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# Tin-selenide as a futuristic material: properties and applications

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SnSe/SnSe<sub>2</sub> is a promising versatile material with applications in various fields like solar cells, photodetectors, memory devices, lithium and sodium-ion batteries, gas sensing, photocatalysis, supercapacitors, topological insulators, resistive switching devices due to its optimal band gap. In this review, all possible applications of SnSe/SnSe<sub>2</sub> have been summarized. Some of the basic properties, as well as synthesis techniques have also been outlined. This review will help the researcher to understand the properties and possible applications of tin selenide-based materials. Thus, this will help in advancing the field of tin selenide-based materials for next generation technology.

## 1. Introduction

The fast developing area of applied material science demands materials to be cheap, non-toxic, environment-friendly, easy to synthesize, and well competitive in performing particular applications. Nowadays, materials with versatility have gained massive attention due to their applicability in almost all fields.<sup>1,2</sup> Various materials have been explored and showed promising versatile applications, *e.g.*, graphene (Gr),<sup>1</sup> TiO<sub>2</sub>,<sup>3</sup> ZnO,<sup>4</sup> Cu<sub>2</sub>SnS<sub>3</sub> (CTS),<sup>2</sup> *etc.* The multifunctional applicability of these materials paves the foundation for interdisciplinary research. Chalcogenide-based materials have also shown such potential and can be seen as the future hope to meet a similar requirement. Among chalcogenides, tin selenide has demonstrated great potential in the applied material science. Tin selenide exists in two phases, *i.e.*, SnSe and SnSe<sub>2</sub>. Some researchers have observed another phase, Sn<sub>2</sub>Se<sub>3</sub>,<sup>5,6</sup> but this phase is the superposition of SnSe and SnSe<sub>2</sub>.<sup>7</sup> Tin selenide has demonstrated versatility in thermoelectric,<sup>8</sup> photodetector,<sup>9</sup> solar cells,<sup>10</sup> photocatalytic,<sup>11</sup> phase change memory,<sup>12</sup> gas sensing,<sup>13</sup> anode material for battery,<sup>14</sup> supercapacitor,<sup>15</sup> and topological insulator (TI).<sup>16</sup> These applications strongly depend upon the properties of SnSe (optical, electrical and micro-structural, *etc.*). Apart from material properties, material synthesis/deposition methods also play an essential role in obtaining high-quality materials.<sup>17–21</sup> Excellent review article

exists on thermoelectric materials consisting of fundamental properties to the thermoelectric device's final design, growth, defects, working environment issues, and applications.<sup>22</sup> Other reviews that focusses on SnSe describes all the aspects mentioned above (like growth, defects, configuration, *etc.*).<sup>23</sup> The aim of this review is to summarize the ongoing progress on SnSe, SnSe<sub>2</sub> synthesis methods, materials properties, and its possible application in various fields. However, some studies focus on pristine tin-selenide phase and its application<sup>8,24</sup> Another good reviews<sup>8,25,26</sup> that concentrates on SnSe describing all aspects mentioned above (like growth, defects, design, *etc.*), provides all-round knowledge to researchers. This review gives insights into the phases, structures, synthesis methods, progress in the tin-selenide, and tin diselenide for various applications. Hence, this article aims to review the tin selenide devices, present status, recent progress in the growth process, related-issues, possible solutions, and their possibility of developing next-generation technology.

## 2. Material properties of tin selenide

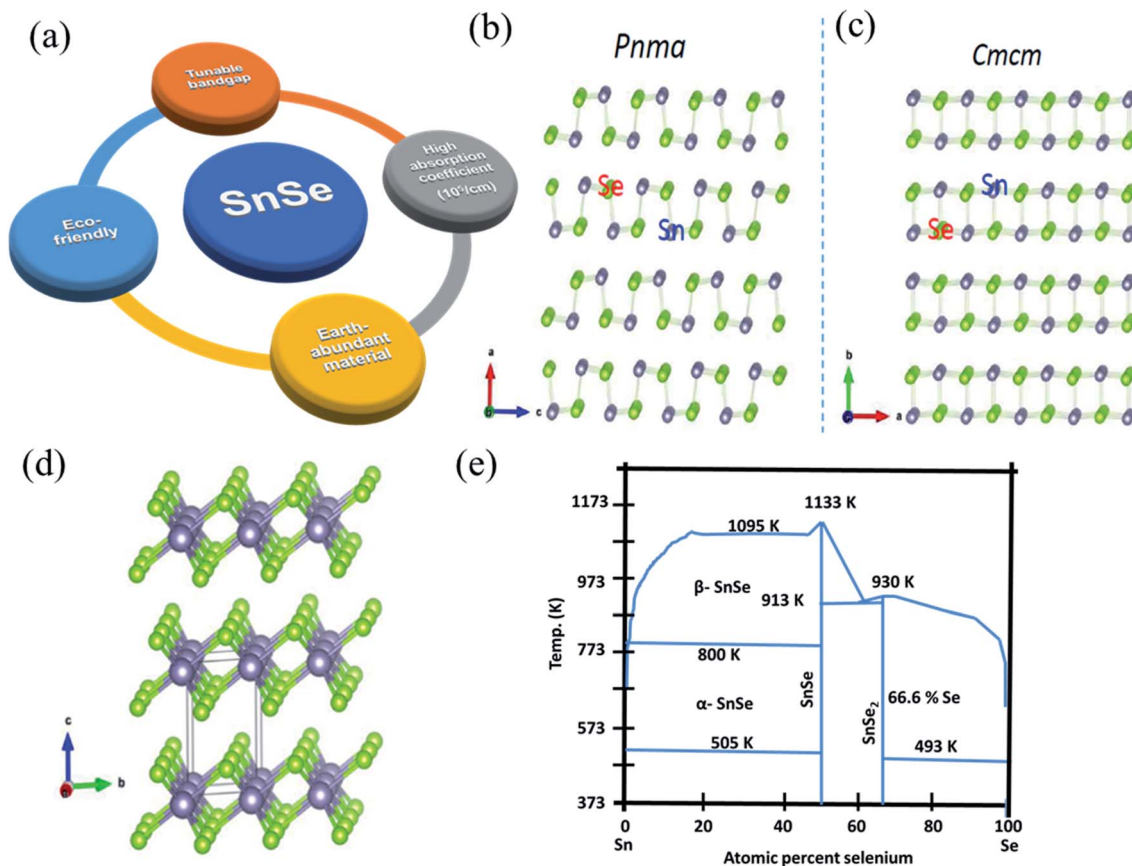
Tin selenide based SnSe/SnSe<sub>2</sub> are binary compound semiconductor materials having p/n-type conductivity.<sup>20,27</sup> Fig. 1(a) shows a salient feature of tin selenide materials. The SnSe exists in two crystallographic phases, *i.e.*, orthorhombic ( $\alpha$ -SnSe)<sup>28</sup> and cubic ( $\pi$ -SnSe),<sup>29</sup> and SnSe<sub>2</sub><sup>30</sup> exists in the hexagonal crystal structure. The  $\pi$ -SnSe material is mechanically stable and energetically comparable to  $\alpha$ -SnSe and has already shown potential in the piezoelectric application.<sup>29</sup> The orthorhombic crystal structure of SnSe showed second-order displacive type phase transition above 750 K, from  $\alpha$ -SnSe (space group *Pnma*) to  $\beta$ -SnSe (space group *Cmcm*) with ( $a = 11.49$  Å,  $b = 4.44$  Å,  $c = 4.135$  Å) to the ( $a = 4.31$  Å,  $b = 11.70$  Å,  $c = 4.31$  Å), respectively (Fig. 1(b)).<sup>31,32</sup> The crystal structure of SnSe (violet, Sn atoms;

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**Fig. 1** (a) Salient feature of tin selenide materials. (b) Crystal structure of SnSe, (c) SnSe having *Pnma* to *Cmcm* phase transition. This figure has been adapted/reproduced from ref. 38 with permission from Elsevier, copyright 2018". (d) Crystal structure of SnSe<sub>2</sub>. This figure has been adapted/reproduced from ref. 33 with permission from ACS, copyright 2016". (e) Equilibrium phase diagram of Sn–Se system. This figure has been adapted/reproduced from ref. 39 with permission from Wiley, copyright 2020".

green Se atoms) (Fig. 1(c)),<sup>32</sup> and crystal structures of SnSe<sub>2</sub> (green, Se atoms; violet, Sn atoms) is shown in Fig. 1(d).<sup>33</sup> Both the tin-selenide phases show the indirect and direct band gaps. The direct band gap of  $\alpha$ -SnSe shows a wide tunable band gap which varies from 0.98 eV (bulk) to 1.43 eV (monolayer).<sup>34</sup> SnSe<sub>2</sub> shows large variation in the band gap from 1.84 eV (bulk) to 2.04 eV (monolayer).<sup>35</sup> This band gap tunability of tin selenide shows its immense application possibilities in optoelectronic device applications.<sup>30,36</sup> Tin selenide exists in two stoichiometric phases, *i.e.*, SnSe and SnSe<sub>2</sub>, as shown by the equilibrium phase diagram (Fig. 1(e)).<sup>7</sup> Various research groups have reported another phase, Sn<sub>2</sub>Se<sub>3</sub>,<sup>37</sup> but this phase is the superposition of the SnSe and SnSe<sub>2</sub> as confirmed by Nuclear Magnetic Resonance (NMR) spectroscopy.

Defects present in the material also influences the properties of materials like electronic, magnetic, and optical properties. SnSe is generally a p-type semiconductor. Density functional theory (DFT) calculation is performed to study the defects during SnSe crystal growth. The Sn vacancy is present as a native defect, which causes the p-type conduction in the SnSe under Se or Sn rich conditions, as confirmed by Scanning Tunneling Microscopy (STM) studies.<sup>40,41</sup> SnSe<sub>2</sub> shows n-type conduction. DFT calculation shows that the vacancy of selenium and interstitial tin led to this n-type conduction in SnSe<sub>2</sub>.<sup>41</sup> The

optoelectronic and physical properties of SnSe/SnSe<sub>2</sub> materials depend on the growth conditions, size, morphology, phase purity, growth techniques, *etc.* Therefore, selection of synthesis/deposition method impacts material as well as its proposed applications.

### 3. Synthesis of SnSe and SnSe<sub>2</sub>

The material's properties also depend on the growth conditions,<sup>42</sup> size, and morphology of the material,<sup>43</sup> phase purity,<sup>43</sup> defects,<sup>44</sup> *etc.* Various physical and chemical methods have been established to fabricate SnSe/SnSe<sub>2</sub> materials on different substrates (glass, flexible metal foil, polymer, *etc.*) for multiple applications. The tin selenide based semiconductor materials have been synthesized using atomic layer deposition (ALD),<sup>17</sup> sputtering,<sup>45</sup> thermal evaporation,<sup>46</sup> hydrothermal,<sup>47</sup> spray pyrolysis,<sup>48</sup> chemical vapor deposition (CVD),<sup>49</sup> *etc.*

Bulk crystals of the SnSe and SnSe<sub>2</sub> can be easily obtained by the physical solid-state reaction method. >99.99 pure elemental powders of the Sn and Se (in the stoichiometric ratio) is taken and sealed in the quartz tube at the pressure of  $\sim 10^{-4}$  torr. It is placed at the temperature according to the Sn–Se system phase diagram, as shown in Fig. 2(a). The Bridgman method can be used to obtain single crystals. The powder was melted into the



furnace and brought in the contacts with the single crystal's seed during its cooling. The material grows along the crystallographic orientation of the seed crystal. Depending upon the geometry it is called the horizontal Bridgman method or vertical Bridgman method. The Bridgman–Stockbarger technique as shown in (Fig. 2(a)) can be used to obtain a single crystal of SnSe

and SnSe<sub>2</sub> with an optimized ampule lowering rate of 7 mm h<sup>-1</sup> and 4 mm h<sup>-1</sup>, respectively, for SnSe and SnSe<sub>2</sub> with an optimized temperature gradient of 15 °C cm<sup>-1</sup>.<sup>50</sup> Another technique to grow single crystals are the direct vapor transport method,<sup>42</sup> and temperature gradient method.<sup>51</sup> A detailed review of the growth of the SnSe single crystals can be found in an earlier study.<sup>25</sup> Various

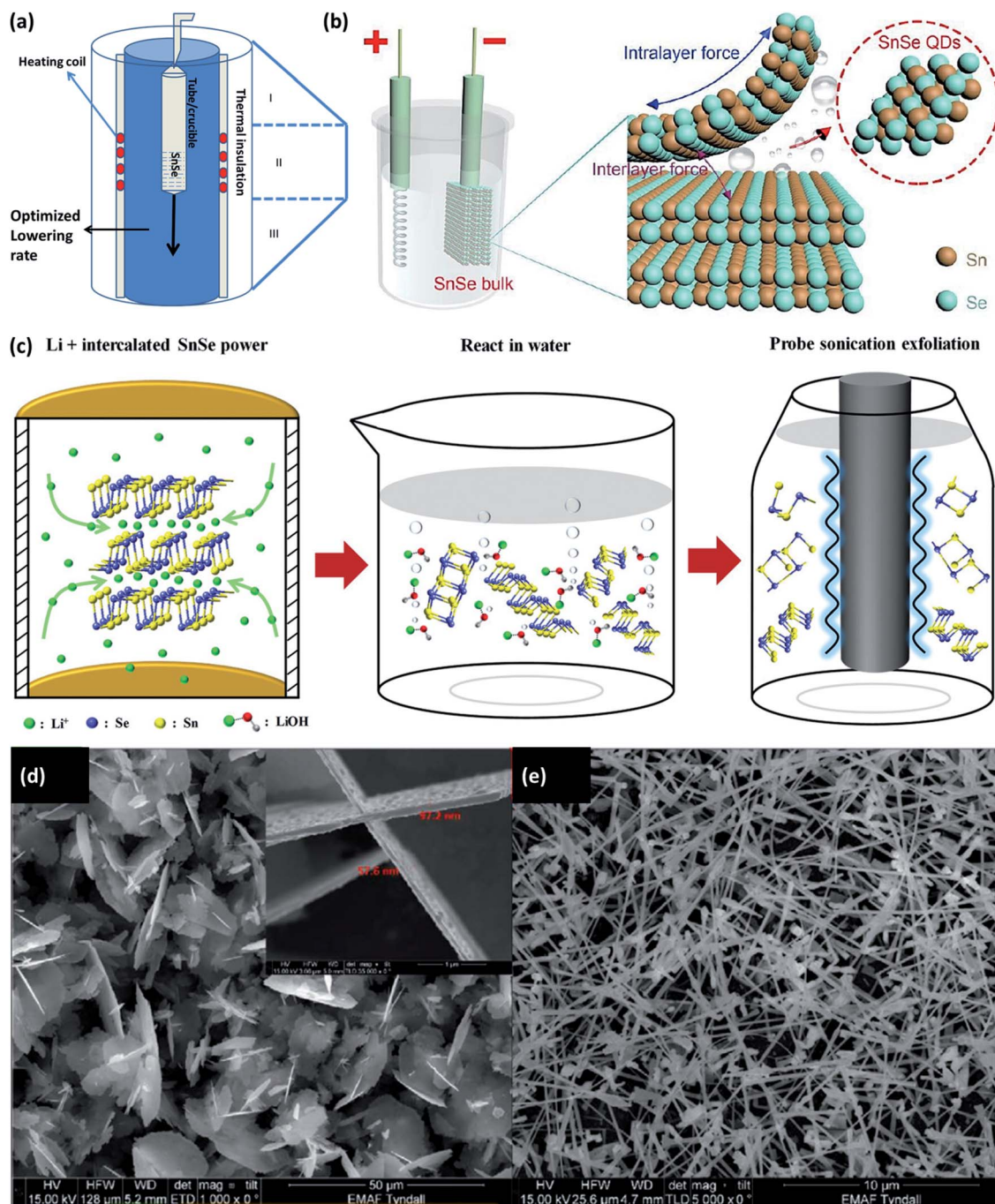


Fig. 2 (a) Bridgman–Stockbarger technique to grow a single crystal. (b) Schematic illustration for the synthesis of SnSe QDs by cathodic exfoliation. This figure has been adapted/reproduced from ref. 52 with permission from RSC, copyright 2019<sup>52</sup>. (c) Schematic synthesis process for the SnSe nanosheets by the liquid-phase exfoliation method includes two main steps: Li<sup>+</sup> hydrothermal intercalation and sonication-assisted exfoliation. This figure has been adapted/reproduced from ref. 54 with permission from Wiley, copyright 2020<sup>54</sup>. SEM images showing the evolution of the SnSe<sub>x</sub> nanostructures concerning the growth temperature, (d) large flakes with diameters of several microns grown at 450 °C, with a thickness of ~60 nm shown in the inset, (e) high yield of nanowire growth at 500 °C. This figure has been adapted/reproduced from ref. 49 with permission from Wiley, copyright 2020<sup>49</sup>.



methods to grow the nanostructure of the SnSe and SnSe<sub>2</sub> is described in this section.

### 3.1 Synthesis of nanostructured SnSe

Li *et al.*<sup>52</sup> devised the facile cathodic exfoliation method to exfoliate SnSe bulk into quantum dots with high yield. Under the organic electrolyte containing 0.2 M of tetrabutylammonium (at  $-7.5$  V for 30 min) bulk SnSe was exfoliated and downsized up to  $\sim 10$  nm in lateral size, which was further downsized to  $\sim 4$  nm by breaking weakly entangled dots with the help of sonication as shown in Fig. 2(b). To synthesize the SnSe flake up to one layer, Jiang *et al.*<sup>53</sup> reported the two-step process, in which first the bulk SnSe flakes were obtained by the atmospheric pressure vapor transport deposition method. Flash evaporation of SnSe powder was done by moving the quartz tube position containing SnSe powder at  $700$  °C under Ar/H<sub>2</sub> gas pressure under atmospheric condition. SnSe is deposited on the Si substrate kept in upside-down position. As obtained flakes were etched with the N<sub>2</sub> gas. Etching time of 5–20 min gave the single layer SnSe flake (thickness of  $6.8$  Å).<sup>53</sup> The SnSe nano-sheet of the best quality can be produced from the SnSe bulk powder by the three-step method.<sup>54</sup> The method comprises the intercalation of Li<sup>+</sup> ion into the SnSe layers by lithification and then reacting the material with water. The rapid expansion of H<sub>2</sub> gas during reaction with water exfoliated the SnSe nano-sheets. To improve the nanosheets' yield, further sonification of the SnSe powders' residuals is carried out (shown in Fig. 2(c)).

Using the precursor diselenoether SnCl<sub>4</sub>[BuSe(CH<sub>2</sub>)<sub>3</sub>-Se<sup>n</sup>SeBu]}, SnSe<sub>x</sub> nanostructures are formed on the Si(100) substrate coated with Au (acted as a catalyst) by the liquid injection chemical vapor deposition method, and growth was studied in the

temperature range  $450$ – $550$  °C. Under the  $1.1$  sccm flow of Ar gas and  $1.5$  mL h<sup>-1</sup> of precursor injection rate, SnSe<sub>2</sub> flake growth occurred at  $450$  °C, and considerable growth of SnSe nanowires occurred at  $500$  °C as shown in Fig. 2(d) and (e).<sup>49</sup>

### 3.2 Synthesis of nanostructured SnSe<sub>2</sub>

Choi *et al.*<sup>55</sup> synthesized the nano-plates of the SnSe<sub>2</sub> by dissolving 5 mL oleylamine in 50 mg, 0.26 mmol SnCl<sub>2</sub> at  $220$  °C.

After that, a solution containing 1,3-dimethylimidazoline-2-selenone (90 mg, 0.53 mmol), dichloromethane (4 mL), and oleylamine (2 mL) were mixed and heated for 2 h at this temperature. After that, it was cooled to room temperature, centrifuged, and the residue was washed using hexane to obtain SnSe<sub>2</sub> nanoplates. The transmission electron microscopy (TEM) image of SnSe<sub>2</sub> nanoplates are shown in Fig. 3(a)–(d).

A simple two-step method, *i.e.*, sonication followed by laser ablation, was used by Li *et al.* to obtain SnSe<sub>2</sub> quantum dots (QD) from bulk powder.<sup>56</sup> Manually grinded SnSe<sub>2</sub> powder (20 mg) was dispersed in 30 mL deionized water, and the solution was sonicated (650 W) for 2 h in an ice bath for 4 s. Sonication resulted in smaller particles and flakes of SnSe<sub>2</sub>. The obtained solution kept in quartz cuvette was irradiated with a 1064 nm laser (Nd:YAG, 2.2 W) for 10 min. Irradiated tiny particles were centrifuged for 30 min at 6000 rpm, resulting in SnSe<sub>2</sub> QD, as shown in Fig. 4(a) and (b).

Ma *et al.*<sup>58</sup> synthesized the nanorods and nanoplates of the SnSe<sub>2</sub> by organic solution phase route. Two different solutions of 0.4 mmol SnCl<sub>2</sub>·2H<sub>2</sub>O, 2 mmol oleic acid, 8.5 mL liquid paraffin oil, and 0.4 mmol selenium powder, 9 mL liquid paraffin were used. The first solution was heated at  $160$  °C till lemon yellow color appeared, and the later solution was heated

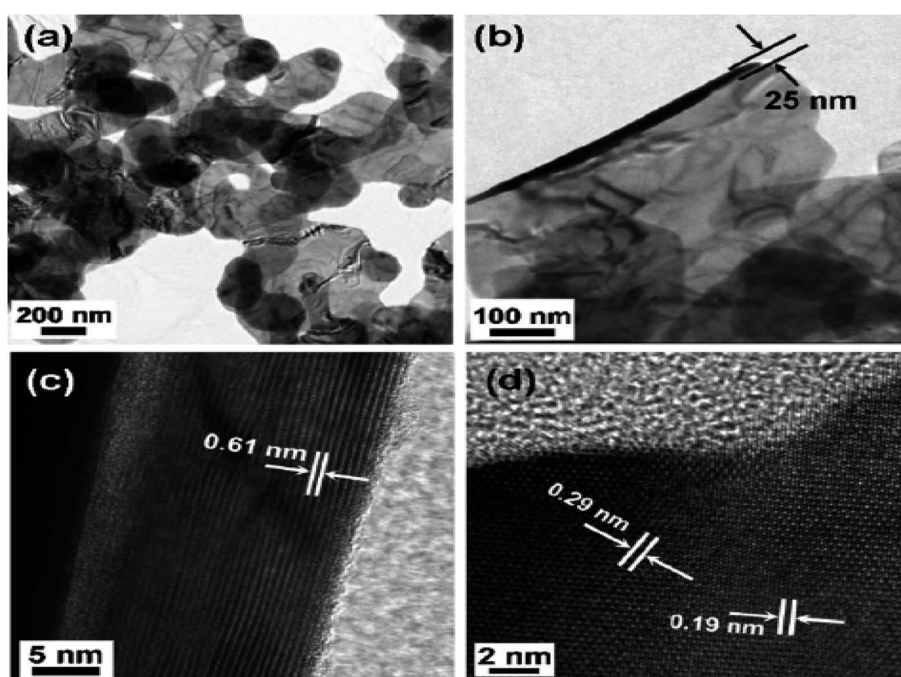


Fig. 3 TEM image (a) and (b) SnSe<sub>2</sub> nanoplates and high resolution TEM (HRTEM) images (c) side, and (d) top views. This figure has been adapted/reproduced from ref. 55 with permission from RSC, copyright 2011\*.



at 240 °C till wine color appeared. The latter solution was injected into first using syringe rapidly and the mixture was heated at 200 °C for 20 min. After cooling, toluene and methanol was added. After centrifugation nanorods and nanoplates were obtained. Chemical vapor deposition was employed by Zhou *et al.* to synthesize ultrathin (1.5 nm) SnSe<sub>2</sub> flakes of high quality.<sup>57</sup> Low melting point precursor SnI<sub>2</sub> (0.01 g) and Se (0.1 g) powders were placed in two different alumina crucible, and mica substrate were placed ~15 cm from the central zone. The central zone was heated at 600 °C for 15 min, with a flow of H<sub>2</sub> (5 sccm) and Ar (20 sccm) at ambient pressure. As-synthesized flakes are shown in Fig. 4(c)–(g).

## 4. Various applications

In recent times, tin selenide has gained significant research attention for solar cells, thermoelectric, photodetectors, photocatalytic, phase change memory, gas sensing, and anode material for battery, supercapacitor, topological insulator *etc.* Various reported applications based on SnSe materials are displayed in Fig. 5.

### 4.1 Thin film solar cells (TFSCs)

Fabrication of low-cost thin-film solar cells is highly needed as they can also be deposited over flexible substrates. The CuInGaSe<sub>2</sub> (CIGSe) and CdTe based solar cells have shown record small cell power conversion efficiency (PCE) of 23.35%

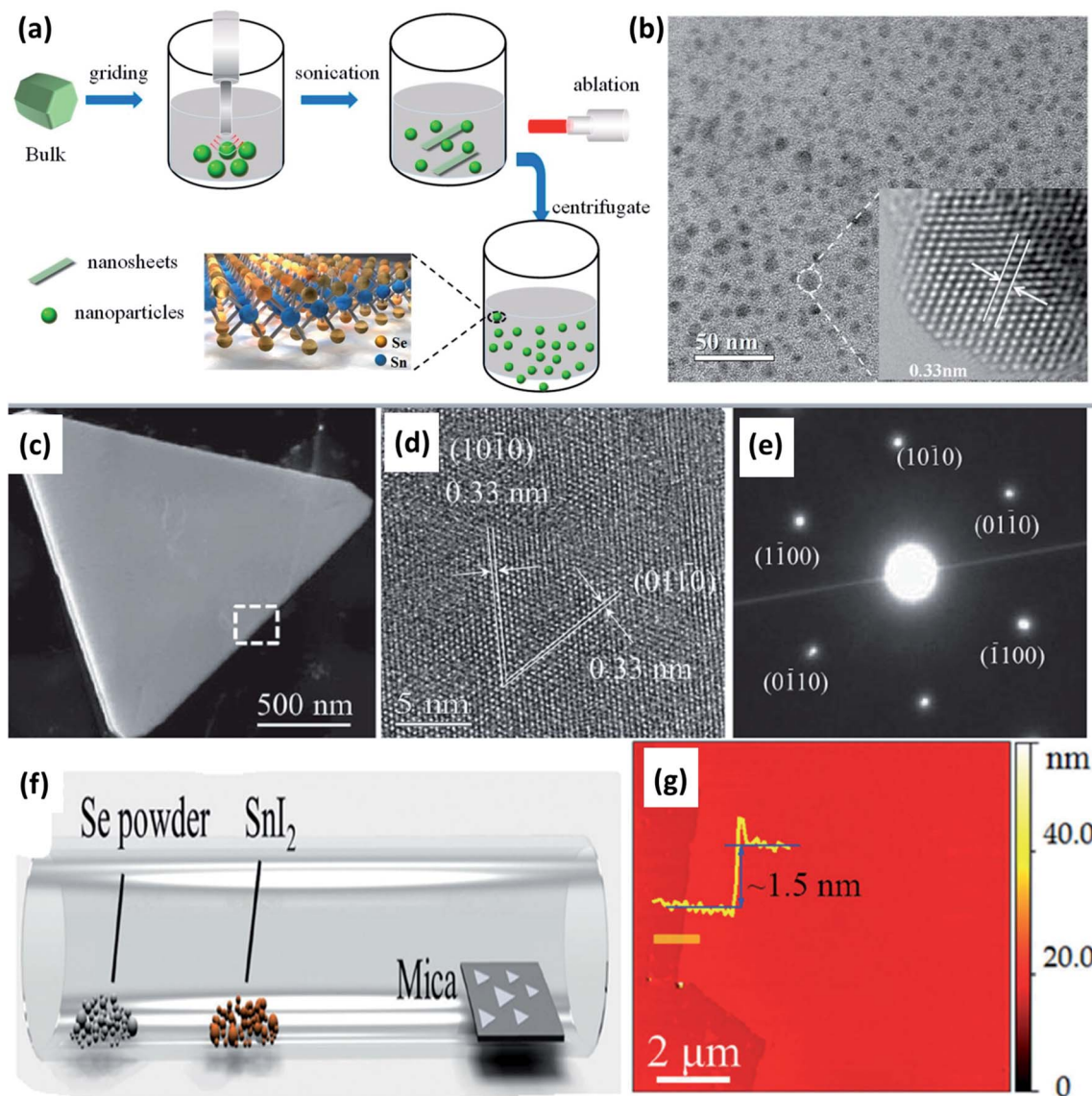


Fig. 4 (a) Schematic show of the SnSe<sub>2</sub> structure and the quantum dot (QD) fabrication process, and (b) TEM image of SnSe<sub>2</sub> QDs with a centrifugal speed of 6000 rpm. These figures has been adapted/reproduced from ref. 56 with permission from MDPI, copyright 2019". (c) Low-magnification TEM image of a SnSe<sub>2</sub> flake, (d) corresponding HRTEM image of the flake, (e) electron diffraction pattern from the same flake, (f) schematic diagram of the chemical vapor deposition method, and (g) a typical atomic force microscope (AFM) image at the flake edge, and the height profile showing a thickness of ~1.5 nm. These figures has been adapted/reproduced from ref. 57 with permission from Wiley, copyright 2015".



and 22.1%, respectively.<sup>66</sup> But in CIGS TFSC, indium is rare, and gallium is very costly. In CdTe TFSC, Cd is toxic, and Te is a rare material. These materials related problems hinder the commercialization of low-cost, earth-abundant TFSC based on CIGS and CdTe. In recent time, earth-abundant and low-cost compound semiconductor materials like  $\text{Cu}_2\text{ZnSnS}_4$  (CZTS),  $\text{Cu}_2\text{ZnSnSe}_4$  (CZTSe), and  $\text{Cu}_2\text{ZnSn}(\text{S,Se})_4$  (CZTSSe) have gained massive consideration as alternate materials instead of CIGS and CdTe as they have shown the efficiency of 10%, 11.95%, and 12.62%, respectively.<sup>67–69</sup> But, kesterite material consists of many elements.<sup>70</sup> The presence of many elements in the absorber layer increases the processing cost, and defects are also formed during the processing, which ultimately reduces the overall PCE of solar cells. To overcome this problem, binary element-based materials with similar optical and electronic properties are more suited than quaternary element-based kesterite materials. Tin sulfide (SnS) is also attracting attention of researchers for optoelectronic device applications. In SnS, only one impurity phase of  $\text{SnS}_2$  is present.<sup>71</sup> Therefore, the formation of phase SnS is easier. Also, being consists of

earth-abundant, inexpensive, and eco-friendly elements, SnSe has attracted significant attention. The SnSe has an optimum bandgap of 1.1 to 1.3 eV, an absorption coefficient of  $\sim 10^5 \text{ cm}^{-2}$ , and p-type conductivity with high carrier concentration ( $10^{17} \text{ cm}^{-3}$ ), making it a fabulous material for solar photovoltaic applications.<sup>8</sup> Due to its optimum bandgap (1.3 eV) and high absorption coefficient ( $10^5 \text{ cm}^{-1}$ ), a thin layer of 300 nm thickness can absorb most of the useful solar spectrum. The tunable bandgap of thin-film SnSe can help absorb more solar radiation and enable more photon absorption, resulting in more electron-hole pair generation. The theoretical efficiency for material with a bandgap of 1.3 eV is 32%.<sup>72</sup> The SnSe thin film can be deposited using various chemical and physical routes like chemical bath deposition (CBD),<sup>73</sup> electrodeposition,<sup>74</sup> spray pyrolysis,<sup>75</sup> atomic layer deposition (ALD),<sup>17</sup> thermal evaporation,<sup>46</sup> sputtering.<sup>76,77</sup>

The device architecture for SnSe TFSC is SLG/Mo/SnSe/CdS/i-ZnO/TCO/metal grid, similar to CIGS,<sup>78</sup> and CZTS/CZTSe/CZTSSe TFSCs.<sup>79</sup> Therefore, a lot of modifications and optimizations are required. Also, many issues need to be mitigated to

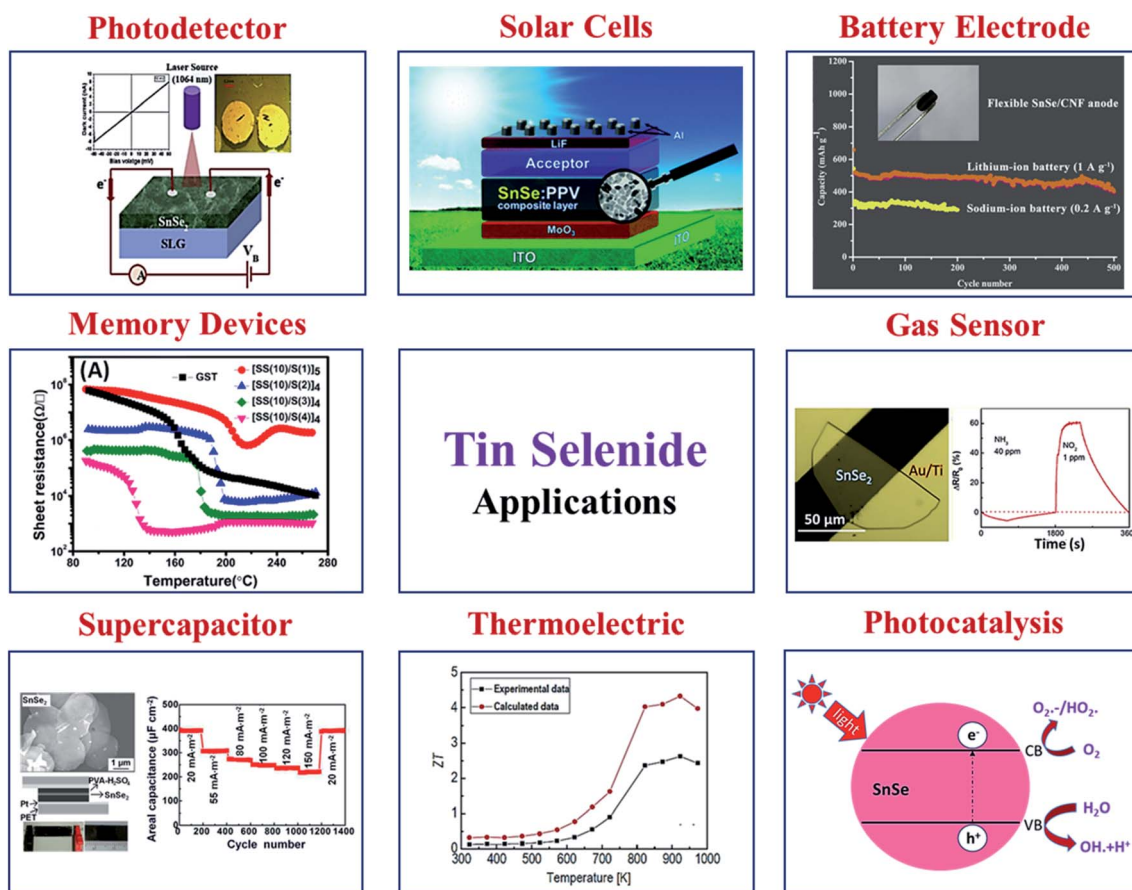


Fig. 5 Various applications are based on SnSe materials. Solar cell, this figure has been adapted/reproduced from ref. 59 with permission from ACS, copyright 2010". Battery electrode, this figure has been adapted/reproduced from ref. 60 with permission from Elsevier, copyright 2020". Gas sensor, this figure has been adapted/reproduced from ref. 61 with permission from ACS, copyright 2019". Photocatalysis, this figure has been adapted/reproduced from ref. 62 with permission from Scielo, copyright 2017". Thermoelectric, this figure has been adapted/reproduced from ref. 38 with permission from Elsevier, copyright 2018". Supercapacitor, this figure has been adapted/reproduced from ref. 63 with permission from ACS, copyright 2014". Memory devices, this figure has been adapted/reproduced from ref. 64 with permission from AIP, copyright 2014". And photodetector, this figure has been adapted/reproduced from ref. 65 with permission from Elsevier, copyright 2020".



improve efficiency.<sup>80</sup> Consequently, it is essential to conduct a study on SnSe material to develop environmentally-friendly, low-cost TFSCs. Any photovoltaic materials should have good optical properties, appropriate band gap, good absorption coefficient, high carrier concentration, and better transport properties. Solar cell efficiency may get affected due to losses, *i.e.*, optical losses, non-absorption, thermalization, reflection loss, transmission loss, area loss, collection losses, and resistance losses.<sup>81</sup> However, the formation of pure phase SnSe thin films has not been reported as the phase lies in a very narrow region, and fine compositional tuning is required to achieve SnSe pure phase. Sn and Se's reaction mechanism plays a crucial role in attaining pure phase SnSe thin films.<sup>19,82</sup> Reddy *et al.* studied the effect of selenization temperature, elemental composition, and selenization pressure on SnSe. They observed that the single-phase SnSe thin films could be achieved during the selenization process in the temperature ranges from 300 to 500 °C.<sup>77</sup> The solar cell efficiency depends on various parameters, like deposition condition, the absorber's crystallinity, carrier concentration, and the p-n junction's nature. The lower solar cell efficiency based on SnSe thin films may be due to SnSe absorber materials' low quality. The parameters that highly influence SnSe thin film quality are phase purity, higher crystallinity, and larger grain size with fewer pinholes. It has been observed that during SnSe phase formation, there is a possibility of the formation of the SnSe<sub>2</sub> phase. Thus, the SnSe<sub>2</sub> phase's impact on the solar cell performance of SnSe thin films needs to be carefully studied and optimized.<sup>8</sup> The electronic properties of p-n junction interfaces are strongly influenced by the discontinuities between the valence band (VB) maxima and conduction band (CB) minima of each material, which restrict the electron transport across the junction interfaces. Band alignment of SnSe with an n-type buffer layer, the density of defect states at the hetero-junction interface, and the nature of back contact also influence the solar cell performance. Therefore, band offsets play a crucial role in reducing the dark current in a diode, reducing the photon-generated carrier losses, and improving the overall solar cell conversion efficiency. The n-type buffer layer with a wider band gap is needed to overcome the optical losses and provide better alignment with SnSe absorber material.

In 1990, Singh *et al.* reported SnSe based solar cells with power conversion efficiency (PCE) of 2.3%.<sup>75</sup> The fabricated solar cell showed an open-circuit voltage ( $V_{oc}$ ) of 410 mV, short-circuit current density ( $J_{sc}$ ) of 9.20 mA cm<sup>-2</sup>, and a fill factor (FF) of 49%, respectively. Rahman *et al.* fabricated heterojunction (p-SnSe/n-Si) solar cells and achieved an efficiency of about 6.44%.<sup>83</sup> They observed improved power conversion efficiency due to the improvement in the  $J_{sc}$  with a graded junction.<sup>83</sup> Shinde *et al.* fabricated SnSe thin-film solar cells using electro-deposition and exhibited device efficiency of 1.4% using an absorber film with 300–400 nm-sized grains, uniform, and dense film morphology.<sup>73</sup> Franzman *et al.* synthesized SnSe/PPV (poly[2-methoxy-5-(3',7'-dimethyl octyl oxy)-1,4-phenylenevinylene]) nanocrystals for solar cells and observed improvement in efficiency from 0.03% to 0.06%.

They observed significant improvement in external quantum efficiency (EQE) and  $J_{sc}$  after adding PPV into the SnSe absorber. Fig. 6(a) and (b) show the high-resolution transmission electron microscopy (HRTEM) image of a single nanocrystal, and selected area electron diffraction (SAED) pattern of SnSe. Fig. 6(c) shows a low-resolution TEM image of SnSe nanocrystals, and Fig. 6(d) shows the schematic device structure of SnSe solar cells.<sup>59</sup> In 2012, Mathews *et al.* reported the first inorganic thin-film solar cell and exhibited a device efficiency of 0.03% ( $V_{oc} = 140$  mV,  $J_{sc} = 0.7$  mA cm<sup>-2</sup>). The low device efficiency is mainly due to poor junction quality.<sup>74</sup> Makori *et al.*<sup>84</sup> fabricated solar cell with device structure glass/Ag/CdO:Sn/SnSe/Ag by using thermal evaporation and reported device efficiency of 0.59% ( $V_{oc} = 273$  mV,  $J_{sc} = 0.993$  mA, and FF = 69%). The observed lower efficiency is mainly due to the thin absorber (148 nm) layer, which attributed to lower  $J_{sc}$  in the device. The above discussion is summarized in Table 1.

Table 1 shows that the efficiency of SnSe based solar cells is well below Cu<sub>2</sub>ZnSn(S,Se)<sub>4</sub>.<sup>85</sup> The observed lower efficiency is mainly due to bulk defects, secondary phase formation in SnSe, non-ideal band-alignment at heterojunction interfaces, and back and front electrode recombination. The absorber quality plays a significant role, and a sound absorber should have a dense, pinhole-free, smooth surface of the film for high-efficiency heterojunction device applications. The low

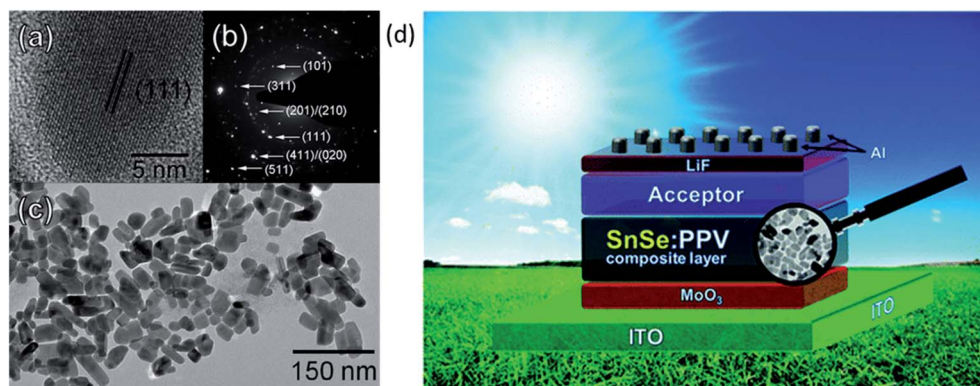


Fig. 6 (a) HRTEM image of a single nanocrystal. (b) Selected area electron diffraction (SAED) pattern of SnSe. (c) Low-resolution TEM image of SnSe nanocrystals. (d) The device structure of SnSe solar cells. These figures has been adapted/reproduced from ref. 59 with permission from ACS, copyright 2010".



Table 1 Summary of SnSe based thin-film solar cells

Deposition method	Absorber	Device architecture	$V_{oc}$ (mV)	$J_{sc}$ (mA cm <sup>-2</sup> )	FF (%)	$\eta$ (%)	Area (cm <sup>2</sup> )	Ref.
Evaporation	SnSe	Si/Al/SnSe/In	425	17.23	44	6.44	0.2	83
Evaporation	SnSe	SLG/FTO/Se/SnSe/Ag	410	9.20	49	2.30	—	75
Electro-deposition	SnSe	ITO/CdS/SnSe/Au	370	5.37	30	1.40	—	73
Co-evaporation	SnSe	SLG/Mo/SnSe/CdS/ i-ZnO/ITO/Ni-Al	172	18.87	31.2	1.02	0.42	10
Sputtering	SnSe	Glass/Ag/CdO:Sn/SnSe/Ag	207	0.90	0.69	0.59	—	84
Electrodeposition	SnSe	Tec15/CdS/SnSe/carbon-paste	140	0.7	31	0.03	0.15	74

interfacial quality could primarily affect the device performance and is responsible for lower  $V_{oc}$  and FF in the solar cells. This is mainly due to CB discontinuity/band offsets at the interfaces, which leads to the CB's energy barrier and restricts carrier transportation across the p-n junction. Apart from this, absorber thickness, grain size, bandgap, and deposition methods also influence the device performance.

## 4.2 Thermoelectric generators

It requires energy to maintain the sustainability and development of humankind. Energy sources are minimal, and there arises a question about the world's power crisis and its ultimate solution. A thermoelectric device that can convert the waste heat into electricity can contribute to a remarkable extent in this.

### 4.2.1 SnSe thermoelectric

*SnSe single crystal.* With the report of a high  $ZT$  value of 2.6 along the  $b$  axis in single-crystal SnSe,<sup>32</sup> this material became the hot topic of research in the field of thermoelectric (TE) domains since 2014. SnSe is anisotropic and exhibits the second-order displacive phase transition from  $Pnma$  (distorted) to  $Cmcm$  (non-distorted) at 750–800 K. It shows a moderate power factor of  $10.1 \mu\text{W cm}^{-1} \text{K}^{-2}$  at  $\sim 850$  K along the  $b$  axis, which is comparable or even lesser than typical thermoelectric materials values.<sup>86–89</sup> The thermal conductivity of SnSe has astonished the researchers. Using the Gruneisen parameter, it has been estimated that SnSe exhibits the ultra-low thermal conductivity ( $<0.25 \text{ W m}^{-1} \text{K}^{-1}$ ) due to the strong anharmonicity in bonding.<sup>32</sup> Theoretical studies<sup>90,91</sup> showed that optimized carrier concentration lies in the range  $10^{19}$ – $10^{20} \text{ cm}^{-3}$  for SnSe for high  $ZT$ . For n-type SnSe, the estimated  $ZT$  is 3.1 along  $a$  axis at 770 K at carrier concentration of  $2.8 \times 10^{19} \text{ cm}^{-3}$ , while the study showed that the n-type SnSe performed better than p-type SnSe.<sup>92</sup>

Motivated by the high  $ZT$  along the  $b$  axis in SnSe, Li *et al.*<sup>38</sup> conducted the theoretical study on the p-type SnSe using the first-principle calculation. A figure of merit ( $ZT$ ) was greatly affected by the carrier concentration and also sensitive to the phases of SnSe ( $Pnma$  and  $Cmcm$ ) (shown in Fig. 7(a) and (b)). Fig. 7(c) shows a comparison of theoretical and experimental data.<sup>32</sup> This study showed that a maximum  $ZT$  of 4.33 could be achieved for the  $Cmcm$  phase of SnSe at 923 K for carrier concentration of  $1.84 \times 10^{19} \text{ cm}^{-3}$ , along the  $b$  axis.

Through another study, for Na doped SnSe single crystal, Zhao *et al.* reported the device with  $ZT_{\text{device}} = 1.34$  in the temperature range 300–773 K. With  $ZT_{\text{max}}$  of 2.0 at 773 K, the device showed the conversion efficiency of  $\sim 16.7\%$ , which was higher than many Pb-based thermo-electric modules.<sup>93</sup> Various researchers showed different  $ZT$  values for the single crystal. For example, Wei *et al.*<sup>94</sup> showed  $ZT$  of  $\sim 1$ ,  $\sim 0.8$ , and  $\sim 0.25$  along  $b$ ,  $c$ , and  $a$  axis, respectively for in-house made fully dense single-crystalline SnSe, Jin *et al.*<sup>95</sup> showed  $ZT \sim 1$  for single crystal grown by a vertical vapor deposition method. Different types of doping are being done to enhance the single-crystal performance. Chang *et al.*<sup>96</sup> doped bromine to make the SnSe conduction n-type and achieved record high  $ZT \sim 2.8 \pm 0.5$  at 773 K along out of the plane direction. With the help of density functional theory and scanning tunneling microscopy, they pointed out that delocalized p electrons of Sn and Se near conduction band minima contribute more to orbital overlapping in out of the plane direction. The S doping in single-crystal SnSe decreased the carrier concentration due to charge trapping at grain interface.<sup>97</sup> Hence, S doping decreased the single crystal's thermoelectric performance rather than enhancing as predicted by theoretical calculation.<sup>98</sup> Pb doped n-type single crystal (*via facile* Sn-flux method) showed a 33% enhanced power factor than its pristine SnSe due to increased carrier concentration.<sup>99</sup> Ag-doped single crystal ( $\text{Sn}_{0.97}\text{Ag}_{0.03}\text{Se}$ ) grown in a horizontal Bridgman furnace showed maximum  $ZT \sim 0.95$  along  $a$  axis at 793 K.<sup>100</sup> Bi doping in single-crystal SnSe resulted in n-type conduction and showed the remarkable  $ZT \sim 2.2$  along the  $b$  axis at 733 K.<sup>51</sup>

*Polycrystalline SnSe thermoelectric.* Polycrystalline SnSe has comparably higher thermal conductivity and lower electrical conductivity than its single crystal counterpart. Though phonon scattering is reduced mainly by the grain boundaries,<sup>101</sup> it has higher thermal conductivity than its single crystal due to tin oxide layer<sup>102</sup> or absorption of oxygen.<sup>103</sup> Highly discrete values of  $ZT$  have been reported for polycrystal SnSe made by different growth techniques. Various researchers have tried to achieve a comparable result to a single-crystal. Many researchers chose to dope (Na, K, Cu, Zn, *etc.*) in poly-crystal SnSe to enhance the electrical properties and reduce the thermal conductivity effectively. In one study, the polycrystalline SnSe thermoelectric material showed  $ZT \sim 0.5$  at 823 K.<sup>104</sup> Several methods improve the device's thermoelectric performance, like large mass fluctuations, band gap engineering, alloying, doping, nano-structuring, *etc.*<sup>105</sup> Gong *et al.*<sup>47</sup> studied the effect of Cu doping





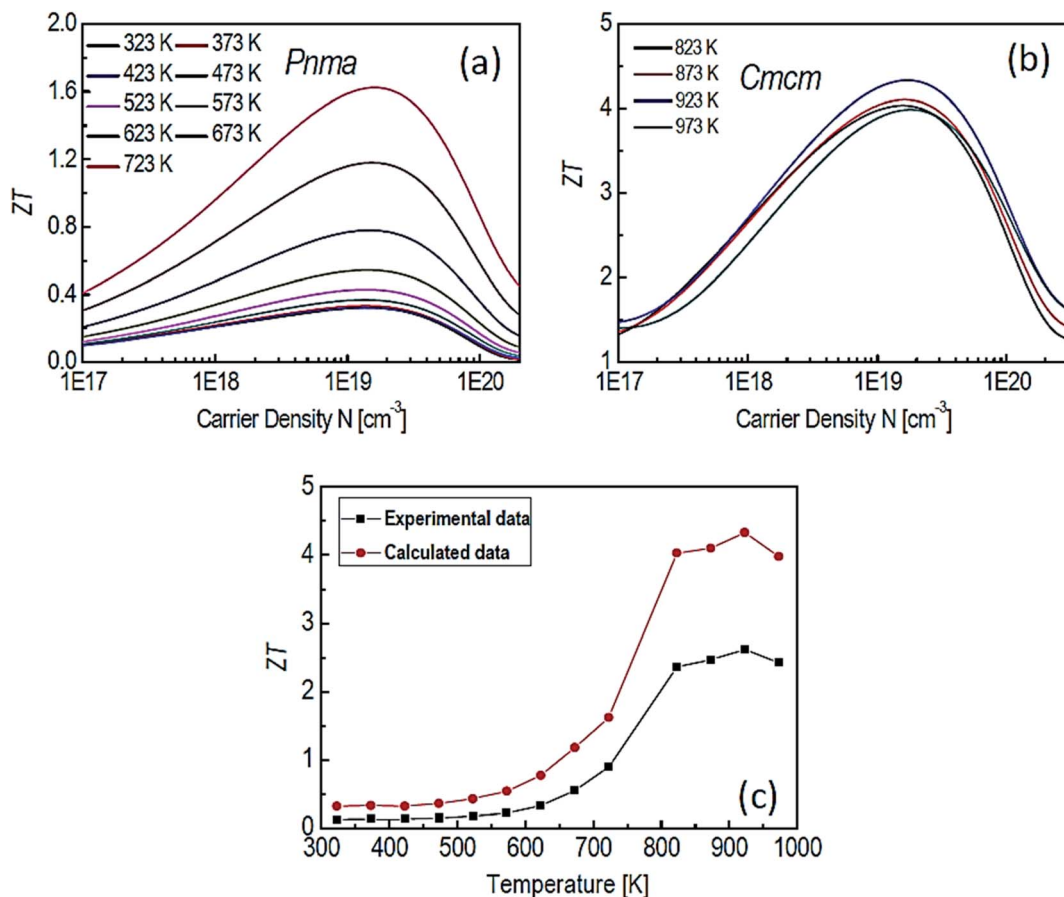


Fig. 7 The ZT of SnSe for *Pnma* (a), *Cmcm* (b) phases and comparison between experimental data<sup>32</sup> and calculated data (c). These figures have been adapted/reproduced<sup>38</sup> with permission from Elsevier, copyright 2018<sup>38</sup>.

in polycrystalline SnSe by varying the concentrations in the  $\text{Cu}_x\text{Sn}_{1-x}\text{Se}$  from  $x = 0.01$  to 0.004. They found that Cu doping enhanced the electrical conductivity due to the increased carrier concentration, as confirmed by Hall measurement. Still, Cu doping acted as a point defect and hence decreased the carrier mobility. The Seebeck coefficient also increased due to doping. Thus, the overall power factor increased due to doping, and the optimized value was obtained for doping  $x = 0.01$ . Around 60% reductions in the lattice thermal conductivity as compared to undoped poly-crystal SnSe is observed. This was due to nanoprecipitation and mesoscale grains (as evidenced by high-resolution transmission electron microscopy (HRTEM)). Finally, they achieved a ZT of 1.2 at 873 K. Shi *et al.*<sup>106</sup> enhanced the Cu's solubility limit in the SnSe and obtained the ZT of  $\sim 1.41$  for the doping level of  $x = 0.118$ .

Lee *et al.*<sup>102</sup> synthesized polycrystalline SnSe ( $\text{Na}_{0.01}(\text{Sn}_{0.95}\text{Pb}_{0.05})_{0.99}\text{Se}$ ), which exhibited  $ZT \sim 2.5$  at 773 K. They noticed that the oxide formation leads to higher thermal conductivity in poly-crystal SnSe ( $\text{SnO}_2$  has  $\sim 140$  times higher thermal conductivity than SnSe).<sup>107</sup> They removed the oxides layers of the tin in SnSe by Ball milling followed by a reduction in 4%  $\text{H}_2/\text{Ar}$  atmosphere for 6 h at 613 K. They pointed out that ball milling (BM) and reduction processes both are critically important to remove the oxides and its residuals from the sample. A remarkably ultralow

thermal conductivity (total) of  $0.20 \text{ W m}^{-1} \text{ K}^{-1}$  (even lower than the single crystal reported by Zhao *et al.*<sup>32</sup>) (Fig. 8(a)) and  $\sim 2$  times enhanced power factor than the pristine SnSe (unball-milled and unreduced). The performance of pristine, reduced, ball-milled and reduced (BR) compared to with the single-crystal SnSe<sup>32</sup> and Na doped SnSe,<sup>93</sup> is shown in Fig. 8(b). Recently Gainza *et al.*<sup>108</sup> reported the highest ZT of 1.8 for the un-doped polycrystalline SnSe, and even with surface oxides layers on the sample.

The pressure-induced effect modifies the textures of the polycrystal SnSe. A sintering pressure of 60 MPa leads to improved electrical properties. There was a meager improvement in thermal conductivity, resulting in a ZT of 0.7.<sup>109</sup> There is a vast literature on doping in the polycrystalline SnSe, like Na,<sup>110</sup> K,<sup>111</sup> Cd,<sup>112</sup> Ag,<sup>113</sup> and co-doping Na/Ag,<sup>114,115</sup> Bi/Cl,<sup>116</sup> Na/CNT,<sup>117</sup> *etc.* which resulted in the enhancement in the thermoelectric performance of the material. Lu *et al.*<sup>118</sup> enhanced the performance of polycrystalline SnSe by introducing large mass fluctuations by doping sulfur, which led to very low lattice thermal conductivity ( $0.13 \text{ W m}^{-1} \text{ K}^{-1}$  at 873 K). Still, this doping enhanced the bandgap of the SnSe, which lowered the electrical conduction and hence the power factor (low carrier concentration due to large bandgap). Pb co-doping improved the power factor, further increasing the carrier concentration by one order of magnitude. Through co-doping a remarkable ZT of 1.85 at 873 K was achieved.



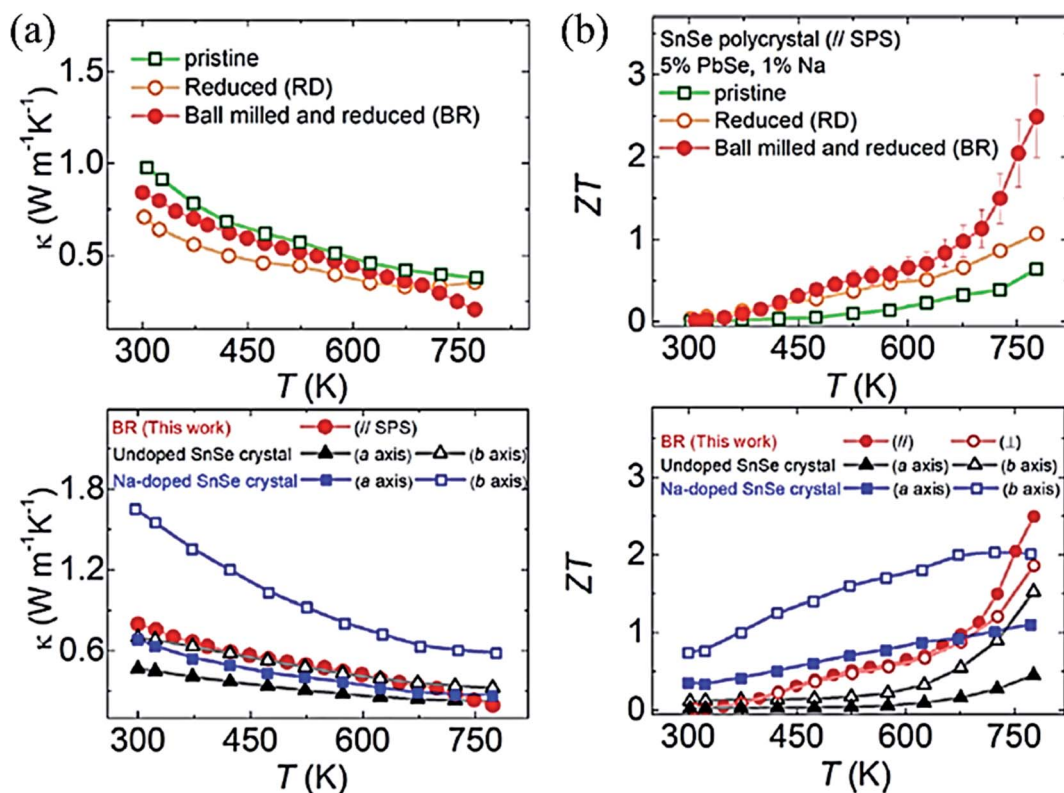


Fig. 8 (a) Total thermal conductivity of polycrystalline SnSe–PbSe (5 mol %) doped with 1 mol% Na for pristine, Reduced (R), and Ball milled then Reduced (BR) samples taken parallel to the press direction of SPS, and total thermal conductivity compared with undoped SnSe<sup>93</sup> and Na doped SnSe,<sup>93</sup> and (b) ZT of polycrystalline SnSe–PbSe (5 mol %) doped with 1 mol% Na for pristine, Reduced (R), and Ball milled then Reduced (BR), and ZT compared with the undoped SnSe<sup>93</sup> and Na doped SnSe.<sup>93</sup> These figures have been adapted/reproduced from ref. 102 with permission from Elsevier, copyright 2019<sup>9</sup>.

The inclusion of micro carbon fibers into the polycrystalline SnSe decoupled the thermal and electrical transport in the host matrix. Carbon fibers acted as good electrical conductors and simultaneously reduced the lattice thermal conductivity ( $0.22 \text{ W m}^{-1} \text{ K}^{-1}$ ) by enhancing the scattering due to the high density of interface, which led to  $ZT \sim 1.3$  at 823 K. This inclusion increased the mechanical stability of the device.<sup>119</sup> The thermoelectric parameters of poly-crystals SnSe are tuned by sintering temperature in the spark plasma sintering (SPS). Zhang *et al.*<sup>120</sup> varied the sintering temperature from 300 to 650 °C in SPS and showed that vacancy defects (Se, Sn, and Se–Sn) were responsible for the performance variation.

The optimized sintering temperature for the best performance of the material was 550 °C. The  $ZT$  of 0.47 was observed at 430 °C. Ge alloying's effect was studied by varying Ge concentration in the range  $x = 0.01$  to  $0.03$  ( $\text{Sn}_{1-x}\text{Ge}_x\text{Se}$ ). The carrier concentration increased from  $3.9 \times 10^{17}$  to  $4.2 \times 10^{19} \text{ cm}^{-3}$  for  $x = 0.03$ , which resulted in a high-power factor of  $\sim 5.10 \mu\text{W cm}^{-2} \text{ K}^{-2}$  at 873 K. A very low lattice thermal conductivity of  $\sim 0.18 \text{ W m}^{-1} \text{ K}^{-1}$  was achieved due to nanoscale grains, precipitation and anharmonicity due to Ge doping. Finally, high  $ZT$  of 1.75 and 2.1 were acquired at 873 K along parallel and perpendicular SPS's pressing direction.<sup>121</sup> The figure of merit ( $ZT$ ) with different doping elements in single and polycrystal SnSe is shown in Fig. 9.

*SnSe thin film.* Low temperature (4–300 K) thermoelectric measurements were carried out by Urmila *et al.*<sup>122</sup> in SnSe thin-

film obtained by the co-evaporation method. The maximum power factor obtained was  $7.2 \times 10^{-4} \text{ W m}^{-1} \text{ K}^{-2}$ . The thermal conductivity was in the range of  $0.023$  to  $0.045 \text{ W m}^{-1} \text{ K}^{-1}$ , which showed a maximum  $ZT$  of 1.2 at 42 K, and thus showed a potential application in a low-temperature TE device. Nair *et al.*<sup>123</sup> fabricated the n–p type thermo-couples with PbSe (p-type), SnSe (p-type), and SnSe<sub>2</sub> (n-type), in the configuration SnSe–SnSe<sub>2</sub>–SnSe–SnSe<sub>2</sub>–SnSe and PbSe–SnSe<sub>2</sub>–PbSe–SnSe<sub>2</sub>–

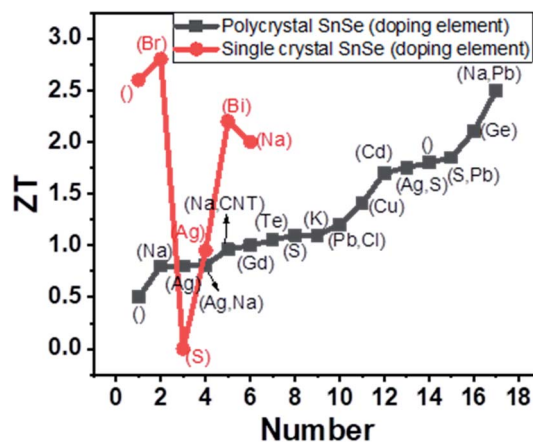


Fig. 9 Variation of  $ZT$  with the doping element given in bracket, where blank ( ) shows the undoped SnSe.



PbSe using chemical method. A voltage was developed for a temperature difference of 20 °C of magnitude 50 mV and 15 mV, respectively (the author did not report the  $ZT$  of the device). Thin-film SnSe deposited by glancing angle (80°) pulsed laser revealed a striking power factor of  $18.5 \mu\text{W cm}^{-1} \text{K}^{-2}$  at 478 K (higher than the single crystal along  $b$  axis) due to enhanced grain boundaries and out of plane thermal conductivity (highest  $0.189 \text{ W m}^{-1} \text{K}^{-1}$  at 340 K). The authors did not report the  $ZT$  and in-plane thermal conductivity.<sup>124</sup> For the first time, Burton *et al.*<sup>125</sup> reported thin film SnSe thermoelectric generator on the glass substrate. Very low thermal conductivity of  $0.08 \text{ W m}^{-1} \text{K}^{-1}$  in the temperature range 375 to 450 K was observed (4 times lower than single-crystal thermal conductivity along  $a$ -axis). Still, it had a low power factor, which resulted in a very low  $ZT$  of 0.055 at 501 K. When the hot side was kept at 618 K, an output of 0.09  $\mu\text{W}$  was achieved. Due to the oxide formation and uncontrolled cold side temperature, the performance degraded for a longer exposure. Song *et al.*<sup>126</sup> synthesized

SnSe thin film of varied thickness in the range of 600–1000 nm by magnetron sputtering and showed effect of post-annealing on the performance. The optimized annealing temperature was 700 K. The power factor for the film with thickness 700 nm was  $2.4 \mu\text{W cm}^{-1} \text{K}^{-2}$  at 675 K. The authors used the thermal conductivity data of bulk polycrystalline SnSe to calculate the  $ZT$  (0.28 at 675 K). Yan *et al.*<sup>127</sup> fabricated a flexible thermoelectric device based on SnSe composite with PEDOT:PSS that showed 13.73 nW output power at 50 K temperature difference at 353 K, and showed possible application in wearable electronics as well due to stability of device even after 1000 cycles of bending. The above discussion is summarized in Table 2.

**4.2.2 SnSe<sub>2</sub> thermoelectric.** A theoretical study on SnSe<sub>2</sub> using the first principal method revealed its anisotropic nature. Due to the bipolar conduction nature (because of the low bandgap) optimal predicted doping values were  $(0.86\text{--}2.03) \times 10^{19}$  and  $(1.71\text{--}2.47) \times 10^{19} \text{ cm}^{-3}$  along with the  $a$  and  $c$  axes, respectively. A maximum power factor of  $11.72 \times 10^{-4} \text{ W m K}^{-2}$

Table 2 Thermoelectric properties of SnSe based materials

Material	Synthesis method	Power factor ( $\mu\text{W cm}^{-1} \text{K}^{-2}$ )	Thermal conductivity ( $\text{W m}^{-1} \text{K}^{-1}$ )	Temp. (K)	$ZT$	Ref.
Pristine SnSe (single crystal)	Bridgman Crystal growth	~9	0.34	923	2.62	32
Single crystal	Vertical Bridgman	8.5	—	850	1	128
Single crystal	Vertical vapor deposition	5.23	0.44	800	1	95
Na doped SnSe single crystal	—	~14	~0.55	773	~2	93
Br doped SnSe single crystal	Temperature gradient	~9	~0.23	773	2.8	96
S doped single crystal	Direct vapor transport	0.0011	—	573	0.0012	97
Pb doped SnSe single crystal	Flux method	1.2	—	300	—	99
Ag-Doped SnSe single crystal	Horizontal Bridgman method	~5.7	0.49	793	0.95	100
Bi-Doped SnSe single crystal	Temperature gradient	~9.75	~0.3	733	2.2	51
Polycrystal SnSe	Solid state reaction (SSR)	~4	<0.5	823	0.5	129
Polycrystal SnSe	Arc melting	~10	~0.4	816	1.8	108
Cu doped polycrystal SnSe	Hydrothermal	3.4	0.26	873	1.2	47
SnSe <sub>0.98</sub> Te <sub>0.02</sub>	—	~4	~0.35	805	1.05	130
Sn <sub>0.97</sub> Gd <sub>0.03</sub> Se	Hydrothermal	6.7	0.4	868	~1	131
Sn <sub>0.985</sub> S <sub>0.25</sub> Se <sub>0.75</sub>	Mechanical alloying	4.5	~0.35	823	1.1	132
Sn <sub>0.978</sub> Ag <sub>0.007</sub> S <sub>0.25</sub> Se <sub>0.75</sub>	—	5.3	~0.25	823	1.75	
Sn <sub>0.90</sub> Pb <sub>0.15</sub> Se <sub>0.95</sub> Cl <sub>0.05</sub>	SSR	6.74	0.5	823	1.2	133
Cu doped polycrystal SnSe	Solvothermal	5.57	0.32	823	1.41	106
Na, Pb doped polycrystal SnSe	SSR	6.85	0.20	773	2.5	134
Polycrystal SnSe	Mechanical alloying	3.9	—	823	0.7	109
Na doped SnSe polycrystal	SSR	4.5	~0.4	773	0.8	110
K doped SnSe polycrystal	Mechanical alloying	~2.9	~0.22	773	1.1	111
Cd doped SnSe polycrystal	Solvothermal	6.9	0.33	823	1.7	112
Ag doped polycrystal SnSe	Facile surfactant-free synthesis	6.34	0.75	850	0.8	113
Ag/Na doped polycrystal SnSe	—	~7.5	~0.5	773	1.33	114
Ag/Na doped polycrystal SnSe	SSR	5	~0.48	773	0.81	115
CNT dispersed Na doped polycrystal SnSe	SSR	4.99	0.40	773	0.96	117
S, Pb doped polycrystal SnSe	Hydrothermal	4.18	~0.20	873	1.85	118
Composite of carbon fiber and polycrystal SnSe	SSR	3.88	~0.22	823	1.3	119
SnSe polycrystal	Solvothermal	~4.4	~0.7	703	0.47	120
Ge doped polycrystal SnSe	Hydrothermal	5.1	~0.21	873	2.1	121
Thin film SnSe	Reactive evaporation	7.2	~0.023	42	1.2	122
SnSe thin film	PLD	18.5	—	478	—	124
SnSe thin film	Thermal evaporation	0.11	~0.11	501	0.055	125
SnSe thin film	Sputtering	2.4	—	675	0.28	126
SnSe thin film	Chemical vapor transport	—	0.7	300	0.16	135
SnSe thin film	Chemical vapor deposition	3.2	1.1	550	0.15	136
SnSe/PEDOT:PSS thin film	Vacuum filtration method	0.24	—	353	—	127



was expected for a doping concentration of  $7.21 \times 10^{19} \text{ cm}^{-3}$  at 800 K along *a* axis. Minimum thermal conductivity predicted along *a* and *c* axes at 300 K were 0.55 and  $0.42 \text{ W m}^{-1} \text{ K}^{-1}$ , respectively. Finally, *ZT* value was expected along *a* direction that showed a value of 0.88 for  $1.94 \times 10^{19} \text{ cm}^{-3}$  at 800 K.<sup>137</sup> Saha *et al.* synthesized the SnSe<sub>2</sub> nanosheets using a solution-based method, and Cl doping was done to enhance the carrier concentration.<sup>138</sup> At room temperature carrier concentration increased from  $7 \times 10^{17}$  to  $2 \times 10^{18} \text{ cm}^{-3}$  due to Cl doping. For the doped sample, almost double power factor of  $1.46 \mu\text{W cm}^{-1} \text{ K}^{-2}$  was achieved at 630 K. As synthesized SnSe<sub>2</sub> nanosheets showed the thermal conductivity in the range of  $\sim 0.45$  to  $0.35 \text{ W m}^{-1} \text{ K}^{-1}$  while doping increased the thermal conductivity in the range of 0.67 to  $0.40 \text{ W m}^{-1} \text{ K}^{-1}$  at 300 to 630 K temperature, respectively. Overall, *ZT* of 0.22 was achieved at 610 K for Cl doped samples. For heavy Cl doping of  $x = 0.12$  by solid-state reaction, *ZT* (out of plane) of  $\sim 0.4$  was achieved by Xu *et al.*<sup>139</sup> Se deficiency and Cl doping's led to enhanced power factor and reduced thermal conductivity resulting in a *ZT* of 0.63 at 673 K along an in-plane direction.<sup>140</sup> A study on Ag doping revealed that it effectively enhances the carrier mobility but decrease the carrier concentration. For 1% doping, the optimized power factor was  $3.50 \mu\text{W cm}^{-1} \text{ K}^{-2}$  at 773 K, which resulted in a *ZT* of  $\sim 0.4$ .<sup>141</sup> In another theoretical study, Ding *et al.* reinvestigated the thermoelectric performance of the SnSe<sub>2</sub> by the first principle method. They considered various models of the phonon-phonon scattering, which may overestimate the thermal conductivity.<sup>33</sup> Room temperature electrical conductivities along *a* and *c* directions were calculated as  $4.97 \times 10^5$  and  $3.39 \times 10^4 \Omega^{-1} \text{ m}^{-1}$ , respectively, at  $10^{20} \text{ cm}^{-3}$  carrier concentration. Overall, for n-type highly doped ( $10^{20} \text{ cm}^{-3}$ ) SnSe<sub>2</sub>, a *ZT* of 3.6 could be achieved at 800 K. Wu *et al.* doped Br to the site of Se in polycrystalline SnSe<sub>2</sub> to enable broad carrier concentrations ( $0.5$  to  $5.6 \times 10^{19} \text{ cm}^{-3}$ , saturation at doping of  $x = 0.01$ ) consistent with the single parabolic model.<sup>142</sup> They reached a *ZT* of 0.6 at 750 K, which was competitive to polycrystalline SnSe. The Cl-doped composite of SnSe–SnSe<sub>2</sub> showed *ZT*  $\sim 0.56$  at 773 K.<sup>143</sup> Dynamical intercalation of Ag into the weak van der Waals

interlayers of SnSe<sub>2</sub> acted as electron donor and introduced line defect, twin boundary dislocation, phase interface, and enhanced phonon–phonon scattering. To further strengthen carrier concentration into Ag-doped SnSe<sub>2</sub>, Liu *et al.* doped it with Cl, which resulted in concentration in the range of  $10^{19}$  due to Fermi leveling up to the conduction band (as shown *via* DFT calculation).<sup>144</sup> Ag interlayer bridge weakened the anisotropy of electrical transport; hence the power factor along parallel and perpendicular directions of pressing were 7.46 and  $8.25 \mu\text{W cm}^{-1} \text{ K}^{-2}$ , respectively, at 789 K, and a maximum *ZT* of  $\sim 1.03$  was achieved at 789 K along the parallel direction of the pressure. Recently, Wang *et al.* showed the room temperature thermoelectric potential of the SnSe<sub>2</sub> embedded with Cu, which enhanced carrier concentration two orders of magnitude. Optimum doping in SnCu<sub>*x*</sub>Se<sub>2</sub> was found to be  $x = 0.01$ , which resulted in an optimized power factor of 1.96 at room temperature and a decrease in thermal conductivity up to  $0.81 \text{ W m}^{-1} \text{ K}^{-1}$  due to increased interface.<sup>145</sup> A very high *ZT* of 0.75 than to pristine SnSe<sub>2</sub> was obtained which was about two orders of magnitude larger than pristine SnSe<sub>2</sub> (Fig. 10(a)).

**Thin-film SnSe<sub>2</sub> thermoelectric.** Thermoelectric performance of SnSe<sub>2</sub> thin films have also been studied. But there are very few reports on this, and the reports are limited to the film's power factor only. Yin *et al.* deposited SnSe<sub>2</sub> thin-film on a glass substrate by spin coating techniques and reported a Seebeck coefficient of  $\sim 126.3 \mu\text{V K}^{-1}$ . The electrical conductivity of the order of  $10^3 \text{ S m}^{-1}$  at room temperature was observed.<sup>146</sup> Chen *et al.* deposited the SnSe thin film on the Si wafer and annealed it in Se vapor. The film was reduced to SnSe<sub>2</sub> and showed enhanced Seebeck coefficient  $631 \mu\text{V K}^{-1}$  for SnSe<sub>2</sub> than  $38.6 \mu\text{V K}^{-1}$  for SnSe. Around 44 times enhanced power factor ( $0.2 \mu\text{W m}^{-1} \text{ K}^{-2}$ ) was obtained.<sup>147</sup> Duong *et al.* deposited SnSe<sub>2</sub> thin film on the Al<sub>2</sub>O<sub>3</sub> substrate by the pulsed laser deposition technique. They reported a power factor of  $8 \mu\text{W m}^{-1} \text{ K}^{-2}$  at 220 K.<sup>148</sup> A detailed summary of the above discussions and *ZT* variation with the doping is provided in Table 3 and Fig. 10(b), respectively.

This material (SnSe) has the highest efficiency among the bulk materials without any doping, but it faces

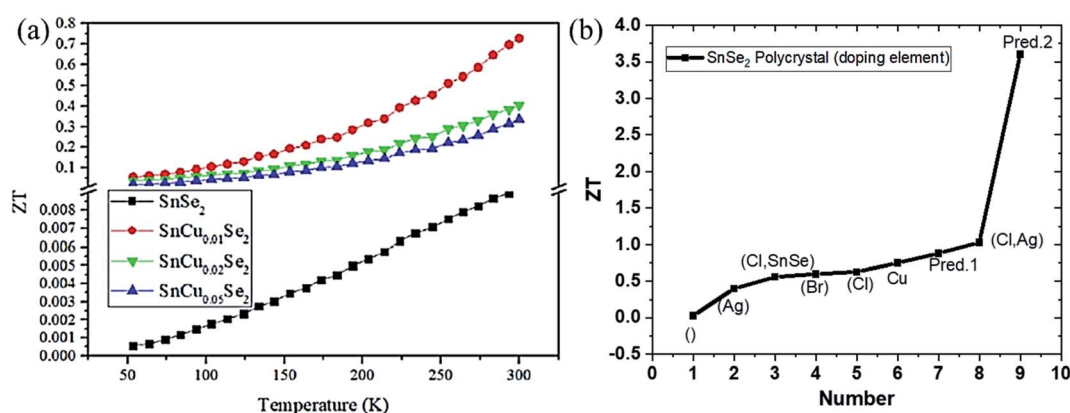


Fig. 10 (a) Figure of merit (*ZT*) of SnCu<sub>*x*</sub>Se<sub>2</sub> ( $x = 0, 0.01, 0.02, 0.05$ ) as a function of temperature. This figure has been adapted/reproduced from ref. 145 with permission from ACS, copyright 2020". (b) Variation of *ZT* with different doping elements, irrespective of operating temperature, where blank () represents undoped SnSe<sub>2</sub>, and Pred. 1 and Pred. 2 mean theoretical predicted values, respectively.



commercialization issues like poor mechanical stability and low thermoelectric property below  $\sim 800$  K. Besides its high  $ZT$ , low cost, easy fabrication, and earth-abundance, SnSe suffers the real possible application likely due to the quick oxides' formation or thermal instability.<sup>150</sup> It readily forms the oxides in the range  $>600$  °C, where it showed maximum  $ZT$ . Protective environment costs hamper its effectiveness as a cheap/high  $ZT$ . Sassi *et al.*<sup>129</sup> showed thermal instability above the phase transition temperature of SnSe, and the working efficiency of the device decreased continuously over cyclic performance due to Se loss. Thus, there should be an optimized temperature range over which there is no loss of performance. Continuous research should go on in the direction to improve the thermal stability up to the temperature for which SnSe showed maximum  $ZT$ . Also, large phonons scattering centers exist in polycrystalline SnSe than in single crystal SnSe due to grain boundaries, yet single crystal SnSe has lower lattice thermal conductivity than the polycrystalline SnSe. Ibrahim *et al.*<sup>151</sup> noted this fact and investigated the reason behind this. They systematically discarded the oxide formation and found many native defects in the single crystal.<sup>152</sup> Their results reflected higher lattice thermal conductivity than the reported one.<sup>32</sup> They also noted that the single crystal grown by Zhao *et al.*<sup>32</sup> was not fully dense and hence may be the reason for low lattice thermal conductivity. But it still requires the study to know the exact causes of the very low lattice thermal conductivity, whether it was due to intrinsic defect or not. SnSe single crystal suffers a low mechanical strength and high growth cost as a single crystal. SnSe single crystal thermoelectric results showed inconsistencies in the different group reports' as shown in Table 2. However, the further study required to unravel the SnSe performance, where it seeks the possibilities of enhancement in  $ZT$ .<sup>38</sup> Another phase, SnSe<sub>2</sub>, was also predicted by theoretical consideration as the best thermoelectric material that can achieve  $ZT$  around 3.5. Till date,  $ZT \sim 1.1$  is achieved in this material. Also, both materials (SnSe and SnSe<sub>2</sub>) showed inferior TE property in thin film form. Thus, this material demands greater attention and research to achieve predicted  $ZT$ . Researchers should pay attention to new approaches to improve the SnSe thermoelectric material's performance like

nano-inclusion,<sup>153</sup> decoupling of interrelated parameters,<sup>119</sup> nano-structuring and texturing, *etc.* Simultaneously, the theoretical study should be carried out to optimize the threshold values of the parameters. Finally, tin-selenide has emerged as a futuristic material that showed the best efficiency. Both materials (SnSe and SnSe<sub>2</sub>) demand more research in the thin film and bulk thermoelectric form because of the inferior TE property and device stability.

### 4.3 Photodetectors (PDs)

**4.3.1 SnSe based PDs.** SnSe act as suitable optoelectronic material due to its high absorption coefficient ( $>10^4$  cm<sup>-1</sup>),<sup>154</sup> low band gap and its tunable band gap over an extensive range.<sup>155</sup> Solvothermal processed SnSe nanorod showed repeatable and stable photoresponse.<sup>156</sup> SnSe thin film grown on the n-Si substrate showed an ultrahigh response/recovery time of 0.9/17.3  $\mu$ s in a position-dependent detector.<sup>157</sup> A flexible device of SnSe on the mica substrate (grown by PLD) showed high responsive behavior of the device in a longer wavelength region (370–808 nm). Also, it reflected its durability in bending test for the flexible device. The device showed the highest photoresponse of 5.5 A W<sup>-1</sup> which is attributed to its shallow potential barrier with Bi<sub>2</sub>Te<sub>3</sub> contacts (perfect band alignment) and low potential fluctuation at the interface.<sup>158</sup> Under white light illumination, SnSe nanoplate and graphene composites showed very high photosensitivity (1110%) and short response time ( $\sim 1$  s).<sup>159</sup> The SnSe film grown on Si substrate by vapor deposition and dispersed with graphene oxide quantum dots showed improved response time and photocurrent than the SnSe film in a larger spectrum.<sup>160</sup> Self-powered photodetection was demonstrated in the SnSe/Si heterostructure device, with very high detection and ultrafast response in  $\mu$ s range. Magnetron sputtered SnSe film on Si(100) showed a nearly ideal diode property with a reported power exponent of 0.92 and ideality factor of  $\sim 1$ , which revealed a trap-free and high-performance of the device.<sup>36</sup> Zhong *et al.*<sup>161</sup> showed the thin film's superiority due to the Marangoni effect over other solution-based methods. The flexible and vertically fabricated Gr/SnSe/Gr device showed a responsivity of 38 mA W<sup>-1</sup>. Magnetron sputtered SnSe thin film on Si substrate

Table 3 Thermoelectric properties of doped-SnSe<sub>2</sub> based materials

Material	Synthesis method	Power factor ( $\mu$ W cm <sup>-1</sup> K <sup>-2</sup> )	Thermal conductivity (W m <sup>-1</sup> K <sup>-1</sup> )	Temp. (K)	$ZT$	Ref.
Cl doped SnSe <sub>2</sub>	Facile low-temperature solution	1.46	$\sim 0.45$	610	0.22	138
Cl doped SnSe <sub>2</sub>	Solid-state reaction followed by hot press	$\sim 7.0$	$\sim 1.6$	673	0.4	139
Cl doped SnSe <sub>2</sub>	Solid-state reaction followed by ball milling and SPS	$\sim 8$	$\sim 1.2$	673	0.63	140
Ag-Doped SnSe <sub>2</sub>	Mechanical alloying followed by SPS	3.50	$< 1$	773	0.4	141
Br doped SnSe <sub>2</sub>	Solid-state reaction followed by quenching and hot press	—	—	750	0.6	142
Cl doped SnSe–SnSe <sub>2</sub>	Solid state reaction	—	0.42	773	0.56	143
Ag and Cl doped SnSe <sub>2</sub>	Bottom up	7.46	0.57	789	1.03	149
Cu embedded SnSe <sub>2</sub>	Solid-state reaction followed by ball mill and SPS	1.96	0.81	300	0.75	145



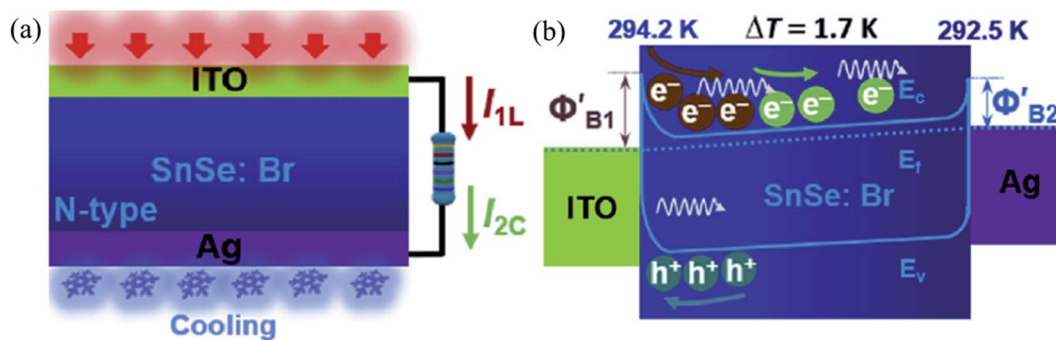


Fig. 11 Schematic working state (a) and energy band diagram (b) of the device under the combined action of light illumination and cooling ( $I_{1L}$  = photovoltaic current,  $I_{2C}$  = thermocurrent). Reprinted with permission from. These figures has been adapted/reproduced from ref. 163 with permission from Elsevier, copyright 2019\*.

showed broadband photoresponse in the 404–980 nm region. A maximum responsivity of  $277.3 \text{ A W}^{-1}$  was reported. This high value of responsivity was attributed to the film's trap states, which was confirmed by the nanosecond order's carrier lifetime observed in the time-resolved photoluminescence studies.<sup>162</sup> Ouyang *et al.*<sup>163</sup> harnessed combined photovoltaic and thermoelectrics effect to enhance the self-biased n-type SnSe:Br (ITO/n-SnSe/Ag) photodetector performance. They coupled the thermoelectric property with the material's photovoltaic property *via* maintaining the temperature difference with Peltier cooler's help at the one contact (Ag contact) of the device and another under the illumination of 760 nm radiation. The thermoelectric diffusive process accelerates photoelectrons generated under photovoltaic conditions. This is due to specific heat capacities of electrons and photonic drags under the thermal gradient resulting in enhanced performance (Fig. 11(a) and (b)). Simultaneously, in the energy band aspect, the thermoelectric effect reduced the Schottky barrier height at the ITO side and increased at the Ag side, which resulted in the de-acceleration of the photoelectrons. The combined effect enhanced the PD parameters. The light current and voltage were

improved by 38.1 and 81.9%, respectively, at a cooling temperature difference of 1.5 K (the Ag side was cooler compared to the ITO side).

When the Ag side is heated w.r.t. ITO with 4.5 K difference, the opposite result was observed, *i.e.*, there was a decrease in photocurrent and voltage. Hence authors concluded that the coupling of the effect is polarity dependent. Record high value (comparable to the commercialize Ge and Si photodetectors) of responsivity ( $R$ ) and detectivity ( $D$ ) were reported in the nanowire SnSe.<sup>164</sup> Above discussions have been summarized in Table 4.

**4.3.2 SnSe<sub>2</sub> based PDs.** The tunable bandgap and high absorption coefficient of the SnSe<sub>2</sub> have made it a potential candidate for optoelectronics applications. For the first time, Zhou *et al.*<sup>173</sup> reported ultrathin SnSe<sub>2</sub> flake (1.5 nm) grown by a CVD method using SnI<sub>2</sub> as a precursor, which exhibited high responsivity ( $R$ )  $1100 \text{ A W}^{-1}$  with meager response time in millisecond. However, the far value of power exponent (0.7), reflected its impurity and defects. Bilayered SnSe<sub>2</sub> of 1T type structure with  $D_{3d}$  point group symmetry showed high responsivity of

Table 4 Performance characteristic of SnSe based PDs

Device structure	$\lambda$ (nm)	Bias (V)	$R$ ( $\text{A W}^{-1}$ )	$D$ (Jones)	Response/recovery time (s)	Ref.
ITO/SnSe/ITO	White light	5	—	—	3	165
Pd/n-Si/SnSe/Pd	780	—	—	—	$0.9/17.3 \times 10^{-6}$	157
Bi <sub>2</sub> Te <sub>3</sub> /SnSe/Bi <sub>2</sub> Te <sub>3</sub>	370	20	5.5	$6 \times 10^{10}$	—	166
Au/SnSe/graphene/Au	White light	15	—	—	$1/<1$	159
Ag/Si/SnSe/graphene oxide QD/Ag	650	—	—	—	0.18/0.75	160
Pd/SnSe/Si/In	850	0	0.566	$4.4 \times 10^{10}$	$1.6/47.7 \times 10^{-6}$	167
Gr/SnSe/Gr	400	1	0.038	—	$\sim 0.18$	161
Ag/Si/SnSe/Ag	404	15	277.3	$7.6 \times 10^{11}$	0.35/1.83	168
ITO/n-SnSe/Ag	760	0	$3.97 \times 10^{-3}$	—	$81/122 \times 10^{-6}$	163
Cr/Au/Si/SnSe (nanowire)/Cr/Au	830	3	$1.0 \times 10^4$	$3.3 \times 10^{12}$	$460/520 \times 10^{-6}$	164
Mica/In <sub>2</sub> Se <sub>3</sub> /SnSe/Au	405	5	0.350	—	0.156/0.139	169
Si/SiO <sub>2</sub> /WSe <sub>2</sub> /SnSe/Ti/Au	1064	5	$6.6 \times 10^{-3}$	—	—	170
	671		$31.8 \times 10^{-3}$		—	
	457		$99 \times 10^{-3}$		0.0082/0.0084	
PCB/Mica/SnSe (nanocrystal)/Ag	(Sunlight)	2	$0.54 \times 10^{-3}$	$1.06 \times 10^9$	1.5/1.7	171
PET/SnSe/Pd	404	20	1745.5	$\sim 4.2 \times 10^{12}$	1.7/4.7	172
	850		78.6	$\sim 9 \times 10^{11}$	0.23/0.27	



$0.5 \text{ A W}^{-1}$  and a swift response time of  $\sim 2 \text{ ms}$ .<sup>174</sup>  $\text{SnSe}_2$  acted as an efficient charge separator in the heterostructure devices. It collected electrons from the  $\text{WSe}_2$ ,<sup>175</sup>  $\text{MoS}_2$ <sup>176</sup> (which have deeper conduction minima than these) so efficiently that it improved performance parameters several times better.  $\text{SnSe}_2$  with  $\text{MoS}_2$  interface resulted in enhanced responsivity from  $37.3$  to  $9.1 \times 10^3 \text{ A W}^{-1}$  under  $500 \text{ nm}$  illumination.<sup>176</sup>

Mukhokosi *et al.* studied the thickness-dependent optical properties of the  $\text{SnSe}_2$  thin film grown by DC sputtering followed by selenization and also studied the photodetection performance (Fig. 12(a) and (b)).<sup>30</sup> Films with thickness  $< 140 \text{ nm}$  did not show any IR response. The film with  $1200 \text{ nm}$  thickness had a slow response (time), and responsivity was  $0.4 \text{ mA W}^{-1}$ .<sup>30</sup> Mukhokosi *et al.*<sup>177</sup> also reported the self-powered, organic-inorganic hybrid heterostructure PD consisting of poly (3,4-ethylene dioxythiophene):poly(styrene sulfonate) (PEDOT:PSS) and  $\text{SnSe}_2$ . The device had responsivity in the range of  $1.4\text{--}2.6 \mu\text{A W}^{-1}$ . The response time improved to  $1.33 \text{ s}$ . Mukhokosi *et al.*<sup>178</sup> reported highly enhanced response time and enhanced responsivity. The device was fabricated on the p-type Si substrate by DC sputtering, followed by selenization. Kumar *et al.*<sup>65</sup> deposited  $\text{SnSe}_2$  thin-film on the soda-lime glass substrate and reported the improved responsivity at very low bias and fast response time in the IR range (Fig. 12(c) and (d)). Although the band gap was in the visible region, photodetection

was observed in the IR region which was attributed to the defect/trap assisted levels lying in the forbidden region.<sup>65</sup> Usually,  $\text{SnSe}_2$  shows high absorption coefficient<sup>30,65</sup> but poor chemical stability against environmental conditions (degrade into  $\text{SnSe}$  after prolonged exposure). Therefore, Gao *et al.* used graphene (Gr) to improve the chemical stability.  $\text{SnSe}_2$  was sandwiched between graphene. Monolayer Gr acted as useful ohmic contacts and also had weak absorption in the UV to IR range.<sup>179</sup> This sandwiched structure of  $\text{SnSe}_2$  with Gr showed responsivity of  $1.09 \times 10^3 \text{ A W}^{-1}$ , and the response time was  $30 \text{ ms}$ . Murali *et al.* studied the substrate-induced effect on  $\text{SnSe}_2$  flake (i) suspended over contacts and (ii) supported on Si substrate. It showed a vast difference in rise and decay times.<sup>180</sup> The suspended structure showed low gain but quicker decay than the supported one because of the device's interface-induced trap states. Field-effect transistors based on multi-layer  $\text{SnSe}_2$  showed very high responsivity and detectivity. The device showed its potential application comparable to the other commercialized photodetector.<sup>9</sup> The above discussion is summarized in Table 5.

#### 4.4 Gas sensors

The gas sensor is a basic need for environment cleaning. A gas sensor is a device, which detects harmful gases and sends an

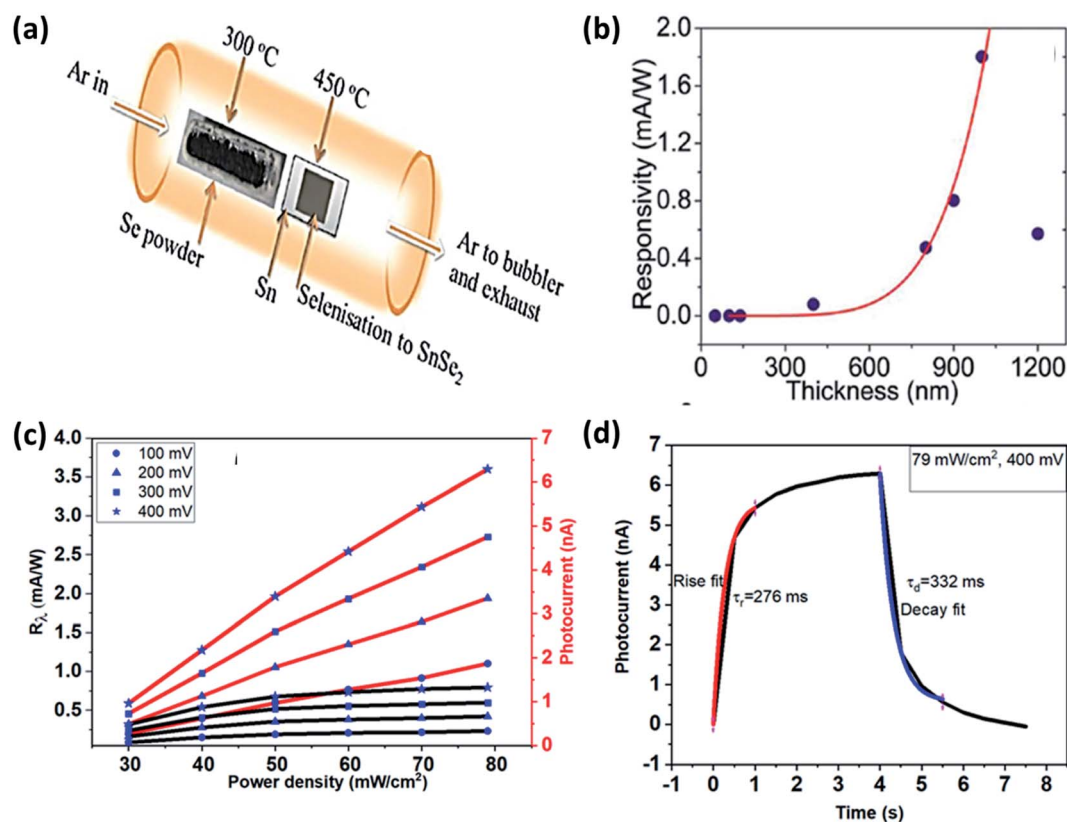


Fig. 12 (a) Schematic of selenization of DC sputtered Sn film and (b) thickness-dependent  $\text{SnSe}_2$  thin films' responsivity. These figures have been adapted/reproduced from ref. 30 with permission from Nature, copyright 2017. (c) Responsivity and photocurrent with power density at different bias voltage and, (d) response and recovery time at  $400 \text{ mV}$  bias, these figures have been adapted/reproduced from ref. 65 with permission from Elsevier, copyright 2020.

Table 5 Summary of performance characteristics of SnSe<sub>2</sub> based photodetectors<sup>a</sup>

Device structure	$\lambda$ (nm)	Bias (V)	$R$ (A W <sup>-1</sup> )	Detectivity (jones)	Rise/decay time (s)	Ref.
Ti/Au/Si/SnSe <sub>2</sub> (flake)/Ti/Au	530	3	1100	10 <sup>10</sup>	14.5/8.1 × 10 <sup>-3</sup>	57
Ti/Au/Si/2LSnSe <sub>2</sub> /Ti/Au	633	0.1	0.5	—	2.2/3.2 × 10 <sup>-3</sup>	174
Ni/Au/ITO/WSe <sub>2</sub> (flake)/SnSe <sub>2</sub> (flake)/SiO <sub>2</sub> /Si/Ni/Au (vertical hetero-structure)	785	0.4	1139	—	8/32 × 10 <sup>-6</sup>	175
Cr/Au/SnSe <sub>2</sub> (flake)/MoS <sub>2</sub> (flake)/Cr/Au	500	1	9.1 × 10 <sup>3</sup>	9.3 × 10 <sup>10</sup>	0.2/0.6	181
Cr/Au/SLG/SnSe <sub>2</sub> (thin film)/Cr/Au	1064	10	0.4 × 10 <sup>-3</sup>	10 <sup>11</sup>	2.5/3.68	30
Cr/Au/SLG/SnSe <sub>2</sub> (thin film)/PEDOT:PSS/Cr/Au	1064	0	1.4–2.6 × 10 <sup>-6</sup>	~10 <sup>8</sup>	1.33/1.22	177
Cr/Au/p-Si/SnSe <sub>2</sub> /Cr/Au	1064	10	0.16	~1.5 × 10 <sup>9</sup>	57/34 × 10 <sup>-6</sup>	182
Au/SLG/SnSe <sub>2</sub> (thin film)/Au	1064	0.4	0.796 × 10 <sup>-3</sup>	5.62 × 10 <sup>7</sup>	0.276/0.332	65
Au/Gr/SnSe <sub>2</sub> (flake)/Gr/Au	532	0.5	1.09 × 10 <sup>3</sup>	1.2 × 10 <sup>12</sup>	30.2/27.2 × 10 <sup>-3</sup>	183
Ni/Au/SnSe <sub>2</sub> (flake)/Ni/Au (suspended structure), Ni/Au/Si/SnSe <sub>2</sub> (flake)/Ni/Au (supported structure)	White light	0.1	115, 8.66 × 10 <sup>4</sup>	~10 <sup>13</sup>	SED-66 s, DED-2.27 s, 53 s	180
Cr/Au/n-Si/SnSe <sub>2</sub> (single crystal flake)/Cr/Au	450	3	5.11 × 10 <sup>5</sup>	2.79 × 10 <sup>13</sup>	—	9

<sup>a</sup> SED = single exponential decay, DED = double exponential decay.

alarm. There are many toxic gases in the environment (CO<sub>2</sub>, NO<sub>2</sub>, NH<sub>3</sub>, SO<sub>2</sub>, H<sub>2</sub>S, CH<sub>4</sub>, etc.). Gas sensing parameters are response time, recovery time, selectivity, gas concentration, temperature, detection limit, etc. Tin based chalcogenide (SnSe and SnSe<sub>2</sub>) have shown promise for gas sensing applications. A brief discussion of tin chalcogenide (SnSe and SnSe<sub>2</sub>) materials for gas sensing applications is described. The gas sensing mechanism of the tin chalcogenide (SnSe and SnSe<sub>2</sub>) gas sensor is based on the adsorption of gas molecule and charge transfer.

Physisorption-based charge transfer in tin chalcogenide (SnSe and SnSe<sub>2</sub>): To evaluate the interaction efficiency of the SnSe/SnSe<sub>2</sub> monolayer and gas molecules, the absorption energy ( $E_a$ ), Hirschfeld charge transfer ( $Q$ ), and the distance ( $d_0$ ) of nearest atoms between the gas molecule and the SnSe/SnSe<sub>2</sub> layer were calculated. The absorbed energy is defined as:

$$E_a = E_{\text{total}} - E_{\text{gas}} - E_{\text{material}} \quad (1)$$

$E_{\text{gas}}$ ,  $E_{\text{material}}$ , and  $E_{\text{total}}$  are the gas molecule's total energy, SnSe/SnSe<sub>2</sub> monolayer, and gas molecule-SnSe/SnSe<sub>2</sub> system, respectively.<sup>184</sup> The adsorption properties, including adsorption energy ( $E_a$ ), equilibrium distance ( $d_0$ ), and Hirshfeld charge transfer ( $Q$ ) for SnSe, are listed in Table 6. The adsorbed CO, CO<sub>2</sub>, CH<sub>2</sub>O, NO<sub>2</sub>, and SO<sub>2</sub> gas molecules on the  $\beta$ -SnSe monolayer have the  $E_a$  values -0.202, -0.175, and -0.322, -0.829, and -0.499 eV, respectively, and the Hirshfeld charge transfer values are -0.033, -0.036, -0.085, -0.279, and -0.279 e, respectively.<sup>185</sup>

The equilibrium distance of CO, CO<sub>2</sub>, CH<sub>2</sub>O, NO<sub>2</sub>, and SO<sub>2</sub> from the  $\beta$ -SnSe layer is 3.293, 3.617, 3.222, 2.531, and 2.692 Å, respectively, which are greater than the C-Sn (2.15 Å), O-Sn (2.22 Å), and N-Sn (2.11 Å) bonds, indicating the process to be physisorption.<sup>184</sup> However, the optimized structure shows some distortion for O<sub>2</sub> on the surface of the  $\beta$ -SnSe sheet. The value of  $E_a$  and  $Q_c$  for O<sub>2</sub> are -1.596 and -0.445 eV, respectively, which is much larger than other gas molecules adsorbed on one SnSe layer, indicating that O<sub>2</sub> molecules (CO, CO<sub>2</sub>, CH<sub>2</sub>O, NO<sub>2</sub>, and SO<sub>2</sub>) are chemically absorbed in the  $\beta$ -SnSe layer.<sup>185</sup> Adsorption

of CO, CO<sub>2</sub>, CH<sub>2</sub>O on the  $\beta$ -SnSe layer showed lower adsorption energy and lower charge transfer values, which indicates that one  $\beta$ -SnSe layer is not suitable for the detection of these three molecules.<sup>184</sup>

The higher absorption energy of NO<sub>2</sub> (-0.829 eV) and SO<sub>2</sub> (-0.499 eV) indicates that NO<sub>2</sub> and SO<sub>2</sub> molecules' absorption behavior for the  $\beta$ -SnSe layer was more potent than that of the absorption system.<sup>185</sup> The charge transfer values for NO<sub>2</sub> and SO<sub>2</sub> are -0.279 and -0.278 eV, respectively, which show a clear charge transfer between the gas molecule and the  $\beta$ -SnSe layer (Fig. 13(a)-(f)).<sup>185</sup> The adsorption distances for NO<sub>2</sub> and SO<sub>2</sub> are 2.531 Å and 2.692 Å, respectively, which is close to the Sn-O bond length range (2.22 Å to 2.66 Å).<sup>186</sup> For the  $\beta$ -SnSe sheet, it is crucial to detect NO<sub>2</sub> and SO<sub>2</sub> in the gas sensor region. All the most stable energy adsorption centers are located on the Sn side of the atom of the  $\beta$ -SnSe sheet, which indicates that the absorbing properties of the metal atom are more substantial than that of the nonmetallic atom.<sup>187</sup>

**4.4.1 SnSe based gas sensors.** Tin selenide (SnSe) is a member of (IVA-VIA) binary semiconductors family and has a high absorption coefficient.<sup>188</sup> The bandgap of SnSe varies from 0.9 eV to 1.3 eV.<sup>189</sup> SnSe has an orthorhombic crystal structure, and it has p-type conductivity with Sn-vacancies.<sup>190</sup> The majority of charge carriers of SnSe are holes and have a good impact on gas sensing applications. Using Ph<sub>3</sub>PSe as

Table 6 The adsorption energy, equilibrium distance, and Hirshfeld charge transfer of different molecules adsorb on a  $\beta$ -SnSe monolayer. Reprinted with permission from<sup>185</sup>

Molecule	$E_a$ (eV)	$d_0$ (Å)	$Q$ (e)
CO	-0.202	3.293 (C-Sn)	-0.033
CO <sub>2</sub>	-0.175	3.617 (C-Sn)	-0.036
CH <sub>2</sub> O	-0.322	3.222 (H-Sn)	-0.085
O <sub>2</sub>	-1.596	2.058 (O-Sn)	-0.445
NO <sub>2</sub>	-0.829	2.531 (O-Sn)	-0.279
SO <sub>2</sub>	-0.499	2.692 (O-Sn)	-0.279





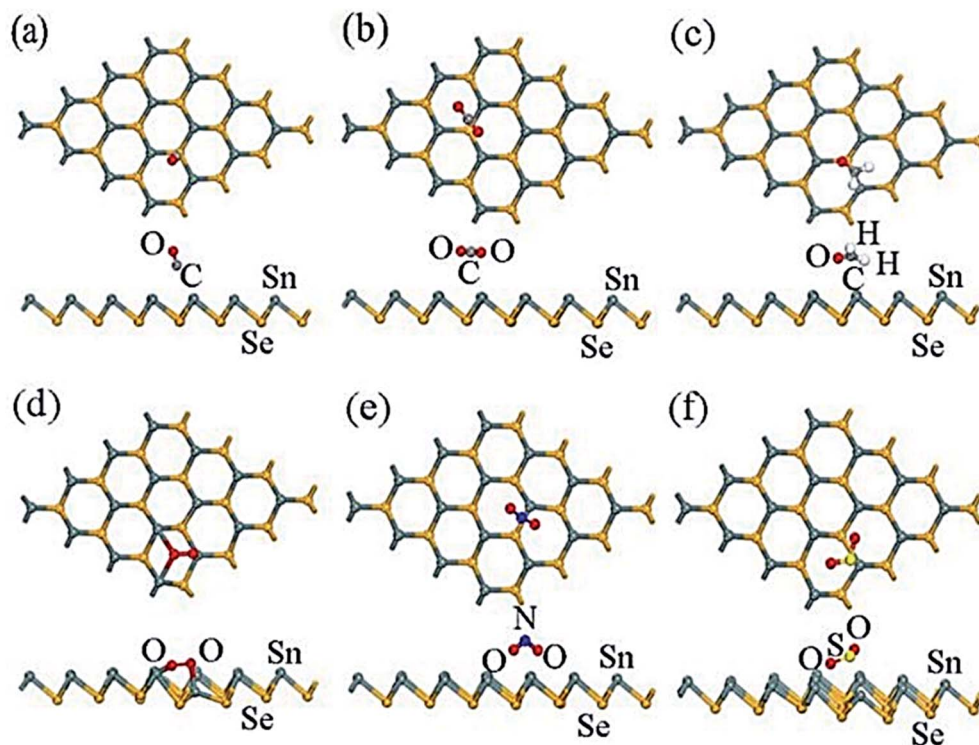


Fig. 13 The most stable sites of optimized configurations of the adsorbate molecules: (a) CO, (b) CO<sub>2</sub>, (c) CH<sub>2</sub>O, (d) O<sub>2</sub>, (e) NO<sub>2</sub>, and (f) SO<sub>2</sub> adsorbed on a  $\beta$ -SnSe monolayer. Most stable sites are exhibited. These figures have been adapted/reproduced<sup>185</sup> with permission from MDPI, copyright 2019\*.

a precursor in chemical vapor deposition, Assili *et al.*<sup>21</sup> deposited orthorhombic tin selenide thin films onto three substrates. A 1% vol concentration of methane gas showed a good sensitivity at the operating temperature of about 200 °C. The sensitivity, response, and recovery times were  $\sim 47\%$ ,  $\sim 52$  s, and  $\sim 220$  s, respectively. They observed that when SnSe was mixed with any n-type materials, it showed excellent response and recovery times at, lower operating temperature. In 2013, Wang *et al.*<sup>191</sup> synthesized SnO<sub>2</sub> nanoparticles decorated SnSe nanosheets *via* a facile, lost-cost, and safe solution method and had studied the gas sensing properties. For 1000 ppm CO gas, this device's response was 1.9 s at the operating temperature of 260 °C (Fig. 14(a)–(c)). Lee *et al.*<sup>192</sup> prepared the SnSe<sub>2</sub>/SnSe heterostructure film from a thick Sn layer onto the glass substrate. The film also showed response for one ppm NO<sub>2</sub> gas at room temperature, and the response value was 75%. In 2020, Wang *et al.*<sup>193</sup> reported a one-step colloid method for making SnSe/SnSe<sub>2</sub> heterostructures, with doping of SnSe  $\approx 30$ , 50, and 70%, respectively. The SnSe (50%)/SnSe<sub>2</sub> (50%) based sensor with an active layer thickness of 2  $\mu\text{m}$  showed the highest sensitivity of 30% to 0.1 ppm NO<sub>2</sub> gas at room temperature (25 °C) with a limit of detection (LOD) down to 69 ppb. The above discussion is summarized in Table 7.

**4.4.2 SnSe<sub>2</sub> based gas sensors.** SnSe<sub>2</sub> is an n-type material due to the presence of Se-vacancies. The majority of charge carriers in SnSe<sub>2</sub> are electrons. The SnSe<sub>2</sub> film is an anisotropic binary material having a hexagonal structure arranged in the form of Se–Sn–Se.<sup>194</sup> SnSe<sub>2</sub> has a small bandgap of 1–2 eV.<sup>30,195</sup> SnSe<sub>2</sub> has the same electronegativity and similar structure as

MoS<sub>2</sub>, thus it is better suited for the detection of NO<sub>2</sub>. 2D materials have a substantial surface to volume ratio, which significantly changes the absorption and release of gases, making them an inexpensive gas detection candidate.<sup>196</sup> Given all these properties of SnSe<sub>2</sub>, various researchers have tried to study the properties of SnSe<sub>2</sub> gas. Lee *et al.*<sup>192</sup> prepared the SnSe<sub>2</sub>/SnSe heterostructure film on the glass substrate from a thick Sn layer. At room temperature (RT), this film showed a better response for one ppm NO<sub>2</sub> gas. The response value of this film was 75%. In 2010, Popescu *et al.*<sup>197</sup> prepared silver doped SnSe<sub>2</sub> and Ge<sub>2</sub>Sb<sub>2</sub>Te<sub>5</sub> thin films by pulsed laser deposition. Ag-Doped SnSe<sub>2</sub> thin film-based gas sensor had a response of 3.26 for 500 ppm of CO gas at 500 °C. In 2018, Chen *et al.*<sup>194</sup> produced SnSe<sub>2</sub> nanoplate arrays and made a sensor that could detect 1% vol. of CH<sub>4</sub> at the operating temperature of 200 °C. This sensor showed a response of 66.7 and response/recovery times of 78/336 s, respectively. In 2019, Moreira *et al.*<sup>61</sup> investigated the NO<sub>2</sub> sensing properties of CVD deposited SnSe<sub>2</sub> binary layer. For one ppm of gas at the operating temperature of 50 °C, this sensor's response was 60%, and the response/recovery times of this sensor were 142/457 s (Fig. 14(d)–(f)). In 2017, Assili *et al.* demonstrated a SnSe<sub>2</sub> thin film-based gas sensor. This sensor had sensitivity around 16% for 200 ppm of methane at 200 °C with the response and recovery times around 75 s and 615 s, respectively.<sup>21</sup> In 2007, Popescu *et al.*<sup>198</sup> reported a SnSe<sub>2</sub> thin film-based gas sensor, which had a response of 300% for 8000 ppm of CH<sub>4</sub> gas at 600 °C. Sanju Rani *et al.* demonstrated a SnSe/SnSe<sub>2</sub> nanostructured thin film-based sensor by using of



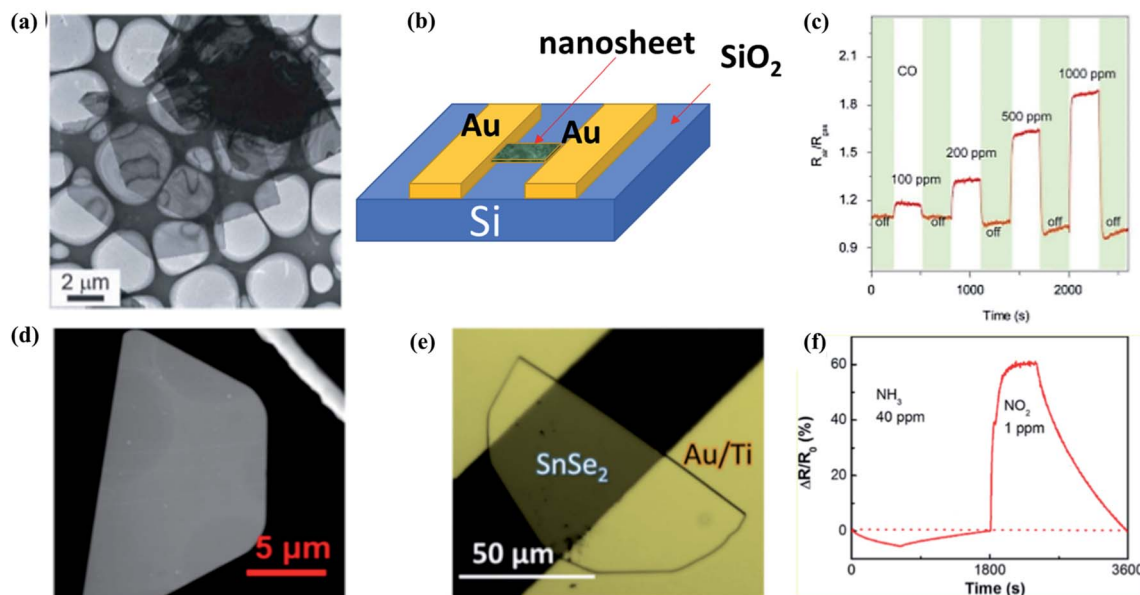


Fig. 14 (a) Low magnification TEM images, (b) schematic structure of the device, (c) transient response of the sensor SnO<sub>2</sub>/SnSe to CO (100–1000 ppm) at 260 μC. These figures has been adapted/reproduced from ref. 191 with permission from Nature, copyright 2013<sup>191</sup>. (d) SEM image of a semi hexagonal nanosheet of SnSe<sub>2</sub>, (e) optical image of a 6 nm thick SnSe<sub>2</sub> gas sensor device, (f) dynamic sensing responses of the 6 nm thick SnSe<sub>2</sub> resistor device measured with 405 nm laser illumination. These figures has been adapted/reproduced from ref. 61 with permission from ACS, copyright 2019<sup>61</sup>.

thermal evaporation technique. This sensor had response of 112% for (5 ppm) NO<sub>2</sub> gas at room temperature and the response/recovery times were 15/10 s, respectively.<sup>199</sup> AuPd coated SnSe<sub>2</sub> thin film based NO<sub>2</sub> sensor showed enhanced response of 117% for 5 ppm NO<sub>2</sub> gas at room temperature. The response/recovery times of this sensor were 10/18.7 s respectively.<sup>200</sup> Thus, efforts have been made to use SnSe<sub>2</sub> for sensing methane as well as NO<sub>2</sub>. The above discussion is summarized in Table 8.

## 4.5 Photocatalysis

**4.5.1 Photocatalytic behavior.** There are several reports on the photocatalytic behavior of Sn-based materials like SnSe and SnS. The possible mechanism reported by Li Cheng and co-workers is based upon the bandgap and absorption of light. The SnS and SnSe nanofibers have a narrow bandgap of 1.01 eV and 0.90 eV, respectively. UV light excites the photoelectron and hole in the material. Now the exciting photoelectron and hole react with dissolved oxygen molecules to form oxide radicals (O<sub>2</sub>). Furthermore, hydroperoxyl HO<sub>2</sub> and hydroxyl radical OH<sup>•</sup>

are formed after the protonation process (hydrogen cation). Simultaneously, the holes (h<sup>+</sup>) can oxidize hydroxyl radicals OH<sup>•</sup> and H<sub>2</sub>O molecules to generate OH<sup>•</sup> and hydroperoxyl HO<sub>2</sub>, which may break the RhB dye molecule to convert them into CO<sub>2</sub> and H<sub>2</sub>O or in other forms as shown in Fig. 15. The dye's degradation rate is defined as

$$D\% = \left[ \frac{A_0 - A_T}{A_0} \right] \times 100\% \quad (2)$$

Here A<sub>0</sub> is the absorbance of dye in the dark condition. A<sub>T</sub> is the absorbance of dye at specific time gaps after irradiation of light of a certain wavelength. Li Cheng *et al.* used the RhB solution under UV radiation to check the SnO<sub>2</sub>, SnSe, and SnS nano-fibers' photocatalytic activity.<sup>62</sup> The degradation curve showed a first-order kinetics equation type curve. The degradation rate for RhB was 85.90%, 92.55% and 92.86% for SnO<sub>2</sub>, SnS and SnSe, respectively. In another experiment, Jing group used thin films of Sn and SnSe to analyze the photodegradation

Table 7 The comparison between the gas sensing parameter of the different structures of SnSe

Material	Structure	Gas	Gas conc. (ppm)	Temp. (°C)	Response	Response time (s)	Recovery time (s)	Ref.
SnSe	Thin film	Methane	1% vol	200	~47% <sup>a</sup>	~52	~220	21
SnO <sub>2</sub> /SnSe	Nanosheet	CO	1000	260	1.9 <sup>b</sup>	—	—	191
SnSe <sub>2</sub> /SnSe	Heterostructure	NO	1	RT	75% <sup>c</sup>	—	—	192
p-SnSe/n-SnSe <sub>2</sub>	Hetero-structure	NO <sub>2</sub>	0.1	RT	30% <sup>b</sup>	—	—	193

<sup>a</sup>  $R = (R_g - R_a)/R_a \times 100$ . <sup>b</sup>  $R = R_a/R_g$ . <sup>c</sup>  $I = (I_g - I_0)/I_0$ . R is called the sensor's response, S is called the sensor's sensitivity, R<sub>g</sub> is the resistance of the sensor in the presence of target gas, and R<sub>a</sub> is the sensor's resistance in the air. I<sub>g</sub> is the current of device in presence of target gas and I<sub>0</sub> is the current of device in presence of air.



Table 8 The comparison between the gas sensing parameter of the different structures of SnSe<sub>2</sub>

Material	Structure	Gas	Gas concentration (ppm)	Temp. (°C)	Response	Response time (s)	Recovery time (s)	Ref.
SnSe <sub>2</sub> /SnSe	Heterostructure	NO <sub>2</sub>	1	RT	75% <sup>a</sup>	—	—	192
SnSe <sub>2</sub> -Ag	Thin film	CO	500	500	3.26 <sup>b</sup>	—	—	197
SnSe <sub>2</sub>	Nano plate	Methane	1% vol	200	66.7 <sup>c</sup>	78	336	194
SnSe <sub>2</sub>	Pristine monolayer	NO <sub>2</sub>	1	50	60% <sup>b</sup>	142	457	61
SnSe <sub>2</sub>	Thin-film	Methane	200	200	~16% <sup>c</sup>	~75	~615	21
p-SnSe/n-SnSe <sub>2</sub>	Hetero-structure	NO <sub>2</sub>	0.1	RT	30% <sup>d</sup>	—	—	193
SnSe <sub>2</sub>	Thin film	CH <sub>4</sub>	8000	600	300% <sup>d</sup>	—	—	198
SnSe-SnSe <sub>2</sub>	Nanostructured thin film	NO <sub>2</sub>	5	RT	112% <sup>e</sup>	15	10	199
Au/Pd/SnSe <sub>2</sub>	Thin film	NO <sub>2</sub>	5	RT	117% <sup>e</sup>	10	18.7	200

<sup>a</sup>  $S = (I_g - I_0)/I_0 \times 100$ . <sup>b</sup>  $R = R_a/R_g$ . <sup>c</sup>  $R = (R_a - R_g)/R_g \times 100$ . <sup>d</sup>  $S = (R_g - R_a)/R_a \times 100$ . <sup>e</sup>  $R = (R_g/R_a) \times 100$ , where  $I_g$  is the current of device in presence of target gas and  $I_0$  is the current of device in presence of air.  $R$  is called the sensor's response,  $S$  is called the sensor's sensitivity,  $R_g$  is the resistance of the sensor in the presence of target gas, and  $R_a$  is the sensor's resistance in the air.

of RhB.<sup>14</sup> Two films SnSe<sub>2</sub> and SnSe<sub>2</sub>/Se were made by the facile solvothermal method. The photocatalytic efficiency was measured from the formula mentioned below,

$$\text{Photocatalytic degradation percentage} = \left[ \left( 1 - \frac{C}{C_0} \right) \times 100 \right] \quad (3)$$

where  $C$  is the concentration at time  $t$  and  $C_0$  is the concentration at time  $t$ . The result showed that SnSe<sub>2</sub>/Se film showed better degradation, i.e., around 94% compared to Sn films in 50 min. This increase in efficiency was attributed to the formation of heterojunction and the production of a larger number of electron-hole pairs and a lower recombination rate.

**4.5.2 SnSe composites for photocatalytic applications.** Zhou Li *et al.* synthesized a composite of SnSe@SnO<sub>2</sub> nanoparticles and used them in the photothermal-photocatalyst mechanism study. The mechanism of electron-hole pair generation is similar to that mentioned earlier.

The photon falls on the surface, and the electron-hole pair is generated. Here, both SnSe and SnO<sub>2</sub> took part in the photothermal and photocatalytic processes synchronously. In the process, the electrons from CB of SnSe travels to the CB of SnO<sub>2</sub>.

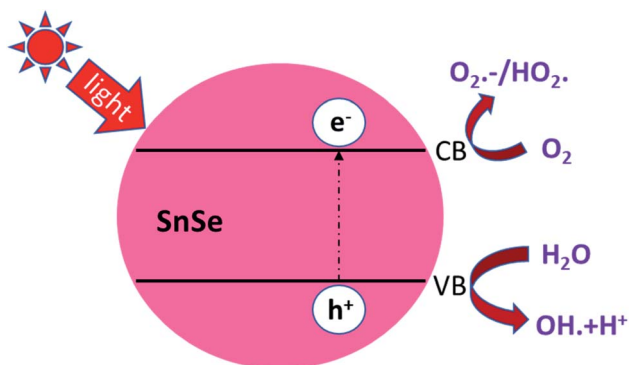


Fig. 15 Indicative representation of photodegradation of RhB dye in the presence of UV light. In a typical experiment, the various dyes' degradation rate is checked using eqn (3). These figures have been adapted/reproduced from ref. 62 with permission from Scielo, copyright 2017<sup>14</sup>.

This increases the oxidation process of photocatalysis and thus improves the photocatalytic behavior. The inner shell SnSe favors the photothermal heat efficiently.<sup>201</sup> The photocatalytic activity of SnSe@SnO<sub>2</sub> nanocomposite was evaluated through the degradation of methyl orange (MO) experiment under solar light's irradiation. The photocatalytic degradation of MO with SnSe, SnO<sub>2</sub>, SnSe@SnO<sub>2</sub>, and P25 (commercially named as TiO<sub>2</sub> nanoparticles) was compared using the time *versus*  $C/C_0$  graph. Here,  $C$  is the concentration at time  $t$  and  $C_0$  is the concentration at time  $t$ . It was observed that SnSe@SnO<sub>2</sub> showed a much higher degradation rate than SnSe, SnO<sub>2</sub>, and TiO<sub>2</sub>. In some other methods, blue SE-2R dye degradation was studied by making a composite.

Karamat *et al.*<sup>202</sup> made SnSe composite with LaNdZr<sub>2</sub>O<sub>7</sub> and compared the photocatalytic efficiency with SnSe and blank dye. The setup composed a 300 W Xe arc lamp (PLS-SXE300, Beijing Trusttech Co. Ltd). The degradation of blue dye was analyzed at 658 nm using a UV-Vis spectrophotometer. Before the experiment, the stability of blue dye was checked under dark and light conditions. It was confirmed that the dye was stable, and the continuous illumination does not show any loss in spectra. The photocatalytic degradation experiment was then carried out, and it was found that LaNdZr<sub>2</sub>O<sub>7</sub>/SnSe showed a better degradation efficiency. The maximum absorption intensity at wavelength 658 nm showed a decrease in absorption as the irradiation time increases. SnSe alone has demonstrated degradation efficiency of just 15.5% while mixing it with LaNdZr<sub>2</sub>O<sub>7</sub> increased it significantly.

## 4.6 Storage devices

**4.6.1 Battery electrode.** Tin-based alloys showed high theoretical capacity and are being considered as a very suitable material for sodium-ion batteries. Tin-based composites have a theoretical storage capacity of nearly 847 mA h g<sup>-1</sup>. Still, a few other composites like SnSe and SnSe<sub>2</sub> showed even higher values than theoretical ones, with additional capacity contributions from conversion reactions.<sup>14</sup> Considering that one molar SnSe<sub>2</sub> can accommodate 7.75 molar sodium, SnSe<sub>2</sub> has shown a theoretical capacity of 756 mA h g<sup>-1</sup>.<sup>203</sup> But during the alloying process, the volume of the material expands, which is dangerous for the cycle



stability; hence some composites of carbon are used to reduce this bottleneck. Tin selenide-based battery application has been divided into two categories, namely  $\text{Li}^-$  ion and  $\text{Na}^+$  ion batteries. The recent advances in  $\text{SnSe}_2$  for sodium-ion batteries and their results are explained in the coming section.

Zhang *et al.* synthesized  $\text{SnSe}_2$  based two-dimensional (2D) nanosheets using the hydrothermal technique. They achieved a theoretical capacity during the first cycle and a stable and reversible specific capacity of  $515 \text{ mA h g}^{-1}$  at  $0.1 \text{ A g}^{-1}$  after 100 cycles, which exhibited excellent performance.<sup>204</sup> The  $\text{SnSe}$  and  $\text{SnSe}_2$  both have been used for both sodium and lithium-ion battery applications. Chen *et al.* (2018) introduced Cu in  $\text{SnSe}$  for sodium-ion battery and exhibited a capacity of  $330 \text{ mA h g}^{-1}$  at  $20 \text{ A g}^{-1}$ . In a study, Xia *et al.* used electrospun  $\text{SnSe}$  with carbon nanofibers for lithium and sodium-ion battery applications and showed improved results. Kim *et al.* investigated  $\text{SnSe}$  alloy as an anode for Na-ion batteries. They exhibited excellent electrochemical performance with a high reversible capacity of  $707 \text{ mA h g}^{-1}$  and stable performance over 50 cycles.<sup>204</sup> In recent time, various studies on  $\text{SnSe}$  composite has been conducted by multiple research groups to find better results due to the layered structure of  $\text{SnSe}$  and  $\text{SnSe}_2$  materials. Table 9 shows a summary of the above discussions.

Tin selenide-based composites have shown an excellent theoretical capacity. After many cycles, real practical uses of tin-

based anodes in both types of batteries are still very limited. Some structural design, preparation advancement, and morphological development are needed to bring these materials to the production level.

**4.6.2 Supercapacitor.** Supercapacitors have gained colossal consideration due to their longer operational life, higher power densities, and better safety tolerances than batteries. Zhang *et al.* synthesized tin selenide ( $\text{SnSe}$ ,  $\text{SnSe}_2$ )-based 2D nanostructures on flexible substrates for supercapacitor applications.<sup>63</sup> Ni *et al.* synthesized  $\text{SnSe}$  based anode using a microwave-assisted method. They exhibited a suitable specific capacitance of  $214.3 \text{ F g}^{-1}$  at  $1 \text{ A g}^{-1}$  and rate capability of  $182.8 \text{ F g}^{-1}$  at  $20 \text{ A g}^{-1}$  with outstanding cyclic stability.<sup>211</sup> Pandit *et al.* synthesized binder-free  $\text{SnSe}$  hexagonal nanosheets using a one-pot colloidal method and achieved high-performing anode material for supercapacitors. It showed a specific capacitance of  $617.9 \text{ F g}^{-1}$  at a scan rate of  $2 \text{ mV s}^{-1}$  and with fast charge-discharge cycles.<sup>15</sup>

#### 4.7 Memory devices

The phase change from amorphous to crystalline induces a sharp change in the optical and electrical properties like reflectivity and resistivity, respectively. This change in the

**Table 9** The tin selenide-based composite used for battery application and their cyclic performance

Material used	Ion transfer	Cycle rate	Ref.
Cu doped $\text{SnSe}$	Sodium	Retains a capacity of $330 \text{ mA h g}^{-1}$ at $20 \text{ A g}^{-1}$ and $304 \text{ mA h g}^{-1}$ after 1000 cycles at $5 \text{ A g}^{-1}$ ( $0.1\text{--}3.0 \text{ V vs. Na/Na}^+$ )	205
Electrospun $\text{SnSe/C}$ nanofibers	Lithium and sodium	The discharge capacity of $405 \text{ mA h g}^{-1}$ at $1000 \text{ mA g}^{-1}$ after 500 cycles in lithium-ion battery and $290 \text{ mA h g}^{-1}$ at $200 \text{ mA g}^{-1}$ after 200 cycles in $\text{Na}^+$ ion battery	60
Tin selenide/N-doped carbon composite	Lithium and sodium	For Li ion-discharge capacity $405 \text{ mA h g}^{-1}$ after 500 cycles at a current density of $1000 \text{ mA g}^{-1}$ for Na ion-discharge capacity of $290 \text{ mA h g}^{-1}$ after 200 cycles at $200 \text{ mA g}^{-1}$	206
$\text{SnSe/SnO}_2$ heterostructure	Lithium	High cyclability having a capacity of $810 \text{ mA h g}^{-1}$ after 200 cycles at a current density value of current density of $400 \text{ mA g}^{-1}$	207
Tin diselenide hexagonal nanosheets	Lithium	Specific capacity of $795 \text{ mA h g}^{-1}$ after 100 cycles at $100 \text{ mA g}^{-1}$	208
$\text{SnSe}_2$ nanoplate-graphene composites	Lithium	Higher storage capacity of $420 \text{ mA h g}^{-1}$ in the first 10 cycles	55
$\text{SnSe}_2/\text{CNTs}$ hybrid nanostructures	Lithium	$\text{SnSe}_2/\text{CNTs}$ electrode showed specific capacity of $457.6 \text{ mA h g}^{-1}$ at $0.1\text{C}$ and $210.3 \text{ mA h g}^{-1}$ after 100 cycles	209
$\text{SnSe}_2/\text{reduced graphene oxide (rGO)}$ composite	Sodium	Exhibits an initial efficiency of 73.7%, showing a high capacity of $402.0 \text{ mA h g}^{-1}$ after 150 cycles at $0.1 \text{ A g}^{-1}$ with retention of 86.2%	210



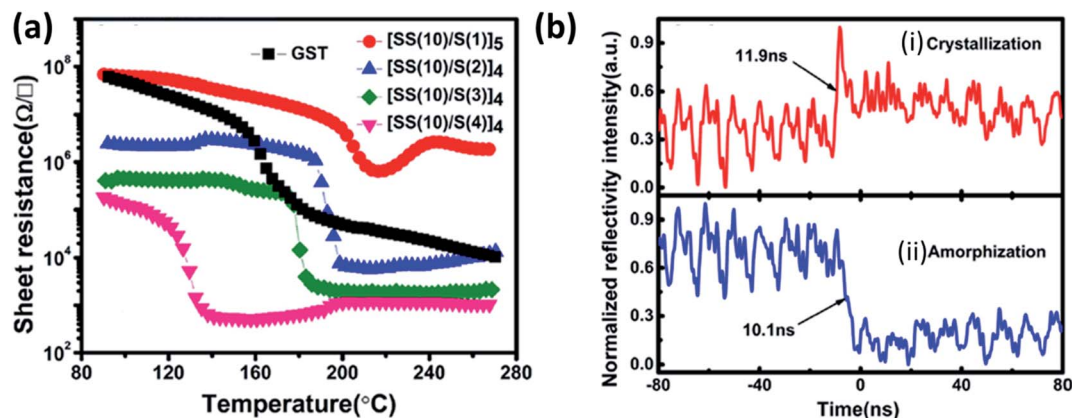


Fig. 16 (a) The temperature dependence of resistance for the SLL SS/S and GST thin films at a constant heating rate of 10 °C min<sup>-1</sup>, and (b) reversible reflectivity evolution of SLL [SS(10)/S(2)]<sub>4</sub> thin film induced by two consecutive picosecond laser pulses with different fluencies: (i) crystallization process and (ii) amorphization process. These figures have been adapted/reproduced from ref. 215 with permission from AIP, copyright 2016".

property is explored to store the data in the form of a bit. As this transition is reversible, this property is utilized in the rewritable data storage. Fast operating speeds, low power consumption, high energy density, large retention time, *etc.*, are the critical parameters for selecting materials for memory applications. Chalcogenides based materials play the central role in this field due to their very short time of the order of few nanoseconds of phase transition, high resistance ratio of the amorphous to the crystalline phase. Chung *et al.* investigated all phases of the tin selenide (SnSe, SnSe<sub>2</sub>, Sn<sub>2</sub>Se<sub>3</sub>)<sup>7</sup> for memory application based on phase change and compared the result with the well-studied GeTe.<sup>212</sup> SnSe<sub>2</sub> is best among all phases, with high crystallization temperature (220 °C), high electrical contrast (resistance ratio of amorphous to crystalline) of  $8.3 \times 10^5$ ,  $T_g/T_m$  of 0.52, which showed its potential for the memory application.<sup>212</sup> Wang *et al.* demonstrated the swift recrystallization time of 20 ns for the solution-processed SnSe<sub>2</sub> thin film,<sup>213</sup> comparable to the commercialized Ge<sub>2</sub>Sb<sub>2</sub>Te<sub>5</sub>. Sun *et al.* studied the multilayer Si/SnSe<sub>2</sub> film by varying the Si's thickness in steps of 4 nm for 0–20 nm, with the film's total thickness nearly constant at 100 nm.<sup>214</sup> The varying thickness of the Si did not show any effect on the crystallization temperature. Still, a remarkable change in the resistance of the film with Si thickness of 12 nm {[Si (12 nm)/SnSe<sub>2</sub> (5 nm)]<sub>6</sub>} was observed, *i.e.*, crystalline resistance increased and amorphous resistance decreased, which led to enough (more than three orders of magnitude) electrical contrast to storing information. Simultaneously increasing crystalline resistance reduced the operating current that led to lower power consumption. Multilayer with Si

thickness of 12 nm (of nearly 100 nm total thickness) showed its high potential for the low power operated phase-change random access memory (PCRAM) device because of almost ten year retention time at 83 °C and low activation energy.<sup>214</sup>

The difficulty in phase-change memory (PCM) application of these materials is due to high amorphous resistivity of SnSe<sub>2</sub>, fast crystallization speed, low thermal stability, retention of Sb rich alloy, and very low phase change temperature of Sb. Optimized design of the super lattice like (SLL) structure of SnSe<sub>2</sub> and Sb ([SS(10)/S(2)]<sub>4</sub>) (SS stands for SnSe<sub>2</sub> and S for Sb with 10 and 2 nm thickness, respectively) proved to be four times better than the commercialized Ge<sub>2</sub>Sb<sub>2</sub>Te<sub>5</sub> (GST)<sup>215</sup> SLL's resistance with varying temperature and phase switching times for optimized SLL is shown in Fig. 16(a) and (b). Table 10 shows the comparison of the supremacy of SLL over GST.

#### 4.8 Topological insulators

Topological insulators are the materials that have insulating bulk accompanied by the protected conducting surface states. These surface states provide an impurity and defect insensitive scattering path to the electrons. The unusual metallic nature of an insulator's surface due to the changes in the insulator's topology is the signature of a topological insulator.<sup>216</sup> These surface states are induced due to the high spin-orbit coupling, inverse parity of bulk band, band inversion,<sup>217,218</sup> and protected by the time-reversal symmetry.<sup>16,218</sup> When the crystal symmetry protects the topological insulator's surface, instead of the time-reversal symmetry protecting it, these insulators are called the topological crystalline insulator (TCI).<sup>217</sup> Sun *et al.*<sup>16</sup> reported

Table 10 The supremacy of SLL over GST

Structure	Crystallization temperature (°C)	Crystallization activation energy (eV)	Year data retention at temperature (°C)	Density variation	SET/RESET operation speed (ns)
Ge <sub>2</sub> Sb <sub>2</sub> Te <sub>5</sub>	160	2.37	10 (75)	8.2	17.7/16.5
[SS(10)/S(2)] <sub>4</sub>	185	3.03	10 (116)	3.5	11.9/10.1



that SnSe is a TCI in its native phase (without any pressure and alloying) based on the theoretical calculation using the first principle along with *ab initio* tight-binding modeling. Chen *et al.*<sup>219</sup> reported the co-existence of topological states and superconducting phase at around 27 GPa of the CsCl type phase of SnSe. *Ab initio* calculation showed that this CsCl phase of SnSe has the topological states' stability against the spin-orbit coupling. SnSe provides the correlation among the superconductivity and topological electronic structures.<sup>219</sup>

## 5. Conclusion

This review summarizes the tin-selenide material's physical properties, phases, defects, growth mechanisms, deposition methods, and various possible applications. Tin-based compound semiconductors have attracted enormous interest in photovoltaic and optoelectronic applications due to favorable bandgap (1.0–1.3 eV) and higher absorption coefficient ( $10^5 \text{ cm}^{-1}$ ). The reported efficiency of tin selenide based solar cells is very low compared to other compound semiconductor materials like CIGS, CdTe, and CZTSSe. The main limiting factor for tin selenide's relative low efficiency might be defects present in absorber materials, interface trap states, and unfavorable hetero-junction interfaces. Moreover, conduction band offset (CBO) between the absorber and buffer layer plays a crucial role. It needs to be engineered in a precise manner for efficient transport of electrons toward a metal electrode. Thus, fine-tuning of elements, precise control over growth conditions, favorable band alignment between absorber and buffer can pave the way towards improving the power conversion efficiency for realizing the development of environmental-friendly low-cost next-generation photovoltaic technology. Tin selenide can be a potential futuristic material due to its versatile applications. It shows leading hands in some applications in its native crystal structure like thermoelectric. Whereas in other fields it shows competitive nature with the other materials. Tin selenide based thermoelectric materials can pave the ways to reuse waste energy. Tin selenide-based photodetector has also exhibited fast response speed with good responsivity and detectivity in the NIR range.

Moreover, flexible photodetector based on tin selenide have also gained attention due to its easy fabrication and fast device response. The device response can be improved by controlling absorber thickness and carrier collection before recombination. Another application based on tin selenide is photocathode in photoelectrochemical water splitting for green hydrogen generation. Lastly, some studies have shown that tin selenide can also be used in gas sensors, storage devices (battery electrode, supercapacitors), memory devices, topological insulators.

## Conflicts of interest

Authors declare no conflict of interest exists.

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