


 Cite this: *RSC Adv.*, 2021, **11**, 5521

Improving the optical and thermoelectric properties of $\text{Cs}_2\text{InAgCl}_6$ with heavy substitutional doping: a DFT insight†

 K. C. Bhamu, ^{*ab} Enamul Haque, ^c C. S. Praveen, ^{de} Nandha Kumar, ^f G. Yumnam, ^g Md. Anwar Hossain ^c and Gautam Sharma^h

The next-generation indium-based lead-free halide material $\text{Cs}_2\text{InAgCl}_6$ is promising for photovoltaic applications due to its good air stability and non-toxic behavior. However, its wide bandgap (>3 eV) is not suitable for the solar spectrum and hence reduces its photoelectronic efficiency for device applications. Here we report a significant bandgap reduction from 2.85 eV to 0.65 eV via substitutional doping and its effects on the optoelectronic and opto-thermoelectric properties from a first-principles study. The results predict that Sn/Pb and Ga and Cu co-doping will enhance the density of states significantly near the valence band maximum (VBM) and thus reduce the bandgap via shifting the VBM upward, while alkali metals (K/Rb) slightly increase the bandgap. A strong absorption peak near the Shockley–Queisser limit is observed in the co-doped case, while in the Sn/Pb-doped case, we notice a peak in the middle of the visible region of the solar spectrum. The nature of the bandgap is indirect with Cu–Ga/Pb/Sn doping, and a significant reduction in the bandgap, from 2.85 eV to 0.65 eV, is observed in the case of Ga–Cu co-doping. We observe a significant increase in the power factor (PF) ($2.03 \text{ mW m}^{-1} \text{ K}^{-2}$) for the n-type carrier after Pb-doping, which is ~ 3.5 times higher than in the pristine case ($0.6 \text{ mW m}^{-1} \text{ K}^{-2}$) at 500 K.

 Received 26th February 2020
 Accepted 6th January 2021

 DOI: 10.1039/d0ra01840f
rsc.li/rsc-advances

1. Introduction

The expeditious increase in the worldwide energy demand has accelerated the quest for alternative fuels based on renewable energy sources, which has emerged as a global challenge.¹ Solar light is considered to be a viable alternative to traditional fossil fuels, as it can be regarded as a clean and inexhaustible source of energy, which effectively benefits the environment and the future. Over the last few decades, large-scale efforts to develop

environment-friendly and cost-effective energy technologies that utilize the solar spectrum using solar cells and waste heat due *via* the Seebeck and Peltier effects have steadily increased.² Therefore, scientists have invested enormous efforts into discovering efficient and low-cost materials that effectively utilize the UV-visible region of the solar spectrum.

Halide perovskites have been studied for the last few decades as potential candidates for solar energy harvesting.³ After the first successful photovoltaic application by Kojima *et al.*,⁴ the organolead perovskites, $\text{CH}_3\text{NH}_3\text{PbBr}_3$, (OHP) have gained significant scientific attention and have emerged as a star material in solid-state solar cells since 2012, but exhibit stability issues.⁵ The available experimental^{6,7} and theoretical studies^{8–11} on halide perovskites mainly focus on their optoelectronic and thermoelectric applications.^{12–14} OHPs have been proposed to be economical and efficient materials like silicon and SnSe for optoelectronic and thermoelectric applications, respectively. Under practical operating conditions, *i.e.*, in the presence of moisture and air and at typical temperatures, organometal halide perovskites decompose into secondary phases.^{15,16} Their long-term stability and toxicity due to Pb still lag behind their high efficiency, and improving these drawbacks remains an ultimate challenge.¹⁷ There have been efforts to replace CH_3NH_3 with metal cations to improve their stability and to replace Pb with non-toxic cations.^{18–20} Recently, double perovskites have emerged as a green alternative to poisonous

^aPMC Division, CSIR-National Chemical Laboratory, Pune 411008, Maharashtra, India. E-mail: kcbhamu85@gmail.com

^bDepartment of Physics, Gramin Mahila P. G. College, Sikar 332024, Rajasthan, India

^cDepartment of Physics, Mawlana Bhashani Science and Technology University, Santosh, Tangail-1902, Bangladesh. E-mail: enamul.phy15@yahoo.com

^dInternational School of Photonics, Cochin University of Science and Technology, University Road, South Kalamassery, Kalamassery, Ernakulam, Kerala 682022, India

^eInter University Centre for Nano Materials and Devices, Cochin University of Science and Technology, University Road, South Kalamassery, Kalamassery, Ernakulam, Kerala 682022, India

^fThe Abdus Salam International Centre for Theoretical Physics, Strada Costiera 11, 34151 Trieste, Italy

^gDepartment of Physics and Astronomy, University of Missouri, Columbia, 65211, USA

^hDepartment of Physics, Indian Institute of Science Education and Research, Pune 411008, India

† Electronic supplementary information (ESI) available. See DOI: 10.1039/d0ra01840f

‡ Author contributed equally to the work.



lead halide perovskites. They also offer a viable engineering route for bandgap tuning *via* various cation and anion choices.^{16,21–26} A successful attempt to use inorganic halide perovskites was made in 2016, in which the first double halide perovskite, $\text{Cs}_2\text{AgBiX}_6$ (X = Br, Cl), was synthesized and tested for photovoltaic applications by Slavney *et al.*²³ and McClure *et al.*²⁵ Later, Xiao *et al.*²⁷ reported that the weak coupling between the 6s and 6p orbitals of Bi and Br, respectively, makes $\text{Cs}_2\text{AgBiX}_6$ less stable than OHPs. The bandgap of $\text{Cs}_2\text{AgBiX}_6$ is also above the Shockley–Queisser limit required for a single-junction solar cell.²⁸ Theoretical studies by Han *et al.*²⁹ based on density functional theory (DFT) contradict the experimental results regarding structural stability reported by Slavney *et al.*²³ and McClure *et al.*²⁵ and predict that its decomposition into binary and ternary compounds occurs above 700 K.

Volonakis *et al.*³⁰ reported $\text{Cs}_2\text{InAgCl}_6$ as a direct band gap (3.3 eV) material with air stability, non-toxicity, and Earth abundance. This double perovskite possesses a highly tuneable direct bandgap in the visible range. However, the measured direct bandgap of $\text{Cs}_2\text{InAgCl}_6$ is too far from the optimum value of the Shockley–Queisser limit required for high-performance solar cells. The large bandgap of $\text{Cs}_2\text{InAgCl}_6$ limits effective carrier transport between the valence band maximum and conduction band minimum, thus limiting the efficiency of the thermoelectric devices.³¹ However, it could be a suitable alternative to OHPs if the bandgap could be reduced by bandgap engineering. Double perovskites have the peculiar feature of wide compositional tunability. They offer the ability to alter their atomic composition by doping or alloying, which are straightforward approaches to engineer their bandgaps. Here, we attempt to tune the bandgap and hence the optoelectronic and opto-thermoelectric properties of $\text{Cs}_2\text{InAgCl}_6$ by substitutional doping.

2. Crystal structure and computational methods

The lead-free halide double perovskite $\text{Cs}_2\text{InAgCl}_6$ crystallizes in the face-centered cubic structure with the space group $Fm\bar{3}m$ (225)³⁰ as shown in the left panel of Fig. SI-1 (ESI†). The Cs/In/Ag/Cl occupies the 8c/4a/4b/24e Wyckoff positions. The optimized lattice parameters of $\text{Cs}_2\text{InAgCl}_6$ are listed in Table SI-1,† along with other available data for comparison, and they are in good agreement with previous reports.³⁰ To model the 25% substitutional doping at different atomic sites, two Cs atoms located at the $(\frac{3}{4}, \frac{3}{4}, \frac{3}{4})$ and $(\frac{1}{4}, \frac{1}{4}, \frac{1}{4})$ positions are substituted by Rb or K. In the Pb and Sn-doped cases, In at (0, 0, 0) and Ag at $(\frac{1}{2}, \frac{1}{2}, \frac{1}{2})$ are substituted by Pb and Sn. For co-doping, we replaced In at (0, 0, 0) with Ga, and Ag at $(\frac{1}{2}, \frac{1}{2}, \frac{1}{2})$ with Cu, resulting in 25% effective doping of the parent system. The doped systems adopt different space groups, as mentioned in Table SI-1,† Further crystallographic details can be obtained in ref. 30.

We performed all the first principle calculations using the full potential augmented plane wave method (FP-LAPW) as implemented in the WIEN2k code.^{32,33} We used the generalized gradient approximation (GGA) with the Perdew–Burke–

Ernzerhof (PBE) functional for the description of electron exchange and correlation.^{34,35} To calculate the electronic properties with an improved band gap, we employed the Becke and Johnson potential³⁶ and its modified versions.^{37,38} Table SI-1,† demonstrates that the modified Beck and Johnson potential (mBJ) proposed by Koller *et al.*³⁹ (hereafter referred to as KTB-mBJ) gives a better estimate of bandgap than other mBJs. Recently, Jishi *et al.*⁴⁰ reported that, with the appropriate tuning of the parameter c (eqn (1) in ref. 36), the experimental band gap of perovskites can be reproduced. However, in the present investigation, we found that the bandgap was overestimated compared to the experimental bandgap, and hence we do not report the results calculated using Jishi's mBJ potential. Therefore, all the relevant properties of interest were computed using KTB-mBJ in this study. We set the size of the k -mesh to be $12 \times 12 \times 12$, the product of the smallest of all the atomic sphere radii (R_{mt}) and the plane wave cut-off (K_{max}) to 8.5, the magnitude of the largest vector (G_{max}) to 12, l_{vnsmax} to 6, and the energy, charge, and force convergence criteria to 10^{-5} Ry, 0.0001e and 0.05 mRy per a.u., respectively. The self-consistent field calculations were performed using $10 \times 10 \times 10$ and $5 \times 5 \times 5$ k -meshes for the pristine and doped cases while all other parameters were kept at the values described above and mentioned in Table SI-1,† For the computation of the electronic and optical properties, $14 \times 14 \times 14$ and $7 \times 7 \times 7$ shifted k -mesh grids were used for pristine and doped systems. We calculated the optical properties using the method described in ref. 41. Thermoelectric (TE) properties, such as the Seebeck coefficient (α), electrical conductivity (σ), electronic thermal conductivity (k_e), power factor (PF) and thermoelectric figure of merit (ZT), were computed using Boltzmann transport theory within the constant relaxation time approximation (CRTA) and rigid band approximation (RBA).^{31,42,43} Convergence was achieved in the TE coefficients (change of less than 1%) with dense k -grids of $28 \times 28 \times 28$ and $14 \times 14 \times 14$ for the pristine and the doped cases, respectively.

3. Results and discussion

3.1. Electronic structure

Fig. 1a–c depicts the band structure of pristine $\text{Cs}_2\text{InAgCl}_6$ computed using PBE (Fig. 1a), PBE + SOC (Fig. 1b) and KTB-mBJ + SOC (Fig. 1c). The overall features of the band structure of $\text{Cs}_2\text{InAgCl}_6$ are similar to those in the previous reports such as in ref. 44. It is to be noted that the band gap (2.85 eV) computed using TB-mBJ agrees well with the values (2.53 eV; TB-mBJ) reported by some of the co-authors⁴⁴ but still underestimates the experimental band gap (3.3 eV) owing to the difference in the mBJ potentials used. Additionally, it is clear from Fig. 1c that KTB-mBJ + SOC ($E_g = 2.85$ eV) improved the PBE band gap (1.0 eV; Fig. 1a), and we observe that the inclusion of SOC (Fig. 1b) slightly changes the band dispersion in the valence band.

Further information about the density of states (DOS) and projected DOS for the studied systems is given in Fig. SI-2,†

To engineer the bandgap, we employed dual substitutional doping at different atomic sites, *i.e.*, isoelectronic (K/Rb at the



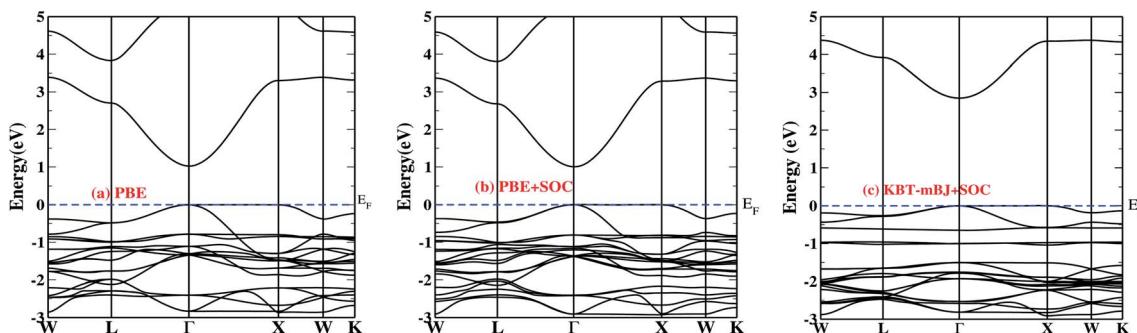


Fig. 1 Electronic dispersion relations of $\text{Cs}_2\text{InAgCl}_6$ using: (a) PBE, (b) PBE + SOC, and (c) PBE with KTB-mBJ + SOC. The blue dashed lines represent the Fermi level (E_F).

Cs site) and non-isoelectronic (Sn, Pb; Cu and Ga co-doping (CGCD)) doping.

For the isoelectronic doping we observed that, for the fully relaxed system, the lattice parameters changed less than 0.5% compared to the pristine case. Since K/Rb is doped in the cage of the octahedra (Cs-site) of AgCl_6 and InCl_6 , their valence electron (ns^1) is not expected to contribute near the VB/CB edges because it is located at the deep valence states near ~ 6 eV. This was confirmed from the electronic structure calculations shown in Fig. SI-(3 and 4).[†] K(Rb) doping increases the bandgap by 0.18 (0.16 eV) compared to the pristine case due to the size effect of the dopant,⁴⁵ and the band dispersion remains intact.

It has been reported⁴⁶ that the band dispersion near the Fermi level can be altered significantly by alloying the octahedra cation site (Ag/In) to enhance the optoelectronic and optothermoelectric properties of the material. Therefore, we chose the Ag/In site for non-isoelectronic doping.

Fig. 2a and b depict the Pb- and Sn-doped band structures, respectively. Our findings reveal that Sn is a better dopant candidate compared to Pb for optoelectronic applications. A detailed description can be found in the next section.

The doped system becomes an indirect semiconductor with a bandgap of 2.40 (2.08) eV after Pb(Sn) doping. The different reduction in the bandgap with Pb doping compared to Sn doping can be understood from the location of In-s states (Fig. 3), which is shifted by ~ 0.3 eV when the system is doped with Pb rather than Sn. The VBM occurs at the Γ point, whereas

in the conduction band there are two minima present at M and R. The minima in the conduction band are separated by ~ 0.07 eV, which gives rise to the valley degeneracy responsible for the high figure of merit [ZT] for n-type carriers.⁴⁷

Interestingly, we observed the presence of Sn/Pb-s states in the proximity of the Fermi level, giving rise to increased charge carriers, which are especially advantageous in the context of thermoelectrics.

Fig. 2c illustrates the band structure of the CGCD system. The top of the valence band is mainly composed of the Cu-d state with slight contributions from the In-s/p states. In the conduction band, the In-s state is shifted towards the Fermi level by ~ 2.2 eV, reducing the bandgap value to 0.65 eV which is the lowest bandgap found among the various systems in the present study. The bandgap is indirect, which prevents fast electron-hole recombination, giving a high exciton lifetime.

The calculated total density of states of the studied alloys are illustrated in Fig. SI-(4 and 5; for K/Rb and CGCD[†]). In the next sections, we focus on the effects of the reduction of the bandgap on the optical and thermoelectric properties.

3.2. Optical properties

After calculating the Kohn-Sham (KS) electronic states using the k -grids mentioned in the crystal structure and computation methods section, we computed the optical absorption spectra using KTB-mBJ + SOC for pristine $\text{Cs}_2\text{InAgCl}_6$ along with its

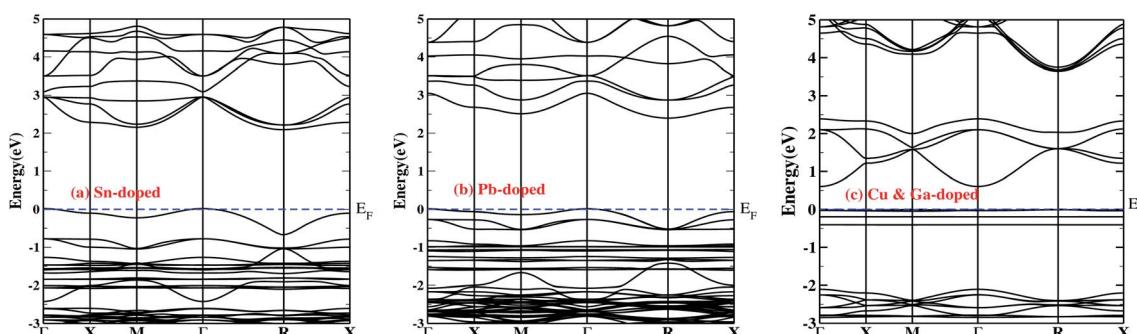


Fig. 2 Electronic dispersion relations of: (a) 25% Sn-doped, (b) 25% Pb-doped, and (c) 12.5% Cu- and 12.5% Ga-doped $\text{Cs}_2\text{InAgCl}_6$ obtained using the PBE functional with the KTB-mBJ potential including the SO effect. The blue dashed lines represent the Fermi level.



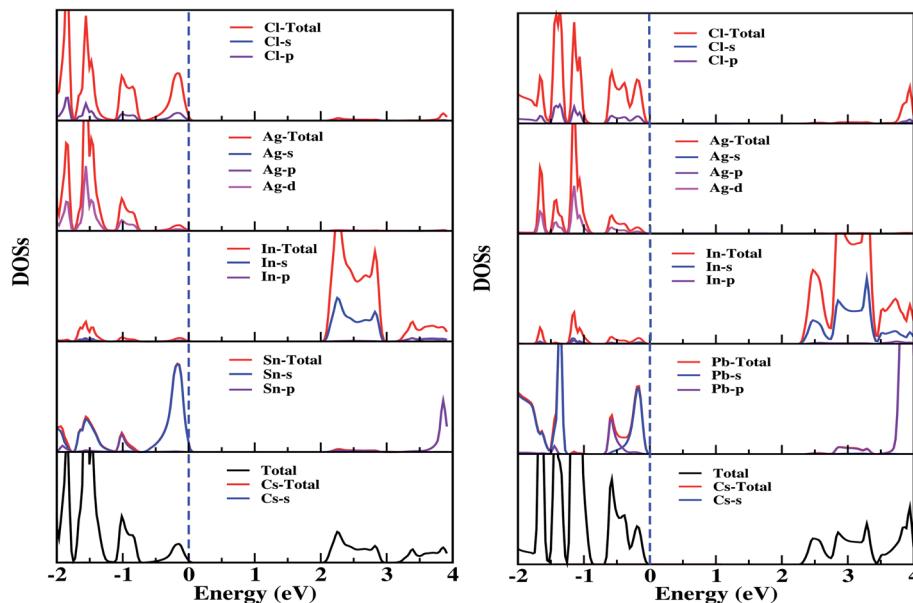


Fig. 3 Calculated total, atom, and orbital decomposed densities of states of Sn- (left panel) and Pb- (right panel) doped $\text{Cs}_2\text{InAgCl}_6$ obtained using the PBE functional with the KTB-mBJ potential.

doped derivatives, which are shown in Fig. 4. A full description of the method and equations used in computing the optical properties is provided in ref. 41. The changes in the bandgap due to the substitutional dopants are clearly reflected in the optical spectra. For photovoltaic applications, a material must show high optical absorbance in the visible region of the solar spectrum.^{48,49} The optical absorption edge was calculated by extrapolation to be ~ 2.80 eV for the pristine case, and gradually shifted towards lower energy with doping. For the CGCD system, the shift was drastic (See Fig. SI-6†).

Our results show that doping with a suitable candidate can significantly improve the optical absorption over the ultraviolet (UV) to IR (infrared) region. The optical absorption strength of the co-doped system in the IR solar spectrum region is stronger than that of the other doped systems, while in the visible region

the Sn-doped system is dominant. The present findings may serve as a guide for experimental studies on these non-isoelectronic doped systems.

3.3. Thermoelectric properties

The thermoelectric (TE) transport properties were calculated using the linearized Boltzmann transport equation (BTE), as implemented in the BoltzTraP code,⁴² which solves the BTE within the constant relaxation time approximation (CRTA). The BoltzTraP code interpolates the band structure obtained from the DFT calculations, carries out the necessary integrations (Fermi integrals) at different temperatures and Fermi levels and provides the required transport coefficients as a function of temperature and carrier concentration (for both p-type and n-type); the results are shown in Fig. 5–7.

The performance of a thermoelectric material is governed by the dimensionless figure of merit (ZT), which is given by^{50,51}

$$ZT = \frac{\alpha^2 \sigma}{\kappa} T$$

where σ is electrical conductivity, $\alpha^2 \sigma$ is the power factor, T is absolute temperature, and κ is the total thermal conductivity, which is composed of an electronic (κ_e) and a lattice part (κ_l). We do not consider the lattice part of the thermal conductivity here.

Tuning the bandgap by substitutional doping changed the dispersion of the electronic bands around the Fermi level. This in turn governed the behavior of $|\alpha|$ and σ , and hence, the power factor was affected.

Fig. 5 depicts the Seebeck coefficient ($|\alpha|$) for p-type and n-type pristine and doped systems. The isoelectronic doping (with K and Rb) (Fig. 5c and e) does not alter the $|\alpha|$ value, as there is no change in the dispersion of bands. We found that $|\alpha|$ shows typical behavior with decreasing carrier

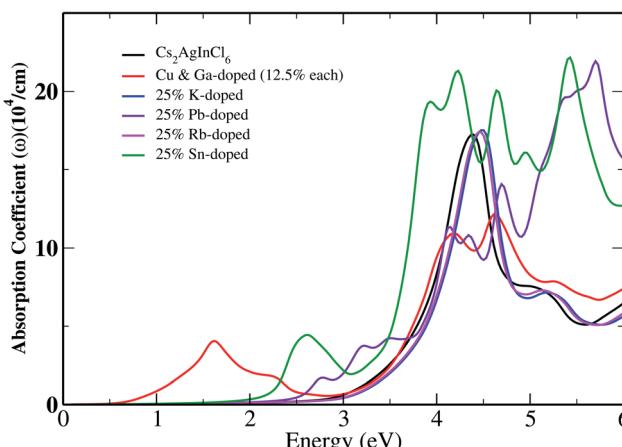


Fig. 4 The absorption coefficient of $\text{Cs}_2\text{InAgCl}_6$ along with those of the doped cases as a function of the photon energy.



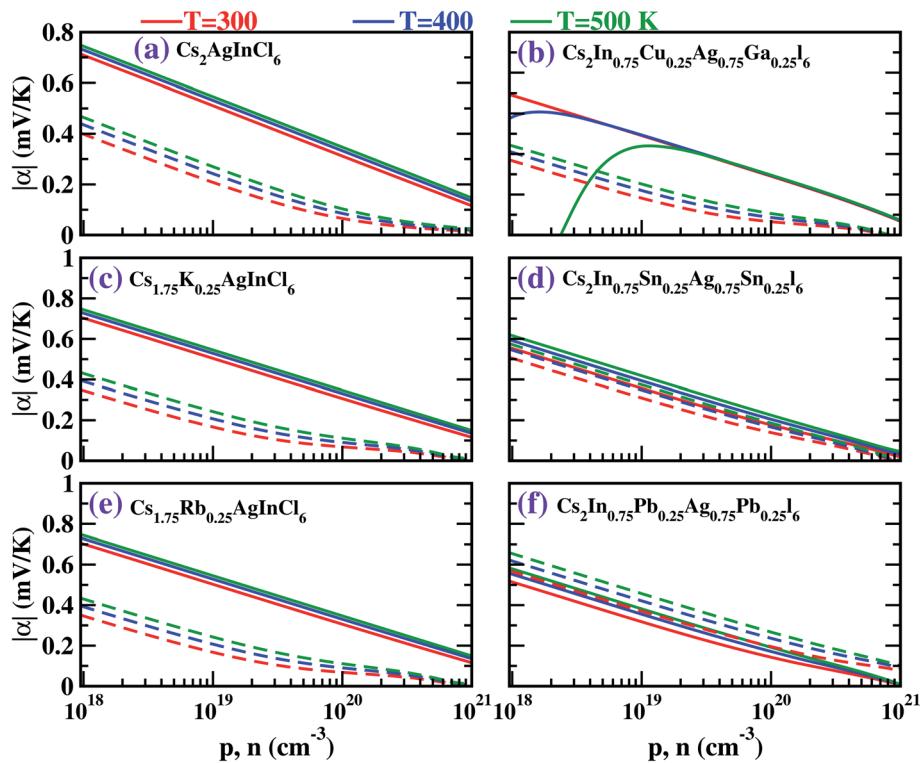


Fig. 5 The carrier-concentration-dependent Seebeck coefficient ($|\alpha|$) values at different temperatures for both p-type (solid) and n-type (dashed) (a) pristine, (b) Cu and Ga co-doped, (c) K-doped, (d) Sn-doped, (e) Rb-doped, and (f) Pb-doped $\text{Cs}_2\text{AgInCl}_6$.

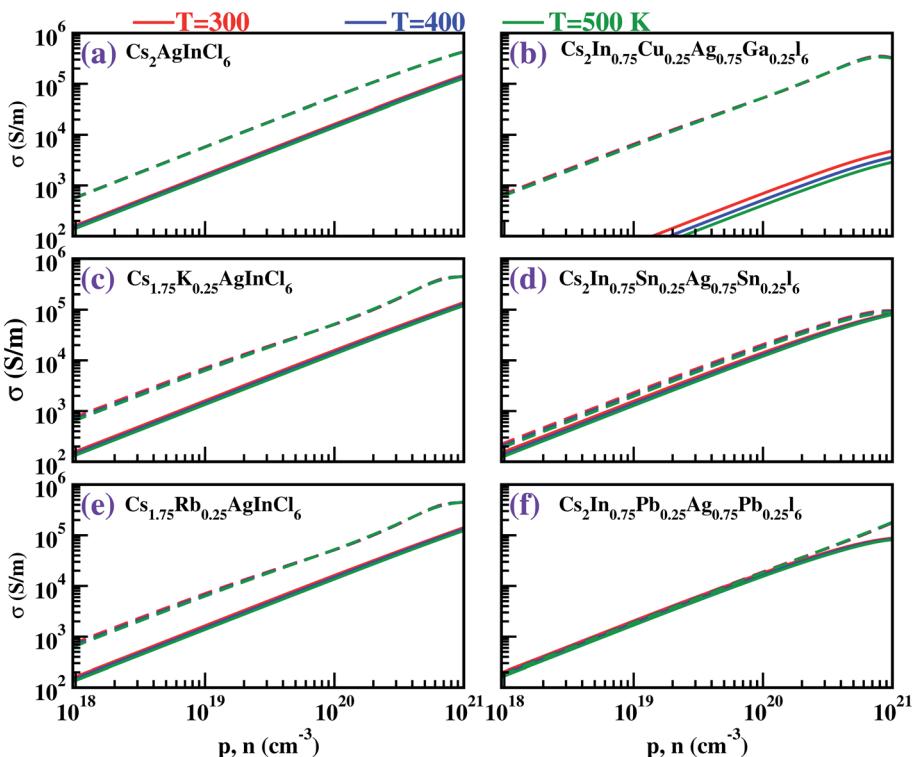


Fig. 6 The electrical conductivity (σ) values at different temperatures for both p-type (solid) and n-type (dashed) (a) pristine, (b) Cu and Ga co-doped, (c) K-doped, (d) Sn-doped, (e) Rb-doped, and (f) Pb-doped $\text{Cs}_2\text{AgInCl}_6$.

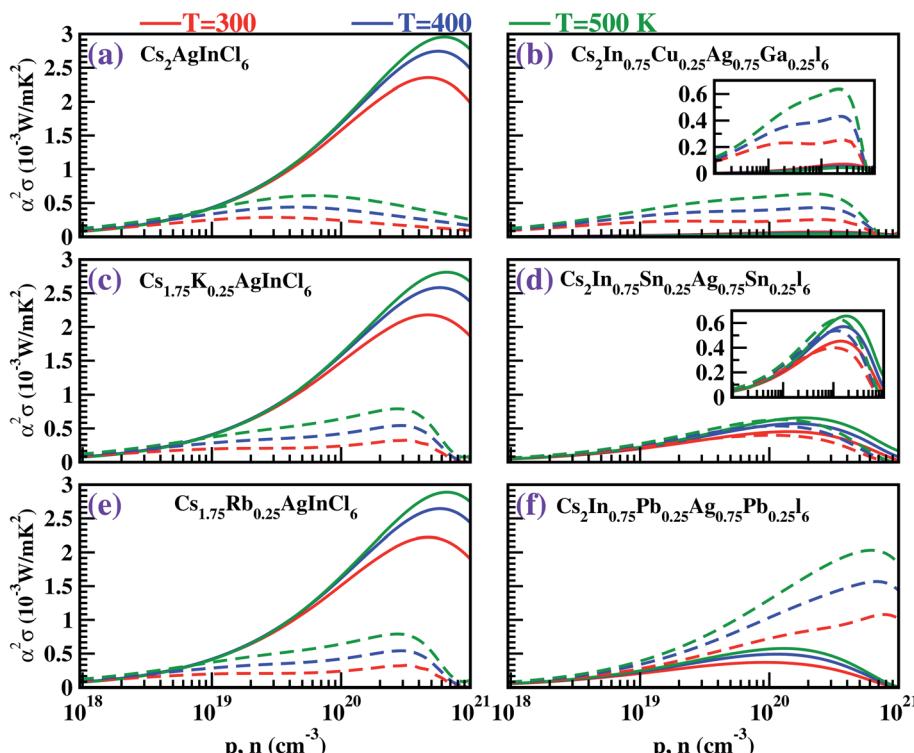


Fig. 7 The power factor ($\alpha^2 \sigma$) values at different temperatures for both p-type (solid) and n-type (dashed) (a) pristine, (b) Cu and Ga co-doped, (c) K-doped, (d) Sn-doped, (e) Rb-doped, and (f) Pb-doped $\text{Cs}_2\text{AgInCl}_6$.

concentration, and also computed that the $|\alpha|$ value for the hole is higher than that for the electrons. Since $|\alpha|$ is directly proportional to the effective mass of the charge carriers, the flat band dispersion (Γ -X) at the valence band edges gives rise to a large hole effective mass (m_h^*). On the other hand, a parabolic band dispersion is observed for the conduction bands, leading to a lighter electron effective mass (m_e^*). For the Pb/Sn-doped (Fig. 5d/f) system, the behavior of $|\alpha|$ can be interpreted based on the band gaps and band dispersion. For all temperatures, the $|\alpha|$ value of the electrons with the Pb-doped system is ~ 3 times higher than that observed in the pristine case at $n = 10^{20} \text{ cm}^{-3}$, which was due to the downward shift of the In-s state. It was found that the Pb-doped system is more favorable than the pristine one. Moreover, we found that the $|\alpha|$ for the electrons exceeds that for the holes at higher carrier concentrations, the effect of which is reflected in the power factor (PF). In the case of CGCD, the narrow band gap triggers bipolar conduction at lower carrier concentrations, and due to the flat band, the electrical conductivity is too low compared to those of the other cases.

For all the cases, we obtained the electronic conductivity (σ) and PF by multiplying by $\tau = 10^{-14} \text{ s}$; this approach has also been adopted by others.^{42,52} The results are displayed in Fig. 6 and 7. We found that σ increases with increasing carrier concentration. In all the cases, σ was found to be greater for electrons than for holes due to the parabolic dispersion of the bands. However, in the Pb-doped system, the σ values for the electrons and holes were found to be of similar order, which is advantageous to boost performance with n-type doping.

The temperature-dependent power factor, which is the product of α^2 and σ , is plotted as a function of carrier concentration in Fig. 7. Because of the isoelectronic doping (Fig. 7c and e), the behavior of the power factor for the K/Rb-doped cases is nearly the same as that for the pristine case with p-type doping, and slightly higher than that for the pristine case for n-type doping. For the co-doped case (Fig. 7b), the optimized value of the power factor with n-type doping is $0.6 \text{ mW m}^{-1} \text{ K}^{-2}$. A remarkable feature in our study is the significantly enhanced PF ($2.03 \text{ mW m}^{-1} \text{ K}^{-2}$) for n-type carrier concentration for the Pb-doped case, in which the PF is ~ 3.5 times higher than in the pristine case ($0.6 \text{ mW m}^{-1} \text{ K}^{-2}$) at 500 K.

4. Conclusions

In summary, we studied the effects of K, Rb, and Cu-Ga/Pb/Sn substitution on the electronic, optical, and thermoelectric transport properties of $\text{Cs}_2\text{InAgI}_6$ using first-principles calculations. We found that K/Rb-substitution increases the bandgap slightly, while Cu-Ga/Pb/Sn-substitution decreases the band gap significantly. The studied K/Rb-based systems are direct bandgap semiconductors, while an indirect bandgap nature is observed for the other doped systems. The absorption peak is shifted from the ultraviolet region to the visible spectrum for the indirect bandgap compounds, but remains similar for the direct bandgap compounds. The optical absorption strength of the co-doped (Cu-Ga-based) system is stronger than that of the other doped systems in the IR solar region, while in the visible



region, the Sn-doped system shows dominant absorption. Significant changes in thermoelectric properties were noted in the Pb- and co-doped cases. The flat valence bands are responsible for the poor thermoelectric performance of the co-doped system. The downward shift of the In-s state in the Pb-doped system enhances the $\text{PF} \sim 3.5$ -fold compared to the pristine case ($0.6 \text{ mW m}^{-1} \text{ K}^{-2}$) at 500 K with n-type doping. Thus, our investigation reveals that the Pb-doped system is a promising candidate for thermoelectric applications.

Conflicts of interest

There are no conflicts to declare.

Acknowledgements

KCB and CSP acknowledge DST and DST-SERB for SERB-NPFD and INSPIRE Faculty fellowships with award numbers IFA-18 PH217 and PDF/2017/002876, respectively. The research used resources at the National Energy Research Scientific Computing Center (NERSC), which is supported by the Office of Science of the USDOE, United States, under contract no. DE-AC02-05CH11231. The high-performance computational facility of the Centre for Development of Advanced Computing, Pune, was also used for the present investigation.

References

- 1 C. R. Lhermitte and K. Sivula, *ACS Catal.*, 2019, **9**, 2007–2017.
- 2 P. Bermel, K. Yazawa, J. L. Gray, X. Xu and A. Shakouri, *Energy Environ. Sci.*, 2016, **9**, 2776–2788.
- 3 S. Hirotsu, J. Harada, M. Iizumi and K. Gesi, *J. Phys. Soc. Jpn.*, 1974, **37**, 1393–1398.
- 4 A. Kojima, K. Teshima, Y. Shirai and T. Miyasaka, *J. Am. Chem. Soc.*, 2009, **131**, 6050–6051.
- 5 M. M. Lee, J. Teuscher, T. Miyasaka, T. N. Murakami and H. J. Snaith, *Science*, 2012, **338**, 643–647.
- 6 M. Neukom, S. Züfle, S. Jenatsch and B. Ruhstaller, *Sci. Technol. Adv. Mater.*, 2018, **19**, 291–316.
- 7 G. Grancini, C. Roldán-Carmona, I. Zimmermann, E. Mosconi, X. Lee, D. Martineau, S. Narbey, F. Oswald, F. De Angelis, M. Graetzel and others, *Nat. Commun.*, 2017, **8**, 15684.
- 8 F. F. Targhi, Y. S. Jalili and F. Kanjouri, *Results Phys.*, 2018, **10**, 616–627.
- 9 N. Hernández-Haro, J. Ortega-Castro, Y. B. Martynov, R. G. Nazmitdinov and A. Frontera, *Chem. Phys.*, 2019, **516**, 225–231.
- 10 G. Giorgi, J.-I. Fujisawa, H. Segawa and K. Yamashita, *J. Phys. Chem. C*, 2014, **118**, 12176–12183.
- 11 P. Umari, E. Mosconi and F. De Angelis, *Sci. Rep.*, 2014, **4**, 4467.
- 12 S.-Y. Yue, X. Zhang, G. Qin, J. Yang and M. Hu, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2016, **94**, 115427.
- 13 X. Qian, X. Gu and R. Yang, *Appl. Phys. Lett.*, 2016, **108**, 63902.
- 14 T. Zhao, D. Wang and Z. Shuai, *Synth. Met.*, 2017, **225**, 108–114.
- 15 Y.-Y. Zhang, S. Chen, P. Xu, H. Xiang, X.-G. Gong, A. Walsh and S.-H. Wei, *Chin. Phys. Lett.*, 2018, **35**, 36104.
- 16 Z. Li, Q. Xu, Q. Sun, Z. Hou and W.-J. Yin, *Adv. Funct. Mater.*, 2019, **29**, 1807280.
- 17 Q. Fu, X. Tang, B. Huang, T. Hu, L. Tan, L. Chen and Y. Chen, *Adv. Sci.*, 2018, **5**, 1700387.
- 18 F. Wei, Z. Deng, S. Sun, F. Xie, G. Kieslich, D. M. Evans, M. A. Carpenter, P. D. Bristowe and A. K. Cheetham, *Mater. Horiz.*, 2016, **3**, 328–332.
- 19 F. Wei, Z. Deng, S. Sun, F. Zhang, D. M. Evans, G. Kieslich, S. Tominaka, M. A. Carpenter, J. Zhang, P. D. Bristowe and others, *Chem. Mater.*, 2017, **29**, 1089–1094.
- 20 Z. Deng, F. Wei, S. Sun, G. Kieslich, A. K. Cheetham and P. D. Bristowe, *J. Mater. Chem. A*, 2016, **4**, 12025–12029.
- 21 F. Giustino and H. J. Snaith, *ACS Energy Lett.*, 2016, **1**, 1233–1240.
- 22 X.-G. Zhao, J.-H. Yang, Y. Fu, D. Yang, Q. Xu, L. Yu, S.-H. Wei and L. Zhang, *J. Am. Chem. Soc.*, 2017, **139**, 2630–2638.
- 23 A. H. Slavney, T. Hu, A. M. Lindenberg and H. I. Karunadasa, *J. Am. Chem. Soc.*, 2016, **138**, 2138–2141.
- 24 A. Jain, O. Voznyy and E. H. Sargent, *J. Phys. Chem. C*, 2017, **121**, 7183–7187.
- 25 E. T. McClure, M. R. Ball, W. Windl and P. M. Woodward, *Chem. Mater.*, 2016, **28**, 1348–1354.
- 26 L. Chu, W. Ahmad, W. Liu, J. Yang, R. Zhang, Y. Sun, J. Yang and others, *Nano-Micro Lett.*, 2019, **11**, 16.
- 27 Z. Xiao, W. Meng, J. Wang and Y. Yan, *ChemSusChem*, 2016, **9**, 2628–2633.
- 28 C. N. Savory, A. Walsh and D. O. Scanlon, *ACS Energy Lett.*, 2016, **1**, 949–955.
- 29 D. Han, T. Zhang, M. Huang, D. Sun, M.-H. Du and S. Chen, *APL Mater.*, 2018, **6**, 84902.
- 30 G. Volonakis, A. A. Haghhighirad, R. L. Miltot, W. H. Sio, M. R. Filip, B. Wenger, M. B. Johnston, L. M. Herz, H. J. Snaith and F. Giustino, *J. Phys. Chem. Lett.*, 2017, **8**, 772–778.
- 31 J. Sun and D. J. Singh, *J. Mater. Chem. A*, 2017, **5**, 8499–8509.
- 32 P. Blaha, K. Schwarz, F. Tran, R. Laskowski, G. K. H. Madsen and L. D. Marks, *J. Chem. Phys.*, 2020, **152**, 74101.
- 33 P. Blaha, K. Schwarz, G. K. H. Madsen, D. Kvasnicka, J. Luitz, R. Laskowski, F. Tran and L. D. Marks, *WIEN2k, An Augmented Plane Wave + Local Orbitals Program for Calculating Crystal Properties*, Karlheinz Schwarz, Techn. Universität Wien, Austria, 2018, ISBN 3-9501031-1-2.
- 34 J. P. Perdew, K. Burke and M. Ernzerhof, *Phys. Rev. Lett.*, 1996, **77**, 3865.
- 35 J. P. Perdew, A. Ruzsinszky, G. I. Csonka, O. A. Vydrov, G. E. Scuseria, L. A. Constantin, X. Zhou and K. Burke, *Phys. Rev. Lett.*, 2008, **100**, 136406.
- 36 F. Tran and P. Blaha, *Phys. Rev. Lett.*, 2009, **102**, 226401.
- 37 D. Koller, F. Tran and P. Blaha, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2012, **85**, 1–10.
- 38 A. D. Becke and E. R. Johnson, *J. Chem. Phys.*, 2006, **124**, 221101.



39 D. Koller, F. Tran and P. Blaha, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2011, **83**, 195134.

40 R. A. Jishi, O. B. Ta and A. A. Sharif, *J. Phys. Chem. C*, 2014, **118**, 28344–28349.

41 C. Ambrosch-Draxl and J. O. Sofo, *Comput. Phys. Commun.*, 2006, **175**, 1–14.

42 G. K. H. Madsen and D. J. Singh, *Comput. Phys. Commun.*, 2006, **175**, 67–71.

43 J. Sun and D. J. Singh, *Phys. Rev. Appl.*, 2016, **5**, 24006.

44 E. Haque and M. A. Hossain, *Comput. Condens. Matter*, 2019, **19**, e00374.

45 K. C. Bhamu, A. Soni and J. Sahariya, *Sol. Energy*, 2018, **162**, 336–343.

46 P. Zhang, J. Yang and S.-H. Wei, *J. Mater. Chem. A*, 2018, **6**, 1809–1815.

47 M. Mukherjee, G. Yumnam and A. K. Singh, *J. Phys. Chem. C*, 2018, **122**, 29150–29157.

48 C. Brabec, U. Scherf and V. Dyakonov, *Organic photovoltaics: materials, device physics, and manufacturing technologies*, John Wiley & Sons, 2011.

49 E. Rephaeli and S. Fan, *Opt. Express*, 2009, **17**, 15145–15159.

50 A. Ioffe, *Semiconductor Thermoelements and Thermoelectric Cooling*, Infosearch, London, 1st edn, 1957.

51 H. J. Goldsmid, *Introduction to thermoelectricity*, Springer, 2010, vol. 121.

52 S.-D. Guo, *J. Mater. Chem. C*, 2016, **4**, 9366–9374.

