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Tin-selenide as a futuristic material: properties and applications

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SnSe/SnSe₂ is a promising versatile material with applications in various fields like solar cells, photodetectors, memory devices, lithium and sodium-ion batteries, gas sensing, photocatalysis, supercapacitors, topological insulators, resistive switching devices due to its optimal band gap. In this review, all possible applications of SnSe/SnSe₂ have been summarized. Some of the basic properties, as well as synthesis techniques have also been outlined. This review will help the researcher to understand the properties and possible applications of tin selenide-based materials. Thus, this will help in advancing the field of tin selenide-based materials for next generation technology.

1. Introduction

The fast developing area of applied material science demands materials to be cheap, non-toxic, environment-friendly, easy to synthesize, and well competitive in performing particular applications. Nowadays, materials with versatility have gained massive attention due to their applicability in almost all fields.^{1,2} Various materials have been explored and showed promising versatile applications, *e.g.*, graphene (Gr),¹ TiO₂,³ ZnO,⁴ Cu₂SnS₃ (CTS),² *etc.* The multifunctional applicability of these materials paves the foundation for interdisciplinary research. Chalcogenide-based materials have also shown such potential and can be seen as the future hope to meet a similar requirement. Among chalcogenides, tin selenide has demonstrated great potential in the applied material science. Tin selenide exists in two phases, *i.e.*, SnSe and SnSe₂. Some researchers have observed another phase, Sn₂Se₃,^{5,6} but this phase is the superposition of SnSe and SnSe₂.⁷ Tin selenide has demonstrated versatility in thermoelectric,⁸ photodetector,⁹ solar cells,¹⁰ photocatalytic,¹¹ phase change memory,¹² gas sensing,¹³ anode material for battery,¹⁴ supercapacitor,¹⁵ and topological insulator (TI).¹⁶ These applications strongly depend upon the properties of SnSe (optical, electrical and microstructural, *etc.*). Apart from material properties, material synthesis/deposition methods also play an essential role in obtaining high-quality materials.^{17–21} Excellent review article

exists on thermoelectric materials consisting of fundamental properties to the thermoelectric device's final design, growth, defects, working environment issues, and applications.²² Other reviews that focusses on SnSe describes all the aspects mentioned above (like growth, defects, configuration, *etc.*).²³ The aim of this review is to summarize the ongoing progress on SnSe, SnSe₂ synthesis methods, materials properties, and its possible application in various fields. However, some studies focus on pristine tin-selenide phase and its application^{8,24} Another good reviews^{8,25,26} that concentrates on SnSe describing all aspects mentioned above (like growth, defects, design, *etc.*), provides all-round knowledge to researchers. This review gives insights into the phases, structures, synthesis methods, progress in the tin-selenide, and tin diselenide for various applications. Hence, this article aims to review the tin selenide devices, present status, recent progress in the growth process, related-issues, possible solutions, and their possibility of developing next-generation technology.

2. Material properties of tin selenide

Tin selenide based SnSe/SnSe₂ are binary compound semiconductor materials having p/n-type conductivity.^{20,27} Fig. 1(a) shows a salient feature of tin selenide materials. The SnSe exists in two crystallographic phases, *i.e.*, orthorhombic (α -SnSe)²⁸ and cubic (π -SnSe),²⁹ and SnSe₂³⁰ exists in the hexagonal crystal structure. The π -SnSe material is mechanically stable and energetically comparable to α -SnSe and has already shown potential in the piezoelectric application.²⁹ The orthorhombic crystal structure of SnSe showed second-order displacive type phase transition above 750 K, from α -SnSe (space group *Pnma*) to β -SnSe (space group *Cmcm*) with ($a = 11.49 \text{ \AA}$, $b = 4.44 \text{ \AA}$, $c = 4.135 \text{ \AA}$) to the ($a = 4.31 \text{ \AA}$, $b = 11.70 \text{ \AA}$, $c = 4.31 \text{ \AA}$), respectively (Fig. 1(b)).^{31,32} The crystal structure of SnSe (violet, Sn atoms;

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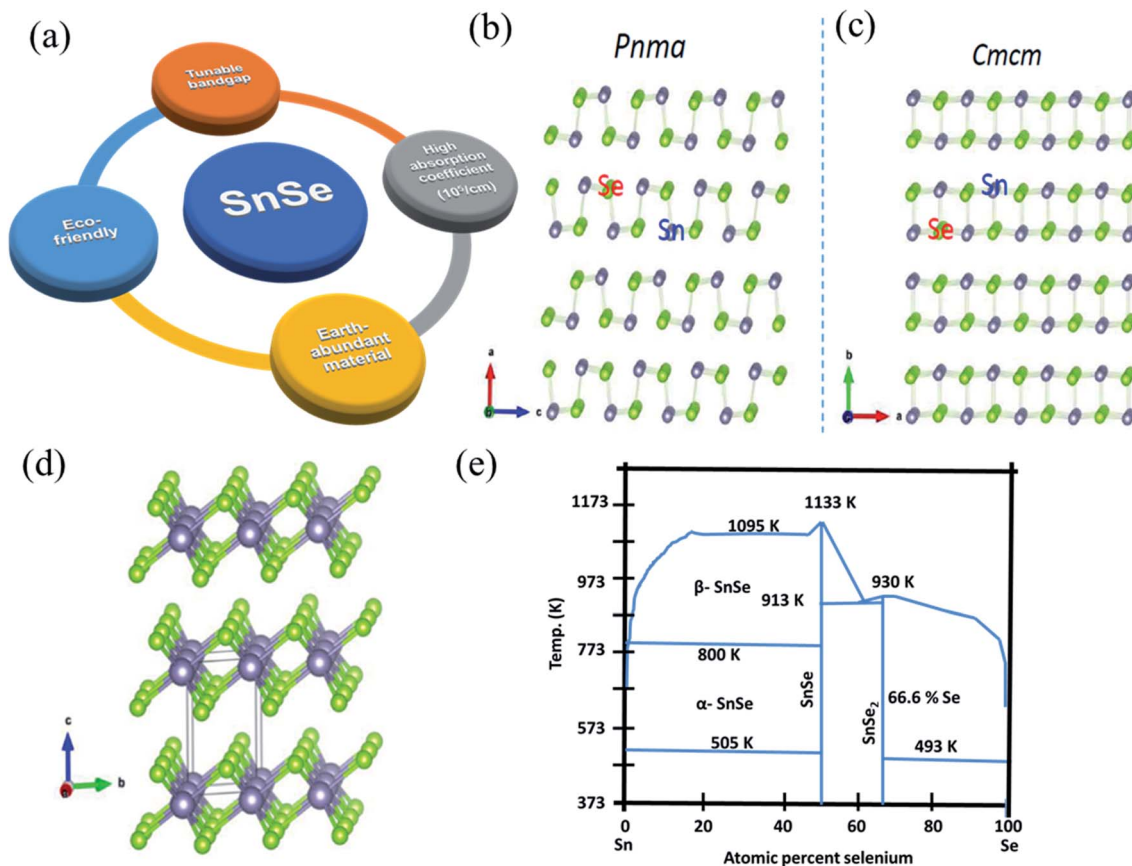


Fig. 1 (a) Salient feature of tin selenide materials. (b) Crystal structure of SnSe, (c) SnSe having *Pnma* to *Cmcm* phase transition. This figure has been adapted/reproduced from ref. 38 with permission from Elsevier, copyright 2018". (d) Crystal structure of SnSe₂. This figure has been adapted/reproduced from ref. 33 with permission from ACS, copyright 2016". (e) Equilibrium phase diagram of Sn–Se system. This figure has been adapted/reproduced from ref. 39 with permission from Wiley, copyright 2020".

green Se atoms) (Fig. 1(c)),³² and crystal structures of SnSe₂ (green, Se atoms; violet, Sn atoms) is shown in Fig. 1(d).³³ Both the tin-selenide phases show the indirect and direct band gaps. The direct band gap of α -SnSe shows a wide tunable band gap which varies from 0.98 eV (bulk) to 1.43 eV (monolayer).³⁴ SnSe₂ shows large variation in the band gap from 1.84 eV (bulk) to 2.04 eV (monolayer).³⁵ This band gap tunability of tin selenide shows its immense application possibilities in optoelectronic device applications.^{30,36} Tin selenide exists in two stoichiometric phases, *i.e.*, SnSe and SnSe₂, as shown by the equilibrium phase diagram (Fig. 1(e)).⁷ Various research groups have reported another phase, Sn₂Se₃,³⁷ but this phase is the superposition of the SnSe and SnSe₂ as confirmed by Nuclear Magnetic Resonance (NMR) spectroscopy.

Defects present in the material also influences the properties of materials like electronic, magnetic, and optical properties. SnSe is generally a p-type semiconductor. Density functional theory (DFT) calculation is performed to study the defects during SnSe crystal growth. The Sn vacancy is present as a native defect, which causes the p-type conduction in the SnSe under Se or Sn rich conditions, as confirmed by Scanning Tunneling Microscopy (STM) studies.^{40,41} SnSe₂ shows n-type conduction. DFT calculation shows that the vacancy of selenium and interstitial tin led to this n-type conduction in SnSe₂.⁴¹ The

optoelectronic and physical properties of SnSe/SnSe₂ materials depend on the growth conditions, size, morphology, phase purity, growth techniques, *etc.* Therefore, selection of synthesis/deposition method impacts material as well as its proposed applications.

3. Synthesis of SnSe and SnSe₂

The material's properties also depend on the growth conditions,⁴² size, and morphology of the material,⁴³ phase purity,⁴³ defects,⁴⁴ *etc.* Various physical and chemical methods have been established to fabricate SnSe/SnSe₂ materials on different substrates (glass, flexible metal foil, polymer, *etc.*) for multiple applications. The tin selenide based semiconductor materials have been synthesized using atomic layer deposition (ALD),¹⁷ sputtering,⁴⁵ thermal evaporation,⁴⁶ hydrothermal,⁴⁷ spray pyrolysis,⁴⁸ chemical vapor deposition (CVD),⁴⁹ *etc.*

Bulk crystals of the SnSe and SnSe₂ can be easily obtained by the physical solid-state reaction method. >99.99 pure elemental powders of the Sn and Se (in the stoichiometric ratio) is taken and sealed in the quartz tube at the pressure of $\sim 10^{-4}$ torr. It is placed at the temperature according to the Sn–Se system phase diagram, as shown in Fig. 2(a). The Bridgman method can be used to obtain single crystals. The powder was melted into the



furnace and brought in the contacts with the single crystal's seed during its cooling. The material grows along the crystallographic orientation of the seed crystal. Depending upon the geometry it is called the horizontal Bridgman method or vertical Bridgman method. The Bridgman–Stockbarger technique as shown in (Fig. 2(a)) can be used to obtain a single crystal of SnSe

and SnSe₂ with an optimized ampule lowering rate of 7 mm h⁻¹ and 4 mm h⁻¹, respectively, for SnSe and SnSe₂ with an optimized temperature gradient of 15 °C cm⁻¹.⁵⁰ Another technique to grow single crystals are the direct vapor transport method,⁴² and temperature gradient method.⁵¹ A detailed review of the growth of the SnSe single crystals can be found in an earlier study.²⁵ Various

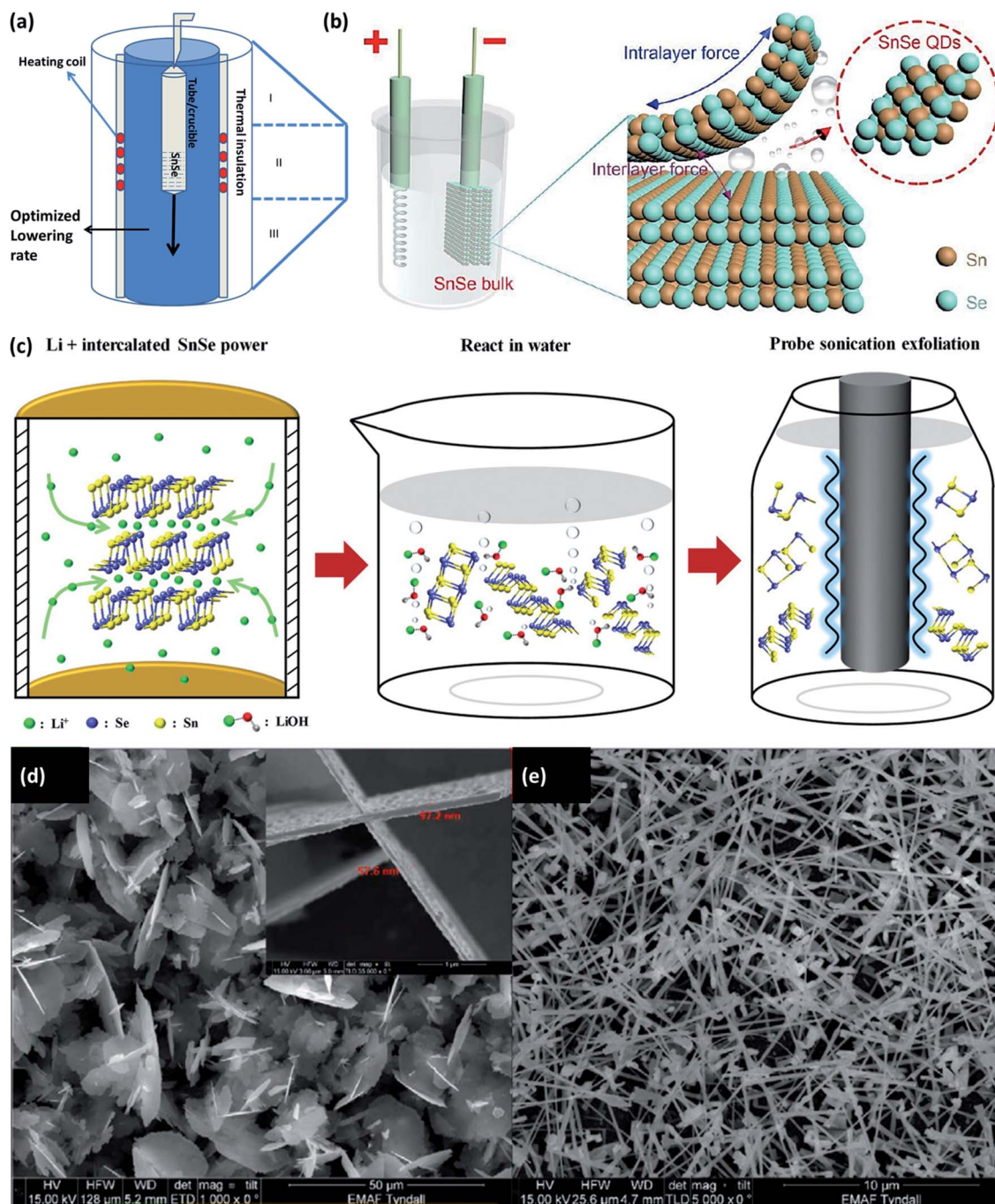


Fig. 2 (a) Bridgman–Stockbarger technique to grow a single crystal. (b) Schematic illustration for the synthesis of SnSe QDs by cathodic exfoliation. This figure has been adapted/reproduced from ref. 52 with permission from RSC, copyright 2019⁵². (c) Schematic synthesis process for the SnSe nanosheets by the liquid-phase exfoliation method includes two main steps: Li⁺ hydrothermal intercalation and sonication-assisted exfoliation. This figure has been adapted/reproduced from ref. 54 with permission from Wiley, copyright 2020⁵⁴. SEM images showing the evolution of the SnSe_x nanostructures concerning the growth temperature, (d) large flakes with diameters of several microns grown at 450 °C, with a thickness of ~60 nm shown in the inset, (e) high yield of nanowire growth at 500 °C. This figure has been adapted/reproduced from ref. 49 with permission from Wiley, copyright 2020⁴⁹.



methods to grow the nanostructure of the SnSe and SnSe₂ is described in this section.

3.1 Synthesis of nanostructured SnSe

Li *et al.*⁵² devised the facile cathodic exfoliation method to exfoliate SnSe bulk into quantum dots with high yield. Under the organic electrolyte containing 0.2 M of tetrabutylammonium (at -7.5 V for 30 min) bulk SnSe was exfoliated and downsized up to ~ 10 nm in lateral size, which was further downsized to ~ 4 nm by breaking weakly entangled dots with the help of sonication as shown in Fig. 2(b). To synthesize the SnSe flake up to one layer, Jiang *et al.*⁵³ reported the two-step process, in which first the bulk SnSe flakes were obtained by the atmospheric pressure vapor transport deposition method. Flash evaporation of SnSe powder was done by moving the quartz tube position containing SnSe powder at 700 °C under Ar/H₂ gas pressure under atmospheric condition. SnSe is deposited on the Si substrate kept in upside-down position. As obtained flakes were etched with the N₂ gas. Etching time of 5–20 min gave the single layer SnSe flake (thickness of 6.8 Å).⁵³ The SnSe nano-sheet of the best quality can be produced from the SnSe bulk powder by the three-step method.⁵⁴ The method comprises the intercalation of Li⁺ ion into the SnSe layers by lithification and then reacting the material with water. The rapid expansion of H₂ gas during reaction with water exfoliated the SnSe nano-sheets. To improve the nanosheets' yield, further sonification of the SnSe powders' residuals is carried out (shown in Fig. 2(c)).

Using the precursor diselenoether SnCl₄[BuSe(CH₂)₃-SeⁿSeBu]}, SnSe_x nanostructures are formed on the Si(100) substrate coated with Au (acted as a catalyst) by the liquid injection chemical vapor deposition method, and growth was studied in the

temperature range 450 – 550 °C. Under the 1.1 sccm flow of Ar gas and 1.5 mL h⁻¹ of precursor injection rate, SnSe₂ flake growth occurred at 450 °C, and considerable growth of SnSe nanowires occurred at 500 °C as shown in Fig. 2(d) and (e).⁴⁹

3.2 Synthesis of nanostructured SnSe₂

Choi *et al.*⁵⁵ synthesized the nano-plates of the SnSe₂ by dissolving 5 mL oleylamine in 50 mg, 0.26 mmol SnCl₂ at 220 °C.

After that, a solution containing 1,3-dimethylimidazoline-2-selenone (90 mg, 0.53 mmol), dichloromethane (4 mL), and oleylamine (2 mL) were mixed and heated for 2 h at this temperature. After that, it was cooled to room temperature, centrifuged, and the residue was washed using hexane to obtain SnSe₂ nanoplates. The transmission electron microscopy (TEM) image of SnSe₂ nanoplates are shown in Fig. 3(a)–(d).

A simple two-step method, *i.e.*, sonication followed by laser ablation, was used by Li *et al.* to obtain SnSe₂ quantum dots (QD) from bulk powder.⁵⁶ Manually grinded SnSe₂ powder (20 mg) was dispersed in 30 mL deionized water, and the solution was sonicated (650 W) for 2 h in an ice bath for 4 s. Sonication resulted in smaller particles and flakes of SnSe₂. The obtained solution kept in quartz cuvette was irradiated with a 1064 nm laser (Nd:YAG, 2.2 W) for 10 min. Irradiated tiny particles were centrifuged for 30 min at 6000 rpm, resulting in SnSe₂ QD, as shown in Fig. 4(a) and (b).

Ma *et al.*⁵⁸ synthesized the nanorods and nanoplates of the SnSe₂ by organic solution phase route. Two different solutions of 0.4 mmol SnCl₂·2H₂O, 2 mmol oleic acid, 8.5 mL liquid paraffin oil, and 0.4 mmol selenium powder, 9 mL liquid paraffin were used. The first solution was heated at 160 °C till lemon yellow color appeared, and the later solution was heated

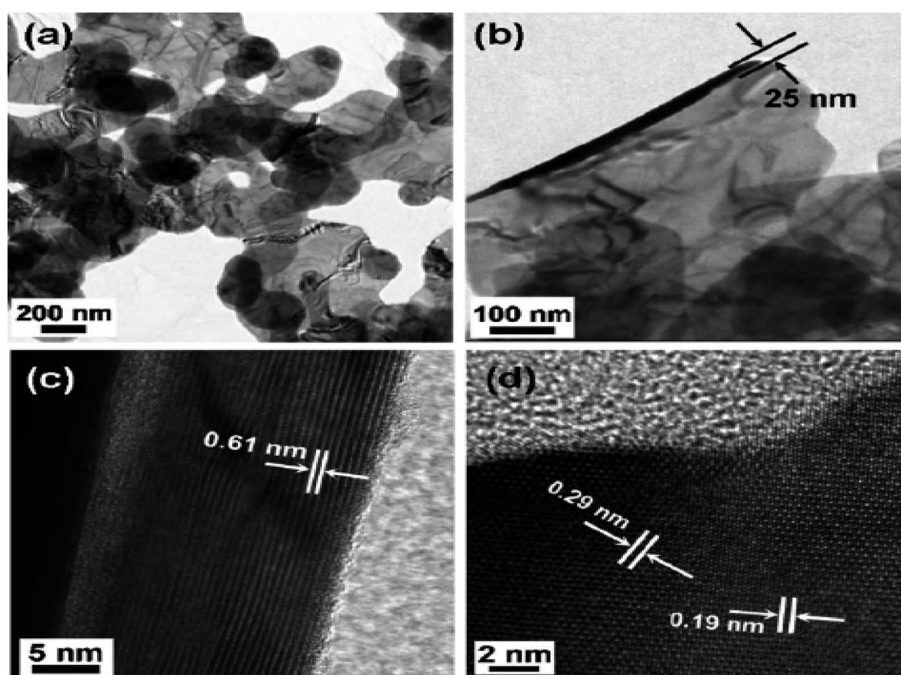


Fig. 3 TEM image (a) and (b) SnSe₂ nanoplates and high resolution TEM (HRTEM) images (c) side, and (d) top views. This figure has been adapted/reproduced from ref. 55 with permission from RSC, copyright 2011*.



at 240 °C till wine color appeared. The latter solution was injected into first using syringe rapidly and the mixture was heated at 200 °C for 20 min. After cooling, toluene and methanol was added. After centrifugation nanorods and nanoplates were obtained. Chemical vapor deposition was employed by Zhou *et al.* to synthesize ultrathin (1.5 nm) SnSe₂ flakes of high quality.⁵⁷ Low melting point precursor SnI₂ (0.01 g) and Se (0.1 g) powders were placed in two different alumina crucible, and mica substrate were placed ~15 cm from the central zone. The central zone was heated at 600 °C for 15 min, with a flow of H₂ (5 sccm) and Ar (20 sccm) at ambient pressure. As-synthesized flakes are shown in Fig. 4(c)–(g).

4. Various applications

In recent times, tin selenide has gained significant research attention for solar cells, thermoelectric, photodetectors, photocatalytic, phase change memory, gas sensing, and anode material for battery, supercapacitor, topological insulator *etc.* Various reported applications based on SnSe materials are displayed in Fig. 5.

4.1 Thin film solar cells (TFSCs)

Fabrication of low-cost thin-film solar cells is highly needed as they can also be deposited over flexible substrates. The CuInGaSe₂ (CIGSe) and CdTe based solar cells have shown record small cell power conversion efficiency (PCE) of 23.35%

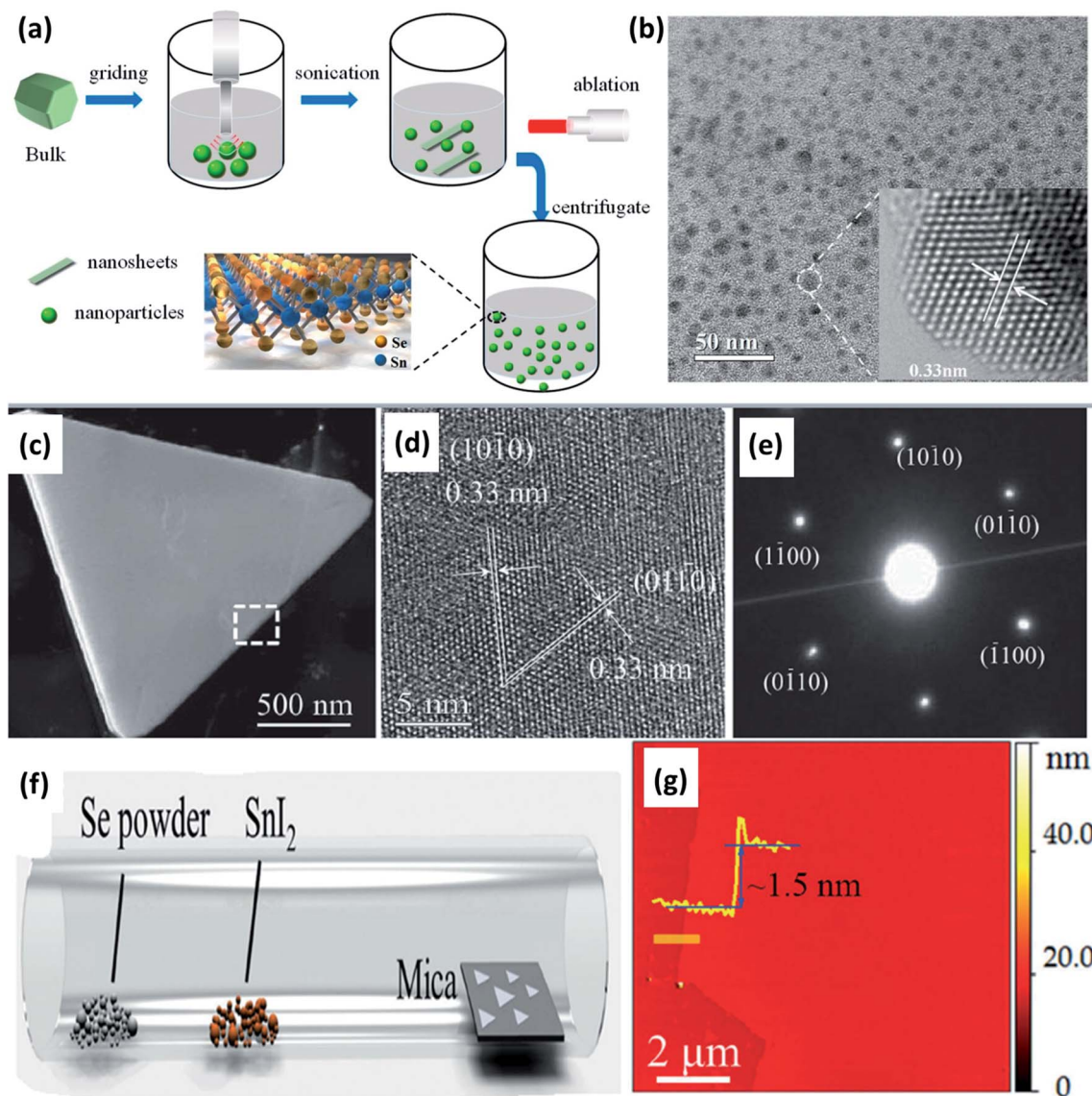


Fig. 4 (a) Schematic show of the SnSe₂ structure and the quantum dot (QD) fabrication process, and (b) TEM image of SnSe₂ QDs with a centrifugal speed of 6000 rpm. These figures has been adapted/reproduced from ref. 56 with permission from MDPI, copyright 2019". (c) Low-magnification TEM image of a SnSe₂ flake, (d) corresponding HRTEM image of the flake, (e) electron diffraction pattern from the same flake, (f) schematic diagram of the chemical vapor deposition method, and (g) a typical atomic force microscope (AFM) image at the flake edge, and the height profile showing a thickness of ~1.5 nm. These figures has been adapted/reproduced from ref. 57 with permission from Wiley, copyright 2015".



and 22.1%, respectively.⁶⁶ But in CIGS TFSC, indium is rare, and gallium is very costly. In CdTe TFSC, Cd is toxic, and Te is a rare material. These materials related problems hinder the commercialization of low-cost, earth-abundant TFSC based on CIGS and CdTe. In recent time, earth-abundant and low-cost compound semiconductor materials like $\text{Cu}_2\text{ZnSnS}_4$ (CZTS), $\text{Cu}_2\text{ZnSnSe}_4$ (CZTSe), and $\text{Cu}_2\text{ZnSn}(\text{S,Se})_4$ (CZTSSe) have gained massive consideration as alternate materials instead of CIGS and CdTe as they have shown the efficiency of 10%, 11.95%, and 12.62%, respectively.^{67–69} But, kesterite material consists of many elements.⁷⁰ The presence of many elements in the absorber layer increases the processing cost, and defects are also formed during the processing, which ultimately reduces the overall PCE of solar cells. To overcome this problem, binary element-based materials with similar optical and electronic properties are more suited than quaternary element-based kesterite materials. Tin sulfide (SnSe) is also attracting attention of researchers for optoelectronic device applications. In SnSe, only one impurity phase of SnSe_2 is present.⁷¹ Therefore, the formation of phase SnSe is easier. Also, being consists of

earth-abundant, inexpensive, and eco-friendly elements, SnSe has attracted significant attention. The SnSe has an optimum bandgap of 1.1 to 1.3 eV, an absorption coefficient of $\sim 10^5 \text{ cm}^{-2}$, and p-type conductivity with high carrier concentration (10^{17} cm^{-3}), making it a fabulous material for solar photovoltaic applications.⁸ Due to its optimum bandgap (1.3 eV) and high absorption coefficient (10^5 cm^{-1}), a thin layer of 300 nm thickness can absorb most of the useful solar spectrum. The tunable bandgap of thin-film SnSe can help absorb more solar radiation and enable more photon absorption, resulting in more electron-hole pair generation. The theoretical efficiency for material with a bandgap of 1.3 eV is 32%.⁷² The SnSe thin film can be deposited using various chemical and physical routes like chemical bath deposition (CBD),⁷³ electrodeposition,⁷⁴ spray pyrolysis,⁷⁵ atomic layer deposition (ALD),¹⁷ thermal evaporation,⁴⁶ sputtering.^{76,77}

The device architecture for SnSe TFSC is SLG/Mo/SnSe/CdS/i-ZnO/TCO/metal grid, similar to CIGS,⁷⁸ and CZTS/CZTSe/CZTSSe TFSCs.⁷⁹ Therefore, a lot of modifications and optimizations are required. Also, many issues need to be mitigated to

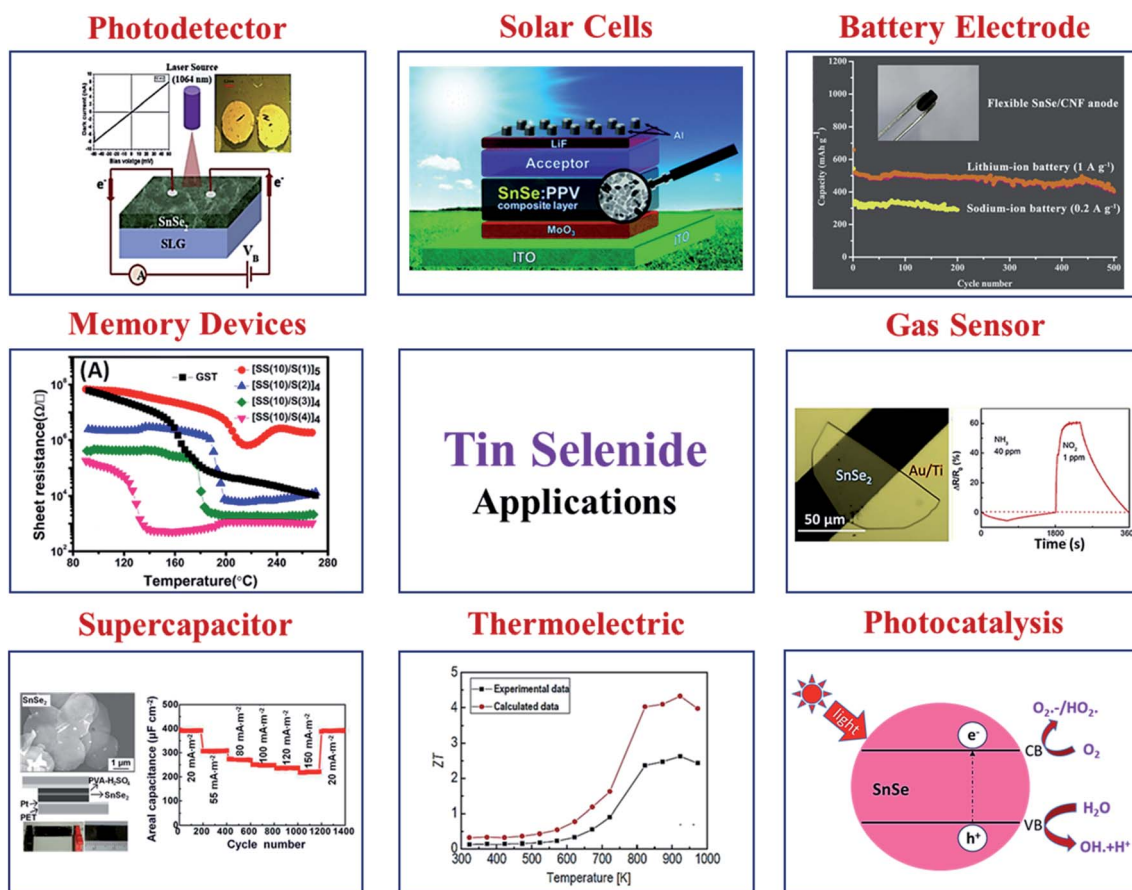


Fig. 5 Various applications are based on SnSe materials. Solar cell, this figure has been adapted/reproduced from ref. 59 with permission from ACS, copyright 2010". Battery electrode, this figure has been adapted/reproduced from ref. 60 with permission from Elsevier, copyright 2020". Gas sensor, this figure has been adapted/reproduced from ref. 61 with permission from ACS, copyright 2019". Photocatalysis, this figure has been adapted/reproduced from ref. 62 with permission from Scielo, copyright 2017". Thermoelectric, this figure has been adapted/reproduced from ref. 38 with permission from Elsevier, copyright 2018". Supercapacitor, this figure has been adapted/reproduced from ref. 63 with permission from ACS, copyright 2014". Memory devices, this figure has been adapted/reproduced from ref. 64 with permission from AIP, copyright 2014". And photodetector, this figure has been adapted/reproduced from ref. 65 with permission from Elsevier, copyright 2020".



improve efficiency.⁸⁰ Consequently, it is essential to conduct a study on SnSe material to develop environmentally-friendly, low-cost TFSCs. Any photovoltaic materials should have good optical properties, appropriate band gap, good absorption coefficient, high carrier concentration, and better transport properties. Solar cell efficiency may get affected due to losses, *i.e.*, optical losses, non-absorption, thermalization, reflection loss, transmission loss, area loss, collection losses, and resistance losses.⁸¹ However, the formation of pure phase SnSe thin films has not been reported as the phase lies in a very narrow region, and fine compositional tuning is required to achieve SnSe pure phase. Sn and Se's reaction mechanism plays a crucial role in attaining pure phase SnSe thin films.^{19,82} Reddy *et al.* studied the effect of selenization temperature, elemental composition, and selenization pressure on SnSe. They observed that the single-phase SnSe thin films could be achieved during the selenization process in the temperature ranges from 300 to 500 °C.⁷⁷ The solar cell efficiency depends on various parameters, like deposition condition, the absorber's crystallinity, carrier concentration, and the p-n junction's nature. The lower solar cell efficiency based on SnSe thin films may be due to SnSe absorber materials' low quality. The parameters that highly influence SnSe thin film quality are phase purity, higher crystallinity, and larger grain size with fewer pinholes. It has been observed that during SnSe phase formation, there is a possibility of the formation of the SnSe₂ phase. Thus, the SnSe₂ phase's impact on the solar cell performance of SnSe thin films needs to be carefully studied and optimized.⁸ The electronic properties of p-n junction interfaces are strongly influenced by the discontinuities between the valence band (VB) maxima and conduction band (CB) minima of each material, which restrict the electron transport across the junction interfaces. Band alignment of SnSe with an n-type buffer layer, the density of defect states at the hetero-junction interface, and the nature of back contact also influence the solar cell performance. Therefore, band offsets play a crucial role in reducing the dark current in a diode, reducing the photon-generated carrier losses, and improving the overall solar cell conversion efficiency. The n-type buffer layer with a wider band gap is needed to overcome the optical losses and provide better alignment with SnSe absorber material.

In 1990, Singh *et al.* reported SnSe based solar cells with power conversion efficiency (PCE) of 2.3%.⁷⁵ The fabricated solar cell showed an open-circuit voltage (V_{oc}) of 410 mV, short-circuit current density (J_{sc}) of 9.20 mA cm⁻², and a fill factor (FF) of 49%, respectively. Rahman *et al.* fabricated heterojunction (p-SnSe/n-Si) solar cells and achieved an efficiency of about 6.44%.⁸³ They observed improved power conversion efficiency due to the improvement in the J_{sc} with a graded junction.⁸³ Shinde *et al.* fabricated SnSe thin-film solar cells using electro-deposition and exhibited device efficiency of 1.4% using an absorber film with 300–400 nm-sized grains, uniform, and dense film morphology.⁷³ Franzman *et al.* synthesized SnSe/PPV (poly[2-methoxy-5-(3',7'-dimethyl octyl oxy)-1,4-phenylenevinylene]) nanocrystals for solar cells and observed improvement in efficiency from 0.03% to 0.06%.

They observed significant improvement in external quantum efficiency (EQE) and J_{sc} after adding PPV into the SnSe absorber. Fig. 6(a) and (b) show the high-resolution transmission electron microscopy (HRTEM) image of a single nanocrystal, and selected area electron diffraction (SAED) pattern of SnSe. Fig. 6(c) shows a low-resolution TEM image of SnSe nanocrystals, and Fig. 6(d) shows the schematic device structure of SnSe solar cells.⁵⁹ In 2012, Mathews *et al.* reported the first inorganic thin-film solar cell and exhibited a device efficiency of 0.03% ($V_{oc} = 140$ mV, $J_{sc} = 0.7$ mA cm⁻²). The low device efficiency is mainly due to poor junction quality.⁷⁴ Makori *et al.*⁸⁴ fabricated solar cell with device structure glass/Ag/CdO:Sn/SnSe/Ag by using thermal evaporation and reported device efficiency of 0.59% ($V_{oc} = 273$ mV, $J_{sc} = 0.993$ mA, and FF = 69%). The observed lower efficiency is mainly due to the thin absorber (148 nm) layer, which attributed to lower J_{sc} in the device. The above discussion is summarized in Table 1.

Table 1 shows that the efficiency of SnSe based solar cells is well below Cu₂ZnSn(S,Se)₄.⁸⁵ The observed lower efficiency is mainly due to bulk defects, secondary phase formation in SnSe, non-ideal band-alignment at heterojunction interfaces, and back and front electrode recombination. The absorber quality plays a significant role, and a sound absorber should have a dense, pinhole-free, smooth surface of the film for high-efficiency heterojunction device applications. The low

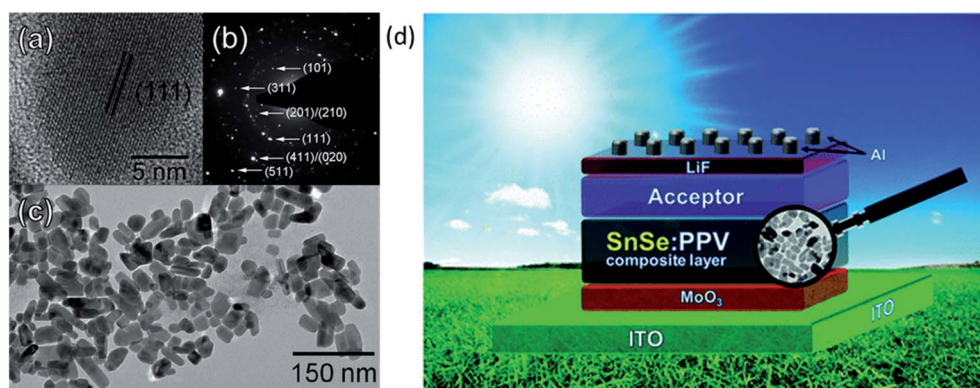


Fig. 6 (a) HRTEM image of a single nanocrystal. (b) Selected area electron diffraction (SAED) pattern of SnSe. (c) Low-resolution TEM image of SnSe nanocrystals. (d) The device structure of SnSe solar cells. These figures has been adapted/reproduced from ref. 59 with permission from ACS, copyright 2010".

Table 1 Summary of SnSe based thin-film solar cells

Deposition method	Absorber	Device architecture	V_{oc} (mV)	J_{sc} (mA cm ⁻²)	FF (%)	η (%)	Area (cm ²)	Ref.
Evaporation	SnSe	Si/Al/SnSe/In	425	17.23	44	6.44	0.2	83
Evaporation	SnSe	SLG/FTO/Se/SnSe/Ag	410	9.20	49	2.30	—	75
Electro-deposition	SnSe	ITO/CdS/SnSe/Au	370	5.37	30	1.40	—	73
Co-evaporation	SnSe	SLG/Mo/SnSe/CdS/ i-ZnO/ITO/Ni-Al	172	18.87	31.2	1.02	0.42	10
Sputtering	SnSe	Glass/Ag/CdO:Sn/SnSe/Ag	207	0.90	0.69	0.59	—	84
Electrodeposition	SnSe	Tec15/CdS/SnSe/carbon-paste	140	0.7	31	0.03	0.15	74

interfacial quality could primarily affect the device performance and is responsible for lower V_{oc} and FF in the solar cells. This is mainly due to CB discontinuity/band offsets at the interfaces, which leads to the CB's energy barrier and restricts carrier transportation across the p-n junction. Apart from this, absorber thickness, grain size, bandgap, and deposition methods also influence the device performance.

4.2 Thermoelectric generators

It requires energy to maintain the sustainability and development of humankind. Energy sources are minimal, and there arises a question about the world's power crisis and its ultimate solution. A thermoelectric device that can convert the waste heat into electricity can contribute to a remarkable extent in this.

4.2.1 SnSe thermoelectric

SnSe single crystal. With the report of a high ZT value of 2.6 along the b axis in single-crystal SnSe,³² this material became the hot topic of research in the field of thermoelectric (TE) domains since 2014. SnSe is anisotropic and exhibits the second-order displacive phase transition from $Pnma$ (distorted) to $Cmcm$ (non-distorted) at 750–800 K. It shows a moderate power factor of $10.1 \mu\text{W cm}^{-1} \text{K}^{-2}$ at ~ 850 K along the b axis, which is comparable or even lesser than typical thermoelectric materials values.^{86–89} The thermal conductivity of SnSe has astonished the researchers. Using the Gruneisen parameter, it has been estimated that SnSe exhibits the ultra-low thermal conductivity ($<0.25 \text{ W m}^{-1} \text{K}^{-1}$) due to the strong anharmonicity in bonding.³² Theoretical studies^{90,91} showed that optimized carrier concentration lies in the range 10^{19} – 10^{20} cm^{-3} for SnSe for high ZT . For n-type SnSe, the estimated ZT is 3.1 along a axis at 770 K at carrier concentration of $2.8 \times 10^{19} \text{ cm}^{-3}$, while the study showed that the n-type SnSe performed better than p-type SnSe.⁹²

Motivated by the high ZT along the b axis in SnSe, Li *et al.*³⁸ conducted the theoretical study on the p-type SnSe using the first-principle calculation. A figure of merit (ZT) was greatly affected by the carrier concentration and also sensitive to the phases of SnSe ($Pnma$ and $Cmcm$) (shown in Fig. 7(a) and (b)). Fig. 7(c) shows a comparison of theoretical and experimental data.³² This study showed that a maximum ZT of 4.33 could be achieved for the $Cmcm$ phase of SnSe at 923 K for carrier concentration of $1.84 \times 10^{19} \text{ cm}^{-3}$, along the b axis.

Through another study, for Na doped SnSe single crystal, Zhao *et al.* reported the device with $ZT_{\text{device}} = 1.34$ in the temperature range 300–773 K. With ZT_{max} of 2.0 at 773 K, the device showed the conversion efficiency of $\sim 16.7\%$, which was higher than many Pb-based thermo-electric modules.⁹³ Various researchers showed different ZT values for the single crystal. For example, Wei *et al.*⁹⁴ showed ZT of ~ 1 , ~ 0.8 , and ~ 0.25 along b , c , and a axis, respectively for in-house made fully dense single-crystalline SnSe, Jin *et al.*⁹⁵ showed $ZT \sim 1$ for single crystal grown by a vertical vapor deposition method. Different types of doping are being done to enhance the single-crystal performance. Chang *et al.*⁹⁶ doped bromine to make the SnSe conduction n-type and achieved record high $ZT \sim 2.8 \pm 0.5$ at 773 K along out of the plane direction. With the help of density functional theory and scanning tunneling microscopy, they pointed out that delocalized p electrons of Sn and Se near conduction band minima contribute more to orbital overlapping in out of the plane direction. The S doping in single-crystal SnSe decreased the carrier concentration due to charge trapping at grain interface.⁹⁷ Hence, S doping decreased the single crystal's thermoelectric performance rather than enhancing as predicted by theoretical calculation.⁹⁸ Pb doped n-type single crystal (*via facile* Sn-flux method) showed a 33% enhanced power factor than its pristine SnSe due to increased carrier concentration.⁹⁹ Ag-doped single crystal ($\text{Sn}_{0.97}\text{Ag}_{0.03}\text{Se}$) grown in a horizontal Bridgman furnace showed maximum $ZT \sim 0.95$ along a axis at 793 K.¹⁰⁰ Bi doping in single-crystal SnSe resulted in n-type conduction and showed the remarkable $ZT \sim 2.2$ along the b axis at 733 K.⁵¹

Polycrystalline SnSe thermoelectric. Polycrystalline SnSe has comparably higher thermal conductivity and lower electrical conductivity than its single crystal counterpart. Though phonon scattering is reduced mainly by the grain boundaries,¹⁰¹ it has higher thermal conductivity than its single crystal due to tin oxide layer¹⁰² or absorption of oxygen.¹⁰³ Highly discrete values of ZT have been reported for polycrystal SnSe made by different growth techniques. Various researchers have tried to achieve a comparable result to a single-crystal. Many researchers chose to dope (Na, K, Cu, Zn, *etc.*) in poly-crystal SnSe to enhance the electrical properties and reduce the thermal conductivity effectively. In one study, the polycrystalline SnSe thermoelectric material showed $ZT \sim 0.5$ at 823 K.¹⁰⁴ Several methods improve the device's thermoelectric performance, like large mass fluctuations, band gap engineering, alloying, doping, nano-structuring, *etc.*¹⁰⁵ Gong *et al.*⁴⁷ studied the effect of Cu doping



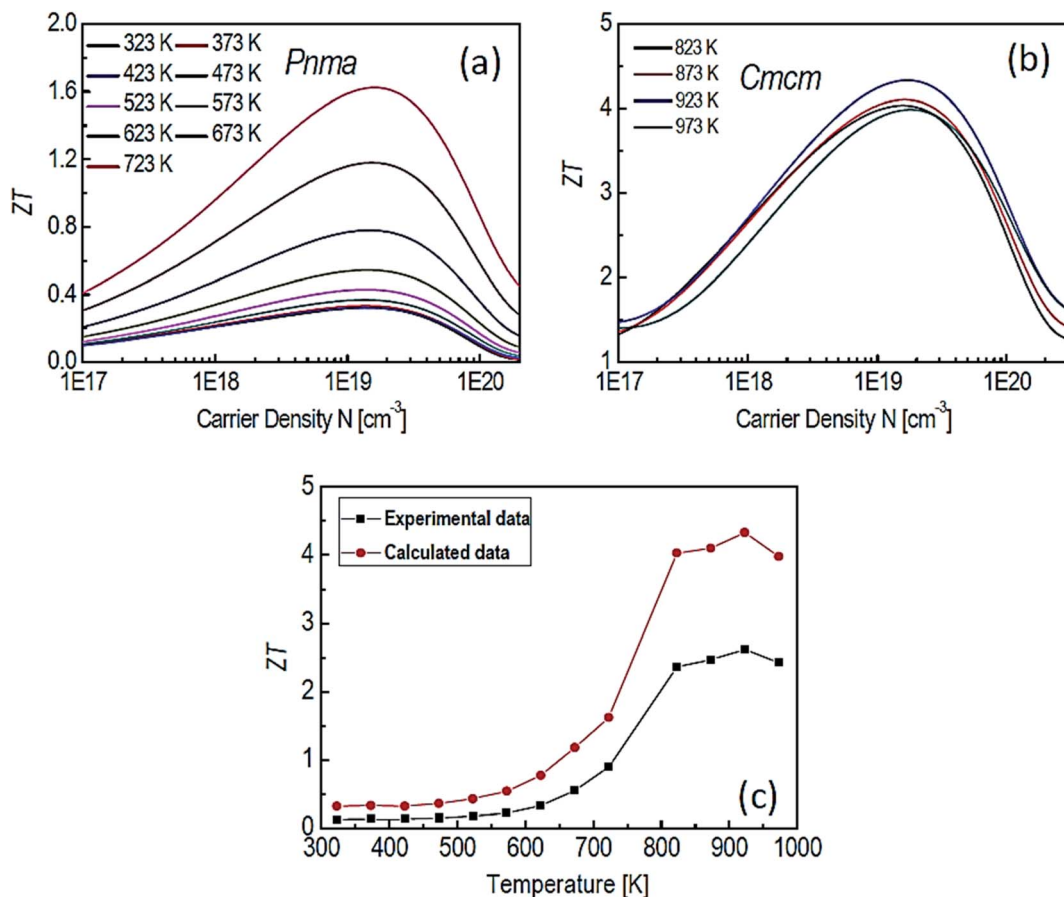


Fig. 7 The ZT of SnSe for *Pnma* (a), *Cmcm* (b) phases and comparison between experimental data³² and calculated data (c). These figures have been adapted/reproduced³⁸ with permission from Elsevier, copyright 2018³⁸.

in polycrystalline SnSe by varying the concentrations in the $\text{Cu}_x\text{Sn}_{1-x}\text{Se}$ from $x = 0.01$ to 0.004. They found that Cu doping enhanced the electrical conductivity due to the increased carrier concentration, as confirmed by Hall measurement. Still, Cu doping acted as a point defect and hence decreased the carrier mobility. The Seebeck coefficient also increased due to doping. Thus, the overall power factor increased due to doping, and the optimized value was obtained for doping $x = 0.01$. Around 60% reductions in the lattice thermal conductivity as compared to undoped poly-crystal SnSe is observed. This was due to nanoprecipitation and mesoscale grains (as evidenced by high-resolution transmission electron microscopy (HRTEM)). Finally, they achieved a ZT of 1.2 at 873 K. Shi *et al.*¹⁰⁶ enhanced the Cu's solubility limit in the SnSe and obtained the ZT of ~ 1.41 for the doping level of $x = 0.118$.

Lee *et al.*¹⁰² synthesized polycrystalline SnSe ($\text{Na}_{0.01}(\text{Sn}_{0.95}\text{Pb}_{0.05})_{0.99}\text{Se}$), which exhibited $ZT \sim 2.5$ at 773 K. They noticed that the oxide formation leads to higher thermal conductivity in poly-crystal SnSe (SnO_2 has ~ 140 times higher thermal conductivity than SnSe).¹⁰⁷ They removed the oxides layers of the tin in SnSe by Ball milling followed by a reduction in 4% H_2/Ar atmosphere for 6 h at 613 K. They pointed out that ball milling (BM) and reduction processes both are critically important to remove the oxides and its residuals from the sample. A remarkably ultralow

thermal conductivity (total) of $0.20 \text{ W m}^{-1} \text{ K}^{-1}$ (even lower than the single crystal reported by Zhao *et al.*³²) (Fig. 8(a)) and ~ 2 times enhanced power factor than the pristine SnSe (unball-milled and unreduced). The performance of pristine, reduced, ball-milled and reduced (BR) compared to with the single-crystal SnSe³² and Na doped SnSe,⁹³ is shown in Fig. 8(b). Recently Gainza *et al.*¹⁰⁸ reported the highest ZT of 1.8 for the un-doped polycrystalline SnSe, and even with surface oxides layers on the sample.

The pressure-induced effect modifies the textures of the polycrystal SnSe. A sintering pressure of 60 MPa leads to improved electrical properties. There was a meager improvement in thermal conductivity, resulting in a ZT of 0.7.¹⁰⁹ There is a vast literature on doping in the polycrystalline SnSe, like Na,¹¹⁰ K,¹¹¹ Cd,¹¹² Ag,¹¹³ and co-doping Na/Ag,^{114,115} Bi/Cl,¹¹⁶ Na/CNT,¹¹⁷ *etc.* which resulted in the enhancement in the thermoelectric performance of the material. Lu *et al.*¹¹⁸ enhanced the performance of polycrystalline SnSe by introducing large mass fluctuations by doping sulfur, which led to very low lattice thermal conductivity ($0.13 \text{ W m}^{-1} \text{ K}^{-1}$ at 873 K). Still, this doping enhanced the bandgap of the SnSe, which lowered the electrical conduction and hence the power factor (low carrier concentration due to large bandgap). Pb co-doping improved the power factor, further increasing the carrier concentration by one order of magnitude. Through co-doping a remarkable ZT of 1.85 at 873 K was achieved.



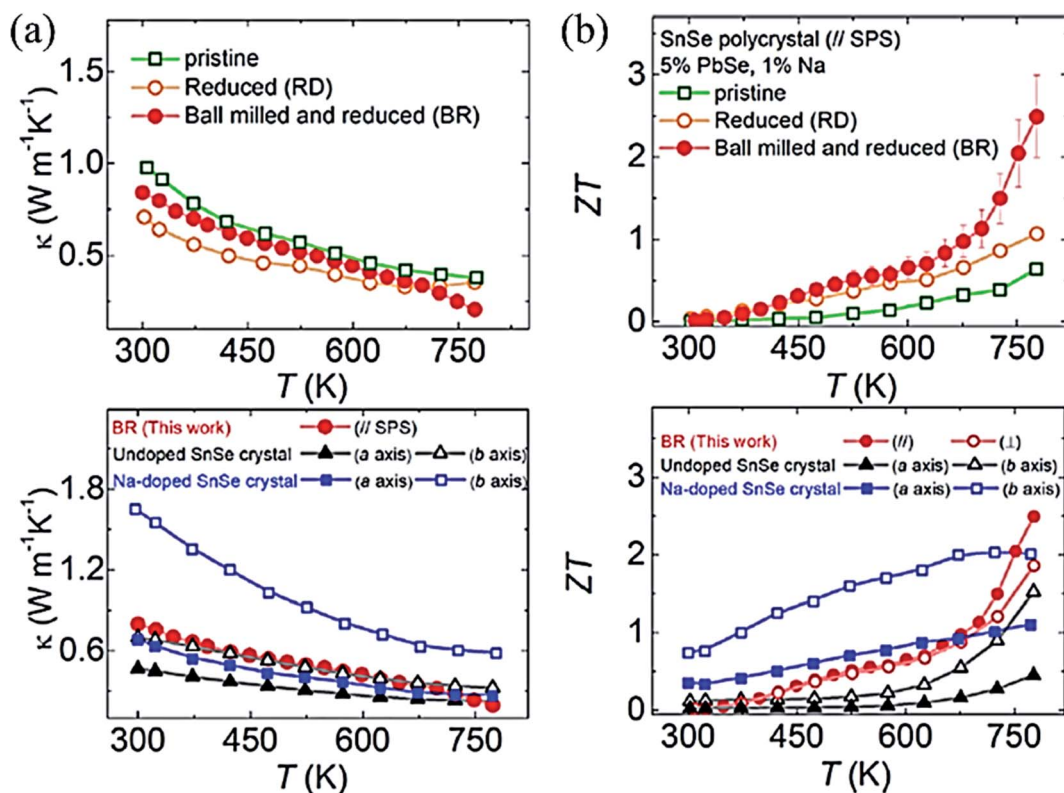


Fig. 8 (a) Total thermal conductivity of polycrystalline SnSe–PbSe (5 mol %) doped with 1 mol% Na for pristine, Reduced (R), and Ball milled then Reduced (BR) samples taken parallel to the press direction of SPS, and total thermal conductivity compared with undoped SnSe⁹³ and Na doped SnSe,⁹³ and (b) ZT of polycrystalline SnSe–PbSe (5 mol %) doped with 1 mol% Na for pristine, Reduced (R), and Ball milled then Reduced (BR), and ZT compared with the undoped SnSe⁹³ and Na doped SnSe.⁹³ These figures have been adapted/reproduced from ref. 102 with permission from Elsevier, copyright 2019⁹.

The inclusion of micro carbon fibers into the polycrystalline SnSe decoupled the thermal and electrical transport in the host matrix. Carbon fibers acted as good electrical conductors and simultaneously reduced the lattice thermal conductivity ($0.22 \text{ W m}^{-1} \text{ K}^{-1}$) by enhancing the scattering due to the high density of interface, which led to $ZT \sim 1.3$ at 823 K. This inclusion increased the mechanical stability of the device.¹¹⁹ The thermoelectric parameters of poly-crystals SnSe are tuned by sintering temperature in the spark plasma sintering (SPS). Zhang *et al.*¹²⁰ varied the sintering temperature from 300 to 650 °C in SPS and showed that vacancy defects (Se, Sn, and Se–Sn) were responsible for the performance variation.

The optimized sintering temperature for the best performance of the material was 550 °C. The ZT of 0.47 was observed at 430 °C. Ge alloying's effect was studied by varying Ge concentration in the range $x = 0.01$ to 0.03 ($\text{Sn}_{1-x}\text{Ge}_x\text{Se}$). The carrier concentration increased from 3.9×10^{17} to $4.2 \times 10^{19} \text{ cm}^{-3}$ for $x = 0.03$, which resulted in a high-power factor of $\sim 5.10 \mu\text{W cm}^{-2} \text{ K}^{-2}$ at 873 K. A very low lattice thermal conductivity of $\sim 0.18 \text{ W m}^{-1} \text{ K}^{-1}$ was achieved due to nanoscale grains, precipitation and anharmonicity due to Ge doping. Finally, high ZT of 1.75 and 2.1 were acquired at 873 K along parallel and perpendicular SPS's pressing direction.¹²¹ The figure of merit (ZT) with different doping elements in single and polycrystal SnSe is shown in Fig. 9.

SnSe thin film. Low temperature (4–300 K) thermoelectric measurements were carried out by Urmila *et al.*¹²² in SnSe thin-

film obtained by the co-evaporation method. The maximum power factor obtained was $7.2 \times 10^{-4} \text{ W m}^{-1} \text{ K}^{-2}$. The thermal conductivity was in the range of 0.023 to $0.045 \text{ W m}^{-1} \text{ K}^{-1}$, which showed a maximum ZT of 1.2 at 42 K, and thus showed a potential application in a low-temperature TE device. Nair *et al.*¹²³ fabricated the n–p type thermo-couples with PbSe (p-type), SnSe (p-type), and SnSe₂ (n-type), in the configuration SnSe–SnSe₂–SnSe–SnSe₂–SnSe and PbSe–SnSe₂–PbSe–SnSe₂–

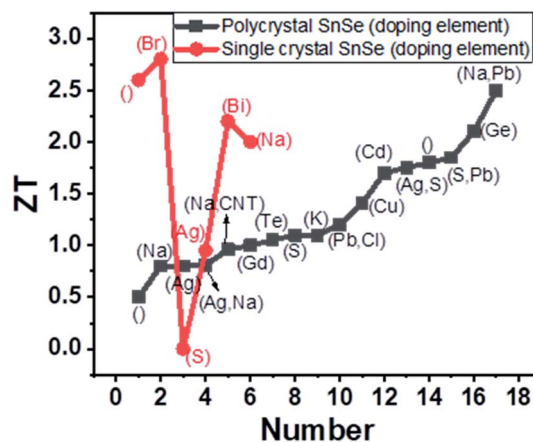


Fig. 9 Variation of ZT with the doping element given in bracket, where blank () shows the undoped SnSe.



PbSe using chemical method. A voltage was developed for a temperature difference of 20 °C of magnitude 50 mV and 15 mV, respectively (the author did not report the ZT of the device). Thin-film SnSe deposited by glancing angle (80°) pulsed laser revealed a striking power factor of $18.5 \mu\text{W cm}^{-1} \text{K}^{-2}$ at 478 K (higher than the single crystal along b axis) due to enhanced grain boundaries and out of plane thermal conductivity (highest $0.189 \text{ W m}^{-1} \text{K}^{-1}$ at 340 K). The authors did not report the ZT and in-plane thermal conductivity.¹²⁴ For the first time, Burton *et al.*¹²⁵ reported thin film SnSe thermoelectric generator on the glass substrate. Very low thermal conductivity of $0.08 \text{ W m}^{-1} \text{K}^{-1}$ in the temperature range 375 to 450 K was observed (4 times lower than single-crystal thermal conductivity along a -axis). Still, it had a low power factor, which resulted in a very low ZT of 0.055 at 501 K. When the hot side was kept at 618 K, an output of 0.09 μW was achieved. Due to the oxide formation and uncontrolled cold side temperature, the performance degraded for a longer exposure. Song *et al.*¹²⁶ synthesized

SnSe thin film of varied thickness in the range of 600–1000 nm by magnetron sputtering and showed effect of post-annealing on the performance. The optimized annealing temperature was 700 K. The power factor for the film with thickness 700 nm was $2.4 \mu\text{W cm}^{-1} \text{K}^{-2}$ at 675 K. The authors used the thermal conductivity data of bulk polycrystalline SnSe to calculate the ZT (0.28 at 675 K). Yan *et al.*¹²⁷ fabricated a flexible thermoelectric device based on SnSe composite with PEDOT:PSS that showed 13.73 nW output power at 50 K temperature difference at 353 K, and showed possible application in wearable electronics as well due to stability of device even after 1000 cycles of bending. The above discussion is summarized in Table 2.

4.2.2 SnSe₂ thermoelectric. A theoretical study on SnSe₂ using the first principal method revealed its anisotropic nature. Due to the bipolar conduction nature (because of the low bandgap) optimal predicted doping values were $(0.86\text{--}2.03) \times 10^{19}$ and $(1.71\text{--}2.47) \times 10^{19} \text{ cm}^{-3}$ along with the a and c axes, respectively. A maximum power factor of $11.72 \times 10^{-4} \text{ W m K}^{-2}$

Table 2 Thermoelectric properties of SnSe based materials

Material	Synthesis method	Power factor ($\mu\text{W cm}^{-1} \text{K}^{-2}$)	Thermal conductivity ($\text{W m}^{-1} \text{K}^{-1}$)	Temp. (K)	ZT	Ref.
Pristine SnSe (single crystal)	Bridgman Crystal growth	~9	0.34	923	2.62	32
Single crystal	Vertical Bridgman	8.5	—	850	1	128
Single crystal	Vertical vapor deposition	5.23	0.44	800	1	95
Na doped SnSe single crystal	—	~14	~0.55	773	~2	93
Br doped SnSe single crystal	Temperature gradient	~9	~0.23	773	2.8	96
S doped single crystal	Direct vapor transport	0.0011	—	573	0.0012	97
Pb doped SnSe single crystal	Flux method	1.2	—	300	—	99
Ag-Doped SnSe single crystal	Horizontal Bridgman method	~5.7	0.49	793	0.95	100
Bi-Doped SnSe single crystal	Temperature gradient	~9.75	~0.3	733	2.2	51
Polycrystal SnSe	Solid state reaction (SSR)	~4	<0.5	823	0.5	129
Polycrystal SnSe	Arc melting	~10	~0.4	816	1.8	108
Cu doped polycrystal SnSe	Hydrothermal	3.4	0.26	873	1.2	47
SnSe _{0.98} Te _{0.02}	—	~4	~0.35	805	1.05	130
Sn _{0.97} Gd _{0.03} Se	Hydrothermal	6.7	0.4	868	~1	131
Sn _{0.985} S _{0.25} Se _{0.75}	Mechanical alloying	4.5	~0.35	823	1.1	132
Sn _{0.978} Ag _{0.007} S _{0.25} Se _{0.75}	—	5.3	~0.25	823	1.75	
Sn _{0.90} Pb _{0.15} Se _{0.95} Cl _{0.05}	SSR	6.74	0.5	823	1.2	133
Cu doped polycrystal SnSe	Solvothermal	5.57	0.32	823	1.41	106
Na, Pb doped polycrystal SnSe	SSR	6.85	0.20	773	2.5	134
Polycrystal SnSe	Mechanical alloying	3.9	—	823	0.7	109
Na doped SnSe polycrystal	SSR	4.5	~0.4	773	0.8	110
K doped SnSe polycrystal	Mechanical alloying	~2.9	~0.22	773	1.1	111
Cd doped SnSe polycrystal	Solvothermal	6.9	0.33	823	1.7	112
Ag doped polycrystal SnSe	Facile surfactant-free synthesis	6.34	0.75	850	0.8	113
Ag/Na doped polycrystal SnSe	—	~7.5	~0.5	773	1.33	114
Ag/Na doped polycrystal SnSe	SSR	5	~0.48	773	0.81	115
CNT dispersed Na doped polycrystal SnSe	SSR	4.99	0.40	773	0.96	117
S, Pb doped polycrystal SnSe	Hydrothermal	4.18	~0.20	873	1.85	118
Composite of carbon fiber and polycrystal SnSe	SSR	3.88	~0.22	823	1.3	119
SnSe polycrystal	Solvothermal	~4.4	~0.7	703	0.47	120
Ge doped polycrystal SnSe	Hydrothermal	5.1	~0.21	873	2.1	121
Thin film SnSe	Reactive evaporation	7.2	~0.023	42	1.2	122
SnSe thin film	PLD	18.5	—	478	—	124
SnSe thin film	Thermal evaporation	0.11	~0.11	501	0.055	125
SnSe thin film	Sputtering	2.4	—	675	0.28	126
SnSe thin film	Chemical vapor transport	—	0.7	300	0.16	135
SnSe thin film	Chemical vapor deposition	3.2	1.1	550	0.15	136
SnSe/PEDOT:PSS thin film	Vacuum filtration method	0.24	—	353	—	127



was expected for a doping concentration of $7.21 \times 10^{19} \text{ cm}^{-3}$ at 800 K along *a* axis. Minimum thermal conductivity predicted along *a* and *c* axes at 300 K were 0.55 and $0.42 \text{ W m}^{-1} \text{ K}^{-1}$, respectively. Finally, *ZT* value was expected along *a* direction that showed a value of 0.88 for $1.94 \times 10^{19} \text{ cm}^{-3}$ at 800 K.¹³⁷ Saha *et al.* synthesized the SnSe₂ nanosheets using a solution-based method, and Cl doping was done to enhance the carrier concentration.¹³⁸ At room temperature carrier concentration increased from 7×10^{17} to $2 \times 10^{18} \text{ cm}^{-3}$ due to Cl doping. For the doped sample, almost double power factor of $1.46 \mu\text{W cm}^{-1} \text{ K}^{-2}$ was achieved at 630 K. As synthesized SnSe₂ nanosheets showed the thermal conductivity in the range of ~ 0.45 to $0.35 \text{ W m}^{-1} \text{ K}^{-1}$ while doping increased the thermal conductivity in the range of 0.67 to $0.40 \text{ W m}^{-1} \text{ K}^{-1}$ at 300 to 630 K temperature, respectively. Overall, *ZT* of 0.22 was achieved at 610 K for Cl doped samples. For heavy Cl doping of $x = 0.12$ by solid-state reaction, *ZT* (out of plane) of ~ 0.4 was achieved by Xu *et al.*¹³⁹ Se deficiency and Cl doping's led to enhanced power factor and reduced thermal conductivity resulting in a *ZT* of 0.63 at 673 K along an in-plane direction.¹⁴⁰ A study on Ag doping revealed that it effectively enhances the carrier mobility but decrease the carrier concentration. For 1% doping, the optimized power factor was $3.50 \mu\text{W cm}^{-1} \text{ K}^{-2}$ at 773 K, which resulted in a *ZT* of ~ 0.4 .¹⁴¹ In another theoretical study, Ding *et al.* reinvestigated the thermoelectric performance of the SnSe₂ by the first principle method. They considered various models of the phonon-phonon scattering, which may overestimate the thermal conductivity.³³ Room temperature electrical conductivities along *a* and *c* directions were calculated as 4.97×10^5 and $3.39 \times 10^4 \Omega^{-1} \text{ m}^{-1}$, respectively, at 10^{20} cm^{-3} carrier concentration. Overall, for n-type highly doped (10^{20} cm^{-3}) SnSe₂, a *ZT* of 3.6 could be achieved at 800 K. Wu *et al.* doped Br to the site of Se in polycrystalline SnSe₂ to enable broad carrier concentrations (0.5 to $5.6 \times 10^{19} \text{ cm}^{-3}$, saturation at doping of $x = 0.01$) consistent with the single parabolic model.¹⁴² They reached a *ZT* of 0.6 at 750 K, which was competitive to polycrystalline SnSe. The Cl-doped composite of SnSe–SnSe₂ showed *ZT* ~ 0.56 at 773 K.¹⁴³ Dynamical intercalation of Ag into the weak van der Waals

interlayers of SnSe₂ acted as electron donor and introduced line defect, twin boundary dislocation, phase interface, and enhanced phonon–phonon scattering. To further strengthen carrier concentration into Ag-doped SnSe₂, Liu *et al.* doped it with Cl, which resulted in concentration in the range of 10^{19} due to Fermi leveling up to the conduction band (as shown *via* DFT calculation).¹⁴⁴ Ag interlayer bridge weakened the anisotropy of electrical transport; hence the power factor along parallel and perpendicular directions of pressing were 7.46 and $8.25 \mu\text{W cm}^{-1} \text{ K}^{-2}$, respectively, at 789 K, and a maximum *ZT* of ~ 1.03 was achieved at 789 K along the parallel direction of the pressure. Recently, Wang *et al.* showed the room temperature thermoelectric potential of the SnSe₂ embedded with Cu, which enhanced carrier concentration two orders of magnitude. Optimum doping in SnCu_{*x*}Se₂ was found to be $x = 0.01$, which resulted in an optimized power factor of 1.96 at room temperature and a decrease in thermal conductivity up to $0.81 \text{ W m}^{-1} \text{ K}^{-1}$ due to increased interface.¹⁴⁵ A very high *ZT* of 0.75 than to pristine SnSe₂ was obtained which was about two orders of magnitude larger than pristine SnSe₂ (Fig. 10(a)).

Thin-film SnSe₂ thermoelectric. Thermoelectric performance of SnSe₂ thin films have also been studied. But there are very few reports on this, and the reports are limited to the film's power factor only. Yin *et al.* deposited SnSe₂ thin-film on a glass substrate by spin coating techniques and reported a Seebeck coefficient of $\sim 126.3 \mu\text{V K}^{-1}$. The electrical conductivity of the order of 10^3 S m^{-1} at room temperature was observed.¹⁴⁶ Chen *et al.* deposited the SnSe thin film on the Si wafer and annealed it in Se vapor. The film was reduced to SnSe₂ and showed enhanced Seebeck coefficient $631 \mu\text{V K}^{-1}$ for SnSe₂ than $38.6 \mu\text{V K}^{-1}$ for SnSe. Around 44 times enhanced power factor ($0.2 \mu\text{W m}^{-1} \text{ K}^{-2}$) was obtained.¹⁴⁷ Duong *et al.* deposited SnSe₂ thin film on the Al₂O₃ substrate by the pulsed laser deposition technique. They reported a power factor of $8 \mu\text{W m}^{-1} \text{ K}^{-2}$ at 220 K.¹⁴⁸ A detailed summary of the above discussions and *ZT* variation with the doping is provided in Table 3 and Fig. 10(b), respectively.

This material (SnSe) has the highest efficiency among the bulk materials without any doping, but it faces

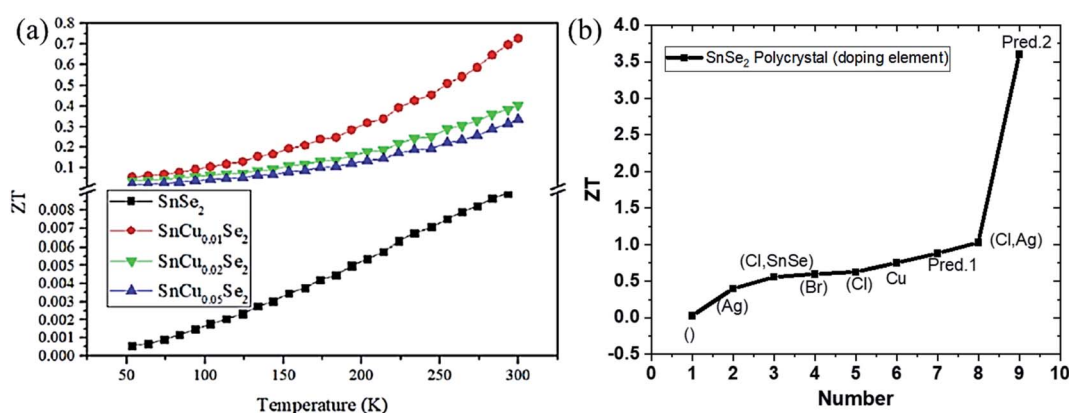


Fig. 10 (a) Figure of merit (*ZT*) of SnCu_{*x*}Se₂ ($x = 0, 0.01, 0.02, 0.05$) as a function of temperature. This figure has been adapted/reproduced from ref. 145 with permission from ACS, copyright 2020". (b) Variation of *ZT* with different doping elements, irrespective of operating temperature, where blank () represents undoped SnSe₂, and Pred. 1 and Pred. 2 mean theoretical predicted values, respectively.



commercialization issues like poor mechanical stability and low thermoelectric property below ~ 800 K. Besides its high ZT , low cost, easy fabrication, and earth-abundance, SnSe suffers the real possible application likely due to the quick oxides' formation or thermal instability.¹⁵⁰ It readily forms the oxides in the range >600 °C, where it showed maximum ZT . Protective environment costs hamper its effectiveness as a cheap/high ZT . Sassi *et al.*¹²⁹ showed thermal instability above the phase transition temperature of SnSe, and the working efficiency of the device decreased continuously over cyclic performance due to Se loss. Thus, there should be an optimized temperature range over which there is no loss of performance. Continuous research should go on in the direction to improve the thermal stability up to the temperature for which SnSe showed maximum ZT . Also, large phonons scattering centers exist in polycrystalline SnSe than in single crystal SnSe due to grain boundaries, yet single crystal SnSe has lower lattice thermal conductivity than the polycrystalline SnSe. Ibrahim *et al.*¹⁵¹ noted this fact and investigated the reason behind this. They systematically discarded the oxide formation and found many native defects in the single crystal.¹⁵² Their results reflected higher lattice thermal conductivity than the reported one.³² They also noted that the single crystal grown by Zhao *et al.*³² was not fully dense and hence may be the reason for low lattice thermal conductivity. But it still requires the study to know the exact causes of the very low lattice thermal conductivity, whether it was due to intrinsic defect or not. SnSe single crystal suffers a low mechanical strength and high growth cost as a single crystal. SnSe single crystal thermoelectric results showed inconsistencies in the different group reports' as shown in Table 2. However, the further study required to unravel the SnSe performance, where it seeks the possibilities of enhancement in ZT .³⁸ Another phase, SnSe₂, was also predicted by theoretical consideration as the best thermoelectric material that can achieve ZT around 3.5. Till date, $ZT \sim 1.1$ is achieved in this material. Also, both materials (SnSe and SnSe₂) showed inferior TE property in thin film form. Thus, this material demands greater attention and research to achieve predicted ZT . Researchers should pay attention to new approaches to improve the SnSe thermoelectric material's performance like

nano-inclusion,¹⁵³ decoupling of interrelated parameters,¹¹⁹ nano-structuring and texturing, *etc.* Simultaneously, the theoretical study should be carried out to optimize the threshold values of the parameters. Finally, tin-selenide has emerged as a futuristic material that showed the best efficiency. Both materials (SnSe and SnSe₂) demand more research in the thin film and bulk thermoelectric form because of the inferior TE property and device stability.

4.3 Photodetectors (PDs)

4.3.1 SnSe based PDs. SnSe act as suitable optoelectronic material due to its high absorption coefficient ($>10^4$ cm⁻¹),¹⁵⁴ low band gap and its tunable band gap over an extensive range.¹⁵⁵ Solvothermal processed SnSe nanorod showed repeatable and stable photoresponse.¹⁵⁶ SnSe thin film grown on the n-Si substrate showed an ultrahigh response/recovery time of 0.9/17.3 μ s in a position-dependent detector.¹⁵⁷ A flexible device of SnSe on the mica substrate (grown by PLD) showed high responsive behavior of the device in a longer wavelength region (370–808 nm). Also, it reflected its durability in bending test for the flexible device. The device showed the highest photoresponse of 5.5 A W⁻¹ which is attributed to its shallow potential barrier with Bi₂Te₃ contacts (perfect band alignment) and low potential fluctuation at the interface.¹⁵⁸ Under white light illumination, SnSe nanoplate and graphene composites showed very high photosensitivity (1110%) and short response time (~ 1 s).¹⁵⁹ The SnSe film grown on Si substrate by vapor deposition and dispersed with graphene oxide quantum dots showed improved response time and photocurrent than the SnSe film in a larger spectrum.¹⁶⁰ Self-powered photodetection was demonstrated in the SnSe/Si heterostructure device, with very high detection and ultrafast response in μ s range. Magnetron sputtered SnSe film on Si(100) showed a nearly ideal diode property with a reported power exponent of 0.92 and ideality factor of ~ 1 , which revealed a trap-free and high-performance of the device.³⁶ Zhong *et al.*¹⁶¹ showed the thin film's superiority due to the Marangoni effect over other solution-based methods. The flexible and vertically fabricated Gr/SnSe/Gr device showed a responsivity of 38 mA W⁻¹. Magnetron sputtered SnSe thin film on Si substrate

Table 3 Thermoelectric properties of doped-SnSe₂ based materials

Material	Synthesis method	Power factor (μ W cm ⁻¹ K ⁻²)	Thermal conductivity (W m ⁻¹ K ⁻¹)	Temp. (K)	ZT	Ref.
Cl doped SnSe ₂	Facile low-temperature solution	1.46	~ 0.45	610	0.22	138
Cl doped SnSe ₂	Solid-state reaction followed by hot press	~ 7.0	~ 1.6	673	0.4	139
Cl doped SnSe ₂	Solid-state reaction followed by ball milling and SPS	~ 8	~ 1.2	673	0.63	140
Ag-Doped SnSe ₂	Mechanical alloying followed by SPS	3.50	< 1	773	0.4	141
Br doped SnSe ₂	Solid-state reaction followed by quenching and hot press	—	—	750	0.6	142
Cl doped SnSe–SnSe ₂	Solid state reaction	—	0.42	773	0.56	143
Ag and Cl doped SnSe ₂	Bottom up	7.46	0.57	789	1.03	149
Cu embedded SnSe ₂	Solid-state reaction followed by ball mill and SPS	1.96	0.81	300	0.75	145



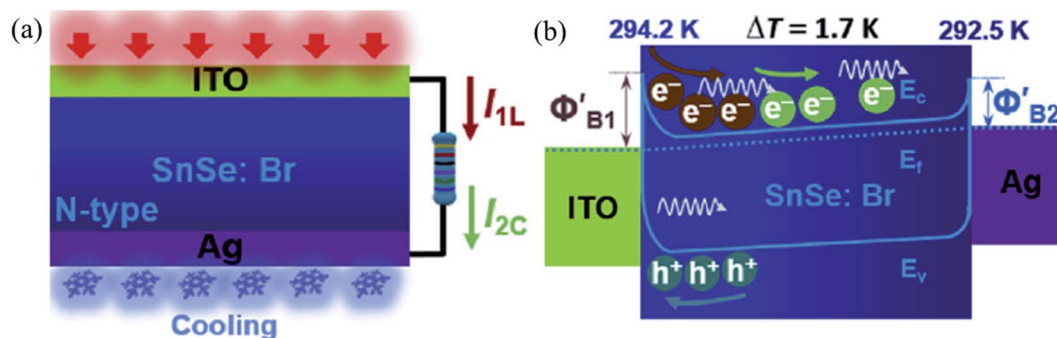


Fig. 11 Schematic working state (a) and energy band diagram (b) of the device under the combined action of light illumination and cooling (I_{1L} = photovoltaic current, I_{2C} = thermocurrent). Reprinted with permission from. These figures has been adapted/reproduced from ref. 163 with permission from Elsevier, copyright 2019".

showed broadband photoresponse in the 404–980 nm region. A maximum responsivity of 277.3 A W^{-1} was reported. This high value of responsivity was attributed to the film's trap states, which was confirmed by the nanosecond order's carrier lifetime observed in the time-resolved photoluminescence studies.¹⁶² Ouyang *et al.*¹⁶³ harnessed combined photovoltaic and thermoelectrics effect to enhance the self-biased n-type SnSe:Br (ITO/n-SnSe/Ag) photodetector performance. They coupled the thermoelectric property with the material's photovoltaic property *via* maintaining the temperature difference with Peltier cooler's help at the one contact (Ag contact) of the device and another under the illumination of 760 nm radiation. The thermoelectric diffusive process accelerates photoelectrons generated under photovoltaic conditions. This is due to specific heat capacities of electrons and photonic drags under the thermal gradient resulting in enhanced performance (Fig. 11(a) and (b)). Simultaneously, in the energy band aspect, the thermoelectric effect reduced the Schottky barrier height at the ITO side and increased at the Ag side, which resulted in the de-acceleration of the photoelectrons. The combined effect enhanced the PD parameters. The light current and voltage were

improved by 38.1 and 81.9%, respectively, at a cooling temperature difference of 1.5 K (the Ag side was cooler compared to the ITO side).

When the Ag side is heated w.r.t. ITO with 4.5 K difference, the opposite result was observed, *i.e.*, there was a decrease in photocurrent and voltage. Hence authors concluded that the coupling of the effect is polarity dependent. Record high value (comparable to the commercialize Ge and Si photodetectors) of responsivity (R) and detectivity (D) were reported in the nanowire SnSe.¹⁶⁴ Above discussions have been summarized in Table 4.

4.3.2 SnSe₂ based PDs. The tunable bandgap and high absorption coefficient of the SnSe₂ have made it a potential candidate for optoelectronics applications. For the first time, Zhou *et al.*¹⁷³ reported ultrathin SnSe₂ flake (1.5 nm) grown by a CVD method using SnI₂ as a precursor, which exhibited high responsivity (R) 1100 A W^{-1} with meager response time in millisecond. However, the far value of power exponent (0.7), reflected its impurity and defects. Bilayered SnSe₂ of 1T type structure with D_{3d} point group symmetry showed high responsivity of

Table 4 Performance characteristic of SnSe based PDs

Device structure	λ (nm)	Bias (V)	R (A W^{-1})	D (Jones)	Response/recovery time (s)	Ref.
ITO/SnSe/ITO	White light	5	—	—	3	165
Pd/n-Si/SnSe/Pd	780	—	—	—	$0.9/17.3 \times 10^{-6}$	157
Bi ₂ Te ₃ /SnSe/Bi ₂ Te ₃	370	20	5.5	6×10^{10}	—	166
Au/SnSe/graphene/Au	White light	15	—	—	$1/<1$	159
Ag/Si/SnSe/graphene oxide QD/Ag	650	—	—	—	0.18/0.75	160
Pd/SnSe/Si/In	850	0	0.566	4.4×10^{10}	$1.6/47.7 \times 10^{-6}$	167
Gr/SnSe/Gr	400	1	0.038	—	~ 0.18	161
Ag/Si/SnSe/Ag	404	15	277.3	7.6×10^{11}	0.35/1.83	168
ITO/n-SnSe/Ag	760	0	3.97×10^{-3}	—	$81/122 \times 10^{-6}$	163
Cr/Au/Si/SnSe (nanowire)/Cr/Au	830	3	1.0×10^4	3.3×10^{12}	$460/520 \times 10^{-6}$	164
Mica/In ₂ Se ₃ /SnSe/Au	405	5	0.350	—	0.156/0.139	169
Si/SiO ₂ /WSe ₂ /SnSe/Ti/Au	1064	5	6.6×10^{-3}	—	—	170
	671		31.8×10^{-3}		—	
	457		99×10^{-3}		0.0082/0.0084	
PCB/Mica/SnSe (nanocrystal)/Ag	(Sunlight)	2	0.54×10^{-3}	1.06×10^9	1.5/1.7	171
PET/SnSe/Pd	404	20	1745.5	$\sim 4.2 \times 10^{12}$	1.7/4.7	172
	850		78.6	$\sim 9 \times 10^{11}$	0.23/0.27	



0.5 A W^{-1} and a swift response time of $\sim 2 \text{ ms}$.¹⁷⁴ SnSe_2 acted as an efficient charge separator in the heterostructure devices. It collected electrons from the WSe_2 ,¹⁷⁵ MoS_2 ¹⁷⁶ (which have deeper conduction minima than these) so efficiently that it improved performance parameters several times better. SnSe_2 with MoS_2 interface resulted in enhanced responsivity from 37.3 to $9.1 \times 10^3 \text{ A W}^{-1}$ under 500 nm illumination.¹⁷⁶

Mukhokosi *et al.* studied the thickness-dependent optical properties of the SnSe_2 thin film grown by DC sputtering followed by selenization and also studied the photodetection performance (Fig. 12(a) and (b)).³⁰ Films with thickness $< 140 \text{ nm}$ did not show any IR response. The film with 1200 nm thickness had a slow response (time), and responsivity was 0.4 mA W^{-1} .³⁰ Mukhokosi *et al.*¹⁷⁷ also reported the self-powered, organic-inorganic hybrid heterostructure PD consisting of poly (3,4-ethylene dioxythiophene):poly(styrene sulfonate) (PEDOT:PSS) and SnSe_2 . The device had responsivity in the range of $1.4\text{--}2.6 \mu\text{A W}^{-1}$. The response time improved to 1.33 s . Mukhokosi *et al.*¹⁷⁸ reported highly enhanced response time and enhanced responsivity. The device was fabricated on the p-type Si substrate by DC sputtering, followed by selenization. Kumar *et al.*⁶⁵ deposited SnSe_2 thin-film on the soda-lime glass substrate and reported the improved responsivity at very low bias and fast response time in the IR range (Fig. 12(c) and (d)). Although the band gap was in the visible region, photodetection

was observed in the IR region which was attributed to the defect/trap assisted levels lying in the forbidden region.⁶⁵ Usually, SnSe_2 shows high absorption coefficient^{30,65} but poor chemical stability against environmental conditions (degrade into SnSe after prolonged exposure). Therefore, Gao *et al.* used graphene (Gr) to improve the chemical stability. SnSe_2 was sandwiched between graphene. Monolayer Gr acted as useful ohmic contacts and also had weak absorption in the UV to IR range.¹⁷⁹ This sandwiched structure of SnSe_2 with Gr showed responsivity of $1.09 \times 10^3 \text{ A W}^{-1}$, and the response time was 30 ms . Murali *et al.* studied the substrate-induced effect on SnSe_2 flake (i) suspended over contacts and (ii) supported on Si substrate. It showed a vast difference in rise and decay times.¹⁸⁰ The suspended structure showed low gain but quicker decay than the supported one because of the device's interface-induced trap states. Field-effect transistors based on multi-layer SnSe_2 showed very high responsivity and detectivity. The device showed its potential application comparable to the other commercialized photodetector.⁹ The above discussion is summarized in Table 5.

4.4 Gas sensors

The gas sensor is a basic need for environment cleaning. A gas sensor is a device, which detects harmful gases and sends an

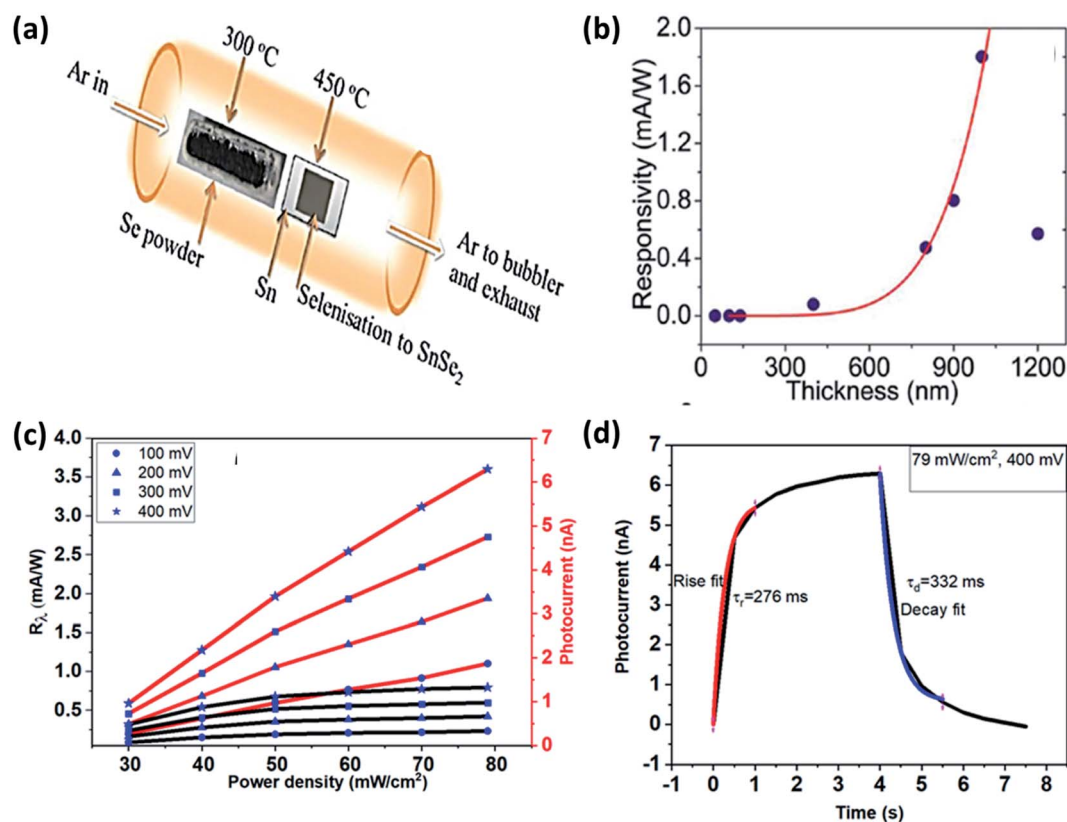


Fig. 12 (a) Schematic of selenization of DC sputtered Sn film and (b) thickness-dependent SnSe_2 thin films' responsivity. These figures have been adapted/reproduced from ref. 30 with permission from Nature, copyright 2017. (c) Responsivity and photocurrent with power density at different bias voltage and, (d) response and recovery time at 400 mV bias, these figures have been adapted/reproduced from ref. 65 with permission from Elsevier, copyright 2020.

Table 5 Summary of performance characteristics of SnSe₂ based photodetectors^a

Device structure	λ (nm)	Bias (V)	R (A W ⁻¹)	Detectivity (jones)	Rise/decay time (s)	Ref.
Ti/Au/Si/SnSe ₂ (flake)/Ti/Au	530	3	1100	10 ¹⁰	14.5/8.1 × 10 ⁻³	57
Ti/Au/Si/2LSnSe ₂ /Ti/Au	633	0.1	0.5	—	2.2/3.2 × 10 ⁻³	174
Ni/Au/ITO/WSe ₂ (flake)/SnSe ₂ (flake)/SiO ₂ /Si/Ni/Au (vertical hetero-structure)	785	0.4	1139	—	8/32 × 10 ⁻⁶	175
Cr/Au/SnSe ₂ (flake)/MoS ₂ (flake)/Cr/Au	500	1	9.1 × 10 ³	9.3 × 10 ¹⁰	0.2/0.6	181
Cr/Au/SLG/SnSe ₂ (thin film)/Cr/Au	1064	10	0.4 × 10 ⁻³	10 ¹¹	2.5/3.68	30
Cr/Au/SLG/SnSe ₂ (thin film)/PEDOT:PSS/Cr/Au	1064	0	1.4–2.6 × 10 ⁻⁶	~10 ⁸	1.33/1.22	177
Cr/Au/p-Si/SnSe ₂ /Cr/Au	1064	10	0.16	~1.5 × 10 ⁹	57/34 × 10 ⁻⁶	182
Au/SLG/SnSe ₂ (thin film)/Au	1064	0.4	0.796 × 10 ⁻³	5.62 × 10 ⁷	0.276/0.332	65
Au/Gr/SnSe ₂ (flake)/Gr/Au	532	0.5	1.09 × 10 ³	1.2 × 10 ¹²	30.2/27.2 × 10 ⁻³	183
Ni/Au/SnSe ₂ (flake)/Ni/Au (suspended structure), Ni/Au/Si/SnSe ₂ (flake)/Ni/Au (supported structure)	White light	0.1	115, 8.66 × 10 ⁴	~10 ¹³	SED-66 s, DED-2.27 s, 53 s	180
Cr/Au/n-Si/SnSe ₂ (single crystal flake)/Cr/Au	450	3	5.11 × 10 ⁵	2.79 × 10 ¹³	—	9

^a SED = single exponential decay, DED = double exponential decay.

alarm. There are many toxic gases in the environment (CO₂, NO₂, NH₃, SO₂, H₂S, CH₄, etc.). Gas sensing parameters are response time, recovery time, selectivity, gas concentration, temperature, detection limit, etc. Tin based chalcogenide (SnSe and SnSe₂) have shown promise for gas sensing applications. A brief discussion of tin chalcogenide (SnSe and SnSe₂) materials for gas sensing applications is described. The gas sensing mechanism of the tin chalcogenide (SnSe and SnSe₂) gas sensor is based on the adsorption of gas molecule and charge transfer.

Physisorption-based charge transfer in tin chalcogenide (SnSe and SnSe₂): To evaluate the interaction efficiency of the SnSe/SnSe₂ monolayer and gas molecules, the absorption energy (E_a), Hirschfeld charge transfer (Q), and the distance (d_0) of nearest atoms between the gas molecule and the SnSe/SnSe₂ layer were calculated. The absorbed energy is defined as:

$$E_a = E_{\text{total}} - E_{\text{gas}} - E_{\text{material}} \quad (1)$$

E_{gas} , E_{material} , and E_{total} are the gas molecule's total energy, SnSe/SnSe₂ monolayer, and gas molecule-SnSe/SnSe₂ system, respectively.¹⁸⁴ The adsorption properties, including adsorption energy (E_a), equilibrium distance (d_0), and Hirshfeld charge transfer (Q) for SnSe, are listed in Table 6. The adsorbed CO, CO₂, CH₂O, NO₂, and SO₂ gas molecules on the β -SnSe monolayer have the E_a values -0.202, -0.175, and -0.322, -0.829, and -0.499 eV, respectively, and the Hirshfeld charge transfer values are -0.033, -0.036, -0.085, -0.279, and -0.279 e, respectively.¹⁸⁵

The equilibrium distance of CO, CO₂, CH₂O, NO₂, and SO₂ from the β -SnSe layer is 3.293, 3.617, 3.222, 2.531, and 2.692 Å, respectively, which are greater than the C-Sn (2.15 Å), O-Sn (2.22 Å), and N-Sn (2.11 Å) bonds, indicating the process to be physisorption.¹⁸⁴ However, the optimized structure shows some distortion for O₂ on the surface of the β -SnSe sheet. The value of E_a and Q_c for O₂ are -1.596 and -0.445 eV, respectively, which is much larger than other gas molecules adsorbed on one SnSe layer, indicating that O₂ molecules (CO, CO₂, CH₂O, NO₂, and SO₂) are chemically absorbed in the β -SnSe layer.¹⁸⁵ Adsorption

of CO, CO₂, CH₂O on the β -SnSe layer showed lower adsorption energy and lower charge transfer values, which indicates that one β -SnSe layer is not suitable for the detection of these three molecules.¹⁸⁴

The higher absorption energy of NO₂ (-0.829 eV) and SO₂ (-0.499 eV) indicates that NO₂ and SO₂ molecules' adsorption behavior for the β -SnSe layer was more potent than that of the absorption system.¹⁸⁵ The charge transfer values for NO₂ and SO₂ are -0.279 and -0.278 eV, respectively, which show a clear charge transfer between the gas molecule and the β -SnSe layer (Fig. 13(a)-(f)).¹⁸⁵ The adsorption distances for NO₂ and SO₂ are 2.531 Å and 2.692 Å, respectively, which is close to the Sn-O bond length range (2.22 Å to 2.66 Å).¹⁸⁶ For the β -SnSe sheet, it is crucial to detect NO₂ and SO₂ in the gas sensor region. All the most stable energy adsorption centers are located on the Sn side of the atom of the β -SnSe sheet, which indicates that the absorbing properties of the metal atom are more substantial than that of the nonmetallic atom.¹⁸⁷

4.4.1 SnSe based gas sensors. Tin selenide (SnSe) is a member of (IVA-VIA) binary semiconductors family and has a high absorption coefficient.¹⁸⁸ The bandgap of SnSe varies from 0.9 eV to 1.3 eV.¹⁸⁹ SnSe has an orthorhombic crystal structure, and it has p-type conductivity with Sn-vacancies.¹⁹⁰ The majority of charge carriers of SnSe are holes and have a good impact on gas sensing applications. Using Ph₃PSe as

Table 6 The adsorption energy, equilibrium distance, and Hirshfeld charge transfer of different molecules adsorb on a β -SnSe monolayer. Reprinted with permission from¹⁸⁵

Molecule	E_a (eV)	d_0 (Å)	Q (e)
CO	-0.202	3.293 (C-Sn)	-0.033
CO ₂	-0.175	3.617 (C-Sn)	-0.036
CH ₂ O	-0.322	3.222 (H-Sn)	-0.085
O ₂	-1.596	2.058 (O-Sn)	-0.445
NO ₂	-0.829	2.531 (O-Sn)	-0.279
SO ₂	-0.499	2.692 (O-Sn)	-0.279



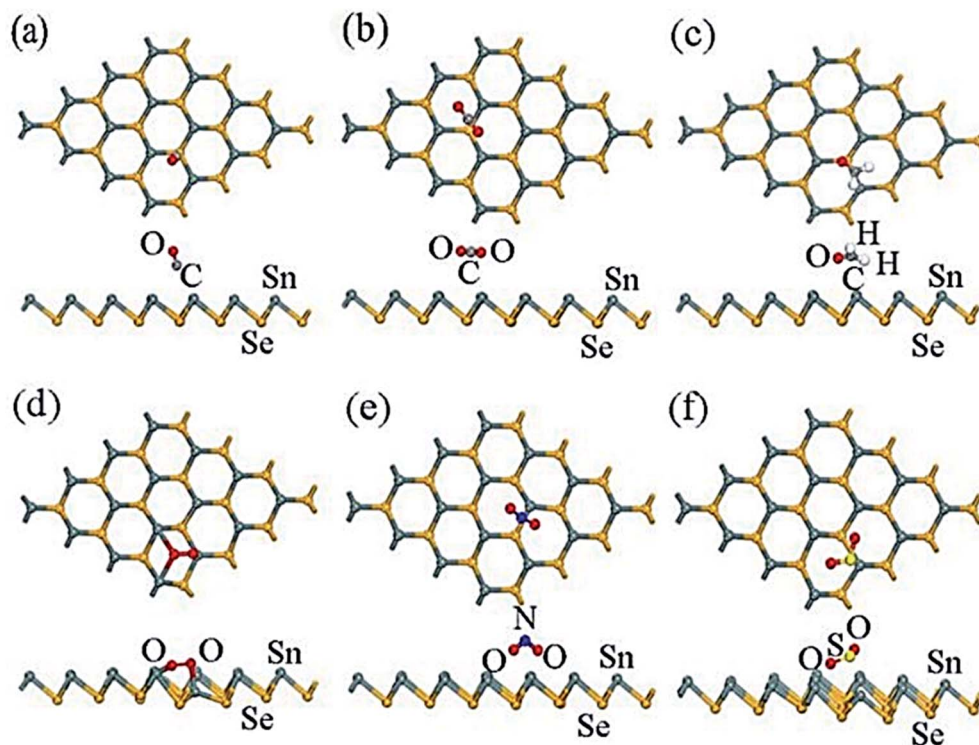


Fig. 13 The most stable sites of optimized configurations of the adsorbate molecules: (a) CO, (b) CO₂, (c) CH₂O, (d) O₂, (e) NO₂, and (f) SO₂ adsorbed on a β -SnSe monolayer. Most stable sites are exhibited. These figures have been adapted/reproduced¹⁸⁵ with permission from MDPI, copyright 2019*.

a precursor in chemical vapor deposition, Assili *et al.*²¹ deposited orthorhombic tin selenide thin films onto three substrates. A 1% vol concentration of methane gas showed a good sensitivity at the operating temperature of about 200 °C. The sensitivity, response, and recovery times were ~47%, ~52 s, and ~220 s, respectively. They observed that when SnSe was mixed with any n-type materials, it showed excellent response and recovery times at, lower operating temperature. In 2013, Wang *et al.*¹⁹¹ synthesized SnO₂ nanoparticles decorated SnSe nanosheets *via* a facile, lost-cost, and safe solution method and had studied the gas sensing properties. For 1000 ppm CO gas, this device's response was 1.9 s at the operating temperature of 260 °C (Fig. 14(a)–(c)). Lee *et al.*¹⁹² prepared the SnSe₂/SnSe heterostructure film from a thick Sn layer onto the glass substrate. The film also showed response for one ppm NO₂ gas at room temperature, and the response value was 75%. In 2020, Wang *et al.*¹⁹³ reported a one-step colloid method for making SnSe/SnSe₂ heterostructures, with doping of SnSe ≈ 30, 50, and 70%, respectively. The SnSe (50%)/SnSe₂ (50%) based sensor with an active layer thickness of 2 μm showed the highest sensitivity of 30% to 0.1 ppm NO₂ gas at room temperature (25 °C) with a limit of detection (LOD) down to 69 ppb. The above discussion is summarized in Table 7.

4.4.2 SnSe₂ based gas sensors. SnSe₂ is an n-type material due to the presence of Se-vacancies. The majority of charge carriers in SnSe₂ are electrons. The SnSe₂ film is an anisotropic binary material having a hexagonal structure arranged in the form of Se–Sn–Se.¹⁹⁴ SnSe₂ has a small bandgap of 1–2 eV.^{30,195} SnSe₂ has the same electronegativity and similar structure as

MoS₂, thus it is better suited for the detection of NO₂. 2D materials have a substantial surface to volume ratio, which significantly changes the absorption and release of gases, making them an inexpensive gas detection candidate.¹⁹⁶ Given all these properties of SnSe₂, various researchers have tried to study the properties of SnSe₂ gas. Lee *et al.*¹⁹² prepared the SnSe₂/SnSe heterostructure film on the glass substrate from a thick Sn layer. At room temperature (RT), this film showed a better response for one ppm NO₂ gas. The response value of this film was 75%. In 2010, Popescu *et al.*¹⁹⁷ prepared silver doped SnSe₂ and Ge₂Sb₂Te₅ thin films by pulsed laser deposition. Ag-Doped SnSe₂ thin film-based gas sensor had a response of 3.26 for 500 ppm of CO gas at 500 °C. In 2018, Chen *et al.*¹⁹⁴ produced SnSe₂ nanoplate arrays and made a sensor that could detect 1% vol. of CH₄ at the operating temperature of 200 °C. This sensor showed a response of 66.7 and response/recovery times of 78/336 s, respectively. In 2019, Moreira *et al.*⁶¹ investigated the NO₂ sensing properties of CVD deposited SnSe₂ binary layer. For one ppm of gas at the operating temperature of 50 °C, this sensor's response was 60%, and the response/recovery times of this sensor were 142/457 s (Fig. 14(d)–(f)). In 2017, Assili *et al.* demonstrated a SnSe₂ thin film-based gas sensor. This sensor had sensitivity around 16% for 200 ppm of methane at 200 °C with the response and recovery times around 75 s and 615 s, respectively.²¹ In 2007, Popescu *et al.*¹⁹⁸ reported a SnSe₂ thin film-based gas sensor, which had a response of 300% for 8000 ppm of CH₄ gas at 600 °C. Sanju Rani *et al.* demonstrated a SnSe/SnSe₂ nanostructured thin film-based sensor by using of



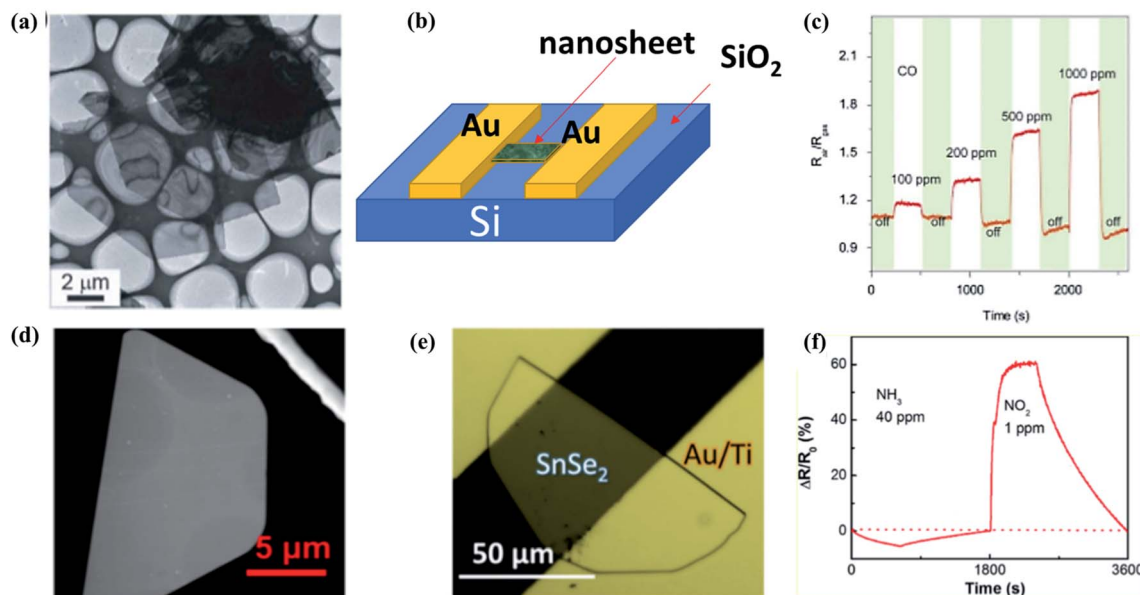


Fig. 14 (a) Low magnification TEM images, (b) schematic structure of the device, (c) transient response of the sensor SnO₂/SnSe to CO (100–1000 ppm) at 260 μC. These figures has been adapted/reproduced from ref. 191 with permission from Nature, copyright 2013¹⁹¹. (d) SEM image of a semi hexagonal nanosheet of SnSe₂, (e) optical image of a 6 nm thick SnSe₂ gas sensor device, (f) dynamic sensing responses of the 6 nm thick SnSe₂ resistor device measured with 405 nm laser illumination. These figures has been adapted/reproduced from ref. 61 with permission from ACS, copyright 2019⁶¹.

thermal evaporation technique. This sensor had response of 112% for (5 ppm) NO₂ gas at room temperature and the response/recovery times were 15/10 s, respectively.¹⁹⁹ AuPd coated SnSe₂ thin film based NO₂ sensor showed enhanced response of 117% for 5 ppm NO₂ gas at room temperature. The response/recovery times of this sensor were 10/18.7 s respectively.²⁰⁰ Thus, efforts have been made to use SnSe₂ for sensing methane as well as NO₂. The above discussion is summarized in Table 8.

4.5 Photocatalysis

4.5.1 Photocatalytic behavior. There are several reports on the photocatalytic behavior of Sn-based materials like SnSe and SnS. The possible mechanism reported by Li Cheng and co-workers is based upon the bandgap and absorption of light. The SnS and SnSe nanofibers have a narrow bandgap of 1.01 eV and 0.90 eV, respectively. UV light excites the photoelectron and hole in the material. Now the exciting photoelectron and hole react with dissolved oxygen molecules to form oxide radicals (O₂). Furthermore, hydroperoxyl HO₂ and hydroxyl radical OH[•]

are formed after the protonation process (hydrogen cation). Simultaneously, the holes (h⁺) can oxidize hydroxyl radicals OH[•] and H₂O molecules to generate OH[•] and hydroperoxyl HO₂, which may break the RhB dye molecule to convert them into CO₂ and H₂O or in other forms as shown in Fig. 15. The dye's degradation rate is defined as

$$D\% = \left[\frac{A_0 - A_T}{A_0} \right] \times 100\% \quad (2)$$

Here A₀ is the absorbance of dye in the dark condition. A_T is the absorbance of dye at specific time gaps after irradiation of light of a certain wavelength. Li Cheng *et al.* used the RhB solution under UV radiation to check the SnO₂, SnSe, and SnS nano-fibers' photocatalytic activity.⁶² The degradation curve showed a first-order kinetics equation type curve. The degradation rate for RhB was 85.90%, 92.55% and 92.86% for SnO₂, SnS and SnSe, respectively. In another experiment, Jing group used thin films of Sn and SnSe to analyze the photodegradation

Table 7 The comparison between the gas sensing parameter of the different structures of SnSe

Material	Structure	Gas	Gas conc. (ppm)	Temp. (°C)	Response	Response time (s)	Recovery time (s)	Ref.
SnSe	Thin film	Methane	1% vol	200	~47% ^a	~52	~220	21
SnO ₂ /SnSe	Nanosheet	CO	1000	260	1.9 ^b	—	—	191
SnSe ₂ /SnSe	Heterostructure	NO	1	RT	75% ^c	—	—	192
p-SnSe/n-SnSe ₂	Hetero-structure	NO ₂	0.1	RT	30% ^b	—	—	193

^a $R = (R_g - R_a)/R_a \times 100$. ^b $R = R_a/R_g$. ^c $I = (I_g - I_o)/I_o$. R is called the sensor's response, S is called the sensor's sensitivity, R_g is the resistance of the sensor in the presence of target gas, and R_a is the sensor's resistance in the air. I_g is the current of device in presence of target gas and I_o is the current of device in presence of air.



Table 8 The comparison between the gas sensing parameter of the different structures of SnSe₂

Material	Structure	Gas	Gas concentration (ppm)	Temp. (°C)	Response	Response time (s)	Recovery time (s)	Ref.
SnSe ₂ /SnSe	Heterostructure	NO ₂	1	RT	75% ^a	—	—	192
SnSe ₂ -Ag	Thin film	CO	500	500	3.26 ^b	—	—	197
SnSe ₂	Nano plate	Methane	1% vol	200	66.7 ^c	78	336	194
SnSe ₂	Pristine monolayer	NO ₂	1	50	60% ^b	142	457	61
SnSe ₂	Thin-film	Methane	200	200	~16% ^c	~75	~615	21
p-SnSe/n-SnSe ₂	Hetero-structure	NO ₂	0.1	RT	30% ^d	—	—	193
SnSe ₂	Thin film	CH ₄	8000	600	300% ^d	—	—	198
SnSe-SnSe ₂	Nanostructured thin film	NO ₂	5	RT	112% ^e	15	10	199
Au/Pd/SnSe ₂	Thin film	NO ₂	5	RT	117% ^e	10	18.7	200

^a $S = (I_g - I_0)/I_0 \times 100$. ^b $R = R_a/R_g$. ^c $R = (R_a - R_g)/R_g \times 100$. ^d $S = (R_g - R_a)/R_a \times 100$. ^e $R = (R_g/R_a) \times 100$, where I_g is the current of device in presence of target gas and I_0 is the current of device in presence of air. R is called the sensor's response, S is called the sensor's sensitivity, R_g is the resistance of the sensor in the presence of target gas, and R_a is the sensor's resistance in the air.

of RhB.¹⁴ Two films SnSe₂ and SnSe₂/Se were made by the facile solvothermal method. The photocatalytic efficiency was measured from the formula mentioned below,

$$\text{Photocatalytic degradation percentage} = \left[\left(1 - \frac{C}{C_0} \right) \times 100 \right] \quad (3)$$

where C is the concentration at time t and C_0 is the concentration at time t . The result showed that SnSe₂/Se film showed better degradation, *i.e.*, around 94% compared to Sn films in 50 min. This increase in efficiency was attributed to the formation of heterojunction and the production of a larger number of electron-hole pairs and a lower recombination rate.

4.5.2 SnSe composites for photocatalytic applications. Zhou Li *et al.* synthesized a composite of SnSe@SnO₂ nanoparticles and used them in the photothermal-photocatalyst mechanism study. The mechanism of electron-hole pair generation is similar to that mentioned earlier.

The photon falls on the surface, and the electron-hole pair is generated. Here, both SnSe and SnO₂ took part in the photothermal and photocatalytic processes synchronously. In the process, the electrons from CB of SnSe travels to the CB of SnO₂.

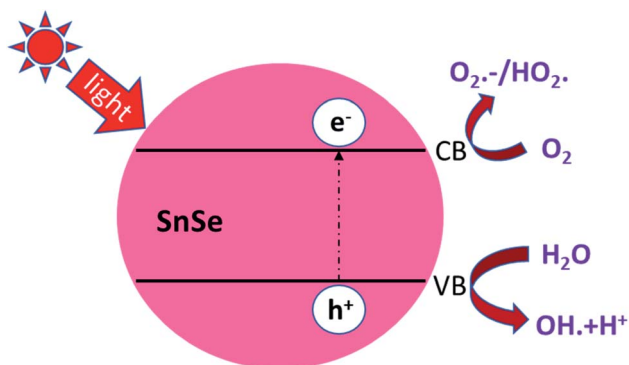


Fig. 15 Indicative representation of photodegradation of RhB dye in the presence of UV light. In a typical experiment, the various dyes' degradation rate is checked using eqn (3). These figures has been adapted/reproduced from ref. 62 with permission from Scielo, copyright 2017".

This increases the oxidation process of photocatalysis and thus improves the photocatalytic behavior. The inner shell SnSe favors the photothermal heat efficiently.²⁰¹ The photocatalytic activity of SnSe@SnO₂ nanocomposite was evaluated through the degradation of methyl orange (MO) experiment under solar light's irradiation. The photocatalytic degradation of MO with SnSe, SnO₂, SnSe@SnO₂, and P25 (commercially named as TiO₂ nanoparticles) was compared using the time *versus* C/C_0 graph. Here, C is the concentration at time t and C_0 is the concentration at time t . It was observed that SnSe@SnO₂ showed a much higher degradation rate than SnSe, SnO₂, and TiO₂. In some other methods, blue SE-2R dye degradation was studied by making a composite.

Karamat *et al.*²⁰² made SnSe composite with LaNdZr₂O₇ and compared the photocatalytic efficiency with SnSe and blank dye. The setup composed a 300 W Xe arc lamp (PLS-SXE300, Beijing Trusttech Co. Ltd). The degradation of blue dye was analyzed at 658 nm using a UV-Vis spectrophotometer. Before the experiment, the stability of blue dye was checked under dark and light conditions. It was confirmed that the dye was stable, and the continuous illumination does not show any loss in spectra. The photocatalytic degradation experiment was then carried out, and it was found that LaNdZr₂O₇/SnSe showed a better degradation efficiency. The maximum absorption intensity at wavelength 658 nm showed a decrease in absorption as the irradiation time increases. SnSe alone has demonstrated degradation efficiency of just 15.5% while mixing it with LaNdZr₂O₇ increased it significantly.

4.6 Storage devices

4.6.1 Battery electrode. Tin-based alloys showed high theoretical capacity and are being considered as a very suitable material for sodium-ion batteries. Tin-based composites have a theoretical storage capacity of nearly 847 mA h g⁻¹. Still, a few other composites like SnSe and SnSe₂ showed even higher values than theoretical ones, with additional capacity contributions from conversion reactions.¹⁴ Considering that one molar SnSe₂ can accommodate 7.75 molar sodium, SnSe₂ has shown a theoretical capacity of 756 mA h g⁻¹.²⁰³ But during the alloying process, the volume of the material expands, which is dangerous for the cycle



stability; hence some composites of carbon are used to reduce this bottleneck. Tin selenide-based battery application has been divided into two categories, namely Li^- ion and Na^+ ion batteries. The recent advances in SnSe_2 for sodium-ion batteries and their results are explained in the coming section.

Zhang *et al.* synthesized SnSe_2 based two-dimensional (2D) nanosheets using the hydrothermal technique. They achieved a theoretical capacity during the first cycle and a stable and reversible specific capacity of 515 mA h g^{-1} at 0.1 A g^{-1} after 100 cycles, which exhibited excellent performance.²⁰⁴ The SnSe and SnSe_2 both have been used for both sodium and lithium-ion battery applications. Chen *et al.* (2018) introduced Cu in SnSe for sodium-ion battery and exhibited a capacity of 330 mA h g^{-1} at 20 A g^{-1} . In a study, Xia *et al.* used electrospun SnSe with carbon nanofibers for lithium and sodium-ion battery applications and showed improved results. Kim *et al.* investigated SnSe alloy as an anode for Na-ion batteries. They exhibited excellent electrochemical performance with a high reversible capacity of 707 mA h g^{-1} and stable performance over 50 cycles.²⁰⁴ In recent time, various studies on SnSe composite has been conducted by multiple research groups to find better results due to the layered structure of SnSe and SnSe_2 materials. Table 9 shows a summary of the above discussions.

Tin selenide-based composites have shown an excellent theoretical capacity. After many cycles, real practical uses of tin-

based anodes in both types of batteries are still very limited. Some structural design, preparation advancement, and morphological development are needed to bring these materials to the production level.

4.6.2 Supercapacitor. Supercapacitors have gained colossal consideration due to their longer operational life, higher power densities, and better safety tolerances than batteries. Zhang *et al.* synthesized tin selenide (SnSe , SnSe_2)-based 2D nanostructures on flexible substrates for supercapacitor applications.⁶³ Ni *et al.* synthesized SnSe based anode using a microwave-assisted method. They exhibited a suitable specific capacitance of 214.3 F g^{-1} at 1 A g^{-1} and rate capability of 182.8 F g^{-1} at 20 A g^{-1} with outstanding cyclic stability.²¹¹ Pandit *et al.* synthesized binder-free SnSe hexagonal nanosheets using a one-pot colloidal method and achieved high-performing anode material for supercapacitors. It showed a specific capacitance of 617.9 F g^{-1} at a scan rate of 2 mV s^{-1} and with fast charge-discharge cycles.¹⁵

4.7 Memory devices

The phase change from amorphous to crystalline induces a sharp change in the optical and electrical properties like reflectivity and resistivity, respectively. This change in the

Table 9 The tin selenide-based composite used for battery application and their cyclic performance

Material used	Ion transfer	Cycle rate	Ref.
Cu doped SnSe	Sodium	Retains a capacity of 330 mA h g^{-1} at 20 A g^{-1} and 304 mA h g^{-1} after 1000 cycles at 5 A g^{-1} ($0.1\text{--}3.0 \text{ V vs. Na/Na}^+$)	205
Electrospun SnSe/C nanofibers	Lithium and sodium	The discharge capacity of 405 mA h g^{-1} at 1000 mA g^{-1} after 500 cycles in lithium-ion battery and 290 mA h g^{-1} at 200 mA g^{-1} after 200 cycles in Na^+ ion battery	60
Tin selenide/N-doped carbon composite	Lithium and sodium	For Li ion-discharge capacity 405 mA h g^{-1} after 500 cycles at a current density of 1000 mA g^{-1} for Na ion-discharge capacity of 290 mA h g^{-1} after 200 cycles at 200 mA g^{-1}	206
SnSe/SnO_2 heterostructure	Lithium	High cyclability having a capacity of 810 mA h g^{-1} after 200 cycles at a current density value of current density of 400 mA g^{-1}	207
Tin diselenide hexagonal nanosheets	Lithium	Specific capacity of 795 mA h g^{-1} after 100 cycles at 100 mA g^{-1}	208
SnSe_2 nanoplate-graphene composites	Lithium	Higher storage capacity of 420 mA h g^{-1} in the first 10 cycles	55
$\text{SnSe}_2/\text{CNTs}$ hybrid nanostructures	Lithium	$\text{SnSe}_2/\text{CNTs}$ electrode showed specific capacity of $457.6 \text{ mA h g}^{-1}$ at 0.1C and $210.3 \text{ mA h g}^{-1}$ after 100 cycles	209
$\text{SnSe}_2/\text{reduced graphene oxide (rGO)}$ composite	Sodium	Exhibits an initial efficiency of 73.7%, showing a high capacity of $402.0 \text{ mA h g}^{-1}$ after 150 cycles at 0.1 A g^{-1} with retention of 86.2%	210



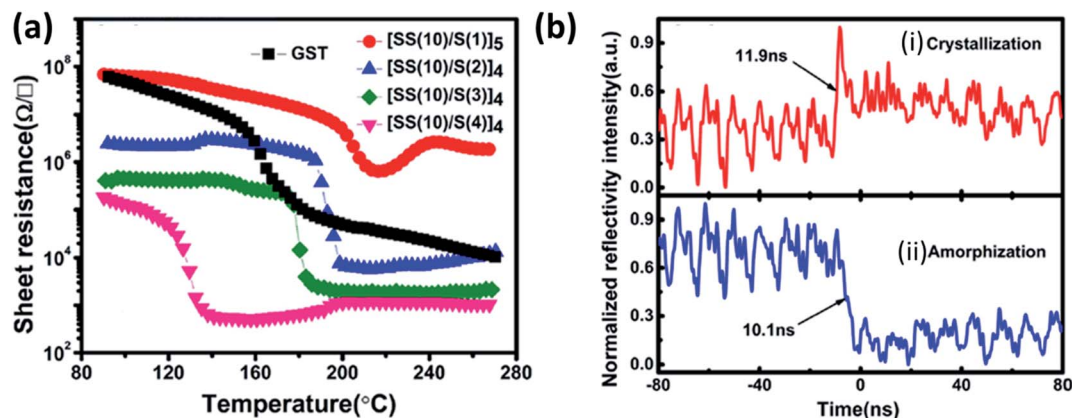


Fig. 16 (a) The temperature dependence of resistance for the SLL SS/S and GST thin films at a constant heating rate of 10 °C min⁻¹, and (b) reversible reflectivity evolution of SLL [SS(10)/S(2)] thin film induced by two consecutive picosecond laser pulses with different fluencies: (i) crystallization process and (ii) amorphization process. These figures has been adapted/reproduced from ref. 215 with permission from AIP, copyright 2016".

property is explored to store the data in the form of a bit. As this transition is reversible, this property is utilized in the rewritable data storage. Fast operating speeds, low power consumption, high energy density, large retention time, *etc.*, are the critical parameters for selecting materials for memory applications. Chalcogenides based materials play the central role in this field due to their very short time of the order of few nanoseconds of phase transition, high resistance ratio of the amorphous to the crystalline phase. Chung *et al.* investigated all phases of the tin selenide (SnSe, SnSe₂, Sn₂Se₃)⁷ for memory application based on phase change and compared the result with the well-studied GeTe.²¹² SnSe₂ is best among all phases, with high crystallization temperature (220 °C), high electrical contrast (resistance ratio of amorphous to crystalline) of 8.3×10^5 , T_g/T_m of 0.52, which showed its potential for the memory application.²¹² Wang *et al.* demonstrated the swift recrystallization time of 20 ns for the solution-processed SnSe₂ thin film,²¹³ comparable to the commercialize Ge₂Sb₂Te₅. Sun *et al.* studied the multilayer Si/SnSe₂ film by varying the Si's thickness in steps of 4 nm for 0–20 nm, with the film's total thickness nearly constant at 100 nm.²¹⁴ The varying thickness of the Si did not show any effect on the crystallization temperature. Still, a remarkable change in the resistance of the film with Si thickness of 12 nm {[Si (12 nm)/SnSe₂ (5 nm)]₆} was observed, *i.e.*, crystalline resistance increased and amorphous resistance decreased, which led to enough (more than three orders of magnitude) electrical contrast to storing information. Simultaneously increasing crystalline resistance reduced the operating current that led to lower power consumption. Multilayer with Si

thickness of 12 nm (of nearly 100 nm total thickness) showed its high potential for the low power operated phase-change random access memory (PCRAM) device because of almost ten year retention time at 83 °C and low activation energy.²¹⁴

The difficulty in phase-change memory (PCM) application of these materials is due to high amorphous resistivity of SnSe₂, fast crystallization speed, low thermal stability, retention of Sb rich alloy, and very low phase change temperature of Sb. Optimized design of the super lattice like (SLL) structure of SnSe₂ and Sb ([SS(10)/S(2)]₄) (SS stands for SnSe₂ and S for Sb with 10 and 2 nm thickness, respectively) proved to be four times better than the commercialized Ge₂Sb₂Te₅ (GST)²¹⁵ SLL's resistance with varying temperature and phase switching times for optimized SLL is shown in Fig. 16(a) and (b). Table 10 shows the comparison of the supremacy of SLL over GST.

4.8 Topological insulators

Topological insulators are the materials that have insulating bulk accompanied by the protected conducting surface states. These surface states provide an impurity and defect insensitive scattering path to the electrons. The unusual metallic nature of an insulator's surface due to the changes in the insulator's topology is the signature of a topological insulator.²¹⁶ These surface states are induced due to the high spin-orbit coupling, inverse parity of bulk band, band inversion,^{217,218} and protected by the time-reversal symmetry.^{16,218} When the crystal symmetry protects the topological insulator's surface, instead of the time-reversal symmetry protecting it, these insulators are called the topological crystalline insulator (TCI).²¹⁷ Sun *et al.*¹⁶ reported

Table 10 The supremacy of SLL over GST

Structure	Crystallization temperature (°C)	Crystallization activation energy (eV)	Year data retention at temperature (°C)	Density variation	SET/RESET operation speed (ns)
Ge ₂ Sb ₂ Te ₅	160	2.37	10 (75)	8.2	17.7/16.5
[SS(10)/S(2)] ₄	185	3.03	10 (116)	3.5	11.9/10.1



that SnSe is a TCI in its native phase (without any pressure and alloying) based on the theoretical calculation using the first principle along with *ab initio* tight-binding modeling. Chen *et al.*²¹⁹ reported the co-existence of topological states and superconducting phase at around 27 GPa of the CsCl type phase of SnSe. *Ab initio* calculation showed that this CsCl phase of SnSe has the topological states' stability against the spin-orbit coupling. SnSe provides the correlation among the superconductivity and topological electronic structures.²¹⁹

5. Conclusion

This review summarizes the tin-selenide material's physical properties, phases, defects, growth mechanisms, deposition methods, and various possible applications. Tin-based compound semiconductors have attracted enormous interest in photovoltaic and optoelectronic applications due to favorable bandgap (1.0–1.3 eV) and higher absorption coefficient (10^5 cm^{-1}). The reported efficiency of tin selenide based solar cells is very low compared to other compound semiconductor materials like CIGS, CdTe, and CZTSSe. The main limiting factor for tin selenide's relative low efficiency might be defects present in absorber materials, interface trap states, and unfavorable hetero-junction interfaces. Moreover, conduction band offset (CBO) between the absorber and buffer layer plays a crucial role. It needs to be engineered in a precise manner for efficient transport of electrons toward a metal electrode. Thus, fine-tuning of elements, precise control over growth conditions, favorable band alignment between absorber and buffer can pave the way towards improving the power conversion efficiency for realizing the development of environmental-friendly low-cost next-generation photovoltaic technology. Tin selenide can be a potential futuristic material due to its versatile applications. It shows leading hands in some applications in its native crystal structure like thermoelectric. Whereas in other fields it shows competitive nature with the other materials. Tin selenide based thermoelectric materials can pave the ways to reuse waste energy. Tin selenide-based photodetector has also exhibited fast response speed with good responsivity and detectivity in the NIR range.

Moreover, flexible photodetector based on tin selenide have also gained attention due to its easy fabrication and fast device response. The device response can be improved by controlling absorber thickness and carrier collection before recombination. Another application based on tin selenide is photocathode in photoelectrochemical water splitting for green hydrogen generation. Lastly, some studies have shown that tin selenide can also be used in gas sensors, storage devices (battery electrode, supercapacitors), memory devices, topological insulators.

Conflicts of interest

Authors declare no conflict of interest exists.

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References

- 1 A. C. Lokhande, I. A. Qattan, C. D. Lokhande and S. P. Patole, *J. Mater. Chem. A*, 2020, **8**, 918–977.
- 2 A. C. Lokhande, P. T. Babar, V. C. Karade, M. G. Gang, V. C. Lokhande, C. D. Lokhande and J. H. Kim, *J. Mater. Chem. A*, 2019, **7**, 17118–17182.
- 3 I. Ali, M. Suhail, Z. A. Allothman and A. Alwarthan, *RSC Adv.*, 2018, **8**, 30125–30147.
- 4 J. Wang, R. Chen, L. Xiang and S. Komarneni, *Ceram. Int.*, 2018, **44**, 7357–7377.
- 5 K. M. Chung, D. Wamwangi, M. Woda, M. Wuttig and W. Bensch, *J. Appl. Phys.*, 2008, **103**, 083523.
- 6 G. Xing, Y. Li, X. Fan, L. Zhang, W. Zheng and D. J. Singh, *J. Appl. Phys.*, 2017, **121**, 225106.
- 7 D. I. Bletskan, *J. Ovonic Res.*, 2005, **1**, 61–69.
- 8 W. Shi, M. Gao, J. Wei, J. Gao, C. Fan, E. Ashalley, H. Li and Z. Wang, *Adv. Sci.*, 2018, **5**, 1700602.
- 9 M. Kang, S. Rathi, I. Lee, L. Li, M. A. Khan, D. Lim, Y. Lee, J. Park, A. T. Pham, A. T. Duong, S. Cho, S. J. Yun and G.-H. Kim, *J. Nanosci. Nanotechnol.*, 2017, **18**, 4243–4247.
- 10 Z. Li, Y. Guo, F. Zhao, C. Nie, H. Li, J. Shi, X. Liu, J. Jiang and S. Zuo, *RSC Adv.*, 2020, **10**, 16749–16755.
- 11 J. Li, H. Zhao, Y. Lei, Q. Yang and Z. Zheng, *Nano*, 2018, **13**, 1850045.
- 12 K. Chung, D. Wamwangi, M. Woda, M. Wuttig and W. Bensch, *J. Appl. Phys.*, 2008, **103**, 083523.
- 13 X. Wang, Y. Liu, J. Dai, Q. Chen, X. Huang and W. Huang, *Chem.–Eur. J.*, 2020, **26**, 3870–3876.
- 14 Z. Li, J. Ding and D. Mitlin, *Acc. Chem. Res.*, 2015, **48**, 1657–1665.
- 15 B. Pandit, C. D. Jadhav, P. G. Chavan, H. S. Tarkas, J. V. Sali, R. B. Gupta and B. R. Sankapal, *IEEE Trans. Power Electron.*, 2020, **35**, 11344–11351.
- 16 Y. Sun, Z. Zhong, T. Shirakawa, C. Franchini, D. Li, Y. Li, S. Yunoki and X. Q. Chen, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2013, **88**(23), 235122.
- 17 V. E. Drozd, I. O. Nikiforova, V. B. Bogevoilov, A. M. Yafyasov, E. O. Filatova and D. Papazoglou, *J. Phys. D: Appl. Phys.*, 2009, **42**(12), 125306.
- 18 Z. Zainal, N. Saravanan, K. Anuar, M. Z. Hussein and W. M. M. Yunus, *Mater. Sci. Eng., B*, 2004, **107**, 181–185.
- 19 M. R. Pallavolu, V. R. Minnam Reddy, P. R. Guddeti and C. Park, *J. Mater. Sci.: Mater. Electron.*, 2019, **30**, 15980–15988.
- 20 J. S. Narro-Rios, M. Ramachandran, D. Martínez-Escobar and A. Sánchez-Juarez, *J. Semicond.*, 2013, **34**, 013001.
- 21 K. Assili, O. Gonzalez, K. Alouani and X. Vilanova, *Arabian J. Chem.*, 2020, **13**, 1229–1246.
- 22 X. Shi, J. Zou and Z. Chen, *Chem. Rev.*, 2020, **120**, 7399–7515.
- 23 E. P. Mukhokosi, G. V. S. Manohar, T. Nagao, S. B. Krupanidhi and K. K. Nanda, *Micromachines*, 2020, **11**, 750.



- 24 V. Q. Nguyen, J. Kim and S. Cho, *J. Korean Phys. Soc.*, 2018, **72**, 841–857.
- 25 Z. Chen, X. Shi, L. Zhao and J. Zou, *Prog. Mater. Sci.*, 2018, **97**, 283–346.
- 26 X.-L. Shi, W.-Y. Chen, X. Tao, J. Zou and Z.-G. Chen, *Mater. Horiz.*, 2020, **7**, 3065–3096.
- 27 B. Qin, W. He and L. D. Zhao, *J. Mater.*, 2020, **6**, 671–676.
- 28 M. Taniguchi, R. L. Johnson, J. Ghijsen and M. Cardona, *Phys. Rev. B*, 1990, **42**, 3634–3643.
- 29 R. E. Abutbul, E. Segev, S. Samuha, L. Zeiri, V. Ezersky, G. Makov and Y. Golan, *CrystEngComm*, 2016, **18**, 1918–1923.
- 30 E. P. Mukhokosi, S. B. Krupanidhi and K. K. Nanda, *Sci. Rep.*, 2017, **7**, 1–10.
- 31 T. Chattopadhyay, J. Pannetier and H. G. Von Schnering, *J. Phys. Chem. Solids*, 1986, **47**, 879–885.
- 32 L. D. Zhao, S. H. Lo, Y. Zhang, H. Sun, G. Tan, C. Uher, C. Wolverton, V. P. Dravid and M. G. Kanatzidis, *Nature*, 2014, **508**, 373–377.
- 33 Y. Ding, B. Xiao, G. Tang and J. Hong, *J. Phys. Chem. C*, 2017, **121**, 225–236.
- 34 M. Zhou, X. Chen, M. Li and A. Du, *J. Mater. Chem. C*, 2017, **5**, 1247–1254.
- 35 J. M. Gonzalez and I. I. Oleynik, *Phys. Rev. B*, 2016, **94**, 1–10.
- 36 L. Hao, Z. Wang, H. Xu, K. Yan, S. Dong, H. Liu, Y. Du, Y. Wu, Y. Liu and M. Dong, *2D Mater.*, 2019, **6**(3), 034004.
- 37 G. Xing, Y. Li, X. Fan, L. Zhang, W. Zheng and D. J. Singh, *J. Appl. Phys.*, 2017, **121**, 225106.
- 38 L. Chunhong, G. Donglin, K. Li, B. Saho, D. Chen, Y. Ma and J. Sun, *Phys. B Phys. Condens. Matter*, 2018, **530**, 264–269.
- 39 X. Shi, X. Tao, J. Zou and Z. Chen, *Adv. Sci.*, 2020, **7**, 1902923.
- 40 Y. Huang, C. Wang, X. Chen, D. Zhou, J. Du, S. Wang and L. Ning, *RSC Adv.*, 2017, **7**, 27612–27618.
- 41 G. Duvjir, T. Min, T. Thi Ly, T. Kim, A. T. Duong, S. Cho, S. H. Rhim, J. Lee and J. Kim, *Appl. Phys. Lett.*, 2017, **110**(26), 262106.
- 42 B. B. Nariya, A. K. Dasadia, M. K. Bhayani, A. J. Patel and A. R. Jani, *Chalcogenide Lett.*, 2009, **6**, 549–554.
- 43 G. Shanmugam, U. P. Deshpande, A. Sharma, P. M. Shirage and P. A. Bhohe, *J. Phys. Chem. C*, 2018, **122**, 13182–13192.
- 44 Y. Zhou, W. Li, M. Wu, L. D. Zhao, J. He, S. H. Wei and L. Huang, *Phys. Rev. B*, 2018, **97**(24), 245202.
- 45 L. Song, J. Zhang and B. B. Iversen, *J. Mater. Chem. A*, 2019, **7**, 17981–17986.
- 46 R. Indirajith, T. P. Srinivasan, K. Ramamurthi and R. Gopalakrishnan, *Curr. Appl. Phys.*, 2010, **10**, 1402–1406.
- 47 Y. Gong, C. Chang, W. Wei, J. Liu, W. Xiong, S. Chai, D. Li, J. Zhang and G. Tang, *Scr. Mater.*, 2018, **147**, 74–78.
- 48 D. Martínez-Escobar, M. Ramachandran, A. Sánchez-Juárez and J. S. Narro Rios, *Thin Solid Films*, 2013, **535**, 390–393.
- 49 F. Davitt, H. G. Manning, F. Robinson, S. L. Hawken, S. Biswas, N. Petkov, M. van Druenen, J. J. Boland, G. Reid and J. D. Holmes, *Adv. Mater. Interfaces*, 2020, **7**(16), 2000474.
- 50 V. P. Bhatt, K. Gireesan and G. R. Pandya, *J. Cryst. Growth*, 1989, **96**, 649–651.
- 51 A. T. Duong, V. Q. Nguyen, G. Duvjir, V. T. Duong, S. Kwon, J. Y. Song, J. K. Lee, J. E. Lee, S. Park, T. Min, J. Lee, J. Kim and S. Cho, *Nat. Commun.*, 2016, **7**(1), 1–6.
- 52 J. Li, W. Liu, C. Chen, X. Zhao, Z. Qiu, H. Xu, F. Sheng, Q. Hu, Y. Zheng, M. Lin, S. J. Pennycook, C. Su and J. Lu, *J. Mater. Chem. A*, 2019, **7**, 23958–23963.
- 53 J. Jiang, C. P. Y. Wong, J. Zou, S. Li, Q. Wang, J. Chen, D. Qi, H. Wang, G. Eda, D. H. C. Chua, Y. Shi, W. Zhang and A. T. S. Wee, *2D Mater.*, 2017, **4**(2), 021026.
- 54 Z. Wang, F. Li, J. Guo, C. Ma, Y. Song, Z. He, J. Liu, Y. Zhang, D. Li and H. Zhang, *Adv. Opt. Mater.*, 2020, **8**(16), 1902183.
- 55 J. Choi, J. Jin, I. G. Jung, J. M. Kim, H. J. Kim and S. U. Son, *Chem. Commun.*, 2011, **47**, 5241–5243.
- 56 X. Li, L. Li, H. Zhao, S. Ruan, W. Zhang, P. Yan, Z. Sun, H. Liang and K. Tao, *Nanomaterials*, 2019, **9**(9), 1324.
- 57 X. Zhou, L. Gan, W. Tian, Q. Zhang, S. Jin, H. Li, Y. Bando, D. Golberg and T. Zhai, *Adv. Mater.*, 2015, **27**(48), 8035–8041.
- 58 D. W. Ma and C. Cheng, *J. Nanosci. Nanotechnol.*, 2013, **13**, 4433–4436.
- 59 M. A. Franzman, C. W. Schlenker, M. E. Thompson and R. L. Brutchey, *J. Am. Chem. Soc.*, 2010, **132**, 4060–4061.
- 60 J. Xia, Y. Yuan, H. Yan, J. Liu, Y. Zhang, L. Liu, S. Zhang, W. Li, X. Yang, H. Shu, X. Wang and G. Cao, *J. Power Sources*, 2020, **449**, 227559.
- 61 Ó. L. C. Moreira, W. Y. Cheng, H. R. Fuh, W. C. Chien, W. Yan, H. Fei, H. Xu, D. Zhang, Y. Chen, Y. Zhao, Y. Lv, G. Wu, C. Lv, S. K. Arora, C. Ó Coileáin, C. Heng, C. R. Chang and H. C. Wu, *ACS Sens.*, 2019, **4**, 2546–2552.
- 62 L. Cheng, D. Li, X. Dong, Q. Ma, W. Yu, X. Wang, H. Yu, J. Wang, G. Liu, L. Cheng, D. Li, X. Dong, Q. Ma, W. Yu, X. Wang, H. Yu, J. Wang and G. Liu, *Mater. Res.*, 2017, **20**, 1748–1755.
- 63 C. Zhang, H. Yin, M. Han, Z. Dai, H. Pang, Y. Zheng, Y. Q. Lan, J. Bao and J. Zhu, *ACS Nano*, 2014, **8**, 3761–3770.
- 64 X. Feng, Y. Hu, J. Zhai, C. Wang, S. Song and Z. Song, *J. Appl. Phys.*, 2014, **115**, 204303.
- 65 M. Kumar, S. Rani, A. Pandey, K. S. Gour, S. Husale, P. Singh and V. N. Singh, *J. Alloys Compd.*, 2020, **838**, 155384.
- 66 M. Powalla, S. Paetel, E. Ahlswede, R. Wuerz, C. D. Wessendorf and T. Magorian Friedlmeier, *Appl. Phys. Rev.*, 2018, **5**, 041602.
- 67 X. Cui, K. Sun, J. Huang, J. S. Yun, C. Y. Lee, C. Yan, H. Sun, Y. Zhang, C. Xue, K. Eder, L. Yang, J. M. Cairney, J. Seidel, N. J. Ekins-Daukes, M. Green, B. Hoex and X. Hao, *Energy Environ. Sci.*, 2019, **12**, 2751–2764.
- 68 X. Li, D. Zhuang, N. Zhang, M. Zhao, X. Yu, P. Liu, Y. Wei and G. Ren, *J. Mater. Chem. A*, 2019, **7**, 9948–9957.
- 69 D. H. Son, S. H. Kim, S. Y. Kim, Y. I. Kim, J. H. Sim, S. N. Park, D. H. Jeon, D. K. Hwang, S. J. Sung, J. K. Kang, K. J. Yang and D. H. Kim, *J. Mater. Chem. A*, 2019, **7**, 25279–25289.
- 70 O. P. Singh, K. S. Gour, R. Parmar and V. N. Singh, *J. Nanosci. Nanotechnol.*, 2018, **18**, 7670–7681.
- 71 B. Pejjai, V. R. Minnam Reddy, K. Seku, M. R. Pallavolu and C. Park, *New J. Chem.*, 2018, **42**, 4843–4853.



- 72 S. Rühle, *Sol. Energy*, 2016, **130**, 139–147.
- 73 D. V. Shinde, S. K. Min, M. M. Sung, N. K. Shrestha, R. S. Mane and S. H. Han, *Mater. Lett.*, 2014, **115**, 244–247.
- 74 N. R. Mathews, *Sol. Energy*, 2012, **86**, 1010–1016.
- 75 J. P. Singh and R. K. Bedi, *Jpn. J. Appl. Phys.*, 1990, **29**, L792–L793.
- 76 S. Delice, M. Isik, H. H. Gullu, M. Terlemezoglu, O. Bayrakli Surucu, M. Parlak and N. M. Gasanly, *J. Phys. Chem. Solids*, 2019, **131**, 22–26.
- 77 V. R. Minnam Reddy, G. Lindwall, B. Pejjai, S. Gedi, T. R. R. Kotte, M. Sugiyama, Z. K. Liu and C. Park, *Sol. Energy Mater. Sol. Cells*, 2018, **176**, 251–258.
- 78 M. Nakamura, K. Yamaguchi, Y. Kimoto, Y. Yasaki, T. Kato and H. Sugimoto, *IEEE J. Photovoltaics*, 2019, **9**, 1863–1867.
- 79 J. Park, H. Yoo, V. Karade, K. S. Gour, E. Choi, M. Kim, X. Hao, S. J. Shin, J. Kim, H. Shim, D. Kim, J. H. Kim, J. Yun and J. hyeok Kim, *J. Mater. Chem. A*, 2020, **8**, 14538–14544.
- 80 K. S. Gour, R. Parmar, R. Kumar and V. N. Singh, *J. Nanosci. Nanotechnol.*, 2019, **20**, 3622–3635.
- 81 M. Dhankhar, O. Pal Singh and V. N. Singh, *Renewable Sustainable Energy Rev.*, 2014, **40**, 214–223.
- 82 P. A. Fernandes, M. G. Sousa, P. M. P. Salomé, J. P. Leitão and A. F. Da Cunha, *CrystEngComm*, 2013, **15**, 10278–10286.
- 83 K. F. Abd El-Rahman, A. A. A. Darwish and E. A. A. El-Shazly, *Mater. Sci. Semicond. Process.*, 2014, **25**, 123–129.
- 84 N. E. Makori, I. A. Amatalo, P. M. Karimi and W. K. Njoroge, *Int. J. Energy Eng.*, 2015, **5**, 1–4.
- 85 W. Wang, M. T. Winkler, O. Gunawan, T. Gokmen, T. K. Todorov, Y. Zhu and D. B. Mitzi, *Adv. Energy Mater.*, 2014, **4**, 1301465.
- 86 Q. Jin, S. Jiang, Y. Zhao, D. Wang, J. Qiu, D. M. Tang, J. Tan, D. M. Sun, P. X. Hou, X. Q. Chen, K. Tai, N. Gao, C. Liu, H. M. Cheng and X. Jiang, *Nat. Mater.*, 2019, **18**, 62–68.
- 87 L. Hu, Y. Zhang, H. Wu, Y. Liu, J. Li, J. He, W. Ao, F. Liu, S. J. Pennycook and X. Zeng, *Adv. Funct. Mater.*, 2018, **28**(35), 1803617.
- 88 M. Markov, X. Hu, H. C. Liu, N. Liu, S. J. Poon, K. Esfarjani and M. Zebarjadi, *Sci. Rep.*, 2018, **8**, 9876.
- 89 H. Nagai, H. Hamada, K. Hayashi and Y. Miyazaki, *J. Electron. Mater.*, 2019, **48**, 1902–1908.
- 90 J. Yang, G. Zhang, G. Yang, C. Wang and Y. X. Wang, *J. Alloys Compd.*, 2015, **644**, 615–620.
- 91 G. Shi and E. Kioupakis, *J. Appl. Phys.*, 2015, **117**, 065103.
- 92 K. Kutorasinski, B. Wiendlocha, S. Kaprzyk and J. Tobola, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2015, **91**(20), 205201.
- 93 L. Zhao, G. Tan, S. Hao, J. He, Y. Pei, H. Chi, H. Wang, S. Gong, H. Xu, V. P. Dravid, C. Uher, G. J. Snyder, C. Wolverton and M. G. Kanatzidis, *Science*, 2016, **351**(6269), 141–144.
- 94 P. C. Wei, S. Bhattacharya, Y. F. Liu, F. Liu, J. He, Y. H. Tung, C. C. Yang, C. R. Hsing, D. L. Nguyen, C. M. Wei, M. Y. Chou, Y. C. Lai, T. L. Hung, S. Y. Guan, C. S. Chang, H. J. Wu, C. H. Lee, W. H. Li, R. P. Hermann, Y. Y. Chen and A. M. Rao, *ACS Omega*, 2019, **4**(3), 5442–5450.
- 95 M. Jin, Z. Tang, J. Jiang, R. Zhang, L. Zhou, S. Zhao, Y. Chen, Y. Chen, X. Wang and R. Li, *Mater. Res. Bull.*, 2020, **126**, 110819.
- 96 C. Chang, M. Wu, D. He, Y. Pei, C. F. Wu, X. Wu, H. Yu, F. Zhu, K. Wang, Y. Chen, L. Huang, J. F. Li, J. He and L. D. Zhao, *Science*, 2018, **360**, 778–783.
- 97 S. Patel, S. H. Chaki and P. C. Vinodkumar, *Mater. Res. Express*, 2019, **6**(8), 085910.
- 98 A. Jayaraman, M. Molli and V. Kamiseti, *AIP Conf. Proc.*, 2015, **1665**(1), 110046.
- 99 Y. Tang, L. Shen, Z. Chen, L. Sun, W. Liu, J. Liu and S. Deng, *Phys. B Phys. Condens. Matter*, 2019, **570**, 128–132.
- 100 M. Jin, H. Shao, H. Hu, D. Li, J. Xu, G. Liu, H. Shen, J. Xu, H. Jiang and J. Jiang, *J. Cryst. Growth*, 2017, **460**, 112–116.
- 101 K. Peng, H. Wu, Y. Yan, L. Guo, G. Wang, X. Lu and X. Zhou, *J. Mater. Chem. A*, 2017, **5**, 14053–14060.
- 102 Y. K. Lee, Z. Luo, S. P. Cho, G. Mercouri, Y. K. Lee, Z. Luo, S. P. Cho, M. G. Kanatzidis and I. Chung, *Joule*, 2019, **3**, 719–731.
- 103 M. Zhang, D. Wang, C. Chang, T. Lin, K. Wang and L. D. Zhao, *J. Mater. Chem. C*, 2019, **7**, 10507–10513.
- 104 S. Sassi, C. Candolfi, J. Vaney, V. Ohorodniichuk, P. Masschelein, A. Dauscher, B. Lenoir, S. Sassi, C. Candolfi, J. Vaney, V. Ohorodniichuk, P. Masschelein and A. Dauscher, *Appl. Phys. Lett.*, 2014, **104**(21), 212105.
- 105 M. Kumar, S. Rani, Y. Singh and V. N. Singh, *J. Nanosci. Nanotechnol.*, 2019, **20**, 3636–3646.
- 106 X. Shi, K. Zheng, M. Hong, W. Liu, R. Moshwan, Y. Wang, X. Qu, Z. G. Chen and J. Zou, *Chem. Sci.*, 2018, **9**, 7376–7389.
- 107 L. D. Zhao, C. Chang, G. Tan and M. G. Kanatzidis, *Energy Environ. Sci.*, 2016, **9**, 3044–3060.
- 108 J. Gainza, A. Alonso, N. M. Nemes, J. Gainza, F. Serrano-sa, F. Simon, L. Marti, A. Alonso and N. M. Nemes, *Cell reports Phys. Sci.*, 2020, **1**(12), 100263.
- 109 J. Y. Cho, M. Siyar, S. H. Bae, J. S. Mun, M. Y. Kim, S. H. Hong and C. Park, *Bull. Mater. Sci.*, 2020, **43**, 63.
- 110 E. K. Chere, Q. Zhang, K. Dahal, F. Cao, J. Mao and Z. Ren, *J. Mater. Chem. A*, 2016, **4**, 1848–1854.
- 111 Y. X. Chen, Z. H. Ge, M. Yin, D. Feng, X. Q. Huang, W. Zhao and J. He, *Adv. Funct. Mater.*, 2016, **26**, 6836–6845.
- 112 X. Shi, A. Wu, T. Feng, K. Zheng, W. Liu, Q. Sun, M. Hong, S. T. Pantelides, Z. G. Chen and J. Zou, *Adv. Energy Mater.*, 2019, **9**(11), 1803242.
- 113 C. H. Chien, C. C. Chang, C. L. Chen, C. M. Tseng, Y. R. Wu, M. K. Wu, C. H. Lee and Y. Y. Chen, *RSC Adv.*, 2017, **7**, 34300–34306.
- 114 Y. Luo, S. Cai, X. Hua, H. Chen, Q. Liang, C. Du, Y. Zheng, J. Shen, J. Xu, C. Wolverton, V. P. Dravid, Q. Yan and M. G. Kanatzidis, *Adv. Energy Mater.*, 2019, **9**(40), 1901800.
- 115 S. Wang, X. Su, T. P. Bailey, T. Hu, Z. Zhang, G. Tan, Y. Yan, W. Liu, C. Uher and X. Tang, *RSC Adv.*, 2019, **9**, 7115–7122.
- 116 X. Zhang, Y. Wang, G. Zhang, C. Wang and Y. Yan, *J. Alloys Compd.*, 2019, **772**, 1061–1066.
- 117 F. Chu, Q. Zhang, Z. Zhou, D. Hou, L. Wang and W. Jiang, *J. Alloys Compd.*, 2018, **741**, 756–764.



- 118 W. Lu, S. Li, R. Xu, J. Zhang, D. Li, Z. Feng, Y. Zhang and G. Tang, *ACS Appl. Mater. Interfaces*, 2019, **11**(48), 45133–45141.
- 119 G. Yang, L. Sang, M. Li, S. M. Kazi Nazrul Islam, Z. Yue, L. Liu, J. Li, D. R. G. Mitchell, N. Ye and X. Wang, *ACS Appl. Mater. Interfaces*, 2020, **12**, 12910–12918.
- 120 Q. K. Zhang, S. T. Ning, N. Qi, Z. Q. Chen, X. F. Tang and Z. Y. Chen, *J. Appl. Phys.*, 2019, **125**(22), 225109.
- 121 S. Chandra and K. Biswas, *J. Am. Chem. Soc.*, 2019, **141**, 6141–6145.
- 122 K. S. Urmila, T. A. Namitha, J. Rajani, R. R. Philip and B. Pradeep, *J. Semicond.*, 2016, **37**(9), 093002.
- 123 P. K. Nair, A. K. Martínez, A. R. G. Angelmo, E. B. Salgado and M. T. S. Nair, *Semicond. Sci. Technol.*, 2018, **33**(3), 035004.
- 124 C. H. Suen, D. Shi, Y. Su, Z. Zhang, C. H. Chan, X. Tang, Y. Li, K. H. Lam, X. Chen, B. L. Huang, X. Y. Zhou and J. Y. Dai, *J. Mater.*, 2017, **3**, 293–298.
- 125 M. R. Burton, T. Liu, J. Mcgettrick, S. Mehraban, J. Baker, A. Pockett, T. Watson, O. Fenwick and M. J. Carnie, *Adv. Mater.*, 2018, **30**(31), 1801357.
- 126 L. Song, J. Zhang and B. B. Iversen, *J. Mater. Chem. A*, 2019, **7**, 17981–17986.
- 127 Z. Yan, Y. Zhao, D. Liu, Y. Zhang, C. Xue, Z. Zhang, Y. Zheng and J. Cui, *RSC Adv.*, 2020, **10**, 43840–43846.
- 128 P. Wei, S. Bhattacharya, Y.-F. Liu, F. Liu, J. He, Y. Tung, C.-C. Yang, C.-R. Hsing, D.-L. Nguyen, C.-M. Wei, M.-Y. Chou, Y.-C. Lai, T.-L. Hung, S.-Y. Guan, C.-S. Chang, H.-J. Wu, C.-H. Lee, W.-H. Li, R. P. Hermann, Y.-Y. Chen and A. M. Rao, *ACS Omega*, 2019, **4**, 5442–5450.
- 129 S. Sassi, C. Candolfi, J.-B. Vaney, V. Ohorodniichuk, P. Masschelein, A. Dauscher and B. Lenoir, *Appl. Phys. Lett.*, 2014, **104**, 212105.
- 130 J. Llorca, D. Cadavid and M. Iba, *ACS Appl. Mater. Interfaces*, 2020, **12**, 27104–27111.
- 131 S. Chandra, P. Dutta and K. Biswas, *ACS Appl. Energy Mater.*, 2020, **3**, 9051–9057.
- 132 B. Cai, H. L. Zhuang, H. Tang and J. F. Li, *Nano Energy*, 2019, 104393.
- 133 J. Cha, C. Zhou, Y. K. Lee, S. P. Cho and I. Chung, *ACS Appl. Mater. Interfaces*, 2019, **11**, 21645–21654.
- 134 Y. K. Lee, Z. Luo, S. P. Cho, M. G. Kanatzidis and I. Chung, *Joule*, 2019, **3**, 719–731.
- 135 C. Ho, W. Lin, L. Chao, K. Lee, J. Inagaki and H. Hsueh, *ACS Appl. Energy Mater.*, 2020, **3**(5), 4896–4905.
- 136 Y. Zhong, L. Zhang, V. Linseis, B. Qin, W. Chen, L. Zhao and H. Zhu, *Nano Energy*, 2020, **72**, 104742.
- 137 B. Z. Sun, Z. Ma, C. He and K. Wu, *Phys. Chem. Chem. Phys.*, 2015, **17**, 29844–29853.
- 138 S. Saha, A. Banik and K. Biswas, *Chem.–Eur. J.*, 2016, **22**, 15634–15638.
- 139 P. Xu, T. Fu, J. Xin, Y. Liu, P. Ying, X. Zhao, H. Pan and T. Zhu, *Sci. Bull.*, 2017, **62**, 1663–1668.
- 140 Y. Luo, Y. Zheng, Z. Luo, S. Hao, C. Du, Q. Liang, Z. Li, K. A. Khor, K. Hippalgaonkar, J. Xu, Q. Yan, C. Wolverton and M. G. Kanatzidis, *Adv. Energy Mater.*, 2018, **8**, 2–7.
- 141 F. Li, Z. Zheng, Y. Li, W. Wang, J. F. Li, B. Li, A. Zhong, J. Luo and P. Fan, *J. Mater. Sci.*, 2017, **52**, 10506–10516.
- 142 Y. Wu, W. Li, A. Faghaninia, Z. Chen, J. Li, X. Zhang, B. Gao, S. Lin, B. Zhou, A. Jain and Y. Pei, *Mater. Today Phys.*, 2017, **3**, 127–136.
- 143 Y. Shu, X. Su, H. Xie, G. Zheng, W. Liu, Y. Yan, T. Luo, X. Yang, D. Yang, C. Uher and X. Tang, *ACS Appl. Mater. Interfaces*, 2018, **10**, 15793–15802.
- 144 C. Liu, Z. Huang, D. Wang, X. Wang, L. Miao, X. Wang, S. Wu, N. Toyama, T. Asaka, J. Chen, E. Nishibori and L.-D. Zhao, *J. Mater. Chem. A*, 2019, **7**(16), 9761–9772.
- 145 J. Wang, X. Jia, S. Lou, G. Li and S. Zhou, *ACS Omega*, 2020, **5**(21), 12409–12414.
- 146 D. Yin, Y. Liu, C. Dun, D. L. Carroll and M. T. Swihart, *Nanoscale*, 2018, **10**, 2533–2541.
- 147 J. Chen, D. M. Hamann, D. Choi, N. Poudel, L. Shen, L. Shi, D. C. Johnson and S. Cronin, *Nano Lett.*, 2018, **18**, 6876–6881.
- 148 A. T. Duong, D. L. Nguyen, M. N. Nguyen, T. M. H. Nguyen, A. D. Nguyen, A. T. Pham, F. Ullah, Z. Tahir, Y. S. Kim, D. Q. Trung, T. Nguyen, H. Van Bui, R. Das, P. T. Huy and S. Cho, *Mater. Res. Express*, 2019, **6**, 066420.
- 149 X. Zhang, D. Liu, L. Yang, L. Zhou and T. You, *J. Mater. Chem. A*, 2015, **3**, 10031–10037.
- 150 Y. Li, B. He, J. P. Heremans, J. Zhao, B. He and J. P. Heremans, *J. Alloys Compd.*, 2016, **669**, 224–231.
- 151 D. Ibrahim, J. Vaney, S. Sassi, C. Candolfi, V. Ohorodniichuk, P. Levinsky, C. Semprimoschnig, A. Dauscher and B. Lenoir, *Appl. Phys. Lett.*, 2017, **110**, 032103.
- 152 P. C. Wei, S. Bhattacharya, J. He, S. Neeleshwar, R. Podila, Y. Y. Chen and A. M. Rao, *Nature*, 2016, **539**, E1–E2.
- 153 R. Gupta and C. Bera, *Nano Express*, 2020, **1**, 030035.
- 154 N. E. Makori, I. A. Amatalo, P. M. Karimi and W. K. Njoroge, *Am. J. Condens. Matter Phys.*, 2014, **4**, 87–90.
- 155 N. Ghobadi and E. Gholami Hatam, *Opt. Quantum Electron.*, 2019, **51**(8), 269.
- 156 A. S. Pawbake, S. R. Jadkar and D. J. Late, *Mater. Res. Express*, 2016, **3**, 105038.
- 157 L. Hao, H. Xu, S. Dong, Y. Du, L. Luo, C. Zhang, H. Liu, Y. Wu and Y. Liu, *IEEE Electron Device Lett.*, 2018, **40**(1), 55–58.
- 158 J. Yao, Z. Zheng and G. Yang, *Adv. Funct. Mater.*, 2017, **27**, 1701823.
- 159 J. Liu, Q. Huang, K. Zhang, Y. Xu, M. Guo, Y. Qian, Z. Huang, F. Lai and L. Lin, *Nanoscale Res. Lett.*, 2017, **12**, 259.
- 160 H. Yao, S. Luo, G. S. Duesberg, X. Qi, D. Lu, C. Yue and J. Zhong, *AIP Adv.*, 2018, **8**, 075123.
- 161 Y. Zhong, L. Zhang, M. Sun, M. Wang, W. Chen, S. Lin, D. Xie and H. Zhu, *Mater. Today Energy*, 2019, **12**, 418–425.
- 162 L. Hao, Y. Du, Z. Wang, Y. Wu, H. Xu, S. Dong, H. Liu, Y. Liu, Q. Xue, Z. Han, K. Yan and M. Dong, *Nanoscale*, 2020, **12**, 7358–7365.
- 163 B. Ouyang, C. Chang, L. Zhao, Z. Lin and Y. Yang, *Nano Energy*, 2019, **66**, 104111.



- 164 D. Zheng, H. Fang, M. Long, F. Wu, P. Wang, F. Gong, X. Wu, J. C. Ho, L. Liao and W. Hu, *ACS Nano*, 2018, **12**, 7239–7245.
- 165 A. S. Pawbake, S. R. Jadhkar and D. J. Late, *Mater. Res. Express*, 2016, **3**, 105038.
- 166 J. Yao, Z. Zheng and G. Yang, *Adv. Funct. Mater.*, 2017, **27**, 1701823.
- 167 L. Hao, Z. Wang, H. Xu, K. Yan, S. Dong, H. Liu, Y. Du, Y. Wu, Y. Liu and M. Dong, *2D Mater*, 2019, **6**, 034004.
- 168 L. Hao, Y. Du, Z. Wang, Y. Wu, H. Xu, S. Dong, H. Liu, Y. Liu, Q. Xue, Z. Han, K. Yan and M. Dong, *Nanoscale*, 2020, **12**, 7358–7365.
- 169 X.-Z. Li, Y.-F. Wang, J. Xia and X.-M. Meng, *Nanoscale Adv.*, 2019, **1**, 2606–2611.
- 170 Z. Jia, J. Xiang, F. Wen, R. Yang, C. Hao and Z. Liu, *ACS Appl. Mater. Interfaces*, 2016, **8**(7), 4781–4788.
- 171 K. Patel, P. Chauhan, A. B. Patel, G. K. Solanki, K. D. Patel and V. M. Pathak, *ACS Appl. Nano Mater.*, 2020, **3**(11), 11143–11151.
- 172 H. Xu, L. Hao, H. Liu, S. Dong, Y. Wu, Y. Liu, B. Cao, Z. Wang, C. Ling, S. Li, Z. Xu, Q. Xue and K. Yan, *ACS Appl. Mater. Interfaces*, 2020, **12**, 35250–35258.
- 173 X. Zhou, L. Gan, W. Tian, Q. Zhang, S. Jin, H. Li, Y. Bando, D. Golberg and T. Zhai, *Adv. Mater.*, 2015, **27**, 8035–8041.
- 174 P. Yu, X. Yu, W. Lu, H. Lin, L. Sun, K. Du, F. Liu, W. Fu, Q. Zeng, Z. Shen, C. Jin, Q. J. Wang and Z. Liu, *Adv. Funct. Mater.*, 2016, **26**, 137–145.
- 175 K. Murali and K. Majumdar, *IEEE Trans. Electron Devices*, 2018, **65**, 4141–4148.
- 176 X. Zhou, N. Zhou, C. Li, H. Song, Q. Zhang, X. Hu, L. Gan, H. Li, J. Lü, J. Luo, J. Xiong and T. Zhai, *2D Mater*, 2017, **4**, 025048.
- 177 E. P. Mukhokosi, S. B. Krupanidhi and K. K. Nanda, *Phys. Status Solidi A*, 2018, **215**, 1800470.
- 178 E. P. Mukhokosi, B. Roul, S. B. Krupanidhi and K. K. Nanda, *ACS Appl. Mater. Interfaces*, 2019, **11**, 6184–6194.
- 179 W. Gao, Z. Zheng, Y. Li, Y. Zhao, L. Xu, H. Deng and J. Li, *Nanoscale*, 2019, **11**, 13309–13317.
- 180 M. Krishna, S. Kallatt and K. Majumdar, *Nanotechnology*, 2018, **29**, 035205.
- 181 X. Zhou, N. Zhou, C. Li, H. Song, Q. Zhang, X. Hu, L. Gan, H. Li, J. Lü, J. Luo, J. Xiong and T. Zhai, *2D Mater.*, 2017, **4**, 025048.
- 182 E. P. Mukhokosi, B. Roul, S. B. Krupanidhi and K. K. Nanda, *ACS Appl. Mater. Interfaces*, 2019, **11**, 6184–6194.
- 183 W. Gao, Z. Zheng, Y. Li, Y. Zhao, L. Xu, H. Deng and J. Li, *Nanoscale*, 2019, **11**, 13309–13317.
- 184 J. Wang, G. F. Yang, J. J. Xue, J. M. Lei, D. J. Chen, H. Lu, R. Zhang and Y. D. Zheng, *IEEE Electron Device Lett.*, 2018, **39**, 599–602.
- 185 T. Liu, H. Qin, D. Yang and G. Zhang, *Coatings*, 2019, **9**(6), 390.
- 186 I. D. Brown, *J. Solid State Chem.*, 1974, **11**, 214–233.
- 187 C. Feng, H. Qin, D. Yang and G. Zhang, *Materials*, 2019, **12**(4), 676.
- 188 X. He, H. Shen, W. Wang, Z. Wang, B. Zhang and X. Li, *J. Alloys Compd.*, 2013, **556**, 86–93.
- 189 G. Shi and E. Kioupakis, *Nano Lett.*, 2015, **15**, 6926–6931.
- 190 I. Lefebvre, M. A. Szymanski, J. Olivier-Fourcade and J. C. Jumas, *Phys. Rev. B*, 1998, **58**, 1896–1906.
- 191 J.-J. Wang, A.-F. Lv, Y.-Q. Wang, B. Cui, H.-J. Yan, J.-S. Hu, W.-P. Hu, Y.-G. Guo and L.-J. Wan, *Sci. Rep.*, 2013, **3**(1), 2613.
- 192 L. Lee, C. W. Chen, A. Manikandan, S. H. Lee, Z. M. Wang and Y. L. Chueh, *Nano Energy*, 2018, **44**, 419–429.
- 193 X. Wang, Y. Liu, J. Dai, Q. Chen, X. Huang and W. Huang, *Chem.–Eur. J.*, 2020, **26**, 3870–3876.
- 194 M. Chen, Z. Li, W. Li, C. Shan, W. Li, K. Li, G. Gu, Y. Feng, G. Zhong, L. Wei and C. Yang, *Nanotechnology*, 2018, **29**(45), 455501.
- 195 B. Subramanian, T. Mahalingam, C. Sanjeeviraja, M. Jayachandran and M. J. Chockalingam, *Thin Solid Films*, 1999, **357**, 119–124.
- 196 S. J. Choi and I. D. Kim, *Recent Developments in 2D Nanomaterials for Chemiresistive-Type Gas Sensors Electronic Mater. Lett.*, 2018, **14**(3), 221–260.
- 197 M. Popescu, A. Velea, F. Sava, A. Lőrinczi, A. Tomescu, C. Simion, E. Matei, G. Socol, I. N. Mihailescu, A. Andonie and I. Stamatina, *Phys. Status Solidi*, 2010, **207**, 516–520.
- 198 M. Popescu, F. Sava, A. Lorinczi, G. Socol, I. N. Mihailescu, A. Tomescu and C. Simion, *J. Non-Cryst. Solids*, 2007, **353**, 1865–1869.
- 199 S. Rani, M. Kumar, Y. Singh, M. Tomar, A. Sharma, V. Gupta and V. N. Singh, *J. Nanosci. Nanotechnol.*, 2021, **11**, DOI: 10.1166/jnn.2021.19278.
- 200 S. Rani, M. Kumar, Y. Singh and V. N. Singh, *J. Nanosci. Nanotechnol.*, 2021, **11**, DOI: 10.1166/jnn.2021.19277.
- 201 Z. Li, L. Sun, Y. Liu, D. Yu, Y. Wang, Y. Sun and M. Yu, *Environ. Sci. Nano*, 2019, **6**, 1507–1515.
- 202 M. Karamat, M. Fahad Ehsan, M. Naeem Ashiq, S. Ijaz, M. Najam-ul-Haq, S. Hamid and H. N. Bahnemann, *Appl. Surf. Sci.*, 2019, **463**, 1019–1027.
- 203 F. Zhang, C. Xia, J. Zhu, B. Ahmed, H. Liang, D. B. Velusamy, U. Schwingenschlögl and H. N. Alshareef, *Adv. Energy Mater.*, 2016, **6**, 1601188.
- 204 Y. Kim, Y. Kim, Y. Park, Y. N. Jo, Y. J. Kim, N. S. Choi and K. T. Lee, *Chem. Commun.*, 2015, **51**, 50–53.
- 205 R. Chen, S. Li, J. Liu, Y. Li, F. Ma, J. Liang, X. Chen, Z. Miao, J. Han, T. Wang and Q. Li, *Electrochim. Acta*, 2018, **282**, 973–980.
- 206 N. Shaji, P. Santhoshkumar, H. S. Kang, M. Nanthagopal, J. W. Park, S. Praveen, G. S. Sim, C. Senthil and C. W. Lee, *J. Alloys Compd.*, 2020, **834**, 154304.
- 207 K. Chen, X. Wang, G. Wang, B. Wang, X. Liu, J. Bai and H. Wang, *Chem. Eng. J.*, 2018, **347**, 552–562.
- 208 Q. Yu, B. Wang, J. Wang, S. Hu, J. Hu and Y. Li, *Front. Chem.*, 2020, **8**, 590.
- 209 H. Chen, B.-E. Jia, X. Lu, Y. Guo, R. Hu, R. Khatoun, L. Jiao, J. Leng, L. Zhang and J. Lu, *Chem.–Eur. J.*, 2019, **25**, 9973–9983.
- 210 T. Wang, K. Yang, J. Shi, S. Zhou, L. Mi, H. Li and W. Chen, *J. Energy Chem.*, 2020, **46**, 71–77.



- 211 D. Ni, Y. Chen, X. Yang, C. Liu and K. Cai, *J. Alloys Compd.*, 2018, **737**, 623–629.
- 212 K. M. Chung, D. Wamwangi, M. Woda, M. Wuttig and W. Bensch, *J. Appl. Phys.*, 2008, **103**(8), 083523.
- 213 R. Y. Wang, M. A. Caldwell, R. Gnana, D. Jeyasingh, S. Aloni, R. M. Shelby, R. Y. Wang, M. A. Caldwell, R. Gnana, D. Jeyasingh, S. Aloni, R. M. Shelby, H. P. Wong and D. J. Milliron, *J. Appl. Phys.*, 2011, **109**(11), 113506.
- 214 A. I. Journal, M. Sun, Y. Hu, B. Shen, J. Zhai and S. Song, *Integr. Ferroelectr.*, 2012, **140**(1), 1–7.
- 215 W. Wu, Z. He, S. Chen, J. Zhai, X. Liu, T. Lai, S. Song and Z. Song, *J. Appl. Phys.*, 2016, **120**(16), 165106.
- 216 J. E. Moore, *Nature*, 2010, **464**, 194–198.
- 217 J. Shen and J. J. Cha, *Nanoscale*, 2014, **6**, 14133–14140.
- 218 L. MÜchler, H. Zhang, S. Chadov, B. Yan, F. Casper, J. Kübler, S. C. Zhang and C. Felser, *Angew. Chem., Int. Ed.*, 2012, **51**, 7221–7225.
- 219 X. Chen, P. Lu, X. Wang, Y. Zhou, C. An, Y. Zhou, C. Xian, H. Gao, Z. Guo, C. Park, B. Hou, K. Peng, X. Zhou, J. Sun, Y. Xiong, Z. Yang, D. Xing and Y. Zhang, *Phys. Rev. B*, 2017, **96**, 165123.

