+0.015
29.44	-0.017	+0.017
-20.74	-0.019	+0.019
0	-0.020	+0.020
12.57	-0.021	+0.021
22.41	-0.025	+0.025
MoS <sub>2</sub> /tetracene		
$P$ (GPa)	MoS <sub>2</sub>	Tetracene
-56.61	+0.062	-0.062
-35.46	+0.060	-0.060
-22.88	+0.052	-0.052
-7.78	+0.045	-0.045
0	+0.048	-0.048
11.68	+0.050	-0.050
23.03	+0.103	-0.103
29.44	+0.363	-0.363



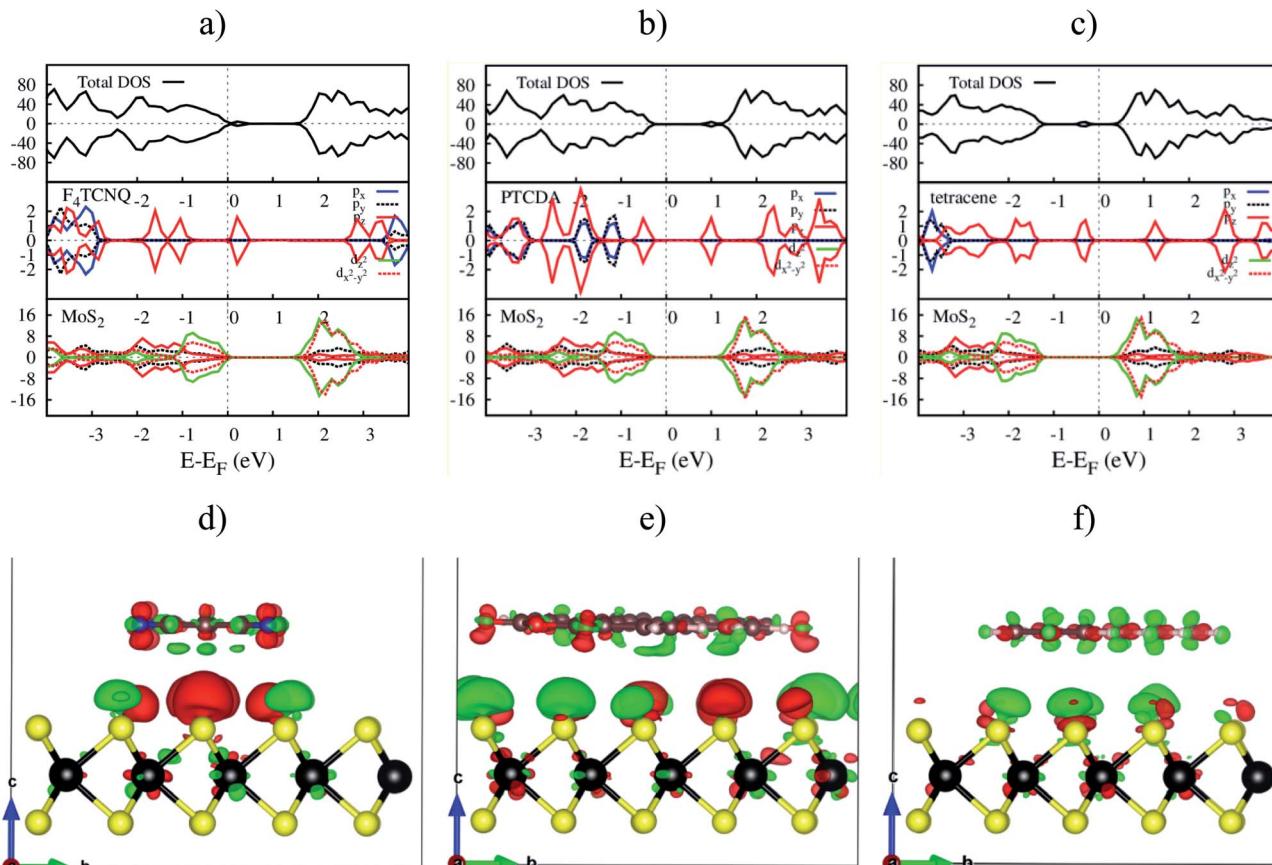


Fig. 7 The upper panel is the total and orbital-projected DOS. The bottom panel is the charge density difference at zero pressure. From left to right is for  $\text{MoS}_2/\text{F}_4\text{TCNQ}$ ,  $\text{MoS}_2/\text{PTCDA}$ , and  $\text{MoS}_2/\text{tetracene}$ , respectively. Occupied and unoccupied states are presented in red and green in that order.

PTCDA, while the charge donation distributes more on tetracene. This result is in agreement with the above analysis about the Bader charge.

Fig. 8 shows the response of the total electronic density of states toward the influence of pressure. The p-type doping  $\text{F}_4\text{TCNQ}$  and PTCDA shift the valence band maximum of  $\text{MoS}_2$

upward, while the n-type doping tetracene moves the conduction band minimum of  $\text{MoS}_2$  downward near to the Fermi level. The variation of the position of the valence band maximum and conduction band minimum leads to modifying the bandgap of the systems.

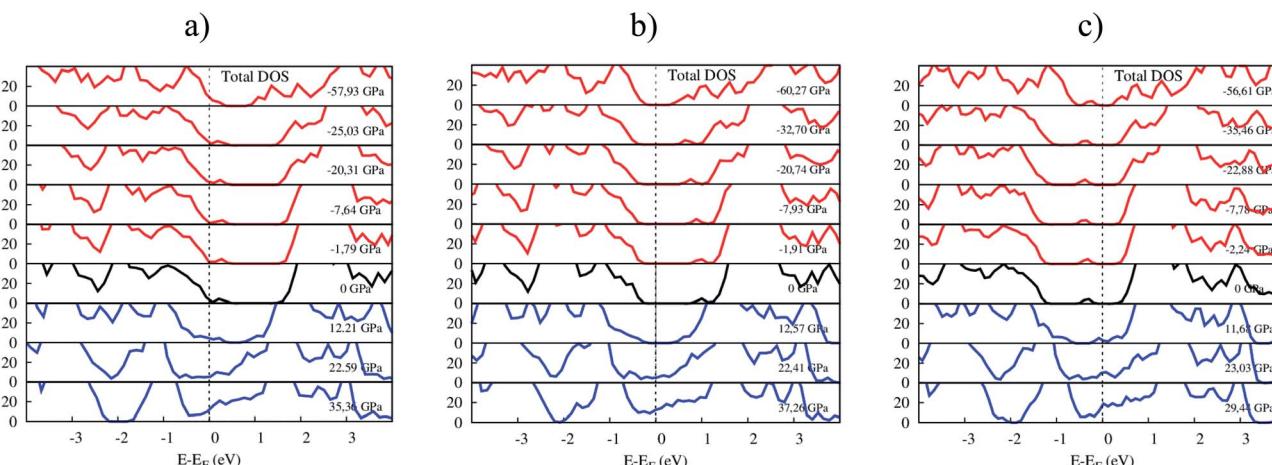


Fig. 8 The total DOS of (a)  $\text{MoS}_2/\text{F}_4\text{TCNQ}$ , (b)  $\text{MoS}_2/\text{PTCDA}$ , and (c)  $\text{MoS}_2/\text{tetracene}$  under pressure.



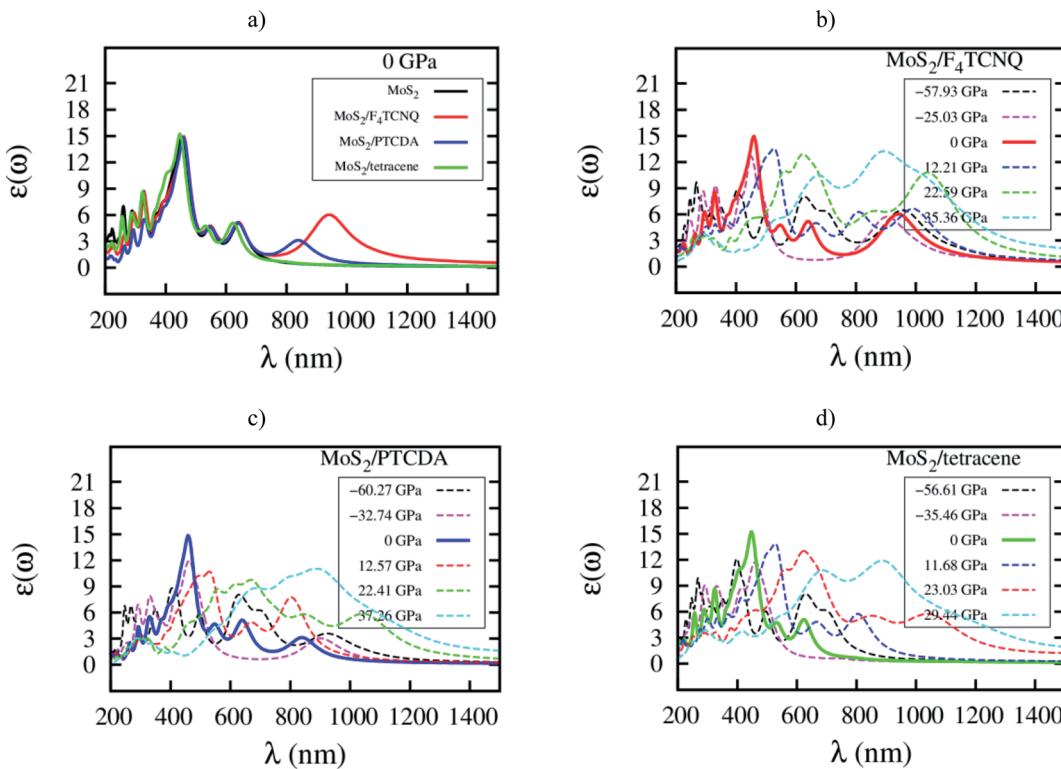


Fig. 9 Imaginary part of dielectric function  $\varepsilon(\omega)$  versus wavelength  $\lambda$  for (a) the systems at zero pressure, (b)  $\text{MoS}_2/\text{F}_4\text{TCNQ}$ , (c)  $\text{MoS}_2/\text{PTCDA}$ , and (d)  $\text{MoS}_2/\text{tetracene}$  under pressure.

### 3.3. Optical properties

Here, we elucidate the influences of the organic molecular adsorption on the optical properties of the monolayer  $\text{MoS}_2$ . The imaginary part of the frequency-dependent dielectric function of the  $\text{MoS}_2$ ,  $\text{MoS}_2/\text{F}_4\text{TCNQ}$ ,  $\text{MoS}_2/\text{PTCDA}$ , and  $\text{MoS}_2/\text{tetracene}$  in the independent-particle approximation including local field effects is presented in Fig. 9. After that, optical parameters as listed in Table 6 can be obtained in two steps: (1)

determining the wavelength  $\lambda$  and the intensity  $\varepsilon(\omega)$  at the maximum peak of the dielectric function, check Fig. 9, (2) calculating the corresponding photon energy  $E_p = \hbar c/\lambda$ . At zero pressure, we found in Fig. 9a that the optical spectrum of the  $\text{MoS}_2/\text{F}_4\text{TCNQ}$ ,  $\text{MoS}_2/\text{PTCDA}$ , and  $\text{MoS}_2/\text{tetracene}$  systems has the maximum peak at almost the same position as that of the monolayer  $\text{MoS}_2$ , which is at about 450 nm. The result indicates that these systems can emit and absorb blue light.<sup>29,37</sup> The optical spectrum of the  $\text{MoS}_2/\text{tetracene}$  has almost the same

Table 6 The wavelength and dielectric function at the maximum intensity of optical spectrum

$\text{MoS}_2$ (ref. 33)	$P$ (GPa)	<b>-56.18</b>	<b>-32.44</b>	<b>0</b>	<b>10.13</b>	<b>24.95</b>	<b>41.14</b>
	$E_p$ (eV) <sup>a</sup>	4.66	2.60	2.73	2.39	2.00	1.42
	$\lambda$ (nm)	266	477	453	521	621	873
	$\varepsilon(\omega)$	11.06	11.58	14.88	13.36	13.18	12.06
$\text{MoS}_2/\text{F}_4\text{TCNQ}$	$P$ (GPa)	<b>-57.93</b>	<b>-25.03</b>	<b>0</b>	<b>12.21</b>	<b>22.59</b>	<b>35.36</b>
	$E_p$ (eV) <sup>a</sup>	4.62	2.75	2.69	2.36	1.99	1.38
	$\lambda$ (nm)	268	451	461	526	624	896
	$\varepsilon(\omega)$	9.65	12.66	14.95	13.54	12.92	13.22
$\text{MoS}_2/\text{PTCDA}$	$P$ (GPa)	<b>-60.27</b>	<b>-32.74</b>	<b>0</b>	<b>12.57</b>	<b>22.41</b>	<b>37.26</b>
	$E_p$ (eV) <sup>a</sup>	3.06	2.67	2.71	2.34	1.86	1.39
	$\lambda$ (nm)	405	463	458	529	668	889
	$\varepsilon(\omega)$	8.82	11.95	14.82	10.72	9.76	11.05
$\text{MoS}_2/\text{tetracene}$	$P$ (GPa)	<b>-56.61</b>	<b>-35.46</b>	<b>0</b>	<b>11.68</b>	<b>23.03</b>	<b>29.44</b>
	$E_p$ (eV) <sup>a</sup>	3.11	2.69	2.77	2.36	1.99	1.40
	$\lambda$ (nm)	398	461	447	526	623	886
	$\varepsilon(\omega)$	12.11	11.72	15.21	13.85	13.03	11.90

<sup>a</sup>  $E_p = \hbar c/\lambda$  is the photon energy,  $\lambda$  is the photon wavelength,  $\hbar$  is Planck constant, and  $c$  is the speed of light in vacuum.



behavior as the monolayer MoS<sub>2</sub> with a little modification at short wavelengths around 300 nm. Contrastingly, the MoS<sub>2</sub>/F<sub>4</sub>TCNQ and MoS<sub>2</sub>/PTCDA exhibits an additional peak at 950 and 850 nm, respectively. It is necessary to remind that the MoS<sub>2</sub>/F<sub>4</sub>TCNQ and MoS<sub>2</sub>/PTCDA are p-type, while the MoS<sub>2</sub>/tetracene is an n-type semiconductor. The behavior of the optical spectrum of the MoS<sub>2</sub>/tetracene was found to be similar to that of the n-type doping of MoS<sub>2</sub> with PEI.<sup>33</sup> Therefore, the auxiliary peak at a longer wavelength was only found for the monolayer MoS<sub>2</sub> with the p-type doping.<sup>5</sup> Furthermore, as shown in Fig. 9b-d, the intensity of the main optical peak at non-zero pressures was suppressed compared to that at 0 GPa.

## 4. Conclusion

This work investigated the effects of the organic molecular adsorption on the physical properties of the monolayer MoS<sub>2</sub> by the DFT calculations. The results have proved that the adsorption of F<sub>4</sub>TCNQ, PTCDA, and tetracene on MoS<sub>2</sub> is physisorption. The organic molecular adsorption reduces the direct bandgap of the monolayer MoS<sub>2</sub> to 0.36, 1.43, and 0.72 eV, respectively. The main cause is due to the emergence of the new energy level at the conduction band minimum for MoS<sub>2</sub>/F<sub>4</sub>TCNQ and MoS<sub>2</sub>/PTCDA, and at the valence band maximum for MoS<sub>2</sub>/tetracene, by the p<sub>z</sub> orbitals of the adsorbates. Under pressure, the bandgap of the MoS<sub>2</sub>/F<sub>4</sub>TCNQ, MoS<sub>2</sub>/PTCDA, and MoS<sub>2</sub>/tetracene systems reaches the maximum value of 0.36, 1.43, and 0.87 eV, respectively. Whilst, it becomes 0 eV implying the transition to the metal at the tensile pressure of 22.59, 22.41, and 23.03 GPa in that order. The adsorption of tetracene changes the optical peak structure of the monolayer MoS<sub>2</sub> at the ultraviolet region around 300 nm; however, the F<sub>4</sub>TCNQ and PTCDA adsorption generate an auxiliary peak at the long wavelengths.

## Conflicts of interest

There are no conflicts of interest to declare.

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