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## Substitutional–interstitial structural transition in Cu–Pt nano-alloys

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Copper–platinum alloys are important binary alloys in catalysis. In this communication, we demonstrate that it is possible to preserve the thermal properties of platinum with a copper–platinum alloy by converting the substitutional alloy into an interstitial one. This conversion occurs when the size of the copper–platinum system is reduced down to the nanoscale. The size-dependent phase diagram of Cu–Pt for a spherical nanoparticle is calculated at various sizes (50, 10 and 5 nm) demonstrating that Cu–Pt alloyed nanoparticles can be formed all over the composition range. Experimentally, the electron microscopy characterization of copper–platinum alloyed nanoparticles synthesized by wet chemistry supports the predicted structural transition.

Platinum is an outstanding catalyst used to produce nitric acid, silicone, and benzene and it is also used to improve the efficiency of fuel cells.<sup>1</sup> Platinum is classified as an energy critical element by the American Physical Society and the Materials Research Society<sup>2,3</sup> due to its low abundance in the Earth's crust which makes its price high and accessibility at risk. Therefore, any strategy to reduce the amount of platinum in applications is highly welcomed. The strategy followed in this paper is first alloying platinum with an abundant metal such as copper because the copper–platinum system displays total miscibility all over its composition range; second reducing the size of the system. In order to achieve this goal experimentally, theoretical guidance is necessary to determine how platinum and copper are going to mix at the nanoscale. Compared to the wealth of research performed on other binary alloys at the nanoscale, Cu–Pt has not been fully explored.<sup>4–8</sup>

To answer this fundamental question, the knowledge of the binary phase diagram at the nanoscale is required. Indeed, the

binary phase diagram is the fingerprint of any binary alloy, displaying a temperature–composition map which indicates the equilibrium phases present at a given temperature and composition.<sup>9–11</sup> Several models (ideal, regular and sub-regular) have been developed to predict theoretically any phase diagram. All those models rely on minimizing the Gibbs free energy of the system. The ideal solution model assumes that the solid solution is completely miscible *i.e.*, not having a miscibility gap, resulting in the classical lens-shape phase diagram. This model is too simple to simulate a binary alloy with ordered phases such as Cu–Pt. By considering explicitly the chemical interaction between Cu and Pt within the solid solution, a composition-dependent mixing term is added to the Gibbs free energy, resulting in the so-called regular solution model. However, this addition is not sufficient to accurately describe the Cu–Pt phase diagram.<sup>12</sup> In order to describe the Cu–Pt system, the mixing enthalpy needs to be temperature–composition-dependent which falls under the sub-regular solution model.<sup>12</sup> In the sub-regular model, the Gibbs free energy of a non-ideal binary system such as Cu–Pt is represented by:<sup>13,14</sup>

$$G = \sum_{i=1}^2 x_i G_i^0 + RT \sum_{i=1}^2 x_i \ln x_i + G_m \quad (1)$$

where  $x_i$  is the molar concentration of the constituent  $i$  ( $i = 1$  represents Cu while  $i = 2$  represents Pt) in solution,  $R$  is the molar gas constant,  $T$  is the temperature, and  $G_i^0$  is the Gibbs free energy of the pure element  $i$ .  $G_m$  is the mixing Gibbs energy describing the interactions between the chemical elements in solution, which is represented by the Redlich–Kister polynomial:<sup>13,15</sup>

$$G_m = x_1 x_2 \sum_{k=0}^n {}^k L_{12} (x_1 - x_2)^k \quad (2)$$

where the  $k^{\text{th}}$  order Redlich–Kister parameters,  ${}^k L_{12}$ , are represented by  ${}^k A + {}^k B T$ , where  ${}^k A$  and  ${}^k B$  are constants.<sup>16</sup> All the parameters used in this sub-regular model are listed in Table 1.

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Table 1 Bulk material properties of the Cu–Pt system

Material property	Cu	Pt	Ref.
$T_{m,\infty}$ (K)	1357	2041	33
$\Delta H_{m,\infty}$ (kJ mol <sup>-1</sup> )	13.3	22.2	33
$\alpha_{\text{shape}}$ (nm)	1.69	1.50	34
${}^0L_{12}^{\text{liq}}$ (J mol <sup>-1</sup> )	-38 932		8
${}^0L_{12}^{\text{sol}}$ (J mol <sup>-1</sup> )	-35 625		8
${}^0L_{12}^{\text{liq}}$ (J mol <sup>-1</sup> )	-24 900 - 3.3T		12
${}^1L_{12}^{\text{liq}}$ (J mol <sup>-1</sup> )	-34 400 + 9.1T		12
${}^2L_{12}^{\text{liq}}$ (J mol <sup>-1</sup> )	-14 500 - 12.1T		12
${}^0L_{12}^{\text{sol}}$ (J mol <sup>-1</sup> )	-44 105 + 6.53T		12
${}^1L_{12}^{\text{sol}}$ (J mol <sup>-1</sup> )	-9100 - 1.24T		12
${}^2L_{12}^{\text{sol}}$ (J mol <sup>-1</sup> )	7255 - 8.65T		12

Experimentally, the bulk binary phase diagram of Cu–Pt does not display any miscibility gap; consequently, both 0<sup>th</sup> order Redlich–Kister parameters ( ${}^0L_{12}^{\text{liq}}$  and  ${}^0L_{12}^{\text{sol}}$ ) need to be negative for both phases (Table 1). Indeed, for miscibility gaps to arise in a phase diagram, both 0<sup>th</sup> order Redlich–Kister parameters ( ${}^0L_{12}^{\text{liq}}$  and  ${}^0L_{12}^{\text{sol}}$ ) need to be positive for both phases, liquid and solid.<sup>9</sup> This is apparent in other binary systems such as Si–Ge,<sup>17</sup> Cu–Ni<sup>18</sup> and Bi–Sb.<sup>19</sup> Cu–Pt is not known to have any regions of immiscibility, *i.e.* displaying miscibility gaps as shown in ref. 8. So, it is curious how the authors of ref. 8 determined that this system would have a miscibility gap since they provided a negative Gibbs free energy of mixing for this alloy. In addition to the analysis concerning the Redlich–Kister parameters, it is also noted that binary systems displaying miscibility gaps typically break at least one of the Hume–Rothery rules.<sup>20</sup> But, Cu–Pt does not break any Hume–Rothery rules as mentioned in ref. 8. Therefore, it makes perfect sense that Cu–Pt does not display any miscibility gap in its binary phase diagram. Moreover, the Cu–Pt system has several ordered structures appearing at low temperature and low platinum concentration namely the L1<sub>1</sub> and L1<sub>2</sub> ordered structures.<sup>12</sup> As a note, the Hume–Rothery rules are used to determine whether a binary system would form a miscible solid solution (negative mixing enthalpy). But, those rules do not inform which model (regular or sub-regular) is the best to calculate the binary phase diagram.

According to Fig. 1, the bulk Cu–Pt system is best modeled by using the sub-regular solution model. Indeed, the sub-regular solution model is in good agreement with all the available experimental data, while the regular solution model from ref. 8 does not fit all the available experimental data very well, except below 20% platinum weight composition. Above 20% platinum weight composition, the solidus–liquidus curves predicted by the regular solution model vastly underestimate the experimental data points provided by ref. 12, 21 and 22. Furthermore, the choice of a sub-regular solution model over a regular solution model is clearly justified by inspecting the mixing enthalpy of the alloy at various temperatures.<sup>12</sup> Indeed, with respect to chemical composition, a sub-regular model is more adequate when the mixing enthalpy of the alloy is non-symmetrical with respect to the chemical composition at 50%, while the regular

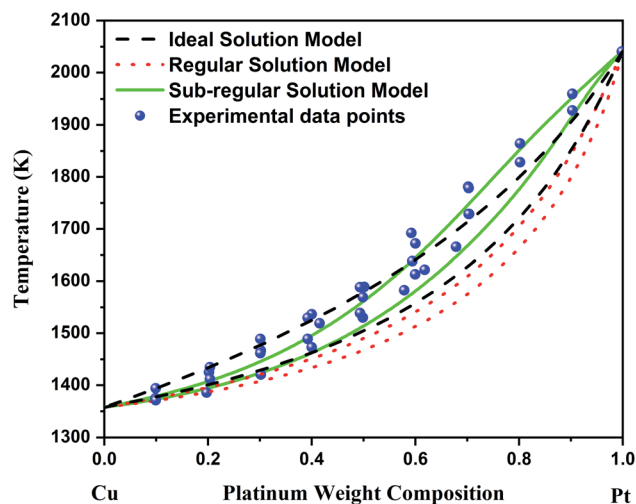


Fig. 1 Bulk phase diagram of Cu–Pt. The black solidus–liquidus curves are obtained by using an ideal solution model. The red solidus–liquidus curves are obtained by using a regular solution model as performed by Che *et al.*<sup>8</sup> The green solidus–liquidus curves are obtained by using a sub-regular solution model as shown by Abe *et al.*<sup>12</sup> The experimental data points come from ref. 12, 21 and 22.

solution model is preferred when the mixing enthalpy is symmetrical with respect to the chemical composition at 50%.<sup>16</sup>

At the nanoscale, Cu–Pt alloys experience size effects just as any other material.<sup>23</sup> Therefore, the nanophase diagram for a Cu–Pt spherical nanoparticle was predicted by using nanothermodynamic concepts; see Fig. 2. The size-dependent parameters used in the sub-regular solution model were varied by using the following relationship:<sup>24–26</sup>

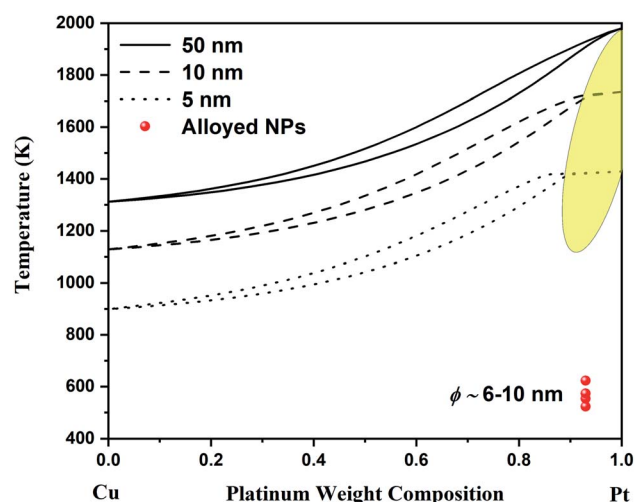


Fig. 2 Nano-phase diagram for spherical Cu–Pt nanoparticles having a diameter of 50, 10 and 5 nm. The solidus–liquidus curves are obtained by using a sub-regular model. The experimental data come from Belenov *et al.*<sup>30,35</sup> The experimental data points symbolize alloyed Cu–Pt nanoparticles after undergoing a thermal treatment. The diameter of the nanoparticles ranged between ~6 and ~10 nm. The yellow region indicates the region where the interstitial Cu–Pt alloy is expected.



$$\frac{T_m(D)}{T_{m,\infty}} = \frac{\Delta H_m(D)}{\Delta H_{m,\infty}} = \frac{G_m(D)}{G_{m,\infty}} = 1 - \frac{\alpha_{\text{shape}}}{D} \quad (3)$$

where  $T_{m,\infty}$ ,  $\Delta H_{m,\infty}$  and  $G_{m,\infty}$  are the bulk melting temperature, bulk melting enthalpy, and bulk Gibbs free energy of mixing, respectively.  $T_m$ ,  $\Delta H_m$  and  $G_m$  are the size-dependent melting temperature, size-dependent melting enthalpy, and size-dependent Gibbs free energy of mixing, respectively, evaluated at a size  $D$ . The  $\alpha_{\text{shape}}$  parameter quantifying the size effect is defined as:<sup>13</sup>

$$\alpha_{\text{shape}} = \frac{(\gamma_s - \gamma_l)D}{\Delta H_{m,\infty}} \frac{S}{V} \quad (4)$$

where  $S/V$  is the surface to volume ratio of the shape being studied,  $\gamma_s$  is the solid surface energy, and  $\gamma_l$  is the liquid surface energy. Since, we are considering a sphere as the shape for the nanoparticle,  $S/V$  is equal to  $6/D$ , where  $D$  is the diameter of the nanoparticle. It is clear from Fig. 2 that the Cu–Pt phase diagram is shifted towards lower temperature as the size is reduced. Furthermore, the liquid–solid region (*i.e.* region between the solidus and liquidus curves) increases as the size reduces due to the size effect on the Gibbs free energy of mixing. It is interesting to note that at small sizes of  $\sim 5$  nm, the Cu–Pt alloyed nanoparticle with a platinum weight composition larger than 90% behaves like a pure Pt nanoparticle. Indeed, congruency occurs in that chemical composition range. Cu atoms seem to occupy interstitial positions within the crystalline structure made of Pt atoms until the chemical composition of Pt drops below the 90% platinum weight composition. This behavior could be explained by Cu–Pt being a substitutional solid solution within 0–90% platinum weight composition and becoming an interstitial solid solution above 90% platinum weight composition due to the size difference between Cu and Pt atoms.

Since, the Gibbs free energy of mixing for Cu–Pt is negative at the bulk scale, the Gibbs free energy of mixing at the nanoscale will also be negative as per eqn (3); see ref. 13 and 24. This is supported by Monte Carlo simulations completed in ref. 27. Consequently, the binary phase diagram of Cu–Pt will not display any miscibility gap at the nanoscale, meaning that the formation of alloyed Cu–Pt nanoparticles is possible for any chemical composition. Experimentally, the synthesis of alloyed Cu–Pt nanoparticles with sizes around 3 nm and around 4–15 nm was successfully achieved by Khanal *et al.*<sup>28</sup> and Peng *et al.*,<sup>29</sup> respectively. Furthermore, Belenov *et al.*<sup>30</sup> demonstrated that Cu–Pt nanoparticles with sizes between  $\sim 6$  and  $\sim 10$  nm formed a solid solution (*i.e.* an alloy) at  $\sim 523$ – $623$  K. Those results are well in agreement with our predicted nano-phase diagram displayed in Fig. 2 as the experimental data points are within the solid region of the phase diagram.

In order to further support the theoretical investigation, Cu–Pt alloyed nanoparticles were synthesized by wet chemistry, which is a common method to synthesize Cu–Pt nanoparticles.<sup>28–31</sup> The chemicals used were purchased from Sigma Aldrich and include platinum(IV) chloride, copper(II) chloride, oleylamine ( $\text{C}_{18}\text{H}_{35}\text{NH}_2$ ), 1-dodecanethiol ( $\text{C}_{12}\text{H}_{26}\text{S}$ ), and hydrochloric acid. Both platinum chloride and copper chloride

were prepared as aqueous solutions at a concentration of 0.1 M. Platinum chloride was prepared in 1 M hydrochloric acid, while copper chloride was in ethanol. 100  $\mu\text{l}$  of the platinum solution was poured in a 15 ml glass vial and heated to 100  $^\circ\text{C}$  to evaporate any water. Once dried, 5 ml of oleylamine was added, the solution was stirred and 100  $\mu\text{l}$  of copper chloride was added. The solution was heated to 250  $^\circ\text{C}$  for 2 hours. There was a color change from golden yellow to deep red. 10  $\mu\text{l}$  of 1-dodecanethiol was then added to the solution. The reaction was finished once the color of the solution changed to a dark brown color, which was achieved after  $\sim 1$  hour.

The nanoparticles were then observed with a NION Scanning Transmission Electron Microscope (STEM) operating at 80 kV with atomic resolution. For STEM observations, samples were prepared by drop-casting one or two droplets of the colloidal solution onto a holey carbon coated TEM grid, which was then let to dry in air at room temperature. A careful observation of the synthesized nanoparticles through Electron Energy Loss Spectroscopy (EELS) confirmed the alloyed structure of the Cu–Pt nanoparticles (Fig. 3). Fig. 3a–c show EELS chemical mapping for Cu and Pt elements as well as the overlay between Cu and Pt, respectively. From the observation of Fig. 3, it is clear that Cu and Pt are well-mixed throughout the nanoparticle. Additionally, it is observed that Pt segregates preferentially to the surface compared to Cu. This is well in agreement with the segregation rules established by Guisbiers *et al.* in ref. 32.

Let's look in detail at one Cu–Pt nanoparticle in particular through the high angle annular dark field image mode (HAADF). Fig. 4a shows a nanoparticle of Cu–Pt. The whole structure of the particle was a superlattice made of intercalated layers of Pt and Cu, with a FCC structure, occupying regular lattice sites. However, it is clearly seen that Cu atoms were also occupying interstitial positions (Fig. 4b). In Fig. 4b, a closer look at the distance between two bright spots is 0.675 nm, corresponding to the separation distance between two Pt atoms. However, this distance needs to be divided by 2 (0.338 nm), to get the distance between two Pt planes in the superlattice

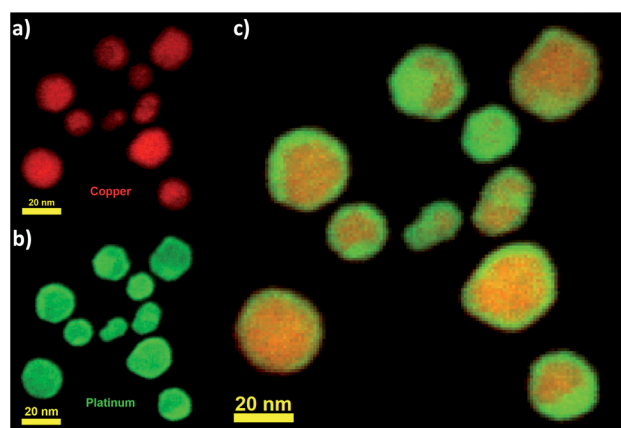


Fig. 3 Chemical mapping of several Cu–Pt nanoparticles acquired using EELS. (a) 2D STEM-EELS map showing the Cu element, (b) 2D STEM-EELS map showing the Pt element and (c) 2D STEM-EELS map showing the overlay of Cu and Pt elements.





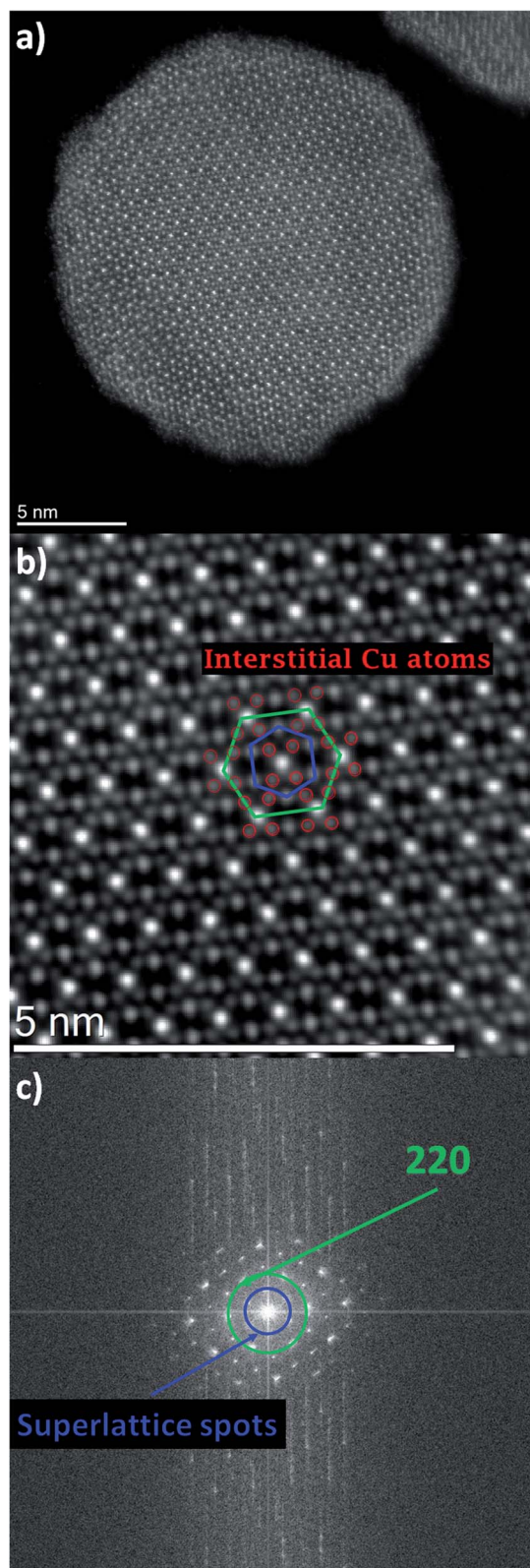


Fig. 4 (a) HAADF-STEM image of a Cu–Pt nanoparticle having a size around 20 nm. Bright spots represent Pt, while the weak spots represent Cu. (b) Zoomed-in image of the core of the nanoparticle shown in (a). The green hexagon connects the Pt atoms. The blue hexagon connects Cu atoms together. The red circles indicate the Cu atoms, which are in interstitial positions. (c) FFT pattern of the nanoparticle shown in (a).

structure. To better characterize the crystallinity of the nanoparticle, Fast Fourier Transform (FFT) was applied on Fig. 4a and is shown in Fig. 4c. The inner most spots denote a superlattice structure in the particle. The brighter spots of the second ring indicate growth along the (220) direction.

The interstitial presence of Cu atoms was also observed with other particles on the same grid with sizes ranging between  $\sim 2$  and  $\sim 20$  nm. Due to the small size, high resolution was not possible; however, FFT information was still gathered from those nanoparticles. Fig. 5 shows the TEM images of two small particles with their corresponding FFT pattern. In Fig. 5a, the particle size is measured to be around  $\sim 7.6$  nm and its FFT pattern showed the same superlattice structure as that observed in larger particles. Fig. 5b shows a smaller particle on the order of  $\sim 3.8$  nm and its FFT also indicates a structure similar to the other nanoparticles as confirmed by the two diffraction rings.

Although theory predicts this transition to occur in a small composition range above 90%, experimentally, it is observed in a much larger composition range. By looking at Fig. 2, it is noticeable that the theoretical region where the transition occurs occupies a larger compositional range as the size of the particle decreases. It is important to remember that nanothermodynamics describes a system (*i.e.* a nanoparticle in this case) at equilibrium, while the synthesis of nanoparticles may not actually be an equilibrium process during the entire synthesis time and consequently depend on kinetic effects. These effects are the source of the difference between theoretical predictions and experimental observations with respect to the chemical composition range. Theory and experiment agree on Cu–Pt nano-alloys having substitutional–interstitial structural transition below  $\sim 20$  nm. Finally, this study points out that nanothermodynamics may be used as a fast theoretical framework to identify phase transitions in nano-alloys. Indeed, molecular dynamics and density functional theory studies typically take a very long time and are limited to very small sizes due to their complexity and the iterative nature of those techniques.

In conclusion, the bulk phase diagram for the Cu–Pt system has been calculated in this work by using a sub-regular solution model which is found to be in excellent agreement with all the available experimental data. Furthermore, the nanophase

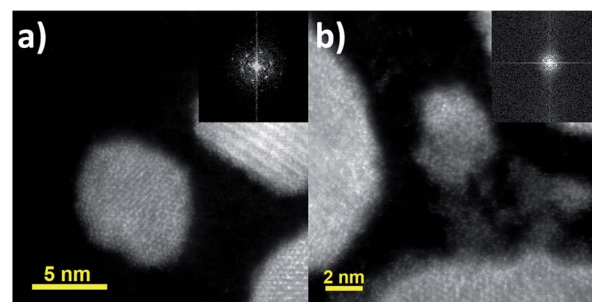


Fig. 5 (a) TEM image of a Cu–Pt nanoparticle having a diameter around  $\sim 7.6$  nm. (b) TEM image of a Cu–Pt nanoparticle having a diameter around  $\sim 3.8$  nm. In both images, the corresponding FFT of the particle is shown in the upper right corner as inset.



diagram for spherical Cu–Pt nanoparticles was provided at various sizes (50, 10 and 5 nm), where solidus–liquidus curves were also obtained by using a sub-regular solution model. From the theoretical calculations, copper and platinum continue to mix fairly well all over the composition range whatever the scale of the system is. Finally, a transition between a substitutional solid solution and an interstitial solid solution is predicted theoretically above 90% platinum weight composition at sizes around 5 nm. This structural transition has been observed experimentally on Cu–Pt nanoparticles having sizes less than ~20 nm in a composition range up to ~50%. Therefore, the theoretical framework used in this paper may be used as a predictive tool to identify structural phase transitions in binary nano-alloys. This discovery is particularly interesting for potential future applications, as the amount of platinum is reduced while preserving the thermal properties of pure platinum.

## Data availability

The data that support the findings of this study are available from the corresponding author upon reasonable request.

## Conflicts of interest

There are no conflicts to declare.

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